

RESEARCH ARTICLE

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11, 3744**RbPbPS₄: a promising IR nonlinear optical material achieved by lone-pair-cation-substitution-induced structure transformation†**A-Yang Wang,^{‡a,b,c} Sheng-Hua Zhou,^{‡a,d} Mao-Yin Ran,^{a,b,e} Bingxuan Li,^{a,b}
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When stereochemically-active-lone-pair (SCALP) cations are introduced into chalcogenide systems, it is beneficial to produce a remarkable second-harmonic generation (SHG) response (d_{eff}), but it also causes a narrow band gap (E_g), which ultimately results in negative two-photon and free-carrier absorption. Hence, how to obtain a wide E_g (>2.33 eV) while maintaining a strong d_{eff} remains a significant challenge in this domain. In this work, Rb is partially replaced by SCALP Pb in the known ternary centrosymmetric (CS) Rb_3PS_4 (space group: $Pnma$). This results in a quaternary non-centrosymmetric (NCS) RbPbPS_4 (space group: $P2_12_12_1$), which possesses distinct 2D $[\text{PbPS}_4]^-$ layers and Rb^+ occupies the interlayer spaces as the counter cations. Notably, RbPbPS_4 exhibits a promising overall performance, including strong d_{eff} ($2.5 \times \text{AgGaS}_2$), wide IR transmittance cutoff edge (up to $18.1 \mu\text{m}$) along with the large E_g (2.75 eV), resulting in an improved laser-induced damage threshold ($7.5 \times \text{AgGaS}_2$). Theoretical calculations further indicated that the favorable balance between strong d_{eff} and large E_g in RbPbPS_4 can be attributed to the synergies of the $[\text{PbS}_7]$ and $[\text{PS}_4]$ nonlinear optical (NLO)-active units. This study not only presents a high-performance Pb-based IR-NLO candidate but also underscores the effectiveness of partial substitution of SCALP cations in inducing a CS-to-NCS structural transformation.

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Introduction

Nonlinear optical (NLO) materials have garnered unprecedented attention in laser science and technology because of their ability to facilitate frequency conversion in solid-state laser devices.¹ As widely recognized, an excellent NLO candidate should meet several crucial prerequisites: a sufficient second-harmonic generation (SHG) intensity (d_{eff}), a large energy gap (E_g), a wide optical transparent window, a moderate birefringence (Δn), chemical stability, and availability for

obtaining large single crystals.² In the infrared (IR) region, numerous vital fields, including optoelectronic instruments, resource exploration, and remote laser communication, has sparked widespread attention and interest. Despite the strong d_{eff} exhibited by commercial IR-NLO crystals such as AgGaQ_2 ($Q = \text{S}, \text{Se}$)³ and ZnGeP_2 ,⁴ which make them suitable for applications in the IR region, they still suffer from limitations in high-power laser systems due to their low laser-induced damage thresholds (LIDT) or detrimental two-photon absorption, primarily attributed to their small E_g . However, integrating these optical performances into a single crystal is extremely challenging because they typically depend on competing structural requirements, such as the trade-off between wide E_g and strong d_{eff} .⁵ Therefore, it is of scientific and technological significance to explore new IR-NLO crystals with outstanding comprehensive performance to overcome these challenges.

In addition to the performance prerequisites mentioned above, a prerequisite of an IR-NLO crystal is that it has a non-centrosymmetric (NCS) structure.⁶ The addition of stereochemically-active-lone-pair (SCALP) cations, like As^{3+} , Sb^{3+} , Bi^{3+} , Sn^{2+} , and Pb^{2+} , to chalcogenide systems has attracted the most attention among all the feasible approaches for creating IR-NLO materials because the majority of these chalcogenides exhibit exceptional IR-NLO performances.⁷ Among these

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cations, Pb^{2+} makes a significant contribution to the d_{eff} but adversely affects the optical E_g .⁸ As is known, narrow energy gaps ($E_g < 2.33$ eV) are unable to mitigate some harmful two-photon or free-carrier absorption under fundamental 1064 nm laser sources.⁹ Therefore, achieving a wide E_g (>2.33 eV) while maintaining a strong d_{eff} remains a significant challenge in Pb-based IR-NLO material design. On the other hand, partial chemical substitution in view of known centrosymmetric (CS) parent structures has shown to be a straightforward but incredibly successful approach in recent years for the design and synthesis of novel high-performance SCALP-based IR-NLO crystals.¹⁰ Successful examples include NCS $\text{Ba}_2\text{As}_2\text{Se}_5$ (parent structure: CS $\text{Ba}_2\text{GaAsSe}_5$),¹¹ NCS $\text{K}_2\text{Ag}_3\text{Sb}_3\text{S}_7$ (parent structure: CS $\text{K}_2\text{Sb}_4\text{S}_7$),¹² NCS ABiP_2S_6 ($A = \text{K}, \text{Rb}$) (parent structure: CS $\text{A}_4\text{P}_2\text{S}_6$),^{13,14} NCS $\text{Sn}_7\text{Br}_{10}\text{S}_2$ (parent structure: CS SnBr_2),¹⁵ $\text{Sr}_{2-x}\text{Pb}_x\text{GeSe}_4$ (parent structure: NCS Sr_2GeSe_4),¹⁶ and NCS Pb_4SeBr_6 (parent structure: CS PbBr_2).¹⁷ Considering these

