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# Participation of electrochemically inserted protons in the hydrogen evolution reaction on tungsten oxides†

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Understanding the mechanisms by which electrodes undergo the hydrogen evolution reaction (HER) is necessary to design better materials for aqueous energy storage and conversion. Here, we investigate the HER mechanism on tungsten oxide electrodes, which are stable in acidic electrolytes and can undergo proton-insertion coupled electron transfer concomitant with the HER. Electrochemical characterization showed that anhydrous and hydrated tungsten oxides undergo changes in HER activity coincident with changes in proton composition, with activity in the order  $H_xWO_3 \cdot H_2O > H_xWO_3 > H_xWO_3 \cdot 2H_2O$ . We used operando X-ray diffraction and density functional theory to understand the structural and electronic changes in the materials at high states of proton insertion, when the oxides are most active towards the HER.  $H_{0.69}WO_3 \cdot H_2O$  and  $H_{0.65}WO_3$  have similar proton composition, structural symmetry, and electronic properties at the onset of the HER, yet exhibit different activity. We hypothesize that the electrochemically inserted protons can diffuse in hydrogen bronzes and participate in the HER. This would render the oxide volume, and not just the surface, as a proton and electron reservoir at high overpotentials. HER activity is highest in  $H_xWO_3 \cdot H_2O$ , which optimizes both the degree of proton insertion and solid-state proton transport kinetics. Our results highlight the interplay between the HER and proton insertion-coupled electron transfer on transition metal oxides, many of which are non-blocking electrodes towards protons.

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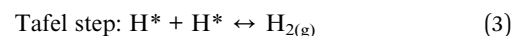
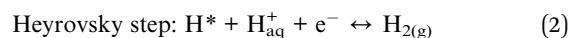
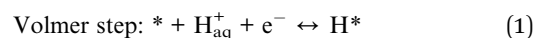
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## 1. Introduction

The hydrogen evolution reaction (HER) is the cathodic reaction in electrochemical water splitting that leads to the formation of  $H_2$  gas, which can be used as an energy dense fuel or precursor for chemical synthesis.<sup>1,2</sup> On the other hand, it can also be a parasitic reaction that limits the potential window of aqueous energy storage devices and competes with other electrochemical reduction reactions, like those involved in the conversion of  $CO_2$ .<sup>3-5</sup> Understanding electrode properties that influence the HER is therefore necessary to advance aqueous energy storage and conversion devices. The HER is part of a broad class of proton-coupled electron transfer (PCET) reactions which are ubiquitous in aqueous electrochemical environments. Based on the possibility of PCET, HER electrocatalysts can be divided into three types (Fig. 1): (a) proton-blocking electrocatalysts, (b) non proton-blocking electrocatalysts, and (c) reconstruction

electrocatalysts (materials undergoing conversion or dissolution and deposition reactions). Gold and platinum represent proton-blocking electrocatalysts, where the bulk composition of the electrocatalyst remains constant and the proton reduction and HER take place at the surface.<sup>6-8</sup> In acidic pH, the elementary steps of the HER on a proton-blocking electrocatalyst surface, where \* indicates a surface adsorption site, proceed as follows:



Importantly, adsorbed hydrogen ( $H^*$ ) produced in (1) is a reactant for  $H_2$  in (2) and (3). The binding energy of adsorbed hydrogen species on different electrode surfaces has long been used as a catalytic activity descriptor to design, for example, different metal alloys for HER electrocatalysts.<sup>9,10</sup>

Palladium represents the most well studied non proton-blocking electrocatalyst (Fig. 1b). In Pd,  $H^*$  can be either a reactant for steps (2) and (3) or diffuse into the electrocatalyst,

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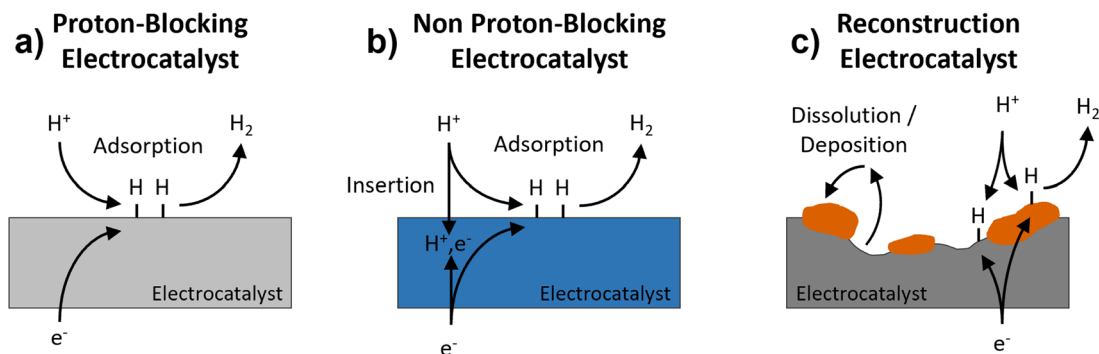


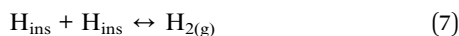
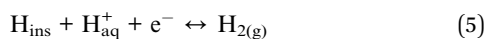
Fig. 1 Schematic representation of three types of electrocatalysts for the HER: (a) proton-blocking electrocatalyst, (b) non proton-blocking electrocatalyst, and (c) reconstruction electrocatalyst.

forming PdH<sub>x</sub>.<sup>11</sup> In this kind of electrocatalyst, the bulk hydrogen composition changes as a function of applied potential which leads to changes in the structure and properties of the electrocatalyst, PdH<sub>x</sub>. A Devanathan–Stachurski electrochemical double cell utilizes a Pd membrane to separate two compartments. Upon cathodic polarization on one side of the membrane, the absorbed hydrogen transports through the Pd membrane and subsequently undergoes reactivity on the other side of the membrane.<sup>12–15</sup> Blom *et al.* proposed that the mobility of hydrogen in palladium allows hydrogen to interact with different active sites, which is critical to understanding its electrochemical reactivity during CO<sub>2</sub> reduction to formate.<sup>16</sup>

A broader class of non proton-blocking electrodes are the transition metal oxides. The search for non-precious metal catalysts involves compounds such as transition metal oxides.<sup>17</sup> The presence of a transition metal means that these materials have tunable redox activity which is necessary for high capacity energy storage electrodes as well as for highly active electrocatalysts. Many transition metal oxides exhibit structures that enable proton-insertion coupled electron transfer.<sup>18–21</sup> Proton-insertion coupled electron transfer reactions with transition metal oxides (MO) lead to changes in the composition, structure, and properties of the oxide host prior to or congruent with the HER:<sup>20,22–24</sup>



This leads to the possible participation of electrochemically inserted protons (H<sub>ins</sub>) in the HER, in parallel to H\* on proton-blocking electrocatalysts. For example, following the Heyrovsky and Tafel steps in eqn (2) and (3):



In this case, H<sub>ins</sub> represents the coupled proton and electron present in the MO after the insertion reaction, and participation of H<sub>ins</sub> in the HER would require electron transfer from the MO to H<sup>+</sup>. The participation of inserted protons in the HER would

dictate that the degree of proton insertion (*e.g.* the state of charge, *x* in H<sub>x</sub>MO) and the solid state insertion kinetics (*e.g.* *D*<sub>H<sup>+</sup>,host</sub>) both influence the HER activity of non proton-blocking transition metal oxides. To investigate this interplay between the HER and proton-insertion coupled electron transfer, we utilized a model transition metal oxide system that allows for proton insertion coupled electron transfer and does not undergo reconstruction in strong acid electrolytes, tungsten oxide and its hydrates (WO<sub>3</sub>·*n*H<sub>2</sub>O; *n* = 0, 1, 2).

Tungsten oxides undergo reversible proton insertion-coupled electron transfer in acidic electrolytes. This reaction leads to a semiconductor-to-metal transition and is accompanied by changes in the crystal structure, octahedral tilting, and surface proton binding energies of the oxide.<sup>25–28</sup> Our group showed that the layered structures of the tungsten oxide hydrates confine structural deformation during proton insertion to two dimensions, as opposed to three dimensions in the anhydrous WO<sub>3</sub>, leading to faster proton insertion kinetics.<sup>26,29</sup> Recently, we showed that proton insertion was necessary to activate tungsten oxides for the HER by comparing the HER activity of WO<sub>3</sub>·H<sub>2</sub>O and octylammonium-pillared WO<sub>3</sub>, which suppressed proton insertion.<sup>30</sup> Furthermore, Miu *et al.* used experimental and computational methods to show that the hydrogen binding energy at different active sites and the corresponding HER activity of WO<sub>3</sub> were dependent on the degree of protonation, *x* in H<sub>x</sub>WO<sub>3</sub>.<sup>31</sup> These outcomes emphasize the importance of the changes in composition, structure, and properties associated with proton insertion that precede the onset of the HER.

Herein, we study the influence of proton composition and degree of hydration on the HER activity of tungsten oxides and propose that electrochemically inserted protons participate in the HER. Transient and steady-state electrochemical characterization techniques showed that anhydrous and hydrated tungsten oxides have differences in their HER activity that were associated with changes in proton composition. H<sub>x</sub>WO<sub>3</sub>·H<sub>2</sub>O had the highest activity, followed by H<sub>x</sub>WO<sub>3</sub> and finally H<sub>x</sub>WO<sub>3</sub>·2H<sub>2</sub>O. Operando X-ray diffraction (XRD) measurements showed that proton insertion into WO<sub>3</sub>·H<sub>2</sub>O and WO<sub>3</sub>·2H<sub>2</sub>O resulted in minor deviations in octahedral tilting, while WO<sub>3</sub> went through two phase transitions: monoclinic → tetragonal



→ cubic, which involve significant changes in octahedral tilting and off-center displacements of the  $W^{6+}$ -ion. At high states of proton insertion,  $H_xWO_3$  and  $H_xWO_3 \cdot H_2O$  trend toward higher symmetry, while  $H_xWO_3 \cdot 2H_2O$  remains in a low symmetry monoclinic phase. This result suggests that higher symmetry crystal structures, which promote electronic conductivity, lead to higher electrocatalytic activity. Given that  $H_{0.69}WO_3 \cdot H_2O$  and  $H_{0.65}WO_3$  have similar surface area and morphology, proton composition, structural symmetry, and electronic properties at the onset of HER activity, an additional factor must explain the difference in their HER overpotential and kinetics. To explain this difference, we propose that the protonated materials can serve as proton reservoirs for the HER, whereby electrochemically inserted protons participate in the HER. Participation of inserted protons would lead to a rate dependence on bulk proton transport, which would favor HER activity in  $H_{0.69}WO_3 \cdot H_2O$  vs.  $H_{0.65}WO_3$ . Our results emphasize the importance and possible involvement of the inserted protons undergoing PCET reactions on the surface of electrocatalysts.

## 2. Experimental methods

### 2.1 Synthesis of tungsten oxides

Monoclinic tungsten oxide dihydrate ( $WO_3 \cdot 2H_2O$ ) was synthesized using the Freedman method.<sup>32</sup> Briefly, 50 mL of 1 M  $Na_2WO_4 \cdot 2H_2O$  (99+% Acros Organics) was added dropwise to 450 mL of 3 M HCl (Certified ACS Plus, Fisher Chemical) that was stirred at 300 rpm at room temperature. This mixture was stirred overnight, and the precipitate was collected *via* vacuum filtration and washed with DI water until the pH of the rinse solution reached ~6. Orthorhombic tungsten oxide monohydrate ( $WO_3 \cdot H_2O$ ) was prepared by heating  $WO_3 \cdot 2H_2O$  at 120 °C overnight in air. Monoclinic tungsten trioxide ( $WO_3$ ) was prepared by heating at 350 °C for 5 hours in air with a heating rate of 5 °C  $min^{-1}$ .

