

## RESEARCH ARTICLE

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# Tm<sup>3+</sup>-mediated energy bridge in lead-free double perovskites: suppressing multiphonon relaxation for multifunctional photonic applications

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The growing demand for efficient and sustainable photonic materials has spurred significant interest in eco-friendly lead-free double perovskites, which exhibit exceptional photoluminescence properties for next-generation sensing and display technologies. However, the emission performance of the material is limited by multiphonon relaxation, which leads to low energy efficiency under complex operating conditions. In this study, the <sup>3</sup>H<sub>5</sub> energy level of Tm<sup>3+</sup> is strategically engineered as an energy bridge within lead-free double perovskites, effectively overcoming the constraints imposed by inefficient multiphonon relaxation in Er<sup>3+</sup> (<sup>4</sup>I<sub>11/2</sub> → <sup>4</sup>I<sub>13/2</sub>). This targeted modulation of energy transfer pathways results in a significant 50-fold enhancement in the red-to-green (R/G) emission ratio of Er<sup>3+</sup>, increasing from 0.09 to 4.58, accompanied by a pronounced chromaticity shift from green (0.278, 0.709) to red (0.678, 0.320). This tunable emission behavior is further modulated by laser power density and excitation wavelength (808 or 980 nm). Moreover, Gd<sub>2</sub>MgTiO<sub>6</sub> (GMTO):Er<sup>3+</sup>,Tm<sup>3+</sup> demonstrates excellent temperature-measuring ability based on the luminescence intensity ratio (<sup>2</sup>H<sub>11/2</sub>/<sup>4</sup>F<sub>9/2</sub> → <sup>4</sup>I<sub>15/2</sub>), achieving a maximum relative sensitivity of 1.03% K<sup>-1</sup> at 313 K. This work not only demonstrates a generalizable strategy to suppress multiphonon relaxation via lanthanide energy bridges but also paves the way for multifunctional photonic devices tailored to emerging applications in smart sensing, secure communication, and next-generation optoelectronics.

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## 1. Introduction

Lanthanide-doped luminescent materials have garnered considerable attention in various fields, including solid-state lighting, bio-imaging, information encryption, and solar energy conversion, due to their unique optical properties, such as narrow emission bands, long lifetimes, and excellent photostability.<sup>1–6</sup> The energy transfer (ET) process between lanthanide ions is a key mechanism that regulates the luminescent performance of these materials, as it enables the manipulation of emission intensity, color, and even multifunctional integration to be manipulated.<sup>7</sup> However, traditional ET strategies often suffer from limitations such as low

efficiency, single functionality and difficulty in achieving dynamic tuning, which restricts the further development of lanthanide-doped materials in advanced applications.<sup>8–10</sup> Therefore, exploring novel ET strategies to overcome these bottlenecks has become a pressing research goal. Recently, double perovskites with the general formula A<sub>2</sub>BB'O<sub>6</sub> have emerged as promising host matrices for lanthanide doping, owing to their superior structural stability, flexible chemical composition, and unique crystal field environment.<sup>11–13</sup> Compared to conventional single perovskites (ABO<sub>3</sub>) or other host materials (e.g., oxides, fluorides), double perovskites offer more adjustable cation sites (A, B, and B') with diverse coordination numbers and ionic radii, providing a versatile platform to design and optimize the spatial distribution and interaction of lanthanide ions.<sup>14,15</sup> This structural flexibility not only facilitates the construction of efficient ET pathways but also allows for the integration of multiple functionalities by tailoring the doping configuration and host composition. Thus, double perovskites hold great potential to serve as ideal hosts for realizing innovative ET strategies.<sup>16,17</sup> Currently, down-conversion (DC) luminescence in rare-earth-doped double perovskites has been extensively studied,<sup>18–20</sup> however, up-conversion

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(UC) luminescence properties remain underexplored. Among halide perovskites, materials such as Cs<sub>2</sub>NaBiCl<sub>6</sub> and other Cs-based compounds exhibit excellent UC luminescent performance, but their synthesis often requires prohibitively harsh conditions (HCl vapors,  $T > 150$  °C) and degrades under standard environmental stressors (RH > 30%,  $T > 85$  °C).<sup>21</sup> These challenges, coupled with potential toxicity concerns associated with certain halide components,<sup>22</sup> create a critical gap in the development of robust, scalable UC materials.

In this study, the oxide-based and lead-free halide perovskite luminescent material Gd<sub>2</sub>MgTiO<sub>6</sub> (GMTO) was selected as a promising host due to its low toxicity, easy preparation, and exceptional optoelectronic properties. Rare earth ions (Er<sup>3+</sup>, Tm<sup>3+</sup> and Yb<sup>3+</sup>) were strategically doped into the GMTO, enabling tunable UC luminescence under two near-infrared light sources (980 nm or 808 nm). Notably, the Tm<sup>3+</sup> ions act as an ET-bridge, thereby effectively modulating the red-to-green (R/G) emission ratio of the Er<sup>3+</sup> ions. Importantly, the UC luminescence can be controlled by adjusting the laser power density, allowing for the innovative design of a series of anti-counterfeiting and information storage patterns. Meanwhile, a series of highly sensitive thermometers based on GMTO:Yb<sup>3+</sup>,Er<sup>3+</sup>,Tm<sup>3+</sup> have been developed based on the luminescence intensity ratio (LIR) temperature sensing technology, which shows promise for real-time monitoring and temperature measurement in extreme working environments. These rare earth-doped GMTO materials exhibit exceptional optical properties and significant application potential, laying a solid foundation for the advancement of multifunctional luminescent materials.

