




Cite this: *Soft Matter*, 2025, 21, 6383

CO₂ transport behavior in poly(fluoroalkyl acrylate) and poly(fluoroalkyl methacrylate): a comparative study of fluoropolymer structure–property relationships†

Sinan Feng,^a Yucheng Zhang,^a Nobuyuki Otozawa,^b Shinichi Murata^a and Atsushi Takahara *^a

This study examines the CO₂ transport behavior of poly(fluoroalkyl acrylate) (PFA) and poly(fluoroalkyl methacrylate) (PFMA) thin films. Using time-resolved attenuated total reflectance Fourier-transform infrared spectroscopy (ATR-FTIR) and a quartz crystal microbalance (QCM), we quantified CO₂ diffusivity and solubility, linking these properties to the polymer side-chain architecture. The results demonstrate that PFMA exhibits lower CO₂ permeability than PFA at comparable side-chain lengths, owing to restricted chain mobility caused by the α -methyl backbone of PFMA. Additionally, longer fluorinated side chains increase CO₂ diffusivity while simultaneously reducing solubility owing to weaker polar interactions with CO₂. Overall, the CO₂ permeability of PFA surpasses that of the PFMA series because of its higher diffusivity. These findings highlight the intricate balance between diffusivity and solubility governed by the molecular structure.

Received 23rd April 2025,
Accepted 23rd June 2025

DOI: 10.1039/d5sm00413f

rsc.li/soft-matter-journal

Introduction

Excessive CO₂ emissions into the atmosphere are a primary driver of global climate change, necessitating urgent strategies to mitigate greenhouse gas emissions. Among these strategies, direct air capture (DAC) has emerged as a promising technology for capturing CO₂ directly from the atmosphere.¹ DAC offers significant advantages, including the ability to capture diffuse CO₂ emissions and retrofit capture units in decentralized locations.² Membrane-based separation technologies have gained attention as a cost-effective and energy-efficient alternative for CO₂ capture.³ These technologies rely on selective materials that facilitate CO₂ transport while minimizing the passage of other gases. Polymers represent a significant class of materials for such applications owing to their tunable properties and ease of fabrication. However, the performance of membrane-based systems depends critically on the interplay between gas permeability and selectivity, which is governed by the chemical structure and physical properties of the polymer.

Fluorinated polymers represent a significant class of high-performance materials, valued for their unique combination of thermal stability, chemical inertness, low surface energy, and tunable optical properties. Polyfluoroalkyl acrylates represent a unique class of functional polymers that combine the advantageous properties of fluorinated side chains with the structural versatility of acrylate backbones. The strong electron-withdrawing nature and low polarizability of the C–F bond impart exceptional thermal stability, chemical resistance, and notably low surface energy to these materials, making them ideal candidates for non-wetting coatings, anti-fouling surfaces, and advanced membranes.^{4,5} Recent advances in controlled radical polymerization techniques, such as ATRP and RAFT, have enabled precise manipulation of molecular weight, architecture, and composition, allowing for the tailored design of polyfluoroalkyl acrylates with tunable surface and bulk properties.^{6,7} Moreover, the self-assembly behavior driven by the incompatibility between fluoroalkyl side chains and hydrocarbon segments leads to microphase-separated structures that further enhance their performance in demanding environments.^{8,9} Fluorinated polymers have emerged as a cornerstone material class for CO₂-selective applications, particularly in supercritical CO₂ (scCO₂) applications, thus revolutionizing various industrial processes in recent decades.^{10–12} They are well known for their affinity for scCO₂ and exhibit higher solubility than non-fluorinated polymers.^{10,13} Their low surface energy and chemical resistance make them ideal for use in aggressive scCO₂ environments, such as in separation, coating, and extraction

^a Research Center for Negative Emissions Technologies, Kyushu University, 744 Motoooka, Nishi-ku, Fukuoka 819-0395, Japan.
E-mail: takahara.atsushi.150@m.kyushu-u.ac.jp

^b New Business Development Department, Development Office, Electronic Materials General Division, Electronics Company, AGC Inc., 1-5-1, Marunouchi, Chiyoda-ku, Tokyo 100-8405, Japan

† Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d5sm00413f>



was then added to the solution. The flask was sealed with a rubber septum, and the mixture was purged with nitrogen for 20 min to remove dissolved oxygen. The flask was then immersed in an oil bath preheated to 60 °C, and the polymerization was allowed to proceed for 24 h with constant stirring. After the reaction, the polymer solution was cooled to room temperature and precipitated in 300 mL of cold methanol. The precipitate was collected *via* filtration, washed thoroughly with methanol, and vacuum-dried at room temperature for 24 h to yield PFMA-C₁ as a white powder (yield of 85%). Identical conditions were applied to synthesize PFMA-C₂, PFMA-C₄, PFA-C₂, and PFA-C₄ using their respective monomers with yields varying between 78 and 90%.