### 2.2 Physical characterization

Scanning electron microscopy (SEM) was performed using a field emission Hitachi SU8700 microscope. *Ex situ* SEM was performed on a field emission FEI Verios 460L microscope. X-ray diffraction was performed on a PANalytical Empyrean diffractometer in the Bragg–Brentano geometry with Cu  $K\alpha$  radiation ( $K\alpha_1/K\alpha_2 \lambda = 1.5406 \text{ \AA}/1.5444 \text{ \AA}$ ).

### 2.3 Electrocatalysis measurements

Electrocatalysis measurements were performed using the rotating disk electrode (RDE) setup in a Teflon cell (both from Pine Research Instrumentation). The Teflon cell was cleaned with detergent, 1 M KOH, and aqua regia (3 : 1 HCl to  $HNO_3$ ), with thorough deionized water rinses between each step. The electrolyte was 0.5 M  $H_2SO_4$  (ACS TraceMetal Grade). A porous glass frit bubbler was used to purge the electrolyte with hydrogen gas (ARC3), and hydrogen was bubbled in the cell throughout each measurement. The counter electrode was a graphite rod in a glass compartment (pine), and the reference electrode was Hg/Hg<sub>2</sub>SO<sub>4</sub> in sat'd K<sub>2</sub>SO<sub>4</sub> (Pine Research Instrumentation). The reference electrode was calibrated using

a reversible hydrogen electrode (RHE). To prepare the RHE, we first electrodeposited platinum on a platinum rod (99.997%, Alfa Aesar). This electrodeposition took place in a 41 mM solution of  $H_2PtCl_4$  (>99.9% trace metals basis, Sigma-Aldrich) in 1 M HCl with a platinum counter electrode. A constant current of 10 mA  $cm^{-2}$  was passed through the working electrode for 30 s. The platinum electrode was rinsed thoroughly with deionized water and transferred to an electrochemical cell with 0.5 M  $H_2SO_4$  saturated with hydrogen gas. Hydrogen gas continuously flowed over the platinized platinum electrode, which maintained a constant concentration of dissolved  $H_2$  and a hydrogen pressure slightly higher than 1 atm since excess hydrogen escaped through an open port. The H/H<sup>+</sup> redox couple that forms as a result of  $H_2$  dissociation on the platinized platinum was therefore suitable for use as a RHE. The OCV was measured *versus* the RHE until the potential reached a plateau (typically 20 minutes), and the average potential during the plateau was used as the calibrated potential of the Hg/Hg<sub>2</sub>SO<sub>4</sub> reference electrode on the RHE scale. The potential measured *vs.* Hg/Hg<sub>2</sub>SO<sub>4</sub> was converted to the RHE using the calibration potential according to eqn (8):

$$E_{RHE} = E_{m,REF} + E_{REF} \quad (8)$$

where  $E_{RHE}$  is the potential *vs.* the reversible hydrogen electrode potential,  $E_{m,REF}$  is the potential measured *vs.* the reference electrode (Hg/Hg<sub>2</sub>SO<sub>4</sub> in sat'd K<sub>2</sub>SO<sub>4</sub>) and  $E_{REF}$  is the calibrated potential of the reference electrode *vs.* the RHE.

The working electrode was a drop cast film deposited on a glassy carbon change-disk RDE (Pine Research Instrumentation). The glassy carbon disk was polished using a 0.05  $\mu m$  alumina slurry. After polishing, it was rinsed with deionized water and sonicated for 10 minutes in detergent and 10 minutes in ethanol. The clean glassy carbon was dried at 60 °C in air for at least 1 hour. The electrocatalyst dispersion was prepared by mixing 20 mg of the active material with 29.45  $\mu g$  of Nafion (Ion Power) dispersed in 5 mL of ethanol. The dispersion was sonicated for 30 minutes in an ice bath. Immediately following sonication and vortexing, 5  $\mu L$  of the dispersion was deposited on the glassy carbon change-disk electrode. The resulting film contained 100  $\mu g cm^{-2}$  active material and 0.15  $\mu g cm^2$  Nafion.

Electrochemical impedance spectroscopy (EIS) was performed on the pristine electrode at open circuit voltage using a frequency range of 500 kHz to 100 mHz and an amplitude of 10 mV. The potential was corrected for  $iR$  losses according to eqn (9):

$$E = E_m - iR \quad (9)$$

where  $E$  is the  $iR$ -corrected potential,  $E_m$  is the measured potential,  $i$  is the measured current and  $R$  is the uncompensated resistance from EIS.

All subsequent measurements were performed using a rotation rate of 1600 rpm. Each electrode underwent 3 cyclic voltammetry (CV) cycles at 10 mV  $s^{-1}$ . The potential range for the conditioning cycles was limited to a lower cutoff potential that did not reach the onset of the HER (this potential cutoff was



dependent on the electrocatalyst). The HER activity was studied using CV and chronoamperometry. The CV measurements were carried out with a scan rate of 10 mV s<sup>-1</sup> with an upper cutoff potential of -0.1 V vs. Hg/Hg<sub>2</sub>SO<sub>4</sub>, and the lower cutoff potential was defined as the potential at which the current reached 10 mA cm<sup>-2</sup> (based on the geometric surface area). Chronoamperometry measurements were performed by stepping the potential in 20 mV increments from 0 V vs. RHE to potentials where the current exceeded 10 mA cm<sup>-2</sup>. Each step was held for 3 minutes unless otherwise stated. All measurements were performed with a Bio-Logic SP-50e potentiostat.

## 2.4 Operando XRD measurements

Synchrotron XRD was performed at the Stanford Synchrotron Radiation Lightsource (SSRL) on beam line 2-1 in Bragg-Brentano geometry using an incident X-ray energy of 17 keV ( $\lambda = 0.72932 \text{ \AA}$ ). The instrument was equipped with a Pilatus 100 K area detector, and the 2D diffraction data was compiled and integrated using a python script specifically developed for angle-resolved measurements at SSRL beam line 2-1.

Slurry electrodes were used for the operando XRD measurements. Slurries consisted of 80 wt% active material, 10 wt% nafion, and 10 wt% acetylene black dispersed in ethanol using 3 rounds of mixing at 2000 rpm for 2.5 min in a Thinky mixer (Thinky USA, model AR-100). Slurries were cast on stainless steel mesh (Fisher Scientific) and carbon paper (Fuel Cell Earth) substrates with a final active material loading of 2–6 mg cm<sup>-2</sup> and dried at 50 °C. The mass loading and substrate are reported for each experiment. Electrodes were cycled using cyclic voltammetry and linear sweep voltammetry using a Bio-Logic SP150 potentiostat in an operando electrochemical cell (Fig. S1† based on a cell design reported in a previous study).<sup>33</sup> The reference electrode was a leakless Ag/AgCl reference electrode (eDAQ), and the counter electrode was carbon paper. The potential recorded from the leakless Ag/AgCl reference electrode was adjusted to the RHE scale by comparing the half wave potential of the WO<sub>3</sub>·H<sub>2</sub>O redox peak in a 3-electrode cell containing a calibrated Hg/Hg<sub>2</sub>SO<sub>4</sub> in saturated K<sub>2</sub>SO<sub>4</sub> reference electrode and graphite counter electrode. Each electrode was cycled in fresh 0.5 M H<sub>2</sub>SO<sub>4</sub> (ACS TraceMetal Grade, Fisher Scientific).

Operando diffraction patterns were collected at a fixed  $2\theta$  position every 1 second when the scan rate was 10 mV s<sup>-1</sup> and every 10 seconds when the scan rate was 1 mV s<sup>-1</sup>. Each  $2\theta$  position provided the diffraction pattern for  $2\theta$  values  $\pm \sim 3.4^\circ$ ,

and diffraction results for different  $2\theta$  ranges were collected during sequential cyclic voltammetry cycles using the same electrode. Results from multiple  $2\theta$  ranges were stitched together to obtain the full diffraction patterns used for structural refinement and analysis. Table 1 summarizes the  $2\theta$  range collected for each sample, along with the electrode substrate, slurry mass loading, scan rate, and potential range. Pawley refinements of the diffraction data were carried out using the TOPAS-Academic (Version 6) software package to determine the lattice parameters.<sup>34,35</sup>

## 2.5 Density functional theory (DFT) calculations

The geometry and electronic structure of anhydrous and hydrated tungsten oxides were studied by DFT calculations, using the periodic plane wave code Vienna *ab initio* simulation package (VASP 6.1).<sup>36,37</sup> Meta-generalized gradient approximation (meta-GGA) functionals, such as the strongly constrained and appropriately normed (SCAN) functional,<sup>38</sup> are more computationally demanding than GGA functionals but provide higher accuracy in properties including band gaps, lattice parameters, formation energies, and energy barriers of non-metals.<sup>39–41</sup> Therefore, we used the second version of the regularized SCAN (r<sup>2</sup>SCAN)<sup>42</sup> to treat electron exchange-correlation with the rVV10 non-local correlation functional ( $b = 11.95$ ) and to semi-empirically describe the van der Waals interactions.<sup>43,44</sup> We employed the “PAW 54” pseudopotentials within the projector augmented wave (PAW) framework to describe the ion–electron interaction.<sup>45</sup> The number of valence electrons treated were 14, 6, and 1, respectively for W, O, and H. For the geometry optimization of tungsten oxides, we sampled the Brillouin zone using Gamma-centered grids with constant spacing density of 0.3 Å<sup>-1</sup> applying Gaussian smearing ( $\sigma = 0.03 \text{ eV}$ ) of partial occupancies, with a plane-wave cut-off energy of 680 eV and energy and force convergence criteria of 10<sup>-6</sup> eV and 0.02 eV Å<sup>-1</sup>, respectively. Spin polarization and Harris–Foulke-type corrections for the calculation of forces were also considered.<sup>46</sup> All the geometries were relaxed using the conjugate gradient algorithm to minimize the total energy. The most stable optimized structures were selected after thorough examination of different hydrogen binding sites (*i.e.* binding configurations) and are presented in Fig. S2† with the lattice parameters in Table S1.† The theoretical bandgap and band structure of the tungsten oxides were evaluated through DFT calculations of fully relaxed structures, employing a denser spacing density of 0.1

**Table 1** Experimental conditions used for each material characterized by operando XRD. All  $2\theta$  values were obtained using an incident X-ray energy of 17 keV ( $\lambda = 0.72932 \text{ \AA}$ )

| Material                           | Mass loading (mg cm <sup>-2</sup> ) | Substrate | Technique    | Scan rate (mV s <sup>-1</sup> ) | Potential range (V vs. RHE) | $2\theta$ range (°) |
|------------------------------------|-------------------------------------|-----------|--------------|---------------------------------|-----------------------------|---------------------|
| WO <sub>3</sub>                    | 2                                   | SS        | CV           | 1                               | -0.308 to 0.642             | 9.63–30.10          |
| WO <sub>3</sub> ·H <sub>2</sub> O  | 6                                   | CP        | CV           | 10                              | -0.493 to 0.644             | 4.13–29.76          |
| WO <sub>3</sub> ·H <sub>2</sub> O  | 6                                   | CP        | Stepwise LSV | 10                              | -0.493 to 0.670             | 3.23–41.89          |
| WO <sub>3</sub> ·2H <sub>2</sub> O | 2                                   | SS        | CV           | 10                              | -0.358 to 0.640             | 2.83–26.61          |



