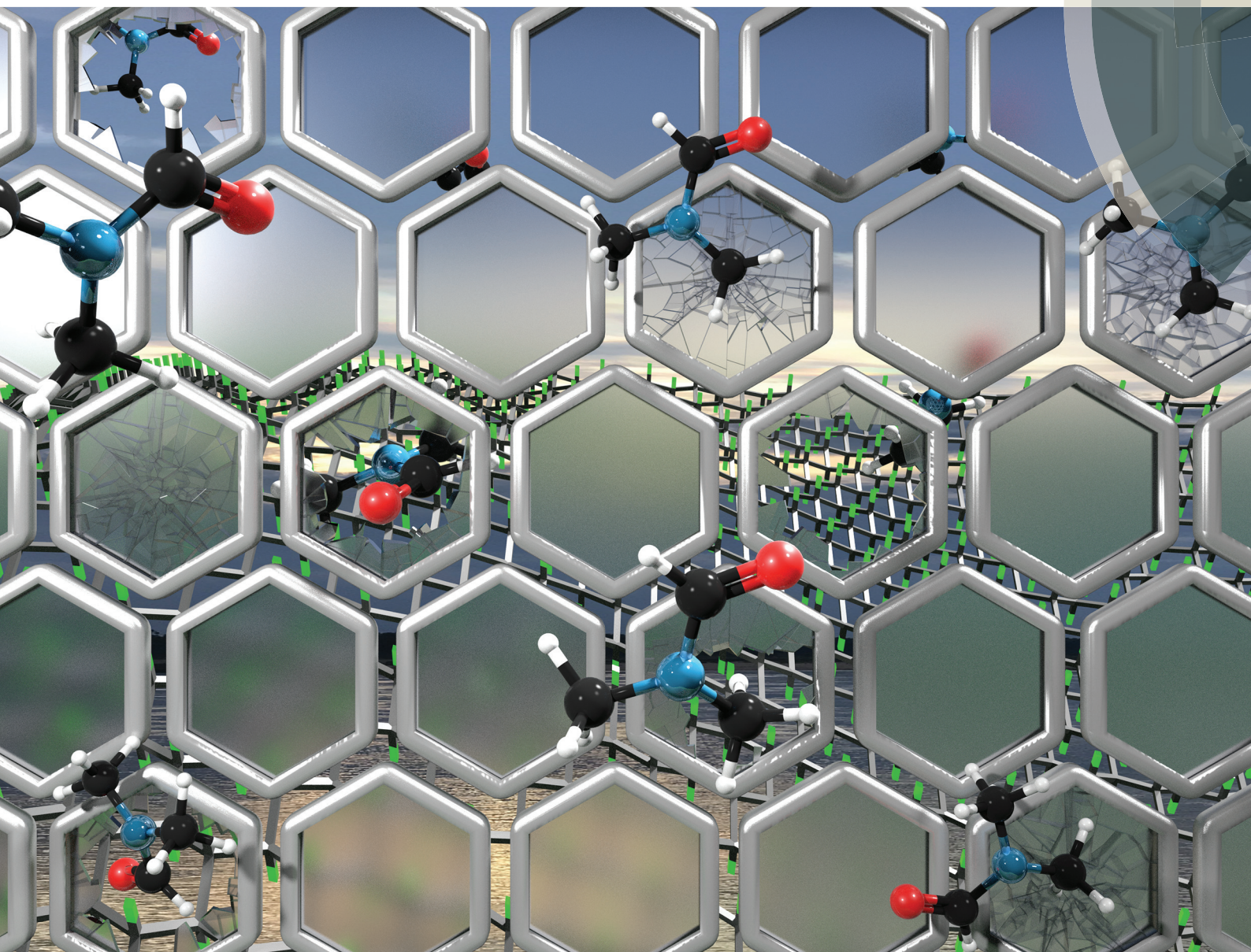


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Reactivity of fluorographene is triggered by point defects: beyond the perfect 2D world



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Reactivity of fluorographene is triggered by point defects: beyond the perfect 2D world†

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Preparation of graphene derivatives using fluorographene (FG) as a precursor has become a key strategy for the large-scale synthesis of new 2-D materials (e.g. graphene acid, cyanographene, allyl-graphene) with tailored physicochemical properties. However, to gain full control over the derivatization process, it is essential to understand the reaction mechanisms and accompanying processes that affect the composition and structure of the final products. Despite the strength of C–F bonds and high chemical stability of perfluorinated hydrocarbons, FG is surprisingly susceptible to reactions under ambient conditions. There is clear evidence that nucleophilic substitution on FG is accompanied by spontaneous defluorination, and solvent-induced defluorination can occur even in the absence of any nucleophilic agent. Here, we show that distributed radical centers (fluorine vacancies) on the FG surface need to be taken into account in order to rationalize the defluorination mechanism. Depending on the environment, these radical centers can react as electron acceptors, electrophilic sites and/or cause homolytic bond cleavages. We also propose a new radical mechanism of FG defluorination in the presence of *N,N'*-dimethylformamide (DMF) solvent. Spin-trap experiments as well as ¹⁹F NMR measurements unambiguously confirmed formation of *N,N'*-dimethylformyl radicals and also showed that *N,N'*-dimethylcarbonyl fluoride plays a key role in the proposed mechanism. These findings imply that point defects in 2D materials should be considered as key factor determining their chemical properties and reactivity.

