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# $\langle 11\bar{2}0 \rangle$ -orientation-dependent crack initiation in a titanium alloy under dwell fatigue

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The crystal orientation of titanium alloys significantly influences the nucleation of dwell fatigue cracks, which deteriorate the service life of aero-engine components. Although soft/hard orientation combinations, depending on the orientation relationship between *c*-axes and loading directions, are widely recognized as favored crack nucleation sites, the effect of  $\langle 11\bar{2}0 \rangle$  orientations of soft grains on the cracking of adjacent hard grains is also crucial, and experimental investigation is desired. Here, we report that soft grains with a "rotation angle", corresponding to the  $\langle 11\bar{2}0 \rangle$  orientation, between 15° and 30° are more prone to induce cracking in hard grains, for the reason that prismatic dislocation slips in these soft grains exert a higher stress normal to the basal plane of adjacent hard grains. In contrast, soft grains with a "rotation angle" between 0° and 15° rarely cause cracking in hard grains. A stress component parameter is proposed according to the "rotation angle", which can accurately predict the nucleation of fatigue cracks in titanium alloys. The findings shed light on the quantitative analysis of crack nucleation and the prediction of fatigue performance.

## 1. Introduction

Titanium alloys have been widely employed in aero-engine components, such as compressor discs and fan discs, owing to a combination of high specific strength, excellent corrosion resistance, and good damage resistance.<sup>1–4</sup> During service, these components are inevitably subjected to dwell cyclic loading, in which the cyclic loads are held at relatively high mean stress levels for a certain period of time, known as dwell fatigue.<sup>5</sup> Dwell fatigue has been proven to significantly deteriorate the service life of titanium alloy components and easily threaten the safe operation of aero-engines.<sup>5–7</sup> Dwell fatigue is one of the critical properties of titanium alloys, and so it is of great significance to investigate the mechanism of dwell fatigue damage for their engineering applications.

In the past few decades, the dwell fatigue of titanium alloys has attracted extensive research efforts to overcome the challenge associated with its damage mechanisms.<sup>5,8–12</sup> The dwell fatigue behavior of titanium alloys is closely related to their microstructure,<sup>9,12–15</sup> such as the crystallographic orientation

and the size and volume fraction of primary  $\alpha$  grains. In addition, the inevitable presence of soft/hard grain pairs in titanium alloys with equiaxed microstructure and bimodal microstructure is considered one of the critical factors in reducing the dwell fatigue life,<sup>16,17</sup> and it has been widely demonstrated that soft/hard grain pairs can accelerate the crack initiation.<sup>5,9,16,17</sup> Moreover, the angle between the *c*-axis of soft grains and the applied stress strongly influences the crack nucleation of adjacent hard grains under dwell fatigue.<sup>18–21</sup> Here, the soft grain is favorably oriented for slips of prismatic  $\langle a \rangle$  type dislocations,<sup>18,22</sup> while the hard grain has its *c*-axis nearly parallel to the applied stress and is poorly oriented for slips.<sup>18,22,23</sup> In recent years, some studies have suggested that the grain conducive to basal slips is also considered soft.<sup>19,24</sup> So far, the definition of soft and hard grains is still unclear.

Based on the characteristics of crack nucleation, the Stroh model for planar slips was proposed to explain the mechanism of dwell fatigue cracking.<sup>5,25,26</sup> The Stroh model suggests that extensive dislocations pile up at the soft/hard grain boundary (GB) in soft grains, inducing shear stress and the formation of slip bands in hard grains.<sup>5,25</sup> Besides, the Stroh model shows that pile-up stresses at the GB cause additional normal stress on the basal plane of the adjacent hard grain,<sup>5,25,26</sup> which, combined with applied stress, ultimately leads to cracking.<sup>5,25</sup> Subsequently, numerous experimental investigations and computational simulations verified the Stroh model.<sup>16–18,22</sup> The results also showed that the stress concentration caused by

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prismatic  $\langle a \rangle$  dislocation ( $\{10\bar{1}0\}\langle 11\bar{2}0 \rangle$ ) pile up of soft grains at the GB is high, activating a large amount of  $\langle a \rangle$  and  $\langle c + a \rangle$  dislocations in hard grains with unfavorable deformation orientation.<sup>16,17,22</sup> The above finding shows that  $\langle a \rangle$  dislocation pile up of soft grains with favorable deformation orientation can easily lead to the cracking of adjacent hard grains. However, the effect of  $\langle 11\bar{2}0 \rangle$  orientation of soft grains on the cracking of adjacent hard grains has yet to be investigated, which is crucial to the research of fatigue damage mechanisms of titanium alloys.

In this work, the quantitative effect of soft grains with different  $\langle 11\bar{2}0 \rangle$  orientations on the cracking of adjacent hard grains in an equiaxed Ti–6.5Al–3.5Mo–1.5Zr–0.3Si alloy under dwell fatigue was investigated using electron backscatter diffraction (EBSD) and transmission electron microscopy (TEM). A stress component parameter was proposed to predict the fatigue crack initiation of hard grains.