## 2. Experimental section

### 2.1 Material and reagents

Gd<sub>2</sub>O<sub>3</sub> (99.8%, Beijing Chemical Factory), TiO<sub>2</sub> (AR, Sino pharm Chemical Reagent Co., Ltd), MgO (99.8%, Beijing Chemical Factory), H<sub>3</sub>BO<sub>3</sub> (99.9%, Produced in Shanghai, China), Yb<sub>2</sub>O<sub>3</sub> (99.99%, Aladdin), Er<sub>2</sub>O<sub>3</sub> (99.99%, Produced in Shanghai, China), Tm<sub>2</sub>O<sub>3</sub> (99.99%, Produced in Shanghai, China), anhydrous ethanol (99.99%, Aladdin). All the chemicals are used without any further purification.

### 2.2 Synthesis of Gd<sub>2</sub>MgTiO<sub>6</sub>:2%Er<sup>3+</sup>,xTm<sup>3+</sup> (x = 0%, 0.5%, 2%, 3%)

A series of Er<sup>3+</sup> and Tm<sup>3+</sup> co-doped GMTO compounds were synthesized by a high-temperature solid-state method. During sample preparation, 2% (by mass) H<sub>3</sub>BO<sub>3</sub> was added as a flux, followed by the addition of an appropriate amount of anhydrous ethanol during grinding. Since MgO was easily volatilized at high temperature, its content in the total mixture was set to be 10% in excess. The remaining raw materials were precisely weighed in accordance with the predetermined stoichiometric ratio. Subsequently, the accurately measured samples were transferred to an agate mortar and subjected to thorough grinding for 30 minutes. The heating process was carried out

in two steps. Initially, the temperature was raised from ambient to 1000 °C at a controlled rate of 10 °C min<sup>-1</sup> in air. Subsequently, it was further increased to 1400 °C at a slower rate of 5 °C min<sup>-1</sup> in air, and the samples were maintained at this elevated temperature for 8 h continuously. After completion of the heat treatment, the samples were left to cool naturally to room temperature. They were then retrieved and placed back into the agate mortar for an additional round of intensive grinding to ensure homogeneity. The well-prepared samples were then ready for subsequent testing and analysis.

### 2.3 Synthesis of GMTO:7%Yb<sup>3+</sup>,2%Er<sup>3+</sup>,yTm<sup>3+</sup> (y = 0.5%, 2%, 3%)

The sample preparation conditions and process were the same as those of GMTO:Er<sup>3+</sup>,Tm<sup>3+</sup>, with raw materials still weighed according to the stoichiometric ratio.

### 2.4 Characterization

Powder X-ray diffraction (XRD) measurements were performed using an X-ray diffractometer equipped with Cu K $\alpha$  radiation ( $\lambda = 0.15405$  nm). To visualize the morphological features and analyze the elemental distribution within the samples, a field-emission scanning electron microscope (FE-SEM, model: Regulus-8100, Hitachi) integrated with an energy-dispersive X-ray spectroscopy (EDS) detector was employed. The Raman spectrum was recorded by the Renishaw *in Via* Raman spectrometer (VERTEX 80v,  $\lambda_{\text{ex}} = 532$  nm). The luminescence spectra were recorded by an Andor SR-500i spectrometer (Andor Technology Co, Belfast, U.K.) equipped with a SR830 DSP lock-in amplifier and a DInGaAs1700-TE detector (1100 nm–1700 nm) and a DSi300-detector (380 nm–800 nm). The copper-constantan thermocouple and a temperature control system (TAP-02, orient-KOJI) were used to conduct temperature sensing. The decay lifetime curves of the samples were measured using an Omni- $\lambda$ 500 (Zolix, China) spectrometer with 808 and 980 nm (MDL-III-808/980-2 W, China) laser excitation.

