

**Self-scrolling MoS₂ Metallic Wires**

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Self-scrolling MoS₂ Metallic Wires

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Two dimensional (2D) van der Waals (vdW) materials with strong in-plane chemical bonds and weak interaction in the out-of-plane direction have been acknowledged as a basic building block for designing dimensional materials in 0D, 1D, 2D and 3D forms. Compared to the explosive research on 2D vdW materials, quasi-one-dimensional (quasi-1D) vdW materials have received rare attention, despite they also present rich physics in electronics and engineering implications. Herein, quasi-1D MoS₂ nanoscrolls are directly fabricated from CVD-grown 2D triangular MoS₂ sheet. The formation, stability and electronic property of quasi-1D MoS₂ nanoscrolls are studied experimentally and theoretically. The formation of nanoscroll always starts from the edge of triangular MoS₂ sheet along its armchair orientation. The electronic properties of MoS₂ nanoscrolls are systemically studied with optical spectroscopy, electrical transport together with density-functional theory (DFT) calculations. Surprisingly, the carrier mobility and contact property of MoS₂ nanoscroll based field effect transistors (FETs) distinct from that of 2D MoS₂ sheet. The transition from 2D semiconductor MoS₂ sheet to 1D metallic MoS₂ nanoscroll is successfully achieved. It is expected that this method of fabricating MoS₂ nanoscrolls will attract wide interest for 1D transition metal dichalcogenides with novel physical and chemical properties.

Introduction

The Beyond graphene, 2D layered vdW materials have been recognized as a basic building block for nanomaterials with all other dimensionalities. For example 2D vdW materials can be wrapped up into 0D, rolled into 1D or stacked into 3D materials.¹⁻³ With the change of dimensionalities, vdW materials have shown unique properties include tunability of energy structure (from indirect to direct bandgap) and variations of many-body quantum criticalities.⁴⁻⁷

As a representative member of vdW family, molybdenum disulfide (MoS₂) displays excellent properties when its dimensionality varies from 0D to 3D. 0D MoS₂ i.e. quantum dots, exhibits special electronic and photo-physical properties due to the quantum confinement and edge effects which enable its application in bio-imaging, luminescent sensing and catalysts.⁸⁻¹¹ While, the high carrier mobility and reasonable bandgap of 2D MoS₂ sheet together with its compatible fabrication make it promising candidate for next diverse electronics in the semiconductor

industry.^{12, 13} In addition, the 3D MoS₂ bulk materials with small friction coefficient makes its excellent application in lubrication.¹⁴⁻¹⁶ However, 1D MoS₂ has received considerably less attention even it is expect to inherit excellent property.¹⁷ In 2001, 1D MoS₂ nanotube has been first synthesized under ultrahigh temperature with long time by bottom-to-up method, but the poor morphology and quality limited its application.^{18, 19} 1D MoS₂ nanoscrolls, as another prototypical structure of 1D MoS₂, is a spirally wrapped 2D MoS₂ sheet.²⁰ Recently, 1D MoS₂ nanoscroll was prepared by plasma treatment and chemical route, however, the poor morphology, atomic defect and unexpected contamination degrade its quality and further limit its application.²¹⁻²⁴ Up to now, all studies are limited on synthesis. The mechanical strength has been only studied by theoretical simulation,²⁵ but the electronic property of 1D MoS₂ nanoscrolls from experimental and theoretical investigation has never been reported impeding by the experimental difficulties in fabrication, isolation, and device integration.

In this study, we introduce a simple and efficient method to prepare 1D MoS₂ nanoscroll directly from chemical vapor deposition (CVD) grown 2D MoS₂ sheet at room temperature. Combining the experiments and theory calculations, the formation of 1D MoS₂ nanoscroll is investigated in detail. A major finding from our experiment is that the 2D MoS₂ sheet always scrolls from its Mo-terminated edge with an armchair orientation. Molecular dynamics simulation and DFT computation provide a sensible explanation to this experimental phenomenon. The band structure of 1D MoS₂ nanoscrolls with different layers of wall is computed, illustrating its unique electronic property. In particular, 1D MoS₂ nanoscroll with three-layer wall exhibits metallic-like

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behavior. The spectroscopic properties of 1D MoS₂ nanoscroll have been also measured using X-ray photoemission spectroscopy and Raman microscopy. Importantly, the electronic transport property of 1D MoS₂ nanoscroll is firstly studied, including electric field effect, carrier mobility and contact property, which show a distinct transport behavior from that of 2D MoS₂ sheet.