## 2. Materials and methods

The Ti–6.5Al–3.5Mo–1.5Zr–0.3Si (wt%) alloy bar was prepared through vacuum arc remelting and forging. The forged alloy was subjected to a two-step heat treatment consisting of holding at 960 °C for 1.5 hours followed by furnace cooling and 6 hours aging at 530 °C. Dwell fatigue specimens with a gauge diameter of 6 mm and a gauge length of 16 mm were cut from the alloy bar. Dwell fatigue tests were conducted on specimens by using servo-hydraulic testing machines at ambient temperature. The tests were carried out until specimen failure. A trapezoidal waveform with a loading/unloading time of 1 s, a 120 s hold at a maximum stress of  $0.91\sigma_{0.2}$  ( $\sigma_{0.2} = 918$  MPa), a 1 s hold at minimum stress and an  $R$  ratio ( $\sigma_{\min}/\sigma_{\max}$ ) of 0 were applied.

The fracture surface of a dwell fatigued specimen was observed using a KEYENCE VHX-1000C stereo-optical microscope and an FEI Apreo field emission scanning electron microscope. All samples for scanning electron microscopy (SEM) and TEM characterizations were taken from near the fracture surface, and the observation direction of these samples was perpendicular to the loading direction. SEM samples were prepared by mechanical grinding and polishing and then etched in a Kroll reagent (2 ml of HF, 6 ml of HNO<sub>3</sub>, and 92 ml of H<sub>2</sub>O). EBSD samples were prepared by grinding and then electropolished in a solution of 30 ml of perchloric acid, 175 ml of *n*-butanol, and 300 ml of methanol at –50 °C to –30 °C. The alloy's initial microstructure and internal cracks were characterized under an FEI Apreo field emission scanning electron microscope. Crystallographic orientations of cracked  $\alpha$  grains (denoted as G1 grains) and adjacent  $\alpha$  grains that may cause cracking (denoted as G2 grains) after dwell fatigue were determined using an EBSD detector equipped on the scanning electron microscope. The mean geometrically necessary dislocation (GND) density with  $\langle a \rangle$  type Burgers vectors in hard grains near soft/hard GBs was calculated from EBSD data with the same step size of 200 nm, and so the comparison of the

GND density under the same conditions was statistically significant. The use of mean GND density<sup>27</sup> within the measurement region can substantially reduce the effect of extreme GND density values<sup>28</sup> and well reflect the trend of GND density variation. The GND density is obtained by determining the dislocation density tensor based on the relationship established by Nye,<sup>29</sup> and the primary calculation approach is provided by the work of Pantleon.<sup>30</sup> Dislocation analysis in soft/hard grain pairs was conducted using a transmission Kikuchi diffraction (TKD) technique and a JEOL JEM-2100Plus transmission electron microscope. TEM samples were ground to a thickness of 50  $\mu\text{m}$  and then electropolished using a conventional twin jet with the same solution as EBSD samples at –30 °C to –20 °C.

## 3. Results and discussion

The Ti–6.5Al–3.5Mo–1.5Zr–0.3Si alloy before dwell fatigue shows an equiaxed microstructure consisting of equiaxed  $\alpha$  grains ( $\sim 85$  vol%) and transformed  $\beta$  regions ( $\sim 15$  vol%), as shown in Fig. 1(a). Transformed  $\beta$  regions are principally distributed at equiaxed  $\alpha$  GBs and triple junctions. Fig. 1(b)–(d) show the fracture surface of a dwell fatigued specimen. The fracture surface appears relatively flat overall, with multiple bright regions indicating quasi-cleavage facets (Fig. 1(b)), suggesting that the fracture initiated inside the specimen and exhibited the characteristics of multiple initiation mode. Notably, the quasi-cleavage facet is approximately perpendicular to the loading direction, as shown in Fig. 1(d).

