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**Resolving the Organization of CO₂ Molecules Confined in
Silica Nanopores Using In-Situ Small-Angle Neutron
Scattering and Molecular Dynamics Simulations**

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Environmental Significance Statement

The societal need to transition to a sustainable low carbon energy and resource future requires rapid advancements in technologies to capture, convert, store, and remove CO₂ from our emissions. In this context, the ability to store CO₂ in geologic formations and in engineered porous materials for utilization as needed is gaining increasing attention. One of the less studied but highly important consideration is the organization of confined CO₂ in nanoporous environments. A fundamental understanding of the organization of CO₂ molecules in nanoconfinement can unlock novel pathways for CO₂ separation or CO₂ storage in geologic materials. In this study, we investigate the organization of pressurized CO₂ molecules in silica materials, MCM-41 and SBA-15 with cylindrical pore geometries and pore diameters of 3.3 nm and 6.8 nm, respectively at pressures ranging from vacuum to about 55 bar, using Small Angle Neutron Scattering (SANS) measurements and classical molecular dynamics (MD) simulations. Our results revealed that the nanoconfined CO₂ molecules are organized into core-shell structures with the shell resulting from CO₂ adsorption on the silica surfaces. The shell thickness increased systematically with the applied pressure. We believe this paper is relevant to the field of low carbon energy and resource utilization technologies in the context of achieving a stable and predictable climate on Earth. Thus, this manuscript is a significant contribution to **Environmental Science: Nano**.

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5 **Resolving the Organization of CO₂ Molecules Confined in Silica**
6 **Nanopores Using *In-Situ* Small-Angle Neutron Scattering and**
7 **Molecular Dynamics Simulations**
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10 *Sohaib Mohammed,¹ Meishen Liu,¹ and Greeshma Gadikota^{1,*}*
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15 ¹School of Civil and Environmental Engineering, Cornell University, Ithaca, NY 14853,
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54 * Corresponding Author. Phone: +1 607-255-4796. E-mail: gg464@cornell.edu
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Abstract

Determining the organization of CO₂ molecules confined in nanoporous environments is essential for unlocking our understanding of the fate of CO₂ stored in nanoporous materials. In this study, we investigate the organization of pressurized CO₂ molecules in silica materials, MCM-41 and SBA-15 with cylindrical pore geometries and pore diameters of 3.3 nm and 6.8 nm, respectively at pressures ranging from vacuum to about 55 bar, using *in situ* Small Angle Neutron Scattering (SANS) measurements and classical molecular dynamics (MD) simulations. The nanoconfined CO₂ molecules are organized into core-shell structures with the shell resulting from CO₂ adsorption on the silica surfaces. The shell thicknesses of the adsorbed CO₂ molecules in MCM-41 pores obtained by SANS measurements are 0.7 ± 0.1 Å, 2.1 ± 0.1 Å, 2.2 ± 0.1 Å, 6.7 ± 0.1 Å, 11.5 ± 0.2 Å, and 12.6 ± 0.1 Å at equilibrated pressures of about 1.0, 14.9, 24.9, 34.7, 45.0 and 54.9 bar, respectively. The shell thicknesses of the adsorbed CO₂ molecules in SBA-15 pores are 2.2 ± 0.1 , 2.9 ± 0.1 , 5.1 ± 0.1 , 8.8 ± 0.1 , 12.4 ± 0.1 , and 20.0 ± 0.1 at pressures of 0.9, 15.4, 24.9, 34.9, 45.0 and 55.5 bar, respectively. Close agreement between the experimental and MD simulation results are obtained. MD simulations also suggest that the adsorption of the CO₂ molecules is primarily driven by van der Waals interactions with minor contribution from the electrostatic interactions and hydrogen bonding with the surface hydroxyl groups. These findings inform the development of novel strategies needed to advance low carbon energy and resource recovery and utilization strategies including the storage of CO₂ in natural and engineered materials.

Keywords: CO₂, MCM-41, SBA-15, Small Angle Neutron Scattering (SANS), molecular dynamics (MD) simulations.

1. Introduction

The societal need to transition to a sustainable low carbon energy and resource future requires rapid advancements in technologies to capture, convert, store, and remove CO₂ from our emissions. In this context, the ability to store CO₂ in geologic formations and in engineered porous materials for utilization as needed is gaining increasing attention. One of the less studied but highly important considerations is the organization of confined CO₂ in nanoporous environments. Fluids confined in nanoporous environments have been shown to exhibit anomalous reactive behavior (1-3). Anomalous reactive behaviors, particularly in nanoconfined pores, influences our ability to predict the fate of CO₂ in geologic environments where CO₂ can mineralize to produce water insoluble calcium or magnesium carbonates (4-6). In the context of engineered carbon removal, there is an emerging interest in harnessing nanoporous minerals and materials to store CO₂ as mineralized inorganic carbonates (7-12). However, one of the more fundamental questions that needs to be addressed is the organization of CO₂ in nanoporous media. Delineating the organization of nanoconfined CO₂ is essential to advance calibrated insights into the nano-scale reactivity of confined CO₂.

One of the less explored hypotheses associated with the transport and adsorption of solutes in pressurized CO₂ is the anomalous partitioning and reactivity due to the core-shell organization of pressurized CO₂ in confinement. With increasing interest in storing CO₂ in subsurface environments, predictions of multi-component ion transport and reactivity in these environments is limited by our understanding of the organization of confined CO₂. Further, there is interest in understanding how multiphase fluids are organized in nano-confinement in the context of long-term CO₂ storage, CO₂ utilization for heat mining, and for using pressurized CO₂ to extract organics in multiphase mixtures. To develop robust prediction capabilities associated with the fate

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3 and transport of CO₂ for these applications, the first step is to characterize the organization of
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5 pressurized CO₂ molecules in nanoconfinement. In this study, we investigate the organization of
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7 CO₂ molecules compressed in silica nanopores using Small Angle Neutron Scattering (SANS)
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9 measurements and classical molecular dynamics (MD) simulations.
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12 The chemical interactions of CO₂ with solid interfaces have been investigated extensively
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14 in the context of the separation of CO₂ from multi-gas streams (13-15). CO₂ adsorption and
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16 separation has been extensively investigated using a variety of porous sorbents such as
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18 carbonaceous materials (16-18), zeolites (19-21), alumina (22), silica (23, 24), titania (25),
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20 aerogels (26) and metal organic frameworks (MOFs) (27-29). Mesoporous silica-based sorbents
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22 have been widely utilized to selectively capture CO₂ from gaseous mixtures owing to the ease in
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24 tailoring their porous structure as well as the flexibility in tuning the surface chemistry by attaching
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26 different functional groups. Amine-incorporated mesoporous silica sorbents showed superior CO₂
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28 uptake extents, high selectivity, high stability on adsorption/desorption performance and low
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30 energy requirements (23, 30). The outstanding performance of amine-silica composites stems from
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32 integrating the high CO₂ affinity of amines with large surface area of the mesoporous silica.
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38 A wide range of characterization approaches were used to determine the performance of
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40 various silica-based sorbents in capturing CO₂. Experimental monitoring of CO₂ capture in silica-
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42 based sorbents have been performed using x-ray diffraction (XRD) (30, 31), neutron diffraction
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44 (32), adsorption isotherms (30), thermogravimetric analysis (TGA) (31) and Fourier transform
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46 infrared spectroscopy (FTIR) (33). In addition, different computational methods have been utilized
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48 to obtain microscopic insights into CO₂-silica interfacial properties, including molecular dynamics
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50 (MD) simulations (34-37). The design basis for developing novel sorbents for CO₂ capture has
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52 conventionally been the separation efficiency and the associated energy needs (38-41). However,
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3 uncertainties associated with the molecular-scale organization of CO₂ molecules as a function of
4 confinement and interfacial interactions in the context of CO₂ separation remain less studied.
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8 Rising interest in storing acid gases in minerals motivated the investigation of the
9 organization of acid gases such as CO₂ and H₂S in porous minerals, clay minerals and silica (35,
10 42-47). These studies predicted the organization of acid gas molecules at solid interfaces and in
11 confinement using classical molecular dynamics simulations (48-50). Compressed fluid storage in
12 engineered and natural materials motivated the investigated of methane (51-54) hydrogen (55-57),
13 oxygen (58) and CO₂ (59).
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21 The feasibility of experimentally validating molecular scale models predicting the
22 organization of confined gases has been made possible by advancements in SANS measurements.
23 Spatial resolution of the organization of confined fluids is possible using SANS measurements.
24 For example, Melnichenko and co-workers (60) showed the formation of a dense adsorbed CO₂
25 phase in pores with diameters smaller than 4 nm at temperatures of 22°C, 35°C, and 60°C in porous
26 fractal silica using SANS measurements. The existence of two phases of CO₂ in silica aerogel
27 nanopores at temperature range of 25-35°C was also determined using SANS measurements (61).
28 Holewinski and co-workers (62) linked the CO₂ adsorption to the polymer morphology in
29 poly(ethyleneimine)/mesoporous silica composites using SANS measurements, and showed an
30 enhanced CO₂ capacity and uptake rates by the polymer supported on hydrophilic silica. Despite
31 these advances, the following research questions remain unresolved in literature, and are therefore
32 addressed in this study: (i) How does the organization of CO₂ in nano-confinement in silica
33 nanopores change as a function of pressure? (ii) What is the level of agreement in the predictions
34 and experiments of the organization of confined CO₂? (iii) What are the energetic interactions
35 underlying the organization of confined CO₂ molecules?
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3 To address these research questions, *in situ* SANS measurements of the organization of
4 pressurized CO₂ molecules in mesoporous silica materials such as MCM-41 and SBA-15
5 mesoporous silica with pore diameters of 3.3 nm and 6.8 nm are probed at CO₂ pressures ranging
6 from 0.93 bar to about 55 bars at ambient temperature. These experimental studies are
7 complimented by classical molecular dynamics (MD) simulations. The hypothesis that pressurized
8 CO₂ molecules under confinement form core-shell structures with the shell arising from the
9 adsorption of CO₂ on the pore surface is investigated. The findings from this study have
10 implications for several applications associated with advancing low carbon energy recovery and
11 utilization, and greenhouse gas storage. In the context of storing CO₂ in depleted gas reservoirs,
12 the organization of CO₂ and methane molecules in nanoconfined environments provides the basis
13 for assessing the fate of these molecules and potential displacement of fluids or gases. The rising
14 need to use anthropogenic CO₂ in distributed locations calls for developing novel engineered CO₂
15 storage materials. In this context, determining the organization of CO₂ in confinement and the ease
16 of storing and recovering CO₂ from these materials is essential. The experimental and simulation
17 methodologies described in this study facilitate the core-shell description of compressed confined
18 fluids.

