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High-entropy materials for electrochemical energy storage devices

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Single phased, high-entropy materials (HEMs) have yielded new advancements as energy storage materials. The mixing of manifold elements in a single lattice has been found to induce synergistic effects leading to superior physicochemical properties. In this review, we summarize recent advances of HEMs in energy storage applications such as metal-ion batteries, supercapacitors, and fuel cells. We begin with defining HE materials (HEMs) and discussion of the synthetic methods and characterization techniques appropriate for evaluating HEMs at various length scales. We also discuss the application of a wide variety of HEMs, including HE alloys, oxides, chalcogenides, Prussian blue analogues, and sodium super ionic conductor (NASICON) materials in energy storage systems. Finally, advantages, challenges, and future perspectives of HEMs in energy storage systems are discussed.

Introduction

In the future, generating usable energy from renewable sources such as solar, hydroelectric, and wind are expected to replace conventional fossil fuels due to their limited reserves and resultant impact issues.[1–3] However, most sources of renewable energy are intermittent, for example, solar panels generate no power at night, which is problematic to maintaining constant energy supplies required for large-scale grid applications. Hence, the development of electrical energy storage systems are of great interest to enabling renewable power generation, coupled with energy stability.[4,5] Among them, electrochemical energy storage (EES) systems such as batteries and supercapacitors have been largely used in electrical devices such as electrical vehicles, mobile phones, and laptops, but their limited energy density and cycling capability have stymied widespread development.[6–9]

To solve these issues, the development of energy storage materials with satisfactory electrochemical performances is key. Recently, high-entropy materials (HEMs) have garnered much attention in this space. HEMs were first manifested as high-entropy alloys (HEA) in 2004,[10,11] consisting of a single-phased, multicomponent solid solution with at least five principal elements in near-equimolar concentrations. The compositional complexity bestows HEAs with emergent properties, including entropy stabilization, lattice distortion, sluggish diffusion and the so-called "cocktail effect".[12] In entropy stabilization, the high configurational entropy within these materials compensates for the enthalpic requirements of structural stabilization. Lattice distortion is caused by atoms with different crystal radii randomly occupying the positions in the lattice of the HEA. A wide range of physiochemical properties such as mechanical, thermal, and electrical behaviors can be conferred by introducing various elements. Sluggish diffusion results from lattice distortion, giving rise to restricted diffusion kinetics within the material. Due to the incorporation of multiple elements in a single phase, with each one playing a critical role in determining the structural properties, the combinatorial and synergistic properties result in the 'cocktail effect', leading to, among others, significantly enhanced catalysis. The concept of HEAs was later extend to other high-entropy (HE) systems such as HE oxides,[13–17] chalcogenides,[18–20] borides,[21–24] carbides,[25–28] nitrides,[29–31] silicides,[32,33] phosphides/phosphates,[34–36] fluorides,[37–39] NASICON-types[40–42] and Prussian blue analogues.[43–45] These HEMs with their interesting and often unexpected properties are expected to be promising active materials for electrochemical energy storage systems.

In this review, we summarize the recent progress on the HEMs related to their electrochemical energy storage applications. Firstly, the concept of HEMs will be introduced. Then, synthetic methods and characterization techniques will be summarized. Next, we provide a review of the reported applications of HEMs in electrochemical energy storage devices, including Li-ion, Na-ion, Li–S, and Zn-ion batteries, supercapacitors, and fuel cells. Overall, this review aims to provide an overview of the wide range of functional HEMs for electrochemical energy storage systems and pinpoint the existing challenges and future directions in this field.

Theoretical concept of HEMs

The general high entropy concept was first introduced by Cantor *et al.* and Yeh *et al.* based on the configurational entropy

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of mixing multiple principal elements with near-equimolar ratios.[10,47] For HEAs, the high entropy term refers to alloys which contain at least five elements in atomic percentages ranging from 5% to 35%. The state of 'high entropy' may then be quantified by the simple entropy of mixing:

$$\Delta S_{\text{conf}} = -R\sum_{i=1}^{n} x_i \ln x_i \quad (1)$$

where R represents the gas constant, x_i is the molar fraction of component i, and n is the total number of constituent elements. Specifically, HEAs with ΔS_{conf} value of $\geq 1.50R$ or larger have been categorized as high-entropy, while medium-entropy and low-entropy alloys refer to alloys with $1.0R < \Delta S < 1.5R$ and $\Delta S < 1.0R$, respectively. For inorganic materials with more than one sub-lattice (such as HE oxides and chalcogenides), overall configurational entropy is calculated by the modified equation to account for contribution of the disordered sub-lattice to the overall material:

$$\Delta S_{\text{conf}} = -R\sum_{S} a^S \sum_{i} y_i^S \ln y_i^S \quad (2)$$

where a^S represents the molar concentration of each sublattice with respect to the whole composition, while y_i represents the molar ratios of each component element within the sublattice. As shown in Fig. 1, Buckingham *et al.* used this equation to calculate for the configurational entropy value of three common metal chalcogenide compositions (MX, MX_2 and M_2X_3, where M is the metals and X is chalcogenides), a value of $1.5R$ cannot be achieved in these systems with a singly disordered sub-lattice, even if the number of elements in that sub-lattice was increased to 15. Therefore, Buckingham *et al.* proposed a new criterion to classify HE chalcogenides by the number of elements present, which is often reported to have occurred at 6. This criterion could be extended to other inorganic HEMs that have two sublattices such as the rock-salt structured HE oxide $(Mg_{0.2}Co_{0.2}Ni_{0.2}Cu_{0.2}Zn_{0.2})O$[48] and the spinel structured $(CrMnFeNiCu)_3O_4$.[49] However, it cannot be applied to materials that have two or more metal sublattices such as HE perovskites of formula ABX_3-type (where A represents electropositive alkaline, alkaline-earth, or rare-earth metal element, and B and X are transition metals and oxygen or halogen anions, respectively)[50–52] and Prussian blue analogues $A_lM_n\left[M'_m(CN)_6\right]H_2O$ (where A represents alkali metal ions (*e.g.*, Na^+ and K^+), and M and M′ correspond to transition metal cations).[53] Due to the conflicting definition of high entropy across multiple materials, we propose a universal criterion to classify HEMs by the number of elements with a singly disordered lattice (or sub-lattice), which needs to be five or higher, with molar fractions of each component greater than 5%. It is worthwhile to note that for inorganic materials with multiple disordered sub-lattices, such as the HE chalcogenide (PbSnSb)(SSeTe), it is possible to achieve a configurational entropy value of $1.5R$[54] even without five principle elements in a disordered sub-lattice.

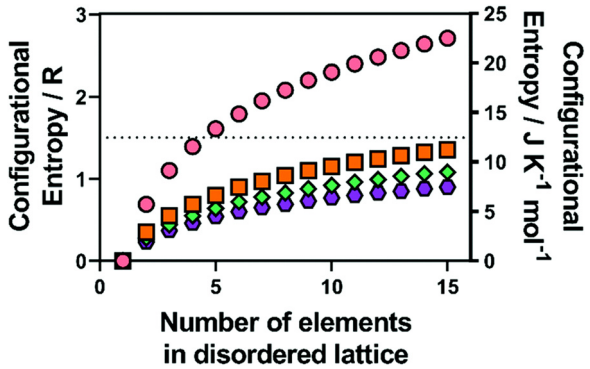

Fig. 1 Calculated configurational entropy (in both R and J K^{-1} mol^{-1}) for species with a single disordered sub-lattice (M) for (pink circles) alloys with a single lattice, and a range of common metal chalcogenide materials with a single chalcogenide sub-lattice (*i.e.* X is all O, S, Se or Te) for (orange squares) MX, (blue diamonds) M_2X_3 and (purple hexagons) MX_2. Copyright permission from Royal Society Chemistry, used from ref. 46.

Synthetic methods and techniques to characterize HEMs

Various physical and chemical synthetic methods have been developed to successfully fabricate HEMs with a single phase and homogenous elemental distributions. Classical solid state reaction is a common method to prepare HEMs, which involves several processing steps including pre-annealing mixing, such as mechanical ball milling to produce an even elemental mixture, followed by high-temperature calcination.[55–58] Generally, re-mixing and re-calcination steps are also required to further improve the elemental homogeneity in the final HEM. One drawback of this method is that the particle size of the solid-state reagents are around 100 nm, even after ball-milling, which inhibits elemental mixing at the nanoscale prior to the annealing process, which leads to the requirement of energy-intensive (usually above 1000 °C)[56] and time-consuming (generally between 10 h and 24 h)[59] procedures to ensure sufficient elemental diffusion at the nanoscale. High energy ball milling has recently been reported to produce a series of HE chalcogenides, which is performed at room temperature and does not require any external heating. However, this method requires a pre-synthetic step to produce each metal sulfide and an extended process time for ball milling, up to 120 h in total, which could potentially preclude scale up.[60]

Another common synthetic route towards HEMs is solvothermal synthesis, where precursors react in an organic solvent in an autoclave. This approach has the advantage of relatively low reaction temperatures (usually between 160 °C and 200 °C).[61–63] In addition, the morphology of the final product can be controlled by adjusting the reaction temperature and reagent concentrations. Bondesgaard *et al.* reported a facile and low-temperature solvothermal synthesis method to fabricate the high entropy alloy $Pt_{0.20}Pd_{0.20}Ir_{0.20}Rh_{0.20}Ru_{0.20}$ by dissolving metal acetylacetonate precursors in an acetone-ethanol

mixture solvent.[64] They found that the choice of reagent precursors and solvent is crucial to produce HEMs. The relatively large metal-acetylacetonate binding constant can significantly promote coprecipitation, which is beneficial to alloy formation.

Cation exchange methods have been reported to synthesize a nanoscale HE metal sulfide (HEMS).[65] As shown in Fig. 2, Schaak *et al.* used pre-fabricated roxbyite copper sulfide ($Cu_{1.8}S$) nanoparticles, which were then suspended in solutions that contain appropriate amount of various metal chlorides to be exchanged with Cu to form (ZnCoCuInGa)S HEMS nanoparticles. The mixture reacts initially at 110 °C under a vacuum for 30 mins and is then heated to 180 °C under Ar for another 30 min. Subsequently, the temperature of the reaction mixture decreases to 140 °C and is held for a further 30 min. After cooling to room temperature, the reaction mixture is centrifuged with a 1 : 1 mixture of IPA and acetone. This method is entropically favorable as the cation exchange reacts in solution, which is advantageous to the formation of HEMs. However, this approach is potentially unscalable as the relatively slow diffusion rate (we estimate *ca.* 115 days to diffuse 100 μm) would preclude bulk HEMS formation.[46]

A very effective synthetic route towards HEMs is *via* molecular precursors.[66–68] This method has significant advantages over other synthetic techniques, including homogenous mixing of the precursors at the atomic scale, which maximizes the disorder and thus entropic stabilisation in the solid-state products, prefabricated M-S bonds in the precursors which enable relatively reduced synthetic times (1–5 h) and relatively low preparation temperatures (300–500 °C) compared to other synthetic approaches.[69] Lewis and co-workers pioneered this route to successfully prepare bulk HE lanthanide oxysulfide materials based on the thermal decomposition (900 °C for 5 h) of five lanthanide dithiocarbamate precursors.[66] Lewis and co-workers also reported the synthesis of quantum confined (HE) lanthanide oxysulfide nanoparticles.[67] This method utilizes solvothermal decomposition of lanthanide dithiocarbamate precursors in the presence of capping agents. The obtained HE nanoparticles had identical elemental composition to that of the bulk HE oxysulfides. Recently, a novel 2D HE metal dichalcogenide $(MoWReMnCr)S_2$ has been fabricated *via* a combination of 'bottom-up' and 'top-down' approaches.[67] The bulk HE material was initially synthesized through solventless thermolysis of a mixture of five individual metal dithiocarbamate precursors at 500 °C for 1 h. The structural similarity between the as-prepared HE disulfide and layered 2H-MoS_2 enabled liquid-phase exfoliation to extract few layered HE $(MoWReMnCr)S_2$ nanosheets (3–7 atomic layers). The few layered nanosheets produced were found to contain a homogeneous elemental distribution of metals at the atomic scale. Overall, these approaches are potentially extendable to the synthesis of manifold HE chalcogenides due to the wide library of available transition, main group, and lanthanide single source precursors.[70] In addition, the ternary transition-metal dichalcogenides $Mo_{1-x}W_xS_2$ or transition-metal oxides $Mo_{1-x}W_xO_3$ have been previously produced *via* single source precursors simply by changing the decomposition atmosphere,[71] this approach can therefore potentially be extended to the synthesis of HE oxides.

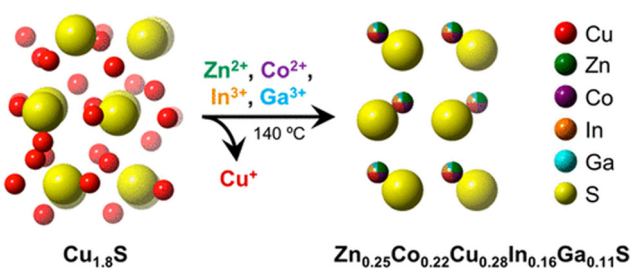

Fig. 2 Figure visually showing the formation of $(Zn_{0.25}Co_{0.22}Cu_{0.28}Ln_{0.16}Ca_{0.11})S$ HE metal sulfide synthesis *via* simultaneous multication exchange. Copyright permission from American Chemical Society, used from ref. 65.

A variety of material characterization techniques can be utilized to deduce HEMs crystal structures, morphologies, elemental distribution, and electronic structures. X-ray diffraction (XRD) is a typical method to characterize the HEM crystal structure, which provides a direct assessment of the phase purity of HEMs.[39] The refined lattice parameters, crystallinity, particle size, and orientation can also be calculated from diffraction data. Scanning electron microscopy (SEM) and transmission electron microscopy (TEM) are frequently used to observe the surface morphologies and particle dimensions of HEMs in the micro-to-nano scales.[65,75] High-angle annular dark-field – scanning transmission electron microscopy (HAADF-STEM) coupled with fast Fourier transform (FFT) analysis can reveal the atomic arrangements, lattice planes, grain boundaries, and dislocations in the structure at nano-to-atomic scale.[76] Yong-sheng *et al.* applied HAADF-STEM to characterize the atomic-scale structure of the HE perovskite $Na(Ni_{0.25}Mg_{0.05}Cu_{0.1}Fe_{0.2}Mn_{0.2}Ti_{0.1}Sn_{0.1})O_2$ (HEO424).[72] As shown in Fig. 3(a) and (b), both $Na(Ni_{0.4}Fe_{0.2}Mn_{0.4})O_2$ (NFM424) and HEO424 show dense structures without grain boundaries, indicating that both samples are a single crystal. Atomic-resolution HAADF-STEM imaging of a HEO424 particle reveals an hexagonally symmetrical arrangement of metals, which is in accordance with imaging down the [001] zone axis. HAADF-STEM used in tandem with energy-dispersive X-ray (EDX) spectroscopy is a powerful tool to confirm the homogenous elemental distribution at near-to atomic scale and discover whether phase separation exists in HEMs or not. Chen *et al.* applied HAADF-STEM elemental mapping to confirm the formation of a heterostructure between an high–entropy oxide and $CuCeO_x$.[73] As shown in Fig. 3(c), the Ce ions did not mix with the bulk HEO solid solution but was grafted onto the surface of the HEO particles. It is interesting to note that Ni, Mg, and Co elements are much diluted in this area compared to that in the core structure, while Cu and Zn elements dissolve with Ce, which indicates that Cu and Zn are incorporated in the CeO_2 lattice to form a $(CuZnCe)O_2$ solid solution on the shell of the HEO.

X-ray photoelectron spectroscopy (XPS) is a tool to analyze chemical environments within a material, including the

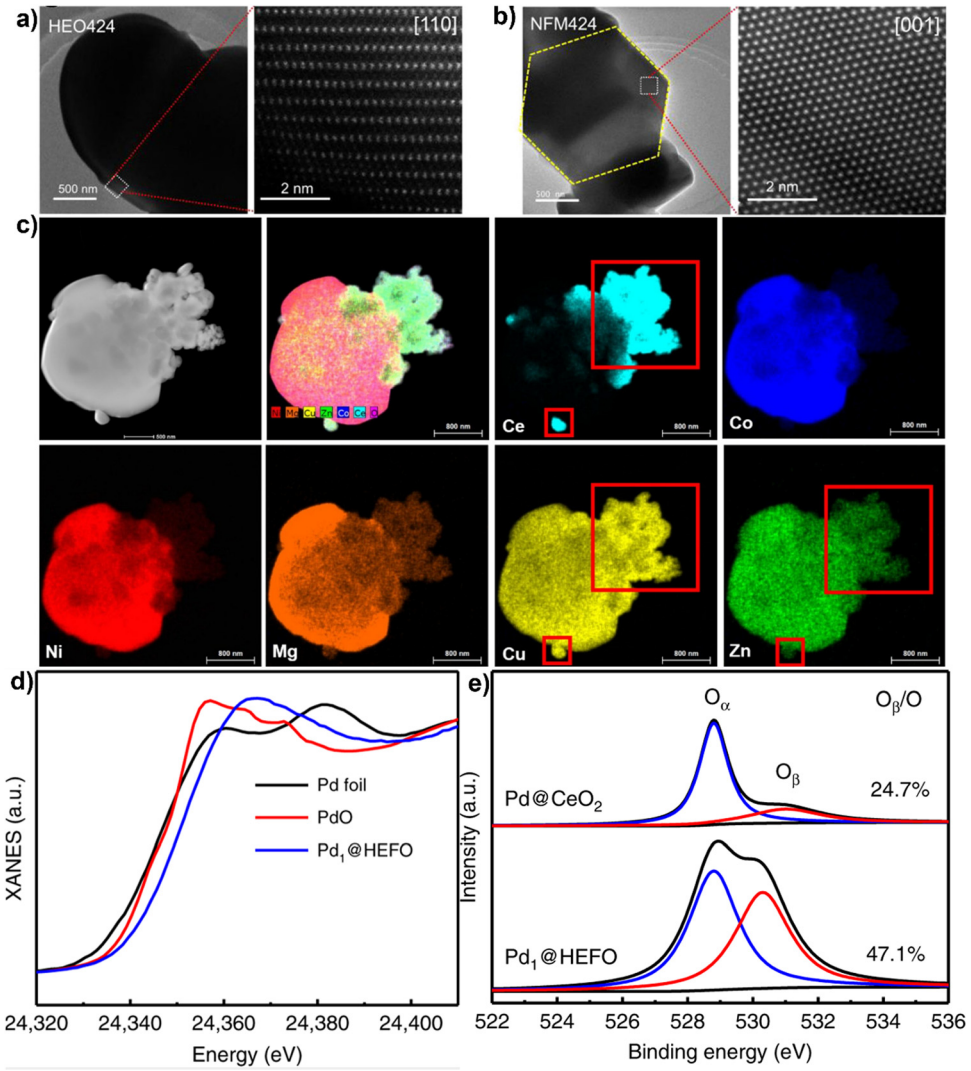

Fig. 3 TEM and HAADF-STEM images of $NaNi_{0.4}Fe_{0.2}Mn_{0.4}O_2$ (a) and $NaNi_{0.25}Mg_{0.05}Cu_{0.1}Fe_{0.2}Mn_{0.2}Ti_{0.1}Sn_{0.1}O_2$. (b) Copyright permission from American Chemical Society, used from ref. 72. (c) Elemental maps and high-resolution HAADF-STEM image results for $CuCeO_x$–(NiMgCuZnCo)O with 20 mol% Ce. Copyright permission from Elsevier, used from ref. 73. (d) XANES spectra at the Pd K-edge and (e) XPS profiles of O 1s for Pd@HEFO and $Pd@CeO_2$. Copyright permission from Springer Nature, used from ref. 74.

valence states and multielement compositions on the surface of HEMs.[77,78] When coupled with X-ray absorption near-edge structure (XANES), it can fundamentally reveal the electronic structure and coordination states of multiple elements in HEMs structure. For example, Xu *et al.* used XANES and XPS measurements in tandem to investigate the chemical surroundings of single Pd atoms in $(CeZrHfTiLa)O_x$ for CO oxidation catalysis.[74] As shown in Fig. 3(d), the Pd K-edge absorption edge spectrum for Pd_1@HEFO is between that of Pd foil and PdO, suggesting that the valence state of Pd is between 0 and +2. The valence state obtained by XANES is lower than the +4 from XPS, which is possibly related to the surface oxidation between Pd atoms and air. In addition, the O 1s XPS spectra shows that the ratio of absorbed oxygen (O_β) to the lattice oxygen (O_α) in Pd_1@HEFO is much higher than that of $Pd@CeO_2$ (Fig. 3(e)), which might be ascribed to the fact that the Pd incorporation in HEFO generates more oxygen vacancies. Also, the formation of more oxygen vacancies is consistent with its improved catalytic activity.

