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Lianhuan Han *et al.*
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Regulation of heterogeneous electron transfer reactivity by defect engineering through electrochemically induced brominating addition†

Lanping Zeng,^{ac} Lianhuan Han,^{id}*^b Wenjing Nan,^a Weiying Song,^a Shiyi Luo,^a Yuan-Fei Wu,^a Jian-Jia Su^a and Dongping Zhan^{id}^a

Enhancing the electrochemical activity of graphene holds great significance for expanding its applications in various electrochemistry fields. In this study, we have demonstrated a facile and quantitative approach for modulating the defect density of single-layer graphene (SLG) via an electrochemically induced bromination process facilitated by cyclic voltammetry. This controlled defect engineering directly impacts the heterogeneous electron transfer (HET) rate of SLG. By utilizing Raman spectroscopy and scanning electrochemical microscopy (SECM), we have established a correlation between the HET kinetics and both the defect density (n_D) and mean distance between defects (L_D) of SLG. The variation of the HET rate (k^0) with the defect density manifested a distinctive three-stage behavior. Initially, k^0 increased slightly with the increasing n_D , and then it experienced a rapid increase as n_D further increased. However, once the defect density surpassed a critical value of about $1.8 \times 10^{12} \text{ cm}^{-2}$ ($L_D < 4.2 \text{ nm}$), k^0 decreased rapidly. Notably, the results revealed a remarkable 35-fold enhancement of k^0 under the optimal defect density conditions compared to pristine SLG. This research paves the way for controllable defect engineering as a powerful strategy to enhance the electrochemical activity of graphene, opening up new possibilities for its utilization in a wide range of electrochemical applications.

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Introduction

Benefitting from its excellent physicochemical properties, including high electron mobility,¹ superior optical transparency,² a zero-point band gap,³ and flexibility, graphene has attracted significant attention as a promising electrode material for a wide range of applications, such as electrochemical sensors,⁴ electrocatalysis⁵ and energy storage devices.⁶ However, despite these remarkable advantages, the electrochemical activity of graphene still requires substantial improvement, which currently limits its application in various electrochemistry fields. The primary problem arises from its highly delocalized π -conjugated structure, rendering graphene inert in heterogeneous electron transfer (HET) reactions. Therefore,

graphene is not yet considered an ideal electrode material for these applications.

In contemporary research, it is widely acknowledged that the introduction of defects into the crystalline lattice of graphene represents an effective strategy for enhancing the HET activity of this remarkable carbon allotrope.^{7,8} Although the precise underlying mechanisms remain the subject of ongoing investigation, it is generally believed that defect engineering, denoting the precise and quantitative regulation of defects, can disrupt the highly delocalized π -conjugated system of graphene. This disruption, in turn, results in an increase in the localized electronic density of states (DOS) near the Fermi level, thereby enhancing the reactivity of graphene in catalytic processes.^{9–11} One notable avenue in defect engineering entails increasing the intrinsic defects within graphene, encompassing edge sites, vacancies, holes, and topological defects. These intrinsic defects have been recognized for their potential to increase the density of active sites, consequently elevating the HET activity of graphene.¹² To this end, various strategies have been proposed, including fabricating graphene nanoribbons by micro-nano machining^{13,14} and ball-milling¹⁵ and chemical syntheses of polycyclic aromatic hydrocarbons.^{16–18} Furthermore, the creation of vacancies or holes on the graphene lattice has been achieved through methods such as electron beam etching,¹⁹ oxidation etching,²⁰ and plasma treatment,^{21,22} while

^aState Key Laboratory of Physical Chemistry of Solid Surfaces (PCOSS), Engineering Research Center of Electrochemical Technologies of Ministry of Education, Department of Chemistry, College of Chemistry and Chemical Engineering, Xiamen University, Xiamen 361005, China

^bDepartment of Mechanical and Electrical Engineering, Pen-Tung Sah Institute of Micro-Nano Science and Technology, Xiamen University, Xiamen 361005, China. E-mail: hanlianhuan@xmu.edu.cn

^cFujian Science & Technology Innovation Laboratory for Energy Materials of China (Tan Kah Kee Innovation Laboratory), Xiamen 361005, China