To investigate the effect of crystallographic orientations on crack nucleation, crystallographic orientations of G1 grains with cracks approximately perpendicular to the loading direction and corresponding G2 grains were analyzed, as shown in Fig. 2 and 3. There is an adjacent  $\alpha$  grain at each end of the crack, and the  $\alpha$  grain adjacent to the end with a larger crack opening is considered as the G2 grain,<sup>11</sup> as shown in Fig. 2. In this work, a total of 32 G1 and G2 grain pairs were measured. The crystallographic orientation of the  $\alpha$  grain was characterized using  $\gamma$  and  $\phi$  angles, as defined in Fig. 3(b). Here,  $\gamma$  is the angle between the  $c$ -axis and the loading direction, and  $\phi$  is the minimum angle between the projection of the loading direction on the basal plane and the  $\langle 11\bar{2}0 \rangle$  crystal direction ( $0 \leq \phi \leq 30^\circ$ ), which corresponds to the “rotation angle”. The statistical results in Fig. 3(a) indicate that  $\gamma$  angles of these G2 grains predominantly fall within the range of  $65^\circ$  to  $90^\circ$ , and  $\phi$  angles mainly range from  $15^\circ$  to  $30^\circ$ . Furthermore, Fig. 4 shows the comparison of  $\gamma$  and  $\phi$  angle distributions of G2 grains and that of  $\alpha$  grains adjacent to uncracked hard grains, and it is evident that  $\gamma$  angles of  $\alpha$  grains adjacent to uncracked hard grains are more distributed between  $20^\circ$  and  $70^\circ$ , while the  $\phi$  angle distribution is relatively uniform between  $0$  and  $30^\circ$  when  $65^\circ \leq \gamma \leq 90^\circ$ . Since these cracked grain pairs for EBSD characterization were randomly selected from samples, the results obtained by characterizing these grain pairs are sufficiently representative. Therefore, there is an actual correlation



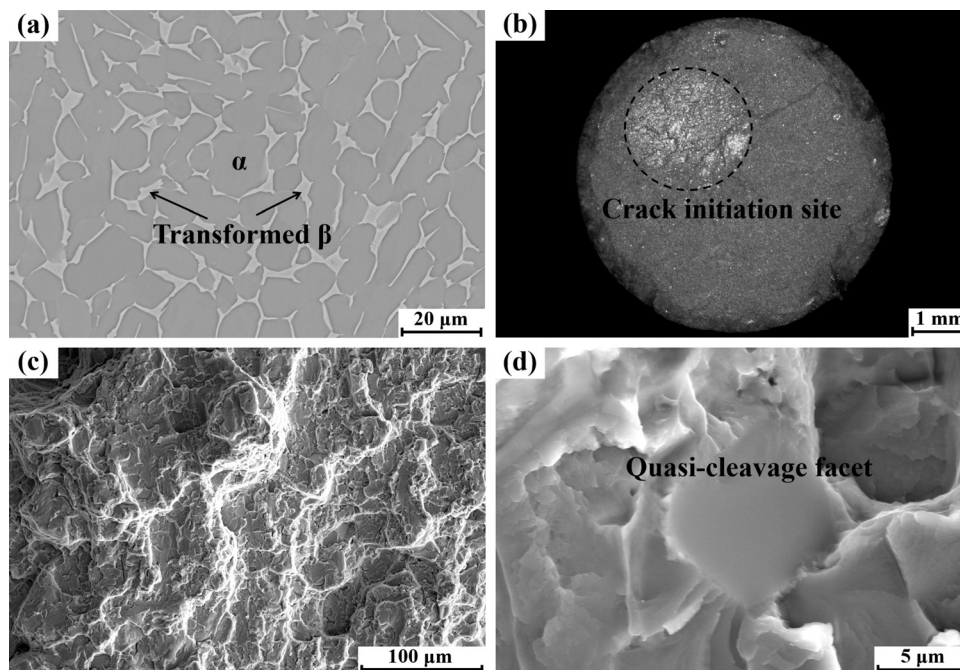


Fig. 1 (a) A SEM backscattered electron image of the initial microstructure of the titanium alloy. (b)–(d) Fractographs of a dwell fatigued specimen: (b) optical microscope image and (c) and (d) SEM secondary electron images, (c) is an amplification of the bright region in (b), and (d) shows a quasi-cleavage facet at the initiation site (bright region in (b)).

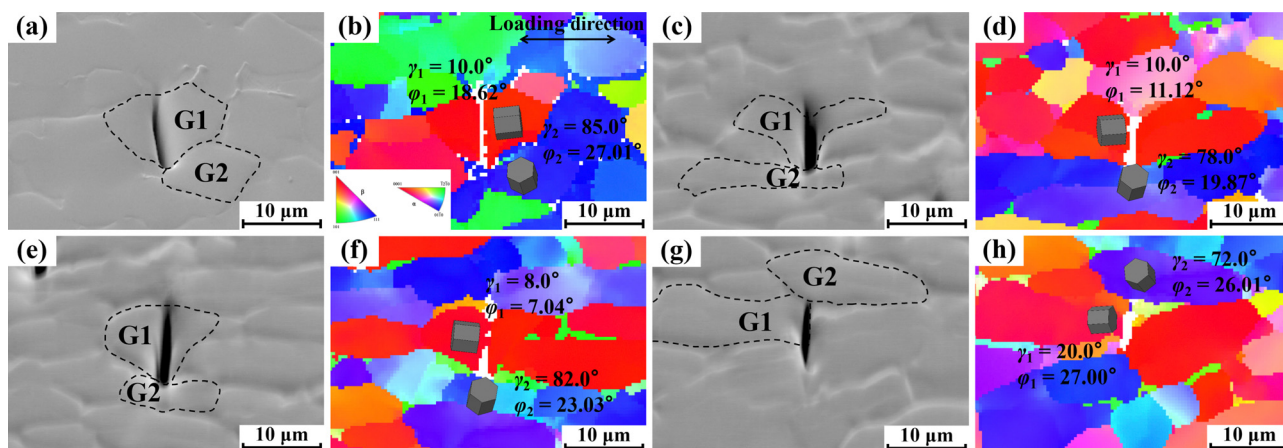


Fig. 2 (a), (c), (e) and (g) Raw SEM secondary electron images after crack initiation and the corresponding (b), (d), (f) and (h) EBSD-IPF maps. The crystallographic orientation is relative to the loading direction.