HEMs for energy storage systems

With the limited resources of fossil fuels and their related environmental issues, the rapid development of alternative energy sources is required.[79–81] This will include energy harvesting from waste materials and energy storage devices.[82–84] Electrochemical energy storage systems have advantages in sustainability and stable energy output. Batteries, supercapacitors, and fuel cells are examples of energy storage technologies.[85] All these devices consist of two electrodes and an electrolyte.[86] In batteries and fuel cells, the chemical energy derived from redox reactions at the two electrodes converts chemical potential to electrical power.[87] The electrode with

lower potential is designated as the anode and the cathode has a relatively higher potential. The major difference between batteries and fuel cells is the energy storage and conversion locations.[88] For batteries, the anode and cathode are not only the charge-transfer mediums, but are also active in the redox reactions. Batteries are therefore closed systems where the energy storage and conversion occur within the electrode materials. For fuel cells, the species participating in the redox reactions are supplied externally *e.g.*, air or a redox reservoir, where the energy storage and conversion are initiated independently. In supercapacitors, energy is not generated by the redox reactions, the electrical energy is stored and released by reversible adsorption and desorption of electrolyte between two electrostatic double layers (EDLs) formed by the electrode/electrolyte interface.[89] As there is no chemical reaction involved, the physical charging mechanism from electrostatic attraction alone enables supercapacitors to have rapid charge–discharge kinetics, cycling durability, and high-power densities.

The Ragone plot (Fig. 4) shows that fuel cells can be characterized by their high energy densities, while supercapacitors are characterized by their high power output.[90] Batteries have intermediary power output and energy storage compared to fuel cells and supercapacitors, with supercapacitors having a low energy density but high power density. Indeed, battery electrodes with high surface areas exhibit similar characteristics to that of supercapacitors,[91–93] where the redox reaction partly contributes to the energy storage, which results in the higher energy densities compared to conventional supercapacitors. At the same time, similar electrochemical storage mechanisms to EDL supercapacitors allow the delivery of better power density compared to batteries. Due to this synergistic effect, the design of hybrid electrochemical storage systems is of interest as they are capable of displaying high energy and power densities in comparison to EDL supercapacitors and batteries alone.

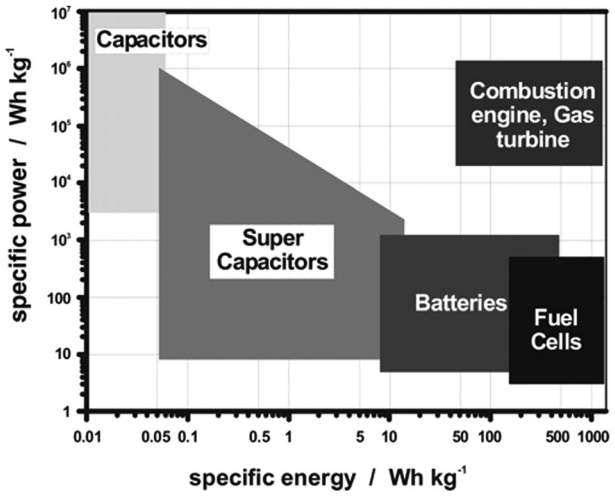

Fig. 4 Ragone plot of the energy density (W h kg^{-1}) and power density (W kg^{-1}) realms for the different electrochemical energy storage systems and conventional capacitors. Copyright permission from American Chemical Society, used from ref. 88.

In recent years, HEMs have been widely applied as electrode materials for electrochemical energy storage systems.[94–96] Due to their complicated compositions, HEMs often exhibit excellent properties that traditional materials cannot access easily, such as superionic conductivity,[97] high reversible capacity[26,98] and extended cycle life.[99] To date, a wide variety of HEMs including HE metal alloys,[100–102] oxides,[14,96,103] chalcogenides,[60] perovskites,[104–107] Prussian blue analogues[44] and NASICON-type materials[42] have showed extraordinary performances in metal-ion batteries,[15,108–110] supercapacitors[102] and fuel cells.[111]

HEMs for lithium-ion batteries (LIBs)

Previous studies have shown that HEMs have excellent performance as both anode and cathode materials for LIBs.[14,95,98,99,112] Owing to the structural stability conferred by entropic stabilization, the materials are more robust with respect to the charging/discharging processes when compared to traditional electrode materials. Furthermore, the introduction of electroactive cations with high redox reaction potentials, such as Ni and Fe, is effective for enhancing specific capacity as well as energy density.[99] Recent reported HEMs as electrode materials for LIBs are summarized in Table 1.

Among HEMs, rock-salt and spinel-type HE oxides have been most widely applied as electrode materials for LIBs.[14,95,140] Transition metal oxides (TMOs) are promising electrodes for LIBs as they have higher operating voltages compared to graphite-based materials.[149] Most TMO electrodes follow the conversion reaction mechanism: $M_xO_y + 2yLi^+ + 2ye^- \leftrightarrow yLi_2O + xM$,[150,151] which possesses higher theoretical specific capacities in comparison to intercalation-type electrodes[152] as more electrons involved in the reaction. The HEO rocksalt-type $(Mg_{0.2}Co_{0.2}Ni_{0.2}Cu_{0.2}Zn_{0.2})O$ was first utilized as an anode for LIBs.[14] In comparison to conventional LIB electrodes, the HEO exhibits significantly improved rate capability, even when cycling at high current densities. It is worthwhile to note that various crystallographic structures and emergent properties are obtained *via* compositional adjustment between multi-metallic cations and anions. For example, spinel HEO anodes $(Mg_{0.2}Ti_{0.2}Zn_{0.2}Cu_{0.2}Fe_{0.2})_3O_4$ and $(Ni_{0.2}Co_{0.2}Mn_{0.2}Fe_{0.2}Ti_{0.2})_3O_4$ have also been developed, with the former showing superior specific capacity of 1235 mA h g^{-1} with excellent retention of 90% after 200 cycles.[140] The electrochemically active Ni and Co in the lattice enable the entropy-stabilized HEO to perform redox reactions during charge/discharging processes, which confers a higher reversible capacity. Electrochemically inert cations, such as Mg^{2+} and Zn^{2+}, could play a critical role in stabilizing HEOs structure during charging/discharging processes and lead to extraordinary cycling performance. Thus, the multi metallic cations in the HEO collectively contribute to the superior electrochemical performances observed.

Due to the incorporation of various metallic cations in entropy-stabilized single-phase HEOs, the elemental interactions may lead to more complex lithiation behaviors compared to traditional metal oxides with one or two metallic cations. Feng *et al.* employed *in situ* electron diffraction (ED) and *ex situ* high-resolution transmission electron microscopy (HR-TEM)

Table 1 Recently reported HEMs as electrode materials for LIBs

HEMs	Synthetic method	Structure	Electrochemical performance	Capacity retention	Ref.
$(Co_{0.2}Ni_{0.2}Mn_{0.2}Zn_{0.2}Fe_{0.2})_3O_{3.2}$	High-energy ball milling	Rock salt	600 mA h g^{-1} after 200 cycles	83.5% after 45 cycles	113
$(NiMnCrCoFe)_3O_4$	Hydrothermal reaction	Spinel-type	527 mA h g^{-1} at 5 A g^{-1} current density	99.8% after 300 cycles	114
$(Mg_{0.2}Co_{0.2}Ni_{0.2}Cu_{0.2}Zn_{0.2})O$	High-energy ball milling	Rock salt	660 mA h g^{-1} at 100 mA g^{-1} current density	80% after 300 cycles	115
$(FeMnCrCoZn)_3O_4$	Co-precipitation method	Spinel-type	548 mA h g^{-1} at 100 mA g^{-1} current density	—	116
(a) (MgFeCoNiZn)O (b) $(TiFeCoNiZn)_3O_4$	High-temperature sintering	(a) Rock salt (b) Spinel-type	548 mA h g^{-1} and 478 mA h g^{-1} at 50 mA g^{-1} current density, respectively.	86% after 300 cycles	116
$Zn_{0.5}Co_{0.5}Mn_{0.5}Fe_{0.5}Al_{0.5}Mg_{0.5}O_4$	Co-precipitation method	Spinel-type	290 mA h g^{-1} at 2 A g^{-1} current density	81% after 5000 cycles	117
$LiNi_{0.2}Mn_{0.2}Co_{0.2}Fe_{0.2}Al_{0.2}O_2$	High-temperature sintering	Layered oxide	200 mA h g^{-1} at 18 mA g^{-1} current density	94.8% after 50 cycles	118
$(FeCoCrMnZn)_3O_4$	Ball milling	Spinel-type	829 mA h g^{-1} at 2 A g^{-1} current density	88% after 2000 cycles	112
$(YDyCeNdLa)_2Sn_2O_7$	Co-precipitation method	Pyro-chlore phase	396 mA h g^{-1} at 0.1 A g^{-1} current density	~100% after 100 cycles	119
$Li_{6.4}La_3Zr_{0.4}Ta_{0.4}Nb_{0.4}Y_{0.6}W_{0.2}O_{12}$	High-temperature sintering	Garnet-type	154 mA h g^{-1} at 17 mA g^{-1} current density	97% after 5 cycles	120
$(Mg_{0.2}Cr_{0.2}Mn_{0.2}Fe_{0.2}Co_{0.2})_3O_4$	Microwave-assisted method	Spinel-type	875 mA h g^{-1} at 1 A g^{-1} current density	87% after 1000 cycles	121
$Mg_{0.2}Co_{0.2}Ni_{0.2}Cu_{0.2}Zn_{0.2}O$	High-temperature sintering	Rock salt	700 mA h g^{-1} at 100 mA g^{-1} current density	—	122
$Co_{0.2}Cu_{0.2}Mg_{0.2}Ni_{0.2}Zn_{0.2})O$	Hydrothermal method	Rock salt	401 mA h g^{-1} at 0.5 A g^{-1} current density	90% after 3200 cycles	123
$(Cr_{0.2}Fe_{0.2}Co_{0.2}Ni_{0.2}Zn_{0.2})_3O_4$	Sol–gel method	Spinel-type	1022 mA h g^{-1} at 1 A g^{-1} current density	96% after 1000 cycles	124
$LiCr_{1/5}Mn_{1/5}Fe_{1/5}Co_{1/5}Ni_{1/5}O_2$	Solution combustion synthesis	Rock-salt	218 mA h g^{-1} at 20 mA g^{-1} current density	51% after 25 cycles	125
$(FeCoNiCrMn)_3O_4$	Ball milling	Spinel-type	597 mA h g^{-1} at 2 A g^{-1} current density	86.2% after 1200 cycles	126
$(CrMnFeNiCu)_3O_4$	Hydrothermal method	Spinel-type	340 mA h g^{-1} at 2 A g^{-1} current density	87.8% after 150 cycles	127
$(Mg_{0.2}Co_{0.2}Ni_{0.2}Cu_{0.2}Zn_{0.2})O$	Solvothermal Method	Rock-salt	480 mA h g^{-1} at 20 mA g^{-1} current density	86.5% after 300 cycles	48
(LiMgCoNiCuZn)O	Molten salt method	Rock-salt	714 mA h g^{-1} at 0.1 A g^{-1} current density	83% after 300 cycles	128
$Mg_{0.2}Co_{0.2}N_{i0.2}Cu_{0.2}Zn_{0.2}O$	Electrospinning method	Rock-salt	365 mA h g^{-1} at 2 A g^{-1} current density	82.8% after 200 cycles	129
$(CrNiMnFeCu)_3O_4$	Hydrothermal method	Spinel-type	480 mA h g^{-1} at 2 A g^{-1} current density	95% after 500 cycles	130
$(Mg_{0.2}Co_{0.2}Ni_{0.2}Cu_{0.2}Zn_{0.2})O$	High-temperature sintering	Rock-salt	1225 mA h g^{-1} at 0.1 A g^{-1} current density	81.7% after 1000 cycles	131
$(Co_{0.2}Cr_{0.2}Fe_{0.2}Mn_{0.2}Ni_{0.2})_3O_4$	Solution combustion synthesis	Spinel-type	428 mA h g^{-1} at 10 A g^{-1} current density	99% after 250 cycles	132
$(CrMnFeNiCu)_3O_4$	Hydrothermal method	Spinel-type	913 mA h g^{-1} at 50 mA g^{-1} current density	83% after 400 cycles	133
(Mg,Co,Ni,Cu,Zn)O	High-temperature sintering	Rock-salt	489 mA h g^{-1} at 0.1 A g^{-1} current density	77% after 100 cycles	134
$(CrMnFeCoNi)_3O_4$	Hydrothermal method	Spinel-type	1225 mA h g^{-1} at 20 mA g^{-1} current density	91% after 200 cycles	135
(Mg, Cu, Ni, Co, Zn)O	Microwave-assisted method	Rock-salt	686 mA h g^{-1} at 0.1 A g^{-1} current density	~100% after 1000 cycles	136
$(Al_{0.2}CoCrFeMnNi)_{0.58}O_{4-d}$	Solution combustion synthesis	Spinel-type	554 mA h g^{-1} at 0.2 A g^{-1} current density	71% after 500 cycles	137
$Mg_{0.2}Co_{0.2}Ni_{0.2}Cu_{0.2}Zn_{0.2}O$	High-temperature sintering	Rock-salt	600 mA h g^{-1} at 0.1 A g^{-1} current density	—	138
$(Ni_{0.2}Co_{0.2}Mn_{0.2}Fe_{0.2}Ti_{0.2})_3O_4$	High-temperature sintering	Spinel-type	594 mA h g^{-1} at 0.05 A g^{-1} current density	~100% after 100 cycles	139
$(Cr_{0.2}Mn_{0.2}Fe_{0.2}Co_{0.2}Ni_{0.2})_3O_4$	Hydro-thermal method	Spinel-type	500 mA h g^{-1} at 2 A g^{-1} current density	90% after 200 cycles	140
$(FeCoNiCrMn)_3O_4$	High-temperature sintering	Spinel-type	182 mA h g^{-1} at 2 A g^{-1} current density	81% after 300 cycles	141
$(MgCoNiZn)_{1-x}Li_xO$	High-temperature sintering	Rock-salt	610 mA h g^{-1} at 1 A g^{-1} current density	31.2% after 130 cycles	142
$(Mg_{0.2}Ti_{0.2}Zn_{0.2}Cu_{0.2}Fe_{0.2})_3O_4$	High-temperature sintering	Spinel-type	272 mA h g^{-1} at 2 A g^{-1} current density	78% after 300 cycles	143
$[(Bi,Na)_{0.2}(La,Li)_{0.2}(Ce,K)_{0.2}Ca_{0.2}Sr_{0.2}]TiO_3$	High-temperature sintering	*Pbnm* (62) space group	70 mA h g^{-1} at 3 A g^{-1} current density	~100% after 300 cycles	144
$(Co_{0.2}Cu_{0.2}Mg_{0.2}Ni_{0.2}Zn_{0.2})O$	Nebulized spray pyrolysis	Rock-salt	950 mA h g^{-1} at 0.02 A g^{-1} current density	94% after 100 cycles	145
$(Co_{0.2}Cu_{0.2}Mg_{0.2}Ni_{0.2}Zn_{0.2})O$	Nebulized spray pyrolysis	Rock-salt	446 mA h g^{-1} at 1.6 A g^{-1} current density	67% after 50 cycles	98
$(Mg_{0.2}Co_{0.2}Ni_{0.2}Cu_{0.2}Zn_{0.2}O)$	High-temperature sintering	Rock-salt	920 mA h g^{-1} at 0.1 A g^{-1} current density	58% after 300 cycles	99
$Li_{1.3}Mn_{0.1}Co_{0.1}Mn_{0.1}Cr_{0.1}Ti_{0.1}Nb_{0.2}O_{1.7}F_{0.3}$	High-temperature sintering	—	170 mA h g^{-1} at 2 A g^{-1} current density	66% after 19 cycles	146
$Li_x(Co_{0.2}Cu_{0.2}Mg_{0.2}Ni_{0.2}Zn_{0.2})OF_x$	High-temperature sintering	Rock-salt	161 mA h g^{-1} at 20 mA g^{-1} current density	71% after 300 cycles	147
$(FeNiCoMnCr)S_2$	High-energy ball milling	$Pa\bar{3}$	716 mA h g^{-1} at 2.5 A g^{-1} current density	67% after 200 cycles	148

imaging to atomically identify phase transformations during the whole lithiation process of a rock-salt HEO electrode $(Co_{0.2}Ni_{0.2}Mn_{0.2}Zn_{0.2}Fe_{0.2})_3O_{3.2}$ ($M_3O_{3.2}$).[113] The pristine ED pattern (Fig. 5(a)) of the HEO electrode shows eight distinct diffraction rings which were indexed to the rock-salt structure (ICDD PDF#01-078-0431). At the initial lithiation stage, an emerging reflection from the Li_2O (220) plane was generated. The eight sharp diffraction rings in Fig. 5(a) gradually vanish and the HEO (311) reflection completely disappeared in Fig. 5(b). Ni, Co, and Mn metal phases also appear sequentially (Fig. 5(c)–(e)). The diffraction intensities of $M_3O_{3.2}$ were further weakened and the (422) and (400) reflections vanish. At the final lithiation stage (Fig. 5(f)), the pristine rock-salt structure is not observed. A new phase of $Li_{0.2}Zn_{0.8}$ with two distinguishable (100) and (101) planes was generated. As shown in Fig. 5(g), the quantitative intensity profiles from time-sequenced ED patterns are plotted to better visualize the structural transformations throughout the lithiation process, the characteristic peaks attributed to the HEO rock-salt phase vanishes, while some broad new peaks indexed as Fe, Ni, Co, Mn, and $Li_{0.2}Zn_{0.8}$ appeared in succession, which is highly consistent with the time-sequenced ED patterns. To further verify the formation of these emerging crystallites, the $M_3O_{3.2}$-based half cells discharged to 0.01 V were disassembled and then were characterized by TEM. As shown in Fig. 5(h)–(l), HRTEM images and their corresponding fast Fourier transformation (FFT) also confirm the occurrence of these new phases. Thus, the detailed *in situ* and *ex situ* TEM observations suggest that HEO anode undergoes complicated conversion and alloying reactions during the lithiation process, and this results in a polycrystalline structure consisting of Fe, Ni, Co, Mn, $Li_{0.2}Zn_{0.8}$ and Li_2O phases.

Fig. 6(a) and (b) shows the comparative TEM images between a fully lithiated and delithiated HEO nanoparticle (NP) in the first discharge–charge process. The fully delithiated NP recovers to its pristine morphology with almost the same particle size observed. The phase transformations during delithiation were identified by *in situ* ED patterns (Fig. 6(c)–(h)). Upon an initial delithiation, $Li_{0.2}Zn_{0.8}$ phase first disappeared, which indicates that the dealloying reaction process occurred. Simultaneously, the diffraction ring of (200) plane re-occurred, suggesting that rock-salt HEO structure is recovered, With the further delithiation process, the diffraction rings arising from Co, Mn, Ni, Li_2O phases are diminished, which is accomplished by emergence of the (420), (311), and (222) reflections attributed to the HEO rock-salt phase. An unexpected symmetrical structural evolution during initial discharge–charge cycle is noticed in Fig. 6(i), demonstrating that the conversion and alloying reactions are highly recoverable and thus enabling reversible transformation of the promoting efficient lithium storage.

