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Highly sensitive and precise optical temperature sensors based on new luminescent Tb³⁺/Eu³⁺ tetrakis complexes with imidazolic counterions†

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In the present work, new Tb³⁺/Eu³⁺ tetrakis(benzoyltrifluoroacetone) complexes containing imidazolic counterions were successfully prepared and characterized via elemental analysis, infrared spectroscopy, thermogravimetry, and X-ray powder diffraction. Photophysical features of the Eu³⁺ ion, such as the intrinsic emission quantum yields, radiative and non-radiative decay rates, and emission lifetimes were dramatically improved when compared with the data from the hydrated tris complex reported in the literature. The values found for the absolute emission quantum yields are up to 0.79 and ratiometric luminescent thermometers were built-up based on the ratio between the spectral areas of the $^5D_0 \rightarrow ^7F_2$ (Eu³⁺) and $^5D_4 \rightarrow ^7F_5$ (Tb³⁺) transitions. The best contactless temperature sensor operates in a wide temperature range (20–225 K) with a relative thermal sensitivity higher than 4% K $^{-1}$ (maximum value of $\sim 7.6\%~{\rm K}^{-1}$ at 20 K) and a temperature uncertainty lower than 0.1 K with a minimal lower than 0.01 K by combining excitation at the ligand (360 nm) and the Tb3+ ion (489 nm), illustrating the potential of lanthanide-based tetrakis complexes in the design of efficient luminescent thermometers.

1. Introduction

Temperature is a fundamental thermodynamic parameter for a vast majority of systems and holds great importance in science and technology. In the most common temperature sensors, direct contact with the sample is required and, therefore, measurements at the sub-micron scale (as intracellular measurements) are almost impossible.^{2,3} This limitation increased the research for accurate thermometers with nanometric resolution such as luminescent thermometers, a topic that has been increasingly explored. 4-12 The first luminescent thermometers developed used mainly non-ratiometric thermometric parameters, such as the integrated area of a single transition. This thermometric feature has enormous disadvantages as the thermometric parameter is compromised by (i) variations

in the sensor concentration, (ii) optoelectronic drifts in the excitation source and detecting system, and (iii) material homogeneity effects. 13-15 Such disadvantages were overcome with the use of ratiometric probes, which use the ratio between two transitions or two spectral regions and, consequently, are selfcalibrating.16,17 It is worth mentioning that several other emission parameters were used as the thermometric parameter (e.g., bandwidth and lifetimes). 10 Nowadays, several Eu³⁺/Tb³⁺ co-doped complexes have been reported as ratiometric luminescent thermometers based on the ${}^5D_4 \rightarrow {}^7F_5$ (Tb³⁺) and ${}^5D_0 \rightarrow {}^7F_2$ (Eu³⁺) transitions, undoubtedly inspired by the pivotal works of Sato et al. and Brites et al. 6,8,10,18-23 Lanthanide complexes are usually used due to their spectroscopic features such as narrow absorption and emission bands, relatively long excited-state lifetimes, and high color purity.²⁴⁻³⁰ The 4f-4f transitions have an atomic-like character owing to the chemical environment shielding by the 5s and 5p filled sub-shells and, thus, the ligand field effect on the 4f levels is small. 31,32 These intraconfigurational transitions are forbidden in the first order by the electric dipole mechanism (Laporte's rule). To overwhelm the low oscillator strengths, chromophores are used as sensitizers to absorb and transfer energy non-radiatively to the Ln³⁺ ions. This effect is known as the antenna effect and is an essential feature in the molecular design of lanthanide emitting materials.³³ It is noteworthy that Weissman first elucidated the antenna effect in 1942.34 β-diketonate compounds and their derivatives are

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reported to be one of the best groups of ligands for the luminescence of Eu^{3+} and Tb^{3+} complexes, due to their first excited triplet state energies (T_1) (and, in some examples, their first excited singlet states S_1). Furthermore, the Ln^{3+} β-diketonate complexes have been studied both in the tris and tetrakis forms. Usually, in the tris complexes, the water molecules can be replaced by neutral ancillary ligands, showing a general formula $[Ln(\beta-diketonate)_3(L)_n]$, where L is a neutral ligand, for example, pyridine, bipyridine, triphenylphosphine oxide, *etc.* However, in the tetrakis complexes, $[cation][Ln(\beta-diketonate)_4]$, there is an additional β-diketonate ligand, which can increase the chance of an intramolecular energy transfer process. 36,37 Besides, the advantage of Ln^{3+} tetrakis complexes is the possibility of energy transfer

example, pyridine, bipyridine, triphenylphosphine oxide, *etc.* However, in the tetrakis complexes, [cation][Ln(β -diketonate)₄], there is an additional β -diketonate ligand, which can increase the chance of an intramolecular energy transfer process. ^{36,37} Besides, the advantage of Ln³⁺ tetrakis complexes is the possibility of energy transfer from organic cations, apart from the diketonate ligands. Also, these compounds are usually anhydrous and the absence of water molecules in the first coordination sphere is known to lead to an increase in the emission quantum efficiency of the Ln³⁺ ions. ^{38,39} It was also reported that organic counterions could improve the stability of the complex and solubility in organic solvents. ⁴⁰ Usually, the ligand T₁ energy is determined by using the Gd³⁺ (4f⁷) ion, due to the large energy gap (\sim 32 000 cm⁻¹) between the ground state ($^{8}S_{7/2}$) and the first excited state ($^{6}P_{7/2}$) and the Gd³⁺ contribution to intersystem crossing (ISC) rates by the paramagnetic and heavy-ion effects. The observation of the ligand phosphorescence broadband in the Gd³⁺ complexes at low temperature is assigned to the T₁ \rightarrow S₀ transition. ^{31,41–43}
In this work, the synthesis, characterization and optical-

thermal properties of the [C₄mim][Tb_{1-x}Eu_x(btfa)₄] complexes (x = 0.01, 0.05, 0.10, 1.00) are reported, where btfa stands for benzoyltrifluoroacetone and [C4mim] stands for 1-butyl-3methyilmidazolium. These luminescent materials were synthesized in a one-pot reaction and characterized via elemental analysis (CHN), Fourier Transform infrared spectroscopy (FTIR), thermogravimetry (TG/DTG), and X-ray powder diffraction (XRD). The photoluminescence properties were investigated based on the excitation and emission spectra, absolute emission quantum yield (QLn), luminescence decay curve of the Eu³⁺ ⁵D₀ level, radiative and non-radiative decay rates, and intrinsic 5D_0 quantum yield (Q_{Eu}^{Eu}) . The thermometric parameters were determined based on the ratio between the integrated intensities of the Tb^{3+ 5}D₄ \rightarrow ⁷F₅ and Eu^{3+ 5}D₀ \rightarrow ⁷F₂ transitions and the relative thermal sensitivity (S_r) , as well as the temperature uncertainty (δT), were discussed.^{8,10} These optical thermometers operate by combining excitations by blue and UV light with minimal temperature uncertainty lower than 0.01 K, indicating that these Tb³⁺/Eu³⁺ tetrakis systems can be used as precise optical temperature sensors.