39 **2. Methods**

40 **2.1. Small-angle Neutron Scattering Measurements**

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42 The MCM-41 and SBA-15 mesoporous silica powders with pore diameters of 3.3 nm and
43 6.8 nm, respectively, were purchased from Sigma-Aldrich. The powders were degassed at 100°C
44 for 1.5 hours to remove the adsorbed water. 0.05 g of the degassed powder was then transferred to
45 a special aluminum pressure cell with a diameter of 11 mm and a thickness of 1 mm. The effective
46 thickness (L_{eff}) of the silica sample is less than the nominal thickness of the aluminum cell due to
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3 the presence of interparticle voids. The effective thickness of the silica powder is calculated as
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5 follows:

$$L_{eff} = L \frac{\rho_{bulk}}{\rho_{app}} \quad (1)$$

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14 Where L is the thickness of the aluminum cell, ρ_{bulk} is the bulk density which is the mass of the
15 silica skeleton divided by the skeleton volume, pores volume and the volume of the interparticle
16 voids, and ρ_{app} is the apparent density which is the mass of the silica skeleton divided by the
17 skeleton volume and pores volume (63).

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23 CO_2 gas was injected to the powder and the neutron scattering intensities were collected at
24 various pressures ranging from vacuum to about 55 bars in MCM-41 and SBA-15 samples (see
25 **Figure 1 (a)**). The Small Angle Neutron Scattering (SANS) measurements were performed at the
26 10m SANS instrument at the National Institute of Standards and Technology (NIST), Center of
27 Neutron Research (NCNR). The incident neutron beam has wavelengths (λ) of 5 Å and 12 Å, and
28 the sample-to-detector distances (SSDs) were 1.2 m and 5.2 m to cover a wavevector (Q) range of
29 0.004 to 0.6 Å⁻¹. SANS data were corrected for the transmission, the background and the detector
30 sensitivity (64). The wavevector is $Q = \frac{4\pi}{\lambda} \sin \theta$, where λ is the incident beam wavelength and θ
31 is half of the scattering angle (2θ). Further information on the SANS measurements and data
32 analysis are included in Mohammed and co-workers (51).

33 34 35 36 37 38 39 40 41 42 43 44 45 46 47 48 **2.2 Classical Molecular Dynamics Simulations**

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51 CO_2 and β -cristobalite silica unit cell were built using Avogadro software. The initial
52 configurations of isolated cristobalite unit cell are optimized using density functional theory
53 algorithm implemented in Quantum Espresso software. The kinetic energy cutoff for
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3 wavefunctions and K-points mesh are converged based on the total energy values. The converged
4 energy cutoff and K-points mesh was used to optimize the silica unit cell (see the initial and
5 optimized unit cell coordinates in **Figure S1** and **Table S1** in the supplementary information).
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Ultrasoft pseudopotentials were used in which the generalized gradient approximation (GGA) (65, 66) is selected for the exchange correlation functional. The structural optimization of the silica unit cell has been performed using the Broyden-Fletcher-Goldfarn-Shanno (BFGS) algorithm.

The optimized silica unit cell is replicated in x , y and z directions and MCM-41 and SBA-15 pores are cleaved in surfaces with dimensions of $77.96 \text{ \AA} \times 76.43 \text{ \AA} \times 39.8 \text{ \AA}$ and $124.45 \text{ \AA} \times 152.85 \text{ \AA} \times 49.78 \text{ \AA}$ in x , y , and z directions, respectively. Cylindrical-shaped pores with diameters of 3.32 nm and 6.83 nm were cleaved in the silica surfaces to mimic the MCM-41 and SBA-15 pores used in the SANS measurements (**see Figure 1(b)**). The pore length of MCM-41 and SBA-15 are 3.54 nm and 4.98 nm, respectively. The nonbridging oxygens on the cleaved pore surfaces were protonated to achieve hydroxyl group density of about 8 OH/nm² (67, 68). CO₂ molecules were distributed randomly in the pore space with the required number of molecules to match the pressure used in the SANS measurements. Silica surfaces are modeled using ClayFF forcefield (69). CO₂ molecules were modeled using the forcefield developed by Potoff and Siepmann (70) (**see Table 1 for the modeling parameters**). The constructed simulation cells are set to be periodic in 3 dimensions.

The initial simulation cells were optimized using steepest descent energy minimization algorithm for 50,000 steps to eliminate the inappropriate geometries. NVT ensemble with a time step of 1 fs was performed on the optimized cells for 50 ns at 298 K using the Nose-Hoover thermostat (71, 72). The bonded interactions are accounted for bonds stretching and angle bending

in CO₂. Bond stretching in -OH groups in silica is also accounted for. The van der Waals and electrostatic nonbonded interactions were modeled by 12-6 Lennard-Jones and Coulomb's functions, respectively. The short range interactions are calculated within a cutoff of 14 Å, while long-range electrostatic interactions were modeled using Particle Mesh Ewald (PME) (73). GROMACS 2018 were used to perform and analyze the MD simulations (74). VMD software were used for visualizing the simulation trajectories (75). The simulations were tested for equilibrium based on the potential energy profile of the system. The simulation was considered to reach equilibration once the potential energy profile become uniform and stable.

The structure of the adsorbed CO₂ on MCM-41 and SBA-15 pores was analyzed by calculating the thickness of the adsorbed layers (shell thickness) and the radial distribution function (g_{AB}). g_{AB} describes the density variations with the distance from a reference particle (r) as follow:

$$g_{AB}(r) = \frac{1}{\langle \rho_B \rangle_{local}} \frac{1}{N_A} \sum_{i \in A} \sum_{j \in B} \frac{\delta(r_{ij} - r)}{4\pi r^2} \quad (2)$$

in the expression above, ρ is the density and N_A is the number of A particles in the system.

3. Results and Discussion

3.1. The adsorption of confined CO₂

The pressure dependent SANS intensities of CO₂ loaded MCM-41 and SBA-15 powders are depicted in **Figure 2**. The intensity curves $I(Q)$ describe the influence of pressure on the form factor [$P(\vec{Q})$] and the structure factor [$S(\vec{Q})$] such that:

$$I(Q) = n \left\langle P(\vec{Q}) S(\vec{Q}) \right\rangle \quad (3)$$

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4 where n is a prefactor related to the number density of silica powder in the neutron beam. $P(\vec{Q})$
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6 (high Q region) represents CO₂ loaded mesopores. $S(\vec{Q})$ (low Q region) provides information on
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8 the hexagonally packed cylindrical pores in MCM-41 and SBA-15 matrices.
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12 $S(\vec{Q})$ modeling at $Q < 0.04 \text{ \AA}^{-1}$ suggests that the scattering intensity follows a power law
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14 trend ($I(Q) \sim Q^{-D}$) with slopes of 3.64 ± 0.13 and 3.80 ± 0.20 in MCM-41 and SBA-15, respectively
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16 (**see Figure 2**). The power law slopes of $3 < D < 4$ indicates scattering from surface geometry or
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18 pore size polydispersity (76, 77). As $S(\vec{Q})$ retains a relatively unchanged power law slope values,
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20 it can be inferred that increasing CO₂ pressure has a negligible impact on the morphology of MCM-
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22 41 and SBA-15 matrices. These observations are consistent with the that of confined methane in
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24 porous silica with similar pore sizes (51).
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30 The pressure effect on the form factor, $P(\vec{Q})$ is more significant compared to that on $S(\vec{Q})$
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32 and was monitored by observing the change in the first-order peak behavior. Under vacuum, the
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34 first-order peak intensity of MCM-41 and SBA-15 powders are located at Q of about 0.17 \AA^{-1} and
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36 0.074 \AA^{-1} , respectively (**see Figure 2**). The peak positions suggest that the mesopores in the silica
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38 matrices are aligned in a $p6mm$ hexagonal packing (78) with pore-center to pore-center distances
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40 of 42.93 \AA and 99.11 \AA in MCM-41 and SBA-15 matrices, respectively. These results are in
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42 agreement with previous studies that discussed methane organization in architected porous silica
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44 materials (51, 79, 80). As the pressure increases, the intensity of the first order peaks decreases
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46 considerably in both MCM-41 and SBA-15 mesopores to reach its minimum at the highest applied
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48 pressures (55 bars). Slight decrease in the peak intensity is observed in MCM-41 as the pressure
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50 increase from vacuum to 24.88 bar followed by a significant drop on increasing the pressure to
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52 44.89 bars (**see Figure 3(a)**). However, in SBA-15, slight decrease was noticed as the pressure
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3 increase to 34.91 followed by a substantial decrease as the pressure increased to 55.52 bar (see
4 **Figure 3 (b)**). This decline is mainly attributed to the reduction of the scattering contrast between
5 the pore surface and the adsorbed CO₂ molecules on that surface. A quantitative analysis of the
6 peak intensities as a function of the pressure is obtained by calculating the full width half maximum
7 (FWHM) using Gaussian fitting (see **Figure 3 (c) and (d)**) as follows, where C is the scale factor
8 and σ is the peak width (FWHM/2.345) (81).