Spinel-type HEOs have been utilized as electrode materials for LIBs. This unique structure with three-dimensional Li^+ transport pathways can potentially facilitate Li^+ conductivity and thus improve lithiation kinetics. The two Wyckoff sites and

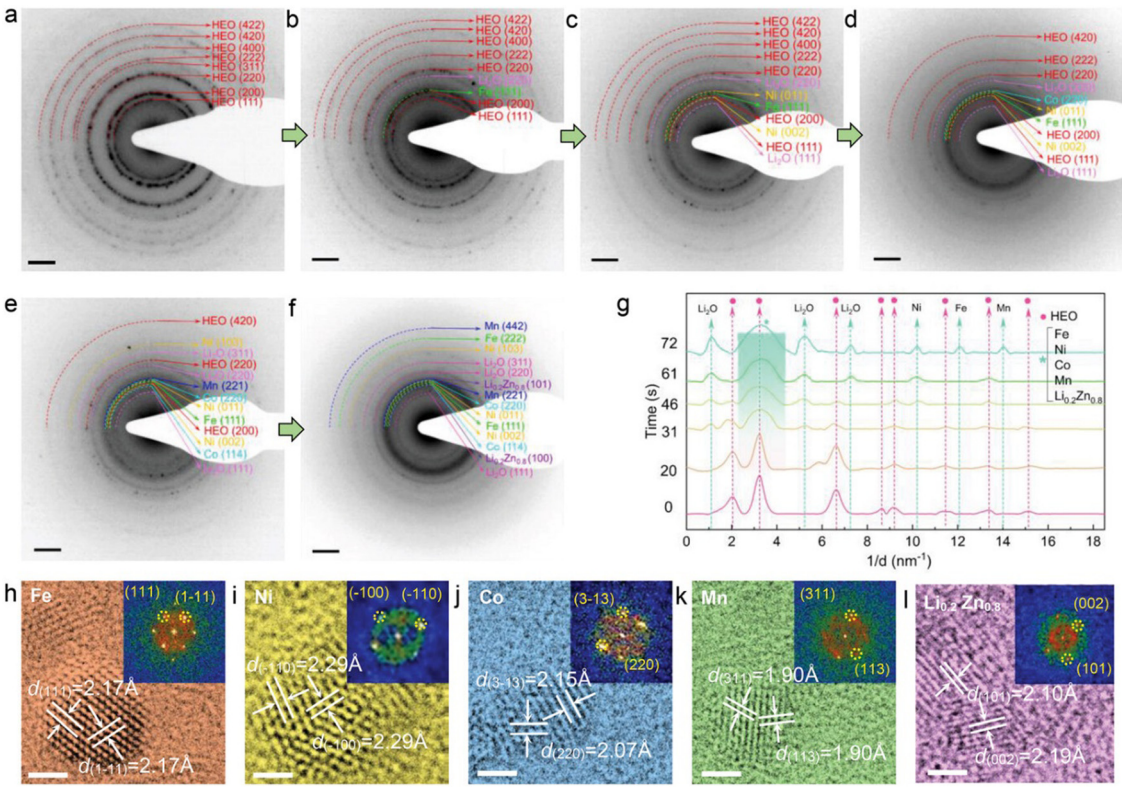

Fig. 5 (a)–(f) *In situ* electron diffraction (ED) and HRTEM imaging during lithiation. (g) The intensity profiles from ED patterns at different lithiation stages. (h)–(l) *Ex situ* HRTEM images and corresponding FFT patterns of lithiated $M_3O_{3.2}$. Scale bar 2 nm. Copyright permission from Wiley, used from ref. 113.

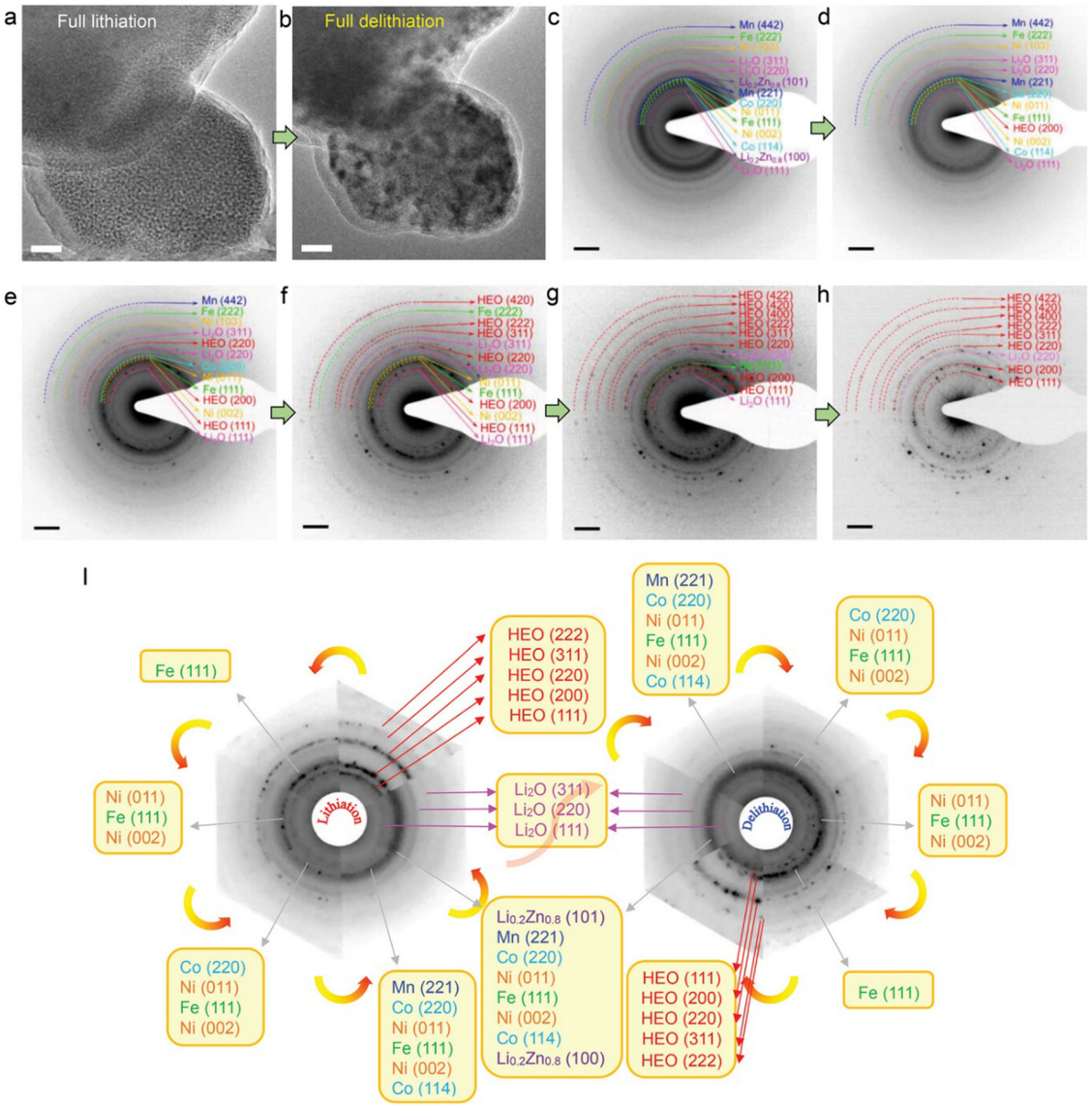

Fig. 6 Structural reversibility of the HEO structure upon delithiation. (a) and (b) TEM morphology evolution of a fully delithiated HEO nanoparticle. (c)–(h) ED patterns from different delithiation stages. (i) Schematic diagram of the ED pattern evolutions during the first discharging/charging cycle. Copyright permission from Wiley, used from ref. 113.

rich oxygen vacancies are favorable for reversible Li^+ storage. Ting *et al.* reported a surfactant-assisted hydrothermal method toward a spinel HEO $(CrMnFeCoNi)_3O_4$ anode for LIBs, which was found to have a desirable specific capacity of 1235 mA h g^{-1} and respectable capacity retention of 90% after 200 cycles.[140] Most of the reported HEO anodes for LIBs contain Co which mainly contributes to high conductivity and reversible capacity. However, Co is expensive, and large quantities of global cobalt is mined in the Democratic Republic of the Congo, which suffers from political instability and ethical production methods (due to civil war, other violence, and child labour).[153,154] Therefore, the development of Co-free HEO anodes is of great interest.

Chang *et al.* reported a Co-free spinel $(CrMnFeNiCu)_3O_4$ HEO and applied *operando* X-ray absorption spectroscopy (XAS) to study ion intercalation mechanisms.[49] In comparison to conventional XAS that is operated in stepping mode, quick-scanning XAS with a monochromator oscillation frequency of 2 Hz has better signal-to-noise ratio sensitivity as well as superior time-resolution. Fig. 7 exhibits the various K-edge spectra for all the metallic elements in the HEO anode during the initial discharging-charging process. Upon Li^+ uptake, the absorption edges of all the spectra shifted towards lower energy, suggesting that reduction happens during lithiation. After full lithiation at 0.01 V, *operando* spectra of all constituent elements are close to the plots of metal standards, indicating that the metallic Mn, Cu, Cr, Fe, and Ni formed at the end of lithiation process. At the delithiation stage, progressive shifts towards the positive side were observed in Mn, Cr, and Fe spectra, suggesting the occurrence of oxidation reactions.

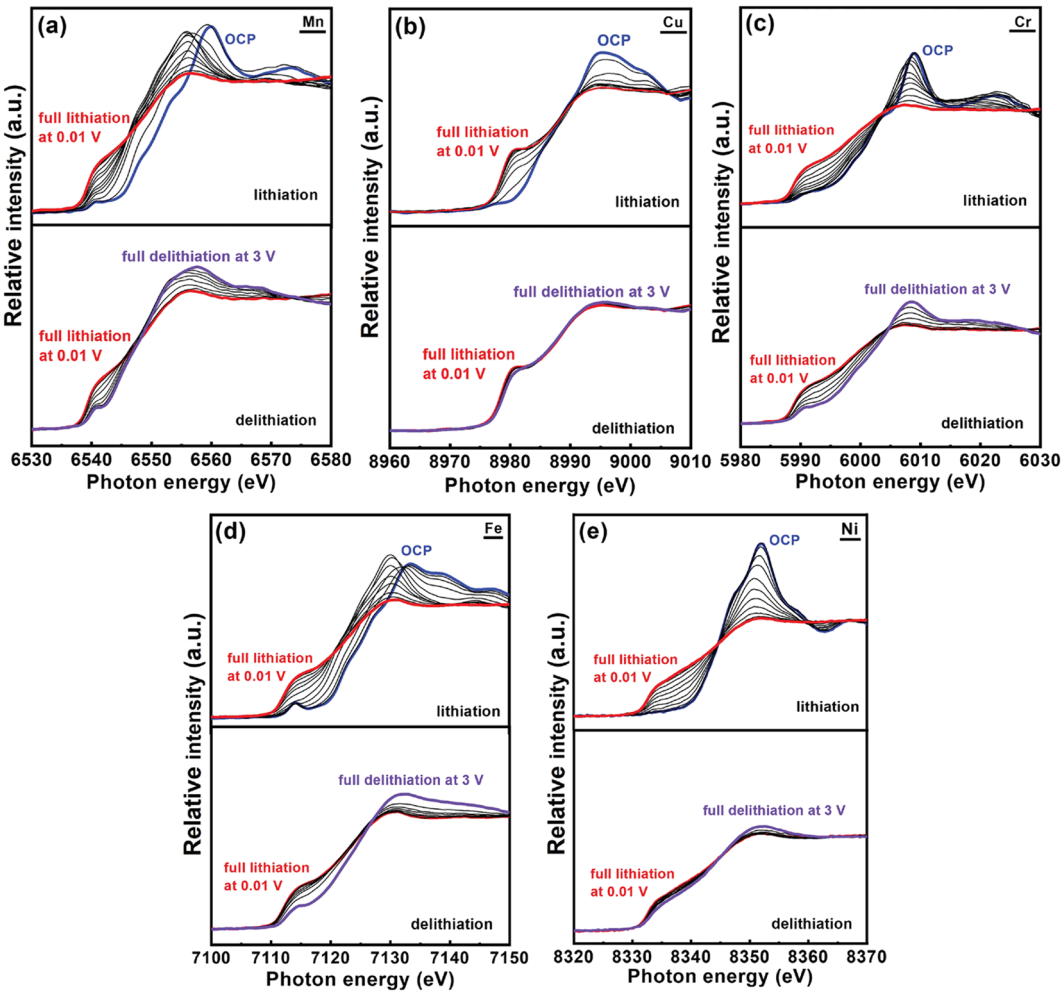

Fig. 7 Operando XAS spectra of HEO anode obtained during first lithiation/delithiation process at 150 mA g^{-1} current density. The interval between spectra is 100 mA h g^{-1}. Copyright permission from Wiley, used from ref. 49.

A minor shift was noticed in the Ni spectra. It is worthwhile to note that the absorption edge of Cu was almost unchanged throughout the delithiation, which is indicative of high irreversibility.

Detailed coordination structure transformations of all the metallic elements present in HEO anodes were examined using first-derivative curves of the K-edge spectra. As shown in Fig. 8(a), the two peaks around 6539 and 6546 eV were ascribed to Mn^0 in metallic Mn and $Mn^{2+/3+}$ in HEO, respectively. After 250 mA h g^{-1}, the Mn^0 signal gradually becomes stronger while the $Mn^{2+/3+}$ peak becomes weaker, suggesting that the conversion reaction between $Mn^{2+/3+}$ and Mn^0 occurs. Cu^{2+} undergoes a single-step conversion to form Cu^0 (Fig. 8(b)). The Cu^{2+} characteristic peak located of 8983.9 eV is diminished up to lithiation of $\approx$250 mA h g^{-1} while the metallic Cu signal at 8979 eV became intensified. After this, the spectral features are almost unchanged. According to Fig. 8(c), two reduction reactions were observed during the charging and discharging processes; Cr^{3+} is reduced to Cr^{2+} before 0.44 V, then Cr^{2+} is reduced to metallic Cr^0 from 675 and 1200 mA h g^{-1}. $Fe^{2+/3+}$ also displays two-step reduction with the reduction of Fe^{3+} to Fe^{2+} followed by reduction of Fe^{2+} to metallic Fe^0 (Fig. 8(d)). Ni underwent a single-step conversion of Ni^{2+} to metallic Ni^0 eventually (Fig. 8(e)). The redox behavior of the elements in the HEO anode is summarized in Fig. 8(f). At the beginning of the lithiation process, reduction reactions including $Mn^{3+} \rightarrow Mn^{2+}$ and $Cu^{2+} \rightarrow Cu^0$ are mainly responsible for the capacity contribution. The transitions $Mn^{2+} \rightarrow Mn^0$, $Cr^{3+} \rightarrow Cr^{2+}$, $Fe^{3+} \rightarrow Fe^{2+}$, $Fe^{2+} \rightarrow Fe^0$, and $Ni^{2+} \rightarrow Ni^0$ give rise to a charging potential plateau at around 0.5 V. The final Li^+ storage charge-compensation was mainly resulted from $Cr^{2+} \rightarrow Cr^0$ conversion. As exhibited in Fig. 8(g), at the initial delithiation process, the measured capacity was largely originated from the oxidation reaction of metallic Mn. After $\approx$1.13 V, both Cr and Fe started to oxidize. The Cu reaction was almost invisible and provided the least capacity. The oxidation of Ni needed a large overpotential of $\approx$1.67 V, contributing to capacity only at the last stage.

HEO materials have also been reported as cathode materials for LIBs. Brezesinski *et al.* reported a high-entropy rock-salt-type fluoride oxide $Li_{0.94}(Co_{0.21}Ni_{0.21}Zn_{0.2}Mg_{0.2}Cu_{0.18})O_1F_{0.87}$ and layered oxide $LiNi_{0.2}Mn_{0.2}Co_{0.2}Fe_{0.2}Al_{0.2}O_2$, which displays

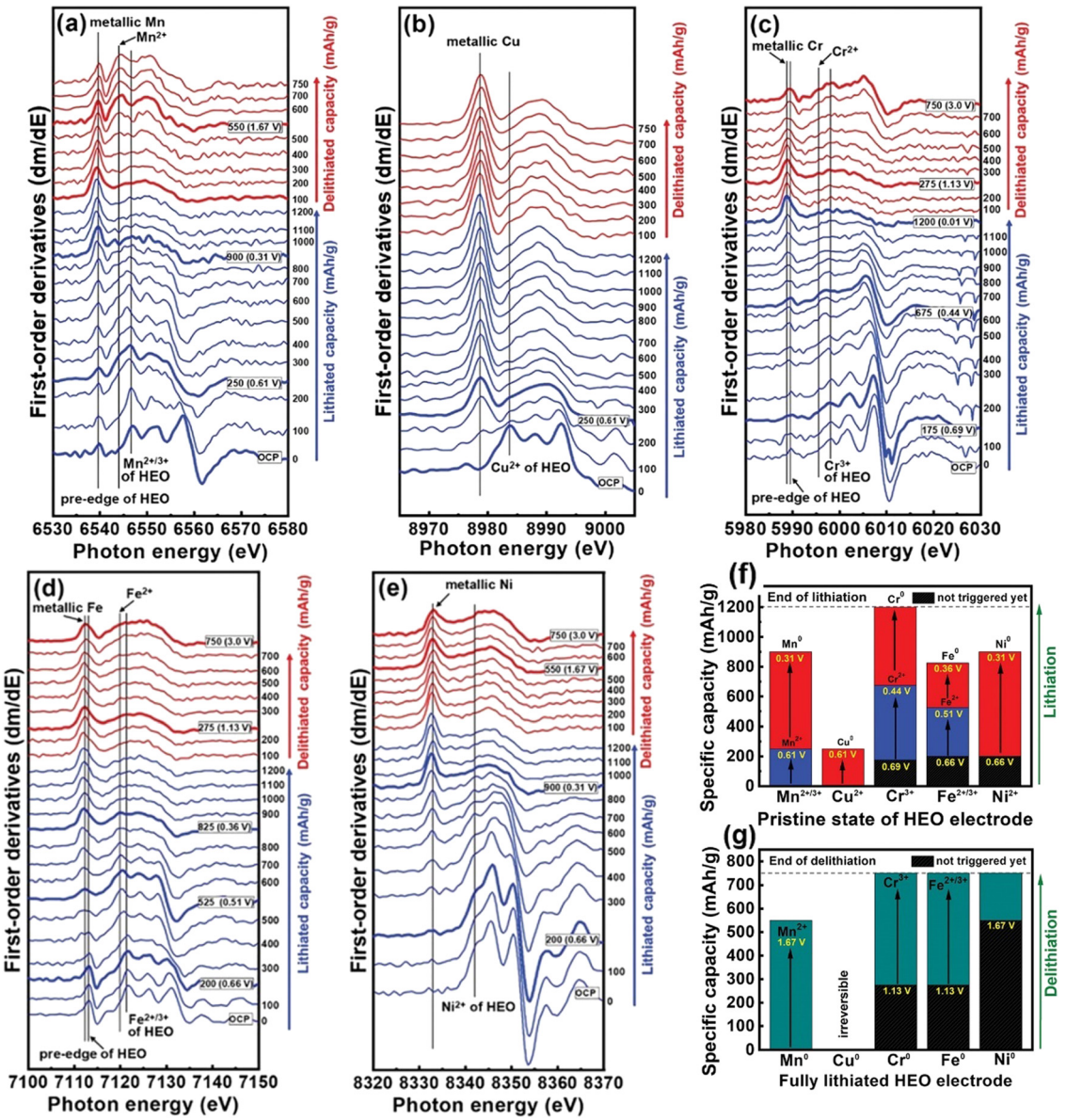

Fig. 8 First-derivative plots of (a) Mn, (b) Cu, (c) Cr, (d) Fe, and (e) Ni XAS operando spectra in $(CrMnFeNiCu)_3O_4$ HEO anode during first lithiation and delithiation cycle at 150 mA g^{-1} current density. (f)–(g) Capacity contribution of all constituent elements in HEO electrode during first lithiation and delithiation cycle. Copyright permission from Wiley, used from ref. 49.

excellent Li^+ storage capability.[147] Ceder *et al.* applied the HE concept to the recently developed cation-disordered rocksalt (DRX) cathodes, which was found to be an effective method to suppress cation short-range order (SRO) due to the mixing of a larger number of TM species.[146] The SRO reduction is favorable for Li^+ storage capability. Three prototype compositions were designed with different numbers of TM species in order to assess their HE DRX strategy. $Li_{1.3}Mn_{0.4}Ti_{0.3}O_{1.7}F_{0.3}$ was regarded as reference sample. Mn, Nb, Co, and Cr elements were successively incorporated into the former structure to form $Li_{1.3}Mn_{0.2}Mn_{0.2}Ti_{0.1}Nb_{0.2}O_{1.7}F_{0.3}$ and $Li_{1.3}Mn_{0.1}Co_{0.1}Mn_{0.1}Cr_{0.1}Ti_{0.1}Nb_{0.2}O_{1.7}F_{0.3}$. The three compositions were referred to as TM2, TM4, and TM6, respectively, based on the various number of TEM species present in each structure.

As shown in Fig. 9(a), 30% Li excess was included in all three compositions to facilitate Li intercalation without undermining the TM redox capacity. 15% fluorine was incorporated to substitute oxygen to ensure a good redox reservoir. Scanning electron microscopy (SEM) revealed that the particle size of TM6 samples is approximately 5–10 μm (Fig. 9(b)) and can be decreased to 200–500 nm by shaker milling with carbon during electrode fabrication. The refined XRD patterns of as-prepared samples are shown in Fig. 9(c). The lattice constants (*a*) of cubic TM2, TM4 and TM6 are 4.1918, 4.2286 and 4.2655 Å, respectively, indicating that lattice parameter increased with more TM species in the structure. The STEM-EDX elemental maps shown in Fig. 9(d) demonstrated that all the incorporated elements are uniformly distributed throughout the material.