2. Experimental section

The lanthanide chlorides $LnCl_3 \cdot 6H_2O$ (Ln^{3+} : Eu and Tb) were prepared by using an aqueous suspension of the respective oxide Eu_2O_3 and Tb_4O_7 (Cstarm, 99.99%) with concentrated HCl acid until the solution reached a pH ~ 6.0 . It is worth mentioning that for the terbium chloride synthesis, hydrogen

peroxide (Synth – 30%) was added for the reduction of Tb^{IV} to Tb^{3+} . The ligand used was benzoyltrifluoroacetone (4,4,4-trifluoro-1-phenyl-1,3-butanedione – btfa) from Sigma-Aldrich (99%). The counterion source, $[C_4mim]Br$ (1-butyl-3-methylimidazolium bromide) was synthesized from the reaction of 1-methylimidazole (Sigma-Aldrich, 99%) and 1-bromobutane (Sigma-Aldrich, 99%), using the microwave method described in the literature. ⁴⁴

Elemental analysis was performed in a PerkinElmer CHN 2400 analyzer. The FTIR measurements were registered in KBr pellets in a PerkinElmer Frontier FTIR from 500 to 4000 cm $^{-1}$. The thermogravimetric curves were recorded in a TA Q500 thermoanalyzer from 25 to 900 °C. The X-ray powder diffraction (XPD) patterns were obtained in Miniflex Rigaku equipment (Cu $K_{\alpha 1}$) from 5 to 65° (2 θ).

The photoluminescence study was based on the excitation and emission spectra recorded at room temperature (300 K) and in a temperature interval from 15 to 350 K. These optical data were recorded using a Fluorolog-3[®] Horiba Scientific (Model FL3-22) fluorometer equipped with a 450 W Xe lamp as an excitation source, a modular double grating excitation spectrometer, and an iHR 320 single emission monochromator. The temperature was varied using a helium-closed cycle cryostat, a vacuum system $(4 \times 10^{-4} \text{ Pa})$, and an autotuning temperature controller (Lakeshore 330, Lakeshore) with a resistance heater. All the measurements began at least 300 s after the temperature indicated in the temperature controller remained constant, thus ensuring the thermalization of the samples and a constant temperature during the measurement. The time-resolved experiments were carried out using a Xe-Hg flash lamp (6 µs pulse half-width and 20-30 µs tail) as the excitation source. The conversion from wavelength to energy units and the Jacobian transformation of the intensity values were performed.45 The emission lifetimes were determined by modeling the emission decay curves with a single exponential function.

The absolute emission quantum yields were measured using a quantum yield measurement system C13534 from Hamamatsu with a 150 W xenon lamp coupled to a monochromator for wavelength discrimination, an integrating sphere as the sample chamber and a multichannel analyzer. An empty quartz Petri plate sample holder was used as a reference, and the samples were placed inside another quartz Petri plate, in the form of a pellet of about 1 cm in diameter. Three measurements were made for each sample and the average is reported. The method is accurate within 10%.

2.1. Synthesis of the $[C_4 mim][Tb_{1-x}Eu_x(btfa)_4]$ (x = 0.01, 0.05, 0.10, 1.00)

First, in the preparation of the complexes, an aqueous solution of NaOH (20 mmol in 10 mL of water) was added to an isopropanol solution of benzoyltrifluoroacetone (20 mmol in 40 mL of isopropanol) with stirring at 60 °C to deprotonate the Hbtfa ligand (p $K_a = 7$). After that, an isopropanol solution of [C₄mim]Br (6 mmol in 10 mL of isopropanol) was added to the mixture. Lastly, an aqueous solution of europium chloride (4 mmol in 10 mL of water) was added to the solution and a precipitate of the [C₄mim][Eu(btfa)₄] complex was formed. 40,46

After stirring at around 70 °C for 2 h, the mixture was cooled down to room temperature and then the complex was filtered, washed with water and dried under reduced pressure for 5 h. The same procedure was used to synthesize $[C_4 \text{mim}][Tb_{1-x}Eu_x(\text{btfa})_4]$ ($x=0.01,\ 0.05,\ 0.10$), the only difference is the addition of terbium chloride to the same solution of EuCl₃. All prepared Ln^{3+} complexes are soluble in the most common organic solvents (acetone, chloroform, dichloromethane, acetonitrile) however, they are insoluble in ethanol and hexane. The experimental and calculated elemental analysis data, CHN (%), of the complexes are very close:

 $[C_4mim][Eu(btfa)_4]$ (pale yellow solid): anal. calcd for $C_{48}H_{39}EuF_{12}N_2O_8$: C 50.05, H 3.41, N 2.43; found: C 50.21, H 3.44, N 2.34.

[C₄mim][Tb_{0.99}Eu_{0.01}(btfa)₄] (pale yellow solid): anal. calcd for C₄₈H₃₉Tb_{0.99}Eu_{0.01}F₁₂N₂O₈: C 49.71, H 3.37, N 2.42; found: C 49.51, H 3.25, N 2.42.

[C₄mim][Tb_{0.95}Eu_{0.05}(btfa)₄] (pale yellow solid): anal. calcd for $C_{48}H_{39}Tb_{0.95}Eu_{0.05}F_{12}N_2O_8$: C 49.72, H 3.37, N 2.42; found: C 49.74, H 3.32, N 2.41.

