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## Strategies to improve the mechanical robustness of metal halide perovskite solar cells†

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We report on the mechanical properties of high-efficiency perovskite solar cells (PSCs) with different chemical components by measuring the fracture energy ( $G_c$ ) of films and devices. With the help of both macroscopic and microscopic techniques, we identify the locations where fracture takes place in the devices (either adhesive or cohesive failure) with various material and device structures. We propose strategies that can improve the fracture energy of PSCs based on the measured  $G_c$ : the use of ozone-nucleated atomic layer deposition to improve charge transport layer robustness and the use of 2D perovskites and morphology control to improve the perovskite robustness. Our findings offer a pathway to rationally study the mechanical properties of PSCs and enable such cells to be more mechanically robust to reach commercial viability.

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### 1. Introduction

The power conversion efficiency (PCE) of perovskite solar cells (PSCs) has substantially increased to a record of 26.08%.<sup>1</sup> In order to achieve operational stability and long-term reliability, however, solar cells are required to be not only highly efficient but mechanically robust. Previous work has shown that environmental stressors, such as light illumination, can induce significant degradation in device efficiency due to mechanical failure.<sup>2,3</sup> Thermal cycling also plays a key role, as PSCs consist of various thin film layers with different thermal expansion coefficients, which leads to the development of stress from thermal processing or heating during operation. These layers are weakly bonded to each other<sup>4</sup> based on low adhesion, making them susceptible to strain-induced delamination. The large stresses and weak bonding in the films and/or interfaces can lead to mechanical failure at both the material

and the device level.<sup>5</sup> Such failures, including fracture and delamination, further accelerate environmental degradation and result in a loss of ohmic contact, which in turn compromises the PCE and ultimately leads to device failure. Other critical factors such as ion migration, hysteresis, and defect formation also play key roles in degradation, however their connection to mechanical failures is still unknown.

There are synergistic environmental and mechanical stresses which occur in operation that are not probed during individual stress tests and the degradation modes in perovskites are different than that of silicon PV. Unfortunately, the only series of accelerated tests are not based on individual factors and are provided by the International Electrotechnical Commission (IEC) and aim to uncover individual factors related to the reduction of the service life for PV modules.<sup>6</sup> Static and dynamic mechanical load tests as well as hail tests are included in the IEC 61646 thin film qualification tests that were designed for silicon PV, which include temperature cycling, humidity freeze, and damp heat accelerated exposures. Such tests currently are the only set of standards for perovskite modules to be mechanically stable enough to consistently operate in the field. Nonetheless, these tests alone are not sufficient to ensure that perovskite modules will reach the targeted 20-year service lifetimes to be competitive with incumbent PV technology. Systematical characterizations to reveal the link between mechanical properties and device stability is relatively scarce, which limit the rational design for stable PSCs.

To better understand and study the mechanical properties of PSCs, fracture energy ( $G_c$ ), an indicator to evaluate the

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robustness of materials, was used.  $G_c$  is the energy required to open a unit area of crack, and the two contributions of  $G_c$  come from plastic deformation and atomic bond breakage. For robust materials, such as PV encapsulants or metals, the former is the main component of the  $G_c$ , often in the range of thousands of  $\text{J m}^{-2}$  and above! However, the energy to break atomic bonds is much lower, and this is the dominant component for PSCs due to their intrinsically fragile and brittle nature. Additionally, given the weak ionic bonding of perovskites, the bond breakage energy of perovskites is comparable to table salt. Previously, the fracture energy of a wide range of PSCs was investigated and the results showed that the  $G_c$  values of PSCs are commonly below  $1.5 \text{ J m}^{-2}$ , which are significantly lower than c-Si cells ( $\sim 10\text{--}200 \text{ J m}^{-2}$ ).<sup>7</sup> In mixed-cation perovskites that incorporated cesium (Cs) and formamidinium (FA), the decreased grain size and impurities between grain boundaries resulted in a reduced  $G_c$  of MA/FA, Cs/FA, and Cs/FA/MA devices.<sup>8</sup> Some strategies for increasing perovskite  $G_c$  included the use of externally reinforcing scaffolds,<sup>9</sup> polymeric additives,<sup>10,11</sup> and grain boundary engineering<sup>12</sup>/toughening approaches.<sup>13</sup> However, these were mostly done for  $\text{MAPbI}_3$ -based perovskites that are no longer favored in the perovskite community due to chemical instability. Additionally, there has been limited progress on improving the highly fragile, small molecule-based (fullerenes, spiro-OMeTAD, Li-based dopants) charge transport layers (CTLs) and additives that often have even lower  $G_c$  values than perovskites.

In this work, we study the mechanical integrity of metal halide PSCs with updated, state-of-the-art structures and

compositions. The goal was to improve the mechanical robustness of PSCs through strategies that incorporate molecules with improved bonding and/or deformability. We found that (1) the widely used fullerene layer in PSCs can be strengthened by adding ozone nucleation-treated tin oxide buffer layer, (2) the vacuum-evaporated MA-based and Cs-FA mixed-cation perovskites have relatively low  $G_c$  values due to their high grain boundaries and small grains, (3) two-dimensional (2D) Ruddlesden–Popper (RP) perovskites with *n*-butylammonium (BA) and benzylammonium (Bnz) spacer cations both have higher  $G_c$  values through the increase in plastic deformation for lower *n*-values. This understanding of how such materials influence the  $G_c$  of PSCs is studied in this work for the first time.