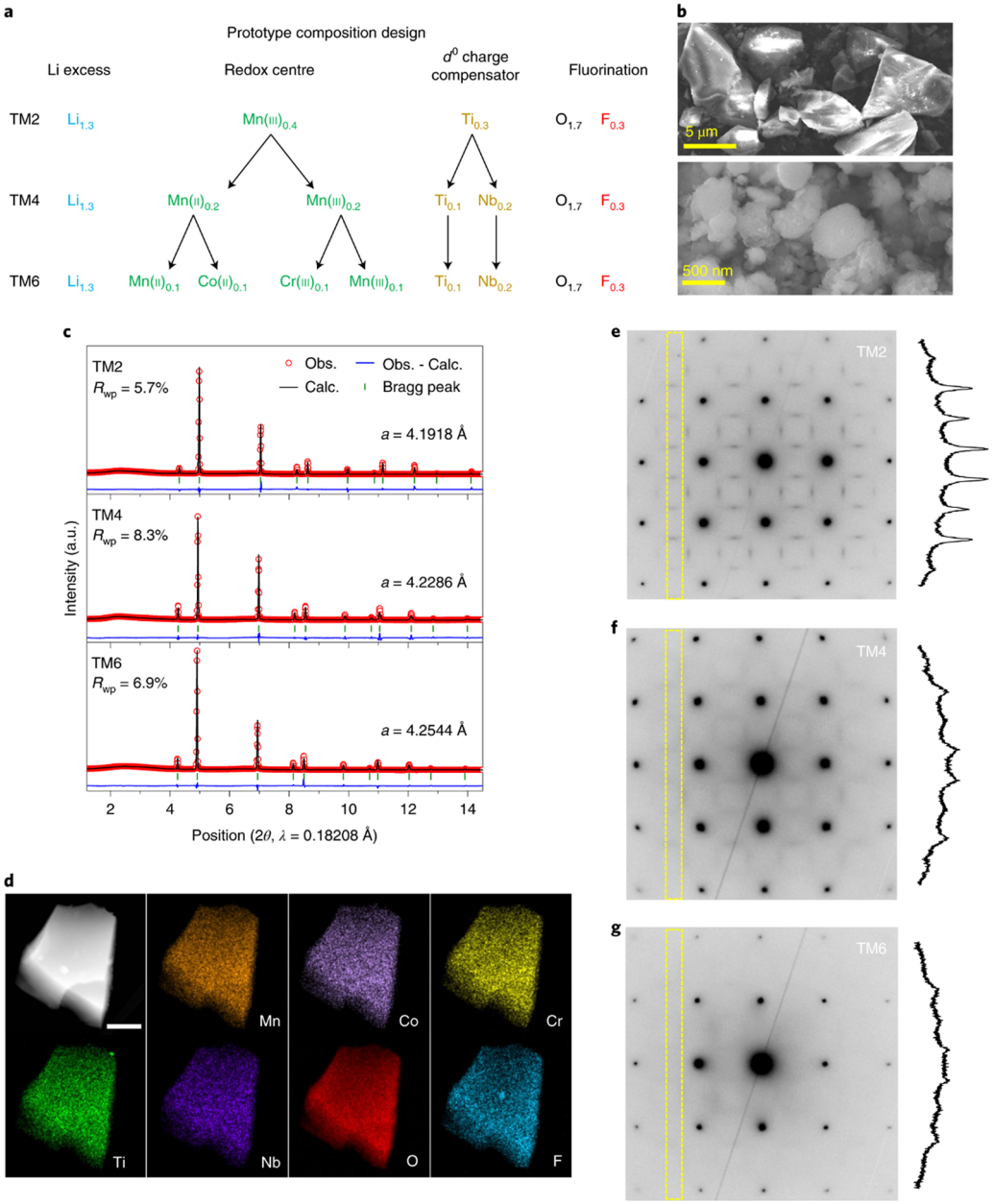

Fig. 9 (a) Composition design of HEO cathodes with variations of TM ions, (b) SEM images of as-synthesized TM6, (c) synchrotron XRD patterns and refined lattice parameters of the as-prepared HEOs, (d) HAADF-STEM EDX mapping of elemental distribution in a particle cluster of TM6. Scale bar, 300 nm. (e)–(g) TEM ED patterns measured on TM2 (e), TM4 (f) and TM6 (g) along the [100] zone axis. Copyright permission from Nature materials, used from ref. 146.

Electron diffraction patterns in the TEM was used to evaluate the cation short-range order (SRO, Fig. 9(e)–(g)). The diffraction spots can be indexed to cubic $Fm\bar{3}m$ rock salt structures. The strength of the SRO is qualitatively associated with the intensity of the diffuse scattering, which is integrated within the dashed rectangular areas. The results clearly show that the intensity of the diffuse scattering decreased with an increase in the number of TM species, indicating that higher entropy is advantageous to suppress SRO in the DRX structure.

The electrochemical performance of the three electrodes was evaluated using galvanostatic cycling. Within a voltage window from 1.5 to 4.7 V, the TM2 electrode delivered a specific capacity of 220 mA h g^{-1}, which increased to 220 mA h g^{-1} for TM4 and to 307 mA h g^{-1} for TM6. These rate capability results are shown in Fig. 10(a)–(f). The TM2 electrode initially delivers a specific capacity of 220 mA h g^{-1} at 20 mA g^{-1} current density, which rapidly decreases to 58 mA h g^{-1} at a higher current density of 2 A g^{-1} with a 74% capacity loss. With more TM ions in the DRX structure, the capacity loss at high

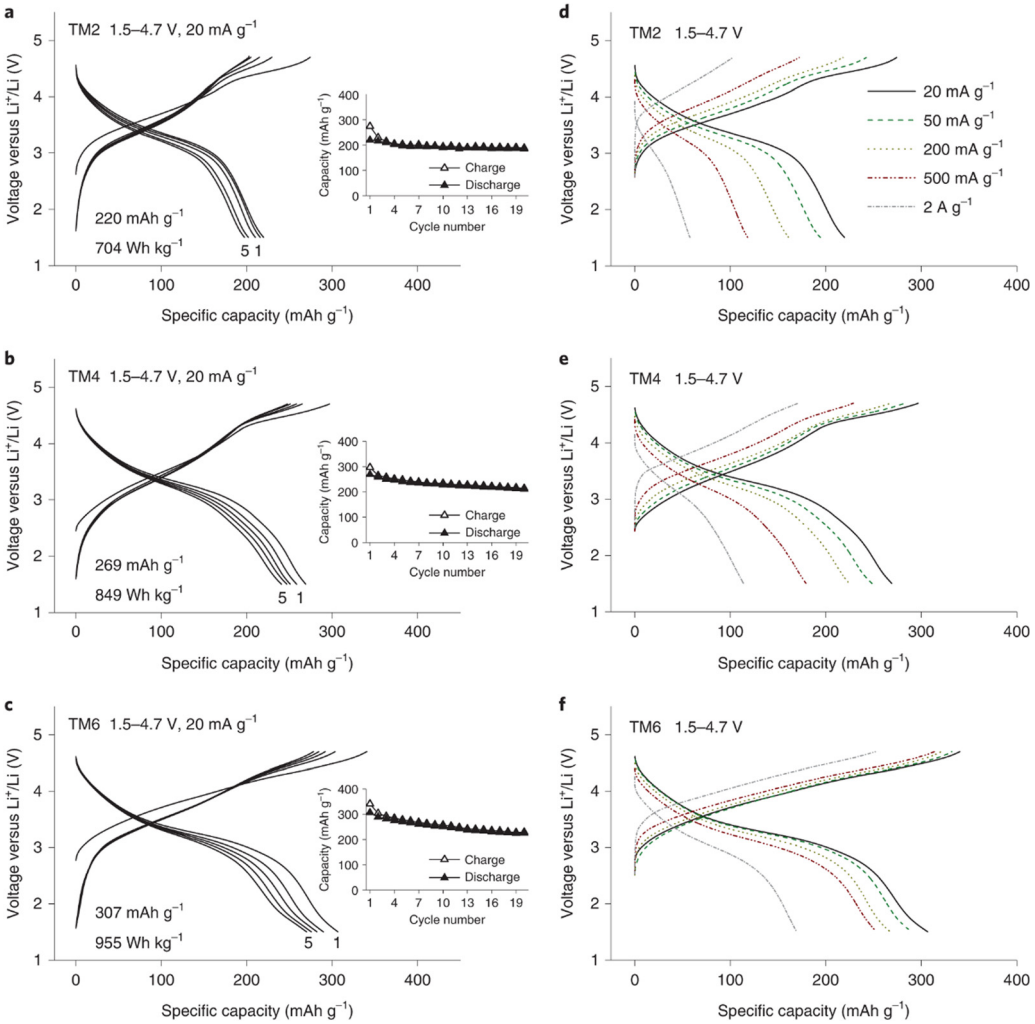

Fig. 10 Voltage curves of TM2 (a), TM4 (b) and TM6 (c) within the potential scope of 1.5–4.7 V at 20 mA g^{-1} current density (insets show the capacity retention plots), rate performance of TM2 (d), TM4 (e) and TM6 (f) exhibiting the 1st voltage curves between 1.5 and 4.7 V at 20, 50, 200, 500 and 2000 mA g^{-1}. Copyright permission from Nature materials, used from ref. 146.

current rate is reduced to 58% for TM4 and to 45% for TM6, respectively. The TM6 electrode still delivers a capacity of 170 mA h g^{-1} even at a high current density of 2 A g^{-1}. The excellent electrochemical performance of TM6 electrode supports the authors' design strategy of the application of the HE concept in the DRX electrode structure, which could potentially reduce SRO and thus improve the Li^+ storage capability in these materials.

Except for HEOs, HE chalcogenides have also displayed great electrochemical performances when used as electrodes for LIBs. Recently, Breitung *et al.* utilized HE sulfides as the anode in LIBs.[60] A one-step mechanochemical reaction was used to produce different HE sulfides including (FeMnNiCoCr)S_2, (FeMnNiCoCu)S_2, (FeMnNiTiCr)S_2 and (FeMnNiCoCr)S. As shown in Fig. 11(a), the (FeMnNiCoCr)S_2 electrode delivers the best electrochemical performance with a high reversible capacity of 716 mA h g^{-1} at a high current density of 2.5 A g^{-1}. In addition, the (FeMnNiCoCr)S_2 electrode exhibited excellent cycling capability (Fig. 11(b)) and long-term cycling performance (Fig. 11(c)).

HEMs for sodium-ion batteries (SIBs)

The layered transition metal (TM) oxide Na_xTMO_2 ($x \leq 1$) is regarded as one of the most promising cathode materials for SIBs due to its high energy density, which can be structurally categorized into O3, P3 and P2 types (where P and O represents prismatic and octahedral crystal structures, respectively, and 2 or 3 represents the minimum number of transition metal layers in the repeating cell unit).[155,156] However, irreversible structural changes during the Na^+ intercalation process detrimentally affects the electrochemical performance, which is the primary cause of capacity fading.[157] It has been reported that the introduction of single or multiple heterogenous metals into Na_xTMO_2 is an effective approach to stabilize the host structure and alleviate the cationic migration from the TM layers to the Na layers.[103] This strategy can be extended to the design of HEMs containing multiple elements in a single-phased structure.[103,158] A selection of previously-reported HEMs as electrode materials for SIBs are summarized in Table 2.

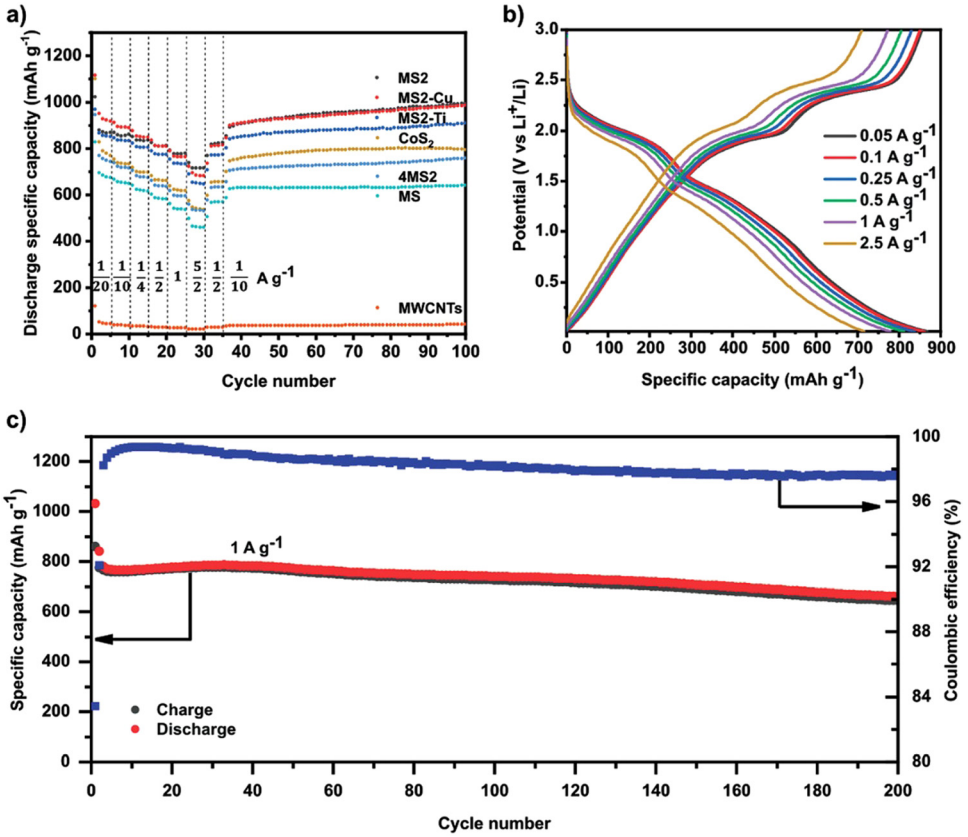

Fig. 11 (a) Galvanostatic rate performance test of all HESs, CoS_2, $4MS_2$, and MWCNTs half-cell at different current densities and 25 °C in the voltage range between 0.01 and 3 V *vs.* Li^+/Li. (b) Voltage profiles of MS_2 half-cell at different currents. (c) Specific charge/discharge capacity of MS_2 half-cell and Coulombic efficiency as a function of cycle number at 1 C and 25 °C in the range between 0.01 and 3 V. Copyright permission from Wiley, used from ref. 60.

Tobola *et al.* utilized a combination of high-temperature solid-state and low-temperature sol–gel methods to produce an O3-type layered six-component HEO $Na_1Ti_{0.16}Mn_{0.16}Fe_{0.16}Co_{0.16}Ni_{0.16}Cu_{0.16}O_2$.[159] It was found that the presence of copper is favorable for improvement of ion conductivity and structural stability. Cheng reported a P2/O3 biphasic $Na_{0.7}Mn_{0.4}Ni_{0.3}Cu_{0.1}Fe_{0.1}Ti_{0.1}O_{1.95}F_{0.1}$ with a remarkable initial coulombic efficiency (ICE), which is also able to deliver a respectable discharge capacity of 86.7 mA h g^{-1} at a current density of 0.8 A g^{-1} with a high ICE of 97.6%.[161] Density functional theory (DFT) calculations revealed that the excellent electrochemical performance potentially originates from reversible structural evolution and fast Na^+ diffusion kinetics.

Xin *et al.* reported the HEO O3-structure $Na_xTi_{0.16}Mn_{0.16}Fe_{0.16}Co_{0.16}Ni_{0.16}Cu_{0.17}O_2$.[171] The XRD pattern (Fig. 12(a)) shows that this material is comprised of a pure O3 phase. Two characteristic peaks at $2\theta \approx 18.5^\circ$ and 44.1° suggested that a honeycomb-type superlattice phase is formed due to the ordering between Li ions and TM ions caused by a large difference in ionic radii. A uniform plate-like morphology was observed in SEM imaging (Fig. 12(c)). As shown in Fig. 12(d) and (e), HAADF-STEM imaging further confirmed the O3 atomic structure. The SAED pattern projected along the [001] zone axis also revealed the diffraction pattern of a layered O3 phase (Fig. 12(f)). Some diffraction spots with relatively weaker contrast around the bright spots can be noticed from the perspective of the [010] zone axis, which is in accordance with superlattice formation (Fig. 12(g)). EDX spectroscopic mapping shows that there is uniform elemental distribution throughout the particles at the nanoscale (Fig. 12(h)).

The charge/discharge curves in Fig. 13(a) show that NaLFCNM cathode delivers considerable capacity retention capability with a slight voltage fading when cycling between 2.0–4.5 V at 0.3 C current density. Long-term cycling performances at higher density are exhibited in Fig. 13(b) and (c), which both delivers excellent cycling capability. The rate performance is very respectable, as shown in Fig. 13(d); even at a high current density of 10 C, a desirable discharge capacity of 78.2 mA h g^{-1} is still achieved, which accounts for 45.3% to the discharge capacity of 0.1 C. This rate capability is potentially facilitated by fast Na^+ diffusion kinetics of about *ca.* 10^{-11} cm^2 S^{-1}[19] obtained from galvanostatic intermittent titration technique (GITT) measurements (Fig. 13(e)).

The structural evolution during the sodium intercalation process was investigated by *in situ* synchrotron XRD experiments. As shown in Fig. 14(a), when charging from open-circuit voltage (OCV) to 3 V during the first cycle, the original O3 phase gradually vanished, and a P3-type structure is observed. At the

Table 2 HEMs as electrode materials for SIBs

HEMs	Synthesis method	Structure	Electrochemical performance	Capacity retention	Ref.
$Na_xTi_{0.16}Mn_{0.16}Fe_{0.16}Co_{0.16}Ni_{0.16}Cu_{0.17}O_2$	Sol–gel method	O3/P3 phase	139 mA h g^{-1} at 0.1 C current density	~100% after 10 cycles	159
$NaNi_{0.4-x}Cu_xFe_{0.2}Mn_{0.2}Ti_{0.2}O_2$	Solid-state reaction	O3-type phase	129 mA h g^{-1} at 24 mA g^{-1} current density	74.83% after 50 cycles	160
$Na_{0.7}Mn_{0.4}Ni_{0.3}Cu_{0.1}Fe_{0.1}Ti_{0.1}O_{1.95}F_{0.1}$	High-temperature sintering	P2/O3 phase	134 mA h g^{-1} at 20 mA g^{-1} current density	93.3% after 200 cycles	161
CoZnCuNiFeZrCeO$_x$-PMA	Solvo-thermal method	Fluorite structure	220 mA h g^{-1} at 300 mA g^{-1} current density	92% after 5000 cycles	162
$Na_{0.67}Li_{0.16}Fe_{0.16}Co_{0.16}Ni_{0.16}Mn_{0.33}O_2$	Solid-state reaction	Superlattice structure	171 mA h g^{-1} at 0.1 C current density	89.3% after 300 cycles	103
$Na[Fe_{0.2}Co_{0.2}Ni_{0.2}Ti_{0.2}Sn_{0.1}Li_{0.1}]O_2$	Solid-state reaction	O3-type phase	81 mA h g^{-1} at 2.0 C current density	81% after 100 cycles	163
$NaCu_{0.1}Ni_{0.3}Fe_{0.2}Mn_{0.2}Ti_{0.2}O_2$	Pechini[a] method	O3-type phase	130 mA h g^{-1} at 0.1 C current density	71% after 500 cycles	164
$Na_{0.62}Mn_{0.67}Ni_{0.23}Cu_{0.05}Mg_{0.07}Ti_{0.01}O_2$	Hydrothermal method	P2-type phase	148 mA h g^{-1} at 12 mA g^{-1} current density	75% after 2000 cycles	165
$NaNi_{0.25}Mg_{0.05}Cu_{0.1}Fe_{0.2}Mn_{0.2}Ti_{0.1}Sn_{0.1}O_2$	High-temperature sintering	O3-type phase	131 mA h g^{-1} at 14 mA g^{-1} current density	75% after 500 cycles	72
$NaMn_{0.2}Fe_{0.2}Co_{0.2}Ni_{0.2}Ti_{0.2}O_2$	High-temp-erature sintering	O3/P3 phase	180 mA h g^{-1} at 0.1 C current density	97% after 100 cycles	166
$Na_{0.6}(Ti_{0.2}Mn_{0.2}Co_{0.2}Ni_{0.2}Ru_{0.2})O_2$	High-temp-erature sintering	P2 phase	164 mA h g^{-1} at 0.1 C current density	99.78% after 40 cycles	167
$NaNi_{0.12}Cu_{0.12}Mg_{0.12}Fe_{0.15}Co_{0.15}Mn_{0.1}Ti_{0.1}Sn_{0.1}Sb_{0.04}O_2$	High-temp-erature sintering	O3-type phase	110 mA h g^{-1} at 0.1 C current density	83% after 500 cycles	103
$Na_3VAl_{0.2}Cr_{0.2}Fe_{0.2}In_{0.2}Ga_{0.2}(PO_4)_3$	Sol–gel method	NASICON-structure	102 mA h g^{-1} at 11 mA g^{-1} current density	86.8% after 5000 cycles	40
$Na_3(Ti_{0.2}V_{0.2}Mn_{0.2}Cr_{0.2}Zr_{0.2})_2(PO_4)_3$	Sol–gel method and solid-state reaction	NASICON-structure	96.5 mA h g^{-1} at 100 mA g^{-1} current density	36.3% after 100 cycles	168
$Na_3V_{1.9}(Ca,Mg,Al,Cr,Mn)_{0.1}(PO_4)_2F_3$	Sol–gel method and solid-state reaction	NASICON-structure	118 mA h g^{-1} at 0.1 C current density	80.4% after 2000 cycles	42
$Na_{1.41}Mn_{0.32}Fe_{0.11}Co_{0.28}Ni_{0.32}Cu_{0.32}[Fe(CN)_6]_{2.89}H_2O$	Co-precipitation	Cubic structure	92.6 mA h g^{-1} at 150 mA g^{-1} current density	95% after 10 000 cycles	43
$Na_{1.70}Fe_{0.2}Mn0_{.2}Co_{0.2}Ni_{0.2}Cu_{0.2}$ $[Fe(CN)_6]_{0.98}\square_{0.02}{\cdot}2{\cdot}35H_2O$ ($\square$ = $[Fe(CN)_6]$ vacancy)	Co-precipitation	Cubic structure	115 mA h g^{-1} at 100 mA g^{-1} current density	79.6% after 1000 cycles	44
$Na_x(FeMnNiCuCo)[Fe(CN)_6]$	Co-precipitation	Cubic structure	120 mA h g^{-1} at 10 mA g^{-1} current density	~100% after 3000 cycles	169

[a] The Pechini method of synthesis is a modified sol–gel method which utilizes cation complexation with citric acid in a water–ethanol solution to form a homogenous mixture of precursors. The pH was adjusted to 5 and ethylene glycol was added and solution annealed at 120 °C for at least 12 hours to leave a porous gel and calcinated twice, once at 480 °C for 6 h and again at 850 °C for 12 h in air.[170]

end of charge (cut-off voltage: 4.5 V), co-existence of the P3/O3 biphasic system can be observed, which was confirmed by the P3 + O3 lattice sequence in the HRTEM image (Fig. 14(d)). Upon the subsequent Na^+ deintercalation process, the P3-type phase is retained until discharging to around 2.6 V. The P3 phase gradually recovers back to O3 phase around 2.0 V, which was also verified by the HRTEM image (Fig. 14(e)). The lattice parameters at different Na^+ (de)intercalation stages which were obtained from the *in situ* XRD patterns are shown in Fig. 14(c), which shows that the deviation of *a* and *c* parameters during initial cycling is 1.84% and 1.97%, respectively. This small volume change of NaLFCNM indicates a minor structural change during the Na^+ intercalation process and is responsible for the excellent cycling performance. Fig. 14(f) schematically illustrates the structural evolution during initial charge/discharge process.