Results and discussion

Triangular 2D MoS₂ sheet is first grown on SiO₂/Si by CVD, where MoO₃ and sulfur are employed as the sources.^{26, 27} Next, 2D MoS₂ sheet is transferred onto a new SiO₂/Si substrate as shown in Fig. 1a. By dropping ~100 μ L isopropyl alcohol (IPA) solution on the surface, the adsorption interaction between MoS₂ sheet and substrate is reduced by inserted IPA solution,²⁸ and then MoS₂ has turned from 2D sheet to quasi-1D nanoscroll structure at room temperature within several minutes. Fig. 1b shows the typical partially scrolled MoS₂ nanoscroll/sheet, and the inset shows the fully scrolled MoS₂ nanoscroll. The length of MoS₂ nanoscroll is defined by the size of 2D MoS₂ sheet itself. The morphology and uniformity is much more perfect than those of previously reported MoS₂ nanoscroll.²¹⁻²⁴ Interestingly, the scrolling directions (black arrow indicated) always keep perpendicular to MoS₂ edge, and the scrolling always starts from its edge rather than corner. This indicates that the scrolling is driven by internal vdW force

rather than external IPA evaporation because the MoS₂ edge is subjected to stronger internal vdW force than the MoS₂ corner. A schematic of MoS₂ atomic model (inset in Fig. 1a) illustrates the Mo-terminated edges of CVD-grown 2D MoS₂ sheet.^{29, 30} Thus, it suggests that the scrolling orientation (black arrow indicated in Fig. 1b) is parallel to armchair direction shown the dashed black arrow in Fig. 1b. The as-grown 2D MoS₂ sheet has also been scrolled as shown in Fig. S1, where one can see that a large amount of short and small MoS₂ nanoscrolls are formed. This is caused by the strong interaction between as-grown MoS₂ and SiO₂ substrate, which makes MoS₂ sheet breaks up into many fragments during scrolling. The scrolling orientation is summarized in Fig. S1. Interestingly, the scrolling orientation is also near perpendicular to MoS₂ edge with period about 60°, which indicates that the scrolling orientation is independent to the size of MoS₂ nanoscrolls. To gain molecular-level insight into the dominant tendency for MoS₂ sheet scrolling via the armchair orientation, we performed a classical MD simulation. As shown in Fig. 1c and 1d, the energy per atom for the nanoscroll scrolled via the armchair orientation is lower than that via the zigzag and other chirality (i.e., (300, 150)) orientation, consistent with the experiment.

Moreover, a MD simulation of the self-scrolling process in real time is performed for two-layer MoS₂ sheet with partially scrolled structure as the initial time (0 ps) (see Fig. 2a; more detailed snapshots are displayed in Fig. S2, one more shorter MoS₂ sheet is also investigated as shown in Fig. S3).³¹ During the ensuring scrolling, the left portion of MoS₂ sheet moves closer to the nanoscroll at ~ 25 ps, and then begins to wrap onto the existing partial MoS₂ nanoscroll at ~ 100 ps (see Fig. S2). This attaching-wrapping process continues until the whole MoS₂ sheet turns into a full nanoscroll at ~ 200 ps (Supporting Video S1). Fig. 2b presents the energy per atom of the system

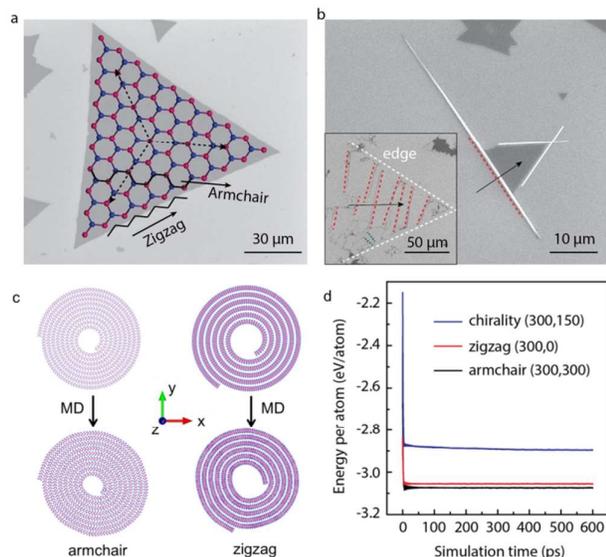


Fig. 1. The formation of MoS₂ nanoscrolls. a) and b) SEM images of MoS₂ sheet and nanoscrolls on SiO₂ substrate. The inset in b) shows the low magnification SEM image. c) The same-sized MoS₂ sheet scrolled via the armchair (left panels) and zigzag (right panels) orientation before (upper panels) and after (lower panels) the structure relaxation at 300 K and 1 bar through classical MD simulation. d) The computed energy per atom of the system versus simulation time for MoS₂ nanoscrolls with armchair, zigzag and chirality (300, 150) orientation, respectively.