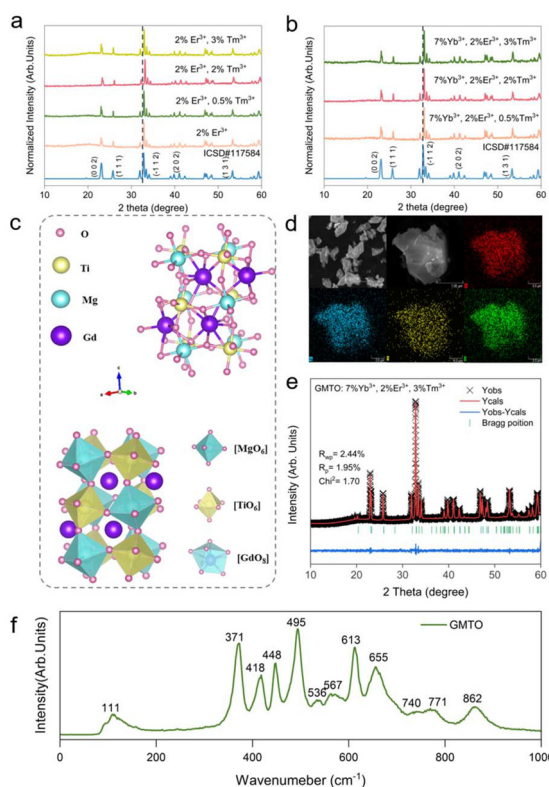
## 3. Results and discussion

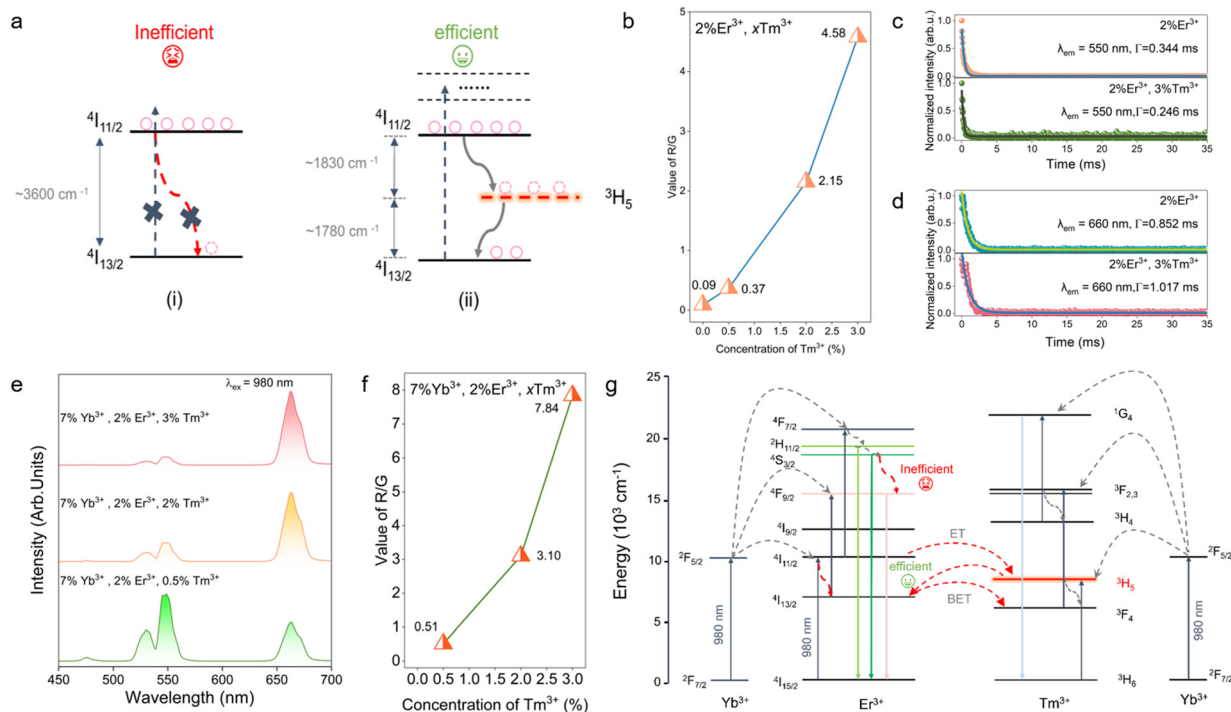
### 3.1 Material characterization

Ultra-bright and tunable double perovskite UC phosphors GMTO:Ln<sup>3+</sup> have been synthesized *via* a high-temperature solid-state method. It can be clearly observed that the reflection peaks of the samples completely match those of ICSD#117584 of GMTO (Fig. 1a and b). Considering that the dopant ions such as Er<sup>3+</sup> (0.89 Å), Tm<sup>3+</sup> (0.88 Å), and Yb<sup>3+</sup> (0.87 Å) have a smaller radius than Gd<sup>3+</sup> (0.94 Å), the ionic radius percentage formula eqn (1)

$$D_p = \frac{R_m(\text{CN}) - R_d(\text{CN})}{R_m(\text{CN})} \times 100\% \quad (1)$$

where  $R_m$  and  $R_d$  denote the ionic radii of the host and dopant ions, respectively, and CN denotes the coordination number of each ionic site (CN = 6 for the above ions), can be used. The





**Fig. 2** (a) A schematic illustration of the regulatory role of the  $\text{Tm}^{3+}\text{-}^3\text{H}_5$  energy level in UC emission of  $\text{Er}^{3+}$ . (b) R/G UC emission intensity ratio of  $\text{GMTO:Er}^{3+}, x\text{Tm}^{3+}$  ( $x = 0\text{--}3\%$ ) under 980 nm excitation. (c) UC decay lifetimes of green emission at 550 nm of  $\text{Er}^{3+}$  in the  $\text{GMTO:2\%Er}^{3+}$  and  $\text{GMTO:2\%Er}^{3+}, 3\%\text{Tm}^{3+}$  under 980 nm excitation, respectively. (d) UC decay lifetimes of red emission at 660 nm of  $\text{Er}^{3+}$  in the  $\text{GMTO:2\%Er}^{3+}$  and  $\text{GMTO:2\%Er}^{3+}, 3\%\text{Tm}^{3+}$  under 980 nm excitation, respectively. (e) UC emission spectra of  $\text{GMTO:7\%Yb}^{3+}, 2\%\text{Er}^{3+}, x\text{Tm}^{3+}$  ( $x = 0.5\text{--}3\%$ ) under 980 nm excitation. (f) R/G UC emission intensity ratio of  $\text{GMTO:Yb}^{3+}, \text{Er}^{3+}, x\text{Tm}^{3+}$  ( $x = 0.5\text{--}3\%$ ). (g) Schematic of UC emissions of  $\text{GMTO:Yb}^{3+}, \text{Er}^{3+}, \text{Tm}^{3+}$  upon 980 nm excitation.

emission intensity ( $I$ ) and the laser power ( $P$ ) is illustrated in Fig. S3b and d. This relationship can be adequately described by the following eqn (2):<sup>35</sup>