 $[C_4 mim][Tb_{0.90}Eu_{0.10}(btfa)_4]$ (pale yellow solid): anal. calcd for $C_{48}H_{39}Tb_{0.90}Eu_{0.10}F_{12}N_2O_8$: C 49.74, H 3.37, N 2.42; found: C 49.96, H 3.39, N 2.42.

3. Results and discussion

The elemental analysis results of the complexes indicate a formula with four benzoyltrifluoroacetone ligands and one 1-butyl-3-methylimidazolium counterion $[C_4 \text{mim}][Tb_{1-x}\text{Eu}_x(\text{btfa})_4]$ (x=0.01, 0.05, 0.10, 1.00) for all the synthesized complexes. The thermogravimetric curves of these complexes were registered in a dynamic synthetic air atmosphere at 10 °C min⁻¹ rate and showed no mass loss in the 80–200 °C temperature interval (Fig. 1), indicating the anhydrous character of the system. The DTG curve profiles (Fig. S1, ESI†) are the same for all complexes, with the second mass loss event being shifted to lower temperatures as the Tb³⁺ concentration increases. Moreover, the residue from the decomposition of the $[C_4 \text{mim}][\text{Eu}(\text{btfa})_4]$

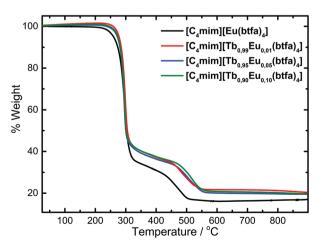


Fig. 1 TGA curves of the $[C_4 \text{mim}][Tb_{1-x}Eu_x(btfa)_4]$ (x = 0.01, 0.05, 0.10, 1.00) complexes recorded under a synthetic air atmosphere at 10 °C min⁻¹.

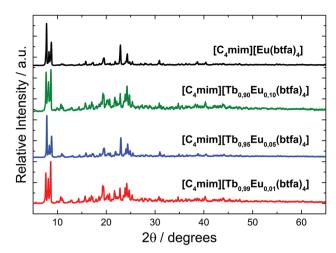


Fig. 2 Powder X-ray diffractograms for the $[C_4 \text{mim}][Tb_{1-x}Eu_x(btfa)_4]$ (x = 0.01, 0.05, 0.10, 1.00) complexes.

is indicated to be europium oxyfluoride and the experimental mass found is very close to the theoretical EuOF residue (calcd: 16.1%, found: 16.2%).

The XRD patterns of the ${\rm Ln}^{3+}$ coordination compounds (Fig. 2) show that the crystalline structure of the materials is very similar, indicating an isomorph character. Since the ionic radius of the ${\rm Tb}^{3+}$ (1.040 Å) and ${\rm Eu}^{3+}$ (1.066 Å) ions are very close for a coordination number 8, the $[{\rm Ln}({\rm btfa})_4]^-$ anionic unities of ${\rm Eu}^{3+}$ and ${\rm Tb}^{3+}$ are isostructural.

FTIR spectra of the $[C_4 mim][Tb_{1-x}Eu_x(btfa)_4]$ (x=0.01, 0.05, 0.10, 1.00) complexes and Hbtfa ligand are reported in Fig. S2 (ESI†). These spectra display a shift and split of the carbonyl stretching absorption from 1603 cm⁻¹ (Hbtfa) to 1574 and 1615 cm⁻¹ for all compounds; there is also a significant change in the 1000–1200 cm⁻¹ range related to C=C and C-O bonds of the enolate anion. 40,46 In addition, it is possible to observe the absorptions of the counterion moiety at around 2960 cm⁻¹ and 2860 cm⁻¹. A weak absorption band observed at ~ 3500 cm⁻¹ is assigned to water absorbed by the KBr pellets once both elemental and thermogravimetric analyses have confirmed the absence of water molecules.

3.1. Photoluminescence

The excitation spectra of the $[C_4 mim][Eu(btfa)_4]$ and $[C_4 mim][Tb_{1-x}Eu_x(btfa)_4]$ (x=0.01, 0.05, 0.10) complexes were recorded in a solid-state from 250 up to 600 nm range monitoring the emission at 610 nm (${}^5D_0 \rightarrow {}^7F_2$ from Eu^{3+}) at 300 K (Fig. 3). These spectra display absorption broadbands between 250 and 450 nm assigned to singlet–singlet transitions ($S_0 \rightarrow S_n$) from the organic moiety (btfa ligands and $[C_4 mim]$ cation). The presence of the ligand absorption band when monitoring at the Eu^{3+} emission indicates an intramolecular ligand-to-metal energy transfer process. 34,35 Furthermore, some narrow absorption bands related to the Eu^{3+} (${}^7F_0 \rightarrow {}^5D_2$) and Tb^{3+} (${}^7F_6 \rightarrow {}^5D_4$) intraconfigurational transitions are also observed. 40,48 Given that a Tb^{3+} absorption transition is present in the excitation spectra while monitoring the Eu^{3+} emission, a metal-to-metal energy transfer may also be present, though the most probable

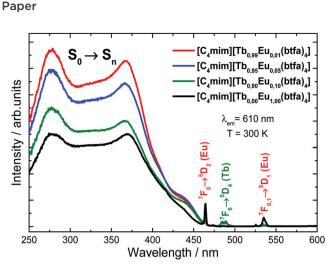


Fig. 3 Excitation spectra for the $[C_4 mim][Tb_{1-x}Eu_x(btfa)_4]$ ($x=0.01,\,0.05,\,0.10$ and 1.00) complexes recorded at room temperature (300 K) monitoring the Eu^{3+} $^5D_0 \rightarrow ^7F_2$ emission (610 nm).

mechanism is ligand-assisted.^{35,49} Notably, the excitation spectrum monitoring at the very weak Tb³⁺ emission has no difference from the Eu³⁺ excitation spectrum, aside from the absence of the Eu³⁺ 4f-4f transitions (not shown).