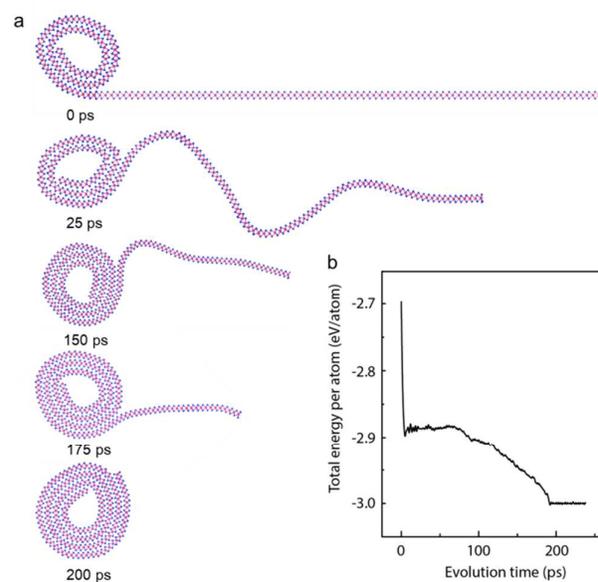


Fig. 2. a) The snapshots of MoS₂ sheet self-scrolling at time 0, 25, 150, 175 and 200 ps; b) the energy of the system versus MD simulation time during the self-scrolling at 300 K.

as a function of time for the self-scrolling process. Clearly, beyond 200 ps the MoS₂ nanoscroll reaches a relatively stable state without showing much energy variation. This MD simulation demonstrates that the self-scrolling process can be a very fast process the formation of the initial partial nanoscrolls.

To see clearer structural details of the fabricated nanoscrolls, MoS₂ nanoscrolls are transferred on grid for high-resolution transmission electron microscope (HR-TEM) characterization. The HR-TEM image (Fig. 3a) shows that it possesses a tube-like structure with a hollow core surrounded by thick walls consisting of MoS₂ multilayer. The HR-TEM image clearly shows the MoS₂ layers are uniformly and compactly stacked. The interlayer distance between adjacent MoS₂ layers is summarized in Fig. 3b (other HRTEM images are shown in Fig. S4), which ranges from 5.5 Å to 6.5 Å, consistent with that of MoS₂ bulk material and MoS₂ nanotubes.^{32, 33} Some interlayer distance is as large as 10 nm (Fig. S5), which should be caused by the uncontrollably scrolling speed. In addition, too large size of MoS₂ sheet will also render to form uniformly nanoscrolls due to the non-uniformly strain during scrolling. To confirm the stability of different structures, MD simulation of MoS₂ nanoscroll with interlayer distance ranging from 4 Å to 10 Å are performed as shown in Fig. 3c and 3d. MoS₂ nanoscroll tends to form highly defected structure or collapsed structure if the interlayer distance is smaller than 4.5 Å or larger than 6.5 Å, respectively. If the interlayer distance increases from 4 Å to 10 Å, the energy per atom (Fig. 3d) will decrease from -2.97 eV to -3.07 eV and then increase to -3.01 eV. The energy per atom as a function of interlayer distance clearly demonstrates that the system is at the energetically favourable state when the

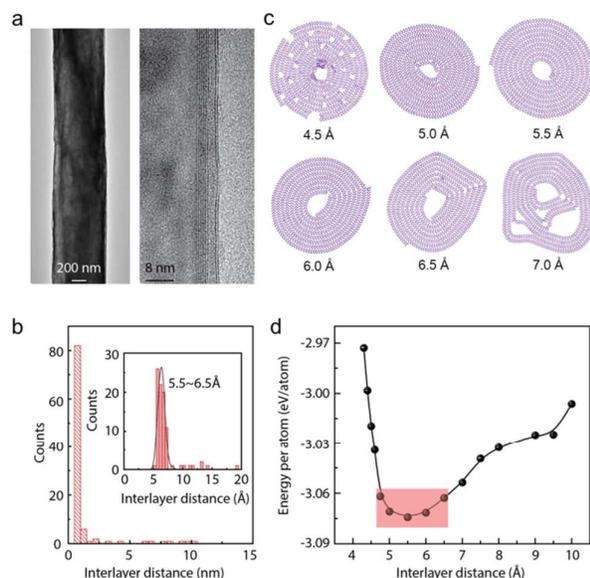


Fig. 3. The stability of MoS₂ nanoscrolls. a) The TEM and HR-TEM images of MoS₂ nanoscroll. b) The histogram of interlayer distance of MoS₂ nanoscrolls, the inset shows the more localized histogram. c) The MD simulation of MoS₂ nanoscrolls with different interlayer distance, ranging from 4.5 Å to 7.0 Å. d) The energy per atom as a function of interlayer distance for the MoS₂ nanoscroll. The red rectangle highlights interlayer

distances that give energetically favorable structures with little defects.

