



# Solvent-induced local environment effect in plasmonic catalysis†

Cite this: *Nanoscale Adv.*, 2023, 5, 5774Tien Le  and Bin Wang \*Received 28th September 2023  
Accepted 30th September 2023

DOI: 10.1039/d3na00835e

rsc.li/nanoscale-advances

Solvents are known to affect the local surface plasmon resonance of metal nanoparticles; however, how solvents can be used to manipulate the interfacial charge and energy transfer in plasmonic catalysis remains to be explored. Here, using NH<sub>3</sub> decomposition on a Ru-doped Cu surface as an example, we report density functional theory (DFT) and delta self-consistent field (SCF) calculations, through which we investigate the effect of different protic solvent molecules on interfacial charge transfer by calculating excitation energy of an electronic transition between the metal and the molecular reactant. We find that the H-bonds between water and NH<sub>3</sub> can alter the direct interfacial charge transfer due to the shift of the molecular frontier orbitals with respect to the metal Fermi level. These effects are also observed when the H-bonds are formed between methanol (or phenol) and ammonia. We show that the solvent possessing stronger basicity induces a more pronounced effect on the excitation energy. This work thus provides valuable insights for tuning the excitation energy and controlling different routes to channel the photon energy into plasmonic catalysis.

relatively high temperatures such as dehydrogenation, ammonia synthesis, CO<sub>2</sub> reduction, and hydrocarbon reforming.<sup>7</sup>

Plasmonic catalysis is stimulated by the intensive charge polarization and the amplified near field due to local surface plasmon resonance (LSPR) of the metal nanoparticles. When LSPR dissipates the energy *via* non-radiative decay, non-equilibrium charge carriers (energetic electrons and holes) can be generated. If there are electron-accepting orbitals available in the proximity, such as unoccupied orbitals of adsorbed molecules, these high energy electrons with adequate energy can be transferred to the adsorbates and induce chemical reactions. Therefore, it is essential to determine the required excitation energy within a complex containing the plasmonic metal and molecular adsorbates. The plasmonic excitation and interfacial energy/charge transfer highly depends on several factors such as size and composition of the metal nanoparticles and the electronic structure of the adsorbed reactants, since these factors define the wavefunction overlap at the interface between reactants and the catalyst surface. Therefore, interfacial engineering provides a venue to modify the interaction between reactants and the catalyst surface and tune the interfacial energy and charge transfer.

This tunability can be enhanced by the presence of a liquid phase. In liquid-phase catalysis, solvents can alter the interfacial binding and even participate directly in the reactions.<sup>8,9</sup> Such effects of solvents were also studied in plasmonic catalysis but mostly focused on the changes in the LSPR of the metal nanoparticles.<sup>10–15</sup> For example, the presence of solvents can cause light scattering, introduce different refractive indices,<sup>15</sup> and manipulate density of surface charges.<sup>10</sup> The energy of LSPR can also be dissipated into the solvent bath through electron–phonon coupling, but its kinetics is normally much slower than that of electron–electron scattering and charge recombination.<sup>16</sup> In addition, solvents, such as methanol as a hole scavenger, and solute species (*e.g.*, OH<sup>−</sup>, O<sub>2</sub>) dissolved in the solvents may also perturb the charge transfer at the interface.<sup>17–19</sup> Furthermore,

## Introduction

Plasmonic catalysis has received extensive interest over the last decade due to its potential to activate chemical bonds by channeling the photonic energy *via* non-equilibrium charge carriers instead of a traditional heat-induced vibrational population.<sup>1–6</sup> This potential motivates researchers to design a catalyst system to achieve desirable product selectivity, which is a challenge in many thermal-driven reactions, especially for processes that are typically performed at

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† Electronic supplementary information (ESI) available: The change in the work function and core level, the DOS of the phenol case, the comparison of the direct and indirect calculation of the HOMO–LUMO gap, and the results of a pristine Cu surface. See DOI: <https://doi.org/10.1039/d3na00835e>



the electrical double layer at the solid–liquid interface can also reduce the Schottky barrier to promote charge injection.<sup>20</sup> Apart from these mentioned multiple effects of the solvents in plasmonic catalysis, the effect of solvents on the interfacial excitation energy and charge transfer between the plasmonic metal and the molecular reactants remains to be explored. This includes how the hydrogen bonds between solvents and reactants change the excitation energy and how such a change can be controlled to achieve different activities and selectivities by using different solvents.

