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# Resistive pulse analysis of chiral amino acids utilizing metal–amino acid crystallization differences†

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Here, we report a proof-of-concept resistive pulse method for analyzing chiral amino acids utilizing metal–amino acid crystallization differences. This method involves introducing an amino acid sample solution into a micropipette through a pressure-driven flow. The sample then mixes with a metal ion solution inside the pipette, forming metal–amino acid crystals. The crystal size depends on the enantiomeric excess ( $x$ ) of chiral amino acid samples. Large  $x$  values lead to large crystals. The crystal size difference is then reflected in the resistive pulse size as they block the ionic transport in a micropipette to different extents. We used Cd–cystine crystallization as a model system and found approximately five times the mean current pulse size difference for racemic ( $x = 0$ ) and L-only ( $x = +1$ ) cystine samples. A similar result was observed for aspartate. Our discovery opens up new opportunities for micro/nanoscale chiral amino acid analysis, which can potentially be used in single-cell analysis.

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## Introduction

Amino acids play a central role in the biosynthesis of proteins, thereby underpinning many biological processes. All common amino acids, except glycine, exist in two forms known as enantiomers: L- and D-. This molecular chirality is highly influential in determining the biological activity of amino acids and the proteins they form.<sup>1</sup> Understanding the chiral nature of amino acids is important for structural biochemistry and enzymatic catalysis, molecular interactions, and the broader physiological implications inherent in the stereochemistry of these fundamental building blocks of life.

L-Amino acids dominate biological systems, but recent studies have shown that there is an increased presence of D-enantiomers in the bodies of aging individuals.<sup>2,3</sup> These D-enantiomers have been suggested as biomarkers for age-related diseases, including schizophrenia and dementia.<sup>4–8</sup> In the quest to combat age-related diseases effectively, early detection of D-amino acids emerges as a strategy. The standard L-/D-amino acid analysis methods require liquid chromatography/tandem mass spectrometry (LC-MS/MS), which is

expensive and complex.<sup>9,10</sup> Therefore, there is a need for cost-effective and user-friendly methods for discriminating and quantifying L- and D-amino acids.

Some progress has been made in developing cost-effective methods for detecting chiral amino acids. For example, Ghasemi *et al.*<sup>11</sup> developed a visual method for naked-eye discriminating chiral cysteines using unmodified CdTe quantum dots, which exhibit different nanoparticle aggregation kinetics in the presence of L- and D-cysteine. The quantum dot aggregation alters the fluorescence emission for chiral amino acid detection. The Pu group developed a fluorescence chiral discrimination method based on 1,1-binaphthyl core fluorescent probes.<sup>12</sup> These probes exhibit high chemoselectivity and enantioselectivity for various amino acids such as glutamic acid, aspartic acid, serine, histidine, lysine, and tryptophan. They show fluorescence for one enantiomeric form of the amino acid while exhibiting little to no fluorescence for the other.

Besides fluorescence methods, electrochemical methods were also developed to distinguish chiral amino acids. These methods typically involved electrode polymerization or modification. For example, in a study conducted by Zhang *et al.*,<sup>13</sup> molecularly imprinted polymers were used on films on gate electrodes to differentiate between tyrosine and tryptophan. Similarly, Erbilin *et al.*<sup>14</sup> used the hierarchy polymerization of a glassy carbon electrode to distinguish tryptophan. However, there is still a need for chiral amino acid detection at the microscopic scale, which can potentially be used for single-cell analysis.

Here, we report a proof-of-concept resistive pulse method that can be used for analyzing chiral amino acids at the micro-

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**Fig. 1** (A) Design of the pulse-resistive method for analyzing chiral amino acids. A pressure is applied to drive the external amino acid (AA) solution into a micropipette, where AA reacts with the metal ions inside to form M-AA crystals. The M-AA crystal size varies depending on the enantiomeric ratio in an amino acid sample. L- or D-amino acid-only solution produces larger M-AA crystals than their mixtures. (B) Schematic illustration of the recorded current pulses for L- or D-amino acid only solution (red) and their mixture solution (green).

scopic scale, as illustrated in Fig. 1. Our method is based on the recent discovery by the Kotov group that metal ions and amino acids react to produce metal–amino acid microcrystals of varying shapes and sizes depending on the amino acid's enantiomeric excess ( $x$ ).<sup>15</sup> The value of  $x$  is calculated as  $(L\text{-enantiomer} - D\text{-enantiomer}) / (L\text{-enantiomer} + D\text{-enantiomer})$ . For instance, pancake-shaped microcrystals are formed when a racemic mixture ( $x = 0$ ) of L- and D-cystine (CST, the disulfide-bridged cysteine) is reacted with  $\text{Cd}^{2+}$ . However, bowtie-shaped microcrystals with left or right-handed directionality appear when  $x$  increases from 0 to +1 or –1. The size of particles formed at  $x = +1$  and –1 also doubles compared to those formed in a racemic CST mixture solution. The crystal shape and size alteration occurs due to the modified hydrogen bonding and geometric frustration among the building units as the metal–amino acid crystals assemble in the presence of different enantiomeric ratios of chiral amino acids.