The emission spectra of the $[C_4\text{mim}][Tb_{1-x}Eu_x(\text{btfa})_4]$ (x = 0.01, 0.05, 0.10 and 1.00) compounds were recorded in the solid-state at 300 K from 450 to 800 nm (Fig. 4) under excitation on the ligand absorption band at 360 nm ($S_0 \rightarrow S_1$). These spectra show the low intensity $Tb^{3+} {}^5D_4 \rightarrow {}^7F_6$ (490 nm)

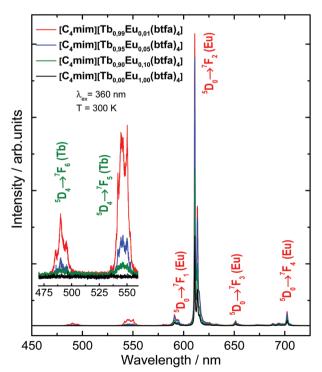


Fig. 4 Emission spectra of the $[C_4 \text{mim}][\mathsf{Tb}_{1-x}\mathsf{Eu}_x(\mathsf{btfa})_4]$ ($x=0.01,\ 0.05,\ 0.10$ and 1.00) complexes recorded at room temperature (300 K), under 360 nm excitation (ligand absorption). The inset shows a magnification of the Tb^{3+} electronic transitions.

Table 1 Spontaneous emission coefficients ($A_{\rm rad}$, $A_{\rm nrad}$, and $A_{\rm tot}$), luminescence lifetimes ($\tau_{\rm obs}$), intrinsic emission quantum yield of the $^5{\rm D}_0$ state ($Q_{\rm EU}^{\rm Eu}$) at 465 nm, and absolute emission quantum yield (at 360 and 465 nm) of the [${\rm C}_4$ mim][Tb $_{1-x}{\rm Eu}_x$ (btfa) $_4$] ($x=0.01,\,0.05,\,0.10,\,1.00$) complexes. The data were collected at room-tempertaure

Complex	$A_{\text{rad}} (s^{-1})$	A_{nrad} (s^{-1})	$A_{\text{tot}} (s^{-1})$	$_{(ms)}^{\tau_{obs}}$	$Q_{\mathrm{Eu}}^{\mathrm{Eu}}$	$\begin{array}{c}Q_{\rm Ln}^{\rm L}\\ \left(360\right)\end{array}$	$\begin{array}{c}Q_{\rm Ln}^{\rm Eu}\\ \left(465\right)\end{array}$
$\begin{split} & \big[\text{Eu}(\text{btfa})_3(\text{H}_2\text{O})_2 \big]^a \\ & \big[\text{C}_4 \text{mim} \big] \big[\text{Eu}(\text{btfa})_4 \big] \\ & \big[\text{C}_4 \text{mim} \big] \big[\text{Tb}_{0.99} \text{Eu}_{0.01}(\text{btfa})_4 \big] \\ & \big[\text{C}_4 \text{mim} \big] \big[\text{Tb}_{0.95} \text{Eu}_{0.05}(\text{btfa})_4 \big] \\ & \big[\text{C}_4 \text{mim} \big] \big[\text{Tb}_{0.90} \text{Eu}_{0.10}(\text{btfa})_4 \big] \end{split}$	775	411 415 507	2632 1250 1190 1282 1190	0.80 0.84 0.78	0.60	0.79	0.12 0.19 0.28
^a Ref. 31, 42, 43.							

and $^5D_4 \rightarrow ^7F_5$ (545 nm) transitions. At room temperature, the Eu $^{3+}$ $^5D_0 \rightarrow ^7F_1$ (590 nm), $^5D_0 \rightarrow ^7F_2$ (610 nm), $^5D_0 \rightarrow ^7F_3$ (650 nm), and $^5D_0 \rightarrow ^7F_4$ (700 nm) transitions are predominant. 48 The low emission intensity of the Tb $^{3+}$ peaks is caused most probably due to back-energy transfer to the btfa ligand, since T $_1$ ($\sim 20\,350~\text{cm}^{-1}$, see Fig. S3, ESI †) is practically resonant with the Tb $^{3+}$ 5D $_4$ level (20 450 cm $^{-1}$). 41 Besides, all the Eu $^{3+}$ transitions have the same emission spectral profiles for all the complexes, indicating the same chemical environment around the Eu $^{3+}$ ion, which is in agreement with the X-ray and FTIR results for the structural similarity.

The radiative and non-radiative decay rate values (Table 1) of the Eu³+ $^5\mathrm{D}_0$ level for the tetrakis complexes were also compared to those reported in the literature for the tris hydrated complex. 31,42,43 The A_{rad} coefficient values were determined by $A_{\mathrm{rad}} = \sum A_{0 \to J}$, where $A_{0 \to J}$ are the spontaneous emission coefficients of $^5\mathrm{D}_0 \to F_J$ transitions (J=1, 2 and 4) of the Eu³+ ion and $S_{0 \to J}$ are the integrated areas under the emission curves (eqn (1)). The spontaneous emission coefficient $A_{0 \to 1}$ of the $^5\mathrm{D}_0 \to F_1$ transition is mostly governed by the magnetic dipole mechanism and depends only on the refractive index of the material. Therefore, this transition is taken as a reference. 31,35

$$A_{0\to J} = A_{0\to 1} \left(\frac{S_{0\to J}}{S_{0\to 1}} \right) \tag{1}$$

The intrinsic emission quantum yield $Q_{\rm Eu}^{\rm Eu}$ of the $^5{\rm D}_0$ emitting level is defined by the ratio between the radiative emission and total decay rates (radiative and non-radiative) for this energy level (eqn (2)). The total decay rate ($A_{\rm rad} + A_{\rm nrad}$) is related to the deactivation process of the $^5{\rm D}_0$ level, mainly by multiphonon relaxation, which is determined by the experimental lifetime values (eqn (3)). The lifetime values reported in Table 1 were calculated from the emission decay curves (Fig. S4, ESI†).

$$Q_{\rm Eu}^{\rm Eu} = \frac{A_{\rm rad}}{A_{\rm rad} + A_{\rm nrad}} = \frac{\tau_{\rm obs}}{\tau_{\rm rad}} \tag{2}$$

$$\tau_{\rm obs} = \frac{1}{A_{\rm rad} + A_{\rm nrad}} = \frac{1}{A_{\rm tot}} \tag{3}$$