interlayer distance is within the range from 4.7 to 6.5 Å, consistent with the experimental observation. When the interlayer distance is beyond 7 Å, the collapsed pore structure will be formed, again accordant with the experimental observations (Fig. S5). This is because either smaller or larger interlayer distance can increase either the repulsive or attractive vdW interaction, thereby causing highly defected structure.

Experimentally, Raman and X-ray photoelectron spectroscopies (XPS) are employed to study the optical property and electronic structure of MoS₂ sheet and nanoscroll, as presented in Fig. 4. The high-resolution Mo 3d and S 2p XPS spectra have been carefully calibrated by C1s peak from surface contamination. As shown, the blue shifts (increase in binding energy) are observed in the binding energy of core levels in MoS₂ after scrolling. After scrolling, the binding energy of Mo 3d_{5/2} shifts from 229.38 eV to 229.84 eV, whereas the peak of S 2p_{3/2} shifts from 162.25 eV to 162.71 eV. Since the Fermi level in material is related to the binding energy,^{26, 34} the blue shifts of binding energy can be interpreted by the move up of the Fermi level in MoS₂ nanoscrolls compared to MoS₂ sheet, suggesting a distinct electronic structure in MoS₂ nanoscrolls. Fig. 4b shows the Raman spectra of MoS₂ sheet and nanoscrolls. It is obvious that E_{2g} and A_{1g} bands red and blue shift, respectively. Furthermore, the position difference between A_{1g} and E_{2g} increases from 19.2 cm⁻¹ to 22.5 cm⁻¹, indicating a strong interlayer interaction which is similar with that of thick MoS₂.³⁵

Because the fabricated MoS₂ nanoscrolls are too large for the first-principles DFT computation, to understand the electronic properties of the nanoscrolls, we computed the band structures of two limiting systems, namely, model multilayered MoS₂ nanosheets and model multi-walled MoS₂ nanotubes (see Fig. S6). Fig. 5a-c show the computed band structure and the corresponding density of states (DOS) for 2D 1 layer (1L)-3 layer (3L) MoS₂ sheet. Clearly, with increasing the number of layers from 1L to 3L, the computed band gap decreases from 1.74 eV to 1.10 eV, a trend consistent with previous results.^{36, 37} On the other hand, for the multi-walled MoS₂ armchair nanotubes with the wall-to-wall distance about 5.5 Å, the computed band structure and corresponding DOS for single-, double- and triple-walled MoS₂ nanotube, respectively, are presented in Fig. 5d-e. Notably, the computed

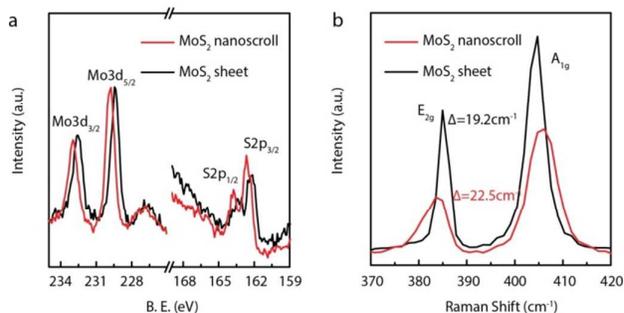


Fig. 4. The spectroscopic characteristic of MoS₂ sheet and nanoscroll. a) the high resolution Mo 3d and S 2p XPS spectra

of MoS₂ sheet and nanoscroll. b) the Raman spectra of MoS₂ sheet and nanoscroll.