Sodium super ionic conductor (NASICON) materials have the advantages of 3D ion transport channels, impressive structural robustness and a large number of Na^+ vacancies, which can significantly alleviate the structural evolution of repeating Na^+-insertion/-extraction processes and facilitate Na^+ diffusion kinetics, enabling them to be competitive as SIBs[172–174] Zhao *et al.* designed the high-entropy NASICON-type material $Na_3VAl_{0.2}Cr_{0.2}Fe_{0.2}In_{0.2}$ $Ga_{0.2}(PO_4)_3$ as a SIB cathode.[40] It was found that the reversible redox reaction of V^{4+}/V^{5+} could be realized and the irreversible phasic transformation above 4.0 V *vs.* Na^+/Na could be effectively depressed. A high energy density of 444 W h kg^{-1} was achieved, which outperformed $Na_3VAl(PO_4)_3$ and $Na_3VCr(PO_4)_3$.

$Na_3(VO_{1-x}PO_4)_2F_{1+2x}$ ($0 \leq x \leq 1$) is of great interest due to its high specific capacity.[175] The electrochemical performance can be tailored by adjusting the F/O ratio.[176] Wu *et al.* reported the high-entropy phosphate cathode $Na_3V_{1.9}(Ca,Mg,Al,Cr,Mn)_{0.1}(PO_4)_2F_3$ (HE-NVPF).[42] Electrochemical tests showed that HE-NVPE delivers remarkably improved energy density, cycling performance and rate capability compared to pristine NVPF (P-NVPF) and other fluorophosphate counterparts.

As shown in Fig. 15(a), three pairs of redox peaks (A1/C1, A2/C2, and A3/C3) in the CV curve can be identified in p-NVPF cathode, while the A1/C1 peaks at low-voltage region disappeared for HE-NVPF cathode, which is in accordance with the corresponding GCD plots in Fig. 15(b). More importantly, the two couples of redox peaks in the CV curves in HE-NVPF are sharper and larger than that of p-NVPF electrode. As shown in Fig. 15(c), the capacity contribution above 3.7 V working voltage for HE-NVPF accounts for 45.7%, which is much higher than the 38.3% for p-NVPF. Indeed, the absence of a 3.4 V plateau region for HE-NVPF cathode improves the working voltage and thus achieves a higher specific capacity, further delivering higher energy density. The HE-NVPF cathode electrode also delivers improved rate performance (Fig. 15(d) and (e)), and the specific capacity for HE-NVPF at different current densities ranging from 0.1 C to 50 C is higher than that of the p-NVPF electrode. Ragone plots (Fig. 15(f)) show that the HE-NVPF displays significant improvement both in energy density and

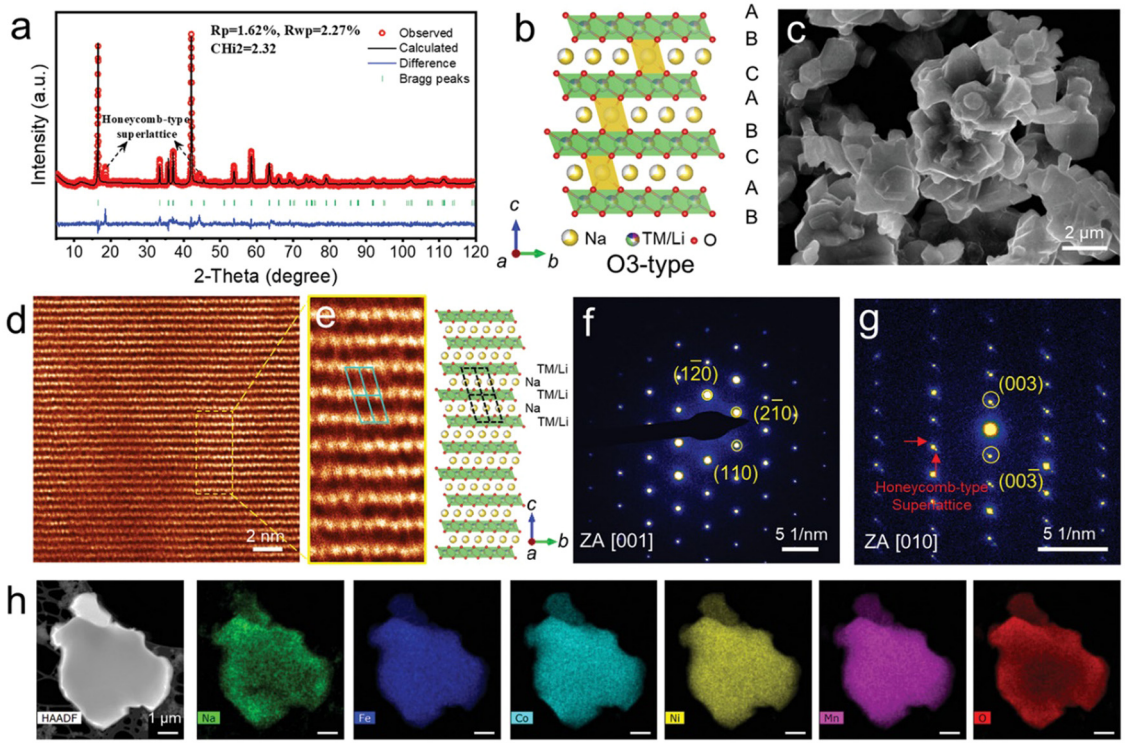

Fig. 12 Structure characterization of pristine NaLFCNM materials. (a) Rietveld refinement of X-ray diffraction (XRD) pattern. (b) Schematic O3-type crystal structure with the presence of superlattice phase. (c) SEM image. (d) HAADF-STEM image of the [001] zone axis. (e) HAADF-STEM image (left) of the area selected from (d) and the corresponding structure model (right). (f) and (g) SAED patterns projected from the zone axis of (f) [001] and (g) [010] zone axis. (h) STEM-EDX spectroscopic mapping. Copyright permission from Wiley, used from ref. 171.

power density in comparison to p-NVPF. In addition, the long-term cycling performance test in Fig. 15(g) and (h) suggests that the capacity ratio of HE-NVPF electrode is significantly higher than p-NVPF, which achieves 90.2% and 80.4% at 0.5 C and 20 C, respectively, after 400 cycles.

Prussian blue analogues (PBAs) are transition metal hexacyanoferrates with a stable open channel structure, ion diffusion paths and redox-active sites.[177] These intrinsic properties make them a competitive cathode material for SIBs, usually having high reversible capacity, excellent rate performance and high working voltage.[178,179] The general chemical formula of PBAs can be represented as $A_xM[Fe(CN)_6]_y\square_m{\cdot}nH_2O$ (with $0 < x \leq 2$, $0 < y \leq 1$, $y + m = 1$), where A is an alkali metal ion, M and □ represent transition-metal ions and $[Fe(CN)_6]$ vacancies, respectively.[180] Brezesinski *et al.* applied the high-entropy approach to the PB structure $Na_x(FeMnNiCuCo)[Fe(CN)_6]$ (simplified as HE-PBA), where five different transition metals are occupied in the M site in equimolar proportions.[181]

The crystal structure of HE-PBA is schematically illustrated in Fig. 16(a), where Fe, Mn, Ni, Cu, and Co share the nitrogen-coordinated M site. The configurational entropy value of HE-PBA is calculated to be 1.61*R* based on the eqn (1), which is higher than medium-entropy PBA (ME-PBAs) with a value of 1.39*R*. The X-ray diffraction pattern and corresponding Rietveld refinement profile is shown in Fig. 16(c), which reveals a single-phase cubic structure in the $Fm\bar{3}m$ space group. Fourier transforms of EXAFS data in Fig. 16(d) show that the radial distribution of all the TM metals is similar, indicating a homogeneous distribution of the cations. The three characteristic peaks are attributed to the M–N/C, M–C/N and M–Fe shells.

As shown in Fig. 17(a), HE-PBA electrode delivers an initial-cycle discharge capacity of *ca.* 100 mA h g^{-1} at 0.1 A g^{-1} current density, which is higher than those ME-PBAs where one of the TM metal cations is removed. In addition, the cycling performance of the HE-PBA electrode shows that it retains around 87 mA h g^{-1} after 150 cycles. All electrodes show similar characteristic redox plateaus at around 3.45 and 3.22 V for charge and discharge, respectively, which is consistent with the redox peaks from the CV curves in Fig. 17(c). The HE-PBA electrode displays highest specific capacities at various current densities ranging from 0.2 to 1.0 A g^{-1}.

Operando XRD data collected during two cycles of the HE-PBA electrode against Na^+/Na were performed to reveal the correlation between structure and electrochemical performance. At a low current density of 3.5 mA g^{-1}, the (200), (220) and (400) characteristic reflections undergo continuous but minor shifts to higher 2θ angles, which is accompanied by minor broadening of the peaks (Fig. 18(a)). Fig. 18(b) shows that the (200) peak intensifies at the end of first charging process. In the subsequent discharge cycle, the three reflections gradually shift back to their pristine positions. Similar behavior is observed in the second cycle, suggesting some reversibility in the structural changes observed. There is no other phase

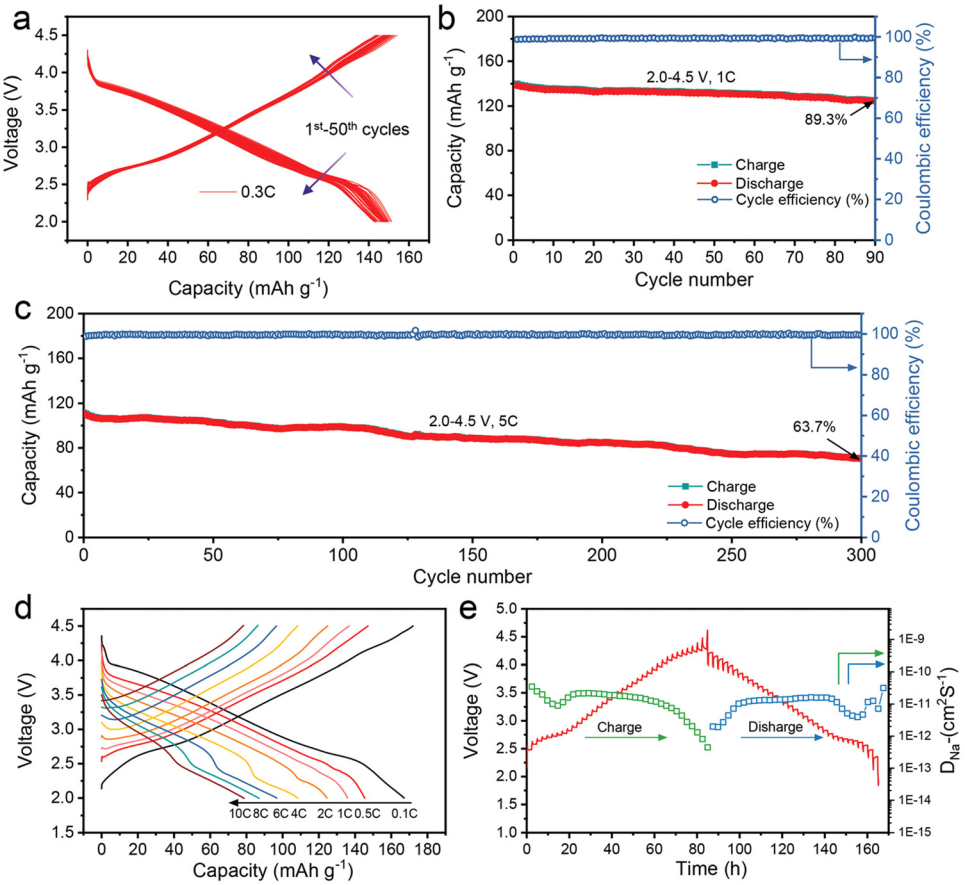

Fig. 13 Electrochemical performance of the NaLFCNM cathode. (a) The charge–discharge profiles at a 0.3 C rate for 50 cycles. (b) and (c) Long-term cycling performance at (b) 1 C rate and (c) at 5 C rate. (d) Charge–discharge curves at different cycling rates. (e) GITT curves for the charge and discharge states of the first cycle and corresponding D_{Na^+}. Copyright permission from Wiley, used from ref. 171.

observed during the whole measurement, which indicates that the sodium intercalation chemistry in HE-PBA occurs *via* a solid-solution mechanism. Compared to the phase-transition mechanism, the solid-solution mechanism can maintain electrode structure during the Na^+ intercalation process, yielding superior cycling performance.[182] The cubic lattice parameters are summarized in Fig. 18(c); the values remain almost unchanged during the two cycles, indicating that the HE-PBA electrode has negligible structural strain upon de/intercalation processes, which agrees with its respectable cycling performance.

HEMs for lithium–sulfur batteries

Lithium–sulfur (Li–S) batteries have high theoretical specific capacities of 1675 mA h g^{-1}.[183] However, sulfur redox reactions often lead to the uncontrollable shuttling of liquid-state lithium polysulfides (LiPSs), resulting in the rapid loss of active material and dramatic capacity fading. A common engineering solution is to utilize metal oxides as additives for the entrapment of LiPSs as they often show excellent adsorption capability owing to the chemical bonding between metal oxides and LiPSs.[184,185] Recently, HE metal oxides have been applied as additives in Li–S batteries due to their vast compositional space and abundant combinatorial active sites, which can potentially provide robust chemical confinement of LiPSs with fast kinetics for polysulfide conversion.[15,186,187] Chen *et al.* utilized a self-sacrificing multi metal–organic framework (*M*-MOFs) to prepare a HEO $(Ni_{0.2}Co_{0.2}Cu_{0.2}Mg_{0.2}Zn_{0.2})O$.[188] When employed as a cathode additive for Li–S batteries, it provides varieties of randomly distributed active metal sites with good electronic conductivity for anchoring polysulfides. A high specific capacity of 1244 mA h g^{-1} was observed, with a capacity retention of 63% after 800 cycles at 0.5 C current density. Gao *et al.* adopted an electrospinning-calcination process to prepare high-entropy perovskite oxide nanofibers of the type $La(Cr_{0.2}Mn_{0.2}Fe_{0.2}Co_{0.2}Ni_{0.2})O_3$ (HE-LMO) and $La_{0.8}Sr_{0.2}(Cr_{0.2}Mn_{0.2}Fe_{0.2}Co_{0.2}Ni_{0.2})O_3$ (HE-LSMO).[189] The as-prepared HEO nanofibers, which have a unique porous structure, could provide rich exposure sites with sulfur. In comparison to $LaMnO_3$ and HE-LMO samples, HE-LSMO plays a bidirectional catalytic role in promoting both the reduction of soluble LiPS and the re-oxidation of insoluble Li_2S. Therefore, it enables fast reaction kinetics for the liquid–solid conversion process. This novel S/HE-LSMO cathode has been reported with a high sulfur loading of 8.4 mg cm^{-2} and a respectable cycling capability at 0.1 C rate.

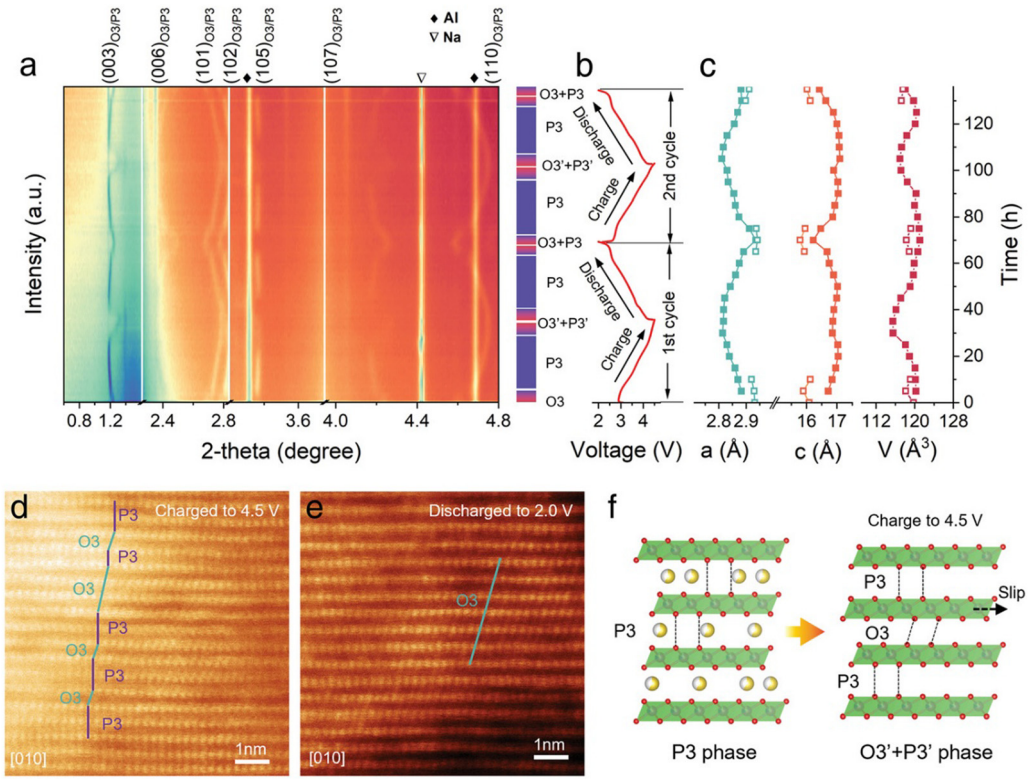

Fig. 14 (a) *In situ* high-energy XRD patterns (λ = 0.15416 nm with a Cu Kα radiation source) of NaLFCNM during the first charge–discharge cycle. (b) Corresponding voltage profile during the *ex situ* XRD experiment. (c) Calculated lattice parameters (*a*-axis, *c*-axis, and unit cell volume [V]) of NaLFCNM at different charge/discharge states. *Ex situ* HR-TEM image of NaLFCNM at (d) 4.5 V state and (e) 2 V state. (f) Schematic illustrations of the whole phase evolution of NaLFCNM upon charge and discharge. Copyright permission from Wiley, used from ref. 171.