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topological defects have been introduced *via* thermal treatment²³ and twisting the graphene superlattice.²⁴ Another strategy involves the incorporation of heteroatoms into graphene through processes such as thermal annealing, plasma irradiation in controlled atmospheres,^{25–28} and chemical modifications such as chemical vapor deposition (CVD),²⁹ solution-phase synthesis,^{30,31} and electrochemical methods.^{32,33} Notably, chemical modification offers the advantage of introducing stable sp^3 -defects by forming covalent bonds between functional groups and graphene. This approach exhibits high stability, good versatility, and cost-effectiveness.³⁴ However, the inherent chemical inertness of graphene poses a challenge in designing suitable chemical reaction systems to introduce stable covalent bonds. Radical reactions, particularly aryl diazonium salt functionalization,^{35,36} and cycloadditions, such as Diels–Alder additions,³⁷ have proven to be effective and widely utilized methods. We have proposed a strategy to open the band gap of single-layer graphene (SLG) by forming highly stable brominated SLG through the brominating addition reaction, which involved bromine radicals generated by electrochemical oxidation from bromide anions.³⁸

In this study, we demonstrated a simple and controllable method to regulate the HET activity of SLG by employing an electrochemical bromination process using cyclic voltammetry in an electrolyte solution containing bromide anions. By combining scanning electrochemical microscopy (SECM) and Raman spectroscopy, we were able to quantitatively investigate the influence of sp^3 defects introduced by the formation of C–Br bonds on the electrochemical activity of SLG. The electrochemical bromination process offers several advantages, including its simplicity, environmental friendliness, and low cost. These advantages make it a promising approach for widespread applications in graphene-based electrochemical devices.

Results and discussion

The inferior activity of graphene as an electrode material for the HET process has limited its application in electrochemistry. To address this issue, we propose a method to enhance graphene's electrochemical activity by grafting bromide to its basal plane through the electrochemical brominating addition reaction. This process interrupts the highly delocalized conjugated π system and thus improves its activity. Fig. 1a illustrates the experimental setup for electrochemical bromination. A CVD-grown SLG substrate with a series of square patterns ($20 \times 20 \mu\text{m}$) is employed as the working electrode. The first pattern was immersed into an electrolyte containing 10 mM KBr and 10 mM H_2SO_4 to initiate the electrochemically induced bromination process by using cyclic voltammetry (CV). Fig. 1b presents a typical CV curve of the SLG electrode in the electrolyte containing bromide anions (Br^-). As previously discussed, when the potential of SLG exceeded 0.7 V, the bromide anions (Br^-) underwent electrochemical oxidation, leading to the formation of highly reactive bromide radicals (Br^\cdot) at the SLG surface. These Br^\cdot radicals then undergo a brominating addition reaction with SLG, resulting in the formation of C–Br bonds and the transformation of the SLG's configuration from sp^2 into sp^3 . Fig. 1c provides a schematic representation elucidating the bromination mechanism. To obtain brominated SLG patterns with varying degrees of bromination on a single workpiece, the SLG substrate is incrementally moved downwards $50 \mu\text{m}$ after every 20 cycles of CV. This design allows one to avoid the influence of different fabrication and pretreatment procedures on the surface properties of SLG and to characterize the degree of bromination correlated electrochemical activity of SLG by Raman spectroscopy and SECM.

XPS experiments were performed to identify the carbon-bromine (C–Br) functionalization of the SLG surface (ESI Fig. S1†). The results demonstrate the occurrence of an



Fig. 1 (a) Schematic diagram of the electrochemically induced bromination process. (b) A typical curve of cyclic voltammograms when the SLG substrate is employed as the working electrode in an aqueous solution with 10 mM KBr and 10 mM H_2SO_4 ; the scan rate was 100 mV s^{-1} . (c) Schematics of the electrochemical bromination mechanism.