$$I \propto P^n \quad (2)$$

where  $n$  represents the number of photons in the UC transitions. Based on the linear fitting results, the  $n$  values for the emissions at 478 nm, 530 nm, 550 nm, and 660 nm in the 7%  $\text{Yb}^{3+}$ , 2%  $\text{Er}^{3+}$ , 3%  $\text{Tm}^{3+}$  system are 3.00, 2.08, 1.75, and 1.67, respectively. In contrast, in the  $\text{Er}^{3+}\text{-Tm}^{3+}$  co-doped system, the fitted  $n$  values for the 530 nm, 550 nm, and 660 nm are 2.01, 1.99, and 1.30, respectively, indicating that the  $\text{Tm}^{3+}\text{-}^1\text{G}_4 \rightarrow ^3\text{H}_6$  transition proceeds *via* a three-photon process, whereas the  $^2\text{H}_{11/2} \rightarrow ^4\text{I}_{15/2}$  and  $^4\text{S}_{3/2} \rightarrow ^4\text{I}_{15/2}$  transitions of  $\text{Er}^{3+}$  follow a two-photon UC mechanism. As shown in Fig. 2g, under 980 nm excitation, the  $\text{Er}^{3+}$  red emission arises from two distinct two-photon UC pathways. First,  $\text{Er}^{3+}$  ions are excited from  $^4\text{I}_{15/2}$  to  $^4\text{I}_{11/2}$ , followed by non-radiative (NR) relaxation to  $^4\text{I}_{13/2}$ ; subsequent excited-state absorption ( $^4\text{I}_{13/2} \rightarrow ^4\text{F}_{9/2}$ ) and radiative decay ( $^4\text{F}_{9/2} \rightarrow ^4\text{I}_{15/2}$ ) generate red emission. Secondly,  $\text{Yb}^{3+}$  ions undergo ground-state absorption ( $^5\text{F}_{7/2} \rightarrow ^5\text{F}_{5/2}$ ) and transfer energy resonantly to the  $^4\text{I}_{11/2}$  state of  $\text{Er}^{3+}$  by resonance energy transfer (RET). Further excited-state absorption (ESA) and NR relaxation populate the higher-lying  $^2\text{H}_{11/2}$ ,  $^4\text{S}_{3/2}$

and  $^4\text{F}_{9/2}$ , yielding green ( $^2\text{H}_{11/2}$ ,  $^4\text{S}_{3/2} \rightarrow ^4\text{I}_{15/2}$ ) and red ( $^4\text{F}_{9/2} \rightarrow ^4\text{I}_{15/2}$ ) emissions.

We therefore propose that introducing  $\text{Tm}^{3+}$  creates an intermediate  $^3\text{H}_5$  level between the  $\text{Er}^{3+}\text{-}^4\text{I}_{11/2}$  and  $^4\text{I}_{13/2}$  states, acting as an energy bridge. The original 3600  $\text{cm}^{-1}$  gap is effectively split into two smaller gaps of  $\sim 1830$  and  $\sim 1780$   $\text{cm}^{-1}$  by the  $^3\text{H}_5(\text{Tm}^{3+})$  level. Subsequent energy transfer (ET) and back energy transfer (BET) between  $\text{Er}^{3+}\text{-}^4\text{I}_{11/2}$  and  $\text{Tm}^{3+}\text{-}^3\text{H}_5$  promote population build-up in the  $\text{Er}^{3+}\text{-}^4\text{I}_{13/2}$  level, enhancing the probability of excitation to  $^4\text{F}_{9/2}$  (and other higher states) and thereby increasing the population density at  $^4\text{F}_{9/2}$ . Overall, incorporation of the  $\text{Tm}^{3+}\text{-}^3\text{H}_5$  level effectively mitigates inefficient multiphonon relaxation between  $^4\text{I}_{11/2}$  and  $^4\text{I}_{13/2}$ , establishes a bidirectional ET channel between  $\text{Er}^{3+}$  and  $\text{Tm}^{3+}$ , and significantly boosts population in the key intermediate  $^4\text{I}_{13/2}$  state, leading to more efficient UC and intensified red emission. Furthermore, as shown in Fig. 2d–f, additional sensitization by  $\text{Yb}^{3+}$  enables  $\text{Tm}^{3+}$  to undergo three-photon UC to yield blue emission ( $^1\text{G}_4 \rightarrow ^3\text{H}_6$ ). Taken together, these results demonstrate a rational strategy for tuning UC luminescence by constructing optimized ET channels that exploit the complementary energy-level structures of  $\text{Er}^{3+}, \text{Tm}^{3+}$ , and  $\text{Yb}^{3+}$ .

### 3.3 Photoluminescence properties upon 808 nm excitation

Owing to the rich ladder-like energy-level structure of  $\text{Er}^{3+}$ , different excitation wavelengths can selectively address distinct

intermediate states, offering an additional degree of freedom for tailoring UC pathways. Building on the 980 nm pumping results, we therefore explored how exploiting this multilevel structure under 808 nm excitation could further regulate the UC behavior. Remarkable UC luminescence was observed in GMTO co-doped with  $\text{Er}^{3+}$ ,  $\text{Tm}^{3+}$  and  $\text{Yb}^{3+}$  under 808 nm excitation. A systematic study of the UC emission spectra and associated ET dynamics revealed significant modulation of the emission color. As shown in Fig. 3a, b and S4a, we observed a systematic increase in the R/G emission ratio. A systematic study of the UC emission spectra and associated ET dynamics revealed significant modulation of the emission. Notably, lifetime analysis (Fig. S5) confirms this finding, under 808 nm excitation, the introduction of  $\text{Tm}^{3+}$  ions regulates the energy level populations of  $\text{Er}^{3+}$ . The lifetime of  $\text{Er}^{3+}$  green emission (550 nm,  $^4\text{S}_{3/2} \rightarrow ^4\text{I}_{15/2}$ ) shortens from 0.546 ms (single-doped) to 0.473 ms (co-doped), attributed to efficient energy transfer from  $\text{Er}^{3+} ^4\text{S}_{3/2}$  to  $\text{Tm}^{3+} ^3\text{H}_5$ . In contrast, the red emission lifetime (660 nm,  $^4\text{F}_{9/2} \rightarrow ^4\text{I}_{15/2}$ ) extends from 0.937 ms to 1.738 ms, as  $\text{Tm}^{3+}$  acts as an “energy bridge” to split the large energy gap between  $\text{Er}^{3+} ^4\text{I}_{11/2}$  and  $^4\text{I}_{13/2}$ , suppressing non-radiative multiphonon relaxation and enhancing the population of the  $^4\text{F}_{9/2}$  level. The power-dependent studies (Fig. S4b) provide critical insights into the UC mechanisms in GMTO:2% $\text{Er}^{3+}$ ,3% $\text{Tm}^{3+}$ . The  $n$  values for  $\text{Er}^{3+}$  emissions at 530 nm ( $n = 2.3$ ), 550 nm ( $n = 2.01$ ) and 660 nm ( $n = 1.75$ ) confirm a two-photon excitation process. Under 808 nm irradiation, the higher photon energy enables direct popu-