Characterization of PFA and PFMA thin films

The synthesized polymers were characterized by gel permeation chromatography (GPC) using a Tosoh HLC-8320 GPC system with a evaporative light scattering detector (ELSD). The separation was performed using an Agilent PLgel 5 μm Mixed-C column (300 × 7.5 mm) at 40 °C. The mobile phase was a mixture of HFE-7300 (1,1,1,2,2,3,4,5,5,5-decafluoro-3-methoxy-4-(trifluoromethyl)pentane) and HFIP (1,1,1,3,3,3-hexafluoroisopropanol) (9:1, v/v) at a flow rate of 1.0 mL min⁻¹. The number-average molecular weight (M_n), weight-average molecular weight (M_w), and polydispersity index (\bar{D}) were determined using PMMA calibration standards. The chemical structure of the synthesized PFMA and PFA was confirmed using ¹H-NMR spectroscopy. Samples were dissolved in a 50/50 vol% mixture of deuterated trifluoroacetic acid and CDCl₃ (5 mg mL⁻¹). ¹H NMR spectra were recorded on a Bruker AVANCE III 600 MHz spectrometer (room temperature). The samples were referenced to tetramethylsilane (TMS) internal standard. Thermal properties of the polymers were analyzed using differential scanning calorimetry (DSC) on a Thermo Plus EVO2 DSC vesta equipped with an automated liquid–nitrogen cooling system (Rigaku Corp., Japan) under an N₂ atmosphere, with a heating/cooling rate of 20 °C min⁻¹ in the range of –80 °C to 150 °C.

The wide-angle X-ray scattering (WAXS) measurements were conducted at the SPring-8 synchrotron radiation facility using beamline BL38B1 with an incident wavelength of 1 Å. A PILATUS3X 2 M detector (Dectris Ltd) with 1475 × 1679 pixels of 172 × 172 μm² was placed 260 mm away from the sample. The samples were mounted between two Kapton films (20 μm each) and exposed to an X-ray beam for 10 s. CeO₂ was used as the standard for calibration of the detector geometry and image integration using the azimuthal integration tool pyFAI.¹⁹ The scattered intensities were azimuthally averaged and recorded as a function of the scattering vector q , where $q = 4\pi\lambda^{-1}\sin\theta$, and θ represents half the scattering angle. The q value was in the 2–36 nm⁻¹ range.

The polymers were analyzed using a Bruker Vertex70 FTIR spectrometer equipped with horizontally attenuated total reflectance accessories (PIKE). Thin films of the polymers were prepared by casting their solutions onto attenuated total reflection (ATR) crystals (45° Ge, length of 8 cm, and thickness of 4 mm). Spectra were recorded in the range of 4000–600 cm⁻¹ with a resolution of 2 cm⁻¹. CO₂ (or N₂) was introduced into the flow cell at a rate of 200 sccm (standard cubic centimeters per

minute) using a mass-flow controller (KOFLOC, 3660). For time-resolved data acquisition, 1000 spectra were collected in an interval of 78 ms with a resolution of 2 cm⁻¹, accumulating 16 scans for each spectrum.

To determine the CO₂ diffusivity in the polymer using ATR–FTIR spectroscopy, the concentration was correlated with the experimental absorbance using the Beer–Lambert law. Barbari *et al.* integrated the equation for the evanescent field strength using the Beer–Lambert law to establish an ATR Fickian diffusion model based on certain assumptions.^{20,21} A comprehensive derivation of the equation was conducted in previous studies.²⁰ The simplified equation for diffusion is as follows:

$$\ln\left(1 - \frac{A_t}{A_\infty}\right) = \ln\left(\frac{4}{\pi}\right) - \frac{D\pi^2}{4L^2}t \quad (1)$$

where L denotes the membrane thickness, and A_t/A_∞ represents the normalized absorbance. D is the only adjustable parameter and can be determined using a linear least-squares regression by plotting $\ln(1 - A_t/A_\infty)$ against t .

The CO₂ sorption capacity was examined at room temperature using a quartz crystal microbalance (QCM922A, Seiko EG&G). CO₂ (or N₂) was introduced into the flow cell at a flow rate of 50 sccm using a mass-flow controller (KOFLOC, BR-2C). The QCM crystals (QA-A9M-AU(M)(SEP), Seiko EG&G) were cleaned through UV/ozone treatment (ASM401OZ, ASUMI GIKEN) for 10 min. Thin films of the polymers were spin-coated (MS-B100 Spin-Coater, Mikasa) onto QCM crystals from their solutions (1 wt% in AK-225). The frequency changes were monitored in response to exposure to certain atmospheres (sorption and desorption). The Sauerbrey equation (eqn (2)) was used to relate the frequency change (ΔF) to the added mass (Δm) on the crystal, either the deposited polymer (m_p) or the gas sorption (m_g).²² The CO₂ sorption capacity is defined as m_g/m_p (mg g⁻¹).

$$-\Delta F = \frac{2F_0^2}{A\sqrt{\mu_q\rho_q}}\Delta m \quad (2)$$

where F_0 is the fundamental resonant frequency, A is the active area (0.196 cm²), μ_q is the shear modulus of AT-cut quartz crystal (2.95×10^{11} g cm⁻¹ s⁻²), and ρ_q is the density of quartz (2.65 g cm⁻³), ΔF is the frequency change, and Δm is mass change.