between angles ( $\gamma$  and  $\phi$  angles) and the distribution of G2 grains. In addition, Schmid factor analysis in Fig. 3(d) indicates that G2 grains are mostly well-oriented for the prismatic slip with a Schmid factor higher than 0.4. For G1 grains,  $c$ -axes of these G1 grains are approximately parallel to the loading direction, with a declination angle of  $0^\circ$ – $20^\circ$ , as shown in Fig. 3(a). Fig. 3(c) shows that basal Schmid factors of G1 grains are lower than 0.32, which are relatively poorly oriented for the basal slip. Besides, the crack surface is along the basal plane of the G1 grain, as shown in Fig. 2. Moreover, there is a blue dot in the grey one group in Fig. 3(a), which may be due to other factors (such as other  $\alpha$  grains nearby) affecting the cracking of

the grain pair.<sup>20,23</sup> In the actual measurements, there may be some incidental deviations due to sampling influences. However, statistical methods were used to investigate and obtain the distribution pattern of G2 grains' orientation ( $\gamma$  and  $\phi$  angles) based on the majority of the data points ( $>96\%$ ). Therefore, the blue dot in the grey one group has little effect on the results and discussions of the paper. To sum up, in addition to the  $\gamma$  angle, the  $\phi$  angle of the adjacent grain plays an essential role in the crack nucleation of the G1 grain.

Dislocation activities in soft/hard grain pairs with different  $\phi$  angles of soft grains were observed, as shown in Fig. 5. Crystallographic orientations of soft/hard grain pairs were



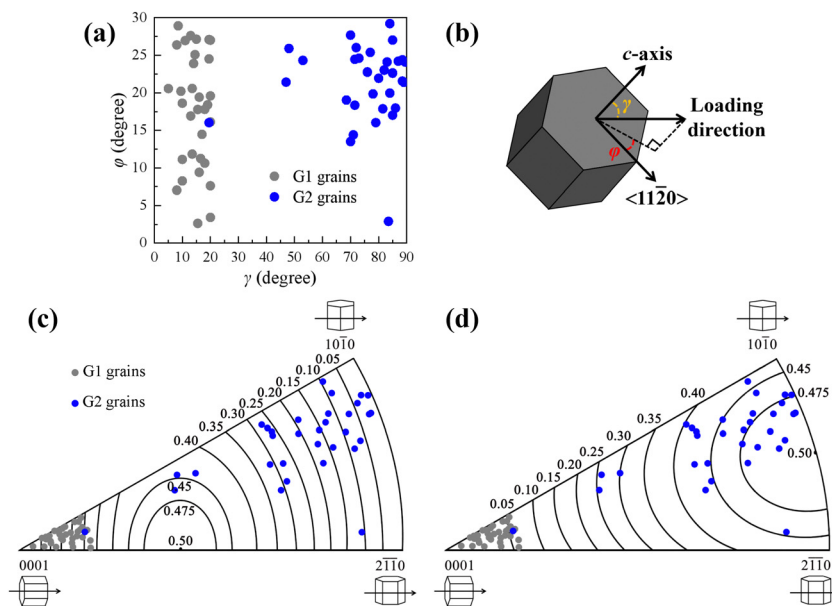


Fig. 3 (a)  $\gamma$  and  $\phi$  angle distributions of G1 and G2 grains. (b) Schematic diagram of  $\gamma$  and  $\phi$  angles. (c) and (d) Crystallographic orientations of G1 and G2 grains, plotted on the IPF map with contours of the (c) basal Schmid factor and (d) prismatic Schmid factor.