$\text{\AA}^{-1}$  for Gamma-centered grids. Transition states were determined by relaxation of solid state nudged elastic band (ssNEB) with the climbing image (CI) method.<sup>47–50</sup>

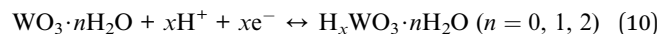
### 3. Results & discussion

#### 3.1 Activation of the HER on tungsten oxides

Anhydrous  $\text{WO}_3$  and hydrated tungsten oxides ( $\text{WO}_3 \cdot \text{H}_2\text{O}$ , and  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$ ) were selected for this study since all three materials undergo proton insertion-coupled electron transfer reactions and are therefore non proton-blocking electrodes in acidic electrolytes. Furthermore, they exhibit unique structural characteristics.  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$  contains layers of corner-sharing  $\text{WO}_5(\text{H}_2\text{O})$  octahedra separated by interlayer water molecules (Fig. 2c). Subsequent dehydration leads to the removal of the interlayer water and redistribution of bound water to form  $\text{WO}_3 \cdot \text{H}_2\text{O}$  (Fig. 2b) and eventually the bound water is removed and the interlayer collapses to form  $\text{WO}_3$  (Fig. 2a). This dehydration mechanism is supported by thermogravimetric analysis, which shows two sequential mass loss regions that correspond to the total removal of  $2\text{H}_2\text{O}$  from  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$ .<sup>25,30</sup> Powder XRD patterns for each hydration state (Fig. S3†) confirm the formation of monoclinic ( $P2_1/n$ )  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$ , orthorhombic ( $Pmnb$ )  $\text{WO}_3 \cdot \text{H}_2\text{O}$ , and monoclinic ( $P2_1/n$ )  $\text{WO}_3$ . SEM shows that the same morphology was maintained for each hydration state (Fig. S4†). The acid-precipitation reaction resulted in nanoplatelets with lateral dimensions on the order of 100–300 nm and thicknesses on the order of 10 s of nm (equivalent to  $\sim 20$ –100 inorganic layers in the hydrated materials). The lack of morphology and particle size changes in SEM suggests the surface areas are comparable between different hydration states.

Therefore, differences in electrochemical reactivity can be attributed to the degree of hydration and the local coordination environment of tungsten (layered vs. interconnected structure, and presence or absence of a coordinated water in the octahedra).

The electrochemical behavior of  $\text{WO}_3 \cdot n\text{H}_2\text{O}$  ( $n = 0, 1, 2$ ) was investigated using cyclic voltammetry of thin films. Fig. 2d–f show cyclic voltammograms of each material in a limited potential range where the current response is predominantly attributed to proton insertion coupled electron transfer, before the onset of the HER:



The shape of a CV can provide insights into the type of electrochemical reaction and its kinetics.<sup>52</sup>  $\text{WO}_3$  (Fig. 2d) shows an asymmetric CV response with three observable redox events (marked by 1/1', 2/2', and 3/3'). The peak potential separation is indicative of the reversibility of the electrochemical reaction. Based on this assessment, the 1/1' redox event with  $\Delta E_p \approx 100$  mV at  $10 \text{ mV s}^{-1}$  is more reversible as compared to the 2/2' redox event that has a  $\Delta E_p \approx 300$  mV at  $10 \text{ mV s}^{-1}$ . It is difficult to assign peak potentials to the broad peaks of the 3/3' redox event, however, this response is typical for a solid-solution insertion reaction (as discussed in the following sections). The CV of  $\text{WO}_3 \cdot \text{H}_2\text{O}$  (Fig. 2e) shows distinct redox events centered at 0.3 V (1/1'),  $-0.08$  V (2/2'), and  $-0.48$  V (3/3') all with a  $\Delta E_p \approx 80$  mV at  $10 \text{ mV s}^{-1}$ . Between the sharp redox peaks, there is also a non-negligible current response that can be attributed to solid-solution insertion of protons. The small peak potential separation suggests facile kinetics for proton insertion and de-insertion in  $\text{WO}_3 \cdot \text{H}_2\text{O}$ .  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$  (Fig. 2f) shows two redox events

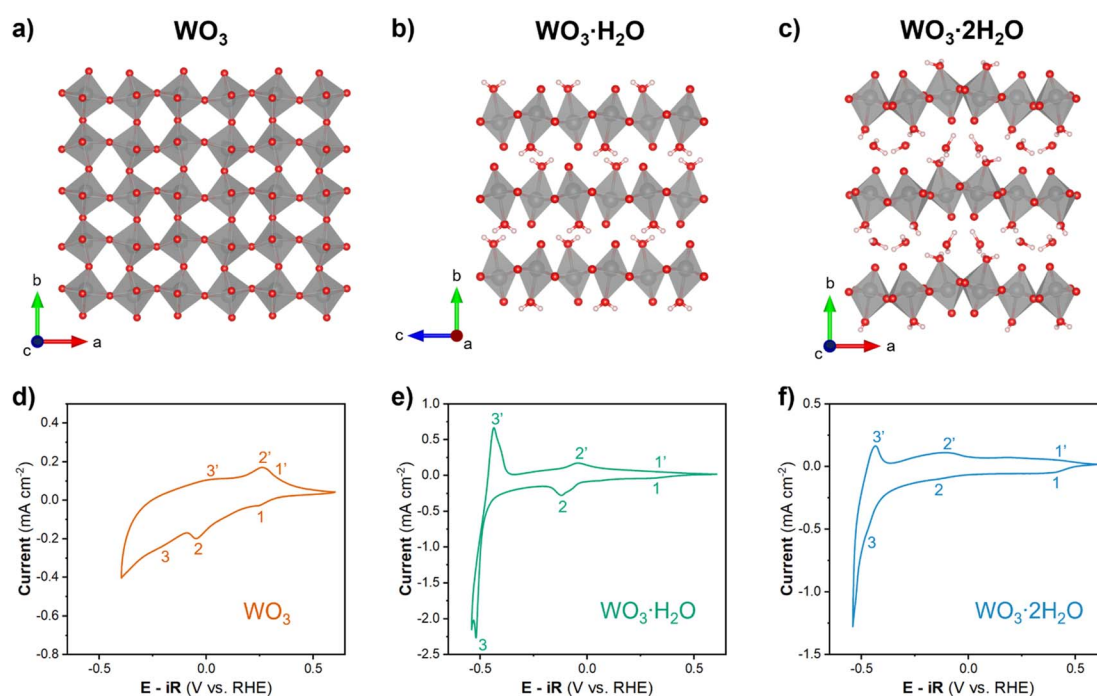


Fig. 2 Crystal structures of (a) monoclinic  $\text{WO}_3$ , (b) orthorhombic  $\text{WO}_3 \cdot \text{H}_2\text{O}$ , and (c) monoclinic  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$ .<sup>51</sup> Cyclic voltammograms of (d)  $\text{WO}_3$ , (e)  $\text{WO}_3 \cdot \text{H}_2\text{O}$ , and (f)  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$  with a rotation rate of 1600 rpm and scan rate of  $10 \text{ mV s}^{-1}$  in  $0.5 \text{ M H}_2\text{SO}_4$ .



(1/1' and 2/2') with broad current responses in the potential range of  $-0.4$  to  $0.5$  V vs. RHE and another redox event (3/3') centered at approximately  $-0.47$  V vs. RHE. The broad redox peaks make it difficult to quantify  $\Delta E_p$  for  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$ , but the approximate values for each redox event are less than 100 mV. Similar to  $\text{WO}_3 \cdot \text{H}_2\text{O}$ , this suggests favorable kinetics for proton insertion and de-insertion into  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$ . In our prior work, we hypothesized that the favorable kinetics of  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$  were due to its more limited structural deformation upon proton insertion.<sup>26,29</sup> The results from cyclic voltammetry reveal that each hydration state of  $\text{WO}_3 \cdot n\text{H}_2\text{O}$  has multiple redox events that take place prior to reaching potentials necessary to catalyze the HER, and the  $\Delta E_p$  trends show that proton insertion is more reversible in  $\text{WO}_3 \cdot \text{H}_2\text{O}$

and  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$  as compared to  $\text{WO}_3$ . In other words, removal of  $\text{H}^+/\text{e}^-$  is kinetically more facile in hydrated tungsten oxides compared to anhydrous tungsten oxide, a factor that may influence HER activity.

To study the HER activity of tungsten oxides, we decreased the cathodic potential limit during cyclic voltammetry, which led to a significant increase in cathodic current density (Fig. 3). In proton-blocking electrodes, the HER current measured during the cathodic and anodic voltammetry scans is equal.<sup>53</sup> However, for  $\text{WO}_3$  and  $\text{WO}_3 \cdot \text{H}_2\text{O}$ , the anodic scan exhibits higher magnitude HER currents at lower overpotentials than the cathodic scan. This crossover between the cathodic and anodic scans indicates an activation process taking place simultaneously with the onset of the HER. In other words, the material that forms at higher overpotentials has a higher activity compared to that which is present prior to the initial onset of HER. Conversely, the anodic and cathodic HER current of  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$  are nearly the same, suggesting there is no further activation process at high overpotentials. The overpotential for HER during the anodic and cathodic scans for each hydration state of tungsten oxide were determined by the potential required to reach  $\sim 10 \text{ mA cm}^{-2}$  based on the geometric area, as summarized in Table 2.  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$  exhibited the highest overpotential of 661 mV during both scans.  $\text{WO}_3 \cdot \text{H}_2\text{O}$  exhibits the lowest overpotential with a cathodic overpotential of 578 mV and anodic overpotential of 547.9 mV (a 30 mV improvement after further activation). The cathodic overpotential of  $\text{WO}_3$  was 626.6 mV and the anodic overpotential was 589.3 mV (a 37 mV improvement after further activation). The proton insertion reactions that precede the onset of the HER reveal that the tungsten oxides are dynamic electrocatalysts. Furthermore, the activation processes that lead to the crossover feature during cyclic voltammetry of  $\text{WO}_3$  and  $\text{WO}_3 \cdot \text{H}_2\text{O}$  are rate and potential dependent.