The permeability (P) of CO₂ was determined at 25 °C at a feed pressure of 150 cmHg and atmospheric permeate pressure using the constant-pressure/variable-volume method.²³ The samples were coated onto microporous polyacrylonitrile (PAN) support membranes (UF 010104, SolSep BV) and sealed inside a stainless steel filter holder (Millipore XX4404700). The permeation flow was measured using a mass flow meter (Agilent ADM Flow Meter, G6691A). P was calculated using the following equation:

$$P = \frac{l}{\Delta p} \frac{273p_{\text{atm}}}{TA76} \left(\frac{dV}{dt}\right) \quad (3)$$

where dV/dt is the permeate-side flow rate (cm³ s⁻¹), T is the operating temperature (K), and A is the effective area of the



membrane (cm^2). All key transport measurements, including permeability, diffusivity, and sorption experiments, were performed in triplicate for each polymer sample. The reported values in figures and tables represent the mean of these measurements, and the associated uncertainty is reported as the standard deviation.

Results and discussion

Characterization of PFA and PFMA thin films

Poly(fluoroalkyl acrylates) (PFA- C_y) and poly(fluoroalkyl methacrylates) (PFMA- C_y) with varying R_f lengths were synthesized *via* radical polymerization. R_f was chosen to be short to yield an amorphous polymer. Polymers with C_2 chains were soluble in THF, whereas longer C_4 chains were dissolved in hydrochlorofluorocarbon solvents (AK-225). As shown in Fig. 1, the ^1H NMR spectra of PFA and PFMA showed similar proton signals corresponding to methylene groups in the backbone (H_a) as well as the sidechain (H_c), which appeared at 2.4 ppm and 4.5 ppm, respectively. PFMA also showed proton signals corresponding to the methyl groups (H_b) at 0.8–1.2 ppm. All the integral results were consistent with the stoichiometric ratio of the molecular structure. The GPC results (Table 1) indicated that all the samples exhibited M_n in the range of 20–100 kDa and $D < 2.0$. This high degree of polymerization suggests that structural integrity is necessary for membrane applications.

The glass transition temperatures (T_g) of the polymer samples were obtained from the DSC traces, as shown in Fig. 2. Our findings indicate that both side-chain length and the presence of an α -methyl group on the backbone played crucial roles in determining T_g , with notable differences between fluorinated polymers and their non-fluorinated counterparts. The side-chain length emerged as a primary factor influencing T_g in these systems. In both the PFA and PFMA series, shorter side chains resulted in higher T_g values, which can be attributed to the stronger intermolecular interactions between the polymer backbones.^{24–26} This trend was also observed in non-fluorinated

Table 1 Number-average (M_n) and weight-average (M_w) molecular weights and glass transition temperatures (T_g) of various PFA and PFMA polymers

	M_n (kDa)	M_w (kDa)	PDI	T_g (K)
PFA- C_2	75.3	147.0	1.95	258
PFA- C_4	94.4	166.7	1.76	249
PFMA- C_2	21.0	28.3	1.35	332
PFMA- C_4	45.2	58.9	1.30	294

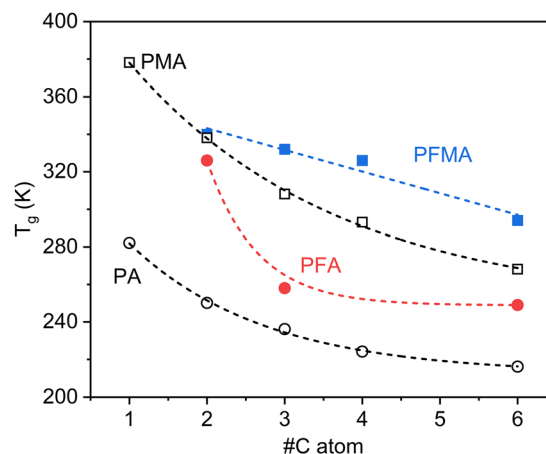


Fig. 2 Glass transition temperature (T_g) of PFA and PFMA as well as nonfluorine counterparts PA and PMA.

analogs.²⁷ Notably, longer side chains lowered T_g by increasing the free volume and acting as internal plasticizers, reducing the energy required for segmental motion. Additionally, the presence of a methyl group on the backbone, which distinguishes methacrylates from acrylates, consistently led to higher T_g values.

A comparison of the fluorinated (meth)acrylates with their non-fluorinated analogs revealed consistently higher T_g values in the fluorinated versions. This can be attributed to a combination of factors. The high electronegativity of fluorine atoms creates strong C-F dipoles, leading to enhanced intermolecular dipole-dipole interactions between chains. Furthermore, the larger atomic radius of fluorine compared to hydrogen results in increased steric hindrance and segmental stiffness, which restricts chain mobility. Together, these effects elevate the energy barrier for segmental motion, resulting in a higher T_g compared to their non-fluorinated analogs. This effect was more pronounced for PFA with shorter side chains, suggesting a complex interplay between fluorination effects and side chain dynamics.

These structure–property relationships have significant implications for the CO_2 transport behavior. Polymers with T_g values near or below the intended operating temperature are likely to exhibit enhanced CO_2 permeability owing to the increased chain mobility in the rubbery state. Tuning the T_g of fluorinated (meth)acrylates *via* careful control of the side chain length, backbone substitution, and degree of fluorination could optimize the CO_2 transport properties.

