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Optimization of promising water/wastewater treatment technologies requires significant resources in terms of time, labor and cost due to complex interactions between flow, microorganisms and reactions. The use of computational fluid dynamic simulations can shrink the possible parameter space, hence decreasing scale-up optimization costs.

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## CFD-accelerated bioreactor optimization: Reducing the hydrodynamic parameter space

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Optimization of bioreactor design can be experimentally challenging because of the complex interactions between hydrodynamic and biological processes. A promising prototyping strategy is the use of computational fluid dynamic (CFD) simulations to identify preferred hydrodynamic parameter spaces. In this work, we describe CFD simulations of flow in anaerobic fluidized-bed reactors (FBRs), with a focus on bed expansion and particle size. The results reveal regimes of putative high mass transfer where the diffusion layer thickness is impacted by a combination of flow velocity and particle collisions. These regimes are observed when bed expansion is narrowed from 10-70% (typically recommended) to 40-60%. Similarly, prospects for short circuiting are minimized by constraining the Archimedes number  $Ar$  of fluidized particles to  $Ar > 1000$  (as opposed to the common wisdom that “smaller is better”). When membranes are added to an FBR design, fluidized particles can effectively scour and clean membranes by constraining  $Ar$  to values  $Ar > 7000$  (a minimum is required). We conclude that CFD can provide valuable insights into reactor design and operation, reducing the hydrodynamic parameter space that must otherwise be explored by laboratory and pilot-scale validation thus decreasing time and cost for system optimization.

### 2 1 Introduction

1  
3 Sustainability is a grand challenge for the  
4 21st century<sup>1</sup>. Current human civilization  
5 is largely supported by linear economies in  
6 which resources are extracted, used, and dis-  
7 carded at end-of-life. This has created enor-  
8 mous challenges, increasing the need for circular  
9 economies based upon recycling and reuse<sup>2-5</sup>.

Microbial processes play an integral role in 10  
the removal of organic carbon and nutrients<sup>6-8</sup>. 11  
The prevailing technology first developed at the 12  
turn of the 20th century is activated sludge, a 13  
process that has since been modified to enable 14  
nutrient removal. Many emerging technologies 15  
cannot cross the “Valley of Death” because they 16  
treat tiny flows (in many cases, just milliliters 17  
per day) while adoption in practical applications 18  
may require treatment of tens of millions of liters 19  
per day. As biological and hydrodynamic com- 20  
plexity increases, the “Valley of Death” becomes 21

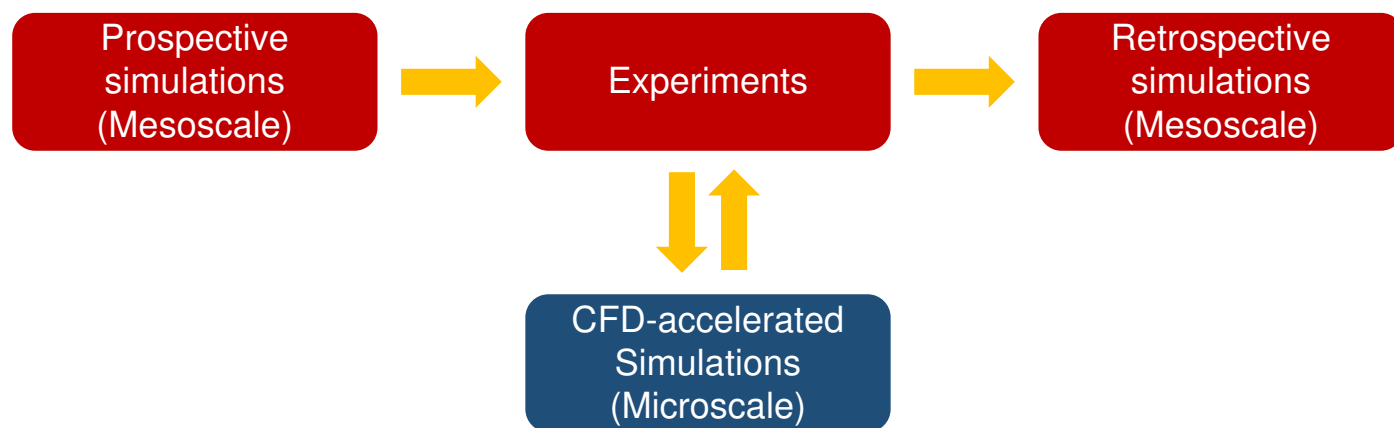
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22 deeper. An example would be bioelectrochemi- 64  
23 cal processes, such as microbial fuel cells<sup>9</sup> and 65  
24 microbial batteries<sup>10,11</sup>, technologies that have 66  
25 been demonstrated at bench- and, in some cases, 67  
26 pilot-scale but not full-scale. Academia is a likely 68  
27 source for such potentially disruptive innova- 69  
28 tion but lacks access to the facilities and fund- 70  
29 ing needed for long-term testing at a meaning- 71  
30 ful scale. To date, microbially-based technolo- 72  
31 gies have largely relied upon experiments for op- 73  
32 timization, but such testing is slow and costly. 74  
33 Not surprisingly, practitioners and utilities tend 75  
34 to innovate incrementally using existing systems. 76  
35 There is thus a great risk of locking-out innova- 77  
36 tion. A pathway for lower cost and more rapid 78  
37 scale-up of promising technology is needed. 79

