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Research progress of Mn doped phosphors

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In this review article, Mn applications have been divided into three parts. We provide an overview of recent progress in developing Mn⁴⁺ doped red phosphors for promising application in warm white light-emitting diodes. In addition, we summarize reports on Mn²⁺ co-doped Eu²⁺ or Ce³⁺ phosphors that can produce white light with outstanding color rendering index (CRI) and color coordinates *via* energy transfer processes, and report the application of Mn²⁺ in green to NIR phosphors.

1 Introduction

The valence electron configuration of transition metal elements is $(n-1)d^{1-8}ns^{1-2}$, with the common characteristic that the d orbitals are not full. Since the energy of the (n-1)d orbital is similar to that of the ns orbital, and the d electrons can completely or partially participate, there are no significant changes in the chemical properties from left to right in the same cycle. In addition, the transition metal elements can lose s electrons as well as d electrons; hence, the same element has a variety of valence states. Mn2+ and Mn4+ ions are commonly used as light-emitting centers. Nowadays, phosphor-converted white light emitting diodes (pc-WLEDs) based solid-state lighting has received much attention, due to the characteristics of high efficiency, high power efficiency, low applied voltage, reliability, long operation life and environmental friendliness, compared to other traditional lighting technologies, such as incandescent, halogen, xenon and fluorescent lamps. 1-5 In order to obtain highly efficient WLED devices, several aspects should be taken into account and optimized, including semiconducting components, phosphors and packaging technologies. Phosphors are indispensable components of pc-WLED devices, which can decide the color rendering index, correlated color temperature, rated average life and luminous flux. One method for obtaining these materials is to combine vellow-emitting phosphors with highly efficient (Ga,In)N blue LEDS. For blue chips, Y₃Al₅O₁₂:Ce³⁺ (YAG:Ce) is the commonly used phosphor; however, due to the lack of the red component of YAG:Ce phosphor, this approach has a poor color rendering index and a high correlated color temperature, which cannot fulfill the requirements for general illumination.6-9 According to Blasse, Mn⁴⁺ absorbs in the entire ultraviolet region and subsequently emits phosphorescence in the deep red (620–680 nm) region, due to transitions from $^{2}E \rightarrow {}^{4}A_{2}$, which indicates that the phosphor doped Mn⁴⁺ can be a good candidate for red fluorescent powder. It has been observed that multi-phase

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phosphor converted systems commonly lead to the strong reabsorption of the blue light by the red and green phosphors. Moreover, such a device is very complicated and expensive, and it is difficult to control the color balance. To solve the abovementioned problems, single-emitting component phosphors with UV LEDs chips, which have excellent color rendering indexes and high luminous efficiencies, have been researched. One of the way to generate white light from single-composition phosphors is by co-doping the sensitizer and activator into a crystalline matrix, with the principle of energy transfer from sensitizer to activator, such as Eu²⁺/Mn²⁺, Ce³⁺/Mn²⁺. Various white light based single-composition phosphors have obtained, such as CaAl₂Si₂O₈:Eu²⁺,Mn²⁺,¹¹ Ba₃MgSi₂O₈:Eu²⁺,Mn²⁺,¹² KMg₄- $(PO_4)_3$: Eu^{2+} , Mn^{2+} , 13 BaSi $_3O_4N_2$: Eu^{2+} , Mn^{2+} , 14 Ca $_2$ Gd $_8$ (SiO $_4$) $_6O_2$:Ce $^{3+}$ / Mn²⁺, 15 and Ca₉La(GeO₄)_{0.75}(PO₄)₆:Ce³⁺,Mn²⁺. 16 Mn²⁺ is influenced by the crystal field, and hence it can easily change luminescence from green to deep red and even near-infrared (NIR). Because the near infrared light has the advantages of large penetration depth, less interference, etc., it has become an important topic in the field of optical fiber communication, solid state lasers, light emitting markers and fluorescence immunoassays. The applications of Mn ion doped phosphors are shown in Fig. 1.

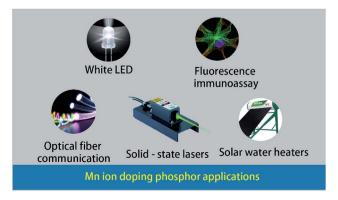


Fig. 1 Applications of Mn ion doped phosphors. 11,17,50

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2 Application of Mn⁴⁺

2.1 Introduction to red phosphors

To date, the most practical phosphor is the blue LED chip excited YAG:Ce³⁺. In order to get white-light emission and a high CRI, improvements must be made by adding a high efficiency red emitting phosphor to the system to solve the problem of the lack of a red component; hence, several red phosphors have been developed. Red fluorescent powder can be divided into three categories: linear emission red phosphor excited by rare earth ions, broadband emission red phosphor excited by rare earth ions, and red phosphor excited by transition metals. With their own characteristics, excellent red phosphors can be obtained by a combination of these.

Linear-emission phosphors excited by the rare ions Eu³⁺ and Sm³⁺ have pure color and high luminous efficiency; however, there is no wide excitation spectrum in the blue and UV regions, which results in a lower luminous efficiency. It is feasible to broaden the excitation spectrum and improve the luminous efficiency, although the effect is not ideal, but this limits the practical application of linear phosphors. Therefore, the application of red emitting Sm³⁺ or Eu³⁺ ions doped phosphors is limited in the present W-LEDs, due to their sharp absorption peaks in the blue and UV region.

In order to create a warm WLED, the development of efficient red phosphors is necessary. The new phosphors are mainly based on rare earth-doped schemes, such as Eu²⁺-doped nitrides, aluminates, oxynitrides and aluminosilicates. Among them, oxynitrides and nitrides have been the subject of maximal studies, due to their excellent performances, *e.g.*, chemical stability and high thermostability; however, these materials must be synthesized at high temperatures and pressures. The strict preparation conditions and high price of raw materials lead to this phosphor being quite expensive. In addition, the reabsorption of the coexistent green emission of WLEDS is not avoided for Eu²⁺-doped red phosphors with blue absorption.¹⁸

2.2 Luminescence properties of transition metal ions

Transition metal ions have unfilled d orbitals, with the electronic configuration of $3d^n$ (0 < n < 10). The Hamilton operator is as follows:

$$H = H_0 + H_{\rm CF} + H_{\rm C} + H_{\rm SO} \tag{1}$$

where H_0 includes the kinetic energy of electrons and the Coulomb energy of the nucleus acting on the electron. $H_{\rm C}$ is the Coulomb energy between electrons, $H_{\rm SO}$ is the spin-orbit coupling energy and $H_{\rm CF}$ is the energy of the crystal field.

The 3d electron is in the outer layer, which is sensitive to the crystal environment; hence the $H_{\rm CF}$ and the Coulomb energy between electrons ($H_{\rm C}$) can be compared. The crystal-field spectral term usually has two types of energy representations with weak crystal field and strong crystal field. We always apply the strong crystal field to deal with problems. The first consideration is how the d orbital is affected by the tetrahedral and octahedral crystal fields. The crystal field causes the five

degenerate d orbitals to split into the doubly degenerate $e_{\rm g}$ and triply degenerate $t_{\rm 2g}$. The distance between $e_{\rm g}$ and $t_{\rm 2g}$ is denoted as 10Dq, and Dq is a parameter to characterize the field strength of the octahedral.

With the single electron orbital wave function of e_g and t_{2g} as basic functions, considering the impact of Hc, a new intrinsic function is characterized by $^{2S+1}\Gamma$, where S is the total spin, and there is an irreducible representation of the octahedral group of the intrinsic state. The $^{2S+1}\Gamma$ spectral term is similar to that used in the study of the free ion ^{2S+1}L spectra of rare earth elements. The spin orbit interaction H_S in transition metal ions is weaker than that of rare earth ions, which is usually treated by the perturbation theory and the deviation of the crystal field in the octahedral symmetry. These interactions cause $^{2S+1}\Gamma$ to split into multiple states. The complete theoretical treatment of the transition metal ions eigenstate was given by S. Sugano and S. Tanabe. Direct coulombic interactions S0 interactions S1 in the deviation of the crystal field in the ransition metal ions eigenstate was given by S2. Sugano and S3.

$$(H_{\rm C})_{\rm dir} = \sum_{k} f_k F^{(k)} \tag{2}$$

where the fitting parameter

$$F^{(k)} = e^2 \int_0^\infty \int_0^\infty \frac{r_{<}^k}{r_{>}^{k+1}} [R_i(r_i)R_j(r_j)]^2 dr_i dr_j$$
 (3)

For the ll' configuration, the k in F(k) is an even number from 2 to min (2l, 2l'). For transition metal ions with 3d electrons, the values can be 0, 2, 4. The relationship between Racah parameters and F(k) is as follows:

$$B = \frac{1}{49}F^{(2)} - \frac{5}{441}F^{(4)} \tag{4}$$

$$C = \frac{35}{441} F^{(4)} \tag{5}$$

$$A = F^{(0)} - \frac{1}{9}F^{(4)} \tag{6}$$

Since the *A* parameter only affects the energy levels, we only consider *B* and *C*. The ratio of B/C for the all the transition metal ions is almost equal to 4.5, which can be seen as approximately constant. The distance between different spectra of each ion depends on the two parameters, Dq and *B*. In the work of S. Sugano and Y. Tanabe, the Tanabe–Sugano (T–S) (Fig. 2)¹⁹ diagram shows the relationship between the energy levels of metal ions E/B and crystal field strength Dq/B of the d^2 - d^8 configurations in the crystal field. From the T–S graph, we can identify each excitation peak corresponding to energy levels in the excitation spectra.

2.3 The transition of Mn⁴⁺ in the host²⁰

Mn⁴⁺ doped materials, such as oxides and fluorides, are commonly used as red emitting phosphors. It was discovered that the Mn⁴⁺ ion can only be stabilized in an octahedral lattice site where it commonly shows strong and broad adsorptions

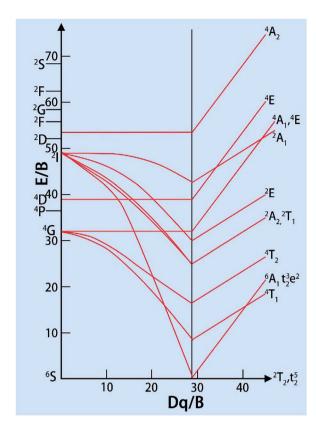


Fig. 2 Tanabe-Sugano diagram for the d⁵ electron configuration.¹⁹

ranging from 300 to 480 nm and emits red light between 600 and 700 nm with 5 sharp peaks. As the lowest energy state 2 E(t₂ 3) hardly changes in different crystal fields, the luminescence properties are very similar from host to host, but intensely depend on the covalence of the Mn⁴⁺-ligand bonding. Moreover, Mn is magnetic and will produce local magnetic ions when introduced into luminescent materials; therefore, Mn may have huge potential application value.