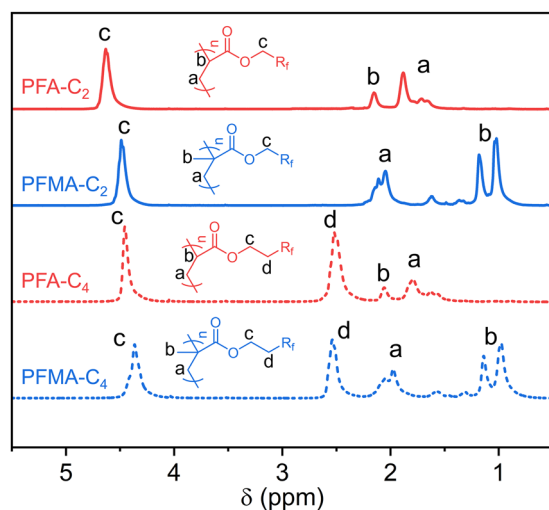


Fig. 1 ^1H NMR spectra of PFA and PFMA.



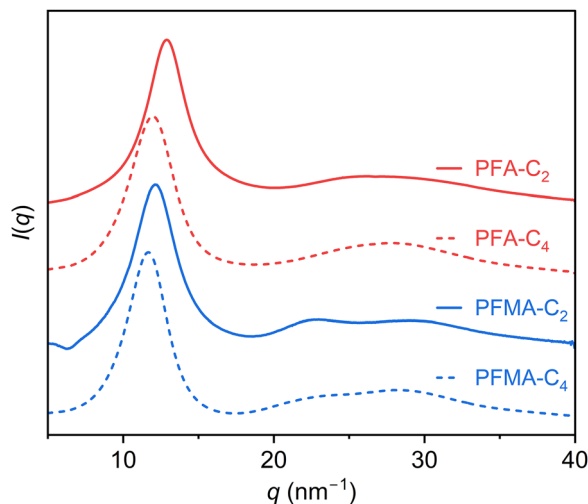


Fig. 3 WAXS 1d profiles of PFA and PFMA.

All samples exhibited broad, diffuse peaks in their WAXS profiles (Fig. 3), with no sharp crystalline peaks, confirming their amorphous nature. This finding is consistent with previous studies and is advantageous for applications requiring gas permeability, as amorphous polymers typically possess a higher free volume than semi-crystalline polymers, facilitating molecular diffusion.⁴

For PFMA, three characteristic peaks were observed at approximately 1.2 \AA^{-1} (q_I), 2.2 \AA^{-1} (q_{II}), and 2.8 \AA^{-1} (q_{III}). The detailed d -spacing (d) values were calculated from q_{peak} using the Bragg equation ($d = 2\pi/q$) and summarized in Table 2. The q_I peak with $d_I \approx 5 \text{ \AA}$ arose from the packing of fluoroalkyl sidechains.²⁸ The q_{II} peak corresponded to an intramolecular methyl–methyl correlation, while the third broad maximum q_{III} peak corresponded to the average repeat-unit correlation, averaged from the C–C distances along the backbone chains, which has the shortest correlation distance.²⁹ This profile resembled that of the nonfluorine counterpart PMMA, but with noticeable variations in peak positions. The q_I peak of PMMA at 0.94 \AA^{-1} arose from the intermolecular backbone–backbone, backbone–ester, and backbone–methyl correlations.^{29,30} The q_{II} and q_{III} peaks were almost identical to those of PFMA, which corresponded to the average repeat-unit correlation in the (meth)acrylate backbone.

In contrast, PFAs exhibited only two prominent broad peaks at $q_I = 1.2 \text{ \AA}^{-1}$ and $q_{III} = 2.8 \text{ \AA}^{-1}$. Similar to PFMA, the q_I peak of PFA arose from the packing of fluoroalkyl side chains. However,

Table 2 Peak positions (q) and d -spacing (d) values of PFA and PFMA from the WAXS 1d profiles

	q_I (\AA^{-1})	d_1 (\AA)	q_{II} (\AA^{-1})	d_2 (\AA)	q_{III} (\AA^{-1})	d_3 (\AA)
PFA-C ₂	1.29	4.88	—	—	2.77	2.27
PFA-C ₄	1.20	5.26	—	—	2.80	2.24
PFMA-C ₂	1.21	5.19	2.23	2.81	2.91	2.16
PFMA-C ₄	1.17	5.38	2.24	2.81	2.81	2.24

d -Spacing (d) is calculated from Bragg equation.

PFA showed a slightly reduced d_1 spacing than PFMA. This is likely due to the presence of the α -methyl group in PFMA, which caused steric hindrance and disrupted close packing of the fluorinated side chains. The absence of a q_{II} peak in PFA indicates the lack of an α -methyl group, eliminating methyl–methyl correlations observed in PFA.

A subtle increase in d_1 was observed for both PFA and PFMA with increasing side-chain length (C₄). This indicates that a longer side chain introduced greater flexibility and potentially higher free volume, which is consistent with the trends observed in T_g . The WAXS data provide insights into the molecular packing and structural organization of these fluorinated polymers, which are crucial for understanding their properties, including their interactions with CO₂ and their potential performance in various applications.³⁰