electrochemically induced brominating addition reaction and that SLG can be brominated to form brominated graphene (SLGBr). To determine the degree of bromination, Raman spectroscopy was performed to characterize the SLG samples with varying durations of CV. The results are depicted in Fig. 2a, showing good agreement with previous literature reports.^{39–41} For the sake of discussion, the SLG patterns were designated from A to K. The Raman spectrum of pristine SLG (spectrum A) represents a G peak at $\sim 1580\text{ cm}^{-1}$, a 2D peak at $\sim 2680\text{ cm}^{-1}$ and a 2D' peak at $\sim 3247\text{ cm}^{-1}$. The presence of a symmetric 2D peak with much stronger intensity than that of the G peak confirms the graphene's single-layer nature. Moreover, the absence of the D peak ($\sim 1345\text{ cm}^{-1}$) and D' peak ($\sim 1620\text{ cm}^{-1}$), which appear only in defective graphene, confirms the high quality of the graphene sample. As the bromination process proceeds, the D and D' peaks begin to emerge and vary from pattern to pattern, indicating the introduction of sp^3 -C defects into graphene due to the formation of C–Br bonds. The intensity of the D peak increases from pattern B to G, followed by an intensity decrease and spectral broadening from pattern H to K. Consistent with our previous research, the structural defects introduced by C–Br bonds are supposed to be initially separated and independent, activating pristine graphene. The D peak gets

stronger as the defects increase and then reaches a maximum (stage 1, shown in Fig. S3 in the ESI†). Subsequently, the “structurally disordered” areas induced by the coalescing defects would dominate in the SLG sheet, leading to a decrease in the D peak intensity (stage 2, Fig. S3 in the ESI†). The peak intensity ratio ($I_{\text{D}}/I_{\text{G}}$) is commonly utilized to characterize the defect density of graphene, which means the ratio of the sp^3 to sp^2 bonding character. The Raman $I_{\text{D}}/I_{\text{G}}$ mapping of the graphene samples with different degrees of bromination is shown in Fig. 2b, and the corresponding average values of $I_{\text{D}}/I_{\text{G}}$ are listed in Table 1.

The mean distance between structural defects in graphene (L_{D} , nm) can be determined from the intensity ratio $I_{\text{D}}/I_{\text{G}}$ according to the quantitative formula expressed as follows:³⁹

$$\frac{I_{\text{D}}}{I_{\text{G}}} = C_{\text{A}} \frac{r_{\text{A}}^2 - r_{\text{S}}^2}{r_{\text{A}}^2 - 2r_{\text{S}}^2} \left[\frac{e^{-\frac{\pi r_{\text{S}}^2}{L_{\text{D}}^2}}}{e^{-\frac{\pi(r_{\text{A}}^2 - r_{\text{S}}^2)}{L_{\text{D}}^2}}} \right] \quad (1)$$

where r_{S} ($\sim 1\text{ nm}$) and r_{A} ($\sim 3\text{ nm}$) are regarded as the radii of the “structurally disordered” area and the “activated” area around the structural defects induced by the C–Br bonds, respectively. C_{A} is defined using the electron–phonon matrix elements, and the value is taken as 4.2 at the green line excitation (532 nm)



Fig. 2 (a) Raman spectra (vertically aligned) and (b) Raman $I_{\text{D}}/I_{\text{G}}$ mapping of the brominated SLG patterns. (c) The probe approaching curves and (d) SECM images obtained on the brominated SLG patterns. In the SECM experiments, the tip electrode was held at a potential of 0.4 V in the electrolyte solution containing 1 mM FcMeOH and 0.1 M KCl, while the SLG substrate was kept unbiased. The tip–substrate distance was kept constant at $2.2\text{ }\mu\text{m}$ ($d/a = 0.44$).



Table 1 Summary of I_D/I_G and the corresponding L_D , n_D and k^0 of the defective graphene patterns in Fig. 2 (labelled as 1A to 1K) and Fig. S7 in the ESI (labelled as 2A to 2G)