$$11 \quad I(Q) = C \exp \left[-\frac{(Q-Q_0)^2}{2\sigma^2} \right] + Background \quad (4)$$

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23 FWHM for MCM-41 and SBA-15 showed a systematic decrease with an increase in the
24 pressure (**Figure 3 (c) and (d)**). The observed profiles of FWHM are resulted from the magnitude
25 of the peak intensity decrease as the pressure increase. Similarly, the normalized peak intensity
26 decreased substantially from 1 to about 0.2 as the pressure increase from 0 to 55 bar, respectively.

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32 As CO₂ molecules are adsorbed on the pore surfaces of MCM-41 and SBA-15, the contrast
33 between the pore surface and the adsorbed gas layer decreases. The scattering length density (SLD)
34 of amorphous silica is $3.21 \times 10^{-6} \text{ \AA}^{-2}$ at a mass density of 2.20 g cm^{-3} (82). SLD of the adsorbed
35 CO₂ in \AA^{-2} units, is proportional to the density of the adsorbed layer on the pore surface such that:

$$36 \quad \rho_{SLD} = (2.49 \times \rho_{CO_2}) \times 10^{-6} \quad (5)$$

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42 where ρ is the mass density of the adsorbed gas (63). The maximum scattering contrast is under
43 vacuum due to the difference between the SLD of the amorphous silica matrices and that of the
44 vacuum inside the pore, which is 0. This contrast decreases as CO₂ molecules adsorb on the pore
45 surface and develop a finite SLD that reduces the density difference ($\Delta\rho$) with the silica matrix.

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The density of the adsorbed CO₂ increases as the pressure increase which contributes to a

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3 significant decrease in the scattering contrast. These observations are consistent with previous
4 studies after accounting for the differences in the adsorbate chemistry (82, 83). It is interesting to
5 note that the decrease in the amplitude of scattering intensity of MCM-41 and SBA-15 is more
6 significant when CO₂ adsorbs on their pore surface compared to CD₄ (51).
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12 The neutron scattering intensity of binary systems, such as CO₂-silica system, is
13 proportional to the difference of the SLD of the scattering object and the surrounding medium
14 (63). Thus, the contrast matching point, the point at which the contrast between the silica matrix
15 and the adsorbed CO₂ becomes zero (i.e., $I(Q) = 0$) (83, 84). can be estimated by plotting the
16 square root of the maximum intensity [$I(Q_{max})^{1/2}$] as a function of the applied pressure (**see Figure**
17 **4**). It is interesting to observe that $I(Q_{max})^{1/2}$ decreases linearly with pressure in MCM-41 while
18 exponential decay is noted in SBA-15 pores. The pressure at which the contrast matching point
19 occurs in MCM-41 is about 110 bar, which corresponds to a density of 0.83 g/cm³. In SBA-15, the
20 contrast matching point occurs at a pressure of about 72 bar which corresponds to 0.75 g/cm³. The
21 lower pressure required to reach the contrast matching point in SBA-15 can be attributed to the
22 larger pore size that results in higher CO₂ partitioning into the pore space compared to MCM-41.
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38 The spatial distribution of the adsorbed CO₂ molecules on the pore surfaces is determined
39 by calculating the 2D density maps over the axis parallel to the pore length (z-axis) using MD
40 simulations (**Figure 5 and 6**). Anisotropic distribution of confined CO₂ is observed in both MCM-
41 41 and SBA-15 pore spaces, with higher densities noted at the pore surfaces compared to the center
42 of the pore which is attributed to confinement effect (35, 85-87). At pressures under 15 bars, almost
43 all the confined molecules were adsorbed on the MCM-41 and SBA-15 pore surfaces. Traces of
44 CO₂ molecules are distributed unevenly outside the adsorbed layers and closer to the pore center.
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3 of the pore were systematically increased as the pressure increased from about 1 bar to about 55
4 bar. The anisotropic distribution of CO₂ molecules in MCM-41 and SBA-15 pores are also evident
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6 from the snapshots taken from the equilibrated trajectories during the last 1 ns of the simulation
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8 time (see **Figure S2**).
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10 11 12 **3.2. Core-shell structure of confined CO₂** 13

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15 The experimental SANS intensities of the confined CO₂ in MCM-41 and SBA-15 are
16 modeled using a core-shell cylindrical geometry with multilevel densities to probe the pressure
17 dependent shell thickness and core radius (79, 80, 88, 89). The shell is defined as the thickness of
18 the adsorbed CO₂ layer on the pore surface while the core is the diameter of the silica pore
19 excluding the CO₂ shell. These data are then compared with the MD simulation data to determine
20 the density of the confined fluids (see **Figure 7 and Table S2**). The shell thickness are obtained
21 from SANS measurements by utilizing the model developed by Chiang and co-workers (79, 80),
22 while those from MD simulations are directly calculated from the 2D density maps. The model
23 developed by Chiang and co-workers involves the shell thickness of the adsorbed gas, the core
24 radius and the scattering length densities of the adsorbed gas and silica matrix (see **Figure S3**).
25 This model has been applied to predict the adsorption extent of CD₄ in MCM-41 and SBA-15 with
26 cylindrical pore geometries. The shell boundaries in MD density maps are the pore surface and the
27 point where the density of the adsorbed CO₂ start to decline steeply (51).
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44 The shell thickness of the adsorbed CO₂ molecules on MCM-41 and SBA-15 pore surfaces
45 increased systematically with the pressure as indicated by SANS measurements and MD
46 simulations (see **Figure 7**). The experimentally determined shell thicknesses of adsorbed CO₂ on
47 MCM-41 using SANS measurements are $0.7 \pm 0.1 \text{ \AA}$, $2.1 \pm 0.1 \text{ \AA}$, $2.2 \pm 0.1 \text{ \AA}$, $6.7 \pm 0.1 \text{ \AA}$, 11.5
48 $\pm 0.2 \text{ \AA}$, and $12.6 \pm 0.1 \text{ \AA}$ at pressures of about 1.0, 14.9, 24.9, 34.7, 45.0 and 54.9 bar, respectively.
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3 The thickness of the adsorbed CO₂ layer on SBA-15 pore surfaces are 1.9 ± 0.1 , 2.6 ± 0.1 , $4.9 \pm$
4 0.1 , 8.3 ± 0.1 , 12.4 ± 0.1 , and 20.0 ± 0.1 at pressures of about 0.9, 15.4, 24.9, 34.9, 45.0 and 55.5
5 bar, respectively. The thickness of the adsorbed CO₂ molecules, also referred to as “shell”
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7 determined experimentally are in close agreement with that of classical molecular dynamics (MD)
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9 simulations (see **Figure 7 (a)**).

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15 The increase in the shell thickness of the adsorbed CO₂ on MCM-41 and SBA-15 surfaces
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17 are more significant than that of CD₄ adsorption even at higher pressures (51). The increase in the
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19 shell thickness is associated with a considerable decrease in the core radius in MCM-41 and SBA-
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21 15 pores (see **Figure 7 (b)**). The core radius in MCM-41 decreased from 16.1 ± 0.2 Å to 4.9 ± 0.1
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23 Å as the pressure increases from 1.0 to 55 bar, respectively, while the core radius in SBA-15
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25 decreased from 31.7 ± 0.3 Å to 9.9 ± 0.1 Å, respectively.