Doped in a solid, the Mn^{4+} ion has a strong crystal field, due to its highly effective positive charge. Consequently, the spin-forbidden $^2E_g \rightarrow ^4A_{2g}$ transition (sharp line) always controls the emission spectrum of Mn^{4+} , while two broad bands corresponding to the $^4A_{2g} \rightarrow ^4T_{1g}(^4F)$ and $^4A_{2g} \rightarrow ^4T_{2g}$ spin-allowed transitions are observed in the excitation spectra. The third spin-allowed transition ($^4A_{2g} \rightarrow ^4T_{1g}(^4P)$) frequently remains under cover by the host lattice absorption and the charge transfer. The $^2E_g \rightarrow ^4A_{2g}$ transition has a wide variation, which is due to the energy of the 2E_g state in the d^3 electronic configuration being independent of crystal field (Fig. 3). 21

2.4 Spectroscopic parameters of Mn⁴⁺ in different hosts²¹

The positions of zero-phonon lines in the $^2\mathrm{E_g} \rightarrow ^4\mathrm{A_{2g}}$ spin-forbidden transition of the $\mathrm{Mn^{4+}}$ ion, 'host cation ligand' bond distance and Racah parameter B have been summarized numerous crystals. On an average, the distance from $\mathrm{Mn^{4+}}$ to ligand in the fluoride crystals (about 1.67–1.86 Å) is the shortest

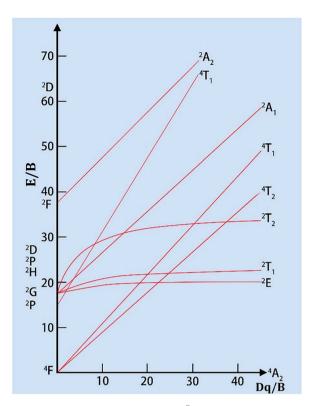


Fig. 3 Tanabe—Sugano diagram for the d³ electron configuration in an octahedral crystal field.²¹

and the following trends (Fig. 4) are formulated from the graphical representations of the data gathered in Table 1, which is suitable for most compounds. The B values are greater in fluorides and thus, the emission lines of ${}^2E_g - {}^4A_{2g}$ are significantly increased in oxides and in materials with the garnet structure (up to 2.11 Å), as shown in Table 1. The energy difference among the ${}^4A_2 \rightarrow {}^4T_2$ and the ${}^4A_2 \rightarrow {}^4T_1$ transitions and the Racah parameter B can be derived from the expression²²⁻²⁴

$$\frac{\text{Dq}}{B} = \frac{15(x-8)}{x^2 - 10x} \tag{7}$$

where x is defined as

$$x = \frac{E(^{4} A_{2g} \rightarrow ^{4} T_{1g}) - E(^{4} A_{2g} \rightarrow ^{4} T_{2g})}{Da}$$
(8)

The Racah parameter C can be evaluated by the following expression, according to the peak energy of ${}^{2}E_{g} \rightarrow {}^{4}A_{2g}$ obtained from the emission spectrum.

$$\frac{E(^{2}E_{g}^{-4}A_{1g})}{B} = 3.05\frac{C}{B} + 7.9 - 1.8\frac{B}{Dq}$$
 (9)

The local crystal-field strength Dq is determined by the mean peak energy of the $^2A_2 \rightarrow ^4T_2$ transition according to the following expression:

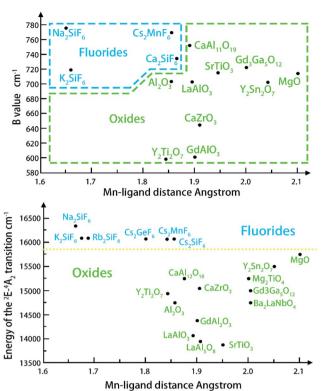


Fig. 4 Relationship between the Mn^{4+} ligand²¹ bond distance and the Racah parameter B (top), and the position of the 2E_a level (bottom).²¹

$$Dq = \frac{E(^{4}T_{2g} - ^{4}A_{2g})}{10}$$
 (10)

The $^2\mathrm{E_g}$ level is always lower for oxides; however, the value of B is also close to that of fluorides in some special cases (Fig. 4). What causes the condition is that the position of the $^2\mathrm{E_g}$ level also depends on the value of the second Racah parameter C.

Such a distinct and explicit separation of fluorides and oxides with respect to the energy position of the $^2\mathrm{E_g}$ level indicates the potential role of the $\mathrm{Mn^{4^+}}$ ion as a reliable indicator of covalency/ionicity of a particular host. On the other hand, $\mathrm{Mn^{4^+}}$ luminescence exhibits sharp emission lines and broad excitation bands in the ultraviolet–blue region, which makes it suitable for excitation by UV and decreases the re-absorption effect when mixing with other yellow and green phosphors. With so many unique features, $\mathrm{Mn^{4^+}}$ can meet the requirement of an ideal red phosphor for warm LEDS. Moreover, the easy to obtain and cheap manganese raw materials are able to reduce the red phosphors. As a consequence, many researchers have focused on developing the $\mathrm{Mn^{4^+}}$ activated red phosphors (Table 2).

2.5 Examples of Mn⁴⁺ doped red phosphors

Numerous researchers have focused on the oxide host, due to their outstanding performance such as high stability. Peng *et al.* have synthesized an efficient red oxide phosphor $Sr_4Al_{14}O_{25}:Mn^{4+},^{28}$ which is prepared by a traditional solid-state reaction in much milder conditions. The red emission intensity

Table 1 Summary of some spectroscopic and structural characteristics

Host	Cation–ligand distance, Å	Racah parameter <i>B</i> , cm ⁻¹	Position of the ${}^2\mathrm{E_g}$ level, cm^{-1}
Na ₂ SiF ₆	1.674	775	16 207
K ₂ SiF ₆	1.6829	719	16 109
Rb ₂ SiF ₆	1.689	625	16 082
Cs ₂ GeF ₆	1.798		16 038
Cs ₂ SiF ₆	1.867	722	16 033
Cs_2MnF_6	1.857	775	16 031
Mg_2TiO_4	1.948	700	15 253
CaZrO ₃	1.9137	650	15 054
$Gd_3Ga_5O_{12}$	2.005	722	14 980
Al_2O_3	1.858	700	14 786
$GdAlO_3$	1.903	600	14 347
LaAlO ₃	1.895	700	14 045
LiAl ₅ O ₈	1.900		13 966
SrTiO ₃	1.952	719	13 827
$Y_2Ti_2O_7$	1.844	600	14 956
$Y_2Sn_2O_7$	2.048	700	15 563
$GaAl_{12}O_{19}$	1.8791	750	15 241
MgO	2.1135	722	15 809
Ba ₂ LaNbO ₆	1.9975	670	14 680

Table 2 Known compounds of red phosphors doped with Mn⁴⁺

Chemical composition	λ_{ex} (nm)	$\lambda_{\mathrm{em}} (nm)$	Incorporated site	Ref
BaTiF ₆ :Mn ⁴⁺	367, 468	634	Ti ⁴⁺	35
Na ₂ SiF ₆ :Mn ⁴⁺	360, 460	600-650	Si ⁴⁺	36
Na ₃ GaF ₆ :Mn ⁴⁺	460	630	Ga ³⁺	37
$Sr_2MgAl_{22}O_{36}:Mn^{4+}$	312	658	Al^{3+}	38
$CaAl_{12}O_{19}:Mn^{4+}$	325, 457	600-700	Al^{3+}	39
Li ₂ MgGeO ₄ :Mn ⁴⁺	323	671	Ge ⁴⁺	40
$(K,Rb)_2Ge_4O_9:Mn^{4+}$	295, 460	620-727	Ge ⁴⁺	41

of $Sr_4Al_{14}O_{25}$: Mn^{4+} is greater than that of the reported commercial phosphors, such as $(K,Rb)_2Ge_4O_9$: Mn^{4+} (ref. 29) and $CaAl_{12}O_{19}$: Mn^{4+} .

A series of novel red emitting Mg3Ga2GeO8 under near-UV excitation was successfully synthesized by Ding et al. with the traditional high-temperature solid-state reaction.31 In this Mg₃Ga₂GeO₈ system, there is one tetrahedral site, which is occupied by Ge4+/Ga3+, and an octahedral site, occupied by Mg²⁺/Ga³⁺. According to radius analysis and luminescence properties of Mg₃Ga₂GeO₈:x%Mn⁴⁺, we can conclude that Mn⁴⁺ reliably occupies the Mg²⁺/Ga³⁺ site in this system. When using MnCO₃ as the raw material, the intensity of the emission is relatively weak, compared to the sample of MnO2 as the raw material and the peak shape does not change. In the MnCO₃, the same Mn²⁺ can just be oxidized to Mn³⁺, rather than Mn⁴⁺, as shown in Fig. 5. The emission properties of Mg₃Ga₂GeO₈:x%Mn⁴⁺ are also shown in Fig. 5. It can be easily seen that the emission intensities reach a maximum at x = 0.5% and the calculated fullwidth at half-maximum (FWHM) of the Mn⁴⁺ emission spectrum is 24 nm, which indicates that the Mg₃Ga₂GeO₈:x%Mn⁴⁺ phosphor

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possesses high color purity. Furthermore, Mg₃Ga₂GeO₈:0.5% Mn⁴⁺ exhibits outstanding quantum efficiency (64.7%).

The oxide matrix with strong covalent, low thermal vibration and weak polarizability is favourable for red light emission. The Mn⁴⁺ doped Y₃Al₅O₁₂ phosphor was prepared by a high temperature solid state reaction method. Huang et al. synthesized the Y3Al5O12:Mn4+ red phosphor and explored its luminescence properties. At 467 nm, this phosphor emits red light.³¹ The main emission peaks are at 643 and 670 nm, all from the ⁵E \rightarrow ⁴A₂ transition of Mn⁴⁺ in Y₃Al₅O₁₂:0.01Mn⁴⁺. The CIE of the $Y_3Al_5O_{12}$:0.01Mn⁴⁺ phosphor is (0.701, 0.299) located in the red area. The excitation spectrum of the emission peak at 670 nm is composed of 400, 450, 468, 474, 482 and 494 nm wide excitation bands.