## II. Methods

The perovskite samples for fracture testing were all deposited on glass substrates as full devices or partial stacks with or without CTLs. To create a double cantilever beam (DCB) fracture sample, a layer of thin epoxy was applied to another identical glass superstrate and then bonded to the device/stack (Fig. 1a) to create a sandwich-like structure with the device layers bonded between glass. The epoxy was cured at room temperature to create the DCB structure in Fig. 1b. It is important to note that all samples had a length of at least 3 cm to ensure enough area to create a beam without reaching the edges of the sample.



Fig. 1 Schematic illustrations of (a) DCB sample fabrication method, (b) General DCB sample structure (c) Pre-crack (d) Load–displacement response of an ideal fracture energy measurement.



By using a DCB configuration in the delaminator adhesion test system (DTS, Menlo Park, CA), measurement of the critical load needed for fracture propagation can be realized along with the value of  $G_c$ . A precrack is introduced to the DCB samples by carefully inserting the tip of a razor in between the two glass substrates to propagate the initial crack (Fig. 1c). In the measurement, as shown in Fig. 1d, DCB samples are pulled apart by an electric actuator with a constant displacement rate. The displacement rate is the velocity of the actuator moving in the direction in which the DCB specimens are pulled apart. The resulting load generated by this change in displacement is measured with the load. At stage I, a uniaxial tensile load is applied to the sample, and a linear loading curve is observed until the initial precrack starts to extend, generating a decrease in the load at stage II. This decrease in load corresponds to the critical load where crack propagation initiates. Once a decrease in the load–displacement plot is observed, the specimen is partially unloaded and then subjected to more loading and unloading cycles at stage III until the crack spreads through the entire length of the thin film structure. It is important to note that the sample is always loaded under uniaxial tensile force until it breaks. The weakest interface or layer in the device is determined after the DCB sample breaks apart fully at stage IV. The value of  $G_c$  is then calculated for each critical load and averaged to obtain multiple data points per specimen in the following equation:<sup>8,14</sup>

$$G_c = \frac{12P_c^2 a^2}{B^2 E' h^3} \left(1 + 0.64 \frac{h}{a}\right)^2 \quad (1)$$

where  $P_c$  is the critical load that deviates from the linear part in the load–displacement plot during the loading cycle;  $a$  is the crack length estimated by the compliance method;  $B$  and  $h$  are the widths and half height of the sample, respectively; and  $E'$  is the plane-strain elastic modulus of the substrate (Fig. S1, ESI<sup>†</sup>). We note one of the benefits of this method is that no elastic properties of the thin films are needed, which greatly simplifies the analysis. Additionally, the process is identical regardless of the number of thin films and the thickness of the films

assuming they remain much thinner than the substrate thickness of 1 mm, which is always the case for PSCs.

The crack length was estimated by a compliance method:

$$a = \left(\frac{d\Delta}{dP} \cdot \frac{BE'h^3}{8}\right)^{\frac{1}{3}} - 0.64h \quad (2)$$

where  $\Delta$  is the displacement and  $P$  is the applied load. The  $G_c$  tests were performed under laboratory air environment.

In this work, we tested a large number of p–i–n devices and film stacks with various chemistries for the CTLs and perovskites to determine strategies for increasing the fracture energy. The key take-aways are reported below.

### III. Results and discussion

#### Challenges in transport layers and improving fullerene mechanical fragility

In a DCB test, the fracture path occurs adhesively at the weakest interface, cohesively within the weakest layer, or through a combination of the two. Here, we first studied the  $G_c$  of p–i–n devices with MAPbI<sub>3</sub> and MAPb(I<sub>0.9</sub>Br<sub>0.1</sub>)<sub>3</sub> as the absorber layers and found that the crack propagated cohesively either in the organic P or N CTLs (Fig. 2a and Fig. S2, ESI<sup>†</sup>), but mainly in C<sub>60</sub> layer, with a  $G_c$  of  $0.23 \pm 0.15 \text{ J m}^{-2}$  and  $0.16 \pm 0.03 \text{ J m}^{-2}$  for MAPbI<sub>3</sub> and MAPb(I<sub>0.9</sub>Br<sub>0.1</sub>)<sub>3</sub> samples, respectively, showing that the CTL tends to be the weakest in the p–i–n device stack. Compared with the  $G_c$  of crystalline silicon ( $10\text{--}200 \text{ J m}^{-2}$ ) and organic PV ( $\sim 5 \text{ J m}^{-2}$ ), the p–i–n devices in this work are considered to be extremely fragile (Fig. 2b). The fragility of organic CTLs has been studied previously,<sup>7</sup> and in most p–i–n devices, there is an increase in  $G_c$  when using inorganic layers (*i.e.*, NiO<sub>x</sub>) and undoped polymers (*i.e.*, PTAA, poly-TPD, PEDOT:PSS) that are comparably more mechanically robust than small molecules.<sup>15–22</sup> The TaTm/F6TCNNQ small molecule structure is therefore not ideal from a mechanical perspective ( $G_c$  of  $0.57 \pm 0.19 \text{ J m}^{-2}$ , Fig. S3, ESI<sup>†</sup>).