I_D/I_G	L_D (nm)	n_D (cm ⁻²)	k^0 (cm s ⁻¹)	Label
—	—	—	1.70×10^{-3}	1A
0.13	27.36	4.25×10^{10}	2.20×10^{-3}	1B
0.15	25.96	4.73×10^{10}	2.50×10^{-3}	2A
0.24	20.4	7.65×10^{10}	3.00×10^{-3}	2B
0.25	20.2	7.80×10^{10}	4.60×10^{-3}	1C
0.28	18.89	8.92×10^{10}	6.50×10^{-3}	1D
0.51	13.85	1.66×10^{11}	2.20×10^{-2}	1E
0.72	11.53	2.39×10^{11}	3.00×10^{-2}	1F
1.34	7.95	5.03×10^{11}	3.50×10^{-2}	2C
2.33	5.31	1.13×10^{12}	5.00×10^{-2}	2D
2.85	4.24	1.77×10^{12}	6.00×10^{-2}	2E
3.08	3.58	2.48×10^{12}	4.00×10^{-2}	1G
3.07	2.96	3.64×10^{12}	2.50×10^{-2}	2F
3.03	2.82	4.00×10^{12}	1.40×10^{-2}	2G
2.94	2.67	4.47×10^{12}	1.00×10^{-2}	1H
2.79	2.48	5.18×10^{12}	3.00×10^{-3}	1I
2.65	2.34	5.81×10^{12}	2.50×10^{-3}	1J
2.34	2.11	7.15×10^{12}	2.20×10^{-3}	1K

which is explicated in ESI Section S2 and Fig. S2.† The defect density n_D (cm⁻²) is then given by the formula⁴¹

$$n_D = \frac{10^{14}}{\pi L_D^2} \quad (2)$$

As shown in Fig. S2,† a specific value of I_D/I_G can correspond to two potential L_D values. To accurately determine the true L_D value, the two-stage spectral changes presented in Fig. 2a are taken into account. During stage 2, the width of the D, G, D' and 2D peaks exhibits significant broadening (Fig. S3†), which helps in discriminating the correct L_D value associated with the particular I_D/I_G ratio. Eventually, both L_D and n_D values are provided in Table 1, taking into consideration the refined L_D values obtained through the two-stage spectral analysis and using the formula (eqn (1) and (2)) mentioned earlier for calculating the defect density (n_D) based on the updated L_D value.

SECM was employed to investigate the electrochemical activity of the SLG patterns with different durations of bromination. A Pt UME with a radius of 5 μm (r) and an RG of 2 was used as the tip electrode, while the brominated SLG was used as the substrate. FcMeOH was used as the redox couple because the outer-sphere reaction involves no specific surface interaction between FcMeOH and graphene. During the SECM measurements, the tip electrode was held at a potential of 0.45 V, while the substrate remained unbiased. The oxidized species (FcMeOH⁺) at the tip electrode underwent reduction to the reduced species (FcMeOH) on the SLG substrate. Subsequently, the reduced species diffused to the tip electrode and were further oxidized, forming positive feedback (see ESI Section S3 and Fig. S4† for the scheme of SECM). The feedback current of the tip is dependent on the kinetics of HET between SLG and the redox (FcMeOH); the higher the feedback current, the faster the local HET rate. The probe approaching curves and SECM imaging of the feedback current on the brominated SLG

are presented in Fig. 2c and d, respectively, corresponding to Fig. 2b depicting the Raman I_D/I_G mapping. It is evident that as the bromination proceeds, the tip current first increases (from pattern A to G) and then decreases (pattern H to K), mirroring the variation in I_D/I_G . This observation indicates a clear correlation between the electrochemical activities of graphene and its defect density.

To quantitatively investigate the HET rate, the probe approaching curves were recorded. FEM modeling was employed to fit the experimental approaching curves as a function of the HET rate constant (k^0) through COMSOL Multiphysics 5.5. The detailed description of the FEM modeling, including solving equations, boundary conditions, initial conditions and necessary parameters, can be found in ESI Section S4.† As depicted in Fig. 2c, the experimental approaching curves (solid lines) exhibited good agreement with the simulated results (symbols), indicating that the parameters used in the FEM modeling were reasonable and appropriate for the system under investigation. The rate constant (k^0) of HET on pristine SLG was calculated to be 1.7×10^{-3} cm s⁻¹ (sample A in Fig. 2b), which is lower than that reported in previous studies.^{21,42–44} The difference may be attributed to the variation in the processing technique of pristine CVD graphene. A perfect graphene is expected to be uniform, smooth and non-defective, thus exhibiting inert behavior in HET reactions due to its highly delocalized conjugated π system. Therefore, the lower value of k^0 implies the high quality of pristine graphene. Another set of SLG patterns with different degrees of bromination was characterized by Raman spectroscopy and SECM, and the results are provided in ESI Section S5 and Fig. S7,† which also show a good correlation between them. All the obtained kinetic results (k^0) are summarized in Table 1 for further systematic analysis.