38 In general, bioreactors can be classified as ei- 80  
39 ther dispersed growth systems, where substrate 81  
40 gradients are minimal, or attached-growth/floc- 82  
41 based systems, where appreciable substrate gra- 83  
42 dients drive diffusion of substrate into floc or at- 84  
43 tached biofilms as products diffuse out. The clas- 85  
44 sic dispersed-growth example is activated sludge 86  
45 (AS), a technology that efficiently removes or- 87  
46 ganic carbon in wastewater by converting solu- 88  
47 ble organics into CO<sub>2</sub>. Active microbial biomass 89  
48 is concentrated by settling or membrane separa- 90  
49 tion. These microorganisms are typically present 91  
50 as discrete cells or as floc within a size range 92  
51 of 2-20  $\mu\text{m}$ <sup>12,13</sup>. This size range confers two 93  
52 benefits: (1) the absence of a significant diffu- 94  
53 sion layer eliminates mass transfer limitations, 95  
54 and (2) small particles follow flow trajectories 96  
55 with negligible disruption of the overall hydro- 97  
56 dynamics. The combined effect of these pro- 98  
57 cesses is to weaken the dependence of biologi- 99  
58 cal activity on diffusion within floc or particles 100  
59 under well-mixed conditions, with minimal en- 101  
60 ergy consumption and minimal short-circuiting. 102  
61 To achieve excellent mixing, vigorous bubbling 103  
62 is used for aeration and/or mechanical mixing. 104  
63 To date, most research has focused on opti-

mization of process-related parameters such as 64  
HRT, solid retention time (SRT), and on micro- 65  
bial kinetic parameters with minimal hydrody- 66  
namic impacts. By excluding hydrodynamics in 67  
such models, reactor design and operation are 68  
greatly simplified. Examples include, but are not 69  
limited to, prediction of biological activity using 70  
ordinary differential equations rather than par- 71  
tial differential equations. In contrast to dis- 72  
persed growth reactions, attached growth and 73  
biofilm reactors are much more complex and 74  
more affected by process hydrodynamics. Exam- 75  
ples would include trickling filters, granular reac- 76  
tors, fluidized-bed reactors, membrane-aerated 77  
reactors, microbial fuel cells, and microbial bat- 78  
teries. For these examples, well-mixed conditions 79  
do not insure a reduction in the hydrodynamic 80  
parameter space. The parameter space in com- 81  
plex systems (such as microbial flocs, biofilm- 82  
coated particles and biofilm-coated porous mate- 83  
rials, and electrically conductive sponge) is much 84  
larger than in dispersed-growth reactors, and a 85  
thicker diffusion layer can increase mass transfer 86  
limitations. In addition, floc and BC-Ps do not 87  
necessarily follow the flow and flow-particle in- 88  
teractions can eventually alter the flow trajecto- 89  
ries, creating more complex hydrodynamics. As 90  
a result, biological activity and overall treatment 91  
efficiency depend upon local hydrodynamics. To 92  
optimize reactor design and operation, a quan- 93  
titative understanding of hydrodynamic-related 94  
parameters such as particle Reynolds number, 95  
porosity and Archimedes number is critical. By 96  
including hydrodynamic-related parameters, the 97  
total number of parameters (both hydrodynamic- 98  
and process-related) increases, resulting in dras- 99  
tic increases in resources in terms of cost and 100  
time and the number of experiments required for 101  
optimization. Simultaneously, the likelihood of 102  
obtaining optimal performance diminishes due to 103  
the high dimensional parameter space. 104



**Fig. 1** Workflow of CFD-accelerated scale-up. The red boxes indicate conventional CFD applications for reactor optimization. The blue box highlights iterative and integrative simulation and experiments. Mesoscale refers to simulations where the size of the computational domain is on the order of meters and the shape resembles an industrial reactor. Microscale refers to simulations at scales much smaller than reactors, on the order of millimeters or even microns.

## 105 2 Computational fluid dynamics

106 Computational fluid dynamics (CFD) uses nu-  
 107 merical methods to study problems that involve  
 108 fluid flows. Over the past few decades, advances  
 109 in computational power and methods have ex-  
 110 panded the range of problems that can be ad-  
 111 dressed using CFD. A review by Karpinska and  
 112 Bridgeman<sup>14</sup> has evaluated different strategies  
 113 and models for optimization of wastewater treat-  
 114 ment. In wastewater treatment (Figure 1), CFD  
 115 can (1) prospectively preview macroscopic reac-  
 116 tor hydrodynamics and (2) retrospectively im-  
 117 prove current design and operation. Studies are  
 118 carried out sequentially by first comparing simu-  
 119 lations with experimental results (i.e., historical  
 120 results from an existing system in retrospective  
 121 applications or from a similar system in prospec-  
 122 tive applications) then conducting simulations by  
 123 varying a parameter of interest. These studies are  
 124 mostly conducted at the mesoscale where the size  
 125 of the computational domain is on the order of  
 126 meters and the shape resembles an industrial re-  
 127 actor. The main disadvantage of this approach is  
 128 loss of microscale information where microscale  
 129 refers to simulations investigating scales that are

much smaller than reactors and in the order of 130  
 micron or millimeters (i.e, interactions between 131  
 small particles on an industrial fluidized-bed re- 132  
 actor). As such, most research adopting this ap- 133  
 proach focuses on macroscopic properties such 134  
 as flow short-circuiting, reactor mixing and oxy- 135  
 gen transfer efficiency using commercial soft- 136  
 ware<sup>15–17</sup>. For dispersed growth reactors, with 137  
 a reduced hydrodynamic parameter space, CFD 138  
 studies have focused on aeration and mixing. An- 139  
 other useful CFD application is disinfection. In 140  
 this case, short-circuiting is minimized to enable 141  
 efficient pathogen removal, a goal that must be 142  
 balanced against the need for minimization of 143  
 disinfection byproducts<sup>18–20</sup>. In these applica- 144  
 tions, the simulations are reactor-specific, so the 145  
 knowledge gained from one system does not nec- 146  
 essarily translate to another. 147

In more complex reactors (attached-growth or 148  
 floc-based), the parameter space increases sig- 149  
 nificantly due to local interactions between hy- 150  
 drodynamics and biological activity. Aerobic flu- 151  
 idized bed reactors require aeration and mixing 152  
 simulations, but also guidelines on bed expan- 153  
 sion, particle size, and other carrier properties of 154