The $Q_{\text{Eu}}^{\text{Eu}}$ values of the anhydrous tetrakis complexes (Table 1) are higher compared to that of the hydrated tris complex. This increment in the tetrakis complexes occurs due to the decrease in the non-radiative process rates (A_{nrad}) since the phonon

coupling between the 5D_0 and O–H oscillators is very strong in hydrated compounds. 38,39,50 Comparing the present values of $A_{\rm rad}$ found in the literature for the [Eu(btfa)₃(H₂O)₂], a contrast in the reported values is observed. 31,42,43

This is (at least partially) because emitted power was used in the calculations. 31,35 The absolute quantum yield values measured for the [C₄mim][Tb_{1-x}Eu_x(btfa)₄] compounds with excitation within the ligands (360 nm) and within the ⁵D₂ level (465 nm) are displayed in Table 1. The larger Q_{Ln}^{L} values were found under ligand excitation being 0.35, 0.79 and 0.72, for x = 0.01, 0.05 and 0.10, respectively. The fact that the more concentrated samples (x =0.05 and 0.10) display larger values may be related to the fact that more Eu³⁺ ions are located near the ligands favoring, therefore, the ligand-to-metal energy transfer. We also note that these optical data suggest no significant Eu³⁺/Tb³⁺ luminescence concentration quenching. Besides, both the emissions from the ⁵D₀ (Eu³⁺) and ⁵D₄ (Tb³⁺) energy levels are contributing to the absolute quantum yield values (Fig. S5, ESI \dagger). To compare the measured Q_{Ln}^{L} values with the $Q_{\rm Eu}^{\rm Eu}$ ones, a selective intra-4f⁶ excitation is used (465 nm, Fig. S5, ESI†). Consistently, we observe that the intrinsic emission quantum yield values (Q_{Eu}^{Eu}) are higher than the absolute quantum yield (Q_{Ln}^{L}), as expected. 35,51

4. Thermometric characterization

The thermometric behavior of the $[C_4 mim][Tb_{1-x}Eu_x(btfa)_4]$ complexes (x = 0.01, 0.05, 0.10) was checked recording their emission spectra from 15 to 300 K. Fig. 5 shows the corresponding CIE (Commission International de l'Éclairage) chromaticity diagram of the $[C_4 mim][Tb_{0.99}Eu_{0.01}(btfa)_4]$ compound (Table S1, ESI†). The emission color presents a shift from monochromatic red

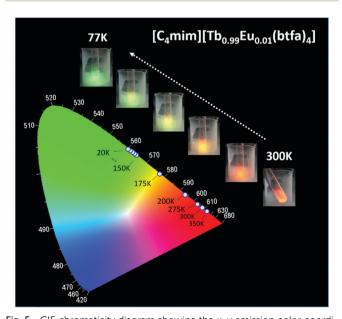


Fig. 5 CIE chromaticity diagram showing the x, y emission color coordinates for $[C_4\text{mim}][Tb_{0.99}Eu_{0.01}(btfa)_4]$ complexes as a function of temperature with excitation at 360 nm. The inset figures are photographs of the complex at 77 to 298 K taken using a digital camera displaying the green to red emission colors under UV irradiation at 360 nm.

(room temperature ~ 300 K) to green (liquid N_2 temperature ~ 77 K). Besides from being excited at the ligand absorption (360 nm), the emission spectra for the $[C_4 \text{mim}][Tb_{1-x}Eu_x(\text{btfa})_4]$ ($x=0.01,\ 0.05,\ 0.10$) complexes (Fig. 6 and Fig. S6, ESI†) were also recorded by exciting selectively at the Tb³⁺ ion (489 nm, $^7F_6 \rightarrow ^5D_4$). Based on these spectra and the integrated emission intensities from 16 259 to 16 499 cm⁻¹ corresponding to the $Eu^{3+}\ ^5D_0 \rightarrow \ ^7F_2$ (I_{Eu}) and from 17 553 to 18 640 cm⁻¹ for the Tb³⁺ $^5D_4 \rightarrow \ ^7F_5$ (I_{Tb}) transitions, the ratiometric thermometric parameter $\Delta = I_{Tb}/I_{Eu}$ was calculated (Fig. 6a–d). 8,10

As in the operating range of the [C₄mim][Tb_{0.90}Eu_{0.10}(btfa)₄] luminescent thermometer under excitation at 360 nm $I_{\rm Tb} > I_{\rm Eu}$ (Fig. S7, ESI†), the temperature dependence of the ratiometric parameter Δ can be approximately described by the classical Mott–Seitz model with one non-radiative channel (eqn (4)):⁵²

$$\Delta \approx \frac{\Delta_0}{1 + B e^{-\frac{E_a}{k_b T}}} \tag{4}$$

where $k_{\rm b}$ is the Boltzmann constant, \varDelta_0 is the limit for the ratio at 0 K, B is a pre-exponential factor and $E_{\rm a}$ is the activation energy for the quenching process. The fitted values are $B=(8\pm5)\times10^5$, $E_{\rm a}=1606\pm80~{\rm cm}^{-1}$, and $\varDelta_0=8.4\pm0.2$ (correlation coefficient, $R^2=0.996$). Remarkably, the $E_{\rm a}$ value is in accordance with a thermal activation ($\sim1050~{\rm cm}^{-1}$) from the $^5{\rm D}_4$ state of the Tb³⁺ ion ($\sim20\,450~{\rm cm}^{-1}$) to the barycenter position ($\sim21\,500~{\rm cm}^{-1}$) of the T₁ state for the btfa ligand via the $^7{\rm F}_5$ excited state. 31,42,43,55,56

A different (ΔT) temperature dependence is observed when the [C₄mim][Tb_{0.90}Eu_{0.10}(btfa)₄] complex is selectively excited at the Tb³⁺ ion (Fig. 6b). The Δ values increase as the temperature increases from 15 to 70 K, decreasing, then, as the temperature increases up to 200 K, and keeping approximately constant for higher temperatures (Fig. 6d). The physical origin of this thermometric behavior is not yet clarified from a theoretical point of view and is not detailed in this work. To allow the characterization of the thermometric performance under excitation at 489 nm, the (ΔT) dependence can be empirically described by:

$$\Delta = \begin{cases}
(12 \pm 2) \left(1 - e^{(46 \pm 2) \times 10^{-3} T} \right), & 20 \le T \le 70 \\
(44 \pm 1) \times 10^2 \times e^{(394 \pm 3) \times 10^{-4} T}, & 70 < T \le 400
\end{cases}$$
(5)

The performance of the luminescent thermometers is made based on the relative thermal sensitivity (S_r), which is an important parameter to compare a thermometer with others, regardless of their origin, and the minimum temperature uncertainty (δT). The relative thermal sensitivity and minimum temperature uncertainty are defined as:

$$S_{\rm r} = \frac{1}{4} \frac{\partial \Delta}{\partial T} \tag{6}$$

and

$$\delta T = \frac{1}{S_{\rm r}} \frac{\delta \Delta}{\Delta} \tag{7}$$

where $\delta \Delta$ is the uncertainty in Δ .