Fig. 3a shows that the I_D/I_G value of the brominated SLG initially increased with the continuation of the bromination process and then decreased. The rate constant (k^0) of HET on the brominated SLG exhibited a similar trend to the I_D/I_G ratio. It is essential to note that the horizontal ordinate in Fig. 3a represents the relative degree of bromination and doesn't strictly represent a proportional relationship between the results. Furthermore, it can be observed that k^0 increased slightly with I_D/I_G at the very beginning of the bromination process and then increased to maximum rapidly, shown in Fig. 3b. After that, k^0 decreased rapidly with the decrease in I_D/I_G and then kept stable relatively.

The correlation between k^0 and the defect parameters, including mean distance between defects (L_D) and defect density (n_D), is shown in Fig. 3c and d, revealing a three-region behavior. In the first region (light green in Fig. 3c), k^0 increased slightly from 1.70×10^{-3} cm s⁻¹ to 6.50×10^{-3} cm s⁻¹ as L_D decreased to 18.89 nm. Simultaneously, n_D increased to 8.92×10^{10} cm⁻². In this variation region, the defect density of SLG is relatively low, and its effect on the electrochemical activity of SLG is limited. As more defects were introduced, a significant enhancement of k^0 from 2.20×10^{-2} cm s⁻¹ to 6.00×10^{-2} cm s⁻¹ was observed in a relatively high defect density range from 1.66×10^{11} cm⁻² to 1.77×10^{12} cm⁻², corresponding to L_D values from ~14.00 to ~4.00 nm, respectively. With the defect





Fig. 3 (a) The I_D/I_G (black line) and the standard HET rate constant k^0 (red line) varying with the continuation of the bromination process. (b) The correlation between I_D/I_G and k^0 . (c and d) k^0 as a function of defect density n_D (cm $^{-2}$) and the mean distance between defects L_D (nm), respectively.

density further increasing, the k^0 decreased dramatically from 6.00×10^{-2} cm s $^{-1}$ to 2.20×10^{-3} cm s $^{-1}$ at a defect density of 7.15×10^{12} cm $^{-2}$ ($L_D = 2.11$ nm). The enhanced electrochemical activity can be attributed to an increase in the density of states (DOS) near the Dirac point in SLG introduced by the sp^3 C–Br bond, as explained in detail in our previous work.²¹ The introduction of defects into graphene could modulate the electronic structure of graphene and consequently enhance the HET rate.

The obtained results demonstrate a clear correlation between defect density introduced by bromination and electrochemical activity of graphene. This implies that precise control of the defect density can lead to the optimization of the electrochemical performance of graphene. Interestingly, our findings follow a similar trend of defect density dependence as observed in previous studies where point vacancy defects were introduced by Ar^+ irradiation. The differences in the absolute values of the defect density (L_D , n_D) and k^0 could be derived from the differences in surface properties of the graphene sample, because the chemical and physical properties of the graphene surface are sensitive to the fabrication method and pretreatment procedures. Nonetheless, compared to plasma irradiation techniques, the electrochemical bromination process offers a simpler and gentler approach for tailoring the defect density and consequently enhancing the electrochemical performance of graphene, making it more suitable for various electrochemical applications.

Experimental

Chemicals and materials

All chemicals used in the experiments are of analytical grade, purchased from Sinopharm Group Co. Ltd (H_2SO_4 , KBr, KCl,

acetone) and Alfa Aesar (hydroxymethyl ferrocene (FcMeOH)). A photoresist (S1805) was purchased from Merck Group. All aqueous solutions were prepared with deionized water (18.2 M Ω cm, Milli-Q, Millipore Co.).