155 interest. Focusing on aeration and mixing alone  
156 is unlikely to result in optimized design and per-  
157 formance. In addition, because most interac-  
158 tions occur at the microscale, ignoring microscale  
159 properties (e.g., floc diffusion layer or floc-floc in-  
160 teractions) with a singular focus on macroscopic  
161 properties (bed expansion or mixing) is also un-  
162 likely to lead to correct conclusions. Understand-  
163 ing of systems that have shared physics can expe-  
164 dited translation across systems. Resources spent  
165 on understanding of one system would benefit  
166 other similar systems. An example is the effect  
167 of Archimedes number (a combination of particle  
168 and fluid properties) in upward flow reactors. A  
169 quantitative understanding of this number would  
170 be beneficial to both non-fluidized granular re-  
171 actors and fluidized-bed reactors. The focus of  
172 this approach is not to identify the exact values  
173 for optimized parameters but rather to reduce  
174 the parameter space within which optimized pa-  
175 rameters fall. By narrowing this space, reactor-  
176 specific experimental studies can be more tar-  
177 geted, enabling more efficient optimization and  
178 scale-up with fewer resources.

179 In this paper, we propose a new framework for  
180 bioreactor optimization: a computational strat-  
181 egy in which CFD is used to understand funda-  
182 mental interactions involving fluid flow, particles,  
183 microorganisms, membranes, and other porous  
184 materials. We envision that this approach will  
185 enable deeper insight into the underlying physics  
186 and accelerated optimization. We use the Staged  
187 Anaerobic Fluidized-bed Membrane bioreactors  
188 (SAF-MBR) as a case study to demonstrate the  
189 feasibility and potential of this framework.

### 190 3 A case study: Staged Anaerobic Fluidized- 191 bed Membrane Bioreactors

192 The SAF-MBR is a recently developed biocarrier-  
193 based anaerobic treatment technology<sup>21,22</sup>. Aer-  
194 ation is eliminated because the active microor-  
195 ganisms are obligate anaerobes that do not toler-

ate oxygen. Energy is recovered as methane, en-196  
abling net energy-positive secondary treatment 197  
of domestic wastewater<sup>23</sup>. Because they are 198  
slow-growing, the anaerobes also generate fewer 199  
biosolids for disposal. These properties make the 200  
SAF-MBR more attractive than conventional aer-201  
obic processes, such as AS<sup>24</sup>. 202

The SAF-MBR consists of two reactors in series 203  
with a conventional anaerobic fluidized-bed reac-204  
tor (AFBR) followed by an anaerobic membrane 205  
bioreactor (AnMBR). AFBRs have been widely 206  
used to treat industrial wastewater where the 207  
Chemical Oxygen Demand (COD) and Biochem-208  
ical Oxygen Demand (BOD) are much greater 209  
than domestic wastewater. In SAF-MBR 1.0, the 210  
AFBR discharges to a particle-sparged membrane 211  
bioreactor (P-MBR), in which fluidized granu-212  
lar activated carbon (GAC) functions as both a 213  
biocarrier of slow-growing microorganisms (in-214  
side the GAC pores) and as a scouring agent for 215  
cleansing of membranes and prevention of bio-216  
fouling<sup>21,22,24,25</sup>. This strategy successfully con-217  
trolled membrane biofouling in a pilot-scale SAF-218  
MBR<sup>21,26</sup>, but also led to particle abrasion and 219  
damage of the polymeric membranes<sup>27</sup>. As noted 220  
by Shin et al.<sup>27</sup>, the GAC used in the P-MBR con-221  
tained two size fractions - one at 1.18 - 1.4mm 222  
(29%) and a second at 1.70 -4.00mm (47%). 223  
Significant membrane damage occurred in the 224  
lower region of the membranes, and this dam-225  
age was attributed to the larger GAC fraction. 226  
In subsequent pilot-scale tests of SAF-MBR 2.0<sup>23</sup>, 227  
membrane sparging was accomplished with bio-228  
gas bubbles rather than solid particles. 229

The hydrodynamics of fluidized-bed reactors 230  
have been investigated experimentally<sup>28,29</sup> and 231  
with simulations<sup>30-34</sup>, but membrane bioreactor 232  
studies of microbial activity have largely focused 233  
on experimental testing<sup>35,36</sup>. These studies do not 234  
track particle dynamics at high volume fraction 235  
(low porosity), but instead focus on macroscopic 236  
behavior such as fluidization stability and expan-237

238 sion<sup>37,38</sup>. The range of bed expansion in flu-  
 239 idized beds fluctuates between 20% and 70%  
 240 with a qualitative understanding that low expan-  
 241 sion leads to flow short-circuiting and high ex-  
 242 pansion leads to biofilm loss. The optimal bed  
 243 expansion or porosity is thus an open question.  
 244 At present, most studies focus on bed expansion  
 245 without considering the impacts of particle prop-  
 246 erties such as Archimedes  $Ar$  (or Galilei  $Ga$ ) num-  
 247 ber on the optimal bed expansion where  $Ar$  is de-  
 248 fined as

$$Ar = Ga^2 = \frac{(s-1)gd_p^3}{\nu^2}, \quad (1)$$

249 where  $s = \rho_p/\rho_f$  is particle-fluid density ratio,  $d_p$   
 250 is the clean particle diameter,  $g$  is the gravita-  
 251 tional acceleration and  $\nu$  is the kinematic viscos-  
 252 ity of water. Qualitatively, small particles are pre-  
 253 ferred to enable more efficient mass transfer and  
 254 enhanced biological activity. Aslam et al.<sup>26</sup> stud-  
 255 ied the effects of three different particles (PET  
 256 beads, silica and GAC) on membrane scouring ef-  
 257 ficiency and concluded that PET beads are best.