factors, our focus lies in identifying suitable CS parent structures with wide E_g and achieving CS-to-NCS structural transformation by introducing SCALP Pb for partial chemical substitution. This approach aims to obtain NCS materials with excellent, well-balanced IR-NLO performance.

The ternary chalcophosphate Rb_3PS_4 has piqued our interest due to its CS space group of $Pmna$ and unique zero-dimensional (0D) cluster structure.¹⁸ By incorporating distorted Pb-based motifs, it is possible to disrupt symmetric centers, leading to structural reconstruction and the formation of NCS compounds. Furthermore, Rb_3PS_4 boasts a sufficiently large E_g (>3.5 eV), which can help offset the potential decrease in E_g resulting from the introduction of SCALP Pb^{2+} cations into new compounds. Following this approach, we have designed and synthesized the quaternary Pb-based thiophosphate RbPbPS_4 , featuring a two-dimensional (2D) layered structure composed of $[\text{PbS}_7]$ polyhedra and $[\text{PS}_4]$ tetrahedra.

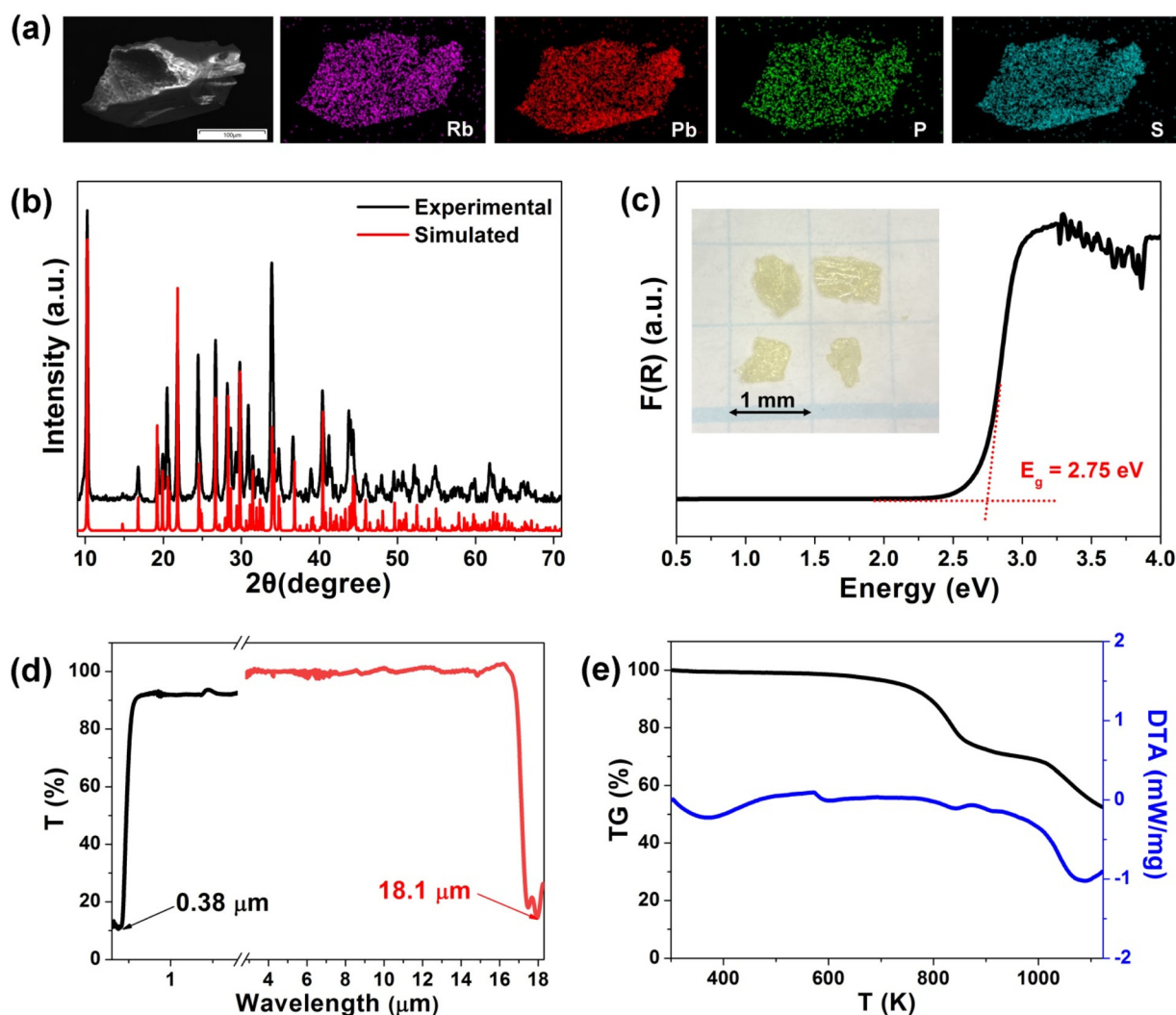


Fig. 1 Experimental characterization results of PbPbPS_4 : (a) SEM image and corresponding elemental distribution maps; (b) experimental (black) and simulated (red) powder XRD patterns; (c) UV-vis-NIR spectrum (inset: photograph of the title crystals); (d) optical transmittance spectrum; (e) TG-DTA test curves.



Remarkably, RbPbPS_4 demonstrates comprehensive IR-NLO performance, including a strong d_{eff} (3.2 times that of AgGaS_2), one of the widest E_g among Pb-based NCS chalcogenides, high LIDT (7.5 times that of AgGaS_2), a broad transmittance range (up to 18.1 μm), and suitable Δn (0.112@2050 nm). In this report, we provide detailed insights into the CS-to-NCS structural transformation and the related optical properties of RbPbPS_4 . Additionally, theoretical calculations are employed to further elucidate its linear optical and NLO performances.

Results and discussion

There are reports of RbPbPS_4 's crystal structure exhibiting two distinct phases.¹⁹ However, we meticulously determined the structure of the title compound to explore the intricate relationship between its NCS crystal structure and NLO properties. Light yellow lamellar crystals of RbPbPS_4 , with sizes reaching the millimeter level, were grown using a simple boron–chalcogen method.²⁰ This method differs from previously reported techniques, which often necessitate complex