lation of the  $^4\text{I}_{9/2}$  state in  $\text{Er}^{3+}$ . Subsequent NR relaxation processes combined with the ET and BET between  $\text{Er}^{3+}$  and  $\text{Tm}^{3+}$ , facilitate efficient energy redistribution. Notably, the  $\text{Tm}^{3+} ^3\text{H}_4$  and  $^3\text{H}_5$  states are found to mediate ET from  $\text{Er}^{3+} ^4\text{I}_{9/2}$  to the  $^4\text{I}_{13/2}$  and  $^4\text{I}_{11/2}$  states. This is supported by the systematic decrease in  $\text{Er}^{3+}$  infrared emission intensity with increasing  $\text{Tm}^{3+}$  concentration (Fig. S6), indicating  $\text{Tm}^{3+}$ -mediated depletion of the  $^4\text{I}_{11/2}$  population through BET processes. The observed concentration-dependent quenching behavior confirms the critical role of  $\text{Tm}^{3+}$  in modifying the  $\text{Er}^{3+}$  excited state population.

In contrast to the absorption and pumping dynamics observed under 980 nm excitation, 808 nm photon absorption shows a significantly higher probability of populating the  $^2\text{H}_{11/2}$  and  $^4\text{F}_{7/2}$  states compared to the  $^4\text{F}_{9/2}$  level in  $\text{Er}^{3+}$  ions. As shown in Fig. 3c, under 808 nm excitation, red emission predominantly occurs *via* the relaxation of electrons from the high-energy  $^4\text{F}_{7/2}$  and  $^2\text{H}_{11/2}$  states rather than through a direct radiative transition from the  $^4\text{F}_{9/2}$  state ( $^4\text{F}_{9/2} \rightarrow ^4\text{I}_{15/2}$ ). The introduction of  $\text{Tm}^{3+}$  ions plays a crucial role in modifying the energy landscape, significantly reducing the energy difference between the  $^4\text{I}_{11/2}$  and  $^4\text{I}_{13/2}$  states of  $\text{Er}^{3+}$ . This reduction facilitates enhanced population at the  $^4\text{I}_{13/2}$  energy level through BET. In addition,  $\text{Tm}^{3+}$  ions efficiently absorb 808 nm photons, promoting transitions from the  $^3\text{H}_6$  ground state to the  $^3\text{H}_4$  excited state. Subsequent RET from  $\text{Tm}^{3+} (^3\text{H}_4)$  to  $\text{Er}^{3+} (^4\text{I}_{9/2})$  and NR relaxation processes populate the  $^4\text{I}_{11/2}$  state of  $\text{Er}^{3+}$ , ultimately increasing the population at the  $^4\text{I}_{11/2}$  energy



**Fig. 3** (a) UC emission spectra of GMTO:2% $\text{Er}^{3+}$ ,  $x\text{Tm}^{3+}$  ( $x = 0-3\%$ ) under 808 nm excitation. (b) R/G UC emission intensity ratio of GMTO:Er<sup>3+</sup>,  $x\text{Tm}^{3+}$  ( $x = 0-3\%$ ). (c) A proposed ET mechanism of GMTO:Er<sup>3+</sup>,  $x\text{Tm}^{3+}$  ( $x = 0-3\%$ ) under 808 nm excitation. (d) UC emission spectra of GMTO:7%Yb<sup>3+</sup>, 2%Er<sup>3+</sup>,  $x\text{Tm}^{3+}$  (0.5–3%) under 808 nm excitation. (e) R/G emission ratio of GMTO:Yb<sup>3+</sup>, 2%Er<sup>3+</sup>,  $x\text{Tm}^{3+}$  ( $x = 0.5-3\%$ ). (f) A proposed ET mechanism of GMTO:Yb<sup>3+</sup>, 2%Er<sup>3+</sup>,  $x\text{Tm}^{3+}$  upon 808 nm excitation.

level. Thus, under 808 nm excitation, the enhanced red luminescence results from the combined effects of (i) direct  $\text{Er}^{3+}$  excitation, (ii) efficient ET from  $\text{Tm}^{3+}$  to  $\text{Er}^{3+}$  and (iii) multi-pathway energy migration processes. These mechanisms promote the population of high-energy emitting states, thereby optimizing red light generation.