To ensure that material transformation was not the rate limiting step when determining the HER activity of the oxides, we obtained quasi-equilibrium HER currents using stepwise chronoamperometry. The potential was stepped in 20 mV intervals from 0 V to the onset of HER activity, and each step was held for 3 minutes. Holding the potential for 3 minutes allowed sufficient time for proton insertion reactions to complete, and the remaining plateau current was attributed to the HER. The average current from the last 30 s of each step was plotted as the absolute value of current density  $|j|$  on a log scale vs. overpotential in Fig. 4a. Although this method does not allow for a direct measurement of overpotential at  $10 \text{ mA cm}^{-2}$ , this value can be interpolated from the linear portion of the  $\log(|j|)$  vs.  $\eta$

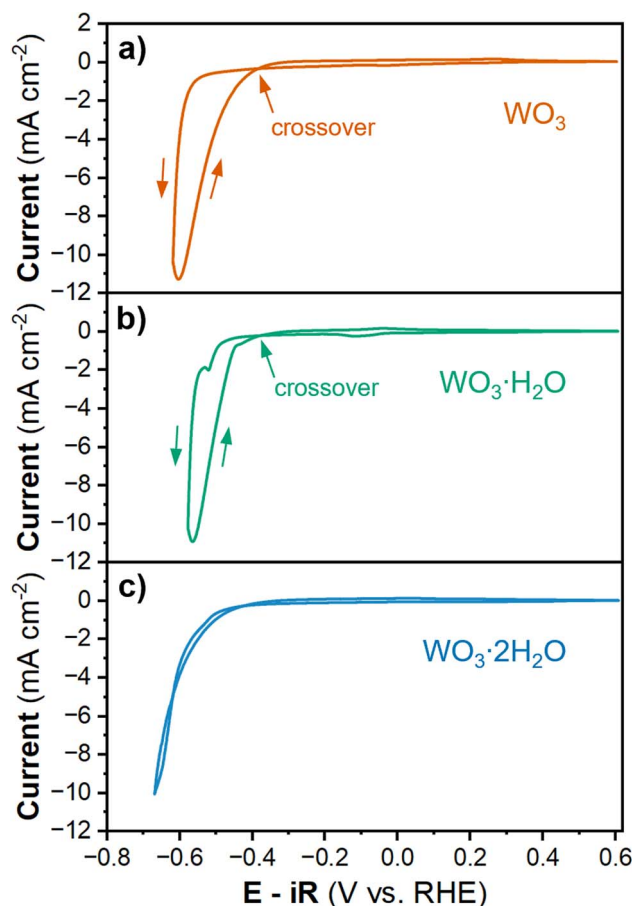


Fig. 3 Cyclic voltammograms of (a)  $\text{WO}_3$ , (b)  $\text{WO}_3 \cdot \text{H}_2\text{O}$ , and (c)  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$  where the cathodic potential limit was set to the potential where current reaches  $10 \text{ mA cm}^{-2}$ . The scan rate was  $10 \text{ mV s}^{-1}$  and rotation rate was 1600 rpm.

Table 2 List of overpotentials and currents for  $\text{WO}_3 \cdot n\text{H}_2\text{O}$  obtained from cyclic voltammetry and overpotentials and Tafel slopes from chronoamperometry

| Material                                | CV $\eta_{\text{cathodic}}$ (mV) | Current ( $\text{mA cm}^{-2}$ ) | CV $\eta_{\text{anodic}}$ (mV) | Current ( $\text{mA cm}^{-2}$ ) | $\Delta\eta$ (mV) | CA $\eta$ (mV) | Tafel slope ( $\text{mV dec}^{-1}$ ) |
|---|----------------------------------|---------------------------------|--------------------------------|---------------------------------|-------------------|----------------|--------------------------------------|
| $\text{WO}_3$                           | 626.6                            | -10.007                         | 589.3                          | -10.009                         | -37.3             | 531.7          | 49.4                                 |
| $\text{WO}_3 \cdot \text{H}_2\text{O}$  | 578.0                            | -10.098                         | 547.9                          | -10.002                         | -30.1             | 501.0          | 36.6                                 |
| $\text{WO}_3 \cdot 2\text{H}_2\text{O}$ | 661.0                            | -10.002                         | 661.3                          | -10.005                         | 0.3               | 621.5          | 95.8                                 |



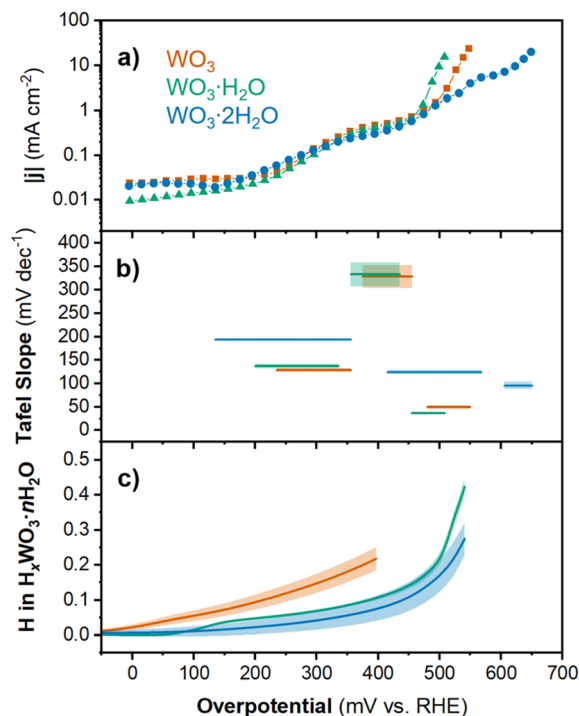


Fig. 4 Chronoamperometry results for WO<sub>3</sub>, WO<sub>3</sub>·H<sub>2</sub>O, and WO<sub>3</sub>·2H<sub>2</sub>O showing the (a) current density and (b) Tafel slope as a function of overpotential (shading represents the standard error of fit for each linear region). (c) Proton composition in H<sub>x</sub>WO<sub>3</sub>·nH<sub>2</sub>O ( $n = 0, 1, 2$ ) as a function of overpotential (calculated using method 2; see Discussion 1 in ESI†). Shading represents the standard deviation from the average of three separate electrodes.

plot. The overpotentials from quasi-equilibrium CA results were 531.7 mV for WO<sub>3</sub>, 501.0 mV for WO<sub>3</sub>·H<sub>2</sub>O, and 621.5 mV for WO<sub>3</sub>·2H<sub>2</sub>O. The overpotentials from CA were all lower than the cathodic or anodic overpotentials measured using cyclic voltammetry. Comparing  $\eta_{\text{cathodic}}$  from CV to  $\eta$  from CA, WO<sub>3</sub> exhibited the most significant improvement of 94.9 mV, followed by WO<sub>3</sub>·H<sub>2</sub>O at 77 mV and WO<sub>3</sub>·2H<sub>2</sub>O at 39.5 mV. We attribute this difference to the quasi-equilibrium conditions of the CA measurement.

The  $\log(|j|)$  vs. overpotential results for each material provide additional information about the electrochemical reactivity and kinetics of the HER. First, we observed that between 0 and 450 mV, the current response was similar for all three materials. The current was negligible up to  $\eta \approx 150$  mV. This was followed by an increase in current in the range of  $150 < \eta < 350$  mV and another plateau between  $350 < \eta < 450$  mV. Finally, the onset of HER occurred for each material at  $\eta > 450$  mV. The current response at intermediate overpotentials suggests that each material exhibits low HER activity until the onset of an activation process. In the high overpotential region ( $\eta > 450$  mV), we observe the onset of HER. The slope of the linear region (Tafel slope) provides mechanistic information about the reaction kinetics.<sup>54</sup> Generally, the lower the Tafel slope the faster the reaction. The materials showed the following Tafel slopes in the linear region where  $|j| > 1$  mA cm<sup>-2</sup>: 49.4 mV dec<sup>-1</sup> for H<sub>x</sub>WO<sub>3</sub>, 36.6 mV dec<sup>-1</sup> for H<sub>x</sub>WO<sub>3</sub>·H<sub>2</sub>O,

and 95.8 mV dec<sup>-1</sup> for H<sub>x</sub>WO<sub>3</sub>·2H<sub>2</sub>O (see ESI Fig. S5† for more details). These Tafel slopes suggest that the HER kinetics decrease in the order of H<sub>x</sub>WO<sub>3</sub>·H<sub>2</sub>O > H<sub>x</sub>WO<sub>3</sub> > H<sub>x</sub>WO<sub>3</sub>·2H<sub>2</sub>O.

The Tafel slope can describe the rate determining step in the HER. The first step in the HER for proton-blocking electrodes is the Volmer reaction with a characteristic Tafel slope of 120 mV dec<sup>-1</sup>. After the Volmer step, the HER can proceed by (1) the Tafel step with a Tafel slope of 30 mV dec<sup>-1</sup>, or (2) the Heyrovsky step with a Tafel slope of 40 mV dec<sup>-1</sup>.<sup>6,55,56</sup> It is also possible to observe Tafel slopes characteristic of a reaction that does not involve electron transfer, which in the HER would occur between the first and second electron transfer reactions and correspond to a Tafel slope of 60 mV.<sup>57</sup> Intermediate values between these “cardinal” values can exist, which are attributed to chemical rate determining steps during the HER. WO<sub>3</sub>·H<sub>2</sub>O has a Tafel slope close to 40 mV dec<sup>-1</sup>, which indicates that the second electron transfer during the HER is the rate determining step (Heyrovsky step following the convention on proton-blocking electrodes). The anhydrous WO<sub>3</sub> has a slightly higher Tafel slope close to 60 mV dec<sup>-1</sup>, suggesting that the rate determining step follows the first electron transfer but does not involve electron transfer. Finally, WO<sub>3</sub>·2H<sub>2</sub>O has a Tafel slope close to 120 mV dec<sup>-1</sup>, where the rate determining step can be attributed to the first electron transfer reaction (Volmer). It is important to note that factors such as the H coverage and symmetry factor of the reaction can influence the Tafel slope, therefore, these Tafel slope assignments should only be considered as the suggested rate determining steps. However, these results demonstrate that the HER kinetics and rate determining steps are dependent not only on the extent of proton insertion, but also the hydration state of tungsten oxide. The overpotentials and Tafel slopes for each hydration state of tungsten oxide were compared to the performance of commercial Pt/C and other non-precious metal catalysts in Table S2 and Fig. S6.† Whereas most catalysts with a high overpotential show high Tafel slopes, the tungsten oxides exhibit high overpotentials with low Tafel slopes, suggesting fast reaction kinetics once sufficient overpotential is reached. To determine the origin of the different onset potentials and Tafel slopes for each hydration state of tungsten oxide, we considered the proton composition, bulk structural transformations, and changes in the electronic structure of each material.

### 3.2 Proton composition as a function of potential

The proton composition of the oxide ( $x$  in H<sub>x</sub>WO<sub>3</sub>·nH<sub>2</sub>O;  $n = 0, 1, 2$ ) determines its structure and thus, properties during the HER. The protonated tungsten oxides, also called hydrogen bronzes, are effectively new electrocatalysts compared to the precatalyst states present in the uncycled electrode (the unprotonated oxides). We performed cyclic voltammetry on high mass loading (6 mg cm<sup>-2</sup>) electrodes where  $Q_{\text{insert}} \gg Q_{\text{HER}}$  in a wide potential range. The amount of proton insertion as a function of potential was determined using the corrected  $Q - Q_0$  approach (see Discussion 1 in ESI† for a detailed explanation of techniques used to calculate the proton composition). Fig. 5a, c and e show the CV results and proton compositions for H<sub>x</sub>WO<sub>3</sub>, H<sub>x</sub>WO<sub>3</sub>·H<sub>2</sub>O, and H<sub>x</sub>WO<sub>3</sub>·2H<sub>2</sub>O slurry electrodes in 0.5 M H<sub>2</sub>SO<sub>4</sub> at