### 258 3.1 Upflow velocity and porosity

259 In fluidized-bed reactors, upflow velocity con-  
 260 trols bed expansion and hence porosity. Under-  
 261 standing particle dynamics as a function of poros-  
 262 ity gives important insights into the biological ac-  
 263 tivity and design and operation of reactors. Re-  
 264 cently, Yao et al.<sup>32</sup> investigated particle dynam-  
 265 ics by varying upflow velocity in simulations of  
 266 a monodispersed/single-size fluidized bed with  
 267 particle properties similar to those of pilot-scale  
 268 and lab-scale reactors<sup>21</sup>. Within FBRs, poros-  
 269 ity controls both the horizontal mixing and col-  
 270 lisions between particles. Since no horizontal  
 271 flow is generated at the inlet, horizontal mixing is  
 272 mainly due to momentum transfer from the verti-  
 273 cal to horizontal directions due to particle fluctu-  
 274 ations. At low porosity, most fluctuations are in-  
 275 duced by weak collisions. At intermediate poros-

ity, collisions and hydrodynamic effects become  
 276 equally important, leading to an increase in par-  
 277 ticle velocity fluctuations and stronger collisions.  
 278 At high porosity, hydrodynamic effects dominate,  
 279 and collisions are diminished. 280

#### 3.1.1 Hypothetical impacts on biofilm de- 281 tachment 282

Accurate quantification of biofilm detachment  
 283 rate provides valuable information in modeling  
 284 biofilm reactor dynamics, such as the height of  
 285 expanded beds and insight into reaction- and  
 286 mass-transfer limitations<sup>39</sup>. The overall biofilm  
 287 detachment rate  $b_t$  is modeled as a combination  
 288 of first-order cell decay and mechanical detach-  
 289 ment: 290

$$b_t = b + b_{det}, \quad (2)$$

where  $b$  is the first-order cell decay constant and  
 291  $b_{det}$  is the mechanical detachment rate. Typically,  
 292  $b \ll b_{det}$  for most the engineered applications  
 293 such that  $b_t \approx b_{det}$ . There are two types of detach-  
 294 ment (continuous and discrete) and three mech-  
 295 anisms (shear stress, abrasion, and sloughing).  
 296 The shear stress is due to flow, while abrasion is  
 297 due to collisions between particles. Since slough-  
 298 ing is typically described as a discrete probabilis-  
 299 tic event that might lead to breakup of the entire  
 300 biofilm<sup>40</sup>, most models do not consider it. Chang  
 301 et al.<sup>41</sup> modeled  $b_{det}$  as 302

$$b_{det} = -3.14 + 0.0335C_p + 19.3Re_{p,b} - 3.46\sigma, \quad (3)$$

where  $C_p$  is the particle concentration in the flu-  
 303 idized bed,  $Re_{p,b} = u_0d_b/\nu$  is the biofilm-covered  
 304 particle Reynolds number,  $u_0$  is the upflow veloc-  
 305 ity in the fluidized bed,  $d_b$  is the diameter of the  
 306 BC-P and  $\sigma$  is the shear stress. The author as-  
 307 sumed that  $C_p$ ,  $Re_{p,b}$  and  $\sigma$  account for abrasion,  
 308 turbulence and shear stress, respectively. The  
 309 main challenge with this model is related to de-  
 310 coupling flow ( $Re_{p,b}$ ) and abrasion effects ( $C_p$ )  
 311