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29 The radial distribution function (RDF) and the associated coordination number (CN) of
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31 carbon atom in CO₂ molecules from the hydroxyl group (-OH) on the pore surface are calculated
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33 to probe the influence of the applied pressure on the organization and the structure of the adsorbed
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35 CO₂ layers on MCM-41 and SBA-15 pore surfaces (see **Figure 8**). The RDF peaks in MCM-41
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37 and SBA-15 decreased systematically with an increase in the pressure, suggesting more substantial
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39 increase in CO₂ density beyond the adsorbed layer. Further, RDF peaks in MCM-41 is higher than
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41 that in SBA-15 at the same pressure which can be attributed to the higher confinement effect in
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43 the narrower pores.
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47 Although RDF peaks intensities decrease with an increase in pressure, the corresponding
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49 CN profiles showed a substantial increase in both MCM-41 and SBA-15 which imply an increase
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51 in the density of carbon atoms in the first coordination shell of -OH group and the adjacent
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53 interfacial layers. The number of carbon atoms in the first coordination shell corresponding to the
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3 pressure is calculated by integrating the coordination number at the end of the RDF peak and
4 depicted in **Figure 9**. The number of carbon atoms in the first coordination shell of -OH on MCM-
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6 41 increased from 0.5 ± 0.1 to 2.8 ± 0.1 as the pressure increased from 1.0 bar to 54.9 bar,
7
8 respectively. Sharper increase in the number of C atoms in the first coordination shell of OH is
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10 noticed in MCM-41 pores as the pressure increase from about 0.98 to 14.92 bars compared to the
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12 change in higher pressures. This trend is attributed to the relatively higher pressure increment (~
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14 14 bar) compared to the following pressure change values (~ 10 bar). In SBA-15, the number of
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16 carbon atoms increased from 0.3 ± 0.1 to 1.3 ± 0.1 as the pressure increased from 0.9 bar to 55.5
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18 bar, respectively. The number of CO₂ carbon atoms in the first coordination shell of OH groups on
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20 MCM-41 and SBA-15 are significantly higher than that of CH₄ carbon atoms due to the higher
21
22 tendency to adsorb CO₂ on silica surfaces, which is in agreement with previous studies (86, 90,
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24 91).

3.3. Energetics of CO₂ adsorption in MCM-41 and SBA-15

33 Further insights on the adsorption and organization of adsorbed CO₂ molecules on MCM-
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35 41 and SBA-15 pore surfaces are obtained by calculating the intermolecular interactions. Here, the
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37 energetics of CO₂ adsorption are obtained by calculating the magnitude of van der Waals and
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39 electrostatic potentials as a function of the applied pressure (**see Figure 10**). Van der Waals and
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41 electrostatic interactions are normalized by the number of -OH group in the MCM-41 and SBA-
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43 15 pore surface, which are 128 and 256, respectively, for a consistent comparison. Both van der
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45 Waals and electrostatic interactions are significantly enhanced with increase in pressure, which
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47 corresponds to higher extents of CO₂ adsorption. In this context, the van der Waals interactions in
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49 MCM-41 are enhanced from -0.8 ± 0.1 to -5.8 ± 0.1 kJ/mol as the pressure increases from 1.0 to
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51 54.9 bar. Similarly, as the pressure increased from 0.9 to 55.9 bar in SBA 15, the favorable van
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3 der Waals interactions are enhanced from -1.2 ± 0.1 to -5.8 ± 0.1 kJ/mol. Similarly, the electrostatic
4 interactions in MCM-41 are enhanced from -0.13 ± 0.01 to -1.3 ± 0.1 kJ/mol as the pressure
5 increases from 1.0 to 54.9 bar. In SBA-15, the electrostatic interactions are enhanced from $-0.4 \pm$
6 0.03 to -2.8 ± 0.05 kJ/mol as the pressure increased from 0.9 to 56.0 bar, respectively.
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12 The intermolecular interactions suggest that CO₂ adsorption on silica surfaces is primarily
13 driven by Lennard-Jones interactions, arising from the atomistic collisions, with smaller
14 contributions from the electrostatic interactions. It is interesting to note that the electrostatic
15 contributions are more significant in the case of CO₂ molecules confined in silica nanopores
16 compared to CH₄ molecules in confinement (51). In additions to the contributions of van der Waals
17 and electrostatic interactions on CO₂ adsorption behavior in MCM-41 and SBA-15 pores, CO₂
18 molecules form hydrogen bonds with the OH groups on the pore surface that contribute positively
19 to the overall intermolecular interactions. The energy of the individual hydrogen bond can vary
20 from 1 to 40 kcal/mol (92). The number of hydrogen bonds in SBA-15 is higher than that in MCM-
21 41 due to the higher number of OH groups on the pore surface and the higher number of CO₂
22 molecules in confinement. The total number of hydrogen bonds between OH group on the pore
23 surface and the confined CO₂ increase with the pressure in both MCM-41 and SBA-15 pores (**see**
24 **Table S3**) due to the higher adsorption extent of CO₂ molecules on the pore surface. Further, the
25 intermolecular interactions between CO₂ molecules confined in MCM-41 and SBA-15 are
26 dominated by van der Waals attractions with a considerable contribution from electrostatic
27 attractions (**see Table S4**). The combination of van der Waals, electrostatic and hydrogen bonds
28 contribute to higher energies of adsorption of CO₂ on MCM-41 and SBA-15 surfaces compared to
29 CH₄ (51).
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4. Conclusions

The structure of confined CO₂ in mesoporous silica, MCM-41 with pore diameter of 3.3 nm and SBA-15 with 6.8 nm, was studied as a function of pressures applied that range from vacuum to about 55 bars using *in situ* small angle neutron scattering and classical molecular dynamics simulations. SANS measurements showed that higher levels of CO₂ are adsorbed with increasing pressure of CO₂. Core-shell organization of the CO₂ molecules in confinement are noted with the adsorbed CO₂ molecules forming a shell-like structure and less dense CO₂ molecules in the center of the pore. The shell thickness increases systematically with pressure, while the corresponding core radius decreases accordingly. The core-shell structures of CO₂ predicted from MD simulations are in close agreement with the experimental SANS data. Our calculations using MD simulations show that CO₂ adsorption is primarily driven by van der Waals interactions with electrostatic interactions having a smaller contribution. These studies demonstrate that the structure of confined CO₂ molecules can be experimentally determined, and these experimental results can be used to validate modeling efforts. The experimental and modeling approaches discussed in this study can be used to determine the organization of confined CO₂ molecules in novel functionalized materials and in natural geologic minerals.

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Table 1. Lennard-Jones parameters and the charges assigned to the silica and CO₂ atoms in the molecular dynamics simulations.

| Atom | σ (nm) | ϵ (kJ/mol) | Charge (q) |
|-----------------------|---------------|------------------------|------------|
| Silica | | | |
| Si | 0.3706 | 7.700×10^{-6} | + 2.1000 |
| O _{bridging} | 0.3553 | 0.6502 | - 1.0500 |
| O _{hydroxyl} | 0.3353 | 0.6502 | - 0.9500 |
| H _{hydroxyl} | 0.0000 | 0.0000 | + 0.4250 |
| CO₂ | | | |
| C | 0.2800 | 0.2245 | + 0.7000 |
| O | 0.3050 | 0.6569 | - 0.3500 |

Captions for Figures

Figure 1. Schematic representation of the experimental setup for small-angle neutron scattering (SANS) measurements for determining the organization of confined CO₂ molecules in the pores of MCM-41 (pore diameter of 3.3 nm) and SBA-15 (pore diameter of 6.8 nm) silica materials is shown in (a). Snapshots of the simulated initial configurations showing CO₂ molecules and the simulated SBA-15 (top), and MCM-41 (bottom) silica pores are shown in (b). The atoms of silica pores and CO₂ molecules (see the color coding for the atoms below (b)) are drawn using VDW drawing method implemented in VMD software.

Figure 2. The small-angle neutron scattering intensities of CO₂-loaded (a) MCM-41 and (b) SBA-15 pores as a function of the wavevector and the applied pressure are shown. The insets represent the power law slope values as a function of the applied pressure.

Figure 3. The fitted first-order peak intensity of MCM-41 and SBA-15 as a function of the applied pressure are represented in (a) and (b), respectively. The full width half maximum (FWHM) and the peak intensity normalized by the intensity of the pore under vacuum for MCM-41 and SBA-15 are shown in (c) and (d), respectively.

Figure 4. The amplitudes of the SANS intensity [$I(Q_{max})^{1/2}$] at the maximum of the first-order peak as a function of the applied pressure in MCM-41 and SBA-15 pores are shown. The pressures at which the intensity reaches zero average contrast in MCM-41 and SBA-15 are about 110 bar and 72 bar, respectively.

Figure 5. The 2D density maps of the CO₂ molecules adsorbed on the pore walls of MCM-41 as a function of the pressure obtained from the classical molecular dynamics simulations are shown. The pore diameter of MCM-41 is 3.3 nm. The density maps are averaged over the last 10 ns of the simulation time.

Figure 6. The 2D density maps of the CO₂ molecules adsorbed on the pore walls of SBA-15 as a function of the pressure obtained from the classical molecular dynamics simulations are shown. The pore diameter of SBA-15 is 6.8 nm. The density maps are averaged over the last 10 ns of the simulation time.

Figure 7. The thickness of the adsorbed CO₂ layer determined from the small-angle neutron scattering measurements and molecular dynamics simulations (a) and the corresponding core radius (b) as a function of the applied pressure are shown. The error bars of the SANS and MD data represent standard deviations based on fittings and simulations performed in triplicates.

Figure 8. The radial distribution functions and the corresponding coordination numbers of -OH-C_{CO₂} in the MCM-41 (a) and SBA-15 (b) as a function of the applied pressure are shown. The simulation data are averaged over the last 10 ns of the simulation time.

Figure 9. The number of carbon atoms in the first coordination shells of the -OH group on the MCM-41 and SBA-15 pore surfaces obtained from the RDFs are shown. The error bars represent the standard deviations from simulations performed in triplicates.