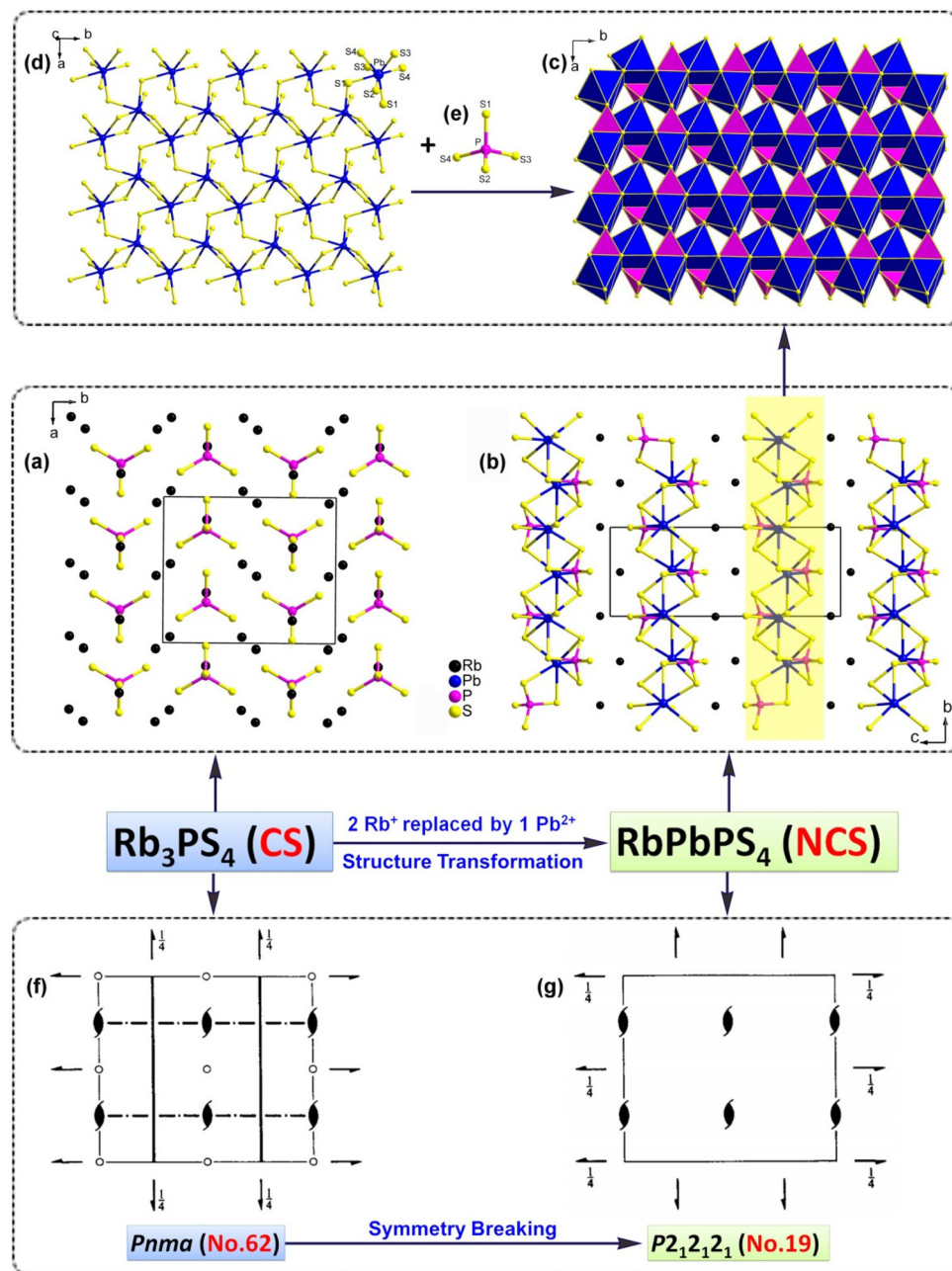


Fig. 2 Structure evolution from ternary CS Rb_3PS_4 to quaternary NCS RbPbPS_4 : view of the 2D layered structures of (a) Rb_3PS_4 and (b) RbPbPS_4 along the ab and bc plane, respectively; (c) a polyhedral 2D layer $[\text{PbPS}_4]^-$ composed of (d) 2D Pb–S layer and (e) discrete tetrahedral $[\text{PS}_4]$ motifs; (f and g) spatial symmetry operation changing from CS $Pnma$ (no. 62) to NCS $P2_12_12_1$ (no. 19).



noteworthy modifications have been made following the partial replacement of Rb by SCALP Pb. First off, the initial highly symmetric structure is broken by the addition of SCALP [PbS₇] BBUs and the reassembly of discrete [PS₄] BBUs, which

(a) Crystal structure of RbPbPS_4 (NCS) showing the arrangement of Rb, Pb, P, and S atoms. The structure is based on the Rb_3PS_4 (CS) structure ($[\text{Pnma}$ (No.62)]). The structure is shown with the following atoms and their positions:

- Rb:4c (1)
- Rb:8d (1)
- P:4c (1)
- S:4c (2)
- S:8d (1)

The structure is shown with the following atoms and their positions:

- Rb:4a (1)
- Pb: 4a (1)
- P: 4a (1)
- S: 4a (4)

The structure is shown with the following atoms and their positions:

- Rb:4c (1)
- Rb:8d (1)
- P:4c (1)
- S:4c (2)
- S:8d (1)

The structure is shown with the following atoms and their positions:

- Rb:4a (1)
- Pb: 4a (1)
- P: 4a (1)
- S: 4a (4)

(b) Plot of $-\text{ICOHP}$ (eV/bond) versus Pb-S distance (Å). The plot shows three regions: Strong (green), Medium (yellow), and Weak (pink). The data points are as follows:

Pb-S distance (Å)	$-\text{ICOHP}$ (eV/bond)
2.92	1.85
2.95	1.75
3.00	1.50
3.05	1.20
3.10	1.20
3.15	1.20
3.25	0.70
3.65	0.30
3.80	0.10

The plot also shows the Pb-S bond lengths (Å) for the different regions:

- Strong: 2.92, 2.95, 3.00, 3.05, 3.10, 3.15
- Medium: 3.25
- Weak: 3.65, 3.80

Quaternary RbPbPS₄ adopts the NCS orthorhombic space group $P2_12_12_1$ (no. 19, Table S1†), according to single-crystal XRD investigation. The cell lattice possesses dimensions of $a = 6.3981(2)$ Å, $b = 6.6888(2)$ Å, $c = 17.2823(5)$ Å, and $V = 739.61(4)$ Å³. These crystal parameters are in good agreement with the NCS phase that was previously published. The compound has a 2D layered structure with discrete counterbalanced Rb⁺ cations inserted in the channels that are arranged in an “ABABAB” mode perpendicular to the ab -plane (Fig. 2b and c). With seven S atoms polyhedral coordinating each Pb atom in the structure, the highly distorted [PbS₇] BBUs with Pb–S bond lengths of 2.916–3.281 Å are formed. As shown in Fig. 2d, adjacent [PbS₇] BBUs are joined by apex-sharing S atoms to generate a 2D Pb–S layer. While the P–S bonds in RbPbPS₄ range in length from 2.031 to 2.069 Å and in bond angles from 106.56 to 111.62°, each P atom is comparable to that of Rb₃PS₄ and is encircled by four S atoms, forming the deformed [PS₄] tetrahedron (Table S2†). Subsequently, apex-sharing between these 2D Pb–S layers (Fig. 2d) and distinct [PS₄] BBUs (Fig. 2e) progressively linked them, resulting in the formation of a 2D [PbPS₄][−] anionic layer (Fig. 2c).