In addition, ATR-FTIR characterization of PFA and PFMA revealed key differences as well as similarities in their chemical structures (Fig. 4). The spectra indicated no presence of unreacted monomers, as evidenced by the absence of significant C=C stretching peaks around 1640 cm^{-1} . This aligns with the GPC results, confirming the high degree of polymerization in both PFA and PFMA. Both polymers exhibited characteristic absorption bands corresponding to their fluorinated side chains and ester functional groups. All the (meth)acrylate samples exhibited a characteristic carbonyl stretching vibration band at $1730\text{--}1760 \text{ cm}^{-1}$ corresponding to the ester functional groups in their backbones. The subtle variation in wavenumber resulted from the α -substituents and R_f groups, probably because of the differences in the chemical environment around the carbonyl group. Furthermore, the presence of spacer methylene groups and the α -methyl group that could alter the electron density and steric interactions may have contributed to the observed shifts. The strong absorption bands at ν_s (CF₂) $\sim 1130 \text{ cm}^{-1}$ and ν_{as} (CF₂) $\sim 1200 \text{ cm}^{-1}$ attributed to the C–F stretching vibrations of fluoroalkyl side chains, were prominent in both PFA and PFMA.³¹ A notable difference emerged when comparing the C₄ and C₂ variants: the C₄ samples exhibited peaks at higher wavenumbers compared with the C₂ samples

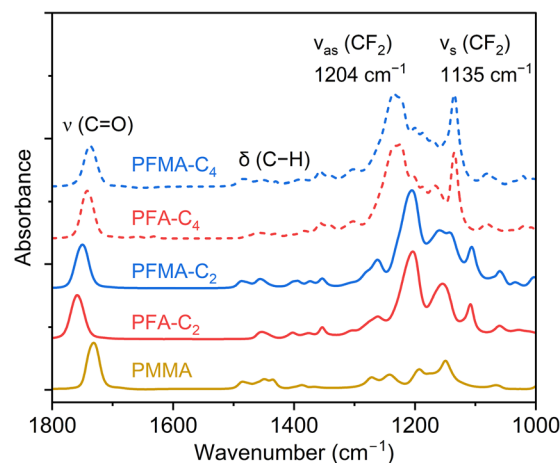


Fig. 4 ATR-FTIR spectra of PFA and PFMA.



(1135 cm^{-1} vs. 1126 cm^{-1}). This shift could be related to side chain aggregation because the increased length favors stable packing.³² Comparable peak profiles suggest that the fluorinated groups dominate the spectra in this region. However, a notable difference between the two polymers was observed in the absorption bands around 2950 cm^{-1} , with PFMA exhibiting more pronounced bands. These bands were attributed to the methyl group associated with the methacrylate structure, which distinguishes PFMA from PFA.

The CO_2 sorption by the polymer at equilibrium was demonstrated by ATR-FTIR spectra (Fig. 5). The CO_2 that diffused into the polymer exhibited a distinct asymmetric stretching vibration (ν_3) band compared to gaseous CO_2 . This result aligns with previous literature on CO_2 dissolved in various solvents and polymers. The ν_3 band for all fluoro(meth)acrylates appeared at an identical wavenumber of 2342 cm^{-1} , which is different from that of PMMA at 2338 cm^{-1} . This discrepancy could be attributed to the different interaction sites of CO_2 with the fluorine and nonfluorine counterparts. For PMMA, CO_2 mainly interacts with the carbonyl groups.³³ In contrast, with the incorporation of R_f groups in PFMA, CO_2 might have interacted with the side chain *via* weak Lewis base (R_f)–Lewis acid (CO_2) interactions or *via* its quadrupole moment interacting with the induced dipoles of the R_f . The identical ν_3 band position at 2342 cm^{-1} across all fluoromethacrylates (regardless of chain length) suggests that the R_f groups dominated the CO_2 interactions, creating a consistent sorption environment distinct from that of PMMA.

In addition, time-resolved ATR-FTIR was used to analyze CO_2 diffusion in both polymers. The concentration of CO_2 was correlated to the experimental absorbance using the Beer-Lambert law. The absorbance (A_t) of the ν_3 band was tracked as a function of time during sorption and desorption. The experimental data were fitted using Fickian diffusion, a model that describes diffusion driven by concentration gradients. The diffusion coefficient (D) was derived from linear regression of $\ln(1 - A_t/A_\infty)$ vs. time (t) plot as shown in Fig. 6a. These plots reveal that CO_2 diffused faster in PFA than in PFMA at

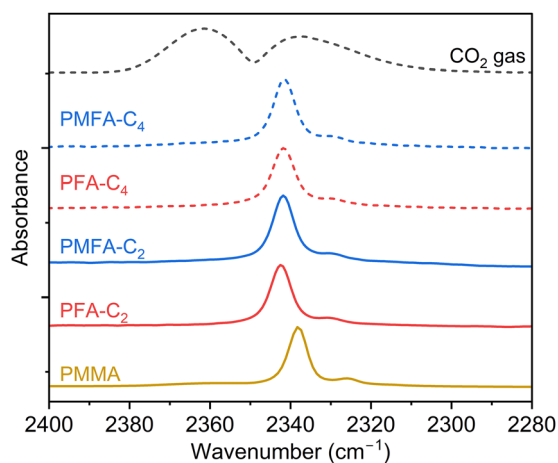


Fig. 5 ATR-FTIR spectra of CO_2 in PFA and PFMA.