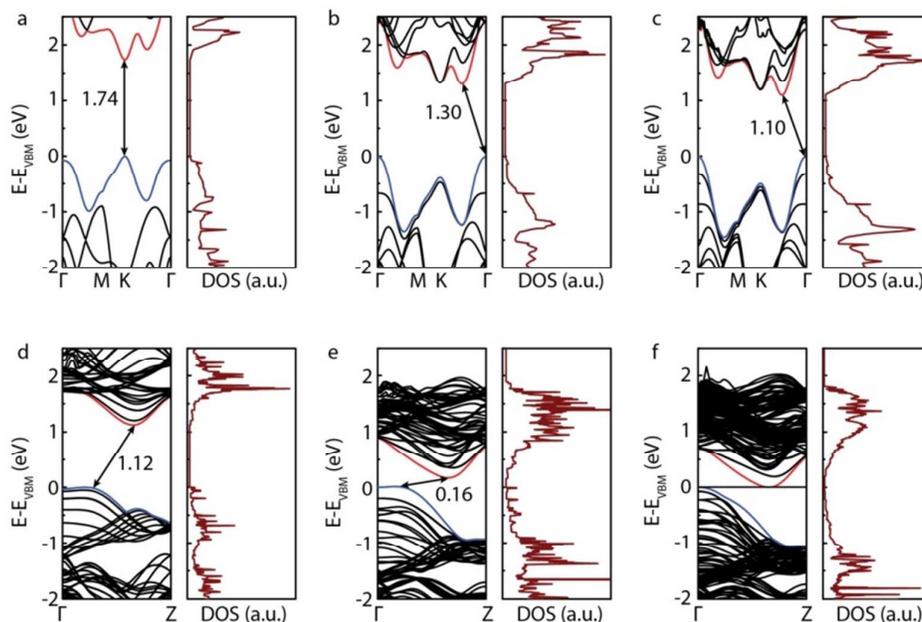


Fig. 5. The band structure of MoS₂ nanostructures. a)-c) The band structure and corresponding DOS distribution of 1L, 2L and 3L MoS₂ nanosheets. d)-f) The band structure and corresponding DOS distribution of MoS₂ nanotube with single-, double- and triple-walled MoS₂ nanotube.

band gap decreases monotonously from 1.12 eV to 0 eV, with thickness increasing from single to triple walled nanotube. In particular, the zero band gap for the triple-walled nanotube indicates that the nanotube already becomes a metal. The two limiting cases indicate that the trend of band-gap change as a function of layer number for 2D MoS₂ sheets is dramatically different from that for 1D MoS₂ nanotubes. The band gap of 2D MoS₂, as tuned by the layer number, is lower bounded by the value 0.88 eV of the bulk MoS₂.³⁸ On the other hand, the trend of band-gap change for 1D MoS₂ nanoscroll is expected to resemble that of 1D MoS₂ nanotubes. Hence, we expect the semiconductor-to-metal transition (or appearance of zero band gap) would occur for 1D MoS₂ nanoscroll at a layer number greater than three.

To explore the electronic transport property, back-gated field effect transistors (FETs) have been fabricated employing MoS₂ sheet and nanoscroll as the channel materials. Ti (10nm)/Au (50nm) as the drain and source electrodes were defined by electron beam lithography, and deposited by electron beam deposition system followed with lift-off. The output curves of MoS₂ nanoscroll FET under V_{GS} varying from -70 V to 70 V is shown in Fig. 6a, which exhibit linearly current-voltage characteristic suggesting an ohmic contact behavior. The transfer curves of MoS₂ nanoscroll under different V_{DS} (Fig. S7) show a well and stable field effect. To compare, the

transfer curves of MoS₂ sheet and nanoscrolls FETs are plotted together in Fig. 6b. As seen, I_{DS} of MoS₂ nanoscroll FET increase with V_{GS} increasing, suggesting that MoS₂ nanoscroll still shows n-type transport characteristic similar with that of MoS₂ sheet FET. However, the current on/off ratio of MoS₂ nanoscroll FET significantly decreases from $\sim 10^4$ to $\sim 10^{-1}$, indicating a difference transport behavior from MoS₂ sheet, which suggest MoS₂ nanoscroll exhibits metallic property consisted with zero bandgap from DFT calculation. Contrary to thick MoS₂ sheet which has on/off ratio of about 10^4 (Fig. S8), the lower on/off ratio ($\sim 10^{-1}$) in MoS₂ nanoscrolls device indicates that the difference transport behavior is originated from its unique structure (continually scrolling structure) rather than the increased thickness.

The carrier mobility in MoS₂ nanoscrolls could be calculated by $\mu_{1D} = (L^2/V_{DS}C_{1D})(dV_{DS}/dV_{GS})$, where L is the channel length. C_{1D} is the capacitance for the case of 1D channel material, which could be calculated by $C_{1D} = (2\pi\epsilon\epsilon_0L)/\cosh^{-1}((r+t_{ox})/r)$ where ϵ is the dielectric constant of SiO₂, t_{ox} is the thickness of SiO₂, r and L are the radius and length of the MoS₂ nanoscroll, respectively.^{28, 39} In the case of MoS₂ sheet, the carrier mobility could be calculated by $\mu_{2D} = (L/WV_{DS}C_{2D})(dV_{DS}/dV_{GS})$, where W is