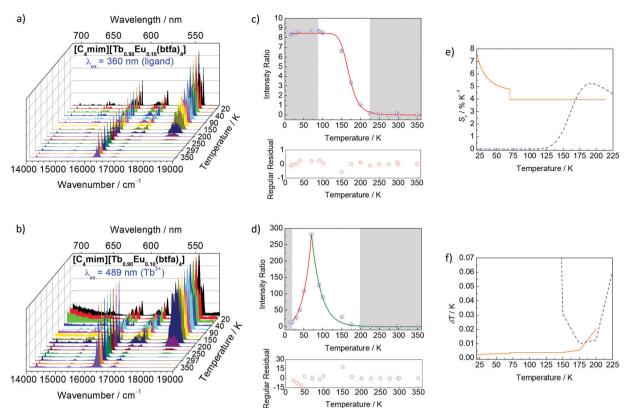


Fig. 6 Emission spectra of the [C₄mim][Tb_{0.90}Eu_{0.10}(btfa)₄] complex recorded as a function of temperature with excitation at 360 nm (a) and 489 nm (b). In (a) the emission spectra are normalized by the intensity at 611 nm (${}^5D_0 \rightarrow {}^7F_2$), while in (b) they are normalized by the maximum intensity in each individual spectrum. The ${}^5D_4 \rightarrow {}^7F_5/{}^5D_0 \rightarrow {}^7F_2$ ratios measured with 360 and 489 nm are shown in (c and d), respectively. The lines represent the fits using egn (5) and (6), respectively. The residual plot is shown on the bottom for a judgment of the fit quality. The shadowed areas mark the regions in which the thermometric parameters change within its uncertainty meaning that for these temperatures the thermometer is out of the so-called operating range. Relative thermal sensitivity (e) and temperature uncertainty (f) obtained at 360 nm (dashed line) and 489 nm (solid line). In (f), the temperature range is limited to $\delta T < 0.07$ K.

The temperature dependence of S_r and δT for the [C₄mim]-[Tb_{0.90}Eu_{0.10}(btfa)₄] luminescent thermometer operating under excitation at 360 and 489 nm are shown in Fig. 6e and f. Under excitation at 360 nm, the thermometer operates in the 150–225 K interval with $S_r > 1\% \text{ K}^{-1}$ (with a maximum value $S_{\rm m} \sim 5\%~{\rm K}^{-1}$ at 200 K) and $\delta T < 1$ K. Similar results were found for $[C_4mim][Tb_{0.99}Eu_{0.01}(btfa)_4]$ and $[C_4mim][Tb_{0.95}-$ Eu_{0.05}(btfa)₄] (Fig. S6a and b, ESI†). Furthermore, under 489 nm [C₄mim][Tb_{0.90}Eu_{0.10}(btfa)₄] can operate in the 20-200 K interval with $S_{\rm r} > 1\%~{\rm K}^{-1}~(S_{\rm m} \sim 7.6\%~{\rm K}^{-1}~{\rm at}~20~{\rm K})$ and $\delta T < 0.03$ K (Fig. 6f). Previous studies on Eu³⁺/Tb³⁺-based luminescent thermometers reported higher $S_{\rm m}$ values although in narrower operating ranges (Table 1 of ref. 10, 44 and 45, as illustrative examples). 10,57,58 Examples are the thermometers with $S_{\rm m}$ = 31% K⁻¹ based on an Ln-HL MOF (Ln = Eu and Tb and HL = 5-hydroxy-1,2,4-benzenetricarboxylic acid) operating between 4 and 50 K,⁵⁷ and $S_{\rm m}$ = 11% K⁻¹ based on a complexpolymer composite of ethyl-diphenylphosphine oxide functionalized poly(dimethylsiloxane) and [Ln(bzac)3] where bzac stands for benzoylacetone (PDMS-eddpo-Ln(bzac)₃; Ln = Eu and Tb) operating between 158 and 248 K.58

An insight into the mechanism for the temperature dependence of the Tb3+ and Eu3+ emissions can be provided by a

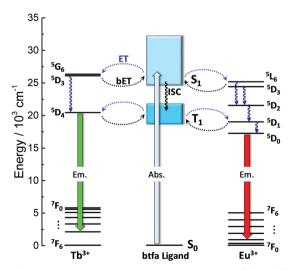


Fig. 7 Schematic partial energy level diagram of the [C₄mim][Tb_{1-x}- $Eu_x(btfa)_4$] (x = 0.01, 0.05 and 0.10) complexes involving the btfa ligand, \mbox{Eu}^{3+} and \mbox{Tb}^{3+} electronic states. \mbox{T}_1 is the lowest ligand triplet level and \mbox{S}_0 and S₁ are the ligand ground and lowest singlet excited states, respectively. Abs. is the initial absorption, ISC is the intersystem crossing and Em is the emission. ET and bET are the forward and backward energy transfer, respectively

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simplified energy level diagram of the system (Fig. 7). The first excited triplet state (T₁) barycenter of the btfa ligand is above the Tb^{3+ 5}D₄ excited state when considering the ⁷F₅ energy level, therefore a back-transfer from the Tb3+ ion to the ligand is expected to occur at room temperature and only Eu³⁺ emission is observed.⁵⁶ On the other hand, at lower temperatures, the back-transfer rate to the ligand should be lower and the Tb³⁺ emission will be predominant.³⁵