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1. Introduction

The discovery of graphene¹ in 2004 opened a route to a new class of carbon-based 2D materials with unique chemical, electronic and magnetic properties. Whereas graphene itself is chemically a low reactive semimetal with high electron conductivity, even dilute concentrations of chemisorbed impurities can dramatically change its band structure and induce local magnetic moments.² Increasing the adsorbate concentration can further disrupt the π -conjugated network, transforming the conducting material to an insulator through a variety of intermediate states.^{3,4} Possibilities for tuning the band gap characteristics and associated electronic and mag-

netic properties by graphene derivatization provide an auspicious base for future technological advancements, including development of batteries,^{5–7} biosensing,⁸ gas sensing on surfaces⁹ and solar cell technologies.¹⁰ Nevertheless, selective and controllable covalent graphene functionalization remains a challenging task due to the limited reactivity of graphene.^{11,12} Many proposed direct derivatization strategies suffer from a low degree of functionalization^{13–16} or require harsh reaction conditions that strongly affect the final stoichiometry and chemical structure of the resulting graphene derivative.^{17,18}

Fluorographene (FG)^{19–21} offers a promising alternative for graphene derivatization because (i) chemisorption of monovalent fluorine atoms prevents formation of complex (non-stoichiometric) structures; (ii) the material is thermally stable up to 400 K; (iii) the pristine 3D material, graphite fluoride, is available in large-scale as an industrial lubricant; and (iv) FG reacts readily as an electrophile under mild conditions.^{22–26} Owing to these features, FG has recently been used as a precursor for the synthesis of several graphene derivatives, including amino-graphenes,^{26–30} cyanographene,³¹ carboxygraphene (graphene acid),³¹ sulfur-linked derivatives^{25,32} and alkylated graphenes.³³ During the chemical transformation of FG, substitution of fluorine atoms proceeds simultaneously with

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In our work, we present a description of the mechanism of the reactivity of FG obtained by combining electron paramagnetic resonance (EPR) and ^{19}F NMR experiments with theoretical calculations. We focus on FG in DMF, which is a widely used solvent for FG exfoliation and chemical treatment. EPR experiments confirmed the presence of radical centers (fluorine vacancies) in FG and allowed their distribution and time-evolution to be studied. We observed defluorination in DMF, despite the fact that it is not a sufficiently strong reducing agent for FG. Based on theoretical calculations, we suggest a new mechanism involving N,N' -dimethylformyl radicals ($\text{DMF}^\bullet$), formation of which was evidenced by a series of spin-trap experiments. In addition, ^{19}F NMR measurements unambiguously confirmed formation of N,N' -dimethylcarbamoyl fluoride, which plays a key role in the proposed mechanism. These findings show that defects present in FG determine the chemical behavior of this 2D material and make it susceptible to defluorination and subsequent substitution and addition reactions. Such defects may be important in determining the reactivity of other 2D materials because they may initiate reactions which then propagate on the surface of 2D materials.

2.1 Chemicals

EPR spectra were recorded on a JEOL JES-X320 electron spin resonance spectrometer operating at the X-band frequency ($\sim 9.16\text{--}9.17$ GHz) and equipped with variable temperature control ES 13060DVT5 apparatus. The cavity Q quality factor was kept above 6000 in all measurements and signal saturation was avoided by working at low-applied microwave powers. Concentrated suspensions of GrF (15 mg GrF per mL solvent, DMF or benzene) thoroughly degassed under N_2 were mechanically stirred at 40°C without sonication and kept in airtight and oxygen-free vials, and EPR spectra of the resulting dense suspensions were recorded during a 10 day time-frame. Use of very concentrated suspensions of GrF was necessary for successful spin-trapping experiments with POBN (see ESI for details, section 1.2 $\dagger$). Generally, 0.1 mL of GrF suspension was loaded into the EPR tubes by using airtight syringes. In the case of GrF powder, 10 mg were used. For the spin-trap experiments, the reaction mixtures (GrF/solvent/POBN reacted together at RT for 40 min) were centrifuged prior to loading 0.1 mL of the supernatants into the EPR tubes. Discussion on the use of GrF instead of FG for the EPR studies and radical formation on the carbon lattice is available in the ESI. $\dagger$ Highly-pure quartz tubes were employed (Suprasil, Wilmad, ≤ 0.5 OD). Simulation of the EPR resonances were carried out using the WinEPR SimFonia software (V.1.25, Dr Ralph T. Weber, EPR Division, Bruker Instruments, Inc. Billerica, MA USA) with second-order perturbation theory. The frequency unit can be converted into magnetic fields units through the following relation: 0.1 mT equals to 2.8025 MHz. Nuclear magnetic resonance (^{19}F) spectra were obtained with a JEOL ECA-500 NMR spectrometer (^{19}F frequency 470 MHz, $\text{CFCl}_3 = 0$ ppm). Further experimental details can be found in the ESI. $\dagger$

The ground state (GS) structures of all the investigated species were optimized by the ω B97X-D method³⁷ using 6-31+G(d) and 6-31++G(d,p) basis sets.³⁸ Solvent effects were included by using the universal continuum solvation model based on solute electron density (SMD).³⁹ Whereas the structures of small and medium size systems (F, F⁻, HF, DMF, F-DMF, *etc.*) were fully relaxed in geometry optimizations, to mimic the semilocal flexibility of FG sheets, FG-like structures were obtained by constrained geometry optimizations, keeping the edge carbon atoms frozen. All calculations were performed with the Gaussian09 program.⁴⁰

The stability of GrF dispersed in two model solvents (benzene and DMF) was probed using electron paramagnetic resonance (EPR) spectroscopy by monitoring the variation of the respective EPR envelopes over time. Fig. 1a shows the frozen matrix ($T = 123$ K) X-band EPR envelope recorded for the freshly prepared GrF/benzene suspension. Although the overall reso-