Novel red-emitting phosphors K₂BaGe₈O₁₈:Mn⁴⁺ were synthesized via the high temperature solid-state reaction method.32 The excitation spectrum of K2BaGe8O18:Mn4+ monitored at 666 nm shows a wide band with two evident peaks at 319 and 468 nm. Upon 468 nm excitation, the emission spectrum presents a narrow band from 600 to 700 nm (red region) with the maximum emission at around 666 nm, which originates from the ²E_g-⁴A₂ transition of Mn⁴⁺ in the GeO₆ octahedral environment, as shown in Fig. 6; the CIE coordinates of the phosphors are (0.702, 0.298) and the QY is 32.9% upon 365 nm excitation. The results indicate that this kind of phosphor can be a supplement of red-components in UV and blue light excited

In the oxide matrix, Mn⁴⁺ ions replace cations such as Al³⁺, Si⁴⁺, Zr⁴⁺, Ti⁴⁺, Ge⁴⁺, and occupy the center of the octahedron. When Mn4+ replaces ions with different valences, charge compensation is required. Therefore, the emission intensity of the phosphor can be enhanced by performing charge

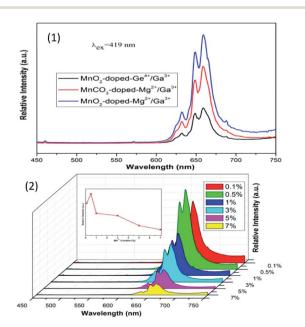


Fig. 5 (1) PL properties of samples with different Mn sources and doped sites. (2) PL under 419 nm excitation of the Mg₃Ga₂GeO₈:x% Mn⁴⁺ series phosphors.³¹

compensation. Mg2+ has been very extensively researched as a highly efficient charge compensator. The substitution of the Mg⁴⁺-Mg²⁺ ion pair to replace the Al³⁺-Al³⁺ ion pair was made under the condition that no additional charge compensation and oxygen atoms are required. Ge^{4+} , Ca^{2+} , Li^+ , Na^+ , K^+ , Cl^- can also be used as charge compensators to enhance the luminescence properties of red phosphors.33,34 Different charge compensators have different effects on the red luminescence properties of Mn⁴⁺.

The chemical stability of fluorides is appropriate, and the phonon energy is relatively low, so fluorides can be used as the host material for various luminescent rare earth ions. Na₂SiF₆:Mn⁴⁺ red fluoride phosphors have been synthesized by the exothermic reduction reaction method, as shown in Fig. 7.36 The phosphors have intense blue-light absorption and strong narrow-band red emission. When the Na2SiF6:Mn4+ red phosphor was fabricated, a prominent improvement in the CRI of white LEDs was observed. The Na₂SiF₆:Mn⁴⁺ phosphors have efficient red luminescence excited by blue light, which is a potential candidate for improving the color reproducibility of white LEDs with simple post-processing, a short reaction time

Mn⁴⁺-activated fluoride phosphors, such as A₂BF₆ (A = Na, K, Rb, Cs; A₂ = Ba, Zn; B = Si, Ge, Ti, Zr, Sn), have attracted

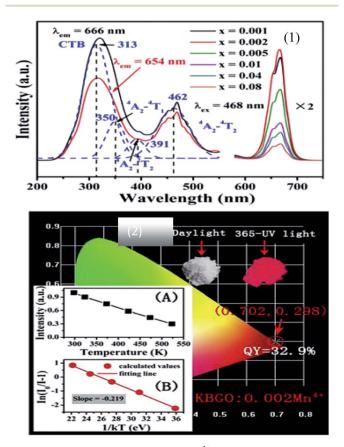


Fig. 6 (1) PL spectra of K₂BaGe₈O₁₈:Mn⁴⁺ with different concentrations, and PL excitation spectra of the K₂BaGe₈O₁₈:0.002Mn⁴⁺ sample monitored at 666 nm and 654 nm. (2) Chromatic CIE coordinates, quantum yields and variation of emission intensity depending on the temperature of the representative K2BaGe8O18:Mn4+.32

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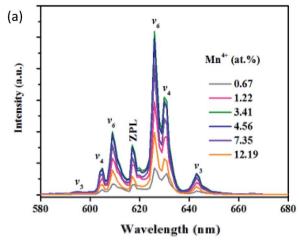
PLE
-h_{em}=631 nm
(1)
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-h_{ex}=460 nm
(c)
-h_{em}=631 nm
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-h_{ex}=460 nm
(d)
-h_{ex}=460 nm
-h_{em}=631 nm

Fig. 7 (1) Excitation and emission spectra of the $Na_2SiF_6:Mn^{4+}$ phosphor with HF concentrations of (a) 10, (b) 20, (c) 30, and (d) 40 wt%. (2) Luminescence spectra based on the $Na_2SiF_6:Mn^{4+}$ phosphor under various drive currents.³⁶

considerable attention as highly promising red phosphors for warm white light emitting diodes (w-LEDs). These fluoride phosphors have been synthesized via traditional chemical routes with HF-solution. In addition to the possible dangers of hypertoxic HF, the uncontrolled precipitation of fluorides and the extensive processing steps produce large morphological variations, limiting its application. Recently, Song made a prototype w-LED with $K_3AlF_6:Mn^{4+}$ as the red light component, via an efficient and water-processable cation-exchange green route. Combined with YAG:Ce³⁺, the prototype of the sample shows an efficient luminous efficacy beyond 190 lm W⁻¹, along with an excellent color rendering index (CRI, $R_a = 84$) and a lower correlated color temperature (CCT = 3665 K), as shown in Fig. 8.

Thus far, Mn^{4+} as an activator doped oxide/fluoride red phosphor has attracted significant attention, due to its special advantages such as relatively simple synthesis methods and low cost. The spectral positions produced by the $^2E_g \rightarrow ^4A_{2g}$ transition can be modified, and the excitation region of the phosphor is located in the ultraviolet and blue band. Therefore, the Mn^{4+} doped red phosphor can be used as a substitute for rare earth doped red phosphors, and has a bright future in the field of warm white LEDs.

Although a large number of Mn⁴⁺ doped red phosphors have been fully studied, there are still some questions worthy of further exploration. For the Mn⁴⁺ doped oxide host phosphor, it is desirable to achieve a blue shift of emission from deep red to orange/red and enhanced absorption in the blue wavelength region. One possible solution is doping to increase the absorption in the blue wavelength region for



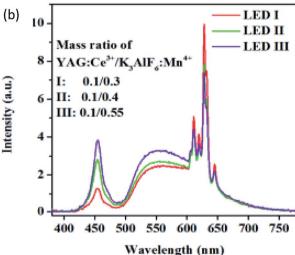


Fig. 8 (a) Concentration dependent emission spectra of K_3AlF_6 :Mn⁴⁺. (b) Electroluminescence spectra of prototype LEDs, employing K_3AlF_6 :Mn⁴⁺ as the red converter component, commercial yellow YAG:Ce³⁺ and a blue InGaN chip at a 40 mA drive current.⁴²

 ${\rm Sr_4Al_{14}O_{25}:Mn^{4+}}$ doped ${\rm Mg^{2+}}$; because the neighboring ${\rm Mn^{4+}-Mn^{4+}}$ is replaced by ${\rm Mn^{4+}-Mg^{2+}}$, the non-radiative decay rate of the ${}^2{\rm E_g}$ state decreases and the luminous intensity of the phosphor is enhanced. For fluoride phosphors, how to improve their thermal stability and resistance to humidity needs further study. Surface treatment by external coating can improve the thermal stability and moisture resistance. Finally, the quantum efficiency of materials can be further improved by optimizing synthesis methods and reducing the lattice defects of materials.

3 Application of Mn²⁺

3.1 Application of Mn²⁺ in green to NIR phosphors

Most phosphors applied in white LEDs are oxides, sulfides or nitrides doped with rare-earth ions, which play a significant role in controlling the color of the light. Most of these rare-earth ions are very expensive and some chloride and oxide doped rare-earth elements are toxic and harmful, which limit their RSC Advances Review

development. Consequently, the search for a cheap and nontoxic phosphor has become an important research area and Mn^{2+} doped oxides and nitrides have become hot research topics. With the d electrons within the 3d shell, Mn^{2+} emission can be affected by the crystal field, which results in the emission changing from green to red. When Mn^{2+} is located in the tetrahedral crystal field, it can emit green light. In other words, through changing the crystal field conditions, the spin-forbidden Mn^{2+} forbidden d–d transition ($^4T_1 \rightarrow {}^6A_1$) can emit green light.

Rare earth ions (Eu²⁺, Eu³⁺, or Pr³⁺) or transition metal elements (Cr3+ or Mn2+) can serve as luminescence centers in NIR phosphors. However, for biological applications, these hygroscopic sulphides of Eu²⁺ show a succession of problems in metal-activated oxide compounds, such as Cr³⁺ or Mn²⁺-doped phosphors, which have outstanding chemical stability and suitable chemical instability. Lately the focus has shifted to the transition-emission range in the optical window (from 650 to 1300 nm) and long-lasting persistent luminescence. With the advantages of high signal to noise ratio, deep tissue penetration, and maximum transparency for silicon-based optical fibers and so on, near infrared (NIR) emitting materials have distinct potential for medical diagnostics, laser systems, and optoelectronic devices. 43-45 On the basis of coordination and the T-S diagram, due to weaker crystal symmetry (angular constraints) and stronger ligand field strength (interionic distance) than that of the four/six-coordinated environments, it is feasible to speculate that the five/seven-coordinated environments of the Mn2+ ion are good for NIR emission generation.46-49 There is a complete list of known green to NIR persistent phosphors activated by Mn²⁺ in Table 3.

The rare-earth-free narrow-band green-emitting KAlSi₂O₆:-Mn²⁺ phosphor has been synthesized by Ding et al. The excitation and emission spectra are shown in Fig. 9. The five distinct peaks of PL excitation are in good agreement with the Mn²⁺ absorption transitions; specifically, it can match blue GaN chips, which make its application in white LEDs possible. Moreover, a green light emission peak at 513 nm with full-width at half-maximum of 30 nm can be seen from Fig. 9, which is attributed to the spin forbidden d-d transition (${}^{4}T_{1} \rightarrow {}^{6}A_{1}$) of Mn²⁺. The electrons could be relaxed from these excited states to the ⁴T₁(G) state by a non-radiative relaxation process and then be transferred back to the ground state ⁶A₁(S), emitting the characteristic green (513 nm) light. The CIE color coordinates of the KAlSi₂O₆:Mn²⁺ phosphor were calculated to be x = 0.27 and y = 0.64, which are much closer to the standard green coordinates (0.21, 0.71). The integrated emission intensity of KAS:3% Mn^{2+} decreased to 67.6% (250 °C) of the initial value (25 °C). The chromaticity coordinates of the fabricated white LED are (0.35, 0.36) with CCT of 4775 K. It is demonstrated that the KAlSi₂O₆:Mn²⁺ phosphor could be simulated to blue-LED for producing efficient white-light.