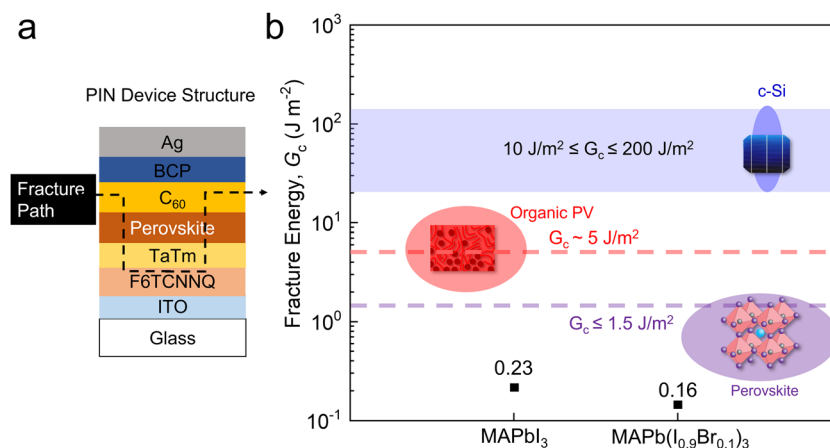


Fig. 2 (a) Device structure and fracture path. (b) The measured average  $G_c$  of the p–i–n samples with MAPbI<sub>3</sub> and MAP(I<sub>0.9</sub>Br<sub>0.1</sub>)<sub>3</sub> in comparison with the  $G_c$  of various PV technologies.



However, almost all p-i-n devices use a  $C_{60}$  layer to prevent perovskite degradation when depositing the N-type CTL.<sup>23–29</sup> There have been only a few reports of fully inorganic CTL devices with high power conversion efficiencies exceeding 20%,<sup>30,31</sup> and these are not possible for stable tandem structures that require a sputtered top transparent conducting electrode. Tin oxide is the most common buffer layer for these top conductive electrodes and to improve device stability. In order to reinforce the  $C_{60}$  layer, adding a tin oxide buffer layer enabled by atomic layer deposition (ALD) with good adhesion can not only boost cell efficiency but improve the operational and mechanical stability of perovskite devices.<sup>32,33</sup> However, previous work showed that peel-off tests on p-i-n perovskite/silicon tandem solar cells delaminated at the interface of  $C_{60}$  and  $SnO_x$ , demonstrating the fragility of the bonding between  $C_{60}$  and  $SnO_x$ .<sup>34</sup>

To address the  $C_{60}$  mechanical fragility, we performed fracture energy measurements on samples consisting of glass/ITO/perovskite/ $C_{60}$ /ALD- $SnO_x$ /Ag. The results confirmed the weak bonding at the  $C_{60}$ /ALD- $SnO_x$  interface with a low  $G_c$  of  $0.68 \pm 0.18 \text{ J m}^{-2}$ . In order to understand the delamination mechanism, the  $C_{60}$ /ALD- $SnO_x$  interface was investigated by scanning transmission electron microscope (STEM) and the cross-sectional STEM images show numerous vertical microscopic cracks at the interface (Fig. S4, ESI<sup>†</sup>), which can readily initiate the delamination of the  $C_{60}$  and ALD- $SnO_x$  under mechanical loads.

By treating the ALD- $SnO_x$  capping layer above  $C_{60}$  with ozone and water,<sup>35</sup> a higher contact area can be achieved due to fewer cracks at the interface (Fig. S4b, ESI<sup>†</sup>). This can be ascribed to the fact that the nucleation of the ALD process on the  $C_{60}$

surface was improved by the ozone treatment, which makes the surface chemistry more reactive to the ALD- $SnO_x$ , allowing for the formation of  $Sn-O-C_{60}$  bonds. Besides, the added ozone may also polymerize and cross-link  $C_{60}$ , enabling ALD- $SnO_x$  and  $C_{60}$  to form stronger bonding. To verify if the improvement in the interface morphology can result in higher mechanical robustness, fracture energy measurements were performed on samples with ALD- $SnO_x$  layers that were treated with or without the aforementioned ozone nucleation process. A relatively low  $G_c$  ( $0.39 \pm 0.11 \text{ J m}^{-2}$ ) and an adhesive failure at the  $C_{60}$ /ALD- $SnO_x$  interface were observed for control samples, while there was a higher  $G_c$  ( $0.58 \pm 0.12 \text{ J m}^{-2}$ ) and a meandering fracture path in the  $C_{60}$  layer for the samples with ozone nucleation treatment (Fig. 3 and Fig. S5, ESI<sup>†</sup>). The higher  $G_c$  value and the meandering fracture path in the  $C_{60}$  layer of ozone-nucleated samples strongly suggest that the ozone nucleation treatment effectively improves the  $C_{60}$ /ALD- $SnO_x$  adhesion.