Here, we applied DFT and the linear expansion  $\Delta$ SCF calculations to investigate the effects of solvent molecules on excitation to induce the interfacial charge and energy transfer. We used  $\text{NH}_3$  decomposition on Ru-doped Cu as a probe reaction since  $\text{NH}_3$  is a promising hydrogen-carrier for  $\text{H}_2$  storage and transportation.<sup>21–23</sup> The Ru-doped Cu was reported to be an effective photocatalyst for  $\text{NH}_3$  decomposition previously: Cu was selected as the plasmonic metal based on its abundance, and well-dispersed Ru is incorporated into the Cu surface to provide an active center for N–H bond activation due to its suitable binding energy with  $\text{NH}_3$ .<sup>22</sup> A few prototype polar solvents, such as water, methanol and phenol, were selected to study the effect of the local environment induced by solvents around the surface-adsorbed  $\text{NH}_3$ . We find that H-bonds formed between these solvents and  $\text{NH}_3$  increase the excitation energy of an electron from the Fermi level of the system to the lowest unoccupied molecular orbital (LUMO) of  $\text{NH}_3$  and reduce the excitation energy of an electron from the highest occupied molecular orbital (HOMO) of  $\text{NH}_3$  to the Fermi level. We also find that changing the basicity of solvents can alter the LUMO of  $\text{NH}_3$  to different levels, resulting in a change in the electronic excitation to the LUMO. That is, the higher the basicity of the solvent, the higher the required energy to populate the LUMO of  $\text{NH}_3$ . The HOMO–LUMO excitation can also be perturbed by the presence of solvents, but this effect is less extensive than the increase in excitation energy of

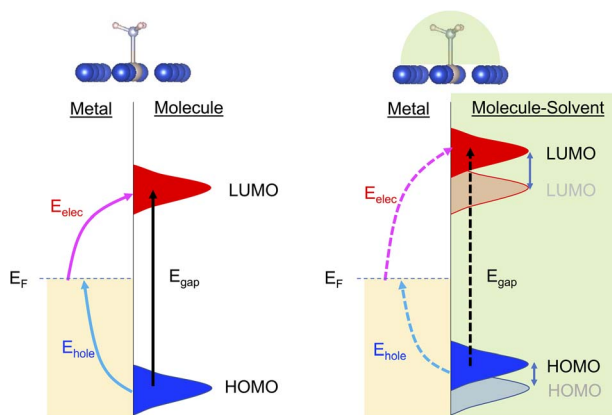
interfacial charge transfer. This observation that the hydrogen bonding between solvents and reactants can modify the plasmonic excitation energy (Scheme 1) provides a valuable route for controlling the selectivity of channeling the photonic energy into a targeted molecular orbital for specific reactions.

## Computational methods

A single Ru atom incorporated into the Cu (111) surface was selected as the model in this study since low loading of Ru was used in previous experiments, and this model was used in the literature to explain the enhanced rates of  $\text{NH}_3$  decomposition over the Ru-doped Cu photocatalyst.<sup>21–23</sup> We adopted this model here to investigate the effect of solvents and illustrate the concept.

The DFT calculations with the PBE-GGA exchange–correlation functional were performed using VASP version 5.4.1.<sup>24,25</sup> The van der Waals (vdW) interaction was included using the DFT-D3 dispersion model.<sup>26</sup> A plane-wave basis set with a kinetic energy cutoff of 500 eV was used for structural optimization, in which the criteria of atomic force was set to be less than  $0.02 \text{ eV } \text{\AA}^{-1}$ . A four-layer  $\sqrt{21} \times \sqrt{21}$  Ru-doped Cu(111) ( $\text{RuCu}_{83}$ ) slab was constructed to model the surface of the Ru-doped Cu nanoparticle. While the two top layers were fully relaxed, the two bottom layers were fixed at the Cu bulk positions. A  $15 \text{ \AA}$  vacuum was added in the  $z$ -direction to separate slabs from their periodic images. The  $k$ -point grid of  $(4 \times 4 \times 1)$  was used to converge the total energy. The energy convergence for electronic relaxation was set to  $10^{-4} \text{ eV}$ . The total energy was extrapolated to 0 K. The  $\text{NH}_3$  molecule was surrounded by one, two, or three molecules of protic solvents. In this study, we focus on solvents that can serve as hydrogen bond acceptors due to their more significant enhancement of  $\text{NH}_3$  adsorption (Table S1†). Both water and methanol have weaker adsorption than  $\text{NH}_3$  over the Ru-doped Cu(111) surface (Table S2†). Phenol may compete for adsorption sites due to its very comparable adsorption energies with that of  $\text{NH}_3$ , and it is included here to show the trend of the excitation energy as a function of the strength of the H-bonds between the solvent molecules and  $\text{NH}_3$ . The effect of solvents on the excitation energy when the solvents are hydrogen bond donors will be investigated in a future study.