Moreover, good performance Mn^{2+} doped Na_2MgGeO_4 phosphors have been reported. Under UV or electron beams, the samples show bright green emission, which is ascribed to the ${}^4T_1(4G) - {}^6A_1(6S)$ transition of the Mn^{2+} ions. Based on this transition, the following phosphors were developed:

Table 3 Known compounds of green to NIR persistent phosphors activated by Mn^{2+}

Host material	Incorporated site	Emission region (nm)	Ref.
CdSiO ₃	Cd ²⁺ (vı)	550-720	50
3	()		
$MgGeO_3$	$Mg^{2+}(vi)$	600-720	51
Li ₂ ZnGe ₃ O ₈	$Zn^{2+}(vi)$	700-900	52
$Ga_3(PO_4)_2$	Ca ²⁺ (v _I)	600-750	53
$Ga_9Ln(PO_4)_7$	Ca ²⁺ (vı)	600-750	54
$(Mg_{1-x}Zn_x)_{2.97}(PO_4)_2$	$Mg^{2+}/Zn^{2+}(vi)$	630-800	55
$KMgBO_3$	Mg(vi)	600-700	56
$MgAl_2O_4$	Mg(iv)	500-560	57
Na ₂ MgGeO ₄	Mg(iv)	500-550	58
KAlSi ₂ O ₆	Al/Si(IV)	500-540	59
$\mathrm{Zn_{2}GeO}$	Zn(II)	500-570	60

Na₂Mg_{1-x}SiO₄:xMn²⁺ green phosphors, with CIE coordinates of (0.163, 0.631) and CTT value of 8643 K,⁶¹ NaAl₁₁O₁₇:Mn²⁺ green phosphors, with CIE chromaticity coordinates of (0.0725, 0.6468), and quantum efficiency of 67.84%, with excitation of 424 nm,⁶² Zn₂SiO₄:Mn²⁺ green phosphors, which exhibit widely tunable CCTs ranging from 3200 to 7000 K and color rendering indices up to 85%, depending on the mixture ratio.⁶³ Those phosphors with good performance indicators can be candidates for high power pcWLED applications. In addition, MgAl₂O₄:Mn²⁺,⁵⁷ Zn₂SiO₄:-Mn²⁺,^{64,65} green phosphors have also been reported.

Song *et al.* have reported the realization of single-band NIR UC emission in a $\mathrm{Mn^{2^+}}$ -doped KZnF₃ nanostructure by a heavy doping strategy. ⁶⁶ Fig. 10 shows the emission spectra of KZn_{1-x}Mn_xF₃ upon 396 nm light excitation. There is a single visible emission located at 585 nm when $\mathrm{Mn^{2^+}}$ concentration is less than x=0.10, corresponding to the ${}^4\mathrm{T_{1g}}(\mathrm{G}) \rightarrow {}^6\mathrm{A_{1g}}(\mathrm{S})$. An anomalous NIR emission band located at 770 nm appears when x=0.10, while at x=0.30, there is a sharp decrease due to the concentration quenching effect, where visible and NIR emission bands are featured from two different emission centers. The luminescence decay curves (Fig. 10(c) and (d)) were studied in order to obtain more insight into the active centers. The visible emission exhibited single exponential decay at a low concentration (x=0.01 and 0.05), and the decay lifetime was estimated by the following equation:

$$I(t) = I_0 + A \exp(-t/\tau) \tag{11}$$

where I(t) and I_0 represent the emission intensity at time t and 0, respectively, A is a constant, and τ is the decay lifetime. For Mn²⁺-heavy doping (x=0.10–0.30), the visible emission exhibited non-exponential decay behavior, probably associated with the prevailing non-radiative transition, which is expressed by the mean decay lifetime ($\tau_{\rm m}$) estimated by

$$\tau_{\rm m} = \int_{t_0}^{\infty} I(t)/I_{\rm max} dt \tag{12}$$

where I(t) is the luminescence intensity at t and $I_{\text{max}} = I(0)$. The non-radiative loss occurring among the Mn^{2+} ions as the Mn^{2+} concentration increases can be demonstrated from the monotonously shortened decay lifetime of the two emission bands.

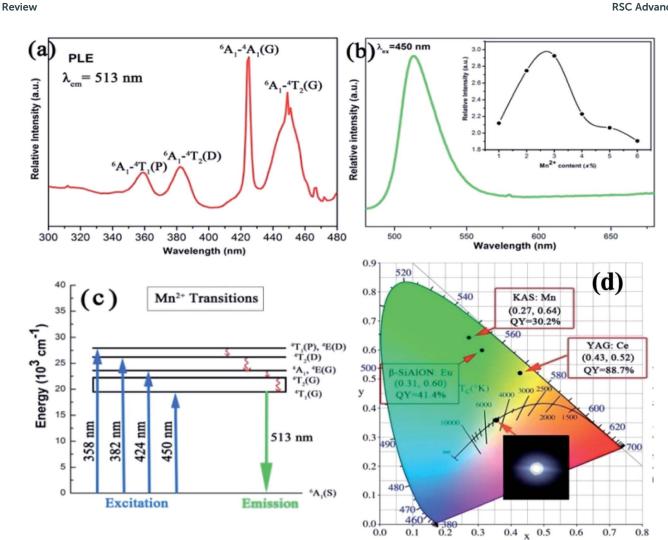


Fig. 9 (a) Photoluminescence excitation of KAlSi₂O₆: Mn^{2+} ; (b) photoluminescence emission spectra of KAlSi₂O₆ phosphors, inset: dependence of the emission intensity of KAISi₂ O_6 :Mn²⁺ phosphors on Mn²⁺ content; (c) the excitation and emission transitions of Mn²⁺ ions; (d) the CIE chromaticity coordinates and their respective quantum yields. The inset shows the phosphor simulated white light using blue (455 nm) LED.⁵⁹

To study the reason for the observed unusual NIR emission in KZn_{1-x}Mn_xF₃, the crystal structure of KZnF₃ and the activation process of Mn²⁺ were further analyzed. KZnF₃ has a typical cubic perovskite structure with a $Pm\bar{3}m$ space group with its lattice constant a = 0.405 nm, as shown in Fig. 10(a). Zn^{2+} is substituted by the incorporated Mn²⁺ ions, leading to the formation of Mn²⁺(Zn²⁺)-Mn²⁺(Zn²⁺) dimmers, especially at a high doping concentration. The ground state (${}^{6}A_{1}$, S = 5/2) and the first excited state $({}^{4}T_{1}, S = 3/2)$ are split due to the exchange interaction between Mn2+ and Mn2+ ions, which are mainly governed by the linear combination of two spin states, S_i and S_i (i.e., $(S_i + S_i)$, ..., $(S_i - S_i)$. Therefore, the spin state is S = 5, 4, 3, 2, 1, and 0 for the ground state $(^{6}A_{1g}(S)^{6}A_{1g}(S))$ and S = 4, 3, 2, and 1 for the first excited state (⁶A_{1g}(S)⁴T_{1g}(G)). The Mn²⁺-Mn²⁺ dimmers can also serve as the chromophoric centers, and the observed NIR luminescence is most likely related to the radiative transition from the excited state (S = 1) to the ground state with various spin components. There could be a shift in the emission peak caused by the changes in the spin components. The reason for decay lifetime shortening may be that the spin selection rule

has changed from spin forbidden ($S = 5/2 \rightarrow S = 3/2$) to spin allowed $(S = 1 \rightarrow S = 1)$, hence, it is reasonable to suggest that the NIR emission in KZn_{1-x}Mn_xF₃ originates from the $^{6}A_{1g}(S)^{4}T_{1g}(G) \rightarrow {^{6}A_{1g}(S)^{6}A_{1g}(S)}$ transitions of coupled Mn²⁺– Mn²⁺ dimmers.

Near-infrared (NIR) emitting Li₂ZnGe₃O₈ (LZG):xMn²⁺ has been synthesized by a conventional solid state reaction.⁵² The excitation spectrum (PLE) and emission spectrum (PL) of the $LZG:xMn^{2+}$ (x = 0.005, 0.01, 0.03, 0.05, 0.07, 0.10, 0.15, 0.20) phosphors are presented in Fig. 11(a) and (b), respectively. The emission spectrum shows a broad emission band from 650 to 900 nm, with the main peak at 832 nm. The luminescence properties can be strongly affected by the crystal field within the 3d shell with an increase in Mn2+ concentration. The PL $(\lambda_{\rm em} = 832 \text{ nm})$ intensity rapidly increases, while above this maximum intensity (0.03), a concentration quenching phenomenon occurs, as shown in Fig. 11(d). Essentially, quenching impurities cause concentration quenching, a decline in the emission intensity, as shown in Fig. 11. Non-radiative energy transfer, which include exchange interactions, radiation

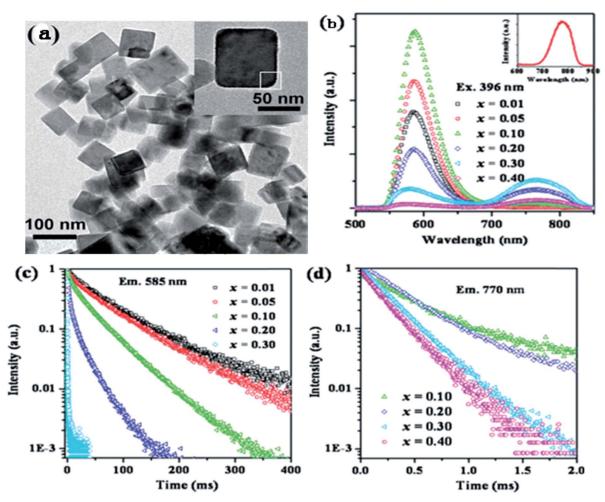


Fig. 10 (a) TEM image of $KZn_{0.80}Mn_{0.20}F_3$ nanocrystals and a typical nanocrystal (inset). (b) Emission spectra and (c) luminescence decay curves at 585 nm and (d) 770 nm for $KZn_{1-x}Mn_xF_3$ (x=0.01-0.40). The inset in (b) shows the emission spectrum of $KMnF_3$. The emission spectra and luminescence decay curves were obtained upon excitation with 396 nm light.⁶⁶

reabsorption and multipolar interactions between Mn^{2+} ions within a certain distance, lead to concentration quenching. According to calculations, the concentration quenching of Mn^{2+} ions originates from quadrupole–quadrupole interactions. Besides, the phosphor has a good thermal stability, which expands the scope of its application.

Energy transfer between ions is a rather common and very important physical phenomenon, mainly through physical processes such as collision, energy exchange, and re-absorption of radiation and radiationless transition processes, which cause the ion energy transfer to other ions. Due to the relatively rich energy levels of transition metal ions, especially in the crystals, the energy splitting is affected by the influence of the crystal field, thus becoming more intensive. At this point, there may be the possibility of energy differences between the two levels of the transition metal ions, equal to the difference between other ions. In order to achieve the energy transfer between ions, radiative or radiationless transition processes occur very easily under the action of multipole moments.