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2
3 **Figure 10.** The electrostatic and van der Waals interaction energy associated with the adsorption
4 of CO₂ molecules on the surfaces of MCM-41 and SBA-15 materials as a function of pressure are
5 shown. These analyses are performed on data averaged over the last 10 ns of the simulation time.
6 The error bars are calculated based on the simulations performed in triplicates.
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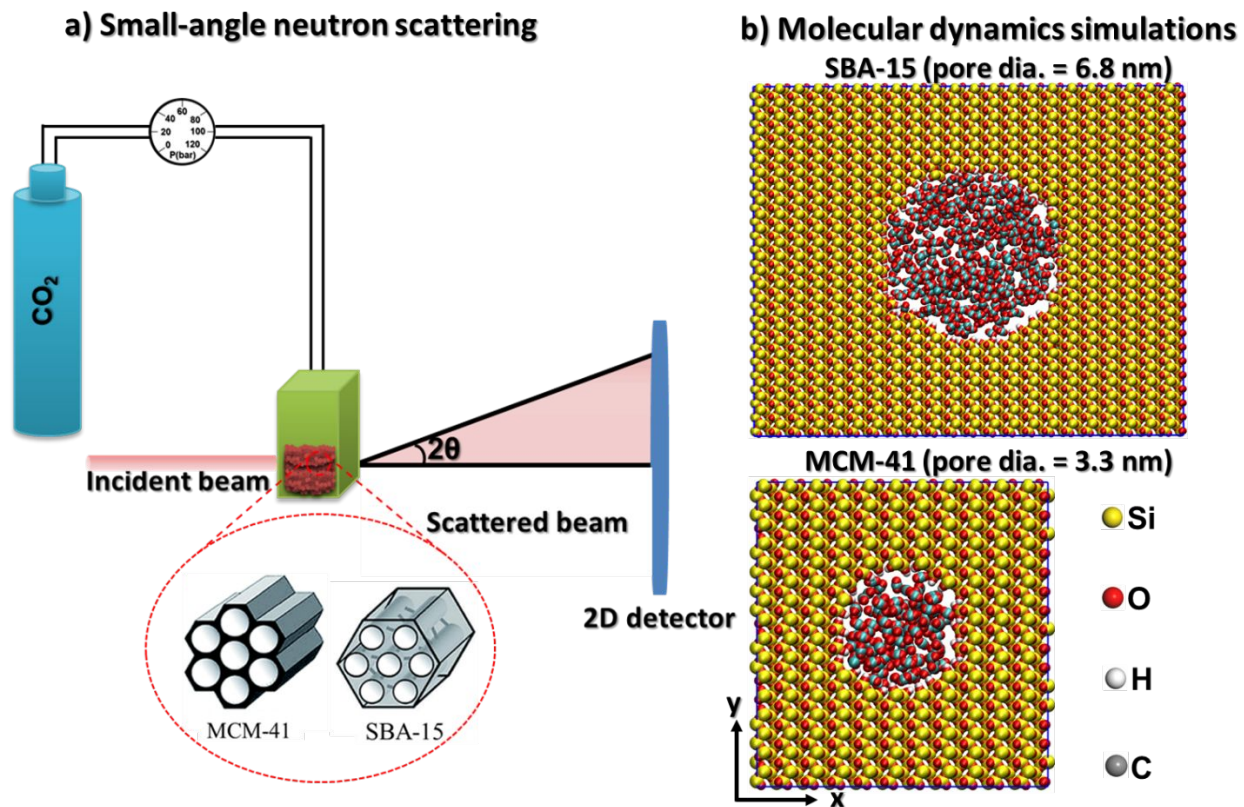


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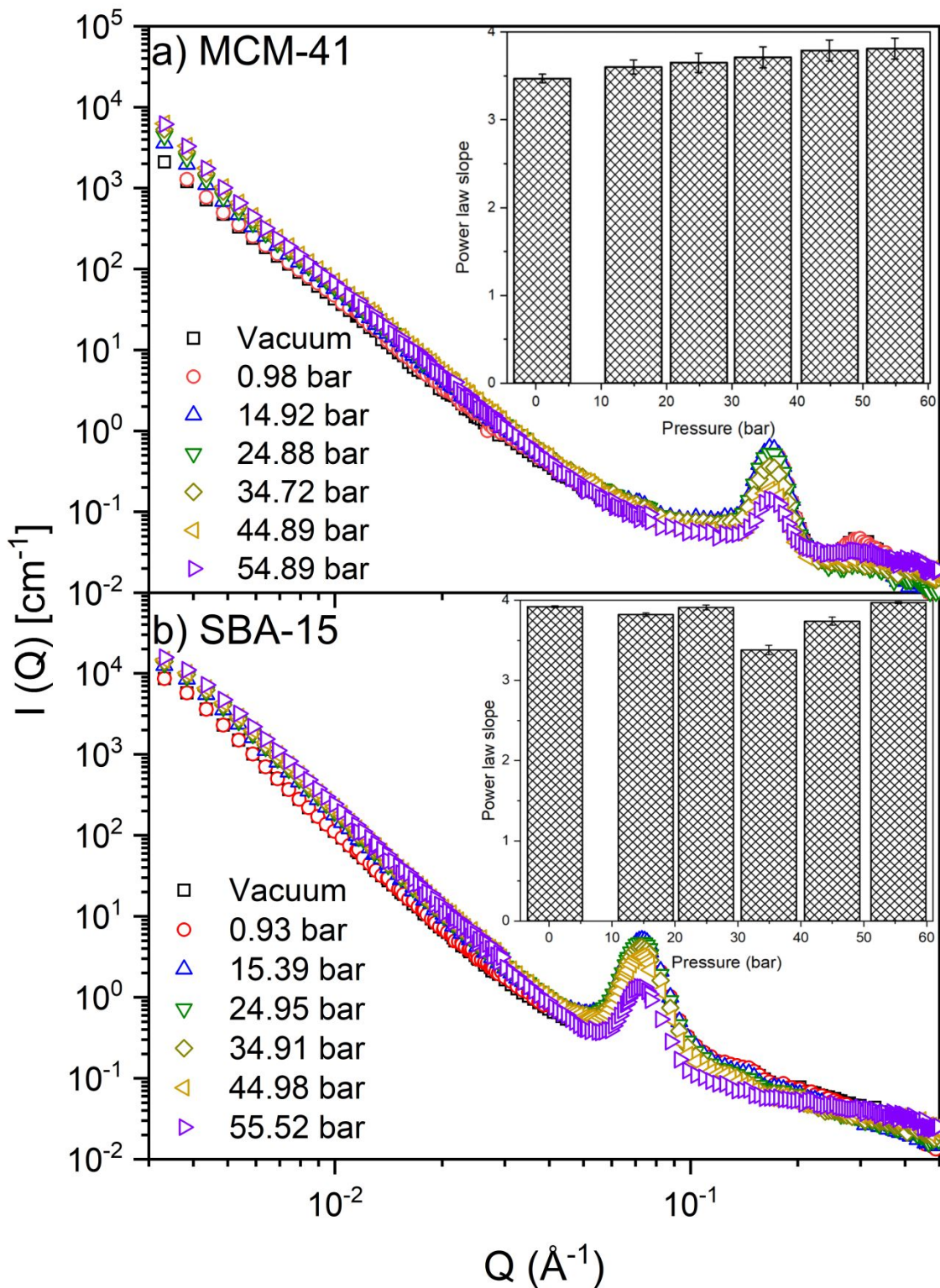


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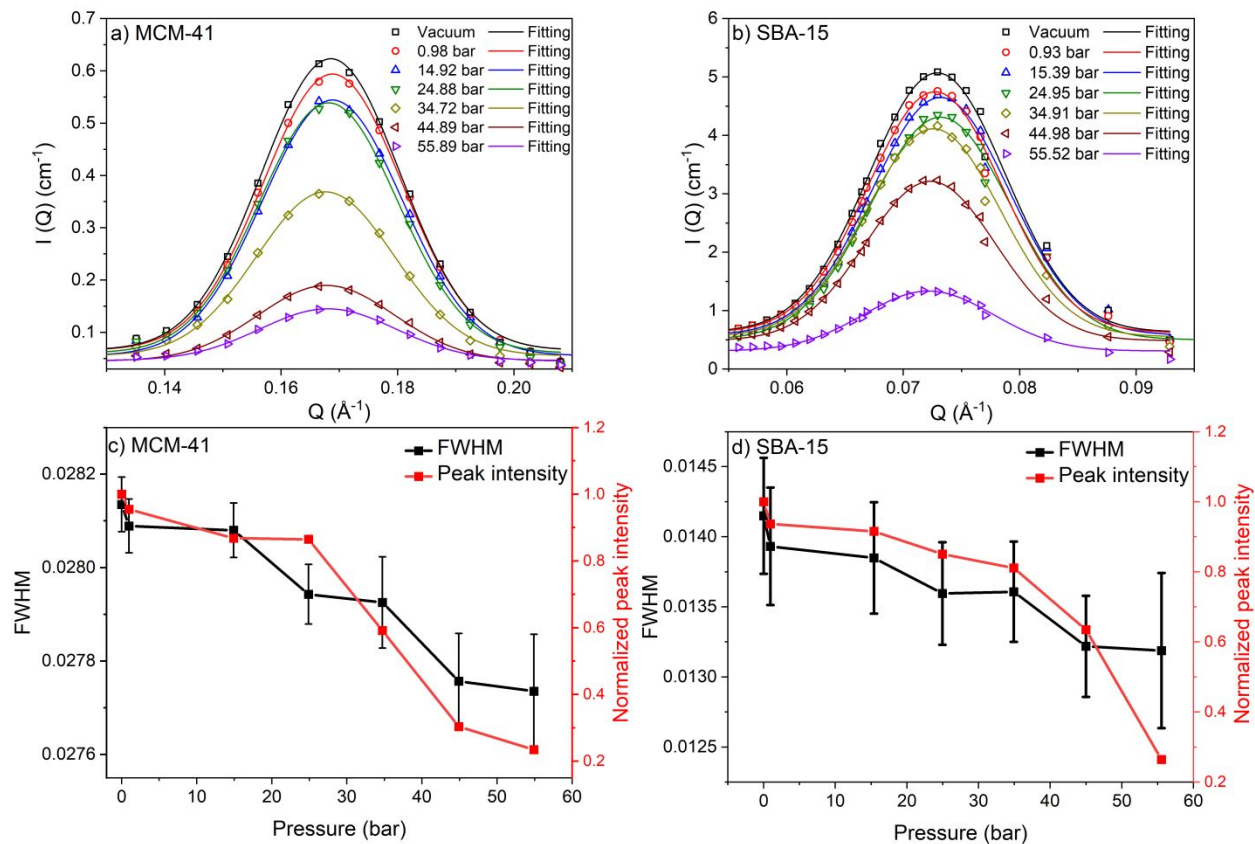