(a) Plot of SHG intensity (a.u.) versus particle size (μm) for PbPbPS_4 (blue diamonds) and AgGaS_2 (black circles) at 2900 nm. The intensity increases with particle size for both materials, with PbPbPS_4 showing higher intensity at larger sizes.

(b) Bar chart showing SHG intensity (a.u.) for AgGaS_2 (1.0) and RbPbPS_4 (2.5) and LIDT (7.5). The LIDT values are significantly higher than the SHG intensities.

(c) Phase diagram of d_{eff} (AgGaS_2) versus E_g (eV) for Pb-based IR-NLO chalcogenides. The diagram shows various compounds plotted as black squares, with RbPbPS_4 highlighted as a red star. The compounds include $\text{PbGa}_2\text{GeSe}_6$, $\text{Pb}_{0.72}\text{Mn}_{2.85}\text{Ga}_{2.95}\text{Se}_8$, $\text{PbGa}_2\text{GeSe}_{12}$, $[\text{Na}_2\text{Pb}][\text{Ga}_2\text{S}_{12}]$, $[\text{K}_2\text{PbX}][\text{Ga}_2\text{S}_{12}]$, $\text{Pb}_{0.65}\text{Mn}_{2.85}\text{Ga}_3\text{S}_9$, AgGaS_2 , PbSeBr_4 , PbGaS_7 , PbZnGaS_{15} , $\beta\text{-PbGa}_2\text{S}_4$, $\text{Pb}_{3.5}\text{GeS}_3\text{Br}_3$, and $\text{PbGa}_2\text{GeSe}_{12}$.

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convert the 0D cluster structure into a 2D layered structure. Fig. 2f and g show the detailed symmetric operation modifications that occur after the substitution of SCALP Pb^{2+} cations, resulting in the removal of all center sites and the realization of the CS-to-NCS structural evolution. In particular, space group $Pnma$ (no. 62) converts to $P2_12_12_1$ (no. 19) from the Bärnighausen tree via a *translationengleiche* reduction in index 2's symmetry (Fig. 3a).²² When Pb is substituted for Rb, the atomic positions (4c and 8d) in Rb_3PS_4 become 4a in RbPbPS_4 (see Table S3† for details). To put it briefly, the novel compound that was produced, RbPbPS_4 , has good structural anisotropy, which helps to get a higher Δn for achieving PM feature. Secondly, the coordination number (CN) of the two crystallographically independent Rb atoms in Rb_3PS_4 differ (CN = 8 for Rb1 and CN = 9 for Rb2), but are identical in RbPbPS_4 (CN = 11) with Rb–S bond distances less than 4.4 Å (refer to Fig. S2 and S3 in ESI†). Finally, considering the flexible CN of Pb^{2+} cations, the integrated crystal orbital Hamilton population (ICOHP) curve was examined.²³ As depicted in Fig. 3b, it is evident that the short Pb–S bond lengths of 2.916–3.114 Å exhibit stronger bonding interactions compared

to the long Pb–S bond lengths (3.647–3.825 Å), while the Pb–S bond distance of 3.281 Å falls between them, representing a medium interaction (see Table S2† for details). Based on the above analysis, it is reasonable to infer that the Pb atom forms a $[\text{PbS}_7]$ mono-capped triangular prism. Research on the distinct $[\text{PbS}_7]$ coordination mode in NCS chalcogenides is notably lacking, and a seven-coordinated example of this has been reported in $\text{Pb}_5\text{Ga}_6\text{ZnS}_{15}$.^{8d} Pb^{2+} cations in NCS chalcogenides are mostly five-coordinated (as in $[\text{Na}_2\text{PbI}][\text{Ga}_7\text{S}_{12}]^{8f}$) or six-coordinated (as in $\text{Pb}_4\text{Ga}_4\text{GeS}_{12}$ ^{8a} and $\text{Pb}_{0.65}\text{Mn}_{2.85}\text{Ga}_3\text{S}_8$ ^{8b}), among other compounds that have been found.