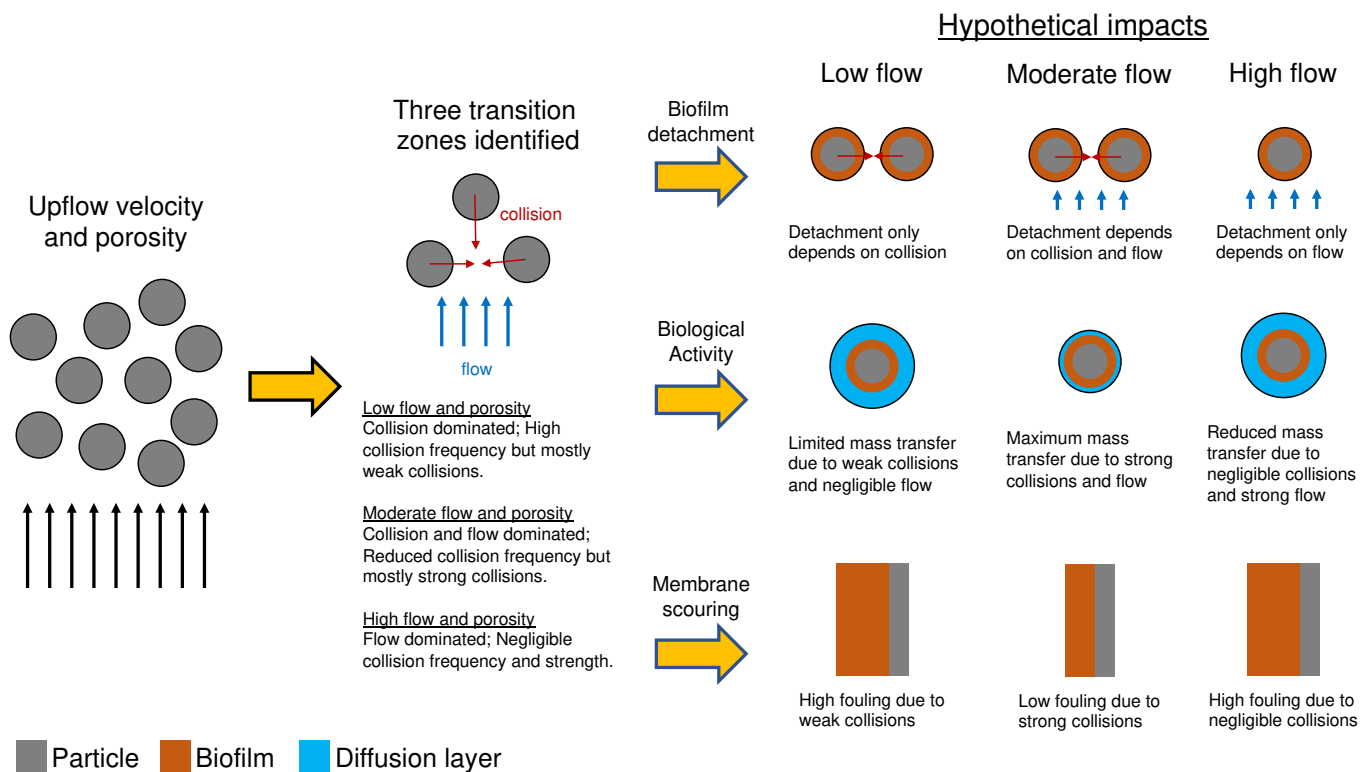


Fig. 2 Hypothetical impacts of upflow velocity and porosity on fluidized-bed reactor modeling, design and operation.

312 where both  $C_p$  and  $Re_{p,b}$  are functions of poros-  
313 ity.

314 Nicoletta et al.<sup>42</sup> constructed an empirical  
315 model for a fluidized-bed reactor based on di-  
316 mensional analysis and showed that the normal-  
317 ized detachment rate  $\hat{b} = d_p \tilde{b}_{det} / \rho_f v$  is given by

$$\hat{b} = 1.95 \times 10^{-10} Re_{p,c}^{1.49} d_*^{2.67}, \quad (4)$$

318 where  $\tilde{b}_{det}$  is the amount of biofilm detached per  
319 unit area and time,  $Re_{p,c} = d_p u_0 / \nu$  is the clean  
320 particle Reynolds number and  $d_* = d_b / d_p$  is the  
321 diameter ratio of a biofilm-covered to a clean par-  
322 ticle. An interpretation of this model is that  $Re_{p,c}$   
323 represents the effects of clean particle-related  
324 flow whereas  $d_*$  includes the effects of biofilm  
325 thickness on the detachment rate. The effect of  
326 flow includes both turbulence and abrasion as  
327 compared to equation 3<sup>41</sup>. Interestingly, detach-  
328 ment rates that were orders of magnitude higher  
329 were observed for  $d_* = 3$  and  $Re_{p,c} = 2.3 - 2.7$

which cannot be explained by this model, al-330  
though one plausible explanation is that detach-331  
ment occurs near the inlet where turbulence is332  
strongest. Overall, this model better parameter-333  
izes the mechanical biofilm detachment rate in334  
the sense that there is much less cross-correlation335  
between parameters. Instead of considering both336  
flow and abrasion, Gjaltema et al.<sup>43</sup> assume that337  
abrasion is the only dominant detachment mech-338  
anism in an airlift reactor and used a model to339  
estimate the energy of abrasion. 340

All of the above models are unable to decouple341  
flow and abrasion or are limited to abrasion. By342  
comparing high- and low-strength FBRs, Shin et343  
al.<sup>44</sup> successfully modeled low-strength FBRs by344  
assuming small  $b_{det}$ . By examining particle dy-345  
namics in a fluidized bed, Yao et al.<sup>32</sup> discovered346  
that collisions dominate over hydrodynamic ef-347  
fects at low porosity. At intermediate porosity,348  
both collisions and hydrodynamic effects are im-349  
portant while hydrodynamic effects dominate at350



351 higher porosity. The combination of effects of  
 352 wastewater strength and hydrodynamics implies  
 353 that the biofilm detachment rate in a fluidized-  
 354 bed reactor can be modeled with a stepwise func-  
 355 tion

$$b_{det} = \begin{cases} b_{col}f(C_{COD}), & \text{for } \varepsilon < \varepsilon_{c1}, \\ (b_{col} + b_{hydro})f(C_{COD}), & \text{for } \varepsilon_{c1} \leq \varepsilon < \varepsilon_{c2}, \\ (b_{hydro})f(C_{COD}), & \text{for } \varepsilon \geq \varepsilon_{c2}, \end{cases} \quad (5)$$

356 where  $f(C_{COD})$  is a function that relates  $b_{det}$  and  
 357 chemical oxygen demand (COD) concentration  
 358 of wastewater,  $\varepsilon$  is the porosity,  $\varepsilon_{c1}$  and  $\varepsilon_{c2}$  are  
 359 critical porosities representing the boundaries of  
 360 the different regimes, and  $b_{col}$  and  $b_{hydro}$  are  
 361 the detachment rates associated with collisional  
 362 and hydrodynamic effects. Biofilm detachment  
 363 is likely maximized to the coexistence of two  
 364 different mechanisms at intermediate Reynolds  
 365 numbers. Although further experimental valida-  
 366 tion is required, observing how particle dynam-  
 367 ics change in fluidized bed simulations provides  
 368 insight into biofilm detachment and how detach-  
 369 ment rates might best be modeled. Furthermore,  
 370 because  $\varepsilon_{c1}$  and  $\varepsilon_{c2}$  vary with particle properties  
 371 such as diameter and density, a universal scaling  
 372 law can be developed that confirms and general-  
 373 izes this approach for different particle diameters  
 374 and densities.