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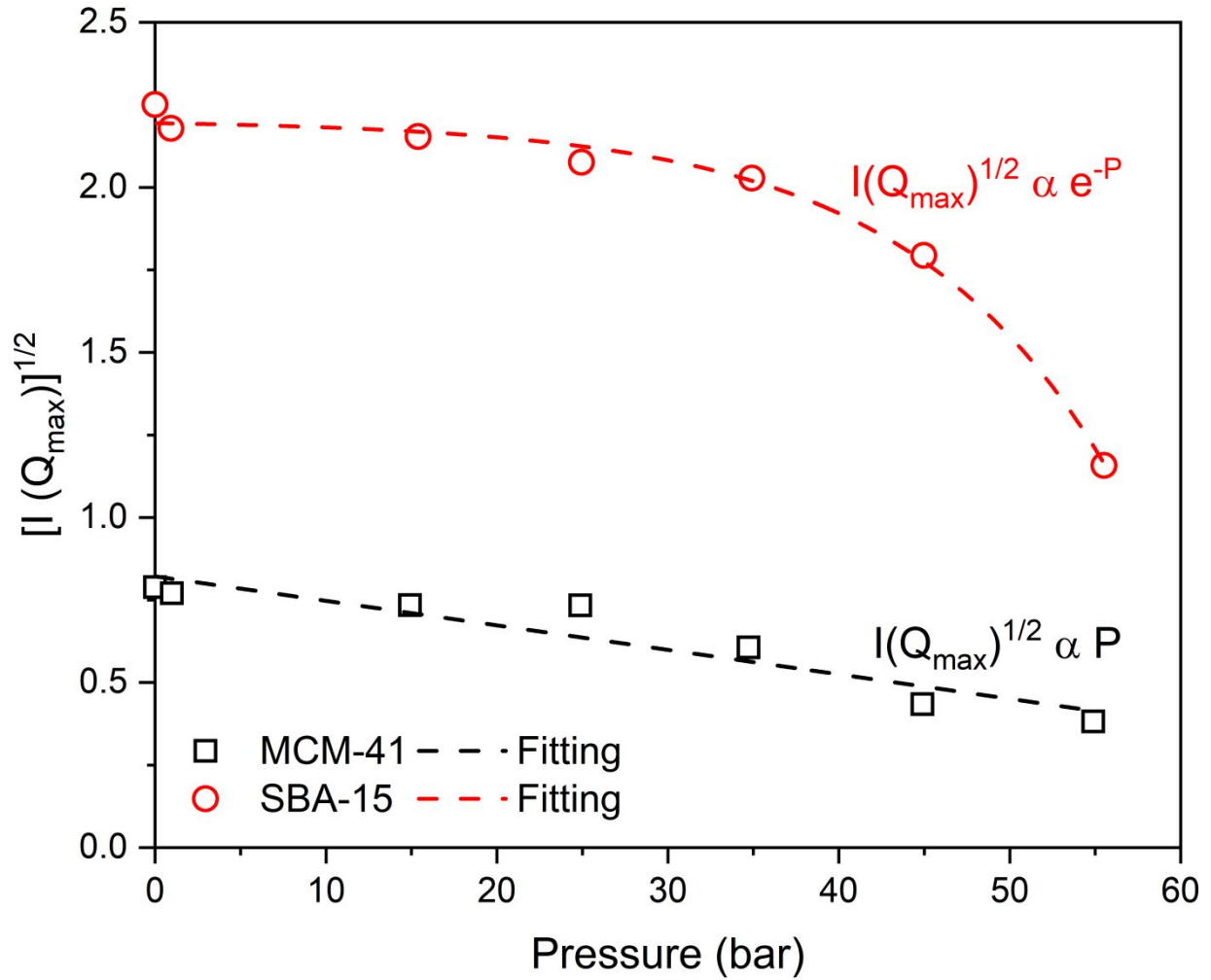


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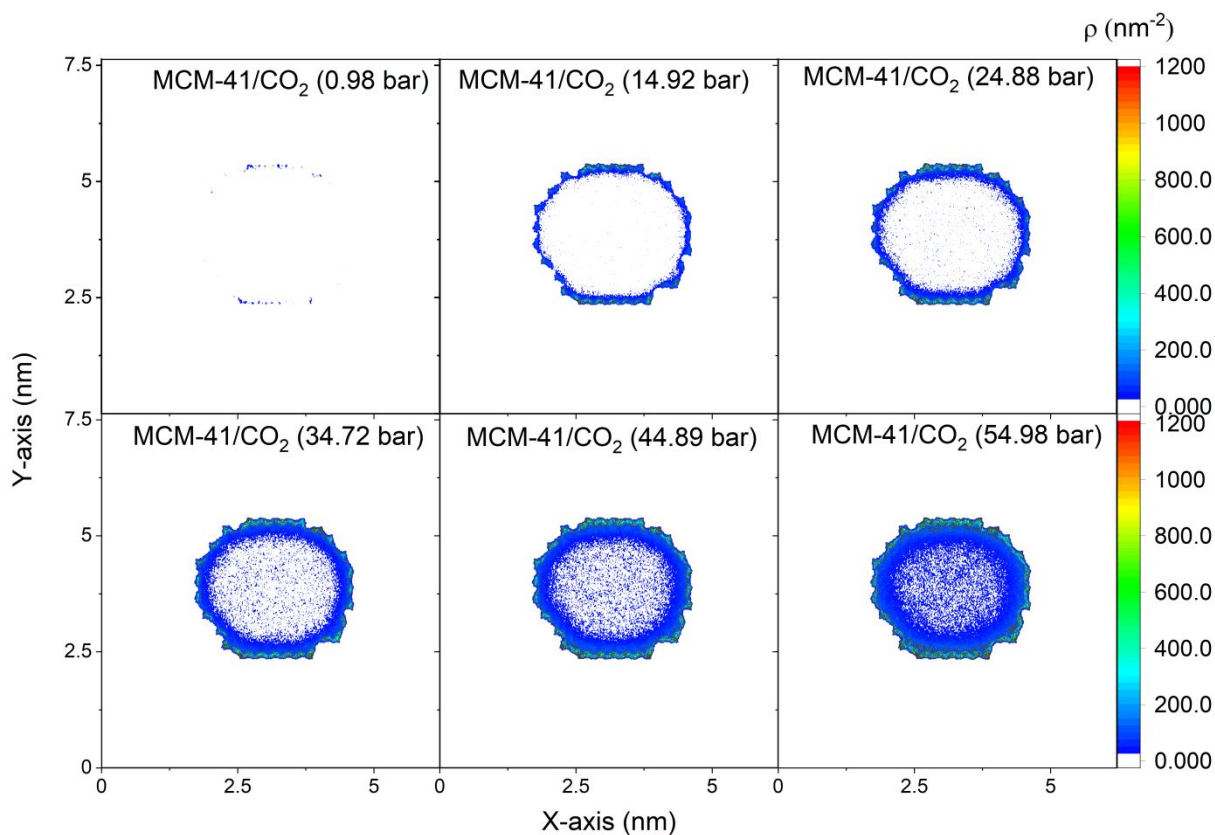