Furthermore, through comparison with known ternary and quaternary chalcophosphates, it is determined that substituting SCALP Pb for partial alkali metal Rb is the most optimal and rational structural approach to accomplish the CS-to-NCS structural transformation of template structure Rb_3PS_4 . Fig. S4† provides a detailed structural evolution. Roughly speaking, these replacements fall into three categories. In the first, three Rb atoms are completely substituted by one trivalent element or three monovalent elements in the absence of

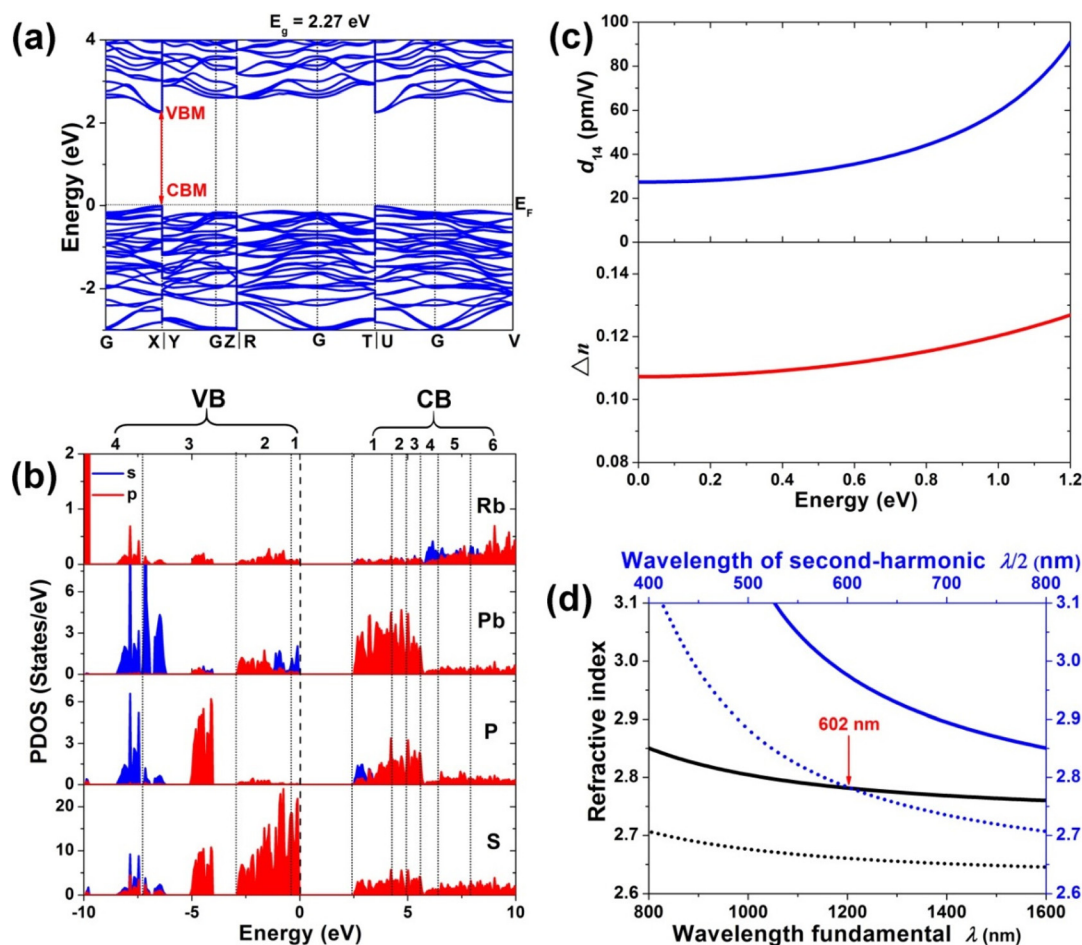


Fig. 5 Theoretical results of electronic structures and optical parameters of PbPbPS_4 : (a) band structure; (b) partial density of states; (c) calculated nonzero independent SHG coefficient d_{14} (pm V^{-1}) and calculated birefringence (Δn); (d) calculated refractive index dispersion curves with the shortest type-I PM cut-off wavelength.



SCALP components. CS structures are displayed by the resultant compounds, including REPS₄ (space group: $I4_1/acd$),²⁴ BPS₄ (space group: $Ibam$),²⁵ GaPS₄ (space group: $P2_1/c$),²⁶ and Tl₃PS₄ (space group: $Pnma$).²⁷ The second category entails complete substitution of SCALP elements, as seen in SbPS₄ (space group: $P\bar{1}$)²⁸ and BiPS₄ (space group: $Ibca$),²⁹ which also belong to the CS structures. The third category involves partial substitution, where Rb is partially replaced by monovalent or divalent elements. The resulting compounds, such as Rb₂AgPS₄ (space group: $P\bar{1}$),³⁰ RbAg₅P₂S₈ (space group: $Pbca$),³¹ RbSrPS₄ (space group: $Pnma$),³² RbBaPS₄ (space group: $Pnma$),³³ RbPdPS₄ (space group: $I4/mcm$),³⁴ and RbHgPS₄ (space group: $P2_1/n$),³⁵ all adopt CS structures. Therefore, it is reasonable to replace partial Rb with SCALP elements to realize a structural transition from CS Rb₃PS₄ to NCS RbPbPS₄. Hence, employing established CS chalcophosphates as the parent structure and partially substituting them with SCALP-cation groups proves to be an effective chemical strategy for accomplishing the CS-to-NCS structural evolution.