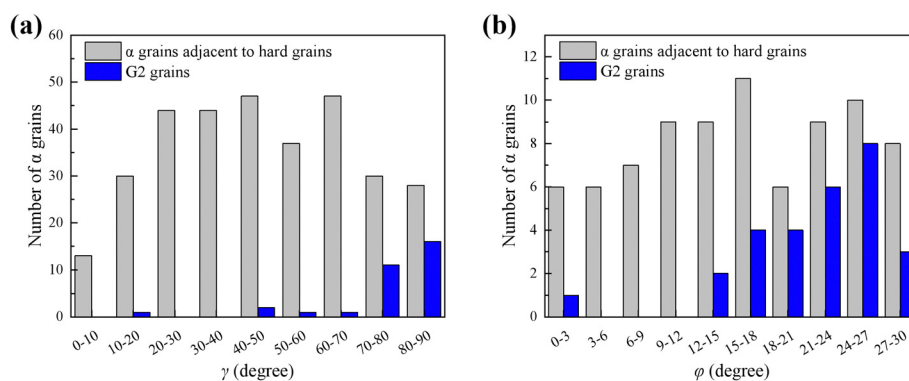


Fig. 4 (a)  $\gamma$  angle distribution of G2 grains and  $\alpha$  grains adjacent to uncracked hard grains. (b)  $\phi$  angle distribution of G2 grains and  $\alpha$  grains adjacent to uncracked hard grains when  $65^\circ \leq \gamma \leq 90^\circ$ .

determined by the TKD technique. The results show that, in both types of grain pairs ( $\phi_1 = 19.52^\circ$  and  $\phi_2 = 1.41^\circ$ ), there are a large number of prismatic slip bands in soft grains (Fig. 5(c) and (f)), and extensive dislocations pile up against soft/hard GBs in these grains, as shown in Fig. 5(b) and (f). Meanwhile, in adjacent hard grains, a large number of  $\langle a \rangle$  and  $\langle c \rangle + \langle a \rangle$  dislocations, especially the formation of basal slip bands, are observed, as shown in Fig. 5(b) and (e). The bright field TEM image in Fig. 5(b) shows that basal slip bands in the hard grain are emitted from the GB where prismatic  $\langle a \rangle$  dislocations of the soft grain pile up, indicating that the basal slip is induced by the dislocation pile up of the soft grain. By comparing dislocation activities in Fig. 5, it can be seen that the dislocation types in soft/hard grain pairs with different  $\phi$  angles of soft grains are the same. Furthermore, the mean GND density in hard grains near soft/hard GBs was measured, where  $\gamma$  angles of soft grains are between  $70^\circ$  and  $90^\circ$ , as shown in Fig. 6(b). Here, the mean standard error is used as the error bar. The result shows

that the GND density within the hard grain rises with increasing  $\phi$  angle of the soft grain when the  $\phi$  angle is between  $0$  and  $22.5^\circ$ , while the GND density drops slightly within the  $\phi$  angle between  $22.5^\circ$  and  $30^\circ$ .

According to the above experimental results, it is found that the cracking of G1 grains is related to  $\phi$  and  $\gamma$  angles of G2 grains. Thus, a stress component parameter  $\kappa$  is proposed to characterize the possibility of cracking. The  $\kappa$  parameter is derived from the stress component, which is the projection of the G2 grain's maximum prismatic shear stress along the normal direction of the hard grain's basal plane. Based on the relationship between the G2 grain orientation and the loading direction, a prismatic Schmid factor is derived:

$$m_{(\phi,\gamma)} = \cos \phi \cdot \sin \phi \cdot \sin^2 \gamma = \tau_{s(\phi,\gamma)} / \sigma \quad (1)$$

where  $\gamma$  is the angle between the  $c$ -axis of the G2 grain and the loading direction,  $\phi$  is the "rotation angle" of the G2