### 375 3.1.2 Hypothetical impacts on mass transfer 376 and biological activities

377 Fluidized-bed reactors are known for their ex-  
 378 cellent mass transfer rate. When applied for  
 379 wastewater treatment, the AFBR can either be  
 380 mass transfer limited or reaction rate limited.  
 381 The latter usually occurs in shallow and fully-  
 382 penetrated biofilms where substrates are me-  
 383 tabolized at a much slower rate than diffusion  
 384 enables. Buffiere et al.<sup>35</sup> discovered that the  
 385 methanogenic step requires deep biofilms while  
 386 acidogenesis only requires shallow biofilms for

treatment of high-strength wastewater. Conflict-387  
 ing results have been reported where increases in 388  
 flow rate can either increase<sup>45</sup> or decrease<sup>46</sup> the 389  
 mass transfer rate. Nicoletta et al.<sup>47,48</sup> discov-390  
 ered that mass transfer of biofilm-covered parti-391  
 cles in airlift reactors is roughly 15% lower than 392  
 that of clean particles. 393

(5) Due to the serial nature of process kinetics, 394  
 with mass transfer preceding biochemical kinet-395  
 ics, overall reactions can be mass transfer-limited 396  
 when the reaction step is fast or they can be 397  
 reaction-limited when the mass transfer step is 398  
 slow<sup>35</sup>. In AFBRs, the particle Reynolds num-399  
 ber based on superficial velocity leads to colli-400  
 sions and hydrodynamic effects that control mass 401  
 transfer. Higher flow rates reduce the thickness 402  
 of the diffusion layer thereby enhancing mass 403  
 transfer. Similarly, more frequent collisions dis-404  
 rupt the diffusion layer reducing its thickness 405  
 in fluidized-bed electrochemical cells<sup>49</sup>. The ef-406  
 fect of collisions alone can be accurately de-407  
 scribed by the collision pressure which is known 408  
 to have a maximum and zeros for both single-409  
 particle ( $\varepsilon \approx 1$ ) and close-packed reactors ( $\varepsilon \approx 410$   
 0.4)<sup>43,50</sup>. After close examination of particle dy-411  
 namics in fluidized bed simulations, Yao et al.<sup>32</sup> 412  
 suggested that mass transfer is most likely maxi-413  
 mized within the intermediate porosity regime at 414  
 which point collisions and hydrodynamic factors 415  
 are equally important, leading to optimal biolog-416  
 ical performance. Although not yet experimen-417  
 tally validated, pilot- and lab-scale reactors oper-418  
 ated at this intermediate porosity (bed expansion 419  
 of 40% to 60%) have achieved optimal treatment 420  
 performance<sup>21,23,51</sup>. 421

### 422 3.1.3 Hypothetical impacts on membrane 423 fouling control

The primary role of the P-MBR is to retain par-424  
 ticulate biodegradable organic matter in the re-425  
 actor because more time is required for hydroly-426  
 sis. The main challenge is to prevent membrane 427

biofouling, which can be accomplished by either particle- or gas-sparging. Particle-sparged operation enables low energy demand<sup>52</sup> but can lead to severe membrane damage in the lower region of the reactor<sup>53</sup>. Moreover, due to non-uniform particle sizes, the fluidized bed in the P-MBR forms segregated layers of particles with larger particles (2-4 mm) located at the bottom of the bed. Yao et al.<sup>32</sup> found that the maximum collision frequency is attained at intermediate porosity for 2 mm particles. Low porosity is characterized by more frequent weak collisions while high porosity is dominated by flow rather than collisions. Comparing the collision frequency as a function of porosity by Yao et al.<sup>32</sup> and the membrane integrity study by Shin et al.<sup>27</sup> with similar particle sizes, the porosity of the lower region in the pilot-scale P-MBR corresponds to the region of maximum effective collisions from the simulations. This result implies that membrane scouring efficiency can be controlled by varying porosity, therefore the bed expansion. Maximum membrane scouring is attained at the porosity with maximum collisions. To avoid membrane damage, varying the porosity to deviate from the maxima say, by reducing or increasing it, is likely to eliminate membrane damage. As discussed below, instead of switching to alternative membrane fouling control methods, studying the effects of the Archimedes number enables a retrospective modification to both new and existing reactors.

### 3.2 Particle properties and the Archimedes number

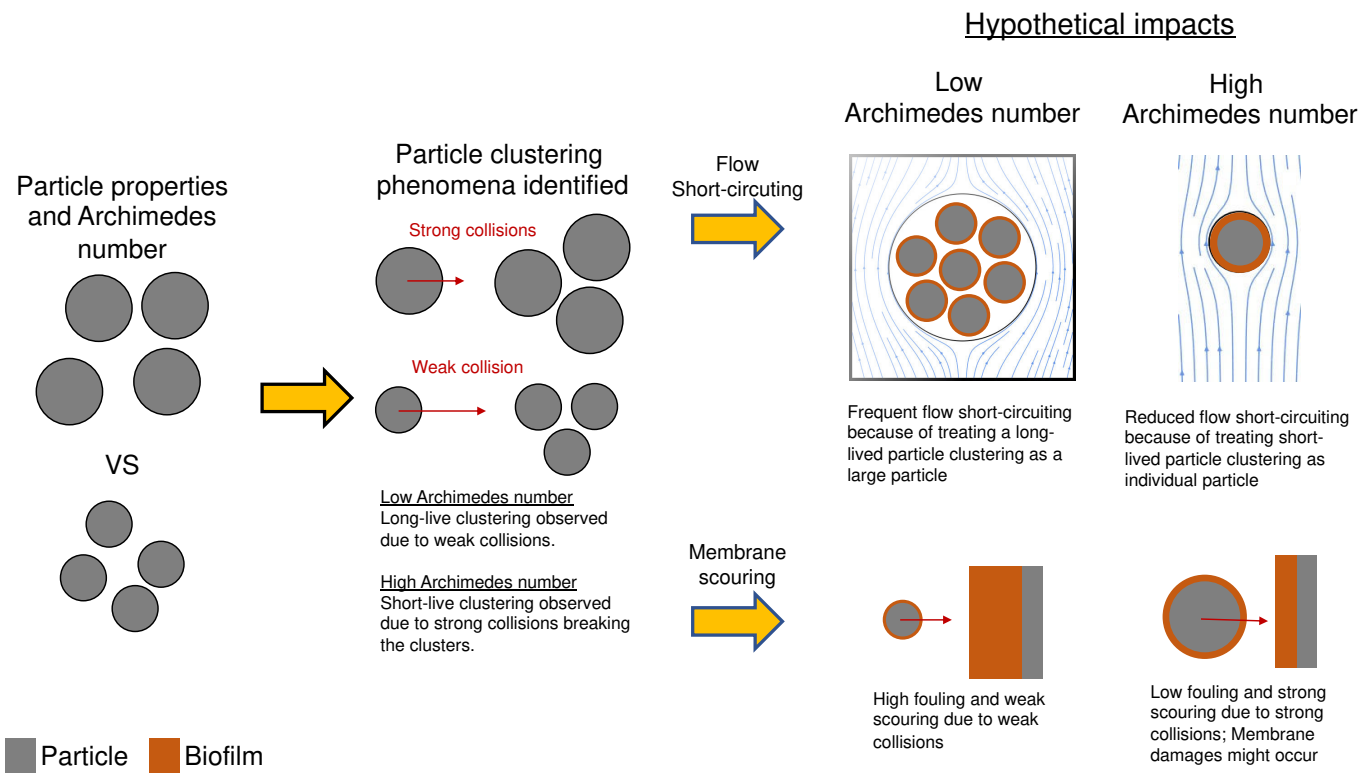
In addition to operating parameters such as up-flow velocity and porosity, choosing optimal or appropriate design parameters (i.e. Archimedes number) is critical. As discussed in the previous section, particles with low Archimedes number are preferred for better mass transfer in the AFBR. In reality, particles with the same prop-