Radiationless transitions play a significant role in photophysics, photobiology and photochemistry, and include internal conversion (IC) (spin-allowed process) and intersystem crossing (ISC) (spin-forbidden process). A simplified Jablonski diagram is presented in Fig. 12, which describes the elementary processes of molecular luminescence.⁶⁸

Absorption

$$S_0 + \hbar w \xrightarrow{k_a} S_1 \tag{13}$$

Fluorescence

$$S_1 \xrightarrow{k_r} S_0 + \hbar w' \tag{14}$$

Internal conversion

$$\mathbf{S}_{1} \xrightarrow{k_{\mathrm{ie}}} \mathbf{S}_{0} \tag{15}$$

Intersystem crossing

$$S_1 \xrightarrow{k_{ise}} T_1 \tag{16}$$

(a) -0.01 -0.03 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05 -0.05

Fig. 11 (a) PLE and (b) PL spectra of the LZG: xMn^{2+} samples (0 < x < 0.2). (c) A diagram of the concentration quenching process.⁵²

Phosphorescence

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$$T_1 \xrightarrow{k'_r} S_0 + \hbar w'' \tag{17}$$

Intersystem crossing

$$T_1 \xrightarrow{k'_{\text{ise}}} S_0 \tag{18}$$

These elementary processes determine whether a molecule can be a good candidate for a dye molecule for organic solar cells or LED. Besides, there are two other significant elementary processes described in the following:

Phonon-induced electron transfer

$$DA \xrightarrow{h\nu} D^*A \to D^+A^-$$
 (19)

Phonon-induced energy transfer

$$DA \xrightarrow{h\nu} D^*A \to DA^*$$
 (20)

All the processes involve two electronic states, *i.e.*, they are non-adiabatic in nature. Vibrational relaxation (VR) exists in molecular luminescence as a non-radiative process, which can happen in either the ground or excited electronic states. The radiative transfer process is also called resonant transmission, if there is a near field force interaction between the two centers and if the energy of light radiation transfer can be one atom to 10 nm, without the help of other nearest neighbor atoms. Dexter first introduced the transfer mechanism to the energy transfer process between the centers of the luminous material, and derived a formula for the probability of resonance energy transfer between the centers. There is an assumption that emission from the ion overlaps another ion, absorbing light, so the light emitting radiation energy is absorbed by another ion, and energy transfer occurs. The former is called the energy sensitizer and the

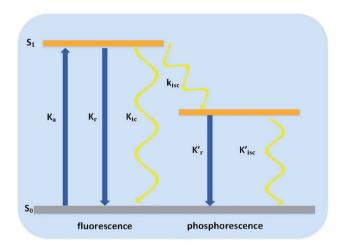


Fig. 12 Elementary photophysical processes of molecular luminescence.⁶⁸

latter is called the activator. Overlap between the emission spectra of the sensitizer and the excitation spectra of the activator is a prerequisite for the occurrence of resonant transmission. Resonant transmission is a quite significant means of energy transfer in the material activated by rare earth or transition metal elements and organic crystals. In this way, there are two centers S and A, where the center initial state is $S^* + A$, *i.e.*, the S center is in the excited state, the A center is in the ground state. The final state is S + A, *i.e.*, the S center is back to the ground state, and the A center is in the excited state. Taking S and A as the dipoles, the transition probability from the initial state to the final state is the S, A resonance energy transfer probability P_{SA} :

$$P_{\rm SA} = \left(\frac{R_0}{R}\right)^6 \times \frac{1}{\tau_{\rm S}^*} \tag{21}$$

where *R* is the distance between the center of S and A and τ_S^* is the measured lifetime of S* state, and the R_0 is defined as:

$$R_0^{\ 6} = \frac{3}{64\pi^5} \times \frac{h^4 c^4}{K^2} \sigma_{\rm A} \eta_{\rm S} \int_0^\infty \frac{\varepsilon_{\rm S}(E) \alpha_{\rm A}(E)}{E^4} dE \tag{22}$$

where h and c are the Planck constant and the velocity of light, respectively, and the K is the dielectric constant of the material. σ_A is the total A center absorption cross section and η_S is the emission efficiency of the S center. $\varepsilon_S(E)$ and $\alpha_A(E)$ are the emission spectrum of S center and absorption spectrum of A centre, respectively (Fig. 13).

3.2 The basic properties of Mn²⁺

In the last few years, non-rare earth-based eco-friendly phosphors prepared under milder conditions have received increasing interest.⁶⁹ The outer electronic structure of Mn^{2^+} is $3d^54s^2$.^{70,71} On the basis of the Mn^{2^+} Tanabe–Sugano diagram, the emission of Mn^{2^+} can be easily changed from green to deep red by transforming the crystal field strength. Because the transition is $^4T_1 \rightarrow ^6A_1$, Mn^{2^+} has broad band emission in the visible range, depending on the crystal field. Octahedral coordinated Mn^{2^+} shows an orange to red emission, while tetrahedral

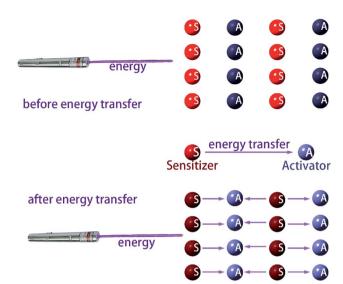


Fig. 13 Schematic diagram of the energy transfer, the dark and light colors indicate the level of energy.

coordinated $\mathrm{Mn^{2^+}}$ shows green emission. The luminescence spectrum of $\mathrm{Mn^{2^+}}$ consists of a structureless band at peak wavelengths of 490–750 nm, with the half width of 1000–2500 cm $^{-1}$. In general, the crystal field at a tetrahedral site is weaker than that of an octahedral site. When $\mathrm{Mn^{2^+}}$ is in a weak crystal field, the splitting of the excited energy in orbitals is small, resulting in a higher energy of $\mathrm{Mn^{2^+}}$ emission, but its energy is low. Generally, when $\mathrm{Mn^{2^+}}$ is substituted for alkaline-earth metals ($\mathrm{Ga^{2^+}}$, $\mathrm{Sr^{2^+}}$,

Ba²⁺), the larger lattice distortion is not conducive to the direct excitation of Mn²⁺, which makes it difficult to achieve direct luminescence. Therefore, Mn²⁺ is not suitable for individually serving as an activator with low luminescence efficiency.

Though Mn^{2+} ions emit red light from $d \rightarrow d$ transitions, the $4T_1(4G) \rightarrow 6A_1(6S)$ emission in the 3d electron configuration is weak because it is forbidden, which leads to weak absorption in the ultraviolet region and disparate fluorescence properties in different hosts. However, on introducing Eu²⁺ or Ce³⁺ as coactivators, the excitation of Mn2+ is covered by the emission of Eu²⁺ or Ce³⁺, which creates the necessary conditions for vibrational transmission, i.e., there is overlap between the emission of the activator and the excitation of the sensitizer. Meanwhile, there may be energy transfer between the sensitizer and activator. On the basis of the resonance energy transfer mechanism, the excitation energy of the sensitizer can induce the activator to emit fluorescence, while the fluorescence intensity decays. White light can be obtained by co-doping Ce³⁺/Mn²⁺ or Eu²⁺/Mn²⁺ under effective resonance-type energy transfer; a schematic diagram of energy transfer is shown in Fig. 14.72-74

3.3 The co-doping of Ce³⁺/Mn²⁺

The outer electronic structure of Mn²⁺ is 3d⁵, which presents typical d-d transitions when excited. Its photoluminescence range changes from 500 to 700 nm, which depends on the crystal field environment; however the d-d transition of Mn²⁺ is forbidden by the selection rules, and the luminescence is very weak since it can only be indirectly excited using sensitized ions or the substrate itself for energy transfer. A large number of

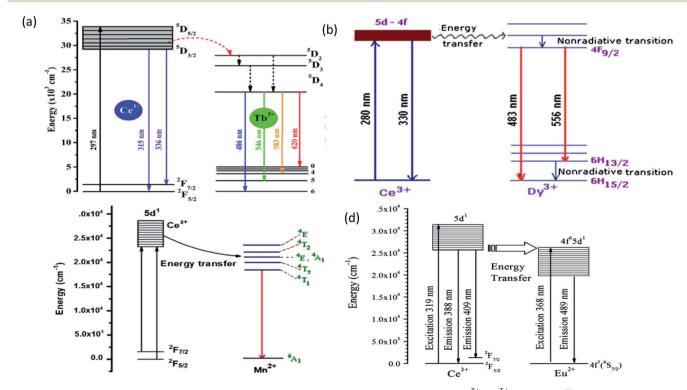


Fig. 14 (a) The schematic energy level diagram for the energy transfer process in the BaMgF₄: Ce_x^{3+} , Tb_y^{3+} phosphors. (b) Energy level diagram of KNaSO₄: Ce_z^{3+} , Ce_z

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studies have been conducted on the mechanism of energy transfer, *e.g.*, $Ce^{3+} \rightarrow Tb^{3+}$, $Ce^{3+} \rightarrow Mn^{2+}$, $Ce^{3+} \rightarrow Dy^{3+} Ce^{3+} \rightarrow Eu^{2+}$, *etc.*, as shown in Table 4, which indicates that the Ce^{3+} ion is a highly efficient ion. The energy level diagram of the energy transfer about Ce is shown in Fig. 14.

The outermost electron configuration of Ce^{3^+} is $4f^1$, which can change to a lower energy level of 5d when absorbing energy. The 5d electronic excited state is different from the 4f electronic state, which is easily shielded by $5s^25p^6$, resulting in the crystal field having a great influence on it, so the 4f–5d transition of Ce^{3^+} in the 200–400 nm spectral region often exhibits strong absorption and excitation spectra.

Ca₂Gd₈(SiO₄)₆O₂:Ce³⁺,Mn²⁺ has been prepared by Li et al., 86 and the emission spectra are shown in Fig. 15 for different concentrations of Ce³⁺ and Mn²⁺. The phosphors emit blue to white light, and with the increased Mn²⁺ concentration, there was complete conversion to yellow light when excited by UV irradiation, due to the 5d-4f transition of Ce3+ and the ${}^{4}\Gamma_{1} \rightarrow {}^{6}A_{1}$ transition of Mn²⁺, respectively. It is noteworthy that the emission peaks of Mn²⁺ have a slightly red shift from 564 nm to 570 nm. The Ce³⁺ ion simultaneously occupies the 4f and 6h sites in the Ca₂Gd₈(SiO₄)₆O₂ host, and gives different blue emission under different UV excitation. Energy transfer from $Ce^{3+} \rightarrow Mn^{2+}$ also exists in the phosphor, hence, via precise control of the contents of Mn²⁺ and Ce³⁺, a wide-range white emission can be obtained. On the other hand, changing excitation sources is another way to obtain the color-tunable emission in a signal phase host. PLE and PL spectra of the LEDs (light emitting diodes) and FEDs (field emission displays) of the Ce3+, Mn2+ co-doped phosphors, NaCaBO3:Ce3+,Mn2+ were also studied, which have potential applications in the field of white EDS. It has been shown that for the crystal structure of NaCaBO₃, the relative PL intensity of blue to yellow emission increases with the increase of the excitation wavelength from 287 to 330 nm. Moreover, $Ca_2Gd_8(SiO_4)_6O_2:Ce^{3+}$, Mn²⁺, has a good CL coefficient (K/S) relation, which was used to calculate the measured reflectance (R) for NaCaBO3, to probe the absorption of the host lattice.87

$$F(R) = \frac{(1-R)^2}{2R} = \frac{K}{S}$$
 (23)

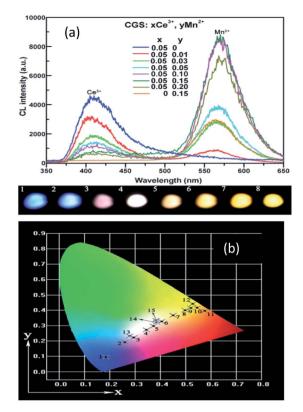


Fig. 15 (a) Emission spectra and luminescence properties of CGS:xCe $^{3+}$,yMn $^{2+}$ samples. (b) The CIE chromaticity diagram for Ca $_2$ Gd $_8$ (SiO $_4$) $_6$ O $_2$:0.05Ce $^{3+}$,yMn $^{2+}$ samples under 287 nm UV excitation. 86

where R, K, and S represent the reflectance, absorption coefficient, and scattering coefficient, respectively. Hence, the band gap of the NaCaBO $_3$ host was calculated to be approximately 4.48 eV by extrapolation, essentially owing to the transition between the valence band and the conduction band in this host.*

From Fig. 16(1) and (2), it is seen that there exists a spectral overlap between the PL spectrum of NaCa_{0.99}BO₃:0.01Ce³⁺ and the PLE spectrum of NaCa_{0.96}BO₃:0.03Mn²⁺, indicating that a resonance-type energy transfer from the Ce³⁺ to Mn²⁺ ions is expected in the co-doped sample. Due to the energy transfer

Table 4 The energy transfer about Ce³⁺ in different hosts and doping ions

Host	Energy transfer	$\lambda_{\mathrm{ex}} \left(\mathrm{nm} \right)$	$\lambda_{\mathrm{em}} (\mathrm{nm})$	Ref.
$BaMgF_4$	$Ce^{3+} \rightarrow Tb^{3+}$	297	488, 545, 584	75
$Ca_{10}(PO_4)_6S$	$Ce^{3+} \rightarrow Tb^{3+}$	280, 430, 545	410, 550	76
Sr ₂ MgSi ₂ O ₇	$Ce^{3+} \rightarrow Tb^{3+}$	228	542	77
$Sr_3Gd_2(Si_3O_9)_2$	$Ce^{3+} \rightarrow Tb^{3+}$	348	540, 550	78
KNaSO ₄	$Ce^{3+} \rightarrow Dy^{3+}$	280	330, 483, 556	79
$(Dy_{0.03}Ce_xY_{0.97-x})_3Al_5O_{12}$	$Ce^{3+} \rightarrow Dy^{3+}$	327, 367	496, 582	80
$Mg_3Ca_3(PO_4)_4$	$Ce^{3+} \rightarrow Mn^{2+}$	362, 410, 418	352, 645	81
$K_2AEP_2O_7$ (AE = Ca, Sr)	$Ce^{3+} \rightarrow Mn^{2+}$	320	534, 539	82
$Ca_4(PO_4)_2O$	$Ce^{3+} \rightarrow Eu^{2+}$	350	460, 630	83
$Ca_4Y_6O(SiO_4)_6$	$Ce^{3+} \rightarrow Eu^{2+}$	356, 280	426, 527	84
$Na_3Ca_6(PO_4)_5$	$Ce^{3+} \rightarrow Eu^{2+}$	319	391, 487	85

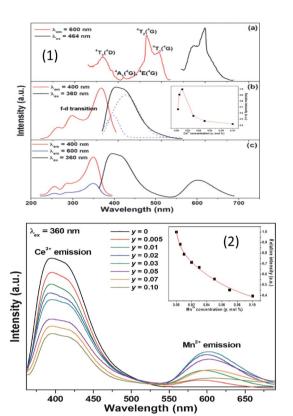


Fig. 16 (1) Emission spectrum of NaCa $_{0.992-y}$ BO $_3$:0.01Ce $^{3+}$,yMn $^{2+}$, and the function between the emission intensity of Ce $^{3+}$ and Mn $^{2+}$ concentration. (2) The emission spectra of NaCa $_{0.992-y}$ BO $_3$:0.01-Ce $^{3+}$,yMn $^{2+}$ phosphors.

between the sensitizer Ce^{3+} and the activator Mn^{2+} , the emission color is adjustable, which can be turned from blue to white and then to orange light. With the increase in Mn^{2+} concentration from 0.005 to 0.10, the emission peak of the Mn^{2+} ions shift toward the long wavelength range from 595 to 610 nm, which can be put down to the change in crystal field strength, as shown in Fig. 16. The relevant optical transitions and energy transfer processes are also demonstrated according to the schematic energy level diagram. The optical properties of a well-packaged WLED lamp by combining the selected $NaCa_{0.96}BO_3:0.01Ce^{3+},0.03Mn^{2+}$ sample, gave a CCT value of 4046 K, CRI value of 90.7, and CIE chromaticity coordinates of (0.326, 0.274), which is superior to those $((0.292, 0.325), R_a = 75, TC = 7756)$ of an InGaN-based WLED relying on a blue LED chip coated with YAG:Ce.⁸⁹

3.4 The co-doping of Eu²⁺/Mn²⁺

Generally speaking, Eu²⁺ presents yellow and blue light, which causes the obtained white light to lack red component, leading to the lower CRI. To increase the CRI of the material, the emission intensity of red light must be improved. Since Mn²⁺ can emit red light, co-doping Eu²⁺ and Mn²⁺ was attempted. For another reason, the emission color of Eu²⁺ singly-doped phosphors can be varied from blue to red in different hosts, nevertheless, the emission color is commonly unchangeable or slightly changeable in certain systems. In order to obtain tuned

emission colors and substantial colors, the energy properties from Eu²⁺ to Mn²⁺ are frequently exploited in phosphors.

It has been proved that there is energy transfer between Eu $^{2+}$ and Mn $^{2+}$, where Eu $^{2+}$ acts as a sensitizer transferring part of its excitation energy into neighboring Mn $^{2+}$ ions. There are many Eu $^{2+}$ /Mn $^{2+}$ co-doped phosphors, such as Sr₂MgSiO₅:-Eu $^{2+}$,Mn $^{2+}$,90 Sr₃La(PO₄)₃,91 NaScSi₂O₆,92 and NaCaBO₃ (ref. 93) all of which have the properties of spectral adjustment, due to the energy transfer between Eu $^{2+}$ and Mn $^{2+}$.

Taking Ca₉Mg(PO₄)₆F₂:Eu²⁺,Mn²⁺ as an example, the excitation emission spectra of Ca₉Mg(PO₄)₆F₂:0.18Eu²⁺ Ca₉Mg(PO₄)₆F₂:0.18Mn²⁺ samples are shown in Fig. 15, respectively. It was observed that there is an overlap of Mn²⁺ excitation and Eu2+ emission in Fig. 17, which proves the possibility for resonance type energy transfer from Eu²⁺ to Mn²⁺ in the host. Except for the intensity in the Ca₉Mg(PO₄)₆F₂:0.18Eu²⁺,0.18Mn²⁺ phosphor from Fig. 17, it can be further demonstrated by the similar excitation spectra monitored at 454 and 565 nm. Due to the emission spectrum covering the emission band of Eu²⁺ and Mn²⁺, the tunable color can be gained *via* adjusting the Mn²⁺ concentration. Moreover, the lifetime of Eu2+ decreases with increasing Mn²⁺, which provides further confirmation of energy transfer from Eu²⁺ to Mn²⁺ ions. Fig. 17 shows the tunable emission color via changing the ratio of Eu2+ and Mn2+ doping concentration. The variation of the CIE chromaticity coordinates from blue to yellow with increasing Mn²⁺ doping concentration from 0 to 0.38 is shown in Fig. 17. Eu2+,Mn2+-activated CMPF phosphors can be potentially applied in UV-pumped white LEDs with outstanding properties.94

Liu *et al.* have synthesized single-phase white light-emitting KCaY(PO₄)₂:Eu²⁺,Mn²⁺ for light-emitting diode (LED) applications. The absolute quantum efficiency of photoconversion was calculated using the following equation:

$$A = \frac{L_{\rm o}(\lambda) - L_{\rm i}(\lambda)}{L_{\rm o}(\lambda)} \tag{24}$$

where $L_{\rm o}(\lambda)$ is the integrated excitation profile when the sample is diffusely illuminated by the integrated sphere's surface; $L_{\rm i}(\lambda)$ is excited by the incident beam. Furthermore, the quantum efficiency (Φ) of KCaY(PO₄)₂:1%Eu²⁺ phosphors can be calculated by

$$\Phi = \frac{E_{\rm i}(\lambda) - (1 - A)E_{\rm o}(\lambda)}{L_{\rm e}(\lambda)A}$$
 (25)

where $E_i(\lambda)$ is the integrated luminescence of the powder upon direct excitation and $E_o(\lambda)$ is the integrated. The internal quantum efficiency of KCaY(PO₄)₂:1%Eu²⁺ was found to be 35.8% and the corresponding external quantum efficiency was 21.4% at the excitation wavelength of 365 nm. Fig. 16 shows the emission spectra of KCaY(PO₄)₂:1%Eu²⁺,x%Mn²⁺ phosphors (x = 0, 1, 2, 4, 5, 7, and 10) excited by 365 nm, and the energy transfer efficiency is shown in the inset in Fig. 18. The phosphors generated red and blue emission bands, centered at 652 nm (${}^4T_1({}^4G) \rightarrow {}^6A_1({}^6S)$ transition of Mn²⁺) and 480 nm (4f⁶5d¹ \rightarrow 4f⁷ transition of Eu²⁺). The intensity of the Eu²⁺ at 480 nm decreased as the Mn²⁺ content increased to x. The intensity at 652 nm increased as the Mn²⁺ content increased,

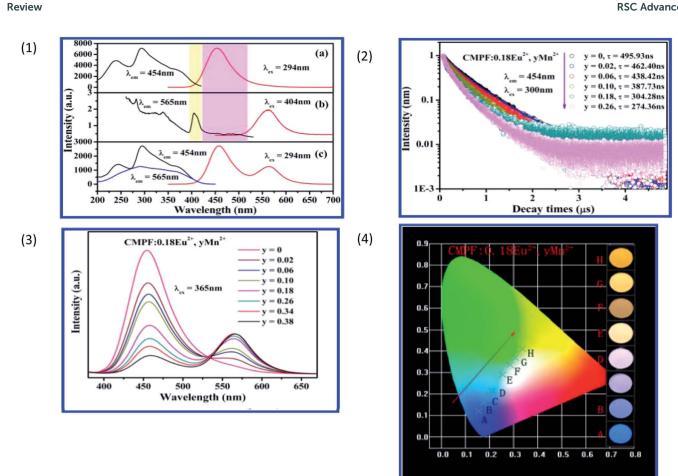


Fig. 17 (1) Excitation and emission spectra of CMPF:0.18Eu²⁺ (a), CMPF:0.18Mn²⁺ (b), and CMPF:0.18Eu²⁺,0.18Mn²⁺ (c) phosphors. The corresponding spectral overlap is also presented in the colored area. (2) Decay curves and lifetimes of Eu²⁺ in representative samples of CMPF: $0.18Eu^{2+}$, vMn^{2+} (monitored at 454 nm and excited at 300 nm). (3) Variation of emission intensity for CMPF: $0.18Eu^{2+}$, vMn^{2+} (v=0-0.38) phosphors on Mn²⁺ doping content excited at 365 nm. (4) CIE chromaticity coordinates of CMPF:0.18Eu²⁺,yMn²⁺ samples (A–H corresponds to y = 0, 0.02, 0.06, 0.10, 0.18, 0.26, 0.34, respectively). The luminescence photographs of corresponding phosphors excited under a 365 nm UV lamp are shown on the right of the picture.94

reaching a maximum at x = 7 mol%, and then decreased when xexceeded 7 mol%, due to the concentration quenching effect. It was demonstrated that the mechanism of energy transfer is an electric dipole-quadrupole interaction from sensitizer Eu²⁺ to an activator Mn2+. Fig. 16 provides a summary of the CIE chromaticity corresponding to the emission spectra excited by 380 nm, which indicates that changing the Mn²⁺ concentration can tune the color from blue through white light and eventually to red in the visible spectral region. The electroluminescence spectrum of white LED lamps fabricated using a NUV 405 nm chip combined with a single-phase white-emitting phosphor KCaY(PO₄)₂:1% Eu²⁺,4%Mn²⁺ driven by a 350 mA current is shown in Fig. 18. The correlated color temperature is 6507 K and CIE color coordinates are (0.314, 0.329), which indicate that the composition-optimized KCaY(PO₄)₂:1%Eu²⁺,4%Mn²⁺ phosphor may have promising applications for white-light NUV LEDs.95

Summary and outlook

Due to the characteristics of the transition metal electronic orbitals, Mn has many valence states, with Mn2+ and Mn4+

being the most common. As activated ions, they have a wide range of applications. Mn⁴⁺ doped materials can work as a complement to the commercial phosphor powder "blue chips + Y₃Al₅O₁₂:Ce³⁺ (YAG:Ce)" to make up for the lack of red component, in order to improve the color rendering indexes and the correlated color temperature. They have the advantages of a simple preparation route and low-cost. Mn⁴⁺ activated oxide and fluoride red phosphors have attracted more and more attention, but low resistance to moisture and unclear structure-related luminescence mechanisms limit their rapid progress in scientific investigation and application of WLED lighting.

Ce³⁺/Mn²⁺,Eu²⁺/Mn²⁺ co-doped materials can produce color tunable phosphors in order to achieve UV excitable single mechanism phosphors to produce white light. In recent years, a single host material emitting white light has made significant progress. In order to further improve its performance, some more studies need to be conducted. Double-doped materials have some deficiencies in the spectral distribution; for example, Eu²⁺/Mn²⁺ or Ce³⁺/Mn²⁺ doped systems in the green emission region are relatively weak, which is certain

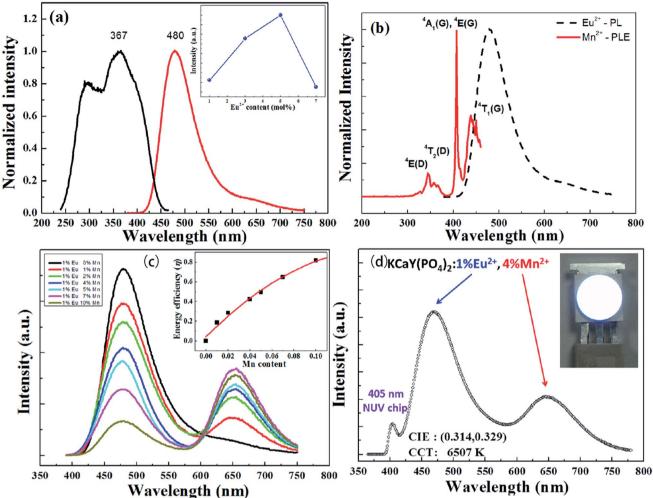


Fig. 18 (a) Photoluminescence excitation/photoluminescence (PLE/PL) spectra of the as-synthesized KCaY(PO₄)₂:1%Eu²⁺ phosphor. The inset shows the PL intensity of KCaY(PO₄)₂:Eu²⁺ phosphors as a function of Eu²⁺ concentration. (b) PLE spectrum of KCaY(PO₄)₂:1%Mn²⁺ and PL spectrum of KCaY(PO₄)₂:1%Eu²⁺. (c) PL spectra of a series of KCaY(PO₄)₂:1%Eu²⁺,x%Mn²⁺ phosphors with different Mn²⁺ concentrations (x = 0, 1, 2, 4, 5, 7, 7), and 10 mol%), excited at 365 nm. Inset shows the energy transfer efficiency from Eu²⁺ to Mn²⁺, as a function of Mn²⁺ content. (d) Electroluminescence spectra of white LED lamps fabricated using a NUV 405 nm chip combined with a white-emitting KCaY(PO₄)₂:1%Eu²⁺,4% Mn²⁺ phosphor driven by a current of 350 mA. Inset shows a photo of the LED package. 95

to affect the luminescence. Further research is being conducted on how to improve the spectral distribution of dual-doping systems, and how to meet the needs of different white phosphors.

In addition, due to the splitting of the Mn²⁺ energy level, the emission spectrum can be red shifted to produce deep red and near infrared light. It has been widely used in the fields of optical fiber communication, solid state lasers, light emitting markers and fluorescence immunoassays.

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References

- H. M. Zhang, H. R. Zhang, Y. L. Liu, J. K. Deng, B. F. Lei,
 L. h. Liu, H. Y. Luo and X. Bai, J. Alloys Compd., 2016, 688,
 1225–1232.
- 2 K. C. Mondal and J. Manam, *J. Mol. Struct.*, 2016, **1125**, 503–513.
- 3 S. X. Li, L. Wang, Q. Q. Zhu, D. M. Tang, X. J. Liu, G. F. Cheng, L. Lu, T. Takeda, N. Hirosaki, Z. R. Huanga and R. J. Xie, *J. Mater. Chem. C*, 2016, 4, 11219–11230.
- 4 B. Yuan, Y. H. Song, Y. Sheng, K. Y. Zheng, X. Q. Zhou, P. C. Ma, X. C. Xu and H. F. Zou, *J. Solid State Chem.*, 2016, 47(2), 169–177.
- D. C. Huang, Y. F. Zhou, W. T. Xu, K. Wang, Z. G. Liu and M. C. Hong, J. Alloys Compd., 2015, 653, 148–155.

Review

6 D. Q. Chen, W. D. Xiang, X. J. Liang, J. S. Zhong, H. Yu, M. Y. Ding, H. W. Lu and Z. G. Ji, *J. Eur. Ceram. Soc.*, 2015, 35, 859–869.

- 7 W. G. Xiao, X. Zhang, Z. D. Hao, G. H. Pan, Y. S. Luo, L. G. Zhang and J. H. Zhang, *Inorg. Chem.*, 2015, 54(7), 3189-31195.
- 8 L. L. Wei, C. C. Lin, M. H. Fang, M. G. Brik, S. F. Hu, H. Jiao and R. S. Liu, *J. Mater. Chem. C*, 2015, 3, 1655–1660.
- 9 L. Chen, R. H. Liu, W. D. Zhuang, Y. H. Liu, Y. S. Hu, X. F. Zhou and X. L. Ma, J. Alloys Compd., 2015, 627(13), 218–221.
- 10 G. Blasse and B. C. Grabmaier, *Luminescent Materials*, Springer, Verlag, Berlin, Heidelberg, 1994, p. 113.
- 11 S. H. Kwon, B. K. Moon, B. C. Choi, J. H. Jeong and J. H. Kim, *J. Korean Phys. Soc.*, 2016, **68**(2), 363–367.
- 12 Z. Y. Mao, J. J. Chen, J. Li and D. J. Wang, *Chem. Eng. J.*, 2016, **284**, 1003–1007.
- 13 J. Chen, Y. G. Liu, L. F. Mei, Z. Y. Wang, M. H. Fang and Z. H. Huang, *J. Mater. Chem. C*, 2015, 3(21), 5516–5523.
- 14 D. C. Huang, Y. F. Zhou, W. T. Xu, K. Wang, Z. G. Liu and M. C. Hong, J. Alloys Compd., 2015, 625, 148–155.
- 15 K. X. Song, J. X. Zhang, Y. F. Liu, C. H. Zhang, J. Jiang, H. C. Jiang and H. B. Qin, J. Phys. Chem. C, 2015, 119(43), 24558-24563.
- 16 W. Xu, G. Zhu, X. Zhou and Y. Wang, *Dalton Trans.*, 2015, 44(19), 9241–9250.
- 17 H. N. Luitel, R. Chand, T. Watari, T. Torikai and M. Yada, *RSC Adv.*, 2015, 5(22), 17034–17040.
- 18 M. Y. Peng, X. W. Yin, P. A. Tanner, C. Q. Liang, P. F. Li, Q. Y. Zhang and J. R. Qiu, J. Am. Ceram. Soc., 2013, 96(9), 2870–2876.
- 19 S. Sugano, Y. Tanabe and H. Kamimura, *Multiplets of Transition Metal Ions in Crystals*, New brk: Academic Press, 1970.
- 20 R. P. Cao, K. N. Sharafudeen and J. R. Qiu, Spectrochim. Acta, Part A, 2014, 117, 402–405.
- 21 M. G. Brik and A. M. Srivastava, J. Lumin., 2013, 133, 69-72.
- 22 H. M. Zhu, C. C. Lin, W. Q. Luo, S. T. Shu, Z. G. Liu, Y. S. Liu, J. T. Kong, E. Ma, Y. G. Cao, R.-S. Liu and X. Y. Chen, *Nat. Commun.*, 2014, 5, 1–10.
- 23 D. Q. Chen, Y. Zhou, W. Xu, J. S. Zhong, Z. G. Ji and W. D. Xiang, *J. Mater. Chem. C*, 2016, 4, 1704–1712.
- 24 M. G. Brik, S. J. Camardello and A. M. Srivastava, *ECS J. Solid State Sci. Technol.*, 2015, 4, R39–R43.
- 25 T. Takahashi and S. Adachi, *J. Electrochem. Soc.*, 2008, **155**, 183–188.
- 26 A. M. Srivastava, Opt. Mater., 2009, 31, 881-885.
- 27 M. G. Brik, A. M. Srivastava and N. M. Avram, *J. Lumin.*, 2011, **131**, 54–58.
- 28 M. Y. Peng, X. W. Yin, P. A. Tanner, M. G. Brik and P. F. Li, *J. Am. Chem. Soc.*, 2015, 27, 2938–2945.
- 29 P. F. Li, L. Wondraczek, M. Y. Peng and Q. Y. Zhang, *J. Am. Ceram. Soc.*, 2016, **99**(10), 3376–3381.
- 30 W. Li, H. R. Zhang, S. Chen, Y. L Liu, J. L Zhuang and B. F. Lei, *Adv. Opt. Mater.*, 2016, **4**(3), 427–434.
- 31 X. Ding, G. Zhu, W. Y. Geng, Q. Wang and Y. H. Wang, *Inorg. Chem.*, 2016, 55(1), 154–162.

- 32 K. Li, D. Zhu and R. V. Deun, Dyes Pigm., 2017, 142, 69-76.
- 33 Y. X. Pan and G. K. Liu, J. Lumin., 2011, 131(3), 465-468.
- 34 M. G. Brik, Y. X. Pan and G. K. Liu, J. Alloys Compd., 2011, 509(5), 1452–1456.
- 35 X. L. Gao, Y. Song, G. X. Liu, X. T. Dong, J. X. Wang and W. S. Yu, *Dalton Trans.*, 2016, 45(44), 17886–17895.
- 36 T. C. Lang, T. Han, L. L. Peng and M. J. Tu, *Mater. Chem. Front.*, 2017, 1, 928–932.
- 37 T. T. Deng, E. H. Song, J. Su, L. Y. Wang, Y. Deng, S. Ye, J. Wang and Q. Zhang, J. Mater. Chem. C, 2017, 5, 2910–2918.
- 38 R. Cao, M. Peng, E. Song and J. Qiu, *ECS J. Solid State Sci. Technol.*, 2012, **1**(4), R123–R126.
- 39 Z. liu, M. Yuwen, J. Liu, C. Yu, T. Xuan and H. Li, *Ceram. Int.*, 2017, 43(7), 5674–5679.
- 40 R. P. Cao, D. Ceng, X. G. Yu, S. L. Guo, Y. F. Wen and G. T. Zheng, *Funct. Mater. Lett.*, 2015, **8**(5), 1550046.
- 41 P. F. Li, L. Wondraczek, M. Y. Peng and Q. Y. Zhang, *J. Am. Ceram. Soc.*, 2016, **99**(10), 3376–3381.
- 42 E. Song, J. Wang, J. Shi, T. Deng, S. Ye, M. Peng, J. Wang, L. Wondraczek and Q. Zhang, *ACS Appl. Mater. Interfaces*, 2017, 9(10), 8805–8812.
- 43 A. Abdukayum, J. T. Chen, Q. Zhao and X. P. Yan, *J. Am. Chem. Soc.*, 2013, **135**(38), 14125–14133.
- 44 Q. H. Zhang, J. Wang, G. G. Zhang and Q. Su, *J. Mater. Chem.*, 2009, **19**, 7088–7092.
- 45 Y. X. Zhuang, Y. Katayama, J. Ueda and S. Tanabe, *Opt. Mater.*, 2014, **36**, 1907–1912.
- 46 J. Yan, M. G. Brik, C. M. Liu, D. J. Hou, W. J. Zhou, B. B. Zhang, Y. Huang, Y. Tao and H. B. Liang, *Opt. Mater.*, 2015, 43, 59-65.
- 47 C. J. Duan, A. C. A. Delsing and H. T. Hintzen, *Chem. Mater.*, 2009, **21**, 1010–1016.
- 48 S. K. Singh, RSC Adv., 2014, 4(102), 58674-58698.
- 49 X. R. Zhong, X. P. Meng, M. Y. Li, X. Q. Wang, X. J. Wen and X. Y. Zhang, *Chem. Phys. Lett.*, 2012, 536, 55–60.
- 50 C. M. Abreu, R. S. Silva, M. E. G. Valerio and Z. S. Macedo, J. Solid State Chem., 2013, 200, 54–59.
- 51 Y. Katayama, J. Ueda and S. Tanabe, *Opt. Mater. Express*, 2014, **4**, 613–623.
- 52 J. G. Cheng, P. L. Li, Y. S. Sun, Q. Y. Bai, Z. L. Li, M. M. Tian, C. Wang and Z. P. Yang, J. Mater. Chem. C, 2017, 5, 127–133.
- 53 A. Bessiére, A. Lecointre, R. A. Benhamou, E. Suard, G. Wallez and B. Viana, *J. Mater. Chem. C*, 2013, 1(6), 1252– 1259.
- 54 A. Bessière, R. A. Benhamou, G. Wallez, A. Lecointre and B. Viana, *Acta Mater.*, 2012, **60**, 6641–6649.
- 55 Y. Y. Ma, J. Q. Hu, E. H. Song, S. Ye and Q. Zhang, *J. Mater. Chem. C*, 2015, 3, 12443–12449.
- 56 L. Wu, B. Wang, Y. Zhang, L. Li, H. R. Wang, H. Yi, Y. F. Kong and J. J. Xu, *Dalton Trans.*, 2014, 43, 13845–13851.
- 57 B. S. Choia, O. G. Jeonga, J. C. Parka, J. W. Kimb, S. J. Leeb, J. H. Ryuc, J. Leec and H. Chob, *J. Ceram. Process. Res.*, 2016, 17(7), 778–781.
- 58 L. Hu, Q. Wang, X. Wang, Y. Li, Y. Wang and X. Peng, *RSC Adv.*, 2015, 5, 104708–104714.
- 59 X. Ding, G. Zhu, Q. Wang and Y. Wang, RSC Adv., 2015, 5, 30001–33004.

60 Y. Pan, L. Li, J. Lu, R. Pang, L. Wan and S. Huang, Dalton

Trans., 2016, 45, 9506–9512.

RSC Advances

- 61 X. Zhu, Y. Yao and Z. Zhou, Opt. Mater., 2016, 62, 104-109.
- 62 R. Cao, D. Peng, H. Xu, S. Jiang, Z. Luo, H. Ao and P. Liu, *J. Lumin.*, 2016, **178**, 388–391.
- 63 K. W. Park, H. S. Lim, S. W. Park, G. Deressa and J. S. Kim, Chem. Phys. Lett., 2015, 636, 141–145.
- 64 B. Chandra Babu, B. Vengla Rao, M. Ravi and S. Babu, J. Mol. Struct., 2017, 1127, 6–14.
- 65 K. Omri and L. El Mir, J. Mater. Sci.: Mater. Electron., 2016, 27(9), 9476-9482.
- 66 E. H. Song, S. Ding, M. Wu, S. Ye, F. Xiao, S. F. Zhou and Q. Y. Zhang, *Adv. Opt. Mater.*, 2014, 2, 670–678.
- 67 C. R. Ronda and T. Amrein, J. Lumin., 1996, 69, 245-248.
- 68 M. S. Kwon, J. H. Jordahl, A. W. Phillips, K. Chung, S. Lee, J. Gierschner, J. Lahann and J. Kim, *Chem. Sci.*, 2016, 7(3), 2359–2363.
- 69 M. Y. Peng, X. W. Yin, P. A. Tanner, C. Q. Liang, P. F. Li, Q. Y. Zhang and J. R. Qiu, J. Am. Ceram. Soc., 2013, 96(9), 2870–2876.
- 70 J. Lü, F. Du, R. Zhu, Y. Huang and H. J. Seo, *J. Mater. Chem.*, 2011, 21, 16398–16405.
- 71 M. Szumera, I. Wacławska and J. Sułowska, *J. Therm. Anal. Calorim.*, 2016, **123**(2), 1083–1089.
- 72 C. Wang, P. L. Li, Z. J. Wang, Y. S. Sun, J. G. Cheng, Z. L. Li, M. M. Tian and Z. P. Yang, *Phys. Chem. Chem. Phys.*, 2016, 18, 28661–28673.
- 73 Z. Zhang and W. Tang, J. Alloys Compd., 2015, 663, 731-737.
- 74 J. Zhou, T. Wang, X. Yu, D. Zhou and J. Qiu, *Mater. Res. Bull.*, 2016, 73, 1–5.
- 75 B. P. Kore, S. Tamboli, N. S. Dhoble, A. K. Sinha, M. N. Singh, S. J. Dhoble and H. C. Swart, *Mater. Chem. Phys.*, 2016, **187**, 233–244.
- 76 C. Liang, H. P. You, Y. B. Fu, X. M. Teng, K. Liu and J. H. He, *Optik*, 2017, **131**, 335–342.
- 77 Y. Hong, J. L. Chen, Y. Pu, T. J. Zhang and S. C. Gan, *J. Rare Earths*, 2015, 33(4), 366–370.

- 78 Y. Zhu, Y. Liang, S. Liu, K. Li, X. Wu and R. Xu, *J. Rare Earths*, 2017, 35(1), 41–46.
- 79 U. Manik, S. C. Gedamb and S. J. Dhoble, *Luminescence*, 2014, **30**(6), 910-913.
- 80 R. L. Zheng, D. W. Luo, Y. Yuan, Z. Y. Wang, Y. Zhang, W. Wei, L. B. Kong and D. Y. Tang, *J. Am. Ceram. Soc.*, 2015, 98(10), 3231–3235.
- 81 J. Zhang, F. Zhang and L. L. Han, *J. Rare Earths*, 2015, 33(8), 820–824.
- 82 L. R. Hatwar, S. P. Wankhede, S. V. Moharil, P. L. Muthald and S. M. Dhopte, *Luminescence*, 2015, 30(6), 904-909.
- 83 Y. Q. Li, H. P. Ma, Y. J. Hua, Q. H. Yang, C. X. Li, D. G. Deng and S. Q. Xu, *J. Rare Earths*, 2016, 34(1), 7–11.
- 84 P. L. Li, Z. J. Wang, Q. L. Guo and Z. P. Yang, J. Am. Ceram. Soc., 2014, 98(2), 495–500.
- 85 F. Zhang and W. J. Tang, Luminescence, 2014, 30(2), 216-220.
- 86 G. G. Li, D. L. Geng, M. M. Shang, C. Peng, Z. Y. Cheng and J. Lin, *J. Mater. Chem.*, 2011, 21, 13334–13344.
- 87 Y. Kim, K. Page, A. Limarga, D. Clarke and R. Seshadri, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2007, **76**, 115204.
- 88 J. F. Sun, W. L. Zhang, Y. M. Shi, D. Z. Shen and J. Y. Sun, *J. Electrochem. Soc.*, 2012, **159**, J5–J12.
- 89 J. F. Sun, Z. P. Lian, G. Q. Shen and D. Z. Shen, *RSC Adv.*, 2013, 3, 18395–18405.
- 90 Z. Chen, J. Zhang, S. Chen, M. Lin, C. He, G. Xu, M. Wang, X. Yu, J. Zou and K. Guo, *J. Alloys Compd.*, 2015, 632, 756–759.
- 91 Z. J. Wang, S. Q. Lou, P. L. Li and Z. G. Lian, *Opt. Mater. Express*, 2017, **56**(4), 1167–1172.
- 92 J. Barzowska, Z. G. Xia, D. Jankowski, D. Włodarczyk, K. Szczodrowski, C. G. Ma, M. G. Brik, Y. Zhydachevskii and A. Suchocki, RSC Adv., 2017, 7, 275–284.
- 93 Y. M. Wang, H. B. Zhang, Q. L. Wei, C. H. Sua and D. Zhang, *Ceram. Int.*, 2016, 42(10), 12422–12426.
- 94 K. Li, D. L. Geng, M. M. Shang, Y. Zhang, H. Z. Lian and J. Lin, J. Phys. Chem. C, 2014, 118, 11026–11034.
- 95 W. R. Liu, C. H. Huang, C. W. Yeh, J. C. Tsai, Y. C. Chiu, Y. T. Yeh and R. S. Liu, *Inorg. Chem.*, 2012, **51**, 9636–9641.