Electrochemical bromination of SLG

SLG was prepared by chemical vapor deposition (CVD, First-Nano, ET2000) on pieces of copper foil and transferred to the SiO_2/Si substrate subsequently. A 500-nm-thick photoresist film (S1805) was spin-coated onto SLG and then patterned by using a laser direct writing photolithography system (Microtech LW405D, Italy) following the drying process. As a result, the exposed SLG patterns with a series of square arrays (20 $\mu\text{m} \times 20 \mu\text{m}$) and a period of 1000 μm were obtained after exposure and development procedures. To serve as current terminals, a 10-nm-thick chrome (Cr) film and a 40-nm-thick gold (Au) film were sequentially deposited at one end of the SLG sheet sequentially by photolithography and magnetron sputtering, as depicted in Fig. 1a. For defect engineering purposes, the patterned SLG substrate, acting as the working electrode, was immersed in an aqueous electrolyte containing 10 mM KBr and 10 mM H_2SO_4 to perform electrochemically induced bromination by cyclic voltammetry. The potential ranged from -0.5 V to 1.4 V (*vs.* MSE) with a scan rate of 100 mV s $^{-1}$. According to our previous study, brominated SLG would be formed due to the occurrence of an electrochemically induced brominating addition reaction, leading to an increase in defects on SLG.

To mitigate errors arising from variations in pretreatment procedures and to enable reliable characterization of electrochemical activity correlated to defect density through imaging



measurement, a series of SLG patterns with varying bromination degrees were prepared. Specifically, the initial pattern located at the end of the SLG substrate, far away from the current terminal, was immersed into the aqueous electrolyte. Subsequently, the substrate was incrementally moved downwards a specific distance after a certain number of CV cycles, allowing the subsequent square pattern to be exposed to the electrolyte for bromination. This process was repeated until the final pattern was brominated. Finally, the SLG patterns with different well-defined defect densities on a single piece were obtained.

Raman characterization of graphene defects

Raman measurements were performed to characterize the defects of SLG patterns on an Xplora (Horiba Jobin Yvon, France) with 532 nm laser excitation. The Raman spectra were recorded with 1200 lines per mm grating, while the laser power was kept at 5 mW with an acquisition time of 2 s.

SECM characterization

All of the electrochemical measurements were performed on a CHI 920c workstation (CH Instruments Inc.). A 5- μm radius (r) Pt UME ($\text{RG} = 2$) was used as the SECM tip, and the as-prepared SLG substrate was used as the substrate electrode. A Pt wire and an Ag/AgCl electrode were used as the counter electrode (CE) and reference electrode (RE), respectively. The Pt UME was polished using nanoscale polishing cloth and ultrasonically cleaned in ultrapure water (Milli-Q) before experiments. The electrolyte was a solution containing 1 mM hydroxymethylferrocene (FcMeOH) and 0.1 M KCl. The probe approaching curves were recorded in feedback mode, where the tip electrode was held at 0.45 V, and the SLG substrate was unbiased. The tip–substrate distance was determined using the positive approach curves obtained with the substrate biased at 0 V (purely diffusion-controlled conditions). The SECM imaging was performed with the unbiased SLG substrate, and the tip–substrate distance was kept at 2.2 μm ($L = d/a = 0.44$).

Finite element method (FEM) modeling

The heterogeneous electron transfer (HET) kinetics of SLG with different defect densities were derived by the finite element method (FEM) through a COMSOL Multiphysics 5.5 (COMSOL AB, Sweden) software package. The details including mass transport equations, the 2D symmetric geometry, and the boundary conditions are provided in ESI Section S4.†

Conclusions

In conclusion, we have successfully developed a straightforward and controllable chemical modification method to optimize the electrochemical activity of graphene. By controllably introducing different sp^3 -defect densities through electrochemical bromination, we were able to enhance the electrochemical activity of SLG. Raman spectroscopy and scanning electrochemical microscopy were utilized to investigate the correlation between the HET rate and sp^3 -defect density of SLG. The results

indicate that the electrochemical activity of SLG is improved at a moderate defect density and then decreases at high defect density. The research demonstrates that the electrochemical activity of SLG can be precisely tuned by introducing different sp^3 -defect densities through electrochemical bromination under ambient conditions. This approach offers a controllable, effective, and environmentally friendly method to optimize the electrochemical activity of graphene, making it more suitable for a wide range of electrochemical applications.

Data availability

The authors declare that all data supporting the findings of this study are available from the corresponding author upon reasonable request.

Author contributions

L. Zeng, L. Han, W. Nan, and J. Su: electrochemical experiments and instrumental characterization; L. Zeng, W. Song, S. Luo, and Y. Wu: microfabrication and electronic characterization; L. Zeng and L. Han: theoretical discussions; D. Zhan: originated, designed and organized the work and wrote the manuscript; all the authors participated in the discussions and preparation of the manuscript.