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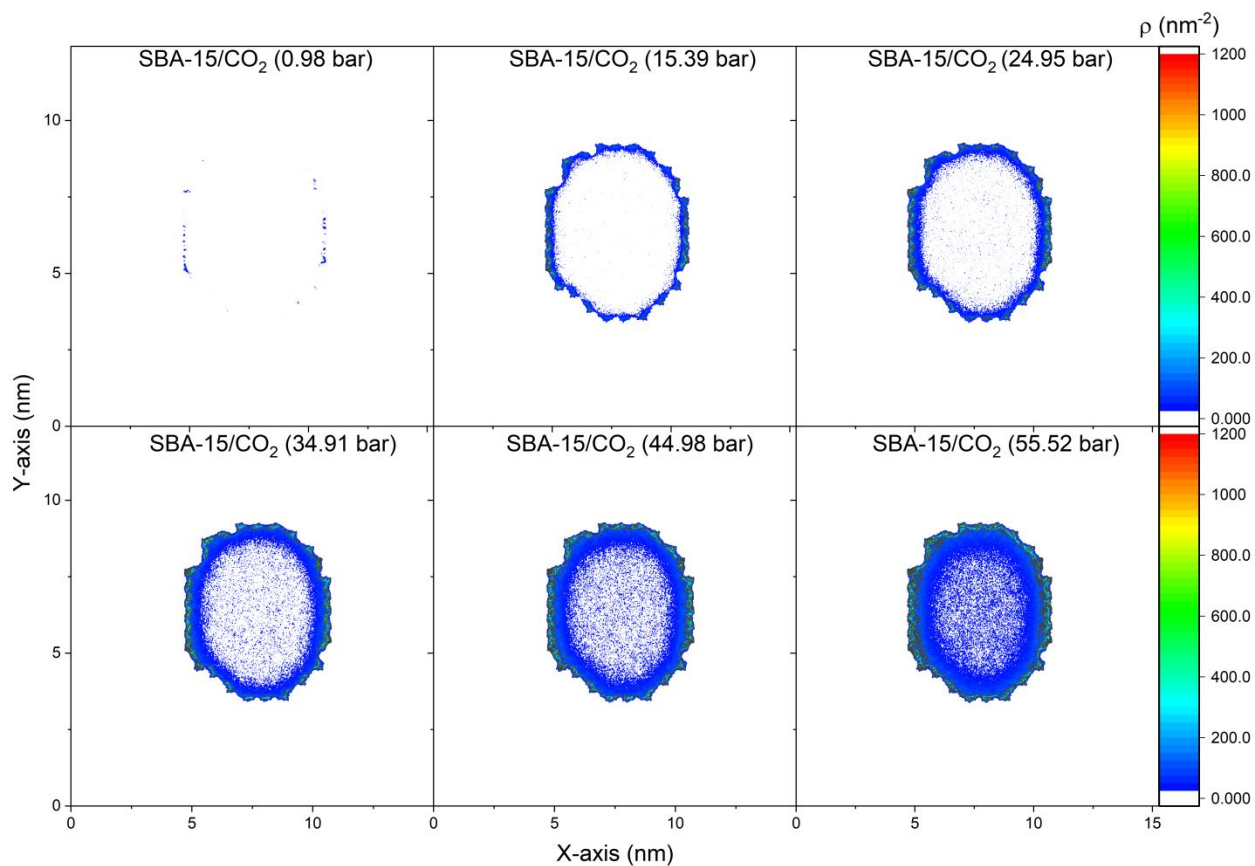


Figure 6. The 2D density maps of the CO₂ molecules adsorbed on the pore walls of SBA-15 as a function of the pressure obtained from the classical molecular dynamics simulations are shown. The pore diameter of SBA-15 is 6.8 nm. The density maps are averaged over the last 10 ns of the simulation time.

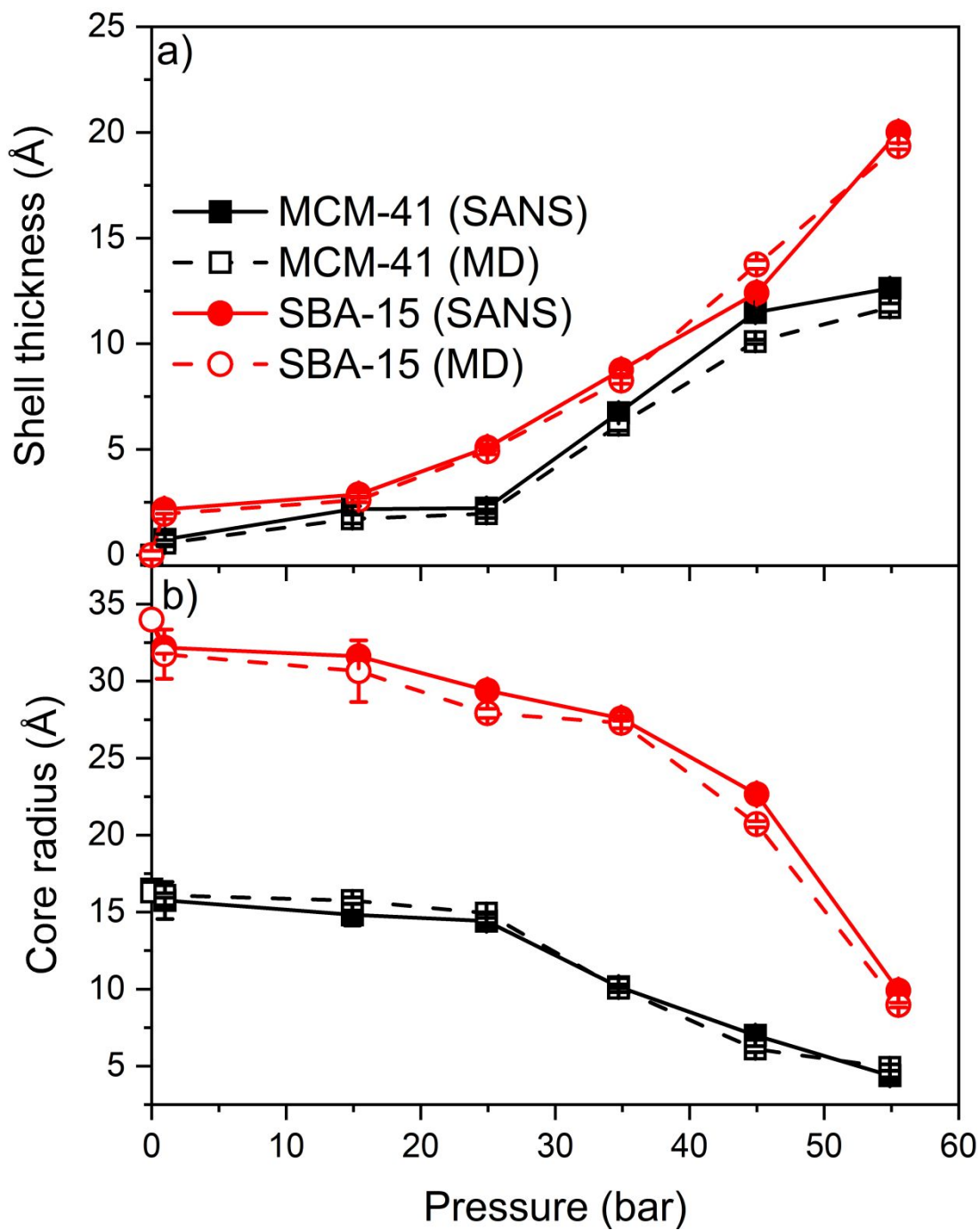


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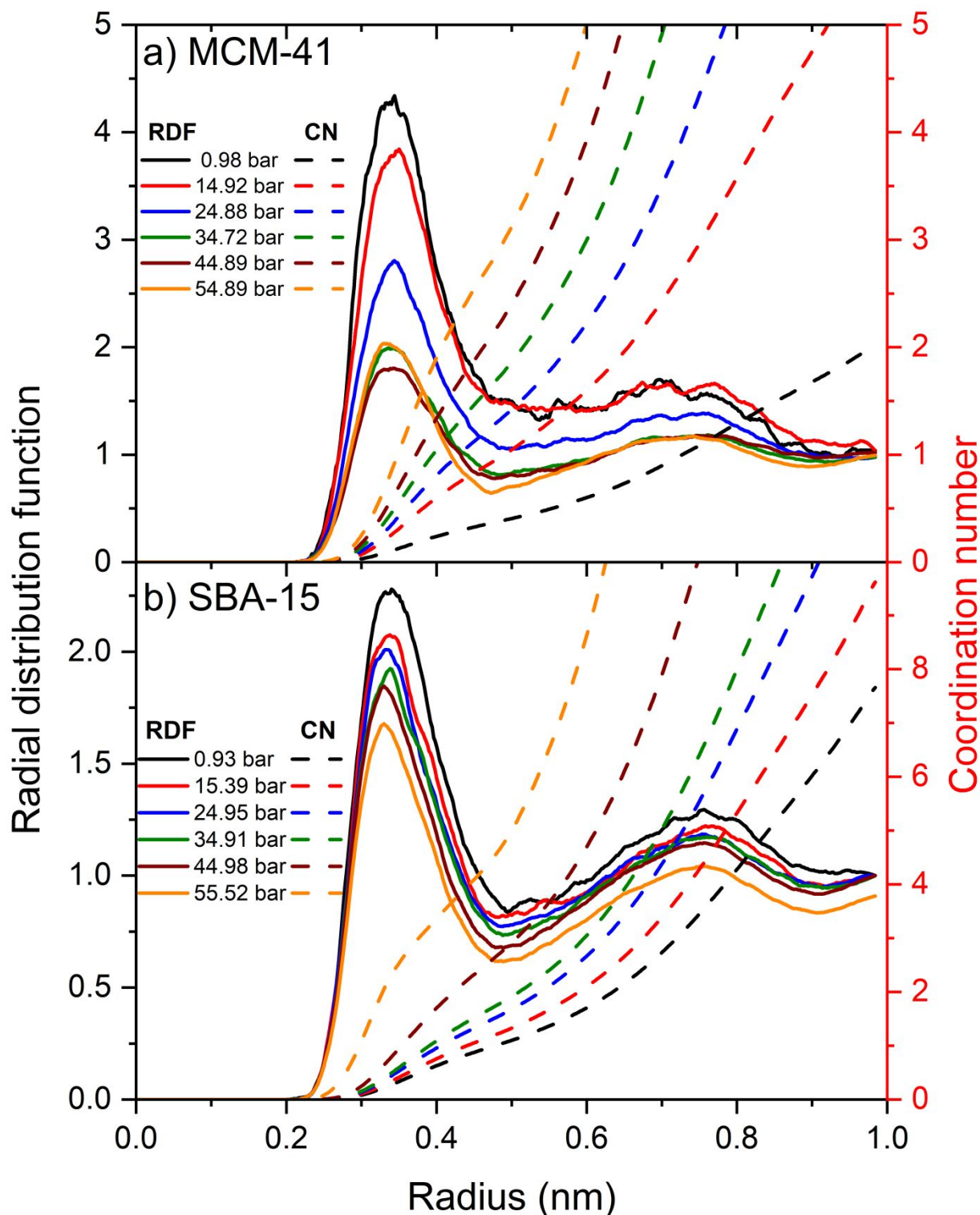