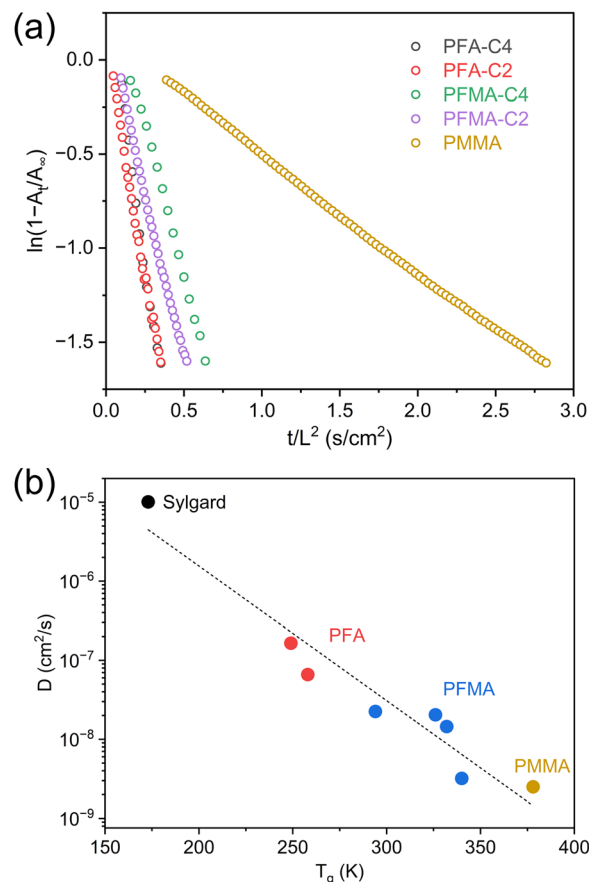


Fig. 6 (a) Plot of $\ln(1 - A_t/A_\infty)$ for the CO_2 asymmetric stretching band (ν_3) as a function of time. (b) Relationship between diffusivity (D) and T_g for various polymers.

comparable side chain lengths. The faster diffusion in PFA is likely due to the absence of the α -methyl group, which may reduce steric hindrance and enhance side chain mobility, thereby increasing free volume and facilitating CO_2 transport. In addition, longer chain length induced higher D values for each polymer series.

The D values were plotted as a function of T_g for various polymer samples in Fig. 6b. This scatter plot, which is accompanied by a downward-sloping trend line, reveals an inverse relationship between D and T_g . Sylgard, with an exceptionally low T_g (170 K), exhibited a high D value ($10^{-5} \text{ cm}^2 \text{ s}^{-1}$), indicating significant chain mobility and free volume. In contrast, PMMA, which has a high T_g (380 K), exhibited a significantly lower D value ($10^{-9} \text{ cm}^2 \text{ s}^{-1}$), indicating its rigid structure with a minimal free volume. Compared with PFMA, PFA chains were more mobile and thus increased the free volume, facilitating faster gas diffusion, as reflected by the higher D values. As discussed earlier, the T_g values of the samples were significantly influenced by the fluoroalkyl (R_f) side-chain length. In both series, longer R_f chains induced higher D values. A side chain with faster dynamics increased the free volume, thus decreasing T_g and inducing faster CO_2 diffusion.

The strong affinity of both polymers for CO_2 , inferred from the pronounced C–F stretching bands, suggests that these materials probably exhibit high CO_2 solubility—a desirable



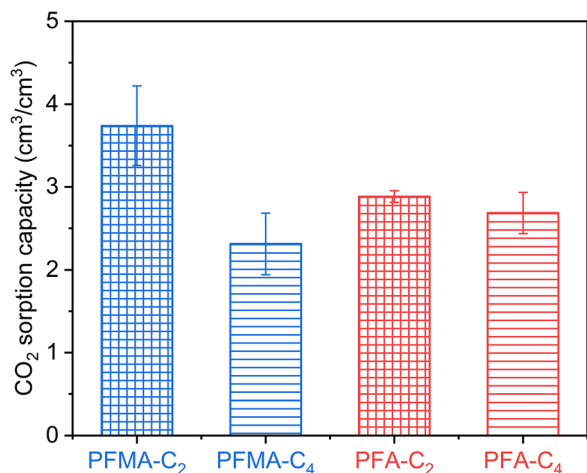
Table 3 CO₂ permeability (*P*), solubility (*S*), and diffusion coefficients (*D*) for PFMA and PFA

	<i>T_g</i> (K)	<i>D</i> /10 ⁻⁸ (cm ² s ⁻¹)	<i>S</i> (cm ³ cm ⁻³)	<i>P</i> (Barrer)
PFA-C ₂	258	6.61 ± 0.28	2.82 ± 0.10	50 ± 3
PFA-C ₄	249	16.5 ± 0.5	2.61 ± 0.25	121 ± 4
PFMA-C ₂	332	1.45 ± 0.05	3.60 ± 0.48	14 ± 1
PFMA-C ₄	294	2.25 ± 0.11	2.57 ± 0.37	20 ± 1

Measured at 297 K.

property for gas separation applications.^{34,35} The CO₂ sorption capacities of PFA and PFMA were evaluated using QCM. The Sauerbrey equation was used to calculate the mass of the deposited polymers and sorbed molecules in a certain atmosphere. The CO₂ sorption data are summarized in Table 3 and Fig. 7. Compared with the *n*-alkyl methacrylates, both PFA and PFMA exhibited higher CO₂ sorption capacities owing to stronger interactions between CO₂ and *R_f* groups.