In our method design, we use a resistive pulse technique to convert the size differences of metal–amino acid crystals into electrical signals. This technique involves a particle blocking the ionic conduction in a micro/nano-channel, which leads to a drop in ionic current.<sup>16–19</sup> This method was chosen mainly for two reasons. Firstly, the resistive pulse analysis produces a cubic function of the particle size signal,<sup>20–22</sup> which amplifies

the size difference of metal–amino acid microcrystals formed at different enantiomeric ratios of chiral amino acids and improves the detection sensitivity. Secondly, the nano/micropipettes used in the resistive pulse analysis provide micro- to nano-meter spatial resolution, allowing for potential analysis of chiral amino acids in a single cell.<sup>23–25</sup>

Specifically, a metal ion solution is placed inside a glass pipette, while an amino acid sample solution is placed outside it (Fig. 1A). The conductance in the pipette is monitored by applying a voltage bias between two Ag/AgCl wire electrodes inside and outside the pipette. As the amino acid solution is driven into the pipette *via* a pressure-driven flow, it mixes with the metal ion solution. The formation of nano- or micron-sized metal–amino acid crystals blocks the pipette orifice, generating a current pulse in the current–time trace. The size difference of the crystals formed in the presence of a single enantiomer *versus* a mixture of both enantiomers is detected from the current pulse size. Specifically, if L- or D-amino acid is alone, it would lead to larger pulses than their mixtures (Fig. 1B). This method can differentiate between a solution containing the L- or D-amino acid and a solution containing both the L- and D-amino acids.

## Results and discussion

### Preparation and characterization of Cd-CST crystals

First, we conducted a study to determine the feasibility of detecting Cd-CST microcrystals prepared *ex situ* through resistive pulse analysis. We created Cd-CST microcrystals using a slightly modified version of an existing protocol.<sup>15</sup> Briefly, we mixed 0.15 M NaCl, 1 mM  $\text{CdCl}_2$ , and 6 mM CST with varying  $x$  values. NaCl was added as the electrolyte. The resulting mixture immediately turned cloudy, indicating the formation of crystals (Fig. S1†). We shook the reaction vial and left it undisturbed for 15 minutes to complete the reaction.

The scanning electron microscopic (SEM) images of Cd-CST microcrystals synthesized from CST solutions with  $x$  values ranging from –1.0 (only D-CST) to 0 (equal L-CST and D-CST or a racemic mixture) and +1.0 (only L-CST) are shown in Fig. 2A. The crystals' shape and size evolved as a function of  $x$ , with the shape changing from chiral bowties at  $x = \pm 1.0$  to achiral disks at  $x = 0$ . The average length of crystals decreased from  $4.28 \pm 0.37 \mu\text{m}$  at  $x = +1.0$  and  $4.42 \pm 0.37 \mu\text{m}$  at  $x = -1.0$  to  $1.93 \pm 0.23 \mu\text{m}$  at  $x = 0$  (Fig. 2B). The average width of the bowtie was  $2.37 \pm 0.47 \mu\text{m}$  (Fig. S2†). Additionally, all microcrystal samples had a zeta potential of  $\sim -33 \text{ mV}$  (Table S1†), indicating negatively charged particles. The as-prepared glass pipettes have a native pore size of  $\sim 200 \text{ nm}$  (Fig. S3†). To match the Cd-CST microcrystal sizes in the resistive pulse analysis, pipette tips were gently cut to increase the pore size gradually. After successful resistive pulse analysis experiments, the pipettes were imaged by scanning electron microscopy. The typical pipette sizes used in our resistive pulse analysis experiments had a diameter of  $\sim 10$  to  $20 \mu\text{m}$  and a half-cone angle of  $\sim 2.4^\circ$  (Fig. 2C).