5. Conclusions

A series of new temperature sensors based on trivalent lanthanide tetrakis complexes were synthesized successfuly by a one-pot procedure. The formula [C₄mim][Tb_{1-x}Eu_x(btfa)₄] was confirmed via CHN and TG analyses. Besides, the XPD and FTIR data alongside the photoluminescence study indicate structural similarity between the Eu³⁺ and Tb³⁺ doped complexes. The compounds exhibited relatively high absolute and intrinsic emission quantum yields (up to 0.79 and 0.68, respectively), indicating that they can also act as efficient lightconverting molecular devices. The temperature dependence of the intensity ratio between the ${}^5D_4 \rightarrow {}^7F_5$ (Tb³⁺) and ${}^5D_0 \rightarrow {}^7F_2$ (Eu³⁺) transitions for excitation at the ligand (360 nm) in the [C₄mim][Tb_{0.9}Eu_{0.1}(btfa)₄] complex was fitted according to the Mott-Seitz model, suggesting that the depopulation of the ⁵D₄ emitting state of the Tb³⁺ ion could be through thermal activation. The mechanism of this energy transfer process is still unclear, principally when the complex is selectively excited at the ⁵D₄ state (489 nm) with this investigation being a current project in our research group. The [C₄mim][Tb_{0.90}Eu_{0.10}(btfa)₄] luminescent thermometer presented a maximum value for the relative thermal sensitivity around 7.6% at 20 K and a minimum temperature uncertainty below 0.03 K. Moreover, by using excitation at the ligand (360 nm) or the Tb³⁺ ion (489 nm), it is possible to shift the operating range of the thermometer with $S_{\rm m} > 1\%$ in the 20-200 K and 150-225 K intervals, respectively. By combining both excitations, the luminescent thermometers can operate in an extended temperature range with a lowtemperature uncertainty. These spectroscopic results show that the optical thermometers can operate under excitations at blue and UV spectral regions with very low minimal temperature uncertainty (below 0.01 K) with high emission quantum yields (\sim 0.70). Therefore, the [C₄mim][Tb_{1-x}Eu_x(btfa)₄] luminescent materials can be applied as precise photonic temperature sensors.

Conflicts of interest

There are no conflicts to declare.

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Notes and references

- 1 P. R. N. Childs, J. R. Greenwood and C. A. Long, Rev. Sci. Instrum., 2000, 71, 2959-2978.
- 2 L. Shi, C. Dames, J. R. Lukes, P. Reddy, J. Duda, D. G. Cahill, J. Lee, A. Marconnet, K. E. Goodson, J. H. Bahk, A. Shakouri, R. S. Prasher, J. Felts, W. P. King, B. Han and J. C. Bischof, Nanoscale Microscale Thermophys. Eng., 2015, 19, 127-165.
- 3 D. Jaque, E. Martín, L. Martínez and P. Haro, Nanomedicine, 2014, 9, 1047-1062.
- 4 A. S. Souza, L. A. O. Nunes, I. G. N. Silva, F. A. M. Oliveira, L. L. Da Luz, H. F. Brito, M. C. F. C. Felinto, R. A. S. Ferreira, S. A. Júnior, L. D. Carlos and O. L. Malta, Nanoscale, 2016, 8, 5327-5333.
- 5 Y. Cui, H. Xu, Y. Yue, Z. Guo, J. Yu, Z. Chen, J. Gao, Y. Yang, G. Qian and B. Chen, J. Am. Chem. Soc., 2012, 134, 3979-3982.
- 6 R. Piñol, C. D. S. Brites, R. Bustamante, A. Martínez, N. J. O. Silva, J. L. Murillo, R. Cases, J. Carrey, C. Estepa, C. Sosa, F. Palacio, L. D. Carlos and A. Millán, ACS Nano, 2015, 9, 3134-3142.
- 7 Z. Wang, D. Ananias, A. Carné-Sánchez, C. D. S. Brites, I. Imaz, D. Maspoch, J. Rocha and L. D. Carlos, Adv. Funct. Mater., 2015, 25, 2824-2830.
- 8 C. D. S. Brites, P. P. Lima, N. J. O. Silva, A. Millán, V. S. Amaral, F. Palacio and L. D. Carlos, Nanoscale, 2012, 4, 4799-4829.
- 9 D. Jaque and F. Vetrone, Nanoscale, 2012, 4, 4301-4326.
- 10 C. D. S. Brites, S. Balabhadra and L. D. Carlos, Adv. Opt. Mater., 2018, 7, 1801239.
- 11 X. D. Wang, O. S. Wolfbeis and R. J. Meier, Chem. Soc. Rev., 2013, 42, 7834-7869.
- 12 S. W. Allison, Meas. Sci. Technol., 2019, 30, 072001.
- 13 M. I. J. Stich, L. H. Fischer and O. S. Wolfbeis, Chem. Soc. Rev., 2010, 39, 3102-3114.
- 14 D. Yang, Z. Li, L. He, Y. Deng and Y. Wang, RSC Adv., 2017, 7, 14314-14320.
- 15 Z. Li, Z. Hou, D. Ha and H. Li, Chem. Asian J., 2015, 10, 2720-2724.
- 16 T. Wang, P. Li and H. Li, ACS Appl. Mater. Interfaces, 2014, 6, 12915-12921.
- 17 X. Y. Xu and B. Yan, J. Mater. Chem. C, 2016, 4, 1543-1549.
- 18 D. Ananias, F. A. A. Paz, D. S. Yufit, L. D. Carlos and J. Rocha, J. Am. Chem. Soc., 2015, 137, 3051-3058.
- 19 C. D. S. Brites, P. P. Lima, N. J. O. Silva, A. Millán, V. S. Amaral, F. Palacio and L. D. Carlos, Adv. Mater., 2010, 22, 4499-4504.
- 20 S. Sato, R. Yamaguchi and T. Nose, J. Inst. Electron., Inf. Commun. Eng., 1989, J72-C2, 906.
- 21 D. Mara, F. Artizzu, B. Laforce, L. Vincze, K. Van Hecke, R. Van Deun and A. M. Kaczmarek, J. Lumin., 2019, 213, 343-355.