erties are usually used for the P-MBR. Aslam et al.<sup>26</sup> attempted to relate particle properties such as materials, diameter and density to membrane scouring efficiency and concluded that larger particles are better at membrane fouling control. Recently, Yao et al.<sup>33</sup> elucidated the role of the Archimedes number on particle dynamics in a fluidized bed. The Archimedes number combines different particle properties into a single dimensionless number. Based on the simulations, the normalized particle velocity fluctuation decreases as the Archimedes number increases, indicating that the particles experience weaker effects of wake interactions in which the particle is weakly affected by neighbouring particles. By using Voronoï tessellation, particle clustering is identified and the results suggest that Archimedes number has a strong inverse relationship on particle clustering lifespan such that an increase in Archimedes number strongly decreases the lifespan. Therefore, applications with low Archimedes number are characterized by long-lived clusters while applications with high Archimedes number are characterized by short-lived clusters. The mechanism governing the lifespan of particle clusters is the collision frequency. Increasing the Archimedes number increases the collision frequency, creating conditions more favorable for cluster breakup, leading to short-lived clusters.

#### 3.2.1 Hypothetical impacts on flow short-circuiting

A common practice in the operation of fluidized bed reactors in wastewater treatment is to use small particles that enhance both mass transfer and surface contact. In analogy to boundary layer thickness, the diffusion layer thickness scales as

$$L \sim \sqrt{\frac{Dd_p}{\bar{u}}}, \quad (6)$$



**Fig. 3** Hypothetical impacts of particle properties and Archimedes number on fluidized-bed reactor modeling, design and operation.

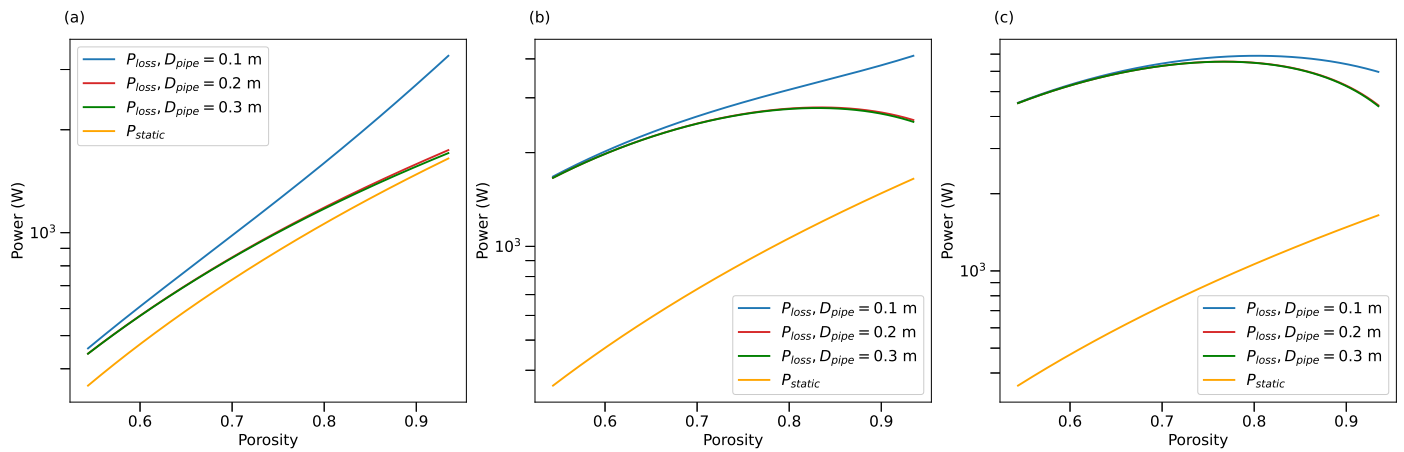
507 where  $D$  is the diffusion coefficient of the tar-  
 508 geted compound and  $\tilde{u}$  is the fluid velocity over  
 509 the particle. From equation 6,  $L$  decreases as  
 510  $d_p$  decreases and  $\tilde{u}$  increases. Therefore, smaller  
 511 particles are less likely to be mass-transfer lim-  
 512 ited due to the reduced diffusion layer thick-  
 513 ness. In practice, particle size is chosen based on  
 514 the minimum particle size or Archimedes number  
 515 that can be easily retained in the system. How-  
 516 ever, contrary to popular opinion, Yao et al.<sup>33</sup>  
 517 found that particles with  $Ar < 1000$  tend to form  
 518 prolonged clusters while particles with  $Ar \geq 1000$   
 519 are more likely to form short-lived clusters. Long-  
 520 lived clusters tend to behave like a single large  
 521 particle, resulting in more fluctuations in particle  
 522 dynamics. This result suggests that flow short-  
 523 circuiting is more likely to occur when fluid flows  
 524 over a large particle cluster rather than each  
 525 individual particle, resulting in reduced surface  
 526 contact and exchange between biofilm and bulk

fluid.