We also tested film stacks with perovskite (glass/ITO/perovskite/ $C_{60}$ /ALD- $SnO_x$ /Ag) and observed that the fracture occurred cohesively within the perovskite, fracturing at  $1.12 \pm 0.34 \text{ J m}^{-2}$ , compared to the control stacks (glass/ITO/perovskite/ $C_{60}$ /Ag) where the fracture involved the  $C_{60}$  layer, failing at  $0.68 \pm 0.18 \text{ J m}^{-2}$  (Fig. S6, ESI<sup>†</sup>). This indicates that the ALD- $SnO_x$  can reinforce the  $C_{60}$  and that the perovskite becomes the weakest layer in the device stack. As such, we focus on strategies to improve perovskite mechanical robustness for the rest of this study.

### Challenges in active layer and improving perovskite mechanical fragility

Vacuum evaporation has been considered one of the most promising solutions to solve PSC scaling challenges.<sup>36–38</sup>

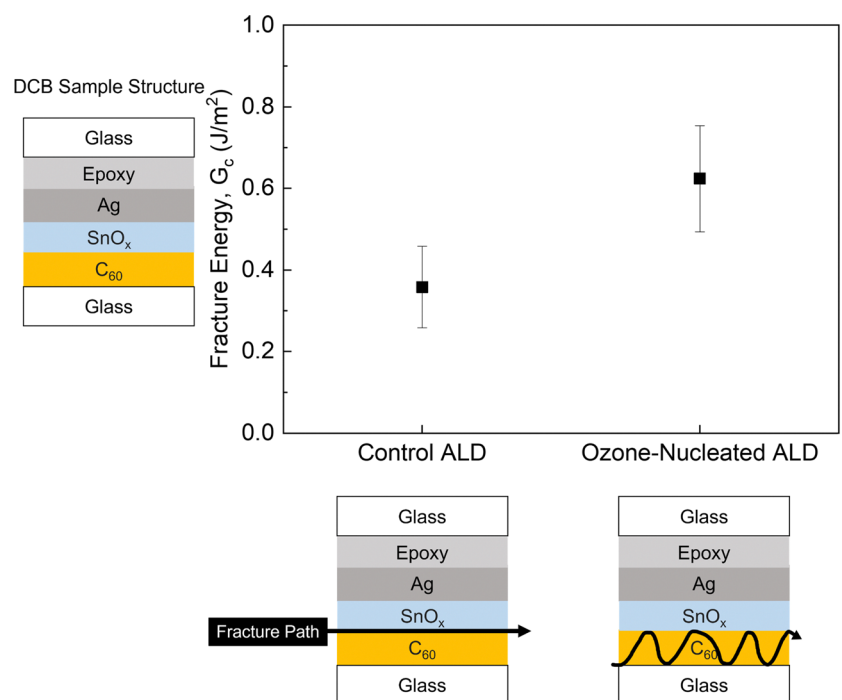


Fig. 3 A comparison between the average  $G_c$  results and fracture paths of control ALD- $SnO_x$  and Ozone-nucleated ALD- $SnO_x$  samples.



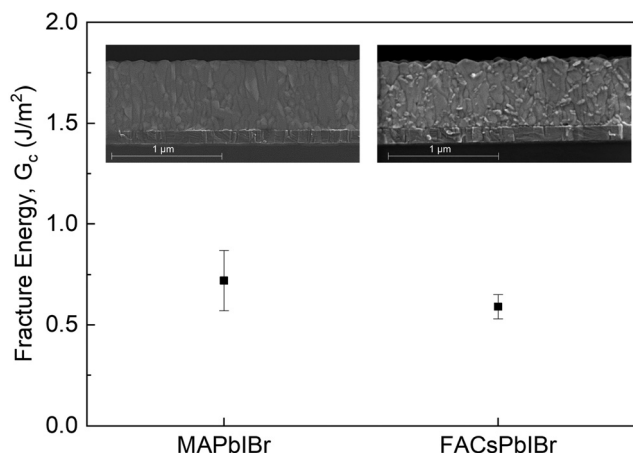


Fig. 4 The measured  $G_c$  of vacuum-evaporated  $\text{MAPb}(\text{I}_{0.9}\text{Br}_{0.1})_3$  and  $\text{FA}_{0.74}\text{Cs}_{0.26}\text{Pb}(\text{I}_{0.86}\text{Br}_{0.14})_3$  perovskites (inset: cross-sectional SEM image of the perovskite layer).