Since the GGA functionals often underestimate the HOMO–LUMO gap, the excitation energy from the HOMO to the LUMO, from the Fermi level to the LUMO, and from the HOMO to the Fermi level was calculated *via* the delta SCF method instead of directly using the HOMO and LUMO energy levels from GGA. The fully relaxed structures in the ground state calculated with VASP were adopted to calculate excited states *via* the vertical excitation using GPAW version 22.1.0.<sup>27,28</sup> The exchange–correlation functional, plane-wave cutoff, and  $k$ -point grid in GPAW were kept the same as those applied in VASP. The linear expansion  $\Delta$ SCF<sup>29</sup> approach was used to calculate the energy of the system at excited states; we recently showed that this approach could provide excitation energy values that were in good agreement with experiments.<sup>30</sup>



**Scheme 1** Illustration of the effect of solvent on the excitation energy. The presence of solvent shifts the energy levels of the molecular frontier orbitals (HOMO and LUMO) with respect to the metal Fermi level ( $E_F$ ). An  $\text{NH}_3$  molecule is used in the atomic structures to illustrate the concept.



The center of the molecular orbital is the mean energy calculated as follows:

$$\bar{E} = \frac{\int \rho(E) E dE}{\int \rho(E) dE}$$

where  $\rho(E)$  represents density of states (DOS) at the corresponding energy  $E$ .<sup>30,31</sup>

It is worth noting that the nuclear quantum effect can influence the hydrogen-bonded system, either increasing the proton sharing or promoting the hydrogen-bond breaking.<sup>32</sup> Here we focused on the NH<sub>3</sub> adsorption and H-bond-induced change in the excitation energy; it is expected that the nuclear quantum effect is less pronounced and that its effect on the excitation energy probably cancels out when adding more solvent molecules. However, for N–H activation and proton shuttling, a further study including the nuclear quantum effect, *e.g.*, using PIMD, is valuable to bring more insights. In addition, since GGA functionals often overestimate the strength of the hydrogen bond,<sup>33,34</sup> instead of using the hydrogen bond length or energy, we used the hydrogen bond acceptor parameter (HBA or  $\beta$ ) in the literature to compare the effects of solvents.

## Results and discussion

Within the mechanism of direct interfacial charge transfer in plasmonic catalysis, the charge excitation can take place from metal nanoparticles to unoccupied frontier orbitals of adsorbates, generating energetic electrons in the molecule. If the reactions are driven by holes, the excitation energy of an electron from the occupied states of the adsorbate to the Fermi level is more relevant. Instead, if the reaction follows the resonance energy transfer (RET) instead of the charge transfer mechanism, the excitation within the molecular adsorbates is important. These scenarios are schematically shown in Scheme 1, in which we illustrate the concept that the shift of the molecular frontier orbitals can lead to change in the excitation energy. In the following section we use NH<sub>3</sub> activation over a Ru-doped Cu surface to reveal the effect of solvents.

As shown in the literature,<sup>21–23</sup> at low surface coverage, the first N–H bond dissociation in NH<sub>3</sub> is the rate-limiting step for light-driven NH<sub>3</sub> decomposition on a Ru-doped Cu surface. In addition, the interfacial charge transfer between the catalyst and adsorbed reactant, *i.e.*, the injection of the hot carriers into the antibonding orbitals of the metal-adsorbate system, was found to be essential for plasmonic ammonia decomposition on a Ru-doped Cu surface.<sup>23</sup> Therefore, we first study the interfacial charge transfer between the metal and the adsorbed reactant as the base case, in which NH<sub>3</sub> adsorbs on a clean Ru-doped Cu surface without solvents. The optimized structure is shown in the top channel in Fig. 1. In this study, the interfacial charge transfer is analyzed *via* both the hole transfer process, which is represented by exciting an electron from the HOMO of NH<sub>3</sub> to the Fermi level (denoted as  $E_{\text{hole}}$ ), and the electron transfer process, in which an electron is excited from the Fermi level to the LUMO of NH<sub>3</sub> (denoted as  $E_{\text{elec}}$ ). In the absence of solvents,  $E_{\text{hole}}$  and  $E_{\text{elec}}$  are 6.0 and 3.2 eV, respectively. In a recent experiment, the plasmonic NH<sub>3</sub> decomposition rate