HEMs for Zn-ion batteries (ZIBs)

Divalent Zn-ion batteries (ZIBs) are of great interest due to their high volumetric energy density.[190] Nevertheless, poor structural stability, slow ion diffusion kinetics, low working voltage, and undesirable energy/power density of cathode electrodes currently hinder the practical application of ZIBs.[191] The development of advanced cathode materials with advantageous properties is key for the optimization of ZIBs. Jin *et al.* reported a high-entropy Prussian blue analogues (HE-PBA) K(MnCuFeNiCo)[Fe(CN)$_6$] for Zn air batteries using a citrate-assisted controlled crystallization method at room temperature.[192]

The HE-PBA delivers best electrochemical performance compared to its low-entropy counterparts. As shown in Fig. 19(a), the HE-PBA electrode attains a high initial discharge capacity of 90 mA h g^{-1} and excellent capacity retention of 82% after 100 cycles. The rate performance was also studied with the HE-PBA electrode delivering average capacities of 89, 80, 69, 51, 53 and 45 mA h g^{-1} at 0.1, 0.2, 0.3, 0.5, 1 and 2 A g^{-1} current densities, respectively. Nyquist plots of as-prepared electrodes consist of two regions: a semicircle and a straight line, which is related to the charge transfer and electron diffusion process, respectively. It is notable that HE-PBA electrode delivers smallest R_{ct} value compared to other electrodes. In addition, the galvanostatic intermittent titration technique (GITT) was performed to acquire Zn^{2+} diffusion kinetics during charging/discharging processes, as shown in Fig. 19(d), the HE-PBA electrode delivers a higher Zn^{2+} diffusion coefficient of about 10^{-9}–10^{-10} cm^2 S^{-1} compared with the KMnHCF electrode. These results show that high-entropy effect may give rise to better charge transfer kinetics in ZIBs.

Rosen *et al.* prepared a HE MXene of formula $Ti_{1.1}V_{0.7}Cr_xNb_{1.0}Ta_{0.6}C_3T_z$ by selective etching of a HE MAX phase.[193] SEM, XRD, and EDX spectroscopic mapping reveal that a layered structure of the freestanding HE MXene film is obtained and elements are homogeneously distributed throughout the material. The HE MXene films were tested as zinc ion hybrid supercapacitors (ZHSCs), and achieved capacities of 77 and 43 mA h g^{-1} at current densities in the range 0.5–10 A g^{-1} respectively, with a respectable capacity retention of *ca.* 90% after 10 000 cycles. This excellent performance is achieved using both $Zn(CF_3SO_3)_2$ and the more cost-effective $ZnCl_2$ as electrolyte. The HE MXene films also achieved capacities of 126 and 48 mA h g^{-1} when applied as a Li-ion cathode, at current densities of 0.01 and 2.0 A g^{-1}, respectively. This study demonstrates that HE MXene material is another promising candidate for energy storage devices.

HEMs for supercapacitors

Recently, high entropy alloys have been explored as electrode materials for supercapacitors.[100] Liqing *et al.* reported an

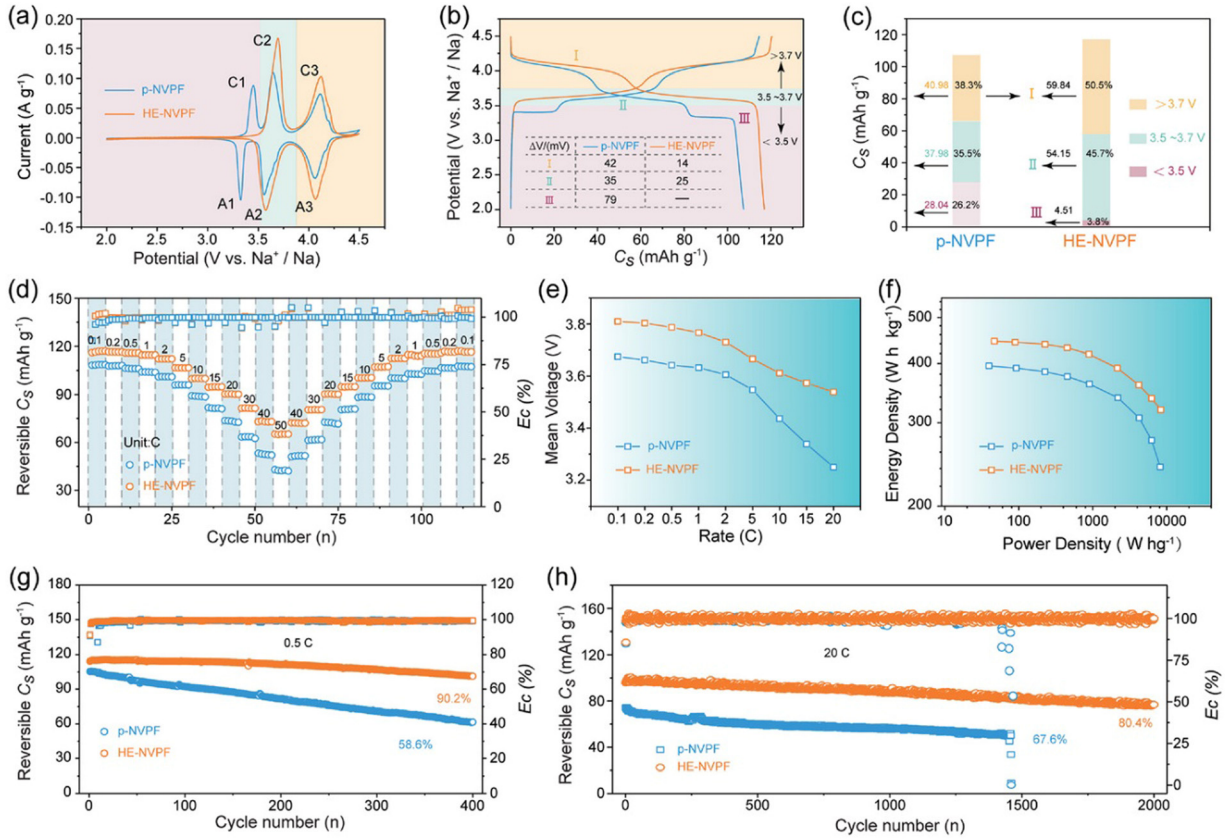

Fig. 15 Electrochemical properties of HE-NVPF and p-NVPF tested in coin-type half-cells. (a) CV patterns at a scan rate of 0.1 mV s^{-1}. (b) GCD curves at 0.1 C in the potential windows of 2.0–4.3 V *versus* Na^+/Na. (c) Discharge C_s values and capacity contribution rates of discharge voltage intervals calculated from the discharge curves in (b). (d) Rate capability. (e) V_m variations calculated based on the discharge curves along with rates. (f) Ragone plots. (g) and (h) Cycling performance at 0.5 and 20 C. Copyright permission from Wiley, used from ref. 42.

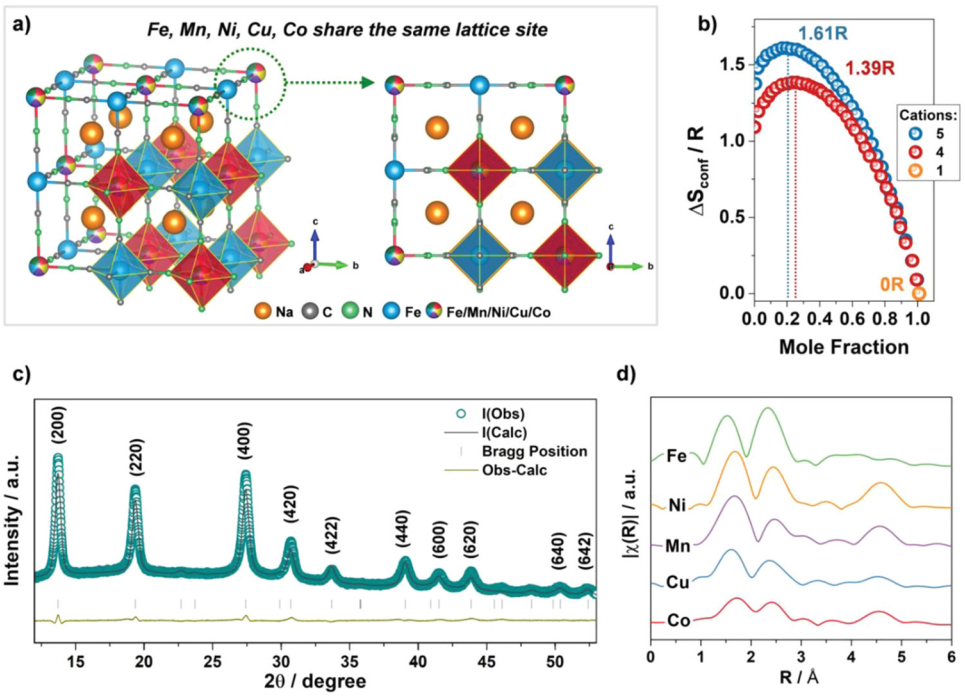

Fig. 16 (a) Schematic illustration of the crystal structure of HE-PBA. (b) Dependence of configurational entropy on the number of elements. (c) XRD pattern of HE-PBA and Rietveld refinement. (d) Fourier transform of the EXAFS signals collected from the HE-PBA material, for each metal edge, plotted with a *y*-axis offset for better clarity. Copyright permission from Wiley, used from ref. 181.

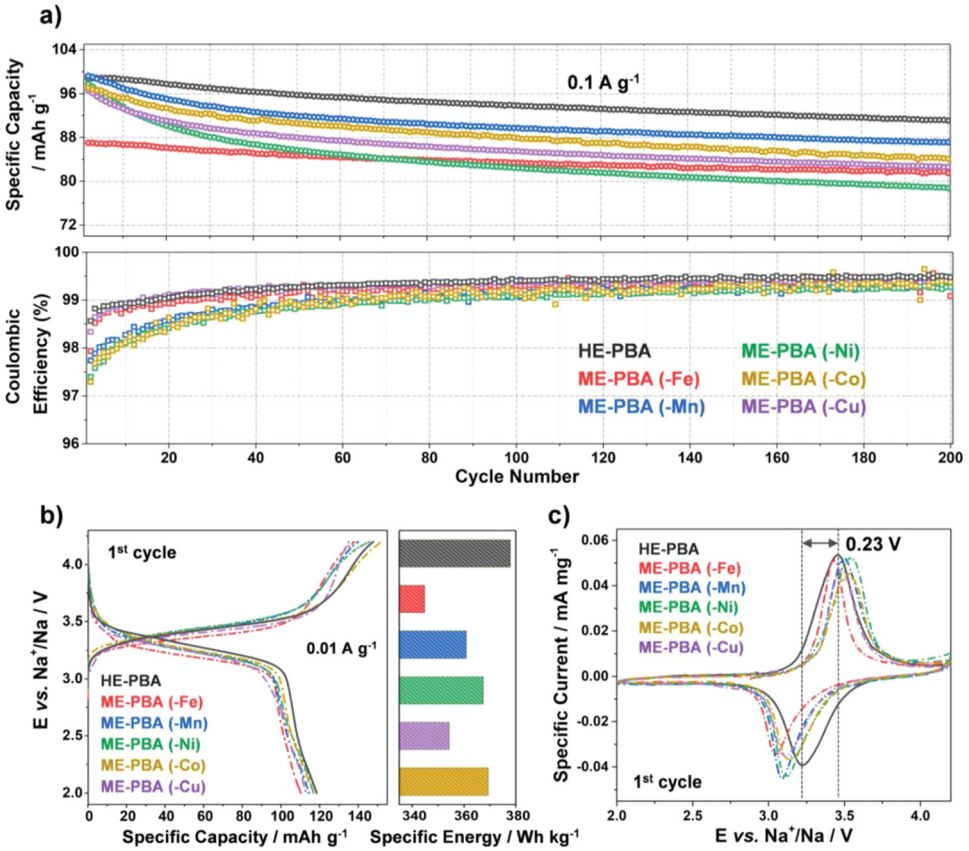

Fig. 17 (a) Galvanostatic cycling at 0.1 A g^{-1} of electrodes based on HE-PBA (in black), ME-PBA(-Fe) (in red), ME-PBA(-Mn) (in blue), ME-PBA(-Ni) (in green), ME-PBA(-Co) (in dark yellow), or ME-PBA(-Cu) (in purple). Specific discharge capacities (top) and Coulombic efficiencies *versus* the cycle number (bottom) are shown (1st cycle is omitted for clarity). (b) First-cycle voltage profiles at 0.01 A g^{-1} (left) and comparison of specific energies (right). (c) Initial cyclic voltammogram at 0.05 mV s^{-1} for HE-PBA and the different ME-PBAs. Copyright permission from Wiley, used from ref. 181.

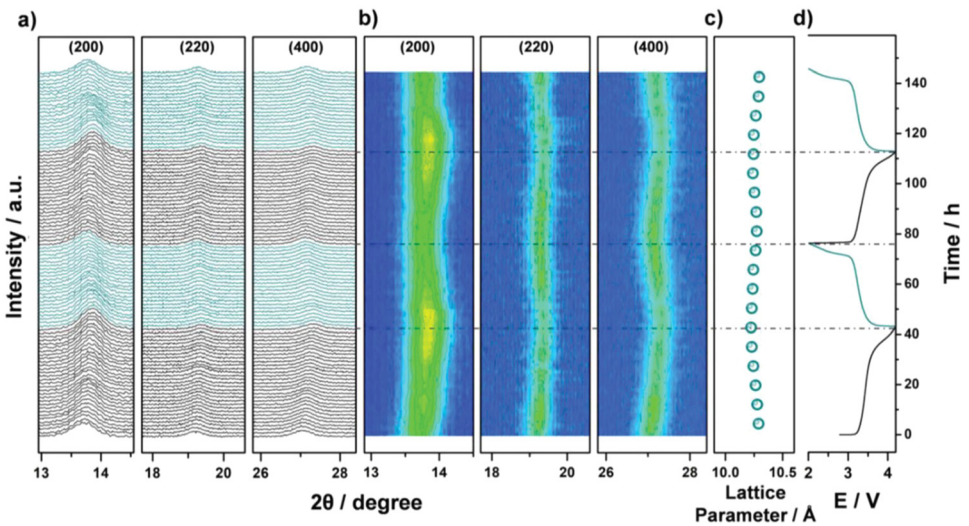

Fig. 18 *Operando* XRD analysis of the electrochemical extraction/insertion of Na ions from/into HE-PBA. (a) Waterfall diagrams, (b) contour plots of consecutively recorded patterns, (c) lattice parameter changes, and (d) the corresponding dis/charge curves for the first two cycles. Copyright permission from Wiley, used from ref. 181.

approach for uniformly depositing HEA-NPs on aligned carbon nanofibers (CNFs) and employed them as electrode materials for supercapacitors.[194] The introduction of CNFs enables high electron transport efficiency, which facilitates the formation of

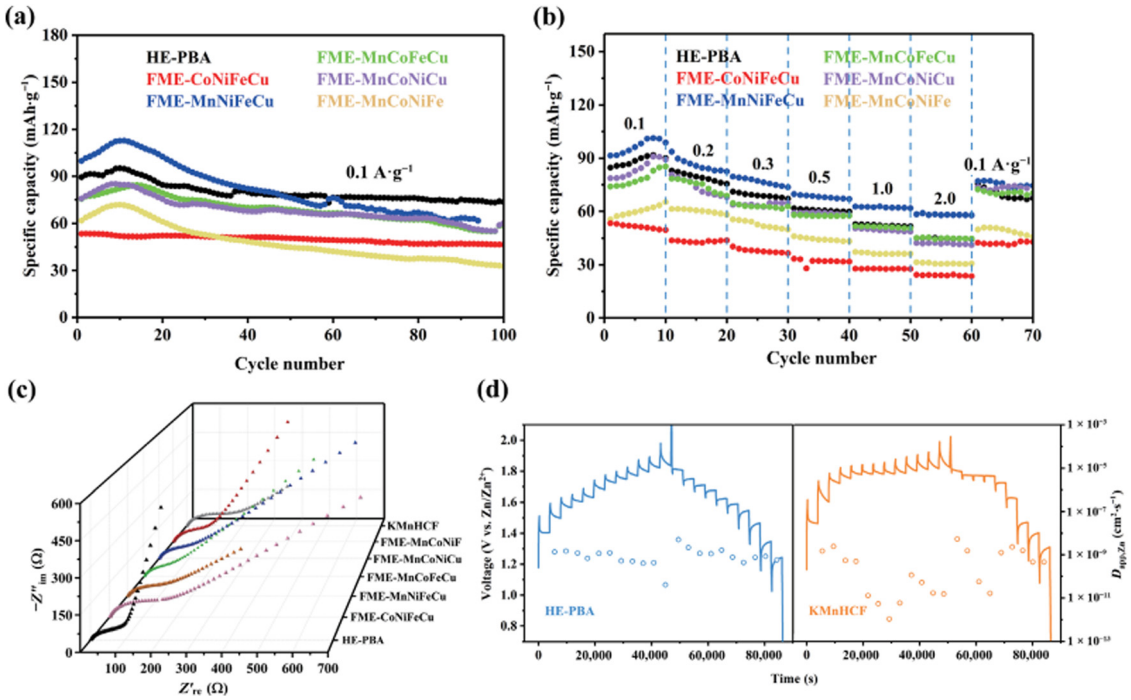

Fig. 19 Electrochemical performance of HE-PBA as ZIB: (a) cycling capability at 0.1 A g^{-1}, (b) rate performance, (c) Nyquist plots of different electrodes, (d) GITT curves and corresponding Zn^{2+} diffusion coefficients of HE-PBA and KMnHCF during charge/discharge process, respectively. Copyright permission from Nano research, used from ref. 192.

HEA-NPs. The FeNiCoMnMg/CNFs hybrid electrode delivers a high capacitance of 203 F g^{-1} and a specific energy density of 21.7 W h Kg^{-1}. In another report, Shen *et al.* incorporated a thin layer of metal oxide (MO) on the outer surface of the HEA FeNiCoCuSn and then deposited this material on hypercrosslinked polymer-based carbon (HCPC).[195] The HEA-NP/HCPCs/MO nanocomposite displayed a high specific capacitance of 495.4 F g^{-1} in 1 M KOH and an excellent capacity retention of 94.7% after 15 000 cycles.

Jiang *et al.* prepared a series of high-entropy Prussian blue analogues (HE-PBAs) (K(MgMnFeNiCu)Fe(CN)$_6$, K(MgMnCoNiCu)Fe(CN)$_6$, K(MnFeCoNiCu)Fe(CN)$_6$, K(MgMnFeCoNi)Fe(CN)$_6$, and K(MgMnFeCoCu)Fe(CN)$_6$.[196] The K(MgMnFeNiCu)Fe(CN)$_6$ electrode displayed the best electrochemical activity, delivering the highest specific capacitance compared to the other HEAs at the same rate. The intrinsic 3D diffusion channels are favorable in electrochemical energy storage devices. Sundra *et al.* synthesized an HEO NP nanocomposite on the surface of carbon nanotubes (HEO-CNT) *via* sol–gel auto-combustion.[102] This composite structure exhibited a high surface area of about 151 m^2 g^{-1}, which is beneficial for electron transport. A HEO-CNT-based full-cell was also assembled, which attains a specific capacity of 286 F g^{-1} at a scan rate of 10 mV s^{-1}. The full-cell also displayed an energy density of 217 W h Kg^{-1} and power density of 24 521 W kg^{-1}.

HEMs for fuel cells

HEMs have also exhibited potential as electrocatalytic electrode materials for fuel cells. Wei *et al.* prepared a HEA-based electrocatalyst PtPdAuNiCo/C for the ethylene glycol oxidation reaction (EGOR) in direct ethylene glycol fuel cells (DEGFC).[197] Testing with cyclic voltammetry (CV), Chronoamperometry (CA), and CO stripping showed that the HEA delivered a lowered EGOR peak potential of 0.55 V, 20 mV lower than a Pt/C electrode with an EGOR activity of 0.482 A mg^{-1}. As the DEGFC anode catalyst, it achieved a power density of 8.38 mW cm^{-2} which was 1.4× that of Pt/C. Xiaofeng *et al.* adopted a sol–gel method to prepare a library of $LaMnO_3$-based high-entropy perovskite oxides (HEPOs) as enhanced cathode materials for solid oxide fuel cells (SOFCs).[111] The optimum electrode $(La_{0.2}Nd_{0.2}Sm_{0.2}Ca_{0.2}Sr_{0.2})MnO_3$ displays increased tolerance to A site cation size differences where higher thermochemical stability and good conductivity can be maintained at high temperatures (215.8 S cm^{-1} at about 970 °C).