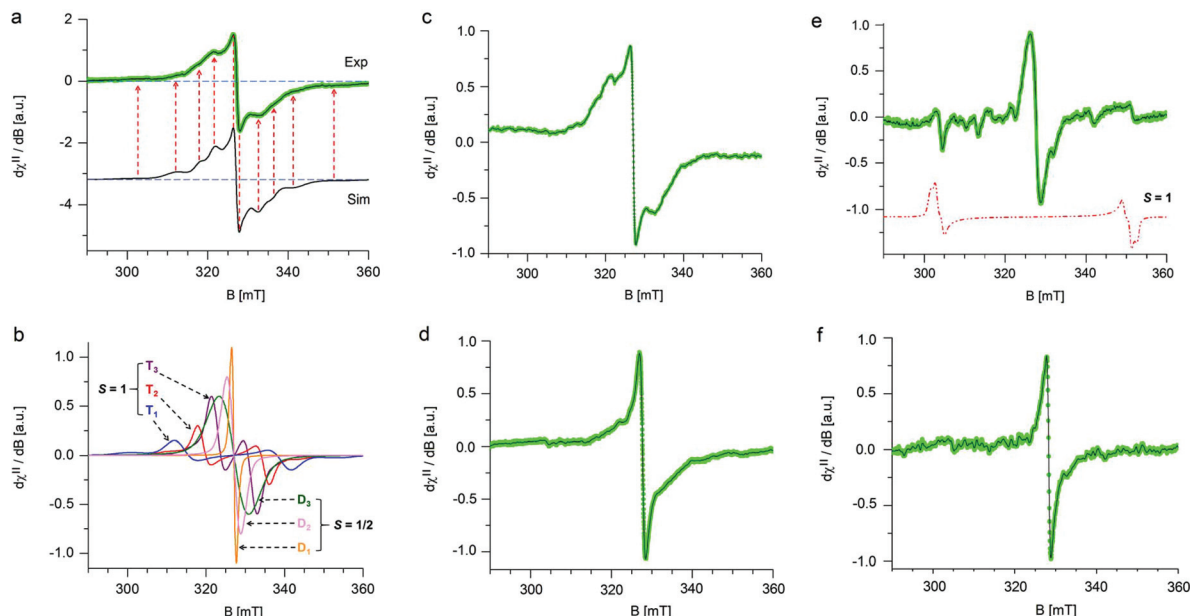


Fig. 1 Low temperature ($T = 123$ K) frozen matrix X-band EPR spectra of GrF dispersed in benzene or DMF recorded at various times. (a) GrF freshly dispersed in benzene (green line, Exp) and its EPR simulation (black line, Sim). (b) Overview of the diverse spin components encoded in the EPR envelope shown in (a, Sim); the labels T_1 , T_2 and T_3 indicate triplet ($S = 1$) species, whereas D_1 , D_2 and D_3 indicate doublet ($S = 1/2$) species. The relative weights were as follows: T_1 (4.3%), T_2 (8.4%), T_3 (16.9%), D_1 (31.0%), D_2 (22.5%), D_3 (16.9%). (c) GrF freshly dispersed in DMF. (d) GrF/DMF solution after 6 h. (e) GrF/DMF solution after 7 days. (f) GrF/DMF solution after 10 days. Experimental parameters: 9.14–9.16 GHz frequency, 100 kHz modulation frequency, 30 ms time constant, 0.5–0.8 mT modulation width, 0.4–0.6 mW applied microwave power. In (a) and (c–f), the solid black line corresponds to the resolution enhanced EPR resonance line (Savitzky–Golay, denoise algorithm).

nance line of the system in benzene was rather broad (~ 30 mT field-spread), it exhibited several small resonances that developed symmetrically with respect to an intense narrow central line, a signature of overlapped contributions of triplet and doublet spin-states. The observed g -averaged (g_{avg}) value was estimated to be 2.000 and validated by recording the spectrum together with the Mn(II)MgO standard (JEOL internal reference standard, Mn(II) , $g_{\text{eff}} = 2.00101 \pm 0.00005$, ESI, Fig. S5 and S6†). A comparable EPR spectrum was observed for the neat GrF powder (ESI, Fig. S1a and S3†). Similar EPR spectra have been previously reported around the $g = 2$ region (~ 280 – 380 mT) by Panich and co-authors⁴¹ on polycrystalline graphite fluoride (CF_n) obtained by fluorination of petroleum coke as well as by Giraudet and co-authors in a later report (~ 310 – 380 mT region).⁴² However, we noticed that by recording the EPR signal of the GrF powder in a broader magnetic-sweep range (50–550 mT), a broad tail emerged in the high field region (360–550 mT) together with a broad resonance signal in the low field region (around $g = 4$) (ESI, Fig. S3†). These features became strongly weakened in the diluted systems (GrF/benzene and GrF/DMF). The ESI provides additional EPR spectra for GrF, including the EPR power-saturation behavior (Fig. S1a, S2–S4†). Panich and co-authors suggested that the fine structure observed around $g = 2.00$ in GrF was due to resolved hyperfine components (A_F) with ~ 5 mT of splitting value, originating from the interaction of an unpaired spin ($S = 1/2$) with six neighboring fluorine nuclei. The authors also