Fig. 1 The effect of H<sub>2</sub>O molecules on the interfacial charge and energy transfer. Top panel: Optimized configurations of NH<sub>3</sub> on RuCu<sub>83</sub> with and without H<sub>2</sub>O. Bottom: the excitation energy required to excite an electron from the HOMO of NH<sub>3</sub> to the Fermi level (blue), from the Fermi level to the LUMO of NH<sub>3</sub> (red), and from the HOMO to the LUMO by adding the two excitation energy values (black).

reaches the maximum at 2.6 eV,<sup>23</sup> which is 0.6 eV lower compared to the calculated  $E_{\text{elec}}$  of the base case. This difference could result from a more complicated conformation of the real catalyst surface of Ru-doped Cu compared to that of the single-atom dopant model used here. The adsorption energy of NH<sub>3</sub> on the clean RuCu<sub>83</sub> surface is about  $-1.4$  eV, which is comparable with literature values.<sup>23</sup>

The effect of the local environment induced by solvents on the interfacial charge transfer is then studied by surrounding NH<sub>3</sub> with H<sub>2</sub>O solvents. In this case we use NH<sub>3</sub>, rather than NH<sub>4</sub><sup>+</sup>, to illustrate the concept. When solvents act as the hydrogen bond acceptors, it is expected that the presence of ammonium ions is not pronounced. A further study on the effect on ammonium ions in water is valuable for capturing the complete picture of this process. Nevertheless, we find that the adsorption energy of NH<sub>3</sub> increases when there is interaction between H<sub>2</sub>O and NH<sub>3</sub> (Table S1 and Fig. S4†); however, a more significant impact is observed when a H-bond is formed between NH<sub>3</sub> and H<sub>2</sub>O at a side-adsorption configuration compared to the on-top adsorption of H<sub>2</sub>O on NH<sub>3</sub> (Fig. 1). The on-top adsorption of H<sub>2</sub>O on NH<sub>3</sub> does not alter the interaction between NH<sub>3</sub> and the metal surface much, leading to a similar value of excitation energy to the one on a clean surface without solvents (Fig. 1). Instead, the side adsorption of H<sub>2</sub>O, through a hydrogen bond, alters the excitation energy for both the electron and the hole transfer process. The more hydrogen bonds formed between H<sub>2</sub>O and NH<sub>3</sub> as the number of surrounding H<sub>2</sub>O molecules increases, the higher the excitation energy needed. We observed that the  $E_{\text{elec}}$  increases by 0.4 eV per H-bond between H<sub>2</sub>O and NH<sub>3</sub>, while the value of  $E_{\text{hole}}$  decreases by 0.2 eV per H-bond (Fig. 1).

This change in the excitation energy can be attributed to the strong electronegativity of the oxygen atom in water, which polarizes the molecule and upshifts the HOMO and LUMO centre as shown in Fig. 2. The more pronounced effect on  $E_{\text{elec}}$  than  $E_{\text{hole}}$  could be related to the fact that the LUMO of NH<sub>3</sub> is mainly located on H atoms, which can be directly influenced by



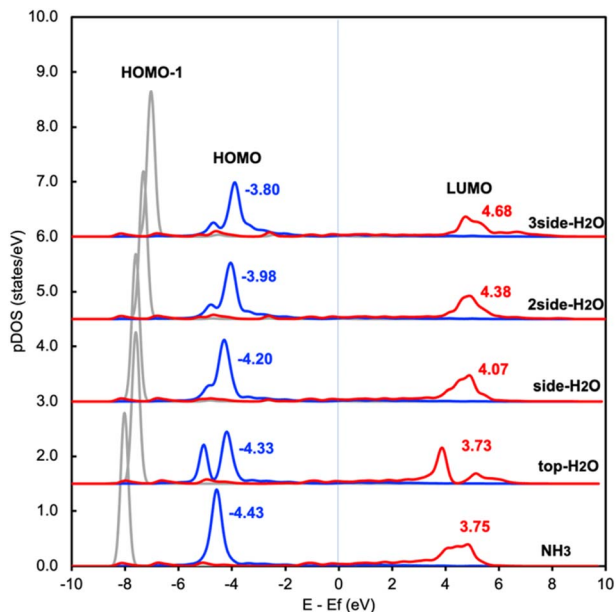


Fig. 2 Density of states projected onto frontier orbitals of  $\text{NH}_3$  adsorbed on  $\text{RuCu}_{83}$  with and without water. The center of the HOMO and LUMO of  $\text{NH}_3$  are shown above the main peak.