Zhao *et al.* applied A_2BO_4 high-entropy oxides to solar oxide fuel cells (SOFCs).[198] The XRD pattern presented showed that the $(La_{0.2}Pr_{0.2}Nd_{0.2}Sm_{0.2}Gd_{0.2})_2CuO_4$ (HE-LCO) had a single tetragonal phase structure with space group *I*4/*mmm*. The HE-LCO delivered respectable conductivity of 64.5 S cm^{-1} at 800 °C with a thermal expansion coefficient of 12.19 × 10^{-6} K^{-1}. The HE-LCO electrode also showed respectable catalytic activity in the oxygen reduction reaction (ORR) with a stable area-specific resistance of 0.52 Ω cm^2 after 70 h in operation. The peak power density of the anode supported fuel cell attained 528 mW cm^{-2} at 700 °C; HE-LCO is thus a fairly promising candidate as a material for SOFC cathodes.

Zhou *et al.* adopted an impregnation method to prepare a nanostructured spinel high-entropy oxide $(Fe_{0.2}Mn_{0.2}Co_{0.2}Ni_{0.2}Zn_{0.2})_3O_4$

(FMCNZ). FMCNZ nanoparticles were uniformly distributed and infiltrated into a network on a $Gd_{0.1}Ce_{0.9}O_{1.95}$ (GDC) skeleton.[199] This novel heterostructured material delivered excellent performance when applied as a cathode material in SOFCs, reaching a maximum power density of 1080 mW cm^{-2} at 800 °C and respectable stability over 100 h operation at 750 °C. This work provides a potential method for designing high-performance SOFCs cathodes by combining the high-entropy concept with nanostructured materials.

Summary and outlook

HEMs have been increasingly explored as active materials in electrochemical energy storage devices. They generally exhibit superior electrochemical performance compared to binary and ternary systems which has been ascribed in part to effects arising from their high configurational entropy. Among them, HE oxides are the most promising electrode materials for energy storage devices. Many HE oxides with rock-salt, spinel, and perovskite structures undergo lithiation, a key feature of modern energy storage materials, and have been found, in general, to possess the dual advantages of high theoretical capacity and excellent reversibility during intercalation processes, which makes them candidates for electrochemical Li storage. In contrast, layered P2, O3 and hybrid P2/O3 HE oxides are primarily applied as cathode materials for electrochemical sodium ion storage. The P2 structure is featured in high ionic conductivity materials while O3 phase generally displays higher initial capacities. HE sodium super ionic conductor (NASICON) and Prussian blue analogue (PBA) materials have also been explored as cathodes for sodium ion batteries. NASICONs possess 3D ion transport channels together with numerous Na^+ vacancies, and thus exhibit rapid Na^+ diffusion kinetics. Due to their rigid open framework, PBAs have the advantages of high operating voltage, excellent rate capability and high specific capacity. In lithium–sulfur batteries, HE oxide additives exhibit superior performance by avoiding the formation of liquid-state lithium polysulfides and thus capacity fading is stymied. For surface charge storage in supercapacitors, HE alloys, and HE oxides deliver high storage capacity and exceptional cyclic stability.

Although significant advancements have been made in HEM-based electrochemical energy storage devices, there are still challenges remaining. These revolve around issues such as understanding phase transformation during synthetic processes, better in-depth understanding of the charge storage mechanisms, and optimal design of the composition and morphology of materials (and also interlayer distances for layered HEMs). A significant challenge remains on the synthesis of HEMs. Atom-up approaches are required in order to obtain maximum mixing of multiple elements into a single phased lattice.[200] It should also be noted again that most synthetic techniques require high temperatures and long timescales in current synthetic procedures, something that may currently stymie production in future especially in the move to sustainable industrial processes. *Ex situ* materials characterization is one method of obtaining further understanding of HEM synthesis on varying timescales. For example, the evolution of colloidal NiPdPtRhIr HE nanoparticles was investigated on the minute timescale and it was found that initial Pd-rich PdNi seeds were formed, followed by the autocatalytic integration of the further metals, eventually yielding roughly equimolar nanoparticles.[201–203] The use of *in situ*/*operando* characterization methods would also greatly enhance our understanding of HEM formation.

The structure–composition–property correlation in HEMs is another possible area that could yield significant benefit in both energy storage capacity, but also in device stability. The individual elements play a critical role in regulating the structure and performance, but it is not yet understood to a high degree how the interactions of different elements in the local environment effect performance. Theoretical calculations on substitutions of individual Mo atoms in 2H MoS_2 slabs with Cr, Mn, W and Re atoms has recently shown a significant benefit towards lowering the absorption energy of protons for hydrogen evolution electrocatalysis, explaining the significantly improved HER performance of HE 2H-(CrMnMoWRe)S_2 over the parent 2H-MoS_2.[20] This calculation substituted up to 3 individual Mo atoms in the slabs, but ideal HEMs are completely randomly distributed and equimolar. The distance of influence was also not understood from this provisional analysis, and is another possibility for future research. Most reported HEMs for energy storage systems emphasize the number of incorporated elements but fail to systematically investigate how the composition of each element influences the structure details and related properties.

The existence of various active sites on the HEMs surfaces is also a significant challenge towards gaining a comprehensive understanding of mechanistic information in energy storage. *In situ* XANES and TEM techniques are powerful tools to investigate the electronic states and structural transformation during charging/discharging process, which is vital for complicated storage mechanism of HEMs; given the structural complexity of many HE materials, these are likely to generate rather complex data that may not even be representative of the whole sample, requiring other complementary and multiscale techniques to fully comprehend. A final advancement in HEMs for energy storage is expanding the scope of investigated HEMs. HE oxides are the most common material utilized to date, but there remains limited investigations into HE chalcogenides, MXenes, oxyfluorides, borides, carbides, nitrides, and phosphides, to name but a few.

Conflicts of interest

There are no conflicts of interest to declare.

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References

1 D. Larcher and J. M. Tarascon, *Nat. Chem.*, 2015, **7**, 19–29.
2 J. Chow, R. J. Kopp and P. R. Portney, *Science*, 2003, **302**, 1528–1531.
3 S. Chu, Y. Cui and N. Liu, *Nat. Mater.*, 2017, **16**, 16–22.
4 C. Chen, Q. Liang, G. Wang, D. Liu and X. Xiong, *Adv. Funct. Mater.*, 2022, **32**, 2107249.
5 D.-c Zuo, S.-c Song, C.-s An, L.-b Tang, Z.-j He and J.-c Zheng, *Nano Energy*, 2019, **62**, 401–409.
6 S.-q Yang, P.-b Wang, H.-x Wei, L.-b Tang, X.-h Zhang, Z.-j He, Y.-j Li, H. Tong and J.-c Zheng, *Nano Energy*, 2019, **63**, 103889.
7 Z. Jiang, Y. Li, C. Han, Z. Huang, X. Wu, Z. He, W. Meng, L. Dai and L. Wang, *J. Electrochem. Soc.*, 2020, **167**, 020550.
8 J. Qu, T. Sheng, Z.-G. Wu, T.-R. Chen, H. Chen, Z.-G. Yang, X.-D. Guo, J.-T. Li, B.-H. Zhong and X.-S. Dou, *J. Mater. Chem. A*, 2018, **6**, 13934–13942.
9 J. Qu, X.-X. Dai, J.-S. Cui, R.-X. Chen, X. Wang, Y.-H. Lin, R. Verduzco and H.-L. Wang, *J. Mater. Chem. A*, 2021, **9**, 16554–16564.
10 J. W. Yeh, S. K. Chen, S. J. Lin, J. Y. Gan, T. S. Chin, T. T. Shun, C. H. Tsau and S. Y. Chang, *Adv. Eng. Mater.*, 2004, **6**, 299–303.
11 B. Cantor, I. T. H. Chang, P. Knight and A. J. B. Vincent, *Mater. Sci. Eng., A*, 2004, **375–377**, 213–218.
12 J.-W. Yeh, *JOM*, 2013, **65**, 1759–1771.
13 C. M. Rost, E. Sachet, T. Borman, A. Moballegh, E. C. Dickey, D. Hou, J. L. Jones, S. Curtarolo and J.-P. Maria, *Nat. Commun.*, 2015, **6**, 8485.
14 A. Sarkar, L. Velasco, D. Wang, Q. Wang, G. Talasila, L. de Biasi, C. Kübel, T. Brezesinski, S. S. Bhattacharya, H. Hahn and B. Breitung, *Nat. Commun.*, 2018, **9**, 3400.
15 A. Sarkar, Q. Wang, A. Schiele, M. R. Chellali, S. S. Bhattacharya, D. Wang, T. Brezesinski, H. Hahn, L. Velasco and B. Breitung, *Adv. Mater.*, 2019, **31**, 1806236.
16 J. Dąbrowa, M. Stygar, A. Mikuła, A. Knapik, K. Mroczka, W. Tejchman, M. Danielewski and M. Martin, *Mater. Lett.*, 2018, **216**, 32–36.
17 M. Stygar, J. Dąbrowa, M. Moździerz, M. Zajusz, W. Skubida, K. Mroczka, K. Berent, K. Świerczek and M. Danielewski, *J. Eur. Ceram. Soc.*, 2020, **40**, 1644–1650.
18 R.-Z. Zhang, F. Gucci, H. Zhu, K. Chen and M. J. Reece, *Inorg. Chem.*, 2018, **57**, 13027–13033.
19 T. X. Nguyen, Y.-H. Su, C.-C. Lin and J.-M. Ting, *Adv. Funct. Mater.*, 2021, **31**, 2106229.
20 J. Qu, A. Elgendy, R. Cai, M. A. Buckingham, A. A. Papaderakis, H. de Latour, K. Hazeldine, G. F. S. Whitehead, F. Alam, C. T. Smith, D. J. Binks, A. Walton, J. M. Skelton, R. A. W. Dryfe, S. J. Haigh and D. J. Lewis, *Adv. Sci.*, 2023, **10**, 2204488.
21 M. Qin, J. Gild, H. Wang, T. Harrington, K. S. Vecchio and J. Luo, *J. Eur. Ceram. Soc.*, 2020, **40**, 4348–4353.
22 M. Qin, Q. Yan, H. Wang, C. Hu, K. S. Vecchio and J. Luo, *Scr. Mater.*, 2020, **189**, 101–105.
23 A. Sekowska, S. Wendel, E. C. Fischer, M. H. H. Nørholm and A. Danchin, *Sci. Rep.*, 2016, **6**, 2.
24 Y. Zhang, Z.-B. Jiang, S.-K. Sun, W.-M. Guo, Q.-S. Chen, J.-X. Qiu, K. Plucknett and H.-T. Lin, *J. Eur. Ceram. Soc.*, 2019, **39**, 3920–3924.
25 B. Ye, T. Wen, M. C. Nguyen, L. Hao, C.-Z. Wang and Y. Chu, *Acta Mater.*, 2019, **170**, 15–23.
26 D. E. McCoy, T. Feo, T. A. Harvey and R. O. Prum, *Nat. Commun.*, 2018, **9**, 1.
27 T. J. Harrington, J. Gild, P. Sarker, C. Toher, C. M. Rost, O. F. Dippo, C. McElfresh, K. Kaufmann, E. Marin, L. Borowski, P. E. Hopkins, J. Luo, S. Curtarolo, D. W. Brenner and K. S. Vecchio, *Acta Mater.*, 2019, **166**, 271–280.
28 V. Braic, M. Balaceanu, M. Braic, A. Vladescu, S. Panseri and A. Russo, *J. Mech. Behav. Biomed. Mater.*, 2012, **10**, 197–205.
29 C.-W. Tsai, S.-W. Lai, K.-H. Cheng, M.-H. Tsai, A. Davison, C.-H. Tsau and J.-W. Yeh, *Thin Solid Films*, 2012, **520**, 2613–2618.
30 K. Johansson, L. Riekehr, S. Fritze and E. Lewin, *Surf. Coat. Technol.*, 2018, **349**, 529–539.
31 K. Balasubramanian, S. V. Khare and D. Gall, *Acta Mater.*, 2018, **159**, 77–88.
32 J. Gild, J. Braun, K. Kaufmann, E. Marin, T. Harrington, P. Hopkins, K. Vecchio and J. Luo, *J. Materiomics*, 2019, **5**, 337–343.
33 Y. Liu, Y. Zhang, H. Zhang, N. Wang, X. Chen, H. Zhang and Y. Li, *J. Alloys Compd.*, 2017, **694**, 869–876.
34 X. Zhao, Z. Xue, W. Chen, Y. Wang and T. Mu, *ChemSusChem*, 2020, **13**, 2038–2042.
35 D. Lai, Q. Kang, F. Gao and Q. Lu, *J. Mater. Chem. A*, 2021, **9**, 17913–17922.
36 H. Qiao, X. Wang, Q. Dong, H. Zheng, G. Chen, M. Hong, C.-P. Yang, M. Wu, K. He and L. Hu, *Nano Energy*, 2021, **86**, 106029.
37 T. Wang, H. Chen, Z. Yang, J. Liang and S. Dai, *J. Am. Chem. Soc.*, 2020, **142**, 4550–4554.
38 X. Chen and Y. Wu, *J. Am. Ceram. Soc.*, 2020, **103**, 750–756.
39 P. A. Sukkurji, Y. Cui, S. Lee, K. Wang, R. Azmi, A. Sarkar, S. Indris, S. S. Bhattacharya, R. Kruk, H. Hahn, Q. Wang, M. Botros and B. Breitung, *J. Mater. Chem. A*, 2021, **9**, 8998–9009.
40 M. Li, C. Sun, Q. Ni, Z. Sun, Y. Liu, Y. Li, L. Li, H. Jin and Y. Zhao, *Adv. Energy Mater.*, 2023, **13**, 2203971.
41 B. Wu, G. Hou, E. Kovalska, V. Mazanek, P. Marvan, L. Liao, L. Dekanovsky, D. Sedmidubsky, I. Marek, C. Hervoches and Z. Sofer, *Inorg. Chem.*, 2022, **61**, 4092–4101.
42 Z.-Y. Gu, J.-Z. Guo, J.-M. Cao, X.-T. Wang, X.-X. Zhao, X.-Y. Zheng, W.-H. Li, Z.-H. Sun, H.-J. Liang and X.-L. Wu, *Adv. Mater.*, 2022, **34**, 2110108.
43 J. Peng, B. Zhang, W. Hua, Y. Liang, W. Zhang, Y. Du, G. Peleckis, S. Indris, Q. Gu, Z. Cheng, J. Wang, H. Liu,

S. Dou and S. Chou, *Angew. Chem., Int. Ed.*, 2023, **62**, e202215865.
44 Y. Huang, X. Zhang, L. Ji, L. Wang, B. B. Xu, M. W. Shahzad, Y. Tang, Y. Zhu, M. Yan, G. Sun and Y. Jiang, *Energy Storage Mater.*, 2023, **58**, 1–8.
45 Y. Ma, Y. Ma, S. L. Dreyer, Q. Wang, K. Wang, D. Goonetilleke, A. Omar, D. Mikhailova, H. Hahn, B. Breitung and T. Brezesinski, *Adv. Mater.*, 2021, **33**, 2170269.
46 M. A. Buckingham, B. Ward-O'Brien, W. Xiao, Y. Li, J. Qu and D. J. Lewis, *Chem. Commun.*, 2022, **58**, 8025–8037.
47 B. Cantor, *Entropy*, 2014, **16**, 4749–4768.
48 C. Triolo, W. Xu, B. Petrovičovà, N. Pinna and S. Santangelo, *Adv. Funct. Mater.*, 2022, **32**, 2202892.
49 X.-F. Luo, J. Patra, W.-T. Chuang, T. X. Nguyen, J.-M. Ting, J. Li, C.-W. Pao and J.-K. Chang, *Adv. Sci.*, 2022, **9**, 2201219.
50 Y. Wang, J. Liu, Y. Song, J. Yu, Y. Tian, M. J. Robson, J. Wang, Z. Zhang, X. Lin, G. Zhou, Z. Wang, L. Shen, H. Zhao, S. Grasso and F. Ciucci, *Small Methods*, 2023, **7**, 2201138.
51 Q. Yang, G. Wang, H. Wu, B. A. Beshiwork, D. Tian, S. Zhu, Y. Yang, X. Lu, Y. Ding, Y. Ling, Y. Chen and B. Lin, *J. Alloys Compd.*, 2021, **872**, 159633.
52 Z. Li, B. Guan, F. Xia, J. Nie, W. Li, L. Ma, W. Li, L. Zhou, Y. Wang, H. Tian, J. Luo, Y. Chen, M. Frost, K. An and X. Liu, *ACS Appl. Mater. Interfaces*, 2022, **14**, 24363–24373.
53 M. B. Zakaria and T. Chikyow, *Coord. Chem. Rev.*, 2017, **352**, 328–345.
54 B. Jiang, Y. Yu, J. Cui, X. Liu, L. Xie, J. Liao, Q. Zhang, Y. Huang, S. Ning, B. Jia, B. Zhu, S. Bai, L. Chen, S. J. Pennycook and J. He, *Science*, 2021, **371**, 830–834.
55 Z. Deng, A. Olvera, J. Casamento, J. S. Lopez, L. Williams, R. Lu, G. Shi, P. F. P. Poudeu and E. Kioupakis, *Chem. Mater.*, 2020, **32**, 6070–6077.
56 A. Mikuła, J. Dąbrowa, A. Kusior, K. Mars, R. Lach and M. Kubowicz, *Dalton Trans.*, 2021, **50**, 9560–9573.
57 A. Yamashita, R. Jha, Y. Goto, T. D. Matsuda, Y. Aoki and Y. Mizuguchi, *Dalton Trans.*, 2020, **49**, 9118–9122.
58 Z. Ma, T. Xu, W. Li, Y. Cheng, J. Li, D. Zhang, Q. Jiang, Y. Luo and J. Yang, *Adv. Funct. Mater.*, 2021, **31**, 2103197.
59 R. Liu, H. Chen, K. Zhao, Y. Qin, B. Jiang, T. Zhang, G. Sha, X. Shi, C. Uher, W. Zhang and L. Chen, *Adv. Mater.*, 2017, **29**, 1702712.
60 L. Lin, K. Wang, A. Sarkar, C. Njel, G. Karkera, Q. Wang, R. Azmi, M. Fichtner, H. Hahn, S. Schweidler and B. Breitung, *Adv. Energy Mater.*, 2022, **12**, 2103090.
61 W. Li, X. Guo, P. Geng, M. Du, Q. Jing, X. Chen, G. Zhang, H. Li, Q. Xu, P. Braunstein and H. Pang, *Adv. Mater.*, 2021, **33**, 2105163.
62 S. Zheng, Q. Li, H. Xue, H. Pang and Q. Xu, *Natl. Sci. Rev.*, 2020, **7**, 305–314.
63 B. Li, P. Gu, Y. Feng, G. Zhang, K. Huang, H. Xue and H. Pang, *Adv. Funct. Mater.*, 2017, **27**, 1605784.
64 M. Bondesgaard, N. L. N. Broge, A. Mamakhel, M. Bremholm and B. B. Iversen, *Adv. Funct. Mater.*, 2019, **29**, 1905933.
65 C. R. McCormick and R. E. Schaak, *J. Am. Chem. Soc.*, 2021, **143**, 1017–1023.
66 B. Ward-O'Brien, E. J. Pickering, R. Ahumada-Lazo, C. Smith, X. L. Zhong, Y. Aboura, F. Alam, D. J. Binks, T. L. Burnett and D. J. Lewis, *J. Am. Chem. Soc.*, 2021, **143**, 21560–21566.
67 B. Ward-O'Brien, P. D. McNaughter, R. Cai, A. Chattopadhyay, J. M. Flitcroft, C. T. Smith, D. J. Binks, J. M. Skelton, S. J. Haigh and D. J. Lewis, *Nano Lett.*, 2022, **22**, 8045–8051.
68 J. Qu, A. Elgendy, R. Cai, M. A. Buckingham, A. A. Papaderakis, H. de Latour, K. Hazeldine, G. F. S. Whitehead, F. Alam, C. T. Smith, D. J. Binks, A. Walton, J. M. Skelton, R. A. W. Dryfe, S. J. Haigh and D. J. Lewis, *Adv. Sci.*, 2023, **10**, 2204488.
69 M. A. Malik, M. Afzaal and P. O'Brien, *Chem. Rev.*, 2010, **110**, 4417–4446.
70 J. C. Sarker and G. Hogarth, *Chem. Rev.*, 2021, **121**, 6057–6123.
71 N. Zeng, Y.-C. Wang, J. Neilson, S. M. Fairclough, Y. Zou, A. G. Thomas, R. J. Cernik, S. J. Haigh and D. J. Lewis, *Chem. Mater.*, 2020, **32**, 7895–7907.
72 F. Ding, C. Zhao, D. Xiao, X. Rong, H. Wang, Y. Li, Y. Yang, Y. Lu and Y.-S. Hu, *J. Am. Chem. Soc.*, 2022, **144**, 8286–8295.
73 H. Chen, K. Jie, C. J. Jafta, Z. Yang, S. Yao, M. Liu, Z. Zhang, J. Liu, M. Chi, J. Fu and S. Dai, *Appl. Catal., B*, 2020, **276**, 119155.
74 H. Xu, Z. Zhang, J. Liu, C.-L. Do-Thanh, H. Chen, S. Xu, Q. Lin, Y. Jiao, J. Wang, Y. Wang, Y. Chen and S. Dai, *Nat. Commun.*, 2020, **11**, 3908.
75 S.-Y. Su, Y.-T. Fan, Y.-J. Su, C.-W. Huang, M.-H. Tsai and M.-Y. Lu, *J. Alloys Compd.*, 2021, **851**, 156909.
76 P. Shi, R. Li, Y. Li, Y. Wen, Y. Zhong, W. Ren, Z. Shen, T. Zheng, J. Peng, X. Liang, P. Hu, N. Min, Y. Zhang, Y. Ren, P. K. Liaw, D. Raabe and Y.-D. Wang, *Science*, 2021, **373**, 912–918.
77 L. Wang, B. Cheng, L. Zhang and J. Yu, *Small*, 2021, **17**, 2103447.
78 P. Xie, Y. Yao, Z. Huang, Z. Liu, J. Zhang, T. Li, G. Wang, R. Shahbazian-Yassar, L. Hu and C. Wang, *Nat. Commun.*, 2019, **10**, 4011.
79 S. Rehman, M. A. Bader and S. A. Al-Moallem, *Renewable Sustainable Energy Rev.*, 2007, **11**, 1843–1857.
80 G. R. Timilsina, L. Kurdgelashvili and P. A. Narbel, *Renewable Sustainable Energy Rev.*, 2012, **16**, 449–465.
81 F. O. Rourke, F. Boyle and A. Reynolds, *Appl. Energy*, 2010, **87**, 398–409.
82 L. Wang, Y. Han, X. Feng, J. Zhou, P. Qi and B. Wang, *Coord. Chem. Rev.*, 2016, **307**, 361–381.
83 M. J. Guan and W. H. Liao, *J. Intell. Mater. Syst. Struct.*, 2007, **19**, 671–680.
84 R. J. M. Vullers, R. van Schaijk, I. Doms, C. Van Hoof and R. Mertens, *Solid-State Electron.*, 2009, **53**, 684–693.
85 P. Simon, Y. Gogotsi and B. Dunn, *Science*, 2014, **343**, 1210–1211.
86 P. Thounthong, S. Raël and B. Davat, *J. Power Sources*, 2006, **158**, 806–814.
87 D. Shin, K. Lee and N. Chang, *Int. J. Hydrogen Energy*, 2016, **41**, 1381–1390.