As RbPbPS₄ belongs to the NCS space group $P2_12_12_1$ and possesses the NLO-active motifs [PbS₇] and [PS₄], a relatively

strong d_{eff} could be expected. Powder SHG testing for RbPbPS₄ was conducted using the modified Kurtz–Perry technology,³⁵ with irradiation from a 2900 nm laser and commercial material AgGaS₂ as a reference. As depicted in Fig. 4a, the positive correlation between the particle size and SHG intensity indicates that RbPbPS₄ exhibits PM behavior. The SHG intensity of RbPbPS₄ is approximately 2.5 times that of benchmark AgGaS₂ in a particle size range of 150–210 μm (Fig. 4b). Based on $d_{\text{eff}} = d_{\text{eff,R}} (I^2\omega/I_R^2\omega)^{1/2}$ ($d_{\text{eff,R}} = 13.4 \text{ pm V}^{-1}$ for AgGaS₂), the effective SHG coefficient d_{eff} of RbPbPS₄ at 2900 nm is 21.2 pm V^{-1} . In addition, powder LIDT measurement was employed to preliminarily evaluate the LIDT value,³⁶ with the increased optical E_g suggesting that the title crystal would exhibit higher LIDT. As illustrated in Fig. 4b, the LIDT of RbPbPS₄ (10.88 MW cm^{-2}) is estimated to be 7.5 times that of AgGaS₂ (1.45 MW cm^{-2}) under the same testing conditions. Compared with other reported Pb-based IR-NLO materials (Fig. 4c and Table S4†), RbPbPS₄ demonstrates a favorable balance between strong d_{eff} and large E_g . Notably, RbPbPS₄ represents the first quaternary Pb-based chalcogenide to overcome the incompatibility between a wide E_g ($>2.56 \text{ eV}$, the E_g of AgGaS₂) and a

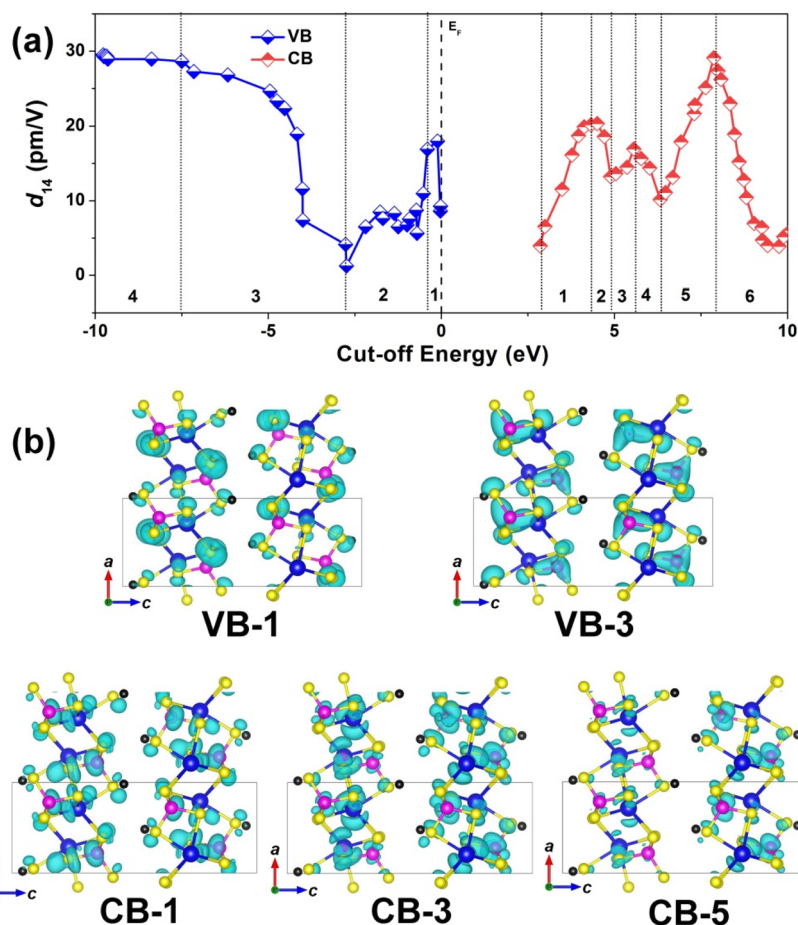


Fig. 6 Theoretical analysis of the intrinsic mechanism of SHG source for PbPbPS₄: (a) cut-off energy (eV) dependence of the static d_{14} (pm V^{-1}); (b) distribution of the partial charge density maps with major contributions in the VB-1, VB-3, CB-1, CB-3 and CB-5 regions. Black atoms: Rb; blue atoms: Pb; pink atoms: P; yellow atoms: S.