Conflicts of interest

There are no conflicts to declare.

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Notes and references

- 1 S. V. Morozov, K. S. Novoselov, M. I. Katsnelson, F. Schedin, D. C. Elias, J. A. Jaszczak and A. K. Geim, *Phys. Rev. Lett.*, 2008, **100**, 016602.
- 2 R. R. Nair, P. Blake, A. N. Grigorenko, K. S. Novoselov, T. J. Booth, T. Stauber, N. M. R. Peres and A. K. Geim, *Science*, 2008, **320**, 1308.
- 3 A. K. Geim, *Science*, 2009, **324**, 1530–1534.
- 4 Y. Xiao, Y. Su, X. Liu and W. Xu, *J. Phys. Chem. Lett.*, 2019, **10**, 5402–5407.
- 5 M. Pumera, *Chem. Soc. Rev.*, 2010, **39**, 4146–4157.
- 6 M. Vaqueiro-Contreras, C. Bartlam, R. S. Bonilla, V. P. Markevich, M. P. Halsall, A. Vijayaraghavan and A. R. Peaker, *Sol. Energy Mater. Sol. Cells*, 2018, **187**, 189–193.
- 7 Y. Han, X. Yan, Q. Wu, H. Xu, Q. Li, A. Du and X. Yao, *Small Struct.*, 2023, **4**, 2300036.
- 8 J. Zhu and S. Mu, *Adv. Funct. Mater.*, 2020, **30**, 2001097.