Fig. 3d demonstrates that the incorporation of  $\text{Yb}^{3+}$  ions enables blue emission from  $\text{Tm}^{3+}$  under 808 nm excitation. The doping concentrations of  $\text{Yb}^{3+}$  and  $\text{Tm}^{3+}$  were found to effectively modulate  $\text{Er}^{3+}$  luminescence, and the variation in the R/G emission ratio (Fig. S4c and Fig. 3e). Notably, the presence of  $\text{Yb}^{3+}$  necessitates consideration of ET processes between  $\text{Yb}^{3+}$  and  $\text{Tm}^{3+}$ . The power-dependent studies (Fig. S4d) provide critical insights into the UC mechanisms. The  $n$  values for  $\text{Er}^{3+}$  emissions at 530 nm ( $n = 2.44$ ), 550 nm ( $n = 2.08$ ) and 660 nm ( $n = 1.87$ ), as well as the  $n$  values for  $\text{Er}^{3+}$  emissions at 475 nm ( $n = 2.06$ ), confirm a two-photon excitation process. Therefore, a proposed ET mechanism for the  $\text{Yb}^{3+}$ ,  $\text{Er}^{3+}$ ,  $\text{Tm}^{3+}$  system under 808 nm excitation is illustrated in Fig. 3f. Our results suggest that the concentration of  $\text{Tm}^{3+}$  plays a dominant role in tuning the color of  $\text{Er}^{3+}$  UC emission. Additionally, we postulate that laser power is another critical parameter for luminescence modulation. Systematic investigations reveal that UC emission in both  $\text{Er}^{3+}$ ,  $\text{Tm}^{3+}$  and  $\text{Yb}^{3+}$ ,  $\text{Er}^{3+}$ ,  $\text{Tm}^{3+}$  systems can be effectively tuned through the selection of excitation wavelength (980 nm or 808 nm) coupled with precise modulation of the laser power (Fig. S7), the R/G emission ratio in the GMTO:2% $\text{Er}^{3+}$ , 0.5% $\text{Tm}^{3+}$  phosphor displays a distinct decreasing trend as the laser power density increases from 0.5 to 4 W, which differs from the concentration-dependent tuning effects. The multimodal and tunable luminescence characteristics of GMTO:Ln $^{3+}$  samples provide a robust theoretical and experimental foundation for future developments in information storage and fluorescence-based anti-counterfeiting technologies.<sup>36–39</sup> CIE coordinate analysis (Fig. S8) further quantified this wavelength-dependent and power-dependent chromaticity shift, revealing that the magnitude of color change was significantly greater under 808 nm excitation than under 980 nm excitation. These results also demonstrate the existence of an excitation wavelength/power-selective ET mechanism, highlighting the crucial role of pump wavelength optimization in tuning UC luminescence.

### 3.4 Temperature sensing and optical information storage

It is well established that lanthanide ions ( $\text{Ln}^{3+}$ ) and transition metal ions exhibit considerable potential for applications in temperature sensing and optical anti-counterfeiting owing to their distinctive luminescence characteristics. The luminescence intensity ratio (LIR) technique relies on the thermally induced redistribution of electrons between thermally coupled levels (TCLs) or non-thermally coupled levels (non-TCLs) of these ions. Nevertheless, such systems are inherently constrained by the relatively narrow energy gaps (typically  $\Delta E = 200\text{--}2000\text{ cm}^{-1}$ ), which can limit their sensitivity and operational range.<sup>40</sup> Although non-TCL-based LIR techniques overcome the limitation of  $\Delta E$ , they introduce another significant

challenge: the absence of a Boltzmann distribution between the two monitored energy levels prevents the establishment of thermal equilibrium. Consequently, the ET efficiency from higher to lower energy levels adversely affects the achievable temperature sensitivity. These fundamental limitations highlight that conventional single-approach LIR thermometry inevitably faces intrinsic drawbacks, underscoring the need for innovative alternative strategies.

Building upon these facts, we conducted a comprehensive investigation of LIR thermometry by synergistic integration of both TCLs and non-TCLs, thereby achieving enhanced temperature sensing accuracy. For experimental validation, we selected GMTO:7% $\text{Yb}^{3+}$ , 2% $\text{Er}^{3+}$ , 0.5% $\text{Tm}^{3+}$  phosphors, and monitored their luminescence over a broad temperature range (313–523 K) under 980 nm excitation (Fig. S9). As depicted in Fig. 4a, the temperature-dependent UC emission spectra revealed distinct thermal responses for different transitions (Fig. 4b). Specifically, the emissions from the  $^1\text{G}_4 \rightarrow ^3\text{H}_6$ ,  $^4\text{S}_{3/2} \rightarrow ^4\text{I}_{15/2}$ , and  $^4\text{F}_{9/2} \rightarrow ^4\text{I}_{15/2}$  transitions exhibited significant thermal quenching with increasing temperature, while the  $^2\text{H}_{11/2} \rightarrow ^4\text{I}_{15/2}$  transition demonstrated remarkable thermal stability. Based on these observations, we developed a triple-mode LIR thermometric model that incorporates both non-TCLs (LIR<sub>1</sub> and LIR<sub>3</sub> to eqn (3)) and TCLs mechanisms (LIR<sub>2</sub> to eqn (4)). In particular, LIR<sub>2</sub> utilized the classic Boltzmann-distribution-based TCL approach ( $^2\text{H}_{11/2}/^4\text{S}_{3/2}$ ):<sup>32,41</sup>

$$\text{LIR} = A + B \exp\left(\frac{C}{T}\right) \quad (3)$$

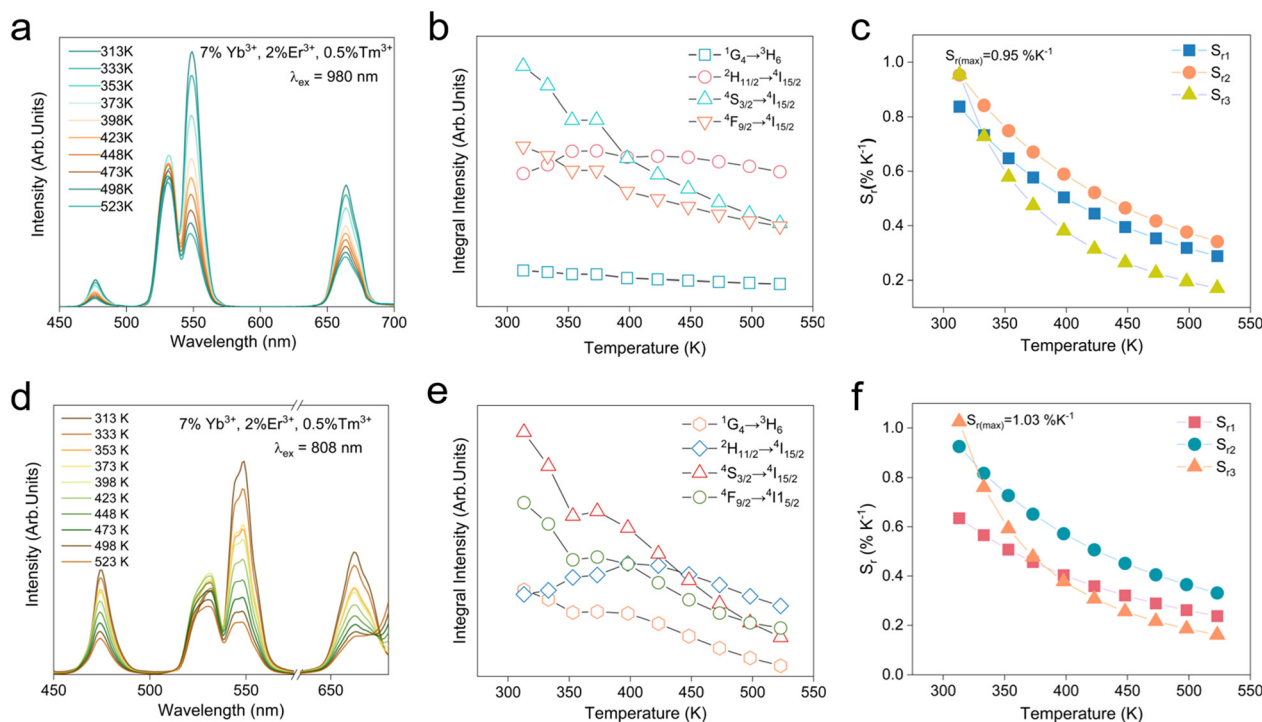
$$\text{LIR} = B \exp\left(\frac{C}{T}\right) \quad (4)$$