sufficient d_{eff} ($>1.0 \times \text{AgGaS}_2$). This balance IR-NLO performance can be compared to the recently reported quaternary NCS chalcogenides, such as $\alpha\text{-Li}_2\text{ZnGeS}_4$,³⁷ $\text{Li}_4\text{CdSn}_2\text{S}_7$,³⁸ $\text{NaMg}_3\text{Ga}_3\text{S}_8$,³⁹ and $(\text{CuI})_3\text{P}_4\text{S}_4$.⁴⁰

To gain further insights into the underlying structure-property relationship of the title compound, systematic theoretical calculations were conducted, encompassing band structure, densities of states (DOSs), and linear optical and NLO parameters, based on DFT. The calculated electronic band structure reveals a direct E_g of 2.27 eV for RbPbPS_4 , as the valence band maximum (VBM) and conduction band minimum (CBM) coincide at the same high-symmetry k -points (Fig. 5a). The smaller calculated E_g can be attributed to underestimation by the DFT calculation.⁴¹ The partial DOSs of RbPbPS_4 are depicted in Fig. 5b, indicating that the tops of the VBs are predominantly composed of S-3p states, with a minor contribution from Pb-5s and P-3p states, while the bottoms of the CBs primarily consist of S-3p, Pb-5p, P-3s/3p, and Rb-5p states. Furthermore, the NLO coefficients of RbPbPS_4 can be calculated according to the electronic structure. In the light of Kleinman symmetry,⁴² RbPbPS_4 exhibits only one nonzero independent SHG coefficient, $d_{14} = 31.8 \text{ pm V}^{-1}$ at 2900 nm (*ca.* 0.43 eV) (Fig. 5c). Furthermore, the calculated d_{eff} of RbPbPS_4 is 26.88 pm V^{-1} , under 2900 nm irradiation,³⁵ which essentially coincides with the experimental findings. Additionally, the calculated Δn value of RbPbPS_4 is 0.109 at 2900 nm, significantly surpassing that of AgGaS_2 ($\Delta n \sim 0.04$).⁴³ This suggests that RbPbPS_4 may be well-suited to achieve the PM condition for the SHG process. Furthermore, the line refractive dispersion diagrams indicate that the shortest type-I phase-matched output wavelength is 602 nm (Fig. 5d).⁴⁴

To visually illustrate the contribution of the constituent BBUs to the SHG effect, the cutoff energy dependence of the static d_{14} for RbPbPS_4 is calculated using a length-gauge formalism method.⁴⁵ As depicted in Fig. 6a, the d_{14} values exhibit an upward trend in the VB-1, VB-3, CB-1, CB-3, and CB-5 intervals. This illustrates that the orbitals in these regions have a significant influence on the overall NLO response.⁴⁶ Moreover, by integrating the partial DOSs (Fig. 5b) with the relevant partial charge density maps (Fig. 6b), it is evident that the enhancement of the d_{eff} in RbPbPS_4 primarily stems from the strong hybridization of electron states at the top of the VBs and the substantial distortion of $[\text{PbS}_7]$ and $[\text{PS}_4]$ groups, indicative of the 2D $[\text{PbPS}_4]^-$ alternating arrangement layer.

Conclusions

In summary, building upon the parent ternary CS compound Rb_3PS_4 , we implemented a strategy involving the partial substitution of Rb with SCALP Pb, resulting in the creation of a new quaternary NCS chalcophosphate, RbPbPS_4 . This substitution also triggered a structural transformation from 0D clusters to 2D layers. As anticipated, experimental findings highlight RbPbPS_4 as a promising candidate for IR-NLO applications, boasting a large d_{eff} ($2.5 \times \text{AgGaS}_2$ at 2900 nm), a broad trans-

mittance cutoff window (0.38–18.1 μm), and a significant E_g (2.75 eV) conducive to a high LIDT ($7.5 \times \text{AgGaS}_2$ at 1064 nm). Detailed structure-property analyses elucidate that the strong calculated d_{eff} (26.88 pm V^{-1} at 2900 nm) and large calculated birefringence (0.109@2900 nm) in RbPbPS_4 primarily stem from the cooperative effects of $[\text{PbS}_7]$ and $[\text{PS}_4]$ NLO-active motifs. These findings underscore the efficacy of the lone-pair-cation substitution strategy in achieving both structural transformation and effectively balancing strong d_{eff} with a large E_g in SCALP-based chalcogenides.

Author contributions

A-Yang Wang: investigation, formal analysis, writing – original draft. Sheng-Hua Zhou: investigation, methodology, validation. Mao-Yin Ran: investigation, formal analysis, validation. Bingxuan Li: formal analysis, validation. Xin-Tao Wu: conceptualization, writing – review & editing. Hua Lin: supervision, conceptualization, writing – review & editing. Qi-Long Zhu: supervision, writing – review & editing.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

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