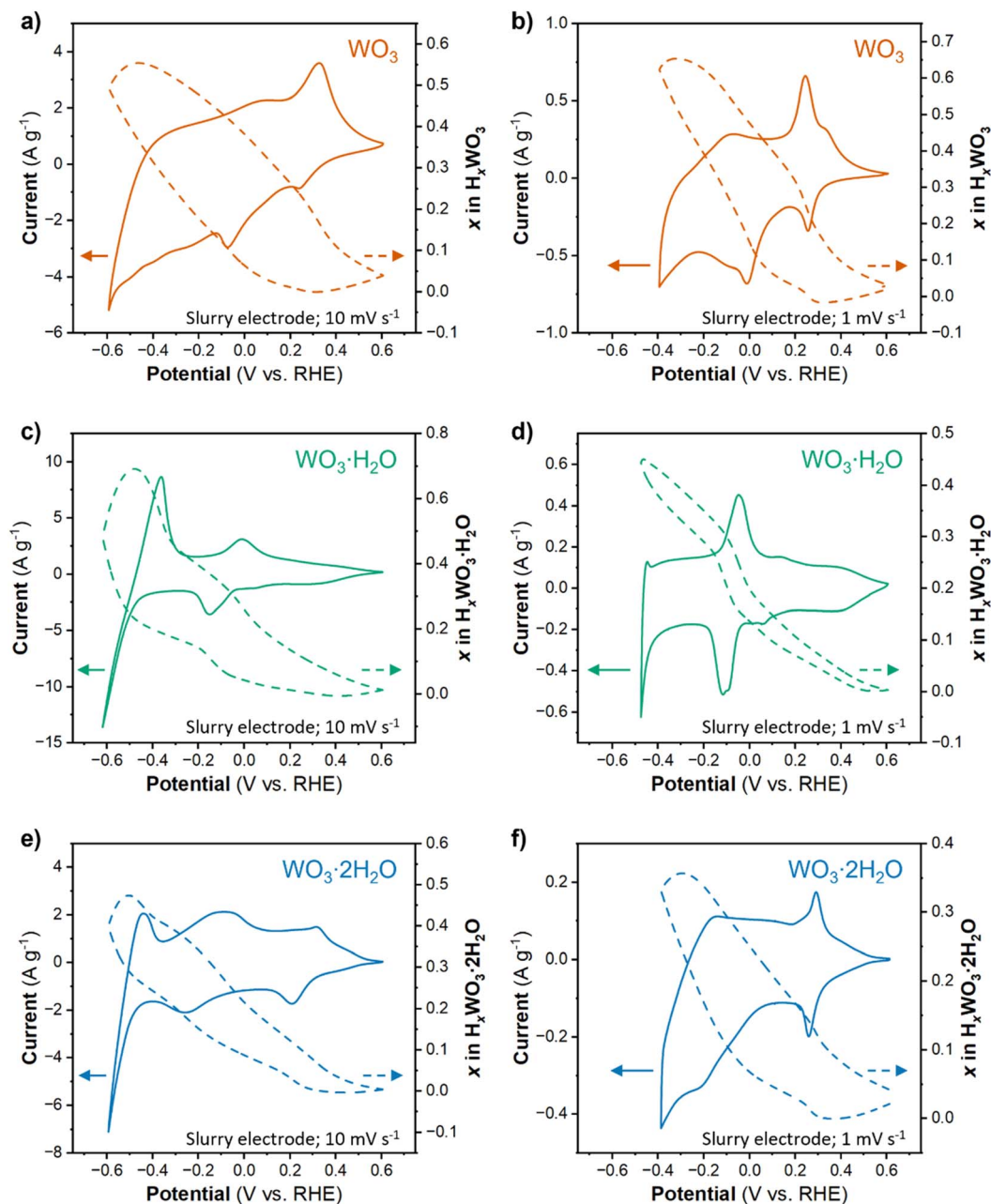


Fig. 5 CVs and proton composition as a function of potential for slurry electrodes of (a and b)  $\text{WO}_3$ , (c and d)  $\text{WO}_3 \cdot \text{H}_2\text{O}$ , and (e and f)  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$  at  $10 \text{ mV s}^{-1}$  (a, c and e) and  $1 \text{ mV s}^{-1}$  (b, d and f) in  $0.5 \text{ M H}_2\text{SO}_4$ . Solid lines indicate the current, and dashed lines indicate the proton composition as a function of potential.

$10 \text{ mV s}^{-1}$ . The maximum proton composition for each material was  $\text{H}_{0.55}\text{WO}_3$ ,  $\text{H}_{0.69}\text{WO}_3 \cdot \text{H}_2\text{O}$ , and  $\text{H}_{0.47}\text{WO}_3 \cdot 2\text{H}_2\text{O}$ .

The electrocatalytic measurements revealed a significant enhancement in HER activity upon switching from a non-equilibrium technique (cyclic voltammetry) to a quasi-equilibrium technique (chronoamperometry). Under quasi-equilibrium conditions, there is more time for proton insertion reactions to reach completion, which particularly impacted  $\text{WO}_3$ . Fig. 5b shows the CV and proton composition of  $\text{H}_x\text{WO}_3$  cycled at  $1 \text{ mV s}^{-1}$ . The maximum capacity reached  $75.3 \text{ mA h g}^{-1}$  (corresponding to  $\text{H}_{0.65}\text{WO}_3$ ) – a difference of  $0.1\text{H}^+$  per W compared to results from

a shorter timescale using a scan rate of  $10 \text{ mV s}^{-1}$ . We also cycled the hydrated tungsten oxides at  $1 \text{ mV s}^{-1}$  (Fig. 5d and f), however, at slow scan rates the final reduction event ( $3/3'$  in Fig. 2e and f) coincides with the onset of HER activity and complicates the quantification of proton insertion at high overpotentials. Although slower scan rates limit the accessible potential window, the slope of the proton composition as a function of potential for  $\text{WO}_3 \cdot \text{H}_2\text{O}$  and  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$  was greater when scan rate =  $1 \text{ mV s}^{-1}$  (Fig. 5d and f) compared to the results when scan rate =  $10 \text{ mV s}^{-1}$  (Fig. 5c and e). Two observations from this analysis of proton composition are (1) the highest HER activity correlates with the





materials containing the highest proton content at potentials close to the onset of the HER, and (2) longer reaction times during quasi-equilibrium conditions lead to higher HER activity (smaller overpotentials) in all three materials. Quantitatively, the results show that  $H_xWO_3 \cdot H_2O$  and  $H_xWO_3$  reach similar proton compositions between 0.69 and 0.65 $H^+$  per W, whereas  $H_xWO_3 \cdot 2H_2O$  only reaches  $\sim 0.5H^+$  per W. Further assessment of the structural and electronic properties is necessary to explain the HER activity trends, especially in the case of  $H_{0.69}WO_3 \cdot H_2O$  and  $H_{0.65}WO_3$  which have similar proton compositions.

### 3.3 Structural transformations as a function of potential/composition

Electrochemical proton insertion into tungsten oxides leads to changes in octahedral tilting and phase transformations. To understand the structures present in  $H_xWO_3 \cdot nH_2O$  ( $n = 0, 1, 2$ ) at the onset of HER activity, we performed operando synchrotron XRD during cyclic voltammetry.  $WO_3 \cdot H_2O$  and  $WO_3 \cdot 2H_2O$  were cycled at 10  $mV s^{-1}$  and  $WO_3$  was cycled at 1  $mV s^{-1}$  to ensure that all possible phase transformations could be observed since the proton insertion kinetics in the anhydrous material are slower than the layered hydrates. Fig. S8† shows

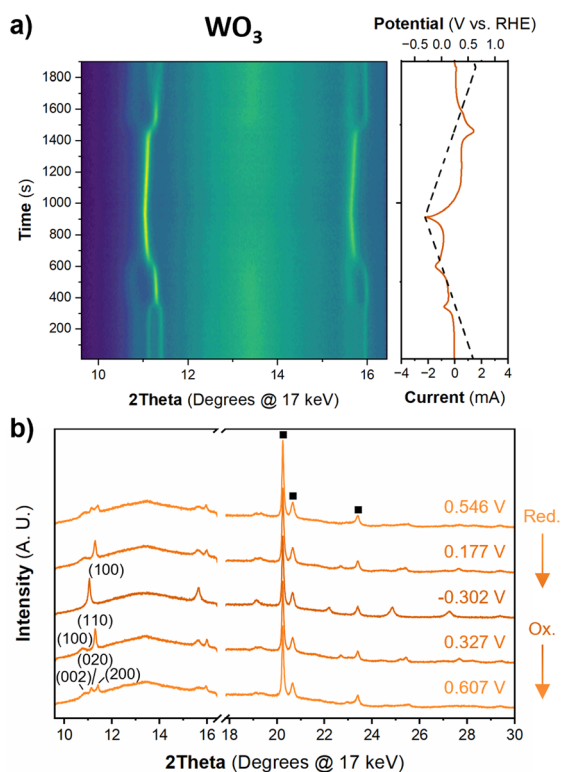


Fig. 6 (a) Operando XRD results for a  $WO_3$  electrode at 1  $mV s^{-1}$  in 0.5 M  $H_2SO_4$ . The contour plot shows diffraction patterns collected every 10 s (10 mV) and the right panel shows the potential and current as a function of time. (b) 1D diffraction patterns showing the full  $2\theta$  region collected and used for Pawley refinements of the structure in regions where no redox events were observed. The direction of the CV scan for the 1D diffraction patterns is marked by reduction (Red.) or oxidation (Ox.). Diffraction peaks associated with the stainless steel mesh substrate are marked by black squares.

the CVs for each electrode in the operando electrochemical cell, in agreement with the features observed using the RDE (Fig. 2d–f) and slurry electrodes in a 3-electrode cell (Fig. 5).

Fig. 6 shows the diffraction patterns and corresponding potential and current as a function of time for  $WO_3$  during cyclic voltammetry in 0.5 M  $H_2SO_4$  at 1  $mV s^{-1}$ . After the first proton insertion event with a peak potential ( $E_p$ ) of 0.26 V vs. RHE, the peaks originally associated with the (002), (020), and (200) indices of the monoclinic cell at 10.84, 11.14, and 11.40°  $2\theta$ , respectively, merge into two peaks indicating a structural transition to higher symmetry. The 1D XRD pattern collected at 0.177 V (vs. RHE) can be well fit to the tetragonal  $P4/nm$  space group symmetry with no observable tilting of the octahedra while the antiparallel displacements of  $W^{6+}$ -ions remain.<sup>58</sup> After the second proton insertion event at  $E_p = 0.0$  V vs. RHE, the (100) and (110) reflections positioned at 10.76 and 11.29°  $2\theta$ , respectively, merge into a single peak suggesting another transition to higher symmetry. The 1D pattern collected under an applied potential of  $-0.30$  V (vs. RHE) was fit to a cubic unit cell with  $Pm\bar{3}m$  space group symmetry.<sup>59</sup> The perovskite-like structure does not demonstrate any observable tilts, nor does it allow for the displacement of the  $W^{6+}$ -ion. The broad current response between  $-0.1$  and  $-0.28$  V vs. RHE is attributed to proton insertion into cubic  $H_xWO_3$  without further structural transitions (solid solution proton insertion). Fig. 7 shows potential dependent proton composition and *pseudo*-cubic *a*-lattice parameter ( $a_p$ ) calculated from the lattice parameters of each phase using eqn (S2)–(S4).† After the pure cubic structure forms,  $a_p$  continues to increase from 3.770 Å at  $-0.083$  V to

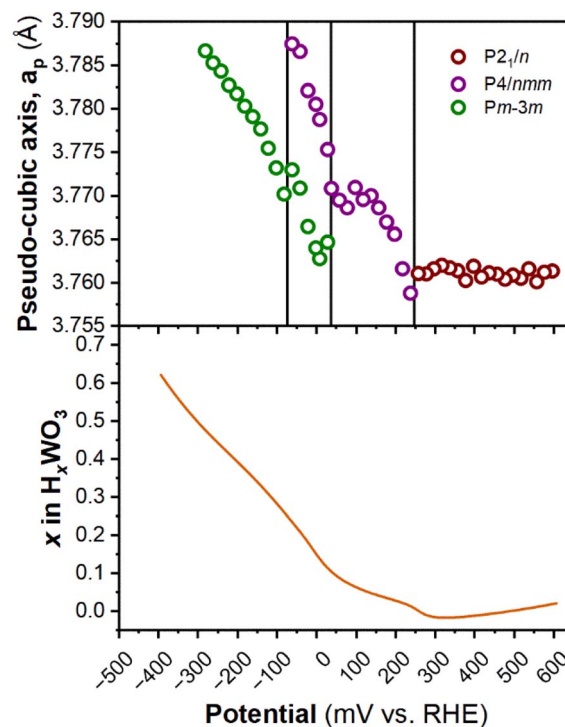


Fig. 7 *Pseudo*-cubic parameter *a* (from operando XRD), and proton composition of  $WO_3$  (from Fig. 5b) as a function of potential during a cathodic cyclic voltammetry scan in 0.5 M  $H_2SO_4$  at 1  $mV s^{-1}$ .