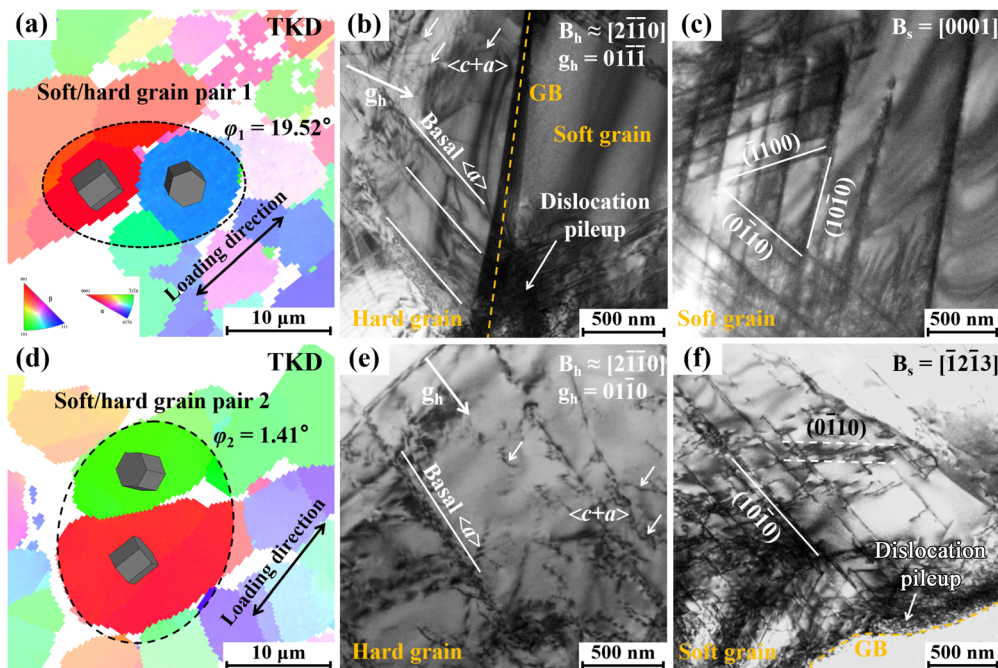


Fig. 5 (a) and (d) TKD-IPF maps of soft/hard grain pairs; crystallographic orientation is relative to the loading direction. Bright field TEM images of dislocation activities in (b) and (c) soft/hard grain pair 1 ( $\phi_1 = 19.52^\circ$ ) and (e) and (f) soft/hard grain pair 2 ( $\phi_2 = 1.41^\circ$ ). GBs are denoted by yellow dashed lines.

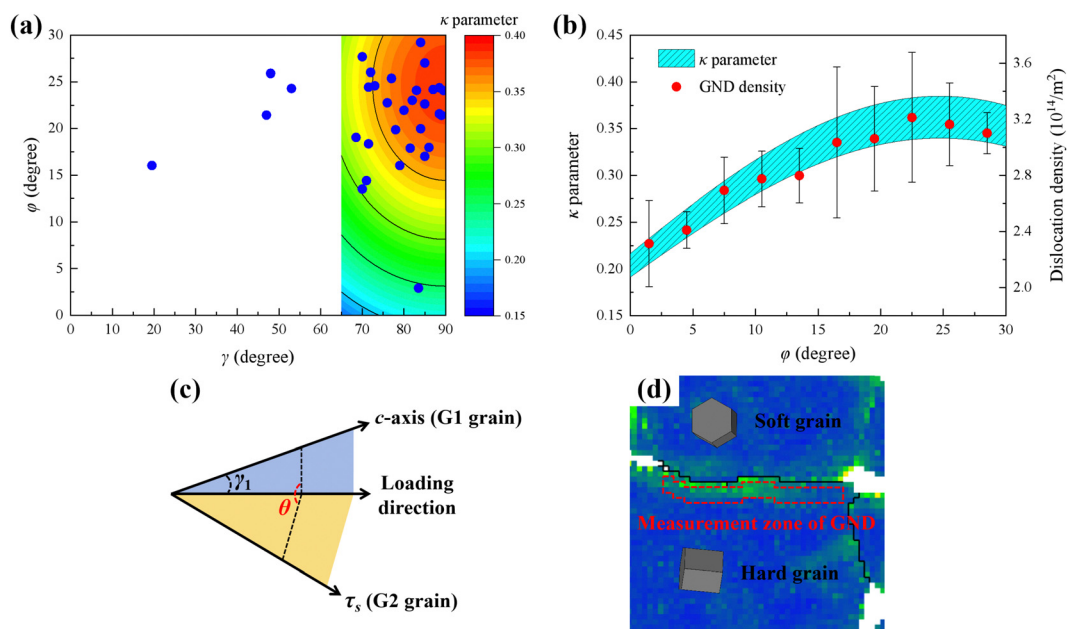


Fig. 6 (a)  $\kappa$  parameter of  $\alpha$  grains related to  $\gamma$  and  $\phi$  angles.  $\gamma$  and  $\phi$  angle distributions of G2 grains (blue solid circle) are plotted in (a). (b)  $\kappa$  parameter of  $\alpha$  grains with  $70^\circ \leq \gamma \leq 90^\circ$  and the mean GND density of hard grains near soft/hard GBs. (c) Schematic diagram of the dihedral angle  $\theta$ . (d) The measurement zone of GND. The mean standard error is used as the error bar.

grain,  $m_{(\phi,\gamma)}$  is the prismatic Schmid factor,  $\tau_{s(\phi,\gamma)}$  is the resolved shear stress, and  $\sigma$  is the applied stress. A G2 grain has three prismatic Schmid factors,  $m_{(\phi,\gamma)}$ ,  $m_{(60-\phi,\gamma)}$  and  $m_{(60+\phi,\gamma)}$ . When  $0 \leq \phi \leq 30^\circ$ , the  $m_{(60-\phi,\gamma)}$  is the maximum, and so the maximum prismatic resolved shear stress is  $\tau_{s(60-\phi,\gamma)}$ . Subsequently, the component of the G2 grain's maximum prismatic

shear stress along the normal direction of the hard grain's basal plane can be calculated using the following formula:

$$\sigma' = \tau_s [\cos \gamma_1 \cdot \cos(60 - \phi) \cdot \sin \gamma + \cos \theta \cdot \sin \gamma_1 \cdot (1 - \cos^2(60 - \phi) \cdot \sin^2 \gamma)^{1/2}] \quad (2)$$