the H-bonds with water. Instead, the dominant part of the HOMO is located on the N atom in  $\text{NH}_3$ . As shown in Fig. 2, the centre of the HOMO shifts from  $-4.43$  eV in the absence of solvents to  $-4.20$ ,  $-3.98$ , and then  $-3.80$  eV with respect to the Fermi level of the system when one, two, and three  $\text{H}_2\text{O}$  molecules are added in the proximity of  $\text{NH}_3$ . The centre of the LUMO also moves to a higher energy, from  $3.73$  to  $4.07$ ,  $4.38$ , and  $4.68$  eV with respect to the Fermi level as the number of surrounding  $\text{H}_2\text{O}$  molecules increases. The local hydrogen-bonding environment around  $\text{NH}_3$  can thus modulate the excitation energy landscape for  $\text{NH}_3$  adsorbed over the Ru-doped Cu surface.

It is worth noting that another possibility of the change in the excitation energy involving the Fermi level is the modification of the work function. We have examined the change in work function of the catalyst surface with and without the surrounding  $\text{H}_2\text{O}$  molecules and find that there is only a slight change in the work function of the metal surface (Fig. S1†). In addition, there is no difference between the calculated core level energy of the  $1s$  orbital of Cu with and without  $\text{H}_2\text{O}$ . These results imply that the change in  $E_{\text{hole}}$  and  $E_{\text{elec}}$  shown in Fig. 1 is related to the adjustment of the HOMO and LUMO due to the H-bond formation between  $\text{NH}_3$  and  $\text{H}_2\text{O}$  rather than the change in the electronic properties of the metal substrate. In addition, we also investigated the effect of  $\text{H}_2\text{O}$  on the Cu surface instead of the Ru-doped Cu surface and find a similar trend between these two catalyst surfaces (Table S1 and S3†).

Next, the concept of solvent-induced local environment change in the interfacial charge transfer is explored with two other solvents – methanol and phenol (Fig. 3 and 4). Both molecules can form hydrogen bonds with the molecular adsorbate. Since the top adsorption configuration of water does

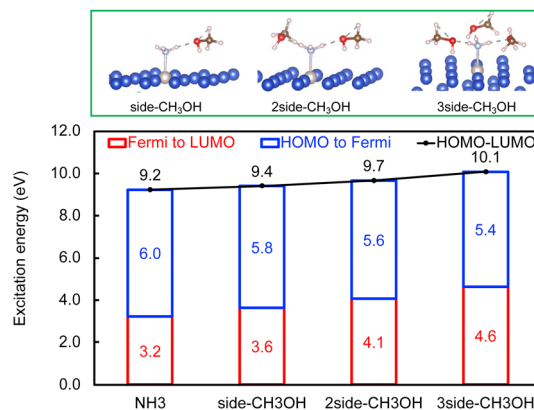


Fig. 3 The effect of  $\text{CH}_3\text{OH}$  on interfacial charge and energy transfer. The energy to excite an electron from the HOMO of  $\text{NH}_3$  to the Fermi level (blue), from the Fermi level to the LUMO of  $\text{NH}_3$  (red), and from the HOMO to the LUMO (black). Optimized configurations of  $\text{NH}_3$  with  $\text{CH}_3\text{OH}$  on  $\text{RuCu}_{83}$  are shown in the top panel.

not affect the interfacial charge transfer much, only the side adsorption of methanol and phenol is considered in the following discussion. As expected for a solvent with slightly higher basicity than that of water (Table S4†), methanol induces an increase of about  $0.5$  eV of  $E_{\text{elec}}$  per hydrogen bond formed between it and  $\text{NH}_3$ , which is higher by  $0.1$  eV per H-bond in the presence of water.

In contrast, the H-bond between phenol and  $\text{NH}_3$  only causes an increase of less than  $0.2$  eV of  $E_{\text{elec}}$  per H-bond (Fig. 4), which is lower than the value with water. A comparison of  $E_{\text{elec}}$  when water or phenol is introduced near  $\text{NH}_3$  is displayed in Fig. 5. When the number of solvent molecules increases, the bond length between N and Ru decreases due to the stronger interaction between  $\text{NH}_3$  and the catalyst surface. This increased interaction results from the upshifts of frontier orbitals and reduction of occupation in the antibonding states of Ru–N bond hybridization