527

### 3.2.2 Hypothetical impacts on scouring fre-528 quency and membrane lifetime 529

Besides porosity, particle properties controlling 530  
 the Archimedes number can affect membrane 531  
 scouring efficiency. As demonstrated in many 532  
 papers<sup>26,50,54</sup>, larger particles (high Archimedes 533  
 number) tend to result in more frequent impact 534  
 collisions that ultimately damage the membrane 535  
 over time while small particles (low Archimedes 536  
 number) do not induce effective collisions and 537  
 hence minimal membrane scouring<sup>33</sup>. Therefore, 538  
 choosing particles with appropriate Archimedes 539  
 number is critical. To ensure effective mem- 540  
 brane biofouling control, particles with  $Ar > 1000$  541  
 are likely to have effective collisions<sup>33</sup>. There- 542  
 fore, the minimum Archimedes number for ef- 543  
 fective membrane scouring is  $Ar \approx 1000$ . For 544  
 better scouring efficiency, particles with higher 545



**Fig. 4** Power requirement as a function of porosity for (a)  $Ar = 2.3 \times 10^4$ , (b)  $Ar = 1.2 \times 10^5$  and (c)  $Ar = 2.3 \times 10^5$ .

546 Archimedes number will be more effective when  
 547 there is a risk of membrane damage.

548 To alleviate membrane damage due to parti-  
 549 cle scouring, both the frequency and energy im-  
 550 pacts of collisions must be reduced. This can be  
 551 achieved by operating the P-MBR at a porosity  
 552 that favors weaker collisions. Collision frequency  
 553 and strength can both be reduced by changing  
 554 porosity (both by increasing and decreasing it).  
 555 Because the expanded fluidized particles must be  
 556 able to access to the entire membrane module in  
 557 order to provide effective scouring, and because  
 558 the fluidized-bed height is predetermined, the to-  
 559 tal mass of particles must change if the upflow ve-  
 560 locity changes. The disadvantage of adding more  
 561 particles (reducing porosity) is that this leads to  
 562 higher headloss and increased pumping costs. To  
 563 increase porosity, a higher flow rate, hence a  
 564 higher power requirement, is essential. Since the  
 565 total headloss is proportional to both the hydro-  
 566 static pressure loss and pipe friction loss, a more  
 567 detailed analysis of power requirements is re-  
 568 quired to determine the optimal flow rate. As an  
 569 example, figure 4 shows the power requirement  
 570 as a function of porosity for different Archimedes  
 571 numbers  $Ar$  (model details can be found in the  
 572 Supplementary Information). As shown, when  
 573 the recirculation pipe diameter  $D_{pipe} > 0.2 \text{ m}$ ,

$D_{pipe}$  is no longer an important parameter. For  
 574 low  $Ar$ , the power requirement is dominated by  
 575 the static head loss, and the wastewater must be  
 576 pumped from the bottom to the top of the reac-  
 577 tor, leading to a monotonically increasing func-  
 578 tion of  $Ar$ . For high  $Ar$ , the power needed to flu-  
 579 idize the particles exceeds static headloss, lead-  
 580 ing to a parabolic function of  $Ar$ . Therefore, to  
 581 reduce high energy collisions, the flow rate must  
 582 be reduced for small  $Ar$  and can be increased or  
 583 decreased for high  $Ar$  depending on the power  
 584 requirements. 585

#### 4 Conclusion and outlook 586

Simulations of particle dynamics in fluidized-bed  
 587 reactors using CFD suggest that the parameter  
 588 space for optimal bed expansion should decrease  
 589 from 10%-70% to 40%-60% because optimal  
 590 mass transfer is more likely to occur when both  
 591 collisional and hydrodynamic forces are compa-  
 592 rably important. To design an efficient fluidized-  
 593 bed reactor, particles with  $Ar > 1000$  should be  
 594 chosen to avoid flow short-circuiting due to parti-  
 595 cle clustering. Similarly, particles with  $Ar > 1000$   
 596 or preferably  $Ar > 7000$  are needed to induce  
 597 appreciable membrane scouring. The impact of  
 598 membrane scouring can be adjusted by varying  
 599 the porosity or flow rate. 600

Overall, high-fidelity CFD simulations enable a 601

602 close examination of fundamental hydrodynam-  
 603 ics within bioreactors. Although optimal de-  
 604 sign and operating conditions cannot be precisely  
 605 identified, the range of parameter space requir-  
 606 ing experimental testing can be significantly re-  
 607 duced, and the likelihood that optimal conditions  
 608 will be identified is greater. CFD simulations pro-  
 609 vide an added tool for study of problems that  
 610 are difficult to investigate experimentally. Exper-  
 611 iments can both validate and build upon CFD re-  
 612 sults to optimize reactor performance.

613 Although CFD-accelerated strategies have  
 614 tremendous potential for acceleration and opti-  
 615 mization of wastewater treatment systems, more  
 616 work is clearly needed. More sophisticated com-  
 617 putational methods are needed that incorporate  
 618 biological reactions. However, the main chal-  
 619 lenge in integrating biological reactions is the dif-  
 620 ference in timescales. For biological reactions,  
 621 the timescales are typically much longer than the  
 622 time to reach hydrodynamic steady-state. As a re-  
 623 sult, the total computational cost increases signif-  
 624 icantly, and new methods are needed to address  
 625 this challenge.

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## Conflicts of interest 648

There are no conflicts to declare. 649

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