3.787 Å at  $-0.282$  V vs. RHE. The same series of transitions occurred in reverse order as protons de-inserted during the anodic cyclic voltammetry scan from  $-0.31$  V to  $0.64$  V vs. RHE.

The structure of  $\text{WO}_3 \cdot \text{H}_2\text{O}$  was reported previously by Szymanski and Roberts as a primitive orthorhombic cell with  $Pmnb$  space group symmetry (see Discussion 2 in ESI† for more details on the structural refinements).<sup>60</sup> Fig. 8 shows the operando XRD data for  $\text{WO}_3 \cdot \text{H}_2\text{O}$  at  $10 \text{ mV s}^{-1}$ . Multiple reduction and oxidation events were observed (right panel in Fig. 8a), however, the diffraction results do not indicate any solid-state phase transitions associated with these events. Moreover, no morphology changes were observed *via ex situ* SEM (Fig. S9†) between the pristine and cycled slurry electrode, when cycled to a cathodic cutoff potential of  $-0.8$  V vs. RHE, suggesting that no dissolution and redeposition process is taking place. Instead, Pawley refinements of 1D diffraction patterns at various potentials (Fig. 8b) could be modeled by a contraction of the interlayer distance and in-plane expansion of the inorganic layers with increasing proton content (Fig. 9). To verify this observation, we turned to DFT calculations. We found the contraction and expansion trends to align with the DFT-computed lattice parameters shown in Table S1.† This behavior is also in agreement with prior literature that reported similar trends up to a proton composition of  $\text{H}_{0.12}\text{WO}_3 \cdot \text{H}_2\text{O}$ .<sup>61</sup>

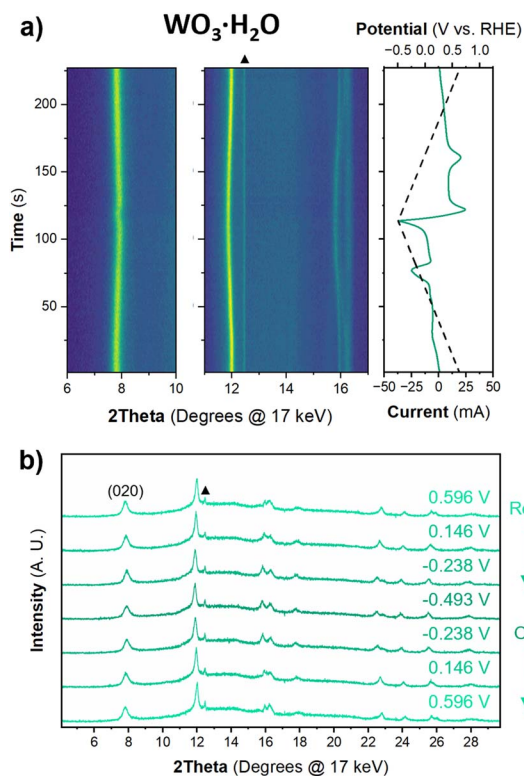


Fig. 8 (a) Operando XRD for  $\text{WO}_3 \cdot \text{H}_2\text{O}$  at  $10 \text{ mV s}^{-1}$  in  $0.5 \text{ M H}_2\text{SO}_4$ . (b) 1D diffraction patterns showing the full  $2\theta$  region collected and used for Pawley refinements of the structure before and after redox features in the CV. The direction of the CV scan for the 1D diffraction patterns is marked by reduction (Red.) or oxidation (Ox.). The diffraction peak associated with the carbon paper substrate is marked with a triangle.

The lack of phase transformations is consistent with the faster proton (de)insertion kinetics compared to  $\text{WO}_3$  observed during cyclic voltammetry.

The structure of  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$  was reported by Li *et al.*, in 2000 from powder XRD. It has a monoclinic cell with  $P2_1/n$  symmetry with a complex pattern of octahedral tilting.<sup>62</sup> During electrochemical proton insertion, operando XRD results (Fig. 10) show that the  $(20\bar{2})$  and  $(202)$  reflections at  $11.16$  and  $11.34^\circ 2\theta$ , respectively, merge as the potential is scanned from  $0.642$  V to  $-0.05$  V vs. RHE, and they proceed to split at higher proton compositions between  $-0.05$  V and  $-0.358$  V vs. RHE. This trend is also observed with the  $(22\bar{2})$  and  $(222)$  reflections at  $12.70$  and  $12.88^\circ 2\theta$ , respectively. Furthermore, the  $(020)$  peak position remains constant, indicating that the interlayer spacing does not change during proton (de)insertion. The opposite transitions occur during deprotonation in the anodic scan of the CV. Pawley refinements were performed on 1D diffraction patterns of the full  $2\theta$  range recorded during the operando measurements of  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$ . All the collected diffraction patterns could be modeled with the monoclinic  $P2_1/n$  structure with changes in the  $a$ - and  $c$ -lattice parameters, with the most significant change in the  $a$ -axis (Fig. 11). Overall, upon proton insertion and de-insertion in  $\text{H}_x\text{WO}_3 \cdot 2\text{H}_2\text{O}$ , the structural transitions are limited to the  $ac$ -plane and may be associated with changes in the degree of tilting and changes in bond lengths within the inorganic layers.

Overall, these results show that  $\text{WO}_3$  undergoes phase transformations from monoclinic to tetragonal and tetragonal

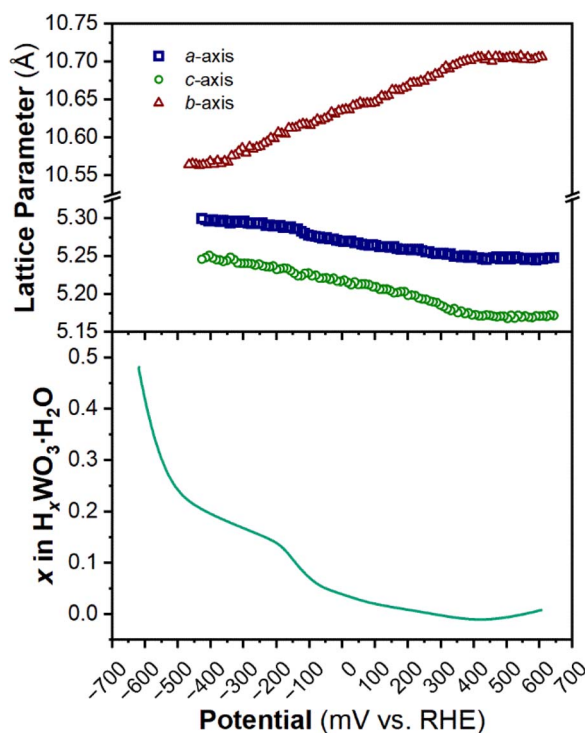


Fig. 9 Trend of lattice parameters (from Pawley refinements of the operando XRD) and proton composition in  $\text{H}_x\text{WO}_3 \cdot \text{H}_2\text{O}$  (from Fig. 5c) as a function of potential during the cathodic cyclic voltammetry scan in  $0.5 \text{ M H}_2\text{SO}_4$  at  $10 \text{ mV s}^{-1}$ .



to cubic accompanied by the loss of octahedral tilting and second order Jahn–Teller distortions upon proton insertion. Conversely, proton insertion into tungsten oxide hydrates causes changes in the degree of tilting without any observed phase transformations.  $\text{WO}_3 \cdot \text{H}_2\text{O}$  trends toward higher symmetry, based on the observation of some diffraction peaks beginning to merge, while  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$  trends toward higher symmetry at low proton compositions but reverts back to lower symmetry and a higher degree of tilting at high proton compositions.

### 3.4 Origin of hydration state dependence of HER activity on protonated tungsten oxides

Taking together the electrocatalytic activity, proton composition, and structure of each electrocatalyst in this study, it is possible to consider the origin of the apparent hydration state dependence of the HER activity of tungsten oxides. The highest HER activity was observed in  $\text{H}_x\text{WO}_3 \cdot \text{H}_2\text{O}$ , followed by  $\text{H}_x\text{WO}_3$ , and finally  $\text{H}_x\text{WO}_3 \cdot 2\text{H}_2\text{O}$  with a significantly higher overpotential and Tafel slope. Here we discuss several factors that may influence electrocatalytic activity: (1) surface area, (2) electronic conductivity, (3) hydrogen adsorption energy, and (4)

participation of inserted protons. Finally, we propose a mechanism for the HER activity differences between each tungsten oxide material characterized here.

We first consider the influence of surface area: the larger the electrochemical interface, the higher the areal current on a geometric basis. SEM images in Fig. S4† show that the particle size and morphology is independent of the hydration state of  $\text{WO}_3 \cdot n\text{H}_2\text{O}$  ( $n = 0, 1, 2$ ), therefore, we do not expect there to be differences in the electrocatalytically active surface area. Second, we consider the influence of the electronic conductivity: electron transport is necessary for the interfacial electrochemical reaction to take place. Electronic conductivity is dependent on the crystal structure and proton composition. Proton insertion is accompanied by electrons that partially fill the W 5d orbitals of the conduction band, leading to a semiconductor-to-metal transition at low proton compositions in tungsten oxides.<sup>25,63–65</sup> By conducting DFT calculations, we observed that the band gap diminishes to zero for all three tungsten oxides when they are protonated, as depicted in Fig. 12. The band structure analysis clearly reveals that proton-insertion coupled electron transfer partially populates energy bands within the conduction band of the tungsten oxides, as indicated by the presence of multiple conduction bands at the Fermi level. For  $\text{WO}_3$  and  $\text{WO}_3 \cdot \text{H}_2\text{O}$ , the theoretical band gaps of their pristine (non-protonated) structures exceeded 1.9 eV, which represents a moderate underestimation compared to the experimentally observed band gap ( $\text{WO}_3$ ; 2.6 eV,  $\text{WO}_3 \cdot \text{H}_2\text{O}$ ; 2.17).<sup>66–68</sup> Note that this discrepancy is a typical observation

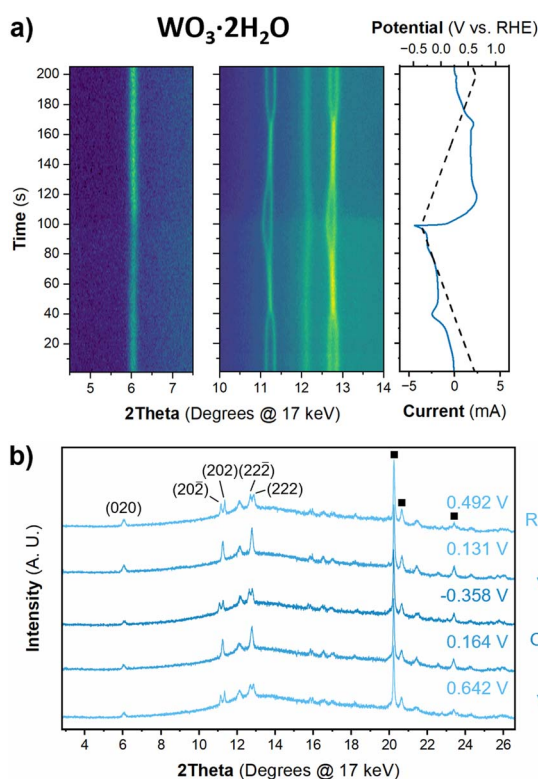


Fig. 10 (a) Operando XRD results for a  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$  slurry electrode at  $10 \text{ mV s}^{-1}$  in  $0.5 \text{ M H}_2\text{SO}_4$ . The contour plot shows diffraction patterns collected every 1 s ( $10 \text{ mV}$ ) and the right panel shows the potential and current as a function of time. (b) 1D diffraction patterns showing the full  $2\theta$  region collected and used for Pawley refinements of the structure before and after redox features in the CV. The direction of the CV scan for the 1D diffraction patterns is marked by reduction (Red.) or oxidation (Ox.). The diffraction peaks associated with the stainless steel mesh substrate are marked with squares.