88 M. Winter and R. J. Brodd, *Chem. Rev.*, 2004, **104**, 4245–4270.
89 A. Kadri, H. Marzougui, A. Aouiti and F. Bacha, *Energy*, 2020, **192**, 116518.
90 P. D. Richards, R. J. Myers, S. M. Swinton and R. T. Walker, *Global Environ. Change*, 2012, **22**, 454–462.
91 H. Wang, H. S. Casalongue, Y. Liang and H. Dai, *J. Am. Chem. Soc.*, 2010, **132**, 7472–7477.
92 K. Zhang, L. L. Zhang, X. S. Zhao and J. Wu, *Chem. Mater.*, 2010, **22**, 1392–1401.
93 S.-M. Paek, E. Yoo and I. Honma, *Nano Lett.*, 2009, **9**, 72–75.
94 D. Bérardan, S. Franger, D. Dragoe, A. K. Meena and N. Dragoe, *Phys. Status Solidi RRL*, 2016, **10**, 328–333.
95 N. Osenciat, D. Bérardan, D. Dragoe, B. Léridon, S. Holé, A. K. Meena, S. Franger and N. Dragoe, *J. Am. Ceram. Soc.*, 2019, **102**, 6156–6162.
96 D. Bérardan, S. Franger, A. K. Meena and N. Dragoe, *J. Mater. Chem. A*, 2016, **4**, 9536–9541.
97 Y. Lee, J. Suntivich, K. J. May, E. E. Perry and Y. Shao-Horn, *J. Phys. Chem. Lett.*, 2012, **3**, 399–404.
98 Q. Wang, A. Sarkar, Z. Li, Y. Lu, L. Velasco, S. S. Bhattacharya, T. Brezesinski, H. Hahn and B. Breitung, *Electrochem. Commun.*, 2019, **100**, 121–125.
99 N. Qiu, H. Chen, Z. Yang, S. Sun, Y. Wang and Y. Cui, *J. Alloys Compd.*, 2019, **777**, 767–774.
100 K. Kong, J. Hyun, Y. Kim, W. Kim and D. Kim, *J. Power Sources*, 2019, **437**, 226927.
101 T. Jin, X. Sang, R. R. Unocic, R. T. Kinch, X. Liu, J. Hu, H. Liu and S. Dai, *Adv. Mater.*, 2018, **30**, 1707512.
102 M. S. Lal and R. Sundara, *ACS Appl. Mater. Interfaces*, 2019, **11**, 30846–30857.
103 C. Zhao, F. Ding, Y. Lu, L. Chen and Y.-S. Hu, *Angew. Chem., Int. Ed.*, 2020, **59**, 264–269.
104 Y. Zhang, Y. Li, X. Xia, X. Wang, C. Gu and J. Tu, *Sci. China: Technol. Sci.*, 2015, **58**, 1809–1828.
105 Y. Shao-Horn, L. Croguennec, C. Delmas, E. C. Nelson and M. A. O'Keefe, *Nat. Mater.*, 2003, **2**, 464–467.
106 J. Zheng, T. Liu, Z. Hu, Y. Wei, X. Song, Y. Ren, W. Wang, M. Rao, Y. Lin, Z. Chen, J. Lu, C. Wang, K. Amine and F. Pan, *J. Am. Chem. Soc.*, 2016, **138**, 13326–13334.
107 T. Ohzuku and Y. Makimura, *Chem. Lett.*, 2001, 642–643.
108 D. A. Vinnik, E. A. Trofimov, V. E. Zhivulin, S. A. Gudkova, O. V. Zaitseva, D. A. Zherebtsov, A. Y. Starikov, D. P. Sherstyuk, A. A. Amirov, A. V. Kalgin, S. V. Trukhanov and F. V. Podgornov, *Nanomaterials*, 2020, **10**, 268.
109 M. Lim, Z. Rak, J. L. Braun, C. M. Rost, G. N. Kotsonis, P. E. Hopkins, J. P. Maria and D. W. Brenner, *J. Appl. Phys.*, 2019, **125**, 055105.
110 R. Djenadic, A. Sarkar, O. Clemens, C. Loho, M. Botros, V. S. K. Chakravadhanula, C. Kübel, S. S. Bhattacharya, A. S. Gandhi and H. Hahn, *Mater. Res. Lett.*, 2017, **5**, 102–109.
111 Y. Shi, N. Ni, Q. Ding and X. Zhao, *J. Mater. Chem. A*, 2022, **10**, 2256–2270.
112 B. Xiao, G. Wu, T. Wang, Z. Wei, Z. Xie, Y. Sui, J. Qi, F. Wei, X. Zhang, L.-B. Tang and J.-C. Zheng, *ACS Appl. Mater. Interfaces*, 2023, **15**, 2792–2803.
113 L. Su, J. Ren, T. Lu, K. Chen, J. Ouyang, Y. Zhang, X. Zhu, L. Wang, H. Min, W. Luo, Z. Sun, Q. Zhang, Y. Wu, L. Sun, L. Mai and F. Xu, *Adv. Mater.*, 2023, **35**, 2205751.
114 X. L. Wang, E. M. Jin, G. Sahoo and S. M. Jeong, *Batteries*, 2023, **9**, 147.
115 O. J. Marques, C. Chen, E. V. Timofeeva and C. U. Segre, *J. Power Sources*, 2023, **564**, 232852.
116 D. O. Bayraktar, E. Lökçü, C. Ozgur, T. Erdil and C. Toparli, *Int. J. Energy Res.*, 2022, **46**, 22124–22133.
117 S. Li, Z. Peng and X. Fu, *J. Adv. Ceram.*, 2023, **12**, 59–71.
118 Q. Zheng, Z. Ren, Y. Zhang, T. Qin, J. Qi, H. Jia, L. Jiang, L. Li, X. Liu and L. Chen, *ACS Appl. Mater. Interfaces*, 2023, **15**, 4643–4651.
119 T. Jiang, F. Wu, Y. Ren, J. Qiu and Z. Chen, *J. Solid State Electrochem.*, 2023, **27**, 763–772.
120 C.-H. Kuo, A.-Y. Wang, H.-Y. Liu, S.-C. Huang, X.-R. Chen, C.-C. Chi, Y.-C. Chang, M.-Y. Lu and H.-Y. Chen, *APL Mater.*, 2022, **10**, 121104.
121 H. Minouei, N. Tsvetkov, M. Kheradmandfard, J. Han, D.-E. Kim and S. I. Hong, *J. Power Sources*, 2022, **549**, 232041.
122 M. Fracchia, D. Callegari, M. Coduri, U. Anselmi-tamburini, M. Manzoli, E. Quartarone and P. Ghigna, *Front. Energy Res.*, 2022, **10**, 883206.
123 J.-Z. Yen, Y.-C. Yang and H.-Y. Tuan, *Chem. Eng. J.*, 2022, **450**, 137924.
124 X. Yang, H. Wang, Y. Song, K. Liu, T. Huang, X. Wang, C. Zhang and J. Li, *ACS Appl. Mater. Interfaces*, 2022, **14**, 26873–26881.
125 T. Kawaguchi, X. Bian, T. Hatakeyama, H. Li and T. Ichitsubo, *ACS Appl. Energy Mater.*, 2022, **5**, 4369–4381.
126 B. Xiao, G. Wu, T. Wang, Z. Wei, Y. Sui, B. Shen, J. Qi, F. Wei and J. Zheng, *Nano Energy*, 2022, **95**, 106962.
127 X.-F. Luo, J. Patra, W.-T. Chuang, T. X. Nguyen, J.-M. Ting, J. Li, C.-W. Pao and J.-K. Chang, *Adv. Sci.*, 2022, **9**, e2201219.
128 X. Liu, Y. Xing, K. Xu, H. Zhang, M. Gong, Q. Jia, S. Zhang and W. Lei, *Small*, 2022, **18**, 2200524.
129 J. Su, Z. Cao, Z. Jiang, G. Chen, Y. Zhu, L. Wang and G. Li, *Int. J. Appl. Ceram. Technol.*, 2022, **19**, 2004–2015.
130 J. Patra, T. X. Nguyen, C.-C. Tsai, O. Clemens, J. Li, P. Pal, W. K. Chan, C.-H. Lee, H.-Y. T. Chen, J.-M. Ting and J.-K. Chang, *Adv. Funct. Mater.*, 2022, **32**, 2110992.
131 H. Guo, J. Shen, T. Wang, C. Cheng, H. Yao, X. Han and Q. Zheng, *Ceram. Int.*, 2022, **48**, 3344–3350.
132 J. Zhao, X. Yang, Y. Huang, F. Du and Y. Zeng, *ACS Appl. Mater. Interfaces*, 2021, **13**, 58674–58681.
133 T. X. Nguyen, C.-C. Tsai, J. Patra, O. Clemens, J.-K. Chang and J.-M. Ting, *Chem. Eng. J.*, 2022, **430**, 132658.
134 S.-Y. Wang, T.-Y. Chen, C.-H. Kuo, C.-C. Lin, S.-C. Huang, M.-H. Lin, C.-C. Wang and H.-Y. Chen, *Mater. Chem. Phys.*, 2021, **274**, 125105.
135 C.-Y. Huang, C.-W. Huang, M.-C. Wu, J. Patra, T. Xuyen Nguyen, M.-T. Chang, O. Clemens, J.-M. Ting, J. Li, J.-K. Chang and W.-W. Wu, *Chem. Eng. J.*, 2021, **420**, 129838.
136 M. Kheradmandfard, H. Minouei, N. Tsvetkov, A. K. Vayghan, S. F. Kashani-Bozorg, G. Kim, S. I. Hong and D.-E. Kim, *Mater. Chem. Phys.*, 2021, **262**, 124265.

137 H.-Z. Xiang, H.-X. Xie, Y.-X. Chen, H. Zhang, A. Mao and C.-H. Zheng, *J. Mater. Sci.*, 2021, **56**, 8127–8142.
138 P. Ghigna, L. Airoldi, M. Fracchia, D. Callegari, U. Anselmi-Tamburini, P. D'Angelo, N. Pianta, R. Ruffo, G. Cibin, D. O. de Souza and E. Quartarone, *ACS Appl. Mater. Interfaces*, 2020, **12**, 50344–50354.
139 T.-Y. Chen, S.-Y. Wang, C.-H. Kuo, S.-C. Huang, M.-H. Lin, C.-H. Li, H.-Y. T. Chen, C. C. Wang, Y. F. Liao, C.-C. Lin, Y.-M. Chang, J.-W. Yeh, S.-J. Lin, T.-Y. Chen and H.-Y. Chen, *J. Mater. Chem. A*, 2020, **8**, 21756–21770.
140 T. X. Nguyen, J. Patra, J.-K. Chang and J.-M. Ting, *J. Mater. Chem. A*, 2020, **8**, 18963–18973.
141 D. Wang, S. Jiang, C. Duan, J. Mao, Y. Dong, K. Dong, Z. Wang, S. Luo, Y. Liu and X. Qi, *J. Alloys Compd.*, 2020, **844**, 156158.
142 E. Lökçü, Ç. Toparli and M. Anik, *ACS Appl. Mater. Interfaces*, 2020, **12**, 23860–23866.
143 H. Chen, N. Qiu, B. Wu, Z. Yang, S. Sun and Y. Wang, *RSC Adv.*, 2020, **10**, 9736–9744.
144 J. Yan, D. Wang, X. Zhang, J. Li, Q. Du, X. Liu, J. Zhang and X. Qi, *J. Mater. Sci.*, 2020, **55**, 6942–6951.
145 B. Breitung, Q. Wang, A. Schiele, Đ. Tripković, A. Sarkar, L. Velasco, D. Wang, S. S. Bhattacharya, H. Hahn and T. Brezesinski, *Batteries Supercaps*, 2020, **3**, 361–369.
146 Z. Lun, B. Ouyang, D.-H. Kwon, Y. Ha, E. E. Foley, T.-Y. Huang, Z. Cai, H. Kim, M. Balasubramanian, Y. Sun, J. Huang, Y. Tian, H. Kim, B. D. McCloskey, W. Yang, R. J. Clément, H. Ji and G. Ceder, *Nat. Mater.*, 2021, **20**, 214–221.
147 Q. Wang, A. Sarkar, D. Wang, L. Velasco, R. Azmi, S. S. Bhattacharya, T. Bergfeldt, A. Düvel, P. Heitjans, T. Brezesinski, H. Hahn and B. Breitung, *Energy Environ. Sci.*, 2019, **12**, 2433–2442.
148 L. Lin, K. Wang, A. Sarkar, C. Njel, G. Karkera, W. Qingsong, R. Azmi, M. Fichtner, H. Hahn, S. Schweidler and B. Breitung, *Adv. Energy Mater.*, 2022, **12**, 2103090.
149 P. G. Bruce, B. Scrosati and J.-M. Tarascon, *Angew. Chem., Int. Ed.*, 2008, **47**, 2930–2946.
150 W. Yang, G. Cheng, C. Dong, Q. Bai, X. Chen, Z. Peng and Z. Zhang, *J. Mater. Chem. A*, 2014, **2**, 20022–20029.
151 R. Malini, U. Uma, T. Sheela, M. Ganesan and N. G. Renganathan, *Ionics*, 2009, **15**, 301–307.
152 M. Armand and J. M. Tarascon, *Nature*, 2008, **451**, 652–657.
153 G. A. Campbell, *Miner. Econ.*, 2020, **33**, 21–28.
154 M. A. Buckingham, K. Laws, E. Cross, A. J. Surman and L. Aldous, *Green Chem.*, 2021, **23**, 8901–8915.
155 C. Delmas, C. Fouassier and P. Hagenmuller, *Phys. B + C*, 1980, **99**, 81–85.
156 S. Guo, P. Liu, H. Yu, Y. Zhu, M. Chen, M. Ishida and H. Zhou, *Angew. Chem., Int. Ed.*, 2015, **54**, 5894–5899.
157 J. Zhang, W. Wang, W. Wang, S. Wang and B. Li, *ACS Appl. Mater. Interfaces*, 2019, **11**, 22051–22066.
158 S. Yan, S. Luo, L. Yang, J. Feng, P. Li, Q. Wang, Y. Zhang and X. Liu, *J. Adv. Ceram.*, 2022, **11**, 158–171.
159 J. Molenda, A. Milewska, W. Zając, K. Walczak, M. Wolczko, A. Komenda and J. Tobola, *J. Mater. Chem. A*, 2023, **11**, 4248–4260.
160 H. Deng, L. Liu and Z. Shi, *Mater. Lett.*, 2023, **340**, 134113.
161 P. Zhou, Z. Che, J. Liu, J. Zhou, X. Wu, J. Weng, J. Zhao, H. Cao, J. Zhou and F. Cheng, *Energy Storage Mater.*, 2023, **57**, 618–627.
162 J. Liu, Y. Li, Z. Chen, N. Liu, L. Zheng, W. Shi and X. Wang, *J. Am. Chem. Soc.*, 2022, **144**, 23191–23197.
163 K. Tian, H. He, X. Li, D. Wang, Z. Wang, R. Zheng, H. Sun, Y. Liu and Q. Wang, *J. Mater. Chem. A*, 2022, **10**, 14943–14953.
164 C.-C. Lin, H.-Y. Liu, J.-W. Kang, C.-C. Yang, C.-H. Li, H.-Y. T. Chen, S.-C. Huang, C.-S. Ni, Y.-C. Chuang, B.-H. Chen, C.-K. Chang and H.-Y. Chen, *Energy Storage Mater.*, 2022, **51**, 159–171.
165 F. Fu, X. Liu, X. Fu, H. Chen, L. Huang, J. Fan, J. Le, Q. Wang, W. Yang, Y. Ren, K. Amine, S.-G. Sun and G.-L. Xu, *Nat. Commun.*, 2022, **13**, 2826.
166 K. Walczak, A. Plewa, C. Ghica, W. Zając, A. Trenczek-Zając, M. Zając, J. Toboła and J. Molenda, *Energy Storage Mater.*, 2022, **47**, 500–514.
167 L. Yang, C. Chen, S. Xiong, C. Zheng, P. Liu, Y. Ma, W. Xu, Y. Tang, S. P. Ong and H. Chen, *JACS Au*, 2021, **1**, 98–107.
168 B. Wu, G. Hou, E. Kovalska, V. Mazanek, P. Marvan, L. Liao, L. Dekanovsky, D. Sedmidubsky, I. Marek, C. Hervoches and Z. Sofer, *Inorg. Chem.*, 2022, **61**, 4092–4101.
169 Y. Ma, Y. Ma, S. L. Dreyer, Q. Wang, K. Wang, D. Goonetilleke, A. Omar, D. Mikhailova, H. Hahn, B. Breitung and T. Brezesinski, *Adv. Mater.*, 2021, **33**, 2170269.
170 A. M. Huízar-Félix, T. Hernández, S. de la Parra, J. Ibarra and B. Kharisov, *Powder Technol.*, 2012, **229**, 290–293.
171 L. Yao, P. Zou, C. Wang, J. Jiang, L. Ma, S. Tan, K. A. Beyer, F. Xu, E. Hu and H. L. Xin, *Adv. Energy Mater.*, 2022, **12**, 2201989.
172 B. Singh, Z. Wang, S. Park, G. S. Gautam, J.-N. Chotard, L. Croguennec, D. Carlier, A. K. Cheetham, C. Masquelier and P. Canepa, *J. Mater. Chem. A*, 2021, **9**, 281–292.
173 Z. Jian, C. Yuan, W. Han, X. Lu, L. Gu, X. Xi, Y.-S. Hu, H. Li, W. Chen, D. Chen, Y. Ikuhara and L. Chen, *Adv. Funct. Mater.*, 2014, **24**, 4265–4272.