Paper

22 R. Boddula, K. Singh, S. Giri and S. Vaidyanathan, *Inorg. Chem.*, 2017, 56, 10127–10130.

- 23 P. K. Shahi, A. K. Singh, S. K. Singh, S. B. Rai and B. Ullrich, ACS Appl. Mater. Interfaces, 2015, 7, 18231–18239.
- 24 S. V. Eliseeva and J. C. G. Bünzli, Chem. Soc. Rev., 2010, 39, 189–227.
- 25 A. De Bettencourt-Dias, P. S. Barber and S. Viswanathan, *Coord. Chem. Rev.*, 2014, 273–274, 165–200.
- 26 M. H. V. Werts, R. T. F. Jukes and J. W. Verhoeven, *Phys. Chem. Chem. Phys.*, 2002, 4, 1542–1548.
- 27 W. M. Faustino, L. A. Nunes, I. A. A. Terra, M. C. F. C. Felinto, H. F. Brito and O. L. Malta, *J. Lumin.*, 2013, 137, 269–273.
- 28 D. Parker, Coord. Chem. Rev., 2000, 205, 109-130.
- 29 N. Shrivastava, L. U. Khan, Z. U. Khan, J. M. Vargas, O. Moscoso-Londoño, C. Ospina, H. F. Brito, Y. Javed, M. C. F. C. Felinto, A. S. Menezes, M. Knobel and S. K. Sharma, J. Mater. Chem. C, 2017, 5, 2282–2290.
- 30 L. C. V. Rodrigues, J. Hölsä, M. Lastusaari, M. C. F. C. Felinto and H. F. Brito, *J. Mater. Chem. C*, 2014, 2, 1612–1618.
- 31 G. F. De Sá, O. L. Malta, C. De Mello Donegá, A. M. Simas, R. L. Longo, P. A. Santa-Cruz and E. F. da Silva, *Coord. Chem. Rev.*, 2000, 196, 165.
- 32 O. L. Malta and L. D. Carlos, Quim. Nova, 2003, 26, 889-895.
- 33 N. Sabbatini, M. Guardigli and J. M. Lehn, *Coord. Chem. Rev.*, 1993, 123, 201–228.
- 34 S. I. Weissman, J. Chem. Phys., 1942, 10, 214-217.
- 35 A. N. Carneiro Neto, E. E. S. Teotonio, G. F. de Sá, H. F. Brito, J. Legendziewicz, L. D. Carlos, M. C. F. C. Felinto, P. Gawryszewska, R. T. Moura, R. L. Longo, W. M. Faustino and O. L. Malta, in *Handbook on the Physics* and Chemistry of Rare Earths, ed. J. C. G. Bünzli and V. K. Pecharsky, Elsevier, Amsterdan, Netherlands, 2019, ch. 310, vol. 56.
- 36 S. Biju, D. B. Ambili Raj, M. L. P. Reddy and B. M. Kariuki, *Inorg. Chem.*, 2006, **45**, 10651–10660.
- 37 A. Mech, M. Karbowiak, C. Görller-Walrand and R. Van Deun, *J. Alloys Compd.*, 2008, 451, 215–219.
- 38 T. R. Faulkner and F. S. Richardson, *Mol. Phys.*, 1978, 35, 1141–1161.
- 39 S. Salama and F. S. Richardson, *J. Phys. Chem.*, 1980, **84**, 512–517.

- 40 S. M. Bruno, R. A. S. Ferreira, F. A. A. Paz, L. D. Carlos, M. Pillinger, P. Ribeiro-Claro and I. S. Gonçalves, *Inorg. Chem.*, 2009, 48, 4882–4895.
- 41 M. Ren, C. D. S. Brites, S. S. Bao, R. A. S. Ferreira, L. M. Zheng and L. D. Carlos, *J. Mater. Chem. C*, 2015, 3, 8480–8484.
- 42 C. De Mello Donegá, S. Alves, O. L. Malta and G. F. De Sá, *Mater. Sci. Forum*, 1999, 315–317, 204–210.
- 43 C. De Mello Donegá, S. Alves and G. F. De Sá, *J. Alloys Compd.*, 1997, **250**, 422–426.
- 44 R. S. Varma and V. V. Namboodiri, *Chem. Commun.*, 2001, 643-644.
- 45 J. Mooney and P. Kambhampati, J. Phys. Chem. Lett., 2013, 4, 3316–3318.
- 46 S. Biju, R. O. Freire, Y. K. Eom, R. Scopelliti, J. C. G. Bünzli and H. K. Kim, *Inorg. Chem.*, 2014, 53, 8407–8417.
- 47 T. Rajkumar and G. Ranga Rao, Mater. Chem. Phys., 2008, 112, 853–857.
- 48 W. T. Carnall, G. L. Goodman, K. Rajnak and R. S. Rana, J. Chem. Phys., 1989, 90, 3443-3457.
- 49 A. N. Carneiro Neto, R. T. Moura, A. Shyichuk, V. Paterlini, F. Piccinelli, M. Bettinelli and O. L. Malta, *J. Phys. Chem. C*, 2020, 124, 10105–10116.
- 50 F. S. Richardson, Chem. Rev., 1982, 82, 541-552.
- 51 L. D. Carlos, R. A. S. Ferreira, V. de, Z. Bermudez and S. J. L. Ribeiro, *Adv. Mater.*, 2009, 21, 509–534.
- 52 C. D. S. Brites, A. Millán and L. D. Carlos, in *Handbook on the Physics and Chemistry of Rare Earths*, ed. J. C. G. Bünzli and V. K. Pecharsky, Elsevier, Amsterdam, Netherlands, 2019, ch. 310, vol. 49.
- 53 N. F. Mott, Proc. R. Soc. London, Ser. A, 1938, 167, 384-391.
- 54 F. Seitz, Trans. Faraday Soc., 1939, 35, 74-85.
- 55 S. Sato and M. Wada, Bull. Chem. Soc. Jpn., 1970, 43, 1955–1962.
- 56 A. S. Souza, L. A. Nunes, M. C. F. C. Felinto, H. F. Brito and O. L. Malta, *J. Lumin.*, 2015, **167**, 167–171.
- 57 X. Liu, S. Akerboom, M. De Jong, I. Mutikainen, S. Tanase, A. Meijerink and E. Bouwman, *Inorg. Chem.*, 2015, **54**, 11323–11329.
- 58 D. A. Gálico, Í. O. Mazali and F. A. Sigoli, New J. Chem., 2018, 42, 18541–18549.