In addition, PFMA exhibited higher CO₂ sorption than PFA, which could be attributed to the additional unoccupied volume of the glassy polymers. The α -methyl group hindered the chain rotation, which limited the efficient packing of the *R_f* side chain. However, both PFMA and PFA exhibited a clear decrease in CO₂ sorption as the *R_f* side-chain length increased from C₂ to C₄. This reduction was attributed to increased van der Waals interactions between the longer *R_f* chains, thereby weakening their interactions with CO₂. Thus, *R_f* chains longer than C₈ are expected to become semi-crystalline, exhibiting significantly lower gas sorption than the amorphous ones. This effect was also observed for the *n*-alkyl methacrylates with increasing side-chain length, as previously reported.^{36,37} The highest CO₂ sorption capacity was observed in PFMA-C₂ (5.8 mg g⁻¹), where the α -methyl group and shorter C₂ *R_f* chains synergistically maximized free volume and CO₂ interaction sites. The α -methyl group disrupted packing efficiency, while the shorter *R_f* chains minimized interchain interactions. In contrast, PFA-C₄ exhibited the lowest sorption (3.5 mg g⁻¹) owing to more efficient packing and stronger *R_f*-*R_f* interactions. It is crucial to note that this high CO₂ affinity is a distinguishing feature of the

**Fig. 7** CO₂ sorption of PFA and PFMA calculated from QCM results.

fluorinated polymers. In their non-fluorinated counterparts, such as poly(*n*-alkyl methacrylates), gas sorption is significantly lower as it is primarily driven by weaker van der Waals forces. The enhanced solubility in PFA and PFMA is a direct consequence of the favorable CO₂-*R_f* interactions, as PFMA surpasses PFA owing to the disruption of packing efficiency by the α -methyl group. Shorter *R_f* chains (C₂) optimized sorption by reducing interchain interactions, whereas longer chains (C₄) diminished sorption due to enhanced *R_f*-*R_f* interactions. The α -methyl group and shorter *R_f* length contributed together in achieving the inefficient packing of *R_f* to realize higher CO₂ sorption.

The gas permeability through a polymer membrane is generally expressed as a product of its solubility and diffusivity ($P = S \times D$). The interplay between solubility, diffusivity, and glass transition temperature (*T_g*) is complex, as structural features—such as the *R_f* length—simultaneously influence the free volume, chain flexibility, and intermolecular interactions. For PFA, *P* rose from 50 Barrer (PFA-C₂) to 121 Barrer (PFA-C₄), driven by a significant *D* increase (6.6 to 16.5 × 10⁻⁸ cm² s⁻¹) despite a slight *S* decrease (2.82 to 2.61 cm³ cm⁻³). This indicates that the increased diffusivity, resulting from lower *T_g* and higher free volume, outweighed any solubility reduction caused by *R_f*-*R_f* interactions. Compared with PFA, PFMA exhibited a significantly lower *P* value at comparable *R_f* lengths. This reveals that the decreased chain mobility outweighed the pronounced *S* increase. This nuanced interplay highlights the challenge of optimizing *D* and *S* to achieve high gas permeability, thereby guiding the design of advanced fluorinated materials for CO₂ capture and gas separation membrane applications.

Conclusions

In this work, we have established critical structure–property relationships governing CO₂ transport in PFA and PFMA polymers. Our findings demonstrate that while both polymer families show a strong affinity for CO₂, their permeability is dictated by a delicate balance between solubility and diffusivity. The absence of an α -methyl group in the PFA backbone enhances chain mobility, leading to higher diffusivity and, consequently, superior CO₂ permeability compared to the PFMA series. Meanwhile, extending the fluoroalkyl side-chain length acts as an internal plasticizer that lowers the glass transition temperature (*T_g*) and boosts diffusivity, albeit at the expense of solubility due to side-chain aggregation. This structure-guided approach provides a foundation for developing next-generation fluorinated membranes with tunable gas transport properties, enabling their application in direct air capture (DAC), post-combustion CO₂ separation, and barrier materials for gas packaging and environmental sensing.

Author contributions

The manuscript was written through contributions of all authors. All authors have given approval to the final version of the manuscript.



Conflicts of interest

There are no conflicts to declare.

Data availability

The authors declare that the data supporting the findings of this study are available within the paper and its ESI.† Should any raw data files be needed in another format they are available from the corresponding author upon reasonable request.

Acknowledgements

This study was based on the results obtained from a project (JPNP18016) commissioned by the New Energy and Industrial Technology Development Organization (NEDO). The synchrotron radiation experiments were performed at the BL38B1 beamline of SPring-8 (RIKEN). The authors thank Dr Hiroyasu Masunaga (Japan Synchrotron Radiation Research Institute) and Dr Sono Sasaki (Kyoto Institute of Technology) for their assistance in X-ray scattering measurement.