Figure 8. The radial distribution functions and the corresponding coordination numbers of -OH-CO_2 in the MCM-41 (a) and SBA-15 (b) as a function of the applied pressure are shown. The simulation data are averaged over the last 10 ns of the simulation time.

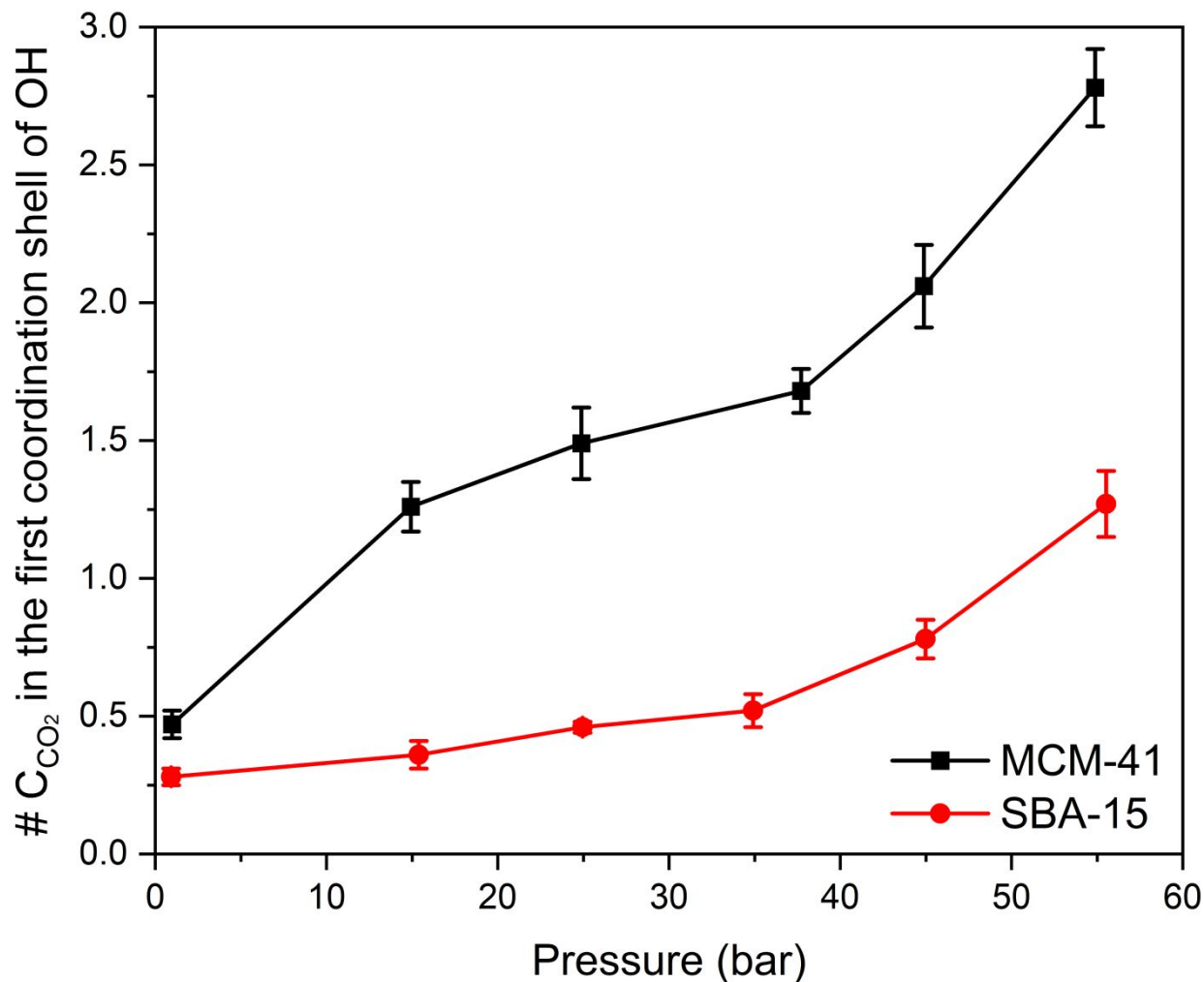


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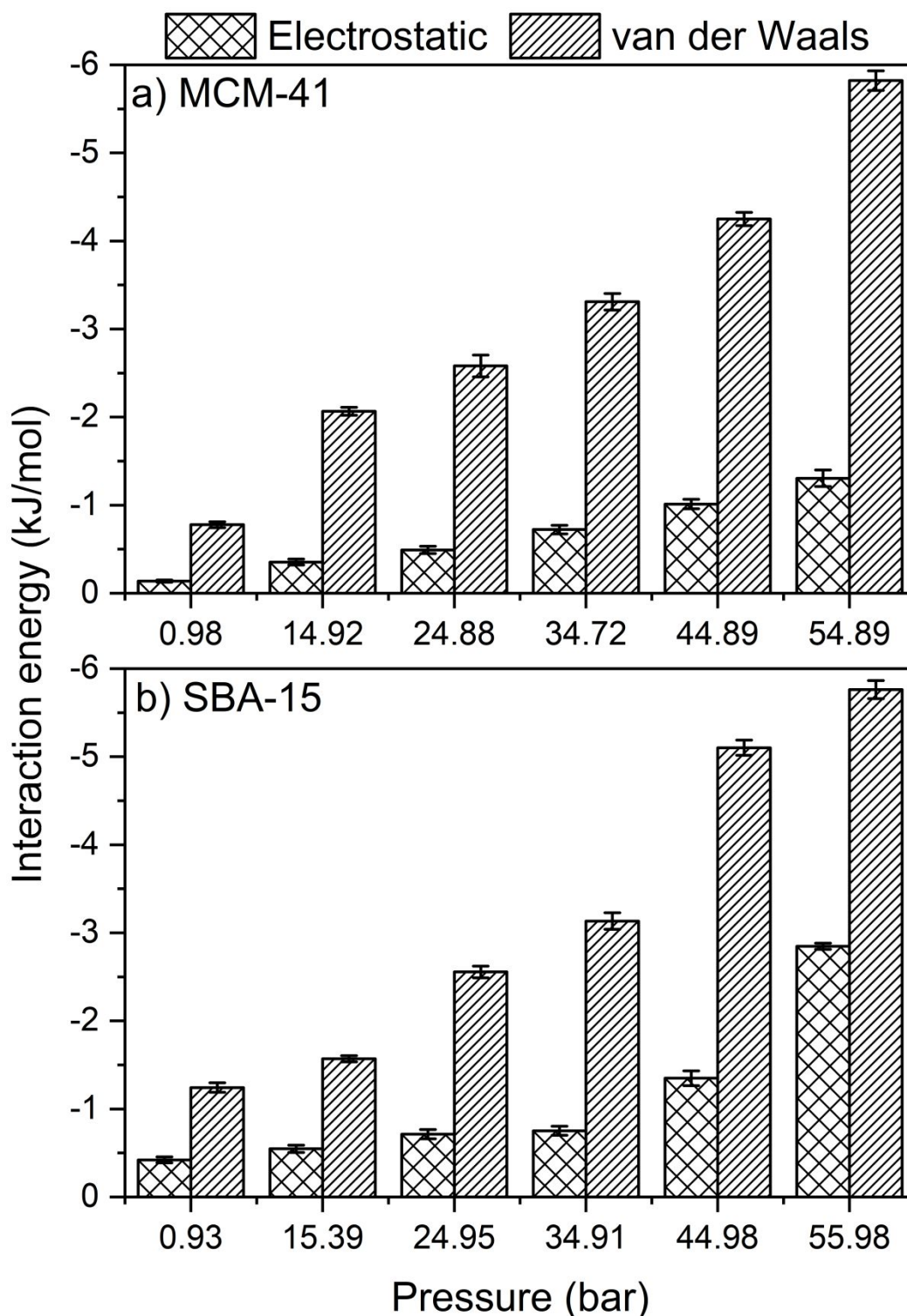


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