where  $\sigma'$  is the component of prismatic shear stress of the G2 grain,  $\gamma_1$  is the angle between the  $c$ -axis of the G1 grain and the loading direction ( $0 \leq \gamma_1 \leq 20^\circ$ ),  $\theta$  is the dihedral angle between the plane composed of the resolved shear stress ( $\tau_s$ ) direction of G2 grains and the loading direction and the plane formed by the  $c$ -axis of G1 grains and the loading direction ( $0 \leq \theta \leq 180^\circ$ ), as shown in Fig. 6(c). Due to the complexity of eqn (2), an approximation is made to the part in square brackets to simplify the equation. Since  $\gamma$  angles of G2 grains predominantly fall within the range of  $65^\circ$  to  $90^\circ$  with a mean value of  $\sim 80^\circ$ , and  $\gamma_1$  angles of G1 grains are distributed between  $0$  and  $20^\circ$  with a mean value of  $\sim 15^\circ$  (Fig. 3(a)), it is reasonable to take the  $\gamma$  angle as  $80^\circ$  and the  $\gamma_1$  angle as  $15^\circ$  in the square bracket. The dihedral angle  $\theta$  is taken as the mean value  $90^\circ$ . Therefore, eqn (2) can be simplified to the following formula:

$$\sigma' = 0.95 \cdot \tau_s \cdot \cos(60 - \varphi) = 0.95 \cdot \sigma \cdot \cos^2(60 - \varphi) \cdot \sin(60 - \varphi) \cdot \sin^2 \gamma \quad (3)$$

Finally, the  $\kappa$  parameter is obtained using the following formula:

$$\kappa = \cos^2(60 - \varphi) \cdot \sin(60 - \varphi) \cdot \sin^2 \gamma \quad (4)$$

Accordingly, it can be considered that the normal stress on the G1 grain's basal plane induced by the prismatic dislocation pile up of the G2 grain at the GB corresponds to the value of the  $\kappa$  parameter. That is, the normal stress is positively correlated with the  $\kappa$  value.

Fig. 6(a) shows the contour plot of the  $\kappa$  parameter related to  $\gamma$  and  $\varphi$  angles, and different colors represent the magnitude of  $\kappa$  values. Considering that eqn (4) has been simplified, the contour plot with  $\gamma$  angles between  $65^\circ$  and  $90^\circ$  is meaningful.  $\gamma$  and  $\varphi$  angle distributions of all G2 grains (blue solid circle) obtained from the experiment are also plotted in Fig. 6(a). It is found that 87.5% of G2 grains have  $\gamma$  angles between  $65^\circ$  and  $90^\circ$ , demonstrating that the simplification of eqn (2) is reasonable. A comparison of experimental and calculated results reveals that 89.3% of G2 grains with  $\gamma$  angles between  $65^\circ$  and  $90^\circ$  have  $\kappa$  values larger than 0.32. Furthermore, Fig. 6(a) shows that the region with a large  $\kappa$  value has a high distribution density of G2 grains, indicating that soft grains with higher  $\kappa$  values preferentially lead to the cracking of adjacent hard grains. In this work, dwell fatigue tests were conducted at ambient temperature, and so orientations of  $\alpha$  grains change little before and after dwell fatigue deformation. The results of quasi *in situ* EBSD reported by Sun *et al.*<sup>31</sup> and Hémery and Villechaise<sup>32</sup> also show that  $\alpha$  grains in titanium alloys did not rotate significantly during room-temperature deformation.

Moreover, Fig. 6(b) shows that the variation trend of the  $\kappa$  value with  $\varphi$  angles is similar to that of the GND density value in hard grains near GBs with  $\varphi$  angles. The GND density value in hard grains near GBs can reflect the magnitude of stress concentration there. The higher stress concentration caused by the dislocations pile up of soft grains can activate more dislocations in hard grains near soft/hard GBs.<sup>16,17</sup> In particular, the  $\kappa$  value is

maximal for  $\gamma = 90^\circ$  and  $\varphi = 25^\circ$ , suggesting that there is the highest stress normal to the basal plane of adjacent hard grains. Fig. 6(a) also shows that the distribution density of G2 grains is high near the region where  $\gamma = 90^\circ$  and  $\varphi = 25^\circ$ . In addition, when the  $\varphi$  angle of soft grains is around  $25^\circ$ , the GND density in hard grains is higher, as shown in Fig. 6(b). Accordingly, the “worst” combination causing cracking is considered as that in which a hard grain is adjacent to an  $\alpha$  grain with the  $c$ -axis normal to the load direction and a  $\varphi$  angle of  $25^\circ$ . In addition, two types of dislocations are formed in crystals during deformation, namely GND and statistically stored dislocation.<sup>33,34</sup> The hard grain is poorly oriented for slips (low basal and prismatic slip Schmid factors), resulting in little plastic deformation<sup>35–37</sup> and thus low statistically stored dislocation density during dwell fatigue. Therefore, the GND density has been mainly determined and discussed in this paper.