**Fig. 2** (A) SEM images of Cd-CST microcrystals formed at varying enantiomeric excess ( $x$ ).  $x = -1.0$  means 100% D-CST and 0% L-CST, while  $x = 0$  means a racemic mixture of 50% D-CST and 50% L-CST. (B) Particle length distributions of Cd-CST crystals at varying  $x$ . The error bars are the standard deviations of particle lengths after analyzing at least 100 particles. (C) SEM image of a micropipette used in the resistive pulse analysis of Cd-CST crystals.

### Resistive pulse analysis of *ex situ* synthesized Cd-CST crystals

With the available Cd-CST microcrystal size and micropipette geometry, we then performed a finite-element simulation to

estimate the theoretical resistive pulse size. For simplification purposes, we approximated the Cd-CST achiral disks ( $x = 0$ ) and bowties ( $x = \pm 1$ ) as 2D axisymmetric spheres and cylindrical rods with similar dimensions. Fig. 3A displays the simu-



**Fig. 3** (A) Simulated electric fields when a Cd-CST microcrystal enters a conical micropipette with an orifice radius of  $8 \mu\text{m}$ . A  $1 \mu\text{m}$ -radius sphere is used to simulate the microdisks synthesized using  $x = 0$ , and a microrod with a length of  $4.4 \mu\text{m}$  and a radius of  $1.25 \mu\text{m}$  is used to simulate the bowties synthesized using  $x = \pm 1$ . (B) Plots of current vs. the particle position for  $x = 0$  and  $\pm 1$ . (C) Typical background-subtracted current-time traces during resistive pulse analysis of *ex situ* synthesized Cd-CST crystals for various  $x$ . (D) Distributions of the current pulse sizes for  $x = -1$ ,  $0$ , and  $+1$ .  $n$  is the number of pulses analyzed, and the errors are the standard deviation of the current pulse sizes.



lated electric fields as a Cd-CST microcrystal enters a conical micropipette with an orifice radius of 8  $\mu\text{m}$  for  $x = 0$  and  $\pm 1$ . The electric field is more significantly disrupted by the bowtie than the disk due to its larger size, producing a current pulse of 0.19 nA, about four times larger than the disk counterpart (Fig. 3B).

Next, the resistive pulse experiments were carried out in a homebuilt setup, as shown in Fig. S4†. A micropipette filled with 0.15 M NaCl was inserted into the Cd-CST crystal solutions. Two Ag/AgCl wires were placed separately inside the micropipette and in the sample solution. A voltage bias of 0.1 V (internal vs. external) was applied between the two electrodes, and the current was monitored throughout the experiment. Since the Cd-CST microcrystals tend to settle down in the solution due to their large sizes (Fig. S5†), the sample solution was vigorously agitated to homogenize it before the measurements. Only the first minute of the chronoamperograms was recorded and analyzed.

Fig. 3C shows typical chronoamperograms during resistive pulse analysis of Cd-CST crystals synthesized from CST solu-

tions with different  $x$  values. All samples showed asymmetric current pulses caused by microcrystals translocating through the pipette. The pulse sizes ranged from  $-0.1$  to  $-0.4$  nA, falling within the range of the simulation result of  $-0.04$  to  $-0.19$  nA. The pulse size for  $x = 0$  was generally slightly smaller than other samples. To quantitatively compare the pulse size difference, a Python code was programmed for background subtraction and peak finding (Fig. S6†). The average current drops for L- and D-enantiomer were  $-0.103 \pm 0.03$  nA and  $-0.105 \pm 0.03$  nA, respectively, which were only slightly higher ( $\sim 8\%$ ) than that for a racemic mixture:  $-0.0956 \pm 0.03$  nA (Fig. 3D). The current pulse difference was smaller than the expected  $\sim 5$  times in Fig. 3B. This discrepancy likely resulted from the sampling bias during the analysis. The unstable microparticle suspension caused the larger particles to precipitate fast, and mostly, the smaller particles that can be driven into the pore by an electric field were sampled (Fig. S7†). However, such sampling bias would be avoided when the crystals were formed *in situ* by mixing the  $\text{Cd}^{2+}$  solution inside the pipette and the CST sample solution outside.



**Fig. 4** (A) A current–time trace during resistive pulse analysis of a racemic CST solution ( $x = 0$ ). The CST solution was mixed with the  $\text{Cd}^{2+}$  solution inside a micropipette *via* a pressure-driven flow. The initial current fluctuation at  $\sim 5$  s was caused by the applied pressure. (B) An expanded view of Panel (A). (C) Shows the enlarged view of the peak indicated by an asterisk. (D)–(F) A current–time trace during resistive pulse analysis of a L-CST solution ( $x = +1$ ) and its corresponding expanded views. (H) Current pulse size distributions for  $x = 0$  and  $+1$ . SE: standard error. (I) Current pulse size as a function of time for  $x = +1$  dataset. The data points are color-coded for each minute.