found that these hyperfine lines were rather broad ($\Delta H = 30$ G by simulation), a factor that pointed towards the joint effect of dipole–dipole and exchange interactions between paramagnetic centers. In the later work, Giraudet and co-authors took the similar approach in the simulation of the observed EPR envelope for GrF, but with the addition of two broad resonance lines. The simulation of the signal was obtained by adding the hyperfine contribution of six fluorine nuclei ($A = 45 \pm 2$ G, line width of 36 ± 2 G, and g factor of 2.003 ± 0.001). As stated by the authors, and similar to the Panich's work, “the large line width is interpreted by the joint effect of dipole–dipole and exchange interactions between paramagnetic centers”. In the case of the GrF/benzene frozen matrix sample, we therefore tried to disentangle these effects previously anticipated by ref. 41 and 42 and carried out an EPR simulation of the resonance signal in the $g = 2.0$ region under the assumption of an overlapped contribution of different spin components (Fig. 1a, black line, Sim, and Fig. 1b showing the individual spin components with their relative weights). The observed EPR signal of GrF/benzene shown in Fig. 1a (Exp) can be interpreted as a superposition of three effective triplet states ($S = 1$) and three sets of spin-doublets ($S = 1/2$), *i.e.*, in GrF exists a distribution of spin-carrying defects, some close enough (through-space distance) to show signatures of high-spin systems and others behaving as isolated $S = 1/2$ defects. This hypothesis is supported by the observation that the low field signal around $g = 4.0$ observed in neat GrF can still be observed in the dilute



GrF/benzene suspension through signal accumulation (ESI, Fig. 1b,† GrF/benzene system, 133 K, 4 mW, 40 scan accumulated and averaged) and may represent the half-field transition ($\Delta m_s = 2$) of these high spin ($S = 1$) components. The triplet signatures (Fig. 1b, marked T_1 , T_2 , and T_3 , with g -tensor (x, y, z) of 2.000, 2.000, 2.000) gave from simulation the following axial values (D): one large with $|D|$ of 701 MHz (marked as T_1) and negligible ^{19}F hyperfine contribution, one intermediate (marked as T_2) with $|D|$ of 462 MHz and with line broadening by hyperfine terms (four ^{19}F nuclei, $A_{x,y,z}$ of 56.0, 14.0, 14.0 MHz), and one much smaller (marked as T_3) with $|D|$ of 202 MHz, with similarly broadened resonances (two ^{19}F nuclei, $A_{x,y,z}$ of 38.0, 12.0, 12.0 MHz). From the point-dipole-approach, these triplet configurations (and D values) indicate a spread of through-space distances among the spin-carrying centers that are effectively magnetically coupled together, with $S_{1/2}$ – $S_{1/2}$ through-space distances that lie within 4.8–7.3 Å. These distances are an estimate, assuming dominant $1/r^3$ dependence on the D value and a dominant spin–spin (SS) contribution.⁴³ The biquadratic zero-field-splitting (ZFS) D -tensor comprised two contributions: (i) a first order term arising from direct magnetic dipoles SS interaction, and (ii) a second-order contribution arising from spin–orbit coupling (SOC). In organic based magnets, the first term usually dominates. It is known that the point-dipole approach gives rather poor estimation of inter-center radical distances (compared to X-ray radical distances) if the radical units interact *via* π -bond pathways. In contrast, when the interacting radicals are connected through an array of σ -bonds, which tends to hamper spin delocalization, the point dipole approach gives fairly accurate distances, comparable to those obtained by X-ray.^{44–46} The other types of spin active species are shown by the three sets of $S = 1/2$ signals in Fig. 1b, marked as D_1 , D_2 and D_3 . One doublet (D_1 , line-width tensor L_{xyz} of 28.0, 28.0, 56.0 MHz) corresponded to the relatively sharp and slightly asymmetric central signal at $g_{\text{eff}} = 2.000$. This fingerprint indicates that the $S = 1/2$ spins experienced the strong exchange regime, *i.e.*, the exchange energy (H_{ex}) of the spin packet was much larger than H_{dip} (the dipolar interaction component), which in turn was greater than H_{hf} (the hyperfine component). Considering the upper limit of the strongest triplet obtained by simulation (701 MHz), this indicates that the exchange term (J) associated to these spins was stronger than 0.02 cm^{-1} . Since the applied microwave power was kept low and under non-saturating conditions ($\ll 1 \text{ mW}$), this sharp central line could not have arisen from a double-quantum transition and very unlikely originated from a quartet spin state ($S = 3/2$). The second broad doublet (D_3 , Fig. 1b) belonged to an ensemble of $S = 1/2$ spins that experienced weak dipolar interaction ($< 100 \text{ MHz}$) but still larger than ^{19}F hyperfine ($H_{\text{dip}} > H_{\text{hf}}$). This effect resulted in a very broad line-width (L_{xyz}) of 250.0, 196.0, 196.0 MHz encoded in the electronic transition. The third doublet (D_2 , Fig. 1b) corresponded to $S = 1/2$ spins that experienced an intermediate regime with $H_{\text{dip}} \sim H_{\text{hf}}$ and a relatively narrow line-width tensor in the frozen matrix (L_{xyz}) of 98.0, 98.0, 98.0 MHz. Overall, our observations support the earlier description of Grigorieva and