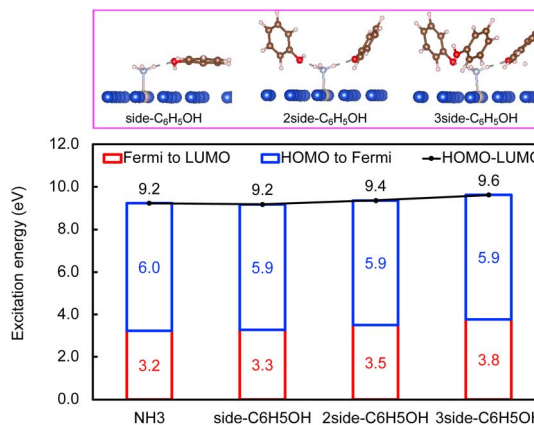


Fig. 4 The effect of  $\text{C}_6\text{H}_5\text{OH}$  on interfacial charge transfer of  $\text{NH}_3$ . The excitation energy required to excite an electron from the HOMO of  $\text{NH}_3$  to the Fermi level (blue), excite an electron from the Fermi level to the LUMO of  $\text{NH}_3$  (red), and from the HOMO to the LUMO (black). Optimized configurations of  $\text{NH}_3$  with  $\text{CH}_3\text{OH}$  on  $\text{RuCu}_{83}$  are shown in the top panel.





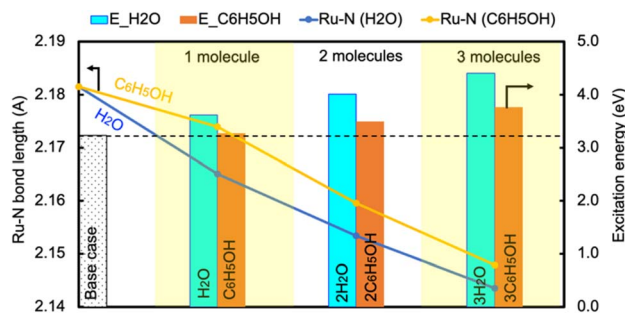


Fig. 5 The change in the excitation energy to excite an electron from the Fermi level to the LUMO of  $\text{NH}_3$  and the Ru–N bond length as the number of  $\text{H}_2\text{O}/\text{C}_6\text{H}_5\text{OH}$  molecules increases.

(Fig. S4†). Despite the similar trend of the bond length and the excitation energy, the calculated values suggest that the impact of water is more significant than that of phenol. Since this effect is determined by the electron donation capability of the solvents, the difference caused by the three protic solvents could be attributed to the varied basicity of solvents, which can be represented by the  $\beta$  (hydrogen bond acceptor) value. The solvent with a higher  $\beta$  value leads to a more significant upshift of the frontier orbitals and a larger change in the excitation energy. On the other hand, there is only a negligible change in  $E_{\text{hole}}$  when  $\text{NH}_3$  forms a H-bond with the surrounding phenol molecules despite the upward shift of the HOMO when more phenol molecules are introduced (Fig. 4 and Fig. S2†). The changes in  $E_{\text{hole}}$  are similar for water and methanol (*i.e.*, decreasing by 0.2 eV per H-bond), implying that the slight difference of basicity between water and methanol is not impactful enough to cause a difference in the generation of hot holes.

In plasmonic catalysis, apart from the interfacial charge transfer mechanism between the metal and adsorbed molecules, the plasmonic RET mechanism has also been reported as a main mechanism for some reactions.<sup>21,35–37</sup> The RET mechanism depends on the electronic excitation within the reactant's frontier orbitals, *e.g.*, the excitation from the HOMO to the LUMO, driven by the strong interfacial dipole–dipole coupling. Therefore, the RET excitation energies can be determined by using the HOMO–LUMO excitation energy ( $E_{\text{homo-lumo}}$ ). We first calculated the  $E_{\text{homo-lumo}}$  simply as the sum of the  $E_{\text{elec}}$  and  $E_{\text{hole}}$ . We find that the effect of solvents on the  $E_{\text{elec}}$  is stronger than that on the  $E_{\text{homo-lumo}}$  for the three solvents. That is, a change of 0.6 eV of the  $E_{\text{homo-lumo}}$  is observed compared to a change of 1.2 eV of  $E_{\text{elec}}$  when the number of water molecules increased from 0 to 3 (Fig. 2). In the case of methanol and phenol, the value of the  $E_{\text{homo-lumo}}$  vs.  $E_{\text{elec}}$  is 0.9 vs. 1.4 eV and 0.4 vs. 0.6 eV, respectively. Moreover, we also explicitly calculated the excitation energy from the HOMO to the LUMO using GPAW (Fig. S3†). Both approaches show the same trend for all studied solvents. The effect of the local environment induced by solvents on plasmonic catalysis thus depends on the mechanism (charge transfer vs. RET) of the chemical reaction.