## Resistive pulse analysis of *in situ* formed Cd-amino acid crystals

After testing with *ex situ* formed Cd-CST crystals, we continued with the experimental design in Fig. 1. Fig. 4A shows the current vs. time trace in the first minute during the resistive pulse analysis of a 6 mM racemic CST solution ( $x = 0$ ). The CST solution was mixed with the 1 mM  $\text{Cd}^{2+}$  solution inside a micropipette *via* a pressure-driven flow (see the apparatus in Fig. S8†). The initial current fluctuation at  $\sim 5$  s was caused by the pressure-driven fluidic flow.<sup>26–28</sup> We observed current pulses with sizes typically less than  $-1$  nA (Fig. 4B). The pulse shape was also different from the asymmetric peaks observed in the *ex situ* experiments (Fig. 4C and 3C) because the crystals form inside the pipette and then exit it rather than translocation into the pipette. In comparison, the current pulse size was significantly larger for 6 mM L-CST solution ( $x = +1$ ) than the racemic sample (Fig. 4E–G). Statistical analysis of the current pulses in Fig. 4H shows that both racemic CST and L-CST-only samples produced a broad distribution of pulse sizes over  $\sim$ two orders of magnitude, possibly due to the stochastic particle nucleation and growth in the pipette. The pulse sizes are a few times to 100 times larger than those observed in the *ex situ* experiment in Fig. 3, possibly because altered crystal nucleation and growth kinetics within confinement compared to bulk solution lead to the formation of large crystals.<sup>29–31</sup> However, the ratio between the mean pulse size for L-CST and racemic CST samples is comparable to the predicted value of  $\sim 5$  by COMSOL simulation (Fig. 3B), indicating that the dependence of crystal growth kinetics on enantiomeric excess is still true in a confined pipette. A similar result was reproduced on a different pipette with a similar pore size (Fig. S9†). In addition, we did not observe a noticeable pulse size change as a function of time (Fig. 4I), suggesting the crystal nucleation and growth should be instant and reach equilibrium quickly. When the amino acid and  $\text{Cd}^{2+}$  solutions swapped positions, the current pulses for L-CST and racemic CST samples showed no noticeable difference (Fig. S10†), possibly because the crystal size is limited by the available metal ions rather than the chirality-dependent growth kinetics.

We further applied our method to other chiral amino acids, such as serine and aspartate. Free D-serine and D-aspartate were found in the mammalian central nervous system and serum that lacked D-amino acid oxidase.<sup>32</sup> Considerable research now indicates that these D-amino acids may be a potential therapeutic agent and biomarker in both schizophrenia and major depressive disorder.<sup>33</sup> Unlike CST, the reactions between  $\text{Cd}^{2+}$  and serine/aspartate do not produce well-defined bowtie or disk crystals (Fig. S11†), but resistive pulse analysis of serine and aspartate samples still produces current pulses (Fig. S12†). Fig. 5A shows the current pulse size distributions of racemic serine and L-serine solutions, which exhibit comparable distributions and mean values ( $-1.43$  vs.  $-1.03$  nA). In contrast, L-aspartate produces noticeably large current pulses ( $>-1$  nA) more frequently than its racemic counterpart (Fig. 5B). These results indicate that the chirality-dependent



Fig. 5 Current pulse size distributions for (A) serine and (B) aspartate samples with  $x = 0$  and  $+1$ . The solid and empty squares show the mean values, and the error bars show the standard errors.

metal–amino acid crystallization is also amino acid-dependent, which could be potentially used as a fingerprint signal for discriminating among amino acids.

## Conclusion

In conclusion, we demonstrated a proof-of-concept resistive pulse method for analyzing chiral amino acids utilizing the dependence of metal–amino acid crystallization on the enantiomeric excess of chiral amino acids. Using Cd-CST crystallization as a model system, we successfully discriminated between L-CST and racemic CST solutions from the resistive pulse size distributions. A similar pulse size difference between L-aspartate and racemic aspartate solution was observed but not for serine. Our work opens a new opportunity for analyzing chiral amino acids at micro or nanoscale, which could be potentially used for single-cell analysis.

## Author contributions

Blake Stringer: conceptualization, methodology, validation, formal analysis, investigation, writing – original draft. Alexandra Schmeltzer: software. C. Hyun Ryu: resources. Hang



Ren: resources. Long Luo: conceptualization, supervision, visualization, writing – review & editing.

## Conflicts of interest

There are no conflicts to declare.

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