co-authors on the dominant spin-half paramagnetism for GrF-based materials.⁴⁷ Following GrF aging in benzene solution, we did not observe severe changes in the overall EPR resonance line, which remained after 6 days and 10 days similar to that shown at the beginning of the process (ESI, Fig. S8†). Thus, benzene can be considered a “nearly inert” solvent for GrF which does not induce substantial degradation of the paramagnetic centers. Chemical stability of GrF in benzene was also confirmed by measurement of the XPS spectrum for a 10-days aged GrF sample in benzene (Fig. S25 in ESI†). After dispersion of GrF in DMF, the freshly prepared suspension showed the same EPR fingerprints in a frozen matrix (Fig. 1c) as witnessed for GrF/benzene (Fig. 1a). Additional EPR spectra for GrF/DMF are given in the ESI† (power saturation behavior of the entire resonance line, from 50 to 550 mT, Fig. S9a–c†). The concentration of spin centers on the GrF surface at this stage was estimated to be 23.2×10^{18} spin per g, which corresponds to *ca.* 1 center per area of $5 \text{ nm} \times 5 \text{ nm}$ (see ESI, section 1.2†). This value closely resembles the estimated spin concentration found to near stoichiometric C/F ratios in GrF ($\text{C}/\text{F}_{0.9-1.0}$) by Grigorieva and co-authors, which were obtained from in-depth analysis of bulk magnetic susceptibility experiments ($\sim 22\text{--}17 \times 10^{18}$ spin per g).⁴⁷ However, already after 6 h stirring at 40°C , the GrF/DMF suspension darkened considerably. The X-band EPR spectrum ($T = 123 \text{ K}$) observed after 6 h revealed some changes in the overall resonance line (Fig. 1d) compared to the initial spectrum. The doublet component ($S = 1/2$, D_1 in the simulation shown in Fig. 1b) appeared to increase strongly in intensity (relative increase) at the expense of the T_1 , T_2 and T_3 triplet components (see $S = 1$ components, Fig. 1b and ESI, Fig. S10†) and the doublet component D_2 . After 7 days of stirring, the EPR spectrum of FG/DMF showed more significant changes (Fig. 1e). Besides still being centered at $g_{\text{eff}} = 2.000$, the doublet components D_1 and D_2 became dominant, several resonances appeared throughout the spectrum and a novel but very minor triplet emerged. The spin-Hamiltonian simulation of this species was very tentative (shown by the dotted red-line in Fig. 1e). This system should exhibit a large axial value ($|D|$ of 1345 MHz, $|E/D|$ of 0.04). In this case, the inter-spin (through-space) distance of the triplet configuration ($S = 1$, arising from two interacting $S = 1/2$) was estimated from the point-dipole approach to fall short, at about 3.87 Å. This value corresponded well with the minimum through-space spin–spin distance of the spin carrying defects estimated by theory (*vide infra*, Fig. 4b). It is important to note that at such later stage of the GrF aging process, an alternative explanation for the appearance of the large triplet signal reported above can be put forward. Namely, it may arise from voids in the GrF sheets, especially in the areas corresponding to sp^2 islands that are formed upon extended defluorination. As discussed by Palacios and co-authors,⁴⁸ the presence of voids, such as type A3B, can lead to non-vanishing local magnetic moments. Consequently, $S = 1$ signatures can arise in the EPR spectrum from the presence of these sublattice imbalances ($N_{\uparrow} = N_{\text{A}} - N_{\text{B}}$). When the Hubbard repulsion is active, Lieb's theorem grants the $S = 1$ spin state to such defect struc-



3.2 Origin of radical centers

EA, the extent of its delocalization increased with the system size. To sum up, although the energetically low-lying C-F σ^* orbital doubtlessly plays an important role in fluorine chemistry (*e.g.*, through orbital mixing, as explained by Borden applying the second-order perturbation theory⁵⁶), it is improbable that mild reducing agents could initiate a radical cascade process of defluorination by direct electron transfer to the C-F σ^* orbital.

Although formation of radical centers on an ideal C_1F_1 layer was expected to be hampered by the low EA, the experimental evidence of observable spin concentrations in C_1F_1 structures clearly proves the presence of isolated sp^3 carbon sites. The origin of such centers depends largely on the preparation process. Recent experimental surface imaging techniques have shown that fluorination of CVD graphene with XeF_2 leads to formation of highly ordered sizable (150 nm) CF , C_2F chair and unfluorinated domains,⁶⁰ whereas fluorination with CF_4 plasma gives rise to inhomogeneous and disordered microscopic configurations.⁶¹ From a radical center formation point of view, in the former case, the carbon atoms at the boundary between these domains can remain unfluorinated as the attachment of fluorine on such sites would lead to a *cis* configurations, which would be energetically unfavorable.⁶² On the other hand, the inhomogeneous spatial distribution of fluorine on CVD graphene treated with CF_4 plasma is accompanied by formation of multilayer islands and structural features such as folds, wrinkles, and ripples that are less fluorinated, which consequently increases the probability of spin center emergence.

The presence of radical centers (even in very low concentrations) can obviously dramatically change the behavior of FG in the presence of reducing agents or nucleophilic species. As shown in Table S1 (see ESI[†]), the electron affinities of radical species derived from parent perfluorinated model systems are very high. For instance, AEA of a fluorinated coronene radical is 80.6 kcal mol⁻¹, *i.e.*, it is even larger than that of a fluorine atom (78.4 kcal mol⁻¹). These extremely large values of EA of perfluorinated radicals are mainly due to the electron-withdrawing effect of highly electronegative fluorine atoms bound to the nearest carbon atoms but also increased delocalization of the negative charge with increasing size of the system. This delocalization is closely related to the presence of the already discussed C-F σ^* orbitals on neighboring carbons, which can stabilize the SOMO orbital in a pyramidal configuration of the radical center by mutual orbital mixing.^{56,63} The extraordinary oxidizing power of highly fluorinated organic radicals can initiate the defluorination process in the presence of even mild reducing agents. Therefore, we probed a possible scenario of initial phases of the cascade reduction of FG in which defluorination spreads outwards from a radical center (Fig. 2).