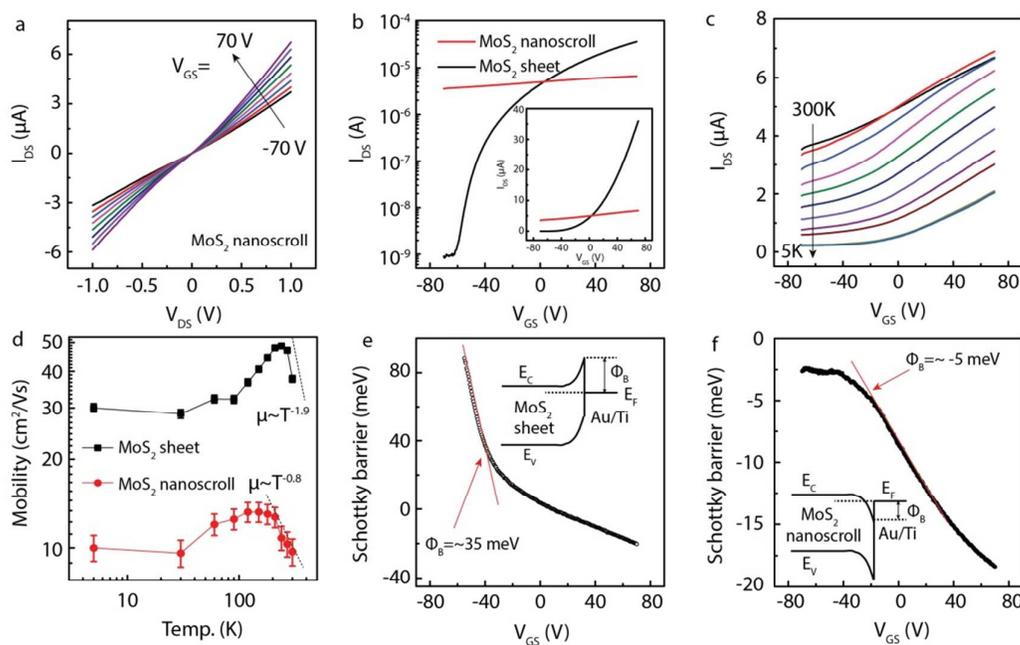


Fig. 6. The electrical transport properties of MoS₂ nanoscrolls. a) The output properties of MoS₂ nanoscroll transistor. b) The transfer curve of MoS₂ sheet and nanoscroll transistors. c) The temperature dependent transfer curves of MoS₂ nanoscroll transistor. d) The mobility as a function of temperature for MoS₂ sheet and nanoscroll. e) and f) The Schottky barrier as a function of gate voltage on MoS₂ sheet and nanoscroll, respectively. The inset shows the corresponded energy diagram structure.

the channel width.^{26, 40} C_{2D} is the capacitance for the case of 2D channel material, which could be calculated by $C_{2D} = \epsilon\epsilon_0/t_{ox}$. According to the estimation, the carrier mobilities are ~ 37.5 cm²/Vs and ~ 9.7 cm²/Vs for MoS₂ sheet and nanoscroll, respectively. The smaller mobility of MoS₂ nanoscrolls is resulted from the much weaker capacitance coupled with the back gated compared to that of flat MoS₂ sheet. Even so, the mobility of MoS₂ nanoscroll is much larger than that of MoS₂ nanotube (0.014 cm²/Vs),⁴¹ possible due to the feature of continuous electron distribution in MoS₂ nanoscrolls. Furthermore, the temperature dependent transfer curves of MoS₂ sheet and nanoscrolls FETs are performed, as shown in Fig. 6c and Fig. S9. The temperature dependent mobilities are plotted in Fig. 6d, where one can clearly see that MoS₂ sheet always has a higher mobility than that of MoS₂ nanoscroll. With temperature cooling, the mobility for MoS₂ sheet and nanoscroll will both increase and then decrease, exhibiting a peak at around ~ 240 K and ~ 180 K, respectively. This phenomenon is related to the dominated carrier scattering mechanism which change from electron-phonon scattering to charged impurity scattering during cooling.⁴² Furthermore,

fitting the plot with $\mu \sim T^{-\gamma}$ in electron-phonon dominated part, one can see that the value of γ for MoS₂ sheet is 1.9 which is similar with previous results.⁴² Significantly, the value of γ for MoS₂ nanoscroll is 0.8, which is much smaller than that of MoS₂ sheet and bulk MoS₂ ($\gamma = 2.6$).^{43, 44} This suggests that the electron-phonon scattering in MoS₂ nanoscroll is weaker than that of MoS₂ sheet even the bulk MoS₂ possibly attributing to its confined structure.