However, little is known about the mechanical properties of vacuum-evaporated perovskites. Here, we studied the fracture resistance of vacuum-evaporated  $\text{MAPb}(\text{I}_{0.9}\text{Br}_{0.1})_3$  and  $\text{FA}_{0.74}\text{Cs}_{0.26}\text{Pb}(\text{I}_{0.86}\text{Br}_{0.14})_3$  with the sample layout that comprised glass/ITO/perovskite/Ag. All samples failed cohesively in the perovskite layers with low  $G_c$  values ranging from  $0.87 \text{ J m}^{-2}$  to  $0.53 \text{ J m}^{-2}$  (Fig. 4 and Fig. S7, ESI<sup>†</sup>), indicating weak bonding in the perovskite films. A relatively lower  $G_c$  value in the FACS-based samples is due to the smaller grain size and higher density of grain boundary in the film. To better understand the delamination process in the vacuum-evaporated perovskite, we investigated the morphology of the perovskite film with scanning electron microscopy (SEM). The cross-sectional SEM images show fine grains and a high density of grain boundaries (GBs) in the vacuum-evaporated perovskite film, indicating that microscopic cracks can easily be initiated by those GBs under stresses and spread to the rest of the film, which causes cohesive failure in the devices. Therefore, coarse, large-grained perovskites (with sizes comparable to the film thickness) are

preferred for mechanically robust PSCs. Moreover, our results are in accordance with a previous study on the relationship between the mechanical reliability of perovskite thin films and their grain morphology.<sup>12</sup>

In previous studies, PSCs with mixed cations of *n*-butylammonium (BA) and MA were more mechanically robust than single-cation perovskites due to the larger size of BA and the possibility of the 2D cation increasing the film plasticity.<sup>8</sup> Here, we tested the pure 2D benzylammonium (Bnz)-based and BA-based Ruddlesden–Popper (RP) perovskite devices that adopt the chemical formula of  $\text{A}_2\text{MA}_{n-1}\text{Pb}_n\text{I}_{3n+1}$  ( $n = 1, 2, 3, 4$  and 5), where A is an organic spacer cation, and the structure as glass/PEDOT:PSS/perovskite/PMMA/Ag. PMMA and Ag are coated on top of the perovskite as a protective layer against epoxy. The images of fractured samples and fracture energy results are shown in Fig. 5 and Fig. S8, S9 (ESI<sup>†</sup>). The fracture energy decreases as the thickness of the inorganic layer ( $n$ ) increases.

For the BA-based samples with  $n = 1$ , the fracture happened at the interface between Ag/PMMA and the fracture energy of  $\text{BA}_2\text{PbI}_4$  is higher than the measured value ( $0.94 \text{ J m}^{-2}$ ), indicating that the cohesion of  $\text{BA}_2\text{PbI}_4$  thin film is more robust than the measured Ag/PMMA adhesion (Fig. S8, ESI<sup>†</sup>). A thin layer of perovskite was attached to the PMMA when fracture propagated in  $n = 2$  and  $n = 3$  samples. Meanwhile, we found a substantial amount of perovskite remained on both sides of the fractured  $n = 4$  samples, suggesting that the perovskite layer was more mechanically fragile, which is in line with the low  $G_c$  of  $0.37 \pm 0.22 \text{ J m}^{-2}$  and can be ascribed to the weak van der Waals bonds in RP perovskites. A downward trend in  $G_c$  values with the increasing  $n$ -value was observed for both BA-based and Bnz-based 2D RP perovskites as the material properties become more similar to the 3D perovskite structure (Fig. S9, ESI<sup>†</sup>).

The incorporation of additives can significantly affect several critical aspects of perovskites, such as crystallization, film formation, defect passivation and crystal structure. Nonetheless, the impact of additives on the mechanical integrity of perovskites has not been thoroughly explored. Fracture energy

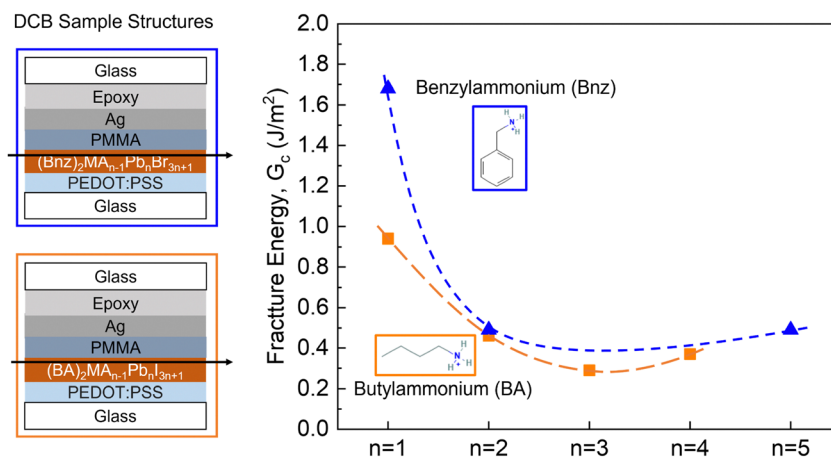


Fig. 5 The measured  $G_c$  values of RP phase perovskites as a function of  $n$ -value.