Although the total reaction enthalpies of the reduction steps depend on the strength of reducing agent as well as solvation enthalpies of the involved species, it can be seen that, in general, they are energetically favorable and polar solvents

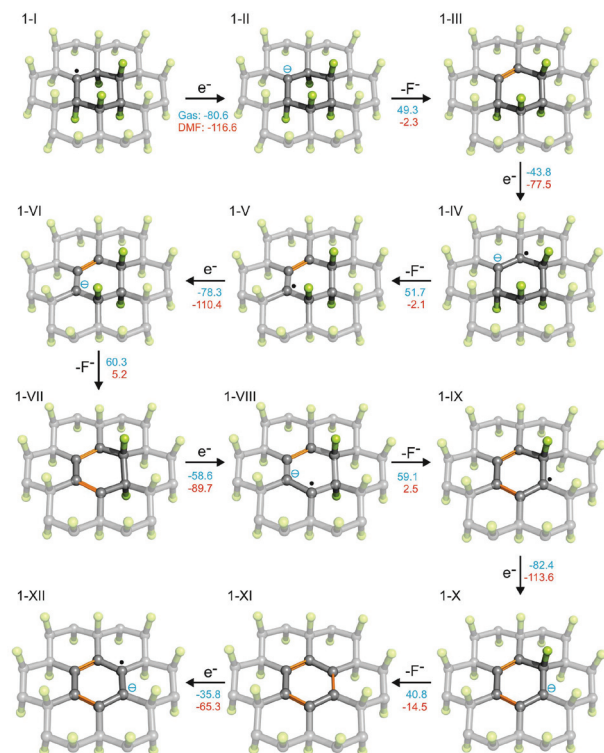


Fig. 2 Possible scenario for the initial phases of reduction of FG (the F atoms in green, the C atoms in grey/black, newly formed C=C bonds in orange). The reaction energies are given in kcal mol⁻¹ (blue in gas phase and red in DMF).

stabilize the charged intermediates. As expected, the oxidizing power of the radical species is much higher (by about a factor of 2) than that of closed-shell systems occurring in the reduction process. However, even for the latter, the EA is rather high as a consequence of the effect of highly polar C–F bonds on the neighboring carbon atoms. Note that EA of the final structure **1-IX**, which is fully π -conjugated, is significantly lower than that of the other structures. It is therefore possible that depending on the power of the reducing agent, the cascade defluorination process may cease after formation of larger π -conjugated domains. Finally, let us consider possible formation of fluorine radicals by charge transfer from fluoride anions to the radical center. Although the gas phase value of the EA of **1-I** is slightly larger than that of an F atom, the stability of a small fluoride anion is significantly increased in polar solvents (much more than **1-II**), and therefore formation of F radicals by this mechanism would be unfavorable. Let us also underline that our results for model systems well reproduce those obtained for larger systems and are also in good correspondence with periodic calculations (see section 2.5 in ESI†).

3.4 Radical defluorination mechanism in DMF

We have shown that radical centers on FG act as efficient electron acceptors. Although such electron transfer can be a trigger for cascade defluorination, in the case of a sufficiently strong reducing agent, it does not explain the defluorination

of FG observed in solvents with ionization potential in the liquid phase significantly higher than the EA of FG point defects. An example of such a solvent is DMF (discussed in section 3.1), which has an ionization potential in the liquid phase of 149 kcal mol⁻¹, compared to the EA of a radical center **1-I** of 117 kcal mol⁻¹ in the same solvent. Let us recall that electron affinities of the subsequent species are even smaller. Clearly, DMF can hardly be considered a direct reductant of FG in the sense of direct electron transfer to such radical centers.

To rationalize the observed spontaneous defluorination of FG in DMF, we suggest a mechanism initiated by hydrogen atom transfer from DMF to the FG radical center (Fig. 3). In

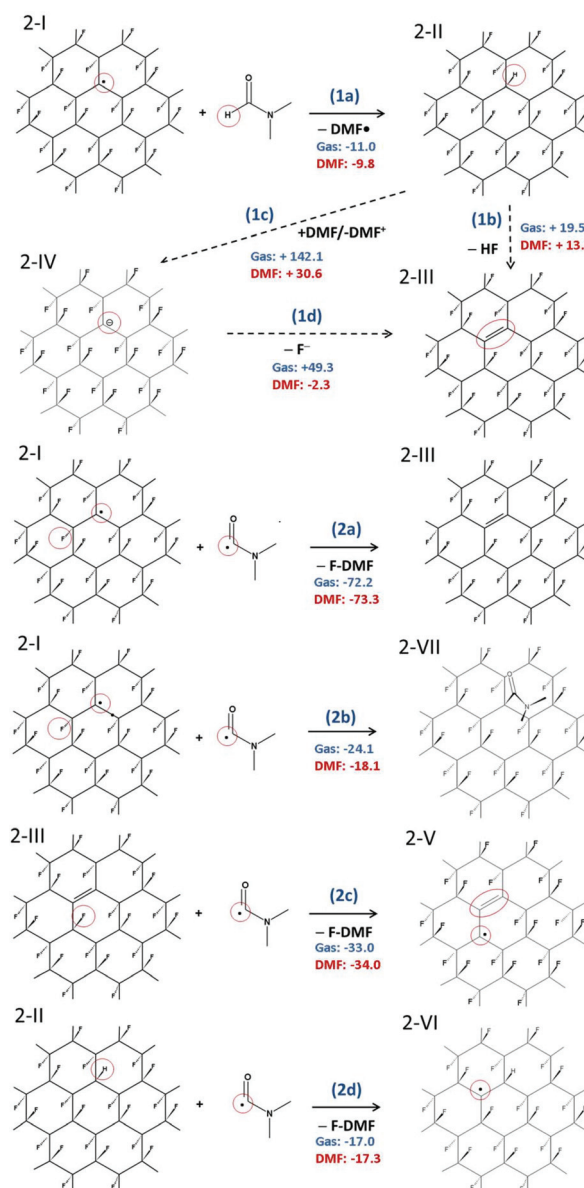
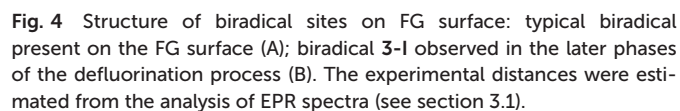