Furthermore, the contact property of MoS₂ nanoscrolls could be studied to explore its energy band structure. To evaluate the Schottky barrier height between metal contact with MoS₂, the temperature dependence of the I_{DS} - V_{GS} curve could be re-plotted with Arrhenius formation with $\ln(I_{DS}/T^2) \sim 1000/T$, as shown in Fig. S10. As seen, the $\ln(I_{DS}/T^2) \sim 1000/T$ has positive correlations for MoS₂ nanoscroll compared with negative correlations for MoS₂ sheet in the temperature range from 300 K to 200 K. Furthermore, the effective Schottky barrier could be obtained from the slope of the Arrhenius plot, as summarized in Fig. 6e and 6f. In the case of MoS₂ sheet (Fig. 6e), the Schottky barrier dramatically decreases from ~ 85 to 0 meV, becomes negative, and finally

saturates near -20 meV with V_{GS} increase. A sudden change with Schottky barrier of 35 meV is observed which should be due to the dominated mechanism switching from thermionic emission to quantum-mechanical tunnelling when the V_{GS} increase.⁴⁵⁻⁴⁷ Therefore, it indicates that the effective Schottky barrier is ~ 35 meV for the particular device, as indicating in Fig. 6e. Different to MoS_2 sheet, the Schottky barrier of MoS_2 nanoscrolls slowly decreases from -2.5 meV to -5 meV, and then rapidly decreases to -17 meV when V_{GS} increase. This indicates that the effective Schottky barrier is ~ -5 meV for the particular device, as indicating in Fig. 6f. According to the Schottky-Mott model,⁴⁸ the Schottky barrier could be estimated based on the work function of the metal contact (Φ_{metal}) relative to the electron affinity of materials ($\chi_{material}$) by $\Phi_B = \Phi_{metal} - \chi_{material}$. Therefore, the negative Schottky barrier suggests that the affinity of MoS_2 nanoscroll is slightly bigger than the work function of metal contact. With the similar metal contact, it can be concluded that the affinity of MoS_2 nanoscroll is ~ 40 meV bigger than that of MoS_2 sheet with $\Delta\chi_{material} = -\Delta\Phi_B$. The energy diagram inset in Fig. 6e and 6f shows the relative location of energy band, where the energy diagram of MoS_2 nanoscroll is quite different from MoS_2 sheet.³³ Hence, we demonstrate that the transition from semiconductor 2D MoS_2 sheet to metallic quasi-1D MoS_2 nanoscroll by self-rolling method.

Experimental

Experimental Section

Synthesis of MoS_2 sheet: The MoS_2 sheet were grown by the chemical vapor deposition (CVD) method. 500 mg sulfur powder was put at the upstream of the furnace with temperature of ~ 200 °C. MoO_3 powder was placed in a quartz boat located in the center of the furnace. A SiO_2/Si substrate was facedown placed on top of the quartz boat. The furnace was first purged with 50 sccm Ar for 30 min and kept the furnace pressure in the range of $1-3$ kPa during the whole growth process, then raised to the growth temperature 850 °C in 30 min. When the furnace reached 750 °C, the temperature of sulfur reached 200 °C. After growing 10 min, the power for heating furnace was terminated, and the system cooled naturally. When the temperature was lower than 600 °C, the power for heating sulfur was switched off.

Transfer of MoS_2 sheet and prepare MoS_2 nanoscrolls: PMMA layer was formed on the surface of $MoS_2/SiO_2/Si$ by spin-coating. The sample was baked at 180 °C for 5 min, and then put into DI water to release PMMA/ MoS_2 from the substrate. After that, a new SiO_2/Si substrate was used to pick up the PMMA/ MoS_2 . After drying in nature, the PMMA was removed by warm acetone for three times. The sample was then heated at 80 °C for 5 min. To form the MoS_2 nanoscrolls, $100\sim 500$ μ L isopropyl alcohol (IPA) was carefully dropped on the surface according the size of substrate. When the solution was vapored, the MoS_2 nanoscrolls were formed. It should be noted that there are some parameters which will influence the scrolling process, for example the clean and size of MoS_2 sheet. The remained contaminations will prevent the scrolling due to

the difficult inserting of IPA solution. Too large size of MoS_2 sheet will results in un-uniform MoS_2 nanoscrolls possible due to the random defect.