where  $T$  is the absolute temperature,  $A$ ,  $B$  and  $C$  are fitting constants. The results in Fig. S10(a–c) clearly showed that the three LIRs were highly fitted to eqn (3) and (4). In these models, LIR<sub>(1–3)</sub> represented the intensity ratio of emissions from different transitions:  $I_{(530/475\text{ nm})}$ ,  $I_{(530/550\text{ nm})}$  and  $I_{(530/660\text{ nm})}$ , respectively. It is well known that the relative sensitivity ( $S_r$ ) is commonly introduced to assess the heat-sensing capabilities of the synthesized materials. This parameter provides a precise metric for comparing thermal sensitivity characteristics among various thermometric materials.  $S_r$  is calculated as follows:<sup>33,34,42,43</sup>

$$S_r = \frac{1}{\text{LIR}} \left| \frac{\partial \text{LIR}}{\partial T} \right| \quad (5)$$

As presented in Fig. 4c, the  $S_{r(\text{max})}$  value of the three modes above reaches up to 0.95%  $\text{K}^{-1}$  in LIR<sub>2</sub> and LIR<sub>3</sub> at 313 K.

The temperature-dependent UC emission spectra of GMTO:7% $\text{Yb}^{3+}$ , 2% $\text{Er}^{3+}$ , 0.5% $\text{Tm}^{3+}$  under 808 nm excitation were systematically investigated (Fig. 4d). Notably, the thermal response trends of the emissions closely resembled those observed under 980 nm excitation (Fig. 4e). Distinct thermal behaviors were identified for the transitions originating from  $^1\text{G}_4 \rightarrow ^3\text{H}_6$  ( $\text{Tm}^{3+}$ ),  $^4\text{S}_{3/2} \rightarrow ^4\text{I}_{15/2}$  ( $\text{Er}^{3+}$ ), and  $^4\text{F}_{9/2} \rightarrow ^4\text{I}_{15/2}$  ( $\text{Er}^{3+}$ ), all of which exhibited pronounced thermal quenching. In con-



**Fig. 4** (a) Temperature-dependent UC emission spectra of GMTO:7%Yb<sup>3+</sup>,2%Er<sup>3+</sup>,0.5%Tm<sup>3+</sup> under 980 nm excitation. (b) The corresponding luminescence integral intensities of  $^1\text{G}_4 \rightarrow ^3\text{H}_6$ ,  $^2\text{H}_{11/2} \rightarrow ^4\text{I}_{15/2}$ ,  $^4\text{S}_{3/2} \rightarrow ^4\text{I}_{15/2}$ ,  $^4\text{F}_{9/2} \rightarrow ^4\text{I}_{15/2}$  under 980 nm excitation. (c) Calculated  $S_r$  values based on LIR model. (d) Temperature-dependent UC emission spectra of GMTO:7%Yb<sup>3+</sup>,2%Er<sup>3+</sup>,0.5%Tm<sup>3+</sup> under 808 nm excitation. (e) The corresponding luminescence integral intensities of  $^1\text{G}_4 \rightarrow ^3\text{H}_6$ ,  $^2\text{H}_{11/2} \rightarrow ^4\text{I}_{15/2}$ ,  $^4\text{S}_{3/2} \rightarrow ^4\text{I}_{15/2}$ ,  $^4\text{F}_{9/2} \rightarrow ^4\text{I}_{15/2}$  under 808 nm excitation. (f) Calculated  $S_r$  values based on the LIR model.

trast, the  $^2\text{H}_{11/2} \rightarrow ^4\text{I}_{15/2}$  ( $\text{Er}^{3+}$ ) transition demonstrated an initial increase in emission intensity with rising temperature, followed by quenching at elevated temperatures. These characteristics enabled the development of a robust multimodal LIR thermometry system under 808 nm excitation, similar to the three-mode LIR temperature sensing achieved under 980 nm excitation (Fig. S11(a-c)). Notably, as shown in Fig. 4f, the optimized  $S_r$  for LIR<sub>3</sub> reached  $1.03 \text{ \% K}^{-1}$  at 313 K. A comparison of the  $S_r$  obtained in this study with those reported in previously published works was provided in Table S2, further confirming the excellent temperature-sensing performance of the proposed thermometer. In addition, GMTO:Yb<sup>3+</sup>,Er<sup>3+</sup>,Tm<sup>3+</sup> exhibits outstanding thermal stability, as demonstrated by its performance over heating-cooling cycles between 333 K and 423 K. And the temperature uncertainty ( $\delta T$ ) is a critical parameter for optical temperature sensing, the formula is as follows:<sup>44</sup>