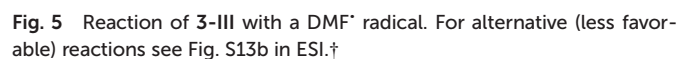
Fig. 3 Radical mechanism of defluorination of FG in DMF. (1a–c) Initial phases of defluorination. (2a–d) Re-creation of FG radical center. The reaction energies are given in kcal mol⁻¹.



As the EPR spectra of FG in DMF confirmed the presence of biradical centers (besides monoradical sites) on the FG surface, we also studied a radical mechanism of defluorination starting from structure **3-I** (Fig. 4), for which the theoretical distance between the spin centers (3.91 Å) was in reasonable



The same trapped radical species by POBN was obtained when the GrF/DMF solution was aged for a longer time (1 day, 7 days), whereas negative results were consistently found for a



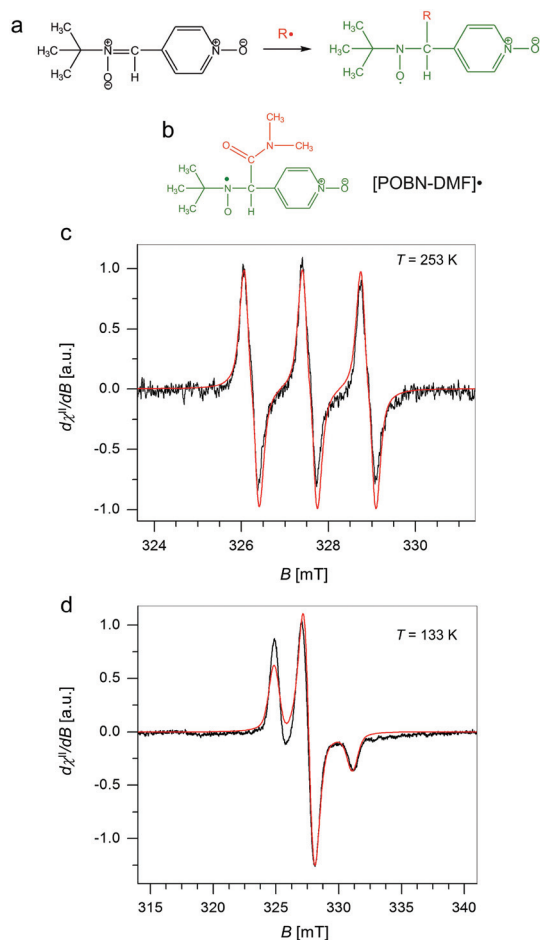


Fig. 6 (a) Scheme for the reaction of α -4-pyridyl-1-oxide-N-tert-butyl-nitron (POBN) with free radicals to generate the corresponding spin-trapped nitroxide radical. (b) Chemical structure of a DMF radical trapped by POBN. (c) X-band EPR spectra obtained from spin-trap (POBN) experiments carried out on the supernatant from a GrF/DMF/POBN mixture ($T = 253$ K) together with its EPR simulation (red-trace), and (d) corresponding spectrum obtained in a frozen matrix ($T = 133$ K) compared with simulation (red-trace). Experimental parameters: 9.16–9.17 GHz frequency, 100 kHz modulation frequency, 30 ms time constant, 0.2 mT (c) or 0.8 mT (d) modulation width, 0.3 mW (d) or 4.0 mW (c) applied microwave power. Simulation parameters are given in the text.

GrF/benzene solution aged in the presence of POBN (see ESI, section 1.3, Fig. S13†). These findings suggest that DMF reacts with GrF *via* a radical mechanism. It is important to note that there is no clear report in literature of direct trapping of carbamoyl radical derived from DMF. However, in the recent study of cross-benzoin and Stetter-type reactions mediated by KOtBu-DMF by Massi *et al.*, the authors observed generation of a weak EPR signal of difficult interpretation in the solution of KOtBu-DMF in presence of spin-trap agent PBN, signal that became well resolved into a doublet of triplet upon addition of diaryl α -diketone.⁶⁸ The EPR signal encoded hyperfine splitting coupling constants similar to those observed here for the GrF/DMF/POBN trapped radical species, with $A_N = 1.33$ mT

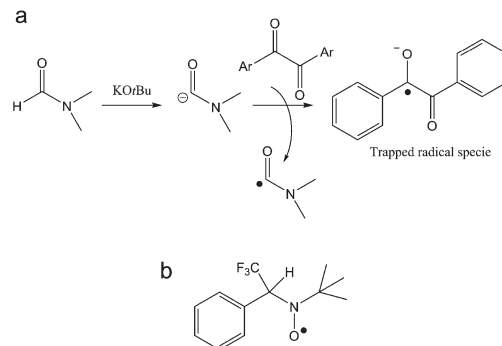


Fig. 7 (a) Possible pathway of the trapped benzyl anion radical by PBN during the cross-benzoin like reaction as adapted from ref. 68. (b) The trifluoromethyl-PBN radical adduct from ref. 69.

and $A_H = 0.177$ mT, which the authors attributed to the successful entrapment of benzil anion radical by PBN (Fig. 7a, trapped radical species). Furthermore, Aggarwal *et al.*⁶⁹ recently reported the successful entrapment of a trifluoromethyl radical by PBN in DMF at 298 K (Fig. 7b), which exhibited $g = 2.0044$, $A_N = 1.411$ mT, $A_H = 0.1664$ mT, $A_F = 0.1781$ mT, hence, similar g_{iso} , A_N and A_H as found for the GrF/DMF/POBN trapped radical system. Thus, because our system contains only GrF, and the POBN radical adducts can be trapped in DMF but not in benzene, the generation of the carbamoyl-radical trapped by POBN (POBN-DMF radical) as shown in Fig. 6b, is plausible.