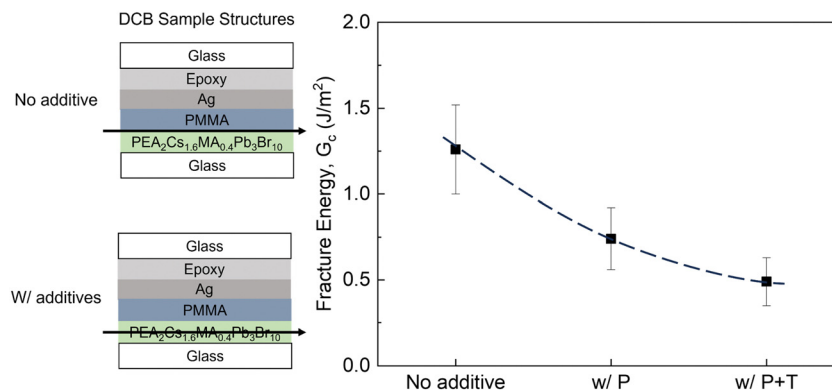


Fig. 6 The measured  $G_c$  values as a function of the incorporation of small molecule additive phosphine oxide (P) and tosylate (T) in the perovskite layer.

tests were performed on the quasi-2D perovskite  $PEA_2CS_{1.6}MA_{0.4}Pb_3Br_{10}$  (PEA = phenethylammonium) with small-molecule additives, such as phosphine oxide (P) and tosylate (T) that have a small molecular weight ( $\leq 1000$  Da). The incorporation process of the aforementioned additives in the 2D perovskite films is schematically shown in the Fig. S10 (ESI†). Adhesive failures at the interface of perovskite/PMMA were observed for the samples with no additive ( $G_c = 1.26 \pm 0.26$   $J m^{-2}$ ), while cohesive failure in the perovskite layers occurred when small-molecule P and T additives were added in the perovskites (Fig. S11, ESI†). The  $G_c$  decreased to  $0.74 \pm 0.18$   $J m^{-2}$  (with P only) and  $0.49 \pm 0.14$   $J m^{-2}$  (with P and T), as shown in Fig. 6. The decrease in the  $G_c$  values with the incorporation of additives, along with the cohesive failures, indicates that the small-molecule additives in perovskites can hinder the robustness of perovskite layers due to their weak bonding and lack of deformability despite a similar morphology with additives (Fig. S12, ESI†). A similar effect has been observed for commonly used Li-based (LiTFSI) dopants added to HTLs in n-i-p PSCs.<sup>39</sup>

## IV. Conclusion and outlook

In summary, we have shown that the CTLs are primarily the weakest element in the p-i-n perovskite device stacks. As a common material used for ETL,  $C_{60}$  has been proven that a layer of ALD-SnO<sub>x</sub> with ozone nucleation process can increase its  $G_c$  value and relocate the fracture to the next weakest layer in the stack, the perovskite. Furthermore, we found that vacuum-evaporated perovskites that have fine grains and a high density of GBs are not ideal for mechanically robust PSCs. Our investigation of the impact of small-molecule additives on the  $G_c$  values of perovskite films demonstrated that those small-molecule additives are fragile and capable of reducing  $G_c$ . Therefore, the key strategies to increase  $G_c$  are (1) to use of ALD-SnO<sub>x</sub> treated with ozone nucleation process to reinforce  $C_{60}$  (with the eventual goal of replacing  $C_{60}$  and any organic small molecules in CTLs and in the perovskite), (2) to use of 2D perovskites to increase plasticity and control of morphology through processing strategies.

This work considered a specific subset of mostly p-i-n perovskite devices where fragility was dominated by the CTL or the perovskite. However, we hope to provide a broader context here that can provide insights into each of the layers and interfaces of a typical perovskite device and module.

### Bottom TCO layer

This is not a source of mechanical fragility; however, there are considerations related to the brittleness of crystalline ITO (source) which leads to fracturing in flexible devices at low bending radii.<sup>40</sup> This is a consideration for flexible devices but not critical for rigid devices. To mitigate this effect, the use of more amorphous TCOs enables higher flexibility.

### TCO-HTL

Not a source of delamination or mechanical fragility to the best of our knowledge.

### HTL-perovskite

Surface roughness plays a role as well as interfacial adhesion; however, we did not observe delamination at this interface because the perovskite layer and perovskite-ETL interfaces were more fragile.

### Perovskite

Grain size is an important consideration (larger grains and more coarse) to improve fracture energy.

### Perovskite-ETL

Needs the most improvement, and ideally  $C_{60}$ /fullerene-based ETLs would be replaced entirely.

### ETL-top electrode

Metals (especially Au) do not bond well to materials and ideally would be replaced by carbon or top TCOs.

### Top electrode-encapsulation

Although adhesion is not a concern, there have been reports of pressure-induced delamination (*e.g.*, encapsulant peeling off the top electrode).<sup>41,42</sup>



Lastly, note that scribe lines present an additional interface and serve as a catalyst for delamination in some package configurations. These effects of monolithic integration are not well understood from a mechanical standpoint and are outside the scope of this work, but this is a key area that we are presently investigating in detail.

## Author contributions

Conceptualization, M. L. and N. R.; methodology, M. L. and N. R.; validation, M. L.; formal analysis, M. L.; investigation, M. L. S. J., L. G. E., M. S., C. A. F. M., H. K. and S. S.; resources, S. J., L. G. E., M. S., C. A. F. M., H. K. and S. S.; data curation, M. L., S. J., and L. G. E.; writing – original draft presentation, M. L. and N. R.; visualization, M. L. and N. R.; supervision, A. M., X. G., L. E., H. J. B., A. P., M. D. M., N. R.; project administration, N. R.; funding acquisition, A. M., X. G., L. E., H. J. B., A. P., M. D. M., N. R. All authors have read and agreed to the published version of the manuscript.