$$\delta T = \frac{1}{S_r} \frac{\delta \text{LIR}}{\text{LIR}} \quad (6)$$

The  $\delta T$  of GMTO:Yb<sup>3+</sup>,Er<sup>3+</sup>,Tm<sup>3+</sup> reflects the precision of its temperature measurement ability. As displayed in Fig. S12, the LIR value is obtained by 50 times consecutive tests at 333 K. The  $\delta T$  values of LIR are 0.53 K. Our findings demonstrated that GMTO:Yb<sup>3+</sup>,Er<sup>3+</sup>,Tm<sup>3+</sup> phosphors not only possessed dual-excitation capability (808/980 nm) but also enabled

high-performance multimode LIR-based temperature sensing under dual-mode excitation.

Leveraging the dual-mode excitation properties and the tunability of the UC luminescence *via* modulation of doping concentration and laser power (Fig. S13), we have developed strategies for advanced optical anti-counterfeiting and information storage applications. As illustrated in Fig. 5a, our design utilizes GMTO:Yb<sup>3+</sup>,Tm<sup>3+</sup> and GMTO:Yb<sup>3+</sup>,Er<sup>3+</sup>,Tm<sup>3+</sup> phosphors with precisely tailored doping concentrations, patterned into clover and sun motifs. These phosphors exhibit distinct chromatic responses under alternating 980 nm and 808 nm laser excitation, enabling dynamic color-switching within the encoded patterns. This functionality provides a robust platform for high-security optical encryption and multiplexed data storage. Fig. 5b schematically depicts the laser power-dependent color modulation of the UC emission in GMTO:2%Er<sup>3+</sup>,3%Tm<sup>3+</sup>, highlighting its exceptional potential for rewritable optical storage and related photonic applications. The capability to tune emission color in real time *via* excitation power allows for programmable spectral signatures, facilitating advanced optical encryption and establishing a new paradigm for dynamic information encoding. Furthermore, Fig. 5c demonstrates a sophisticated anti-counterfeiting application based on phosphor-mediated optical contrast. In this configuration, phosphors “I” (background) and “IV” (characters “J L U”) are selectively deposited to create a dual-responsive



**Fig. 5** Frontier applications of emission color-switchable double perovskite micro-particles. (a) Schematic of lucky clover for information encryption by using the samples with different Tm<sup>3+</sup> concentrations of GMTO:Er<sup>3+</sup>,Tm<sup>3+</sup> (0–3%). (b) Schematic of UC emission color-switching of GMTO:2%Er<sup>3+</sup>,3%Tm<sup>3+</sup> with different laser power. (c) Coating the phosphor of “I” and “IV” on the boards. Using “I” as background, “IV” as characters of “J L U”. Relying on the different changes of emission color with different excitation modes, selecting 980 nm as a pump source, vague green “J L U” will be displayed, while, selecting 808 nm as a pump source, clear yellow “J L U” will be displayed.

authentication label. Under 980 nm excitation, a faint green “J L U” pattern becomes visible, while switching to 808 nm irradiation induces a vivid red emission. This reversible high-contrast transition enables dynamic optical verification and establishes a new paradigm for secure, stimuli-responsive anti-counterfeiting technologies.

## 4. Conclusion

In conclusion, we have engineered a lead-free Gd<sub>2</sub>MgTiO<sub>6</sub>:Yb<sup>3+</sup>,Er<sup>3+</sup>,Tm<sup>3+</sup> double perovskite system centered on a Tm<sup>3+</sup>-mediated energy bridge, where the Tm<sup>3+</sup>-<sup>3</sup>H<sub>5</sub> state strategically mitigates the inefficient multiphonon relaxation of Er<sup>3+</sup> (<sup>4</sup>I<sub>11/2</sub> → <sup>4</sup>I<sub>13/2</sub>). This targeted ET modulation yields a remarkable 50-fold enhancement in the R/G emission ratio of Er<sup>3+</sup> (from 0.09 to 4.58) compared to Er<sup>3+</sup>-single-doped GMTO under 808 nm excitation, accompanied by a dramatic chromaticity shift from green (0.278, 0.709) to red (0.678, 0.320). Integrating Yb<sup>3+</sup> further enables multidimensional optical control: by tuning excitation wavelength (808/980 nm) and laser power, we achieve programmable multicolor emission, which exhibits robust potential for laser-encoded anti-counterfeiting and rewritable optical storage. Moreover, the optimized GMTO:7%Yb<sup>3+</sup>,2%Er<sup>3+</sup>,0.5%Tm<sup>3+</sup> phosphor demonstrates exceptional thermometric performance, with a maximum relative sensitivity of 1.03% K<sup>-1</sup> at 313 K. This work establishes a generalizable strategy to suppress multiphonon relaxation *via* lanthanide energy bridges, demonstrates a single lead-free matrix

integrating dynamic upconversion tuning, high-sensitivity thermal sensing, and optical encryption, and formulates design principles for phonon-engineered energy transfer in perovskite hosts. These advances that bridge spectral manipulation with multifunctional photonic applications provide a versatile platform for next-generation compact, high-efficiency optoelectronic devices in smart sensing and secure communication.

## Conflicts of interest

The authors declare no competing financial interest.

## Data availability

All data generated or analysed during this study are included in this published article and its supplementary information (SI). Supplementary information: detailed descriptions of LIR model, crystal structure, FE-SEM image and elemental mapping images, CIE chromatic coordinates and NIR emission spectrum. See DOI: <https://doi.org/10.1039/d5qj02407b>.

Data are available from the corresponding author on reasonable request.

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