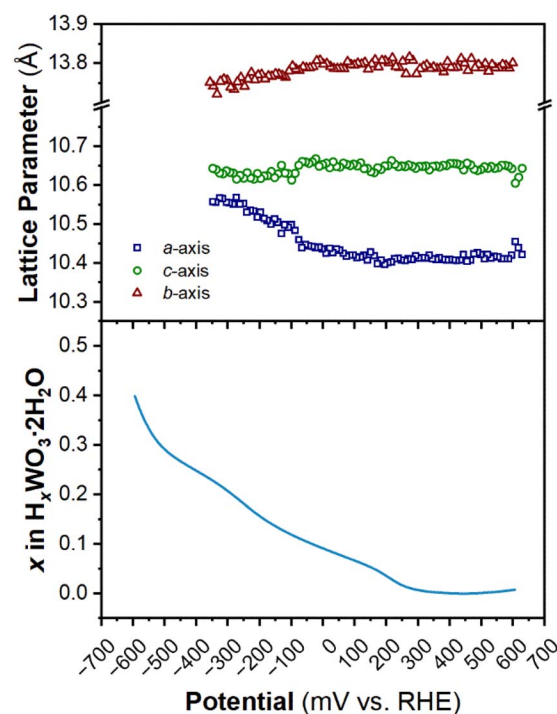


Fig. 11 Trend of lattice parameters (from Pawley refinements of the operando XRD) and proton composition in  $\text{H}_x\text{WO}_3 \cdot 2\text{H}_2\text{O}$  (from Fig. 5e) as a function of potential during the cathodic cyclic voltammetry scan in  $0.5 \text{ M H}_2\text{SO}_4$  at  $10 \text{ mV s}^{-1}$ .



associated with the use of meta-GGA functionals, as is the prediction of a direct bandgap of  $\text{WO}_3$ .<sup>69</sup> The energy difference between the direct and indirect band gaps is very small ( $\sim 0.01$  eV) with hybrid functionals. However, the band gap width most closely aligns with experimental results and is similarly reflected in our calculations. Notably, the protonation of  $\text{WO}_3$  and  $\text{WO}_3 \cdot \text{H}_2\text{O}$  to  $\text{H}_{0.5}\text{WO}_3$  and  $\text{H}_{0.5}\text{WO}_3 \cdot \text{H}_2\text{O}$  leads to a semiconductor-to-metal transition. On the other hand, protonated  $\text{H}_{0.5}\text{WO}_3 \cdot 2\text{H}_2\text{O}$  exhibits two thermodynamically competitive configurations ( $< 0.001$  eV of energy difference), both involving hydrogen binding to bridging oxygen sites. One configuration exhibits a vanishing band gap, while the other remains a semiconductor with a band gap of 0.43 eV (Fig. S10<sup>†</sup>). This observation provides a plausible explanation for the relatively lower electrocatalytic activity observed in  $\text{H}_x\text{WO}_3 \cdot 2\text{H}_2\text{O}$  compared to  $\text{H}_x\text{WO}_3$  and  $\text{H}_x\text{WO}_3 \cdot \text{H}_2\text{O}$ . Orthorhombic  $\text{H}_{0.69}\text{WO}_3 \cdot \text{H}_2\text{O}$  and cubic  $\text{H}_{0.65}\text{WO}_3$  exhibit similar maximum observable proton compositions, have relatively high symmetry crystal structures, and DFT calculations confirm that both undergo the semiconductor-to-metal transition. Further inspection of the

calculated band structures of Fig. 12 reveals that the  $\text{H}_{0.5}\text{WO}_3 \cdot \text{H}_2\text{O}$  exhibits a wider dispersion of bands than the  $\text{H}_{0.5}\text{WO}_3$ , which could facilitate higher electronic conductivity.<sup>70</sup> Although DFT does not provide a quantitative measure of electronic conductivity, the band structure analysis is consistent with the experimentally observed HER trends and could possibly provide a link between the oxide's conductivity and catalytic activity.

Third, we consider the hydrogen adsorption energy. The hydrogen adsorption energy on tungsten oxides is dependent on the proton composition. Miu *et al.* demonstrated computationally that  $\text{H}_{0.5}\text{WO}_3$  and  $\text{H}_{0.625}\text{WO}_3$  were responsible for HER activity and that the most energetically favorable hydrogen adsorption site becomes the W metal center as opposed to a bridging or terminal oxygen (shown schematically in Fig. 13).<sup>31</sup> Our results show that  $\text{H}_x\text{WO}_3$  and  $\text{H}_x\text{WO}_3 \cdot \text{H}_2\text{O}$  reach proton compositions greater than 0.5, therefore, these electrocatalysts are expected to have reactive surface sites.

Finally, we consider the role of the inserted protons. A possible explanation for the origin of the hydration state

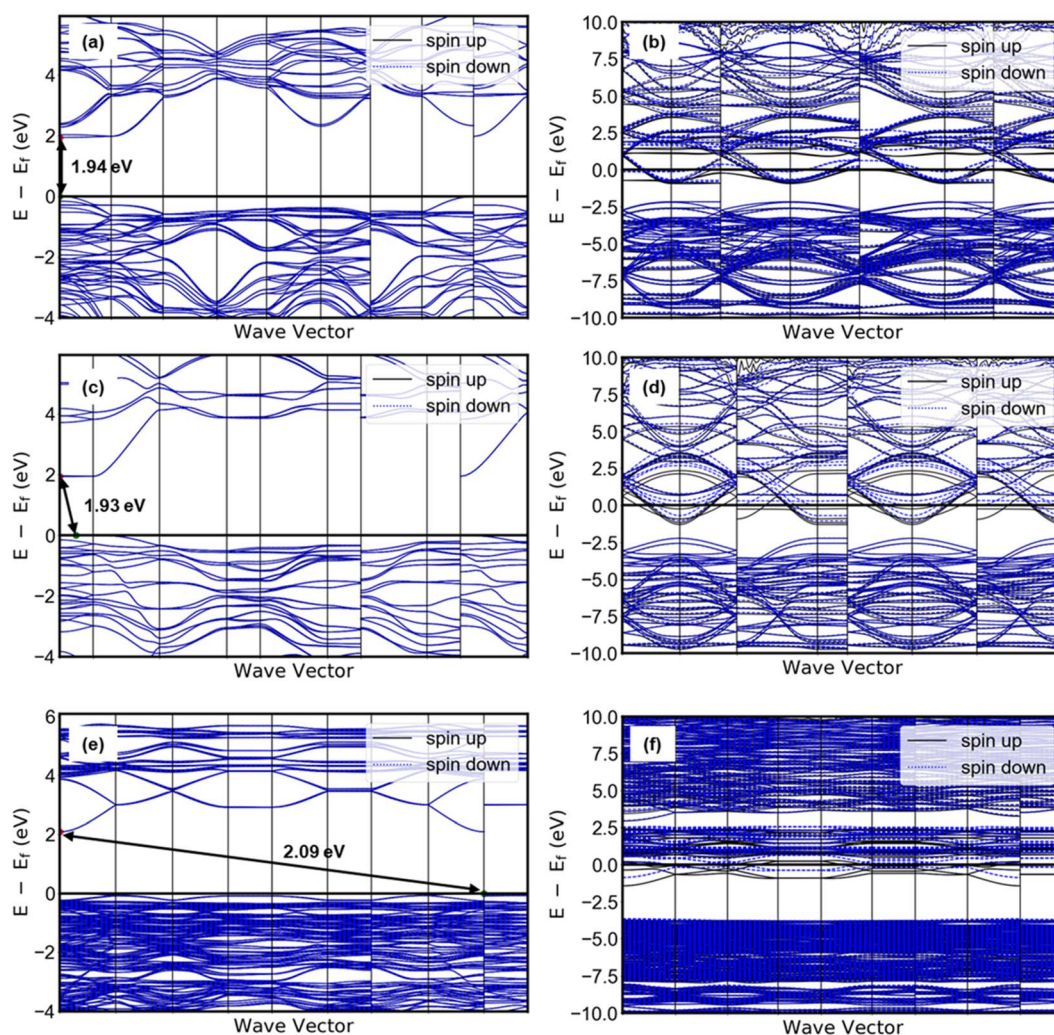


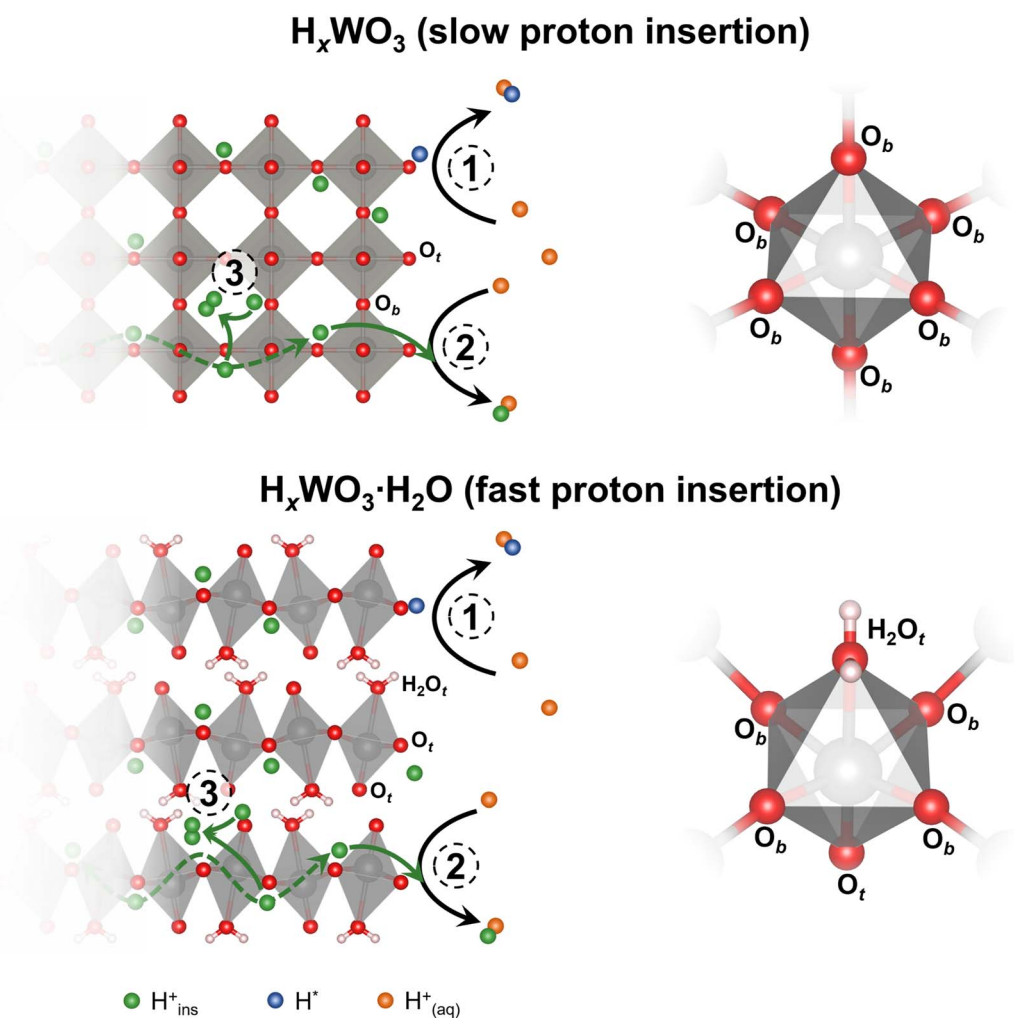
Fig. 12 Band structure of pristine (left) and protonated (right) tungsten oxides: (a)  $\text{WO}_3$ , (b)  $\text{H}_{0.5}\text{WO}_3$ , (c)  $\text{WO}_3 \cdot \text{H}_2\text{O}$ , (d)  $\text{H}_{0.5}\text{WO}_3 \cdot \text{H}_2\text{O}$ , (e)  $\text{WO}_3 \cdot 2\text{H}_2\text{O}$ , (f)  $\text{H}_{0.5}\text{WO}_3 \cdot 2\text{H}_2\text{O}$ .