## Conflicts of interest

There are no conflicts to declare.

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## References

- 1 J. Park, J. Kim, H.-S. Yun, M. J. Paik, E. Noh, H. J. Mun, M. G. Kim, T. J. Shin and S. I. Seok, *Nature*, 2023, **616**, 724–730.
- 2 Q. Dong, C. Zhu, M. Chen, C. Jiang, J. Guo, Y. Feng, Z. Dai, S. K. Yadavalli, M. Hu, X. Cao, Y. Li, Y. Huang, Z. Liu, Y. Shi, L. Wang, N. P. Padture and Y. Zhou, *Nat. Commun.*, 2021, **12**, 973.
- 3 Z. Dai, S. K. Yadavalli, M. Chen, A. Abbaspourtamijani, Y. Qi and N. P. Padture, *Science*, 2021, **372**, 618–622.
- 4 Y. Rakita, S. R. Cohen, N. K. Kedem, G. Hodes and D. Cahen, *MRS Commun.*, 2015, **5**, 623–629.
- 5 C. Ramirez, S. K. Yadavalli, H. F. Garces, Y. Zhou and N. P. Padture, *Scr. Mater.*, 2018, **150**, 36–41.
- 6 R. Arndt and R. Puto, *TÜV SÜD Prod. Serv.*, 2010.
- 7 N. Rolston, B. L. Watson, C. D. Bailie, M. D. McGehee, J. P. Bastos, R. Gehlhaar, J. E. Kim, D. Vak, A. T. Mallajosyula, G. Gupta, A. D. Mohite and R. H. Dauskardt, *Extreme Mech. Lett.*, 2016, **9**, 353–358.
- 8 N. Rolston, A. D. Printz, J. M. Tracy, H. C. Weerasinghe, D. Vak, L. J. Haur, A. Priyadarshi, N. Mathews, D. J. Slotcavage, M. D. McGehee, R. E. Kalan, K. Zielinski, R. L. Grimm, H. Tsai, W. Nie, A. D. Mohite, S. Gholipour, M. Saliba, M. Grätzel and R. H. Dauskardt, *Adv. Energy Mater.*, 2017, **8**(9), DOI: [10.1002/aeam.201702116](https://doi.org/10.1002/aeam.201702116).
- 9 B. L. Watson, N. Rolston, A. D. Printz and R. H. Dauskardt, *Energy Environ. Sci.*, 2017, **10**, 2500–2508.
- 10 A. Giuri, N. Rolston, S. Colella, A. Listorti, C. Esposito Corcione, H. Elmaraghi, S. Lauciello, R. H. Dauskardt and A. Rizzo, *ACS Appl. Energy Mater.*, 2021, **4**, 11194–11203.
- 11 M. Gutwald, N. Rolston, A. D. Printz, O. Zhao, H. Elmaraghi, Y. Ding, J. Zhang and R. H. Dauskardt, *Sol. Energy Mater. Sol. Cells*, 2020, **209**, 110433.
- 12 Z. Dai, S. K. Yadavalli, M. Hu, M. Chen, Y. Zhou and N. P. Padture, *Scr. Mater.*, 2020, **185**, 47–50.
- 13 M. Q. Hovish, N. Rolston, K. Brüning, F. Hilt, C. Tassone and R. H. Dauskardt, *J. Mater. Chem. A*, 2020, **8**, 169–176.
- 14 M. F. Kanninen, *Int. J. Fract.*, 1973, **9**(1), 83–92.
- 15 W. Yan, S. Ye, Y. Li, W. Sun, H. Rao, Z. Liu, Z. Bian and C. Huang, *Adv. Energy Mater.*, 2016, **6**, 1600474.
- 16 E. Rezaee, X. Liu, Q. Hu, L. Dong, Q. Chen, J.-H. Pan and Z.-X. Xu, *Sol. RRL*, 2018, **2**, 1800200.
- 17 Z. Yu and L. Sun, *Small Methods*, 2018, **2**, 1700280.
- 18 F. Ma, Y. Zhao, J. Li, X. Zhang, H. Gu and J. You, *J. Energy Chem.*, 2021, **52**, 393–411.
- 19 S. Yoon and D. W. Kang, *Ceram. Int.*, 2018, **44**, 9347–9352.
- 20 D. S. Mann, P. Patil, S. N. Kwon and S. I. Na, *Appl. Surf. Sci.*, 2021, **560**, 149973.
- 21 M. Hu, C. Bi, Y. Yuan, Y. Bai and J. Huang, *Adv. Sci.*, 2016, **3**, 1500301.
- 22 Md Bodiul Islam, M. Yanagida, Y. Shirai, Y. Nabetani and K. Miyano, *ACS Omega*, 2017, **2**, 2291–2299.
- 23 J. Y. Liu, Y. Guo, M. H. Zhu, Y. C. Li and X. Li, *J. Power Sources*, 2020, **476**, 228648.
- 24 X. Liu, Y. Cheng, B. Tang, Z. G. Yu, M. Li, F. Lin, S. Zhang, Y. W. Zhang, J. Ouyang and H. Gong, *Nano Energy*, 2020, **71**, 104556.
- 25 C. Hanmandlu, S. Swamy, A. Singh, H.-A. Chen, C.-C. Liu, C.-S. Lai, A. Mohapatra, C.-W. Pao, P. Chen and C.-W. Chu, *J. Mater. Chem. A*, 2020, **8**, 5263–5274.
- 26 J. Wang, J. Zhang, Y. Zhou, H. Liu, Q. Xue, X. Li, C.-C. Chueh, H.-L. Yip, Z. Zhu and A. K. Y. Jen, *Nat. Commun.*, 2020, **11**, 177.