References

- 1 S. Fujikawa and R. Selyanchyn, *MRS Bull.*, 2022, **47**, 416–423.
- 2 S. Fujikawa, R. Selyanchyn and T. Kunitake, *Polym. J.*, 2021, **53**, 111–119.
- 3 M. Liu, M. D. Nothling, P. A. Webley, Q. Fu and G. G. Qiao, *Acc. Chem. Res.*, 2019, **52**, 1905–1914.
- 4 K. Honda, M. Morita, H. Otsuka and A. Takahara, *Macromolecules*, 2005, **38**, 5699–5705.
- 5 Y. Liu, Y. Higaki, M. Mukai, N. Ohta, T. Kabe and A. Takahara, *Polym. J.*, 2019, **51**, 189–198.
- 6 W. Yao, Y. Li and X. Huang, *Polymer*, 2014, **55**, 6197–6211.
- 7 S. Banerjee, B. V. Tawade, V. Admiral, L. X. Dupuy, M. P. MacDonald and B. Améduri, *Polym. Chem.*, 2017, **8**, 1978–1988.
- 8 J.-M. Corpart, S. Girault and D. Juhué, *Langmuir*, 2001, **17**, 7237–7244.
- 9 F. Montefusco, R. Bongiovanni, A. Priola and B. Ameduri, *Macromolecules*, 2004, **37**, 9804–9813.
- 10 J. L. Kendall, D. A. Canelas, J. L. Young and J. M. DeSimone, *Chem. Rev.*, 1999, **99**, 543–564.
- 11 M. Zhou, R. Ni, Y. Zhao, J. Huang and X. Deng, *Soft Matter*, 2021, **17**, 5107–5115.
- 12 B. Améduri, *Molecules*, 2023, **28**, 7564.
- 13 J. M. DeSimone, Z. Guan and C. S. Elsbernd, *Science*, 1992, **257**, 945–947.
- 14 E. J. Beckman, *Chem. Commun.*, 2004, 1885–1888.
- 15 B. Ameduri, *Macromolecules*, 2025, **58**, 2781–2791.
- 16 I. T. Cousins, G. Goldenman, D. Herzke, R. Lohmann, M. Miller, C. A. Ng, S. Patton, M. Scheringer, X. Trier, L. Vierke, Z. Wang and J. C. DeWitt, *Environ. Sci.: Processes Impacts*, 2019, **21**, 1803–1815.
- 17 K. Honda, M. Morita, O. Sakata, S. Sasaki and A. Takahara, *Macromolecules*, 2010, **43**, 454–460.
- 18 H. Yamaguchi, P. Gin, H. Arita, M. Kobayashi, S. Bennett, S. K. Satija, M. Asada, T. Koga and A. Takahara, *RSC Adv.*, 2013, **3**, 4778–4785.
- 19 G. Ashiotis, A. Deschildre, Z. Nawaz, J. P. Wright, D. Karkoulis, F. E. Picca and J. Kieffer, *J. Appl. Crystallogr.*, 2015, **48**, 510–519.
- 20 G. T. Fieldson and T. A. Barbari, *Polymer*, 1993, **34**, 1146–1153.
- 21 Y. A. Elabd, M. G. Baschetti and T. A. Barbari, *J. Polym. Sci., Part B: Polym. Phys.*, 2003, **41**, 2794–2807.
- 22 G. Sauerbrey, *Z. Phys. Chem.*, 1959, **155**, 206–222.
- 23 H. Lin and B. Freeman, *Springer Handbook of Materials Measurement Methods*, Springer, Berlin, Heidelberg, 1st edn, 2006, pp. 371–387.
- 24 A. Arbe, A.-C. Genix, S. Arrese-Igor, J. Colmenero and D. Richter, *Macromolecules*, 2010, **43**, 3107–3119.
- 25 F. Rindfleisch, T. P. DiNoia and M. A. McHugh, *J. Phys. Chem.*, 1996, **100**, 15581–15587.
- 26 S. Koizumi, K. Tadano, Y. Tanaka, T. Shimidzu, S. Kutsumizu and S. Yano, *Macromolecules*, 1992, **25**, 6563–6567.
- 27 S. Rogers and L. Mandelkern, *J. Phys. Chem.*, 1957, **61**, 985–991.
- 28 V. V. Volkov, N. A. Platé, A. Takahara, T. Kajiyama, N. Amaya and Y. Murata, *Polymer*, 1992, **33**, 1316–1320.
- 29 S. Eim, S. Jo, J. Kim, S. Park, D. Lee, T. P. Russell and D. Y. Ryu, *ACS Macro Lett.*, 2024, **13**, 1490–1494.
- 30 R. Lovell and A. H. Windle, *Polymer*, 1981, **22**, 175–184.
- 31 M. Matsunaga, T. Suzuki, K. Yamamoto and T. Hasegawa, *Macromolecules*, 2008, **41**, 5780–5784.
- 32 T. Hasegawa, in *Quantitative Infrared Spectroscopy for Understanding of a Condensed Matter*, ed. T. Hasegawa, Springer, Japan, Tokyo, 2017, pp. 165–193.
- 33 S. G. Kazarian, M. F. Vincent, F. V. Bright, C. L. Liotta and C. A. Eckert, *J. Am. Chem. Soc.*, 1996, **118**, 1729–1736.
- 34 M. J. Muldoon, S. N. V. K. Aki, J. L. Anderson, J. K. Dixon and J. F. Brennecke, *J. Phys. Chem. B*, 2007, **111**, 9001–9009.
- 35 P. Raveendran and S. L. Wallen, *J. Phys. Chem. B*, 2003, **107**, 1473–1477.
- 36 Z. Mogri and D. R. Paul, *Polymer*, 2001, **42**, 7781–7789.
- 37 Z. Mogri and D. R. Paul, *Polymer*, 2001, **42**, 7765–7780.