- 9 E. Santos, R. Nazmutdinov and W. Schmickler, *Curr. Opin. Electrochem.*, 2020, **19**, 106–112.
- 10 A. K. Verma, A. M. Verma and A. Govind Rajan, *Curr. Opin. Electrochem.*, 2022, **36**, 101116.
- 11 X. Yan, Y. Jia and X. Yao, *Chem. Soc. Rev.*, 2018, **47**, 7628–7658.
- 12 Y. Jia, J. Chen and X. Yao, *Mater. Chem. Front.*, 2018, **2**, 1250–1268.
- 13 M. López-Suárez, F. Torres, N. Mestres, R. Rurali and G. Abadal, *Microelectron. Eng.*, 2014, **129**, 81–85.
- 14 M. F. El-Kady, V. Strong, S. Dubin and R. B. Kaner, *Science*, 2012, **335**, 1326–1330.
- 15 C. K. Chua, Z. Sofer, B. Khezri, R. D. Webster and M. Pumera, *Phys. Chem. Chem. Phys.*, 2016, **18**, 17875–17880.
- 16 X. L. Li, X. R. Wang, L. Zhang, S. W. Lee and H. J. Dai, *Science*, 2008, **319**, 1229–1232.
- 17 Y. B. Fu, H. Yang, Y. X. Gao, L. Huang, R. Berger, J. Z. Liu, H. L. Lu, Z. H. Cheng, S. X. Du, H. J. Gao and X. L. Feng, *Angew. Chem., Int. Ed.*, 2020, **59**, 8873–8879.
- 18 A. Narita, X.-Y. Wang, X. Feng and K. Müllen, *Chem. Soc. Rev.*, 2015, **44**, 6616–6643.
- 19 A. W. Robertson, C. S. Allen, Y. A. Wu, K. He, J. Olivier, J. Neethling, A. I. Kirkland and J. H. Warner, *Nat. Commun.*, 2012, **3**, 1144.
- 20 X. R. Wang and H. J. Dai, *Nat. Chem.*, 2010, **2**, 661–665.
- 21 J.-H. Zhong, J. Zhang, X. Jin, J.-Y. Liu, Q. Li, M.-H. Li, W. Cai, D.-Y. Wu, D. Zhan and B. Ren, *J. Am. Chem. Soc.*, 2014, **136**, 16609–16617.
- 22 L. Tao, Q. Wang, S. Dou, Z. Ma, J. Huo, S. Wang and L. Dai, *Chem. Commun.*, 2016, **52**, 2764–2767.
- 23 Q. Wang, Y. Lei, D. Wang and Y. Li, *Energy Environ. Sci.*, 2019, **12**, 1730–1750.
- 24 E. Y. Andrei and A. H. MacDonald, *Nat. Mater.*, 2020, **19**, 1265–1275.
- 25 D.-Y. Yeom, W. Jeon, N. D. K. Tu, S. Y. Yeo, S.-S. Lee, B. J. Sung, H. Chang, J. A. Lim and H. Kim, *Sci. Rep.*, 2015, **5**, 9817.
- 26 M. Rybin, A. Pereyaslavtsev, T. Vasilieva, V. Myasnikov, I. Sokolov, A. Pavlova, E. Obraztsova, A. Khomich, V. Ralchenko and E. Obraztsova, *Carbon*, 2016, **96**, 196–202.
- 27 Z. H. Sheng, H. L. Gao, W. J. Bao, F. B. Wang and X. H. Xia, *J. Mater. Chem.*, 2012, **22**, 390–395.
- 28 Y. Wang, Y. Y. Shao, D. W. Matson, J. H. Li and Y. H. Lin, *ACS Nano*, 2010, **4**, 1790–1798.
- 29 L. Chen, J. Han, Y. Ito, T. Fujita, G. Huang, K. Hu, A. Hirata, K. Watanabe and M. Chen, *Angew. Chem., Int. Ed.*, 2018, **57**, 13302–13307.
- 30 Q. Wang, Y. Ji, Y. Lei, Y. Wang, Y. Wang, Y. Li and S. Wang, *ACS Energy Lett.*, 2018, **3**, 1183–1191.
- 31 S. Boopathi, T. N. Narayanan and S. S. Kumar, *Nanoscale*, 2014, **6**, 10140–10146.
- 32 F. Zhou, H. Huang, C. Xiao, S. Zheng, X. Shi, J. Qin, Q. Fu, X. Bao, X. Feng, K. Müllen and Z.-S. Wu, *J. Am. Chem. Soc.*, 2018, **140**, 8198–8205.
- 33 W. Li, Y. Li and K. Xu, *Nano Lett.*, 2021, **21**, 1150–1155.
- 34 C. Wetzl, A. Silvestri, M. Garrido, H.-L. Hou, A. Criado and M. Prato, *Angew. Chem., Int. Ed.*, 2023, **62**, e202212857.
- 35 K. C. Knirsch, R. A. Schäfer, F. Hauke and A. Hirsch, *Angew. Chem., Int. Ed.*, 2016, **55**, 5861–5864.
- 36 Y. Xia, C. Martin, J. Seibel, S. Eyley, W. Thielemans, M. van der Auweraer, K. S. Mali and S. De Feyter, *Nanoscale*, 2020, **12**, 11916–11926.
- 37 J. Li, M. Li, L.-L. Zhou, S.-Y. Lang, H.-Y. Lu, D. Wang, C.-F. Chen and L.-J. Wan, *J. Am. Chem. Soc.*, 2016, **138**, 7448–7451.
- 38 L. Zeng, W. Song, X. Jin, Q. He, L. Han, Y.-f. Wu, C. Lagrost, Y. Leroux, P. Hapiot, Y. Cao, J. Cheng and D. Zhan, *Chem. Sci.*, 2023, **14**, 4500–4505.
- 39 M. M. Lucchese, F. Stavale, E. H. M. Ferreira, C. Vilani, M. V. O. Moutinho, R. B. Capaz, C. A. Achete and A. Jorio, *Carbon*, 2010, **48**, 1592–1597.
- 40 A. Eckmann, A. Felten, I. Verzhbitskiy, R. Davey and C. Casiraghi, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2013, **88**, 035426.
- 41 L. G. Cancado, A. Jorio, E. H. M. Ferreira, F. Stavale, C. A. Achete, R. B. Capaz, M. V. O. Moutinho, A. Lombardo, T. S. Kulmala and A. C. Ferrari, *Nano Lett.*, 2011, **11**, 3190–3196.
- 42 J. Zhang, J. Guo, D. Chen, J.-H. Zhong, J.-Y. Liu and D. Zhan, *Appl. Surf. Sci.*, 2019, **491**, 553–559.
- 43 W. Li, C. Tan, M. A. Lowe, H. D. Abruña and D. C. Ralph, *ACS Nano*, 2011, **5**, 2264–2270.
- 44 N. L. Ritzert, J. Rodriguez-Lopez, C. Tan and H. D. Abruña, *Langmuir*, 2013, **29**, 1683–1694.