Finally, we discuss the experimental implication derived from the calculated solvent effects on plasmonic excitation energy. The tunability of excitation energy achieved *via*

hydrogen bonding of solvents with the adsorbed reactants can be used to shift the molecular frontier orbitals to match the excitation energy with the supplied photon energy, especially in the case of a fixed photon energy source such as light-emitting diodes (LEDs). The solid–liquid interface also provides an opportunity for optimizing the selectivity/efficiency of plasmonic catalysis based on the uneven shift of the frontier orbitals (as demonstrated by the difference of LUMO and HOMO excitation). That is, these orbitals compete for hot carriers' population in the absence of solvents can in principle be lifted beyond the range of the photon energy by introducing a solvent so that the photon energy can be more selectively channeled into targeted orbitals. The solvents also provide dielectric screening that helps reduce the recombination of the hot carriers, increasing the lifetime of these carriers and thus the efficiency of plasmonic catalysis.

## Conclusions

By calculating the excitation energy of the direct interfacial charge transfer and resonance energy transfer in plasmonic catalysis at the solid–liquid interfaces, we show that the hydrogen bond formed between reactants and protic solvents can affect the excitation energy by adjusting the molecular frontier orbitals. Water, methanol, and phenol surrounding the adsorbed ammonia can increase the energy required for exciting an electron from the Fermi level to the LUMO and decrease that for HOMO-to-Fermi excitation. We find that the higher the basicity of the solvents, the stronger the impact on the electron transfer process. The solvents can also alter the HOMO–LUMO intramolecular excitation within the adsorbates, but to a lower extent than the Fermi-to-LUMO excitation. Therefore, the solvents may have effects on plasmonic excitation at varying extents depending on the plasmonic catalysis mechanism. Our results provide insights for understanding how solvents impact plasmonic catalysis and how to tune the reaction environment to enhance the efficiency of a photocatalyst. It is expected that this concept can be validated experimentally by measuring rates of a prototype reaction as a function of photon energy in different solvents.

## Author contributions

Tien Le performed the calculations and prepared the manuscript, and Bin Wang conceptualized the research, supervised the student to conduct this research, and revised the manuscript.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

This work was supported by the U.S. Department of Energy, Basic Energy Sciences, Chemical Sciences, Geosciences, and Biosciences Division, Condensed Phase and Interfacial



Molecular Science (Grant DE-SC0020300). The computations were performed at the OU Supercomputing Center for Education & Research and the National Energy Research Scientific Computing Center (NERSC), a U.S. Department of Energy Office of Science User Facility. Some calculations were performed with the Texas Advanced Computing Center (TACC) at The University of Texas at Austin supported by the NSF ACCESS/XSEDE.