- 27 G. M. Wu, C. C. Tseng, T. W. Chang, W. S. Feng and D. W. Chen, *Surf. Coat. Technol.*, 2021, **405**, 126550.
- 28 L.-B. Chang, C.-C. Tseng, J.-H. Lee, G.-M. Wu, M.-J. Jeng, W.-S. Feng, D. W. Chen, L.-C. Chen, K.-L. Lee, E. Popko, L. Jacak and K. Gwozdz, *Vacuum*, 2020, **178**, 109441.
- 29 R. Prasanna, T. Leijtens, S. P. Dunfield, J. A. Raiford, E. J. Wolf, S. A. Swifter, J. Werner, G. E. Eperon, C. de Paula, A. F. Palmstrom, C. C. Boyd, M. F. A. M. van Hest, S. F. Bent, G. Teeter, J. J. Berry and M. D. McGehee, *Nat. Energy*, 2019, **4**, 939–947.
- 30 B. Yang, R. Ma, Z. Wang, D. Ouyang, Z. Huang, J. Lu, X. Duan, L. Yue, N. Xu and W. C. H. Choy, *ACS Appl. Mater. Interfaces*, 2021, **13**, 27179–27187.
- 31 S. Kohnepoushi, P. Nazari, B. A. Nejand and M. Eskandari, *Nanotechnology*, 2018, **29**, 205201.
- 32 W. Ke, D. Zhao, C. Xiao, C. Wang, A. J. Cimaroli, C. R. Grice, M. Yang, Z. Li, C.-S. Jiang, M. Al-Jassim, K. Zhu, M. G. Kanatzidis, G. Fang and Y. Yan, *J. Mater. Chem. A*, 2016, **4**, 14276–14283.
- 33 B. Chen, P. Wang, N. Ren, R. Li, Y. Zhao and X. Zhang, *J. Semicond.*, 2022, **43**(5), DOI: [10.1088/1674-4926/43/5/052201](https://doi.org/10.1088/1674-4926/43/5/052201).
- 34 M. De Bastiani, G. Armaroli, R. Jalmoed, L. Ferlauto, X. Li, R. Tao, G. T. Harrison, M. K. Eswaran, R. Azmi, M. Babics, A. S. Subbiah, E. Aydin, T. G. Allen, C. Combe, T. Cramer, D. Baran, U. Schwingenschlögl, G. Lubineau, D. Cavalcoli and S. De Wolf, *ACS Energy Lett.*, 2022, **7**, 827–833.
- 35 S. A. Johnson, K. P. White, J. Tong, S. You, A. Magomedov, B. W. Larson, D. Morales, R. Bramante, E. Dunphy, R. Tirawat, C. L. Perkins, J. Werner, G. Lahti, C. Velez, M. F. Toney, K. Zhu, M. D. McGehee, J. J. Berry and A. F. Palmstrom, *Joule*, 2023, DOI: [10.1016/j.joule.2023.10.009](https://doi.org/10.1016/j.joule.2023.10.009).
- 36 J. Ávila, C. Momblona, P. P. Boix, M. Sessolo and H. J. Bolink, *Joule*, 2017, **1**, 431–442.
- 37 L. Tan, J. Zhou, X. Zhao, S. Wang, M. Li, C. Jiang, H. Li, Y. Zhang, Y. Ye, W. Tress, L. Ding, M. Grätzel and C. Yi, *Adv. Mater.*, 2023, **35**, 2205027.
- 38 L. Gil-Escrig, C. Momblona, M.-G. La-Placa, P. P. Boix, M. Sessolo and H. J. Bolink, *Adv. Energy Mater.*, 2018, **8**, 1703506.
- 39 I. Lee, J. Hoon Yun, H. Jung Son and T.-S. Kim, *ACS Appl. Mater. Interfaces*, 2017, **9**, 7029–7035.
- 40 H. S. Jung, G. S. Han, N. G. Park and M. J. Ko, *Joule*, 2019, **3**, 1850–1880.
- 41 R. Cheacharoen, N. Rolston, D. Harwood, K. A. Bush, R. H. Dauskardt and M. D. McGehee, *Energy Environ. Sci.*, 2018, **11**, 144–150.
- 42 F. Toniolo, H. Bristow, M. Babics, L. Loiola, J. Liu, A. Said, L. Xu, E. Aydin, T. Allen, M. Meneghetti, S. Nunes, M. Bastiani and S. De Wolf, *Nanoscale*, 2023, **15**(42), 16984–16991.