Herein a schematic diagram of the crack initiation mechanism is shown in Fig. 7. During dwell fatigue, extensive prismatic  $\langle a \rangle$  dislocations are readily activated in the soft grain and accumulate at the soft/hard GB, which will cause the stress redistribution (namely load shedding) from soft grains to hard grains, resulting in a local stress concentration at the soft/hard GB.<sup>17,22,38,39</sup> Furthermore, the stress concentration at the GB can induce basal slips in the hard grain,<sup>5,25,40</sup> which the measured GND density in the work can reflect. The stress component along the hard grain's basal plane can be released by basal slips, while the stress component normal to the hard grain's basal plane still exists. Therefore, it is considered that the GND density of hard grains near soft/hard GBs can be directly proportional to the magnitude of stress normal to hard grains' basal planes caused by dislocation pile up of soft grains. Moreover, the formed basal slip band may become a preferential crack nucleation site.<sup>5,16,25</sup> When the  $\varphi$  angle of the soft grain is between  $15^\circ$  and  $30^\circ$ , the normal stress is relatively high. Besides, the high normal stress combined with the applied stress is sufficient to overcome the cohesive strength of  $\alpha$  grains, which assists the cracking of the adjacent hard grain along the basal plane with “quasi-cleavage”, as shown in Fig. 7. When the  $\varphi$  angle is between  $0^\circ$  and  $15^\circ$ , the normal

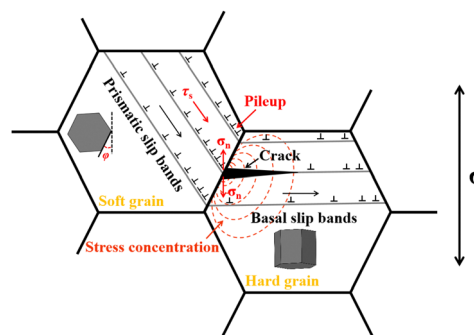


Fig. 7 Schematic diagram of dwell fatigue cracking involving dislocation pile up at the soft/hard GB. Stress concentration induces dislocation activation and normal stress ( $\sigma_n$ ) on the basal plane of the hard grain. Combined with the applied stress ( $\sigma$ ), the normal stress leads to the crack nucleation.



stress is relatively low. Therefore, it is difficult to form cracks in the hard grain.

Moreover, the time-dependent deformation of  $\alpha$  grains also affects the initiation of dwell fatigue cracks.<sup>37,41–45</sup> Numerous investigations show that the time-dependent deformation behavior (or strain rate sensitivity) of  $\alpha$  grains is closely relevant to their crystallographic orientation.<sup>37,43,44</sup> The strain rate sensitivity of  $\alpha$  grains varies with the declination angle ( $\gamma$  angle) between their  $c$ -axis and the applied stress, with a larger  $\gamma$  angle associated with stronger strain rate sensitivity.<sup>44</sup> Meanwhile, a strong strain rate sensitivity can increase the stress redistribution from soft grains to hard grains caused by the load shedding, thereby increasing the local stress of hard grains and promoting the cracking of hard grains.<sup>37,45</sup> The above results<sup>37,44,45</sup> considering time-dependent deformation show that the soft grain with a large  $\gamma$  angle is more likely to cause the cracking of the adjacent hard grain, which is consistent with the results we obtained based on the calculated stress component. Furthermore, the time-dependent deformation behavior is related to the  $\phi$  angle in addition to the  $\gamma$  angle. The crystallographic orientation favorable for slips (*i.e.* large macroscopic Schmid factor associated with  $\gamma$  and  $\phi$  angles) can intensify the magnitude of stress drop, resulting in a large stress oscillation.<sup>43</sup> Although the time-dependent deformation cannot be effectively characterized by our current research methods, the issue is indeed worthy of further research.

## 4. Conclusion

In this study, the quantitative effect of soft grains with different  $\langle 11\bar{2}0 \rangle$  orientations on the cracking of hard grains was investigated in the equiaxed titanium alloy after dwell fatigue testing, showing that, in addition to the angle between the  $c$ -axis and the loading direction, the “rotation angle”, corresponding to the  $\langle 11\bar{2}0 \rangle$  orientation, of adjacent grains plays a vital role in the crack nucleation of hard grains. The change in the “rotation angle” affects the normal stress on the basal plane of hard grains, leading to the difference in the crack nucleation probability of hard grains. Soft grains with a “rotation angle” of  $15^\circ$ – $30^\circ$  are prone to cause cracking of adjacent hard grains. A stress component parameter  $\kappa$  was proposed to predict the crack nucleation. The hard grains preferentially initiate cracks when the  $\kappa$  value is larger than 0.32.

## Author contributions

Zhichun Zhou and Chao Fang conducted the experiments, analyzed the data and prepared the manuscript. Kui Du and Binbin Jiang revised the paper and directed the entire research. The project was supervised by Kui Du and Jianke Qiu. All authors contributed to the discussion.

## Conflicts of interest

There are no conflicts to declare.

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