## References

- 1 S. Ezendam, M. Herran, L. Nan, C. Gruber, Y. Kang, F. Gröbmeyer, R. Lin, J. Gargiulo, A. Sousa-Castillo and E. Cortés, *ACS Energy Lett.*, 2022, **7**, 778–815.
- 2 Y. Zhang, S. He, W. Guo, Y. Hu, J. Huang, J. R. Mulcahy and W. D. Wei, *Chem. Rev.*, 2017, **118**, 2927–2954.
- 3 M. L. Brongersma, N. J. Halas and P. Nordlander, *Nat. Nanotechnol.*, 2015, **10**, 25–34.
- 4 M. Sayed, J. Yu, G. Liu and M. Jaroniec, *Chem. Rev.*, 2022, **122**, 10484–10537.
- 5 S. B. Ramakrishnan, F. Mohammadparast, A. P. Dadgar, T. Mou, T. Le, B. Wang, P. K. Jain and M. Andiappan, *Adv. Opt. Mater.*, 2021, **9**, 2101128.
- 6 U. Aslam, V. G. Rao, S. Chavez and S. Linic, *Nat. Catal.*, 2018, **1**, 656–665.
- 7 E. Cortés and P. H. Camargo, *Plasmonic Catalysis: from Fundamentals to Applications*, John Wiley & Sons, 2021.
- 8 G. Li, B. Wang and D. E. Resasco, *ACS Catal.*, 2019, **10**, 1294–1309.
- 9 G. Li, B. Wang and D. E. Resasco, *Surf. Sci. Rep.*, 2021, **76**, 100541.
- 10 K. P. Rice, E. J. Walker Jr, M. P. Stoykovich and A. E. Saunders, *J. Phys. Chem. C*, 2011, **115**, 1793–1799.
- 11 A. Sanchez-Gonzalez, S. Corni and B. Mennucci, *J. Phys. Chem. C*, 2011, **115**, 5450–5460.
- 12 V. N. Nemykin, R. G. Hadt, R. V. Belosludov, H. Mizuseki and Y. Kawazoe, *J. Phys. Chem. A*, 2007, **111**, 12901–12913.
- 13 J. Neugebauer, C. Curutchet, A. Munoz-Losa and B. Mennucci, *J. Chem. Theory Comput.*, 2010, **6**, 1843–1851.
- 14 J. Moskowitz, R. Sindi and C. D. Geddes, *J. Phys. Chem. C*, 2020, **124**, 5780–5788.
- 15 S. K. Ghosh, S. Nath, S. Kundu, K. Esumi and T. Pal, *J. Phys. Chem. B*, 2004, **108**, 13963–13971.
- 16 K. O. Aruda, M. Tagliazucchi, C. M. Sweeney, D. C. Hannah, G. C. Schatz and E. A. Weiss, *Proc. Natl. Acad. Sci.*, 2013, **110**, 4212–4217.
- 17 A. Y. Bykov, D. J. Roth, G. Sartorello, J. U. Salmón-Gamboa and A. V. Zayats, *Nanophotonics*, 2021, **10**, 2929–2938.
- 18 L. R. Baker, C.-M. Jiang, S. T. Kelly, J. M. Lucas, J. Vura-Weis, M. K. Gilles, A. P. Alivisatos and S. R. Leone, *Nano Lett.*, 2014, **14**, 5883–5890.
- 19 V. G. Rao, U. Aslam and S. Linic, *J. Am. Chem. Soc.*, 2018, **141**, 643–647.
- 20 S. W. Lee, H. Kim and J. Y. Park, *Nano Lett.*, 2023, **23**, 5373–5380.
- 21 J. L. Bao and E. A. Carter, *J. Am. Chem. Soc.*, 2019, **141**, 13320–13323.
- 22 L. Zhou, D. F. Swearer, C. Zhang, H. Robotjazi, H. Zhao, L. Henderson, L. Dong, P. Christopher, E. A. Carter and P. Nordlander, *Science*, 2018, **362**, 69–72.
- 23 Y. Yuan, L. Zhou, H. Robotjazi, J. L. Bao, J. Zhou, A. Bayles, L. Yuan, M. Lou, M. Lou and S. Khatiwada, *Science*, 2022, **378**, 889–893.
- 24 G. Kresse and J. Furthmüller, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1996, **54**, 11169.
- 25 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865.
- 26 S. Grimme, S. Ehrlich and L. Goerigk, *J. Comput. Chem.*, 2011, **32**, 1456–1465.
- 27 J. J. Mortensen, L. B. Hansen and K. W. Jacobsen, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2005, **71**, 035109.
- 28 J. Enkovaara, C. Rostgaard, J. J. Mortensen, J. Chen, M. Dułak, L. Ferrighi, J. Gavnholt, C. Glinsvad, V. Haikola and H. Hansen, *J. Phys.: Condens. Matter*, 2010, **22**, 253202.
- 29 J. Gavnholt, T. Olsen, M. Engelund and J. Schiøtz, *Phys. Rev. B*, 2008, **78**, 075441.
- 30 T. Salavati-fard and B. Wang, *ACS Catal.*, 2022, **12**, 12869–12878.
- 31 T. Le, T. Salavati-fard and B. Wang, *ACS Catal.*, 2023, **13**, 6328–6337.
- 32 T. E. Markland and M. Ceriotti, *Nat. Rev. Chem*, 2018, **2**, 0109.
- 33 B. Civalleri, E. Garrone and P. Ugliengo, *J. Mol. Struct.: THEOCHEM*, 1997, **419**, 227–238.
- 34 A. D. Boese, *ChemPhysChem*, 2015, **16**, 978–985.
- 35 J. L. Bao and E. A. Carter, *ACS Nano*, 2019, **13**, 9944–9957.
- 36 J. M. P. Martirez and E. A. Carter, *J. Am. Chem. Soc.*, 2017, **139**, 4390–4398.
- 37 V. A. Spata and E. A. Carter, *ACS Nano*, 2018, **12**, 3512–3522.