dependence of HER activity on tungsten oxides is that they are non proton-blocking electrocatalysts. Prior literature demonstrates that bulk protons can participate in surface reactions. For example, during coulometry to study proton conduction in solids, bulk protons are reduced and released as hydrogen gas.<sup>71</sup> Coulometric measurements performed on potassium dihydrogen phosphate and lithium hydrazinium sulfate revealed that the volume of hydrogen gas was linearly proportional to the charge passed through the crystals.<sup>71</sup> This mechanism is similar to the concept of exsolution of metal nanoparticles on perovskite transition metal oxides in reducing conditions, with the difference being that the exsolved substance here, hydrogen, is a gas.<sup>72</sup> In non-aqueous electrolytes, proton-containing transition metal oxides such as  $\text{H}_2\text{Ti}_6\text{O}_{13}$  and  $\text{H}_2\text{WO}_4$  can convert to lithium or sodium-containing TMOs, with protons proposed to leave the structure in the form of hydrogen.<sup>73–75</sup> Sub-surface protons are also shown to be more reactive than surface

adsorbed protons, leading to the participation of sub-surface protons in surface catalytic reactions.<sup>76,77</sup> Similarly, a reverse hydrogen spillover mechanism was proposed for a  $\text{Ru}@\text{H}_x\text{WO}_3$  electrocatalyst, whereby protons transferred from  $\text{H}_x\text{WO}_3$  to Ru, thus increasing the hydrogen coverage and the overall HER activity.<sup>78</sup>

We hypothesize that bulk, electrochemically inserted protons participate in the HER on  $\text{H}_x\text{WO}_3 \cdot n\text{H}_2\text{O}$  electrocatalysts. In this scenario, we propose three possible reaction mechanisms for the formation of hydrogen (shown schematically in Fig. 13): (1) combination of two aqueous protons at a surface active site following the Volmer–Heyrovsky or Volmer–Tafel steps (eqn (1)–(3)); (2) combination of one inserted proton with one aqueous proton that each transport to a surface active site (eqn (5) or (6)); and (3) combination of two inserted protons at a confined active site to form  $\text{H}_2$  confined in the oxide lattice (eqn (7)). The participation of bulk



**Fig. 13** Proposed HER mechanisms on proton insertion electrocatalysts  $\text{H}_x\text{WO}_3$  (top) and  $\text{H}_x\text{WO}_3 \cdot \text{H}_2\text{O}$  (bottom): (1) combination of two aqueous protons on active sites at the electrocatalyst/electrolyte interface following the Volmer–Heyrovsky or Volmer–Tafel steps; (2) combination of one inserted proton with one aqueous proton that each migrate to an active site at the electrocatalyst/electrolyte interface; and (3) combination of two inserted protons at a confined active site to form  $\text{H}_2$  confined in the oxide lattice. Isolated octahedra show the different W coordination environments in the bulk lattice of anhydrous and hydrated tungsten oxides:  $\text{H}_x\text{WO}_3$  contains W coordinated with 6 bridging oxygens ( $\text{O}_b$ ), whereas  $\text{H}_x\text{WO}_3 \cdot \text{H}_2\text{O}$  contains W coordinated to 4 bridging oxygens, 1 terminal oxygen ( $\text{O}_t$ ) and 1 terminal water molecule ( $\text{H}_2\text{O}_t$ ). Electrochemically inserted protons (green spheres), aqueous protons (orange spheres), adsorbed hydrogen (blue spheres).



protons in the HER means that electrochemically protonated tungsten oxides serve as a proton reservoir for the HER, whereby the kinetics of proton transport from the bulk insertion site to the electrocatalytically active site will influence the HER kinetics/activity. Our DFT calculations suggest that the surface mediated mechanisms are likely to be dominating the HER of protonated tungsten oxides, albeit, with the involvement of lattice protons. As depicted in Fig. S11,† proton conduction in the bulk is highly favored due to barrierless activation for  $\text{H}_{0.5}\text{WO}_3$  and very low activation (0.14 eV) for  $\text{H}_{0.5}\text{WO}_3 \cdot \text{H}_2\text{O}$ . These low barriers for bulk proton conduction are accessible at room temperature. We note that in the case of  $\text{H}_{0.5}\text{WO}_3$ , the barrierless activation is a result of symmetric oxygen sites (due to the specific 0.5 stoichiometry in hydrogen) that exchange protons (Fig. S11a†) and we should not generalize that the anhydrous oxide would show a barrierless proton conduction (e.g. at different degrees of protonation). However, low activation energies for hydrogen spillover in bulk have been reported in  $\text{WO}_3$ ,<sup>79</sup> in agreement with our observations of facile proton conduction. In contrast, Fig. S12† demonstrates that hydrogen formation within the oxide lattice presents activation barriers of 2.25 and 1.98 eV, for  $\text{H}_{0.5}\text{WO}_3$  and  $\text{H}_{0.5}\text{WO}_3 \cdot \text{H}_2\text{O}$  respectively (additional higher energy barriers for  $\text{H}_{0.5}\text{WO}_3 \cdot \text{H}_2\text{O}$  are presented in Fig. S13†). These barriers are significantly higher than previously reported energy barriers of 0.50 eV for water-mediated hydrogen formation on the surface of  $\text{H}_{0.5}\text{WO}_3$ .<sup>31</sup> For  $\text{H}_{0.69}\text{WO}_3 \cdot \text{H}_2\text{O}$  and  $\text{H}_{0.65}\text{WO}_3$ , we further attribute the difference in activity to the difference in bulk proton transport kinetics. Dickens *et al.* experimentally determined the proton diffusion coefficients of  $\text{H}_x\text{WO}_3$  and  $\text{H}_x\text{WO}_3 \cdot \text{H}_2\text{O}$  (where  $0 < x < 0.2$ ) using a current-pulse relaxation method.<sup>80</sup> Proton transport in  $\text{H}_x\text{WO}_3 \cdot \text{H}_2\text{O}$  is fast ( $D \sim 10^{-6} \text{ cm}^2 \text{ s}^{-1}$ ) due to the low energy transport pathways that exist in the layered structure. However, proton transport in  $\text{H}_x\text{WO}_3$  is slower ( $D \sim 10^{-11} \text{ cm}^2 \text{ s}^{-1}$ ). Although both structures exhibit low activation energies for proton transport as shown in the DFT calculations, the monohydrated structure exhibits faster proton transport kinetics possibly due to the multiple diffusion pathways facilitated by the hydrogen bonds of lattice water (Fig. S11†). Based on this hypothesized mechanism, we propose that the difference in proton transport kinetics of electrochemically inserted protons can further justify the difference in the HER activity of  $\text{H}_{0.65}\text{WO}_3$  and  $\text{H}_{0.69}\text{WO}_3 \cdot \text{H}_2\text{O}$ , where faster proton transport in  $\text{H}_{0.69}\text{WO}_3 \cdot \text{H}_2\text{O}$  leads to a larger HER current contribution from the inserted protons and a higher overall HER activity.

## 4. Conclusions

We systematically studied the HER activity of three hydration states of tungsten oxide ( $\text{WO}_3 \cdot n\text{H}_2\text{O}$ ;  $n = 0, 1, 2$ ), and discussed the possible origin of the hydration state dependent electrocatalytic activity of these proton insertion type electrocatalysts. Cyclic voltammetry revealed the critical role of proton insertion in the activation of tungsten oxides for the HER. Stepwise chronoamperometry measurements were used to measure the electrocatalytic activity of the equilibrium composition of each

electrocatalyst.  $\text{H}_{0.69}\text{WO}_3 \cdot \text{H}_2\text{O}$  was the most active electrocatalyst with an overpotential of 501.0 mV, followed by  $\text{H}_{0.65}\text{WO}_3$  ( $\eta = 531.7$  mV) and  $\text{H}_{0.47}\text{WO}_3 \cdot 2\text{H}_2\text{O}$  ( $\eta = 621.5$  mV). Using slurry electrodes, we calculated the equilibrium proton composition in the maximum accessible potential window, and we found that proton composition scaled positively with HER activity. Operando XRD measurements revealed the structural changes that occur during proton (de)insertion using cyclic voltammetry. Upon proton insertion into  $\text{H}_x\text{WO}_3$ , the monoclinic structure transitions to tetragonal and then to cubic. Conversely,  $\text{H}_x\text{WO}_3 \cdot \text{H}_2\text{O}$  (orthorhombic) and  $\text{H}_x\text{WO}_3 \cdot 2\text{H}_2\text{O}$  (monoclinic) do not undergo phase transitions. These results show that higher symmetry corresponds with higher HER activity. The trends in activity, proton composition, and structural symmetry suggest that the bulk properties of these electrocatalysts can directly influence the electrocatalytic activity. Furthermore, we propose that bulk, electrochemically inserted protons can participate in the HER, whereby the inserted protons serve as a proton reservoir and thus the proton transport properties of each hydration state can explain the differences in electrocatalytic activity. The band structure calculated by DFT suggests a semiconductor-to-metal transition in the electronic structure of all tungsten oxides by proton intercalation and a competition from a state with a gap in the dihydrate case, rationalizing the experimental HER catalytic activity trends. In addition, DFT calculations demonstrate that proton transport in bulk of the oxides is facile with very small activation barriers. The outcomes of this study emphasize the importance of systematically studying dynamic processes that take place on non proton-blocking electrocatalysts. Furthermore, when an inserting species also serves as a reactant in the electrocatalytic reaction of interest, it is possible that the reactant directly participates in the reaction thus contributing to the observed activity.

## Data availability

Data from this work is openly available at <https://doi.org/10.5281/zenodo.10802721>.

## Author contributions

M. A. S.: conceptualization, investigation, formal analysis, methodology, writing, review & editing; N. P. H.: conceptualization, formal analysis, writing, review & editing; K.-E. Y.: investigation, formal analysis, writing; G. M.: funding acquisition, supervision, resources, writing, review & editing; V. A.: funding acquisition, resources, conceptualization, supervision, writing, review & editing.

## Conflicts of interest

There are no conflicts to declare.

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