The use of other experimental setups, *e.g.*, GrF sonication, as often employed for its fast dispersion in organic solvents, leads to generation of an admixture of radical species due to the fast temperature gradients induced in the solution. Due to the number of other experimental variables to dissect (*e.g.* sonication time, power applied) such study will be devoted for future work. As an example of other EPR signatures detected following the sequence of a specific sonication-time/applied power conditions (30 min, 40 kHz applied power) for the GrF/DMF/POBN mixture is given in the ESI (Fig. S20–S24†). However, the present findings wish to stress the notion that radical defluorination reactions of GrF occurs even under mild conditions in DMF.

3.6 Evidence of *N,N'*-dimethylcarbamoyl fluoride

As can be seen in Fig. 3, *N,N'*-dimethylcarbamoyl fluoride (F-DMF) plays a key role in the proposed mechanism since its formation leads directly to defluorination and also generation of new radical centers. To confirm formation of F-DMF, we performed a series of ^{19}F NMR measurements on FG/DMF samples treated in different ways (see details in ESI†). The first sample was sonicated for 2 hours at room temperature and stirred overnight. The day after, the suspension was still pale-grey, the color of the starting suspension. ^{19}F NMR analysis of the supernatant solution did not show the presence of any F-containing compounds. A second sample of FG/DMF was treated at 120 °C for 5 hours and then left stirring overnight at



Finally, we will briefly comment on recent experimental work by Wang *et al.*,³⁴ in which defluorination of FG in the presence of selected solvents was studied by various instrumental techniques, including XRD, FTIR, TEM, TGA and UV-

3.7 Consequences for the reactivity of FG

We have shown that, depending on the chemical composition of the sample, the presence of radical centers on FG can initiate two important processes. On the one hand, direct reduction of FG (*i.e.*, formation of anionic centers) can occur in the presence of mild reducing agents. On the other hand, a radical mechanism initiated by hydrogen atom transfer can take place provided that species sensitive to homolytic R—H bond dissociation are present in the environment. In both cases, a cascade defluorination process can commence, leading preferentially to compact motifs stabilized by π -conjugation of C=C double bonds. It can be expected that the defluorinated chains of carbons carrying partial positive charge due to surrounding C—F bonds will be highly sensitive to nucleophilic substitutions. In fact, the idea of competitive substitution and reductive defluorination has already been applied to account for experimental observations during the covalent functionalization of graphene based on reactions of FG with various nucleophilic agents, including OH^- ,²³ CN^- ,³¹ Grignard reagents³³ and others. In such reactions, the main indication of concurrent substitution and defluorination was

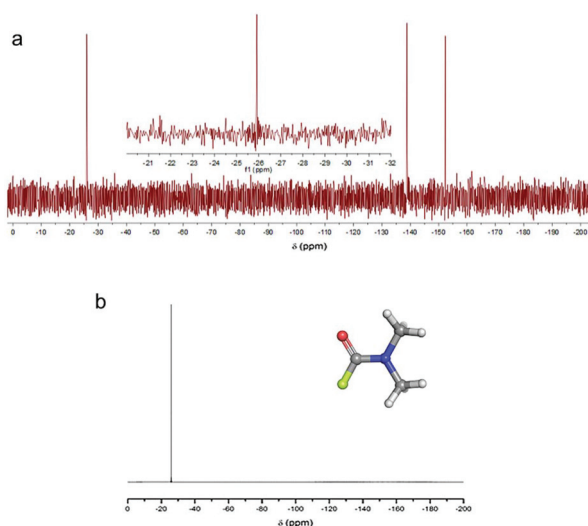


Fig. 8 (a) ^{19}F NMR spectrum (in DMF, 470 MHz) of the supernatant liquid obtained from a FG/DMF suspension heated at 120 °C for 5 hours and stirred at room temperature overnight (in the inset, a close-up of the peak at -26 ppm is shown). (b) ^{19}F NMR spectrum (in DMF, 470 MHz) of the separately prepared compound F-DMF (structure shown in inset). (c) ^{19}F NMR spectrum (in DMF, 470 MHz) of the separately prepared compound F-DMF (structure shown in inset).

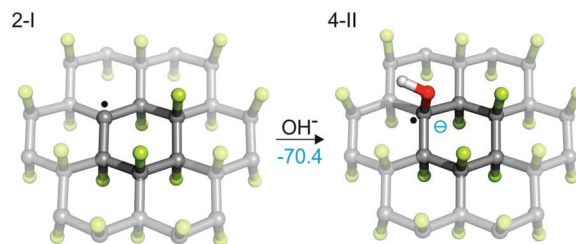


Fig. 9 Direct attack of a nucleophilic species (here OH^-) on a radical site is energetically very favorable (energy in kcal mol^{-1} in the gas phase).

Conclusions

Conflicts of interest

Acknowledgements

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