Device fabrication and Characterization: The electrodes were patterned by electron beam lithography. 10 nm Ti/ 50 nm Au layers were deposited by electron beam deposition with 0.2 Å/s. The source and drain electrodes were formed after lift-off. The electrical property of field effect transistor was measured in vacuum probe station with Keithley 4200 SCS. The Raman spectroscopy measurements were carried out on Renishaw inVia Raman Spectroscope with 514 nm and 1.0 mW laser excitation. The XPS measurements were carried out on Kratos Axis Ultra DLD spectrometer under a vacuum of 1×10^{-9} torr. To prepare the sample for TEM measurement, a PMMA layer was spin-coated on MoS_2 nanoscrolls/ SiO_2/Si . After baked at 100 °C for 5 min, PMMA/ MoS_2 nanoscrolls/ SiO_2/Si was immersed into DI water, and then PMMA/ MoS_2 nanoscrolls member would release from SiO_2/Si substrate. The member was picked up by TEM grid and dry in air. The PMMA was removed by warm acetone leaving MoS_2 nanoscrolls on TEM grid.

Computation Methods

We perform all the classical molecular dynamics (MD) simulation by using large-scale atomic/molecular massively parallel simulator (LAMMPS) within the isothermal-isobaric (NPT) ensemble.⁴⁹ The Stillinger-Weber (SW) potential developed by Kandemir et al⁵⁰ is employed to describe the interatomic interactions of MoS_2 nanoscrolls. The formula of SW potential is divided into a two-body bond stretching term and a three-body bond bending term. To ensure an accurate description of the three-body bond bending term, the top layer S atoms are treated as a different type of atoms from the bottom layer S atoms. Each nanoscroll model is constructed with the axial direction under periodic boundary conditions. Our MD runs consist of at least $1,200,000$ time steps with each time step of 0.5 fs. Each run time is at least 600 ps. All the classical MD simulations are conducted at room temperature 300 K. For the density functional theory (DFT) calculations, we adopt the Perdew-Burke-Ernzerhof functional within the general gradient approximation, as implemented in the Vienna ab initio simulation package (VASP 5.4).^{51,52} A plane-wave energy cutoff of 500 eV is used. The lattice constants and all atomic positions are fully relaxed by the conjugate gradient procedure until all the atomic forces are less than 0.01 eV. The tolerance for energy convergence is set to 10^5 eV. A Monkhorst-Pack mesh of $12\times 12\times 1$ and $1\times 1\times 6$ k -points are used to sample the Brillouin zone of the MoS_2 multi-layered nanosheet and multi-walled nanotube, respectively. The index for the single-, two- and triple-walled MoS_2 nanotubes are (12, 12), (6, 6)-(12, 12) and (6, 6)-(12, 12)-(18, 18), respectively.

Conclusions

In summary, the MoS_2 nanoscrolls are fabricated from CVD-grown MoS_2 sheet successfully. The formation of nanoscrolls always starts from triangular MoS_2 edge and the scrolling direction is perpendicular to MoS_2 edge along armchair direction, which suggests scrolling is driven by the internal

vdW interaction. Classical MD simulation is applied to investigate the dynamic formation of MoS₂ nanoscroll. Furthermore, the band structure of MoS₂ nanosheets and nanotubes, two limiting dimensional systems for the nanoscrolls, is calculated by DFT method. The electronic structure is also discussed based on its optical spectroscopies. According to the electronic transport property from low temperature electrical measurement, the carrier mobility and Schottky barrier are systematically investigated. In contrary to semiconducting 2D MoS₂ sheet, the quasi-1D MoS₂ nanoscroll shows a metallic-like behavior with low current on/off ratio, higher carrier mobility, negative Schottky barrier, possibly due to its notable feature of continuous electron distribution and strong electron-electron interaction. Such self-scrolling method can be easily applied to fabricate other 1D TMDs with unique functionality and extend potential application of TMDs.

Conflicts of interest

There are no conflicts to declare.

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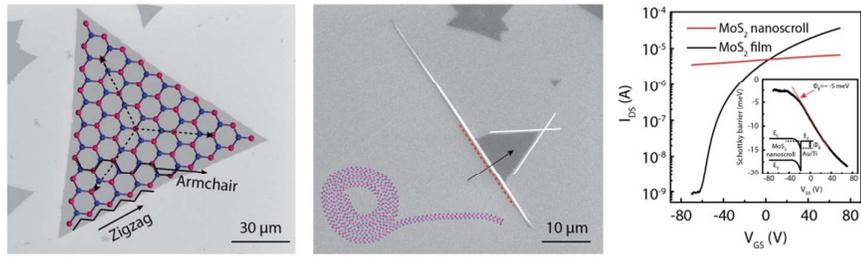
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Herein, the properties of quasi-1D MoS₂ nanoscrolls are systemically studied via experiment and theoretical simulation demonstrating that MoS₂ can transmit from semiconductor to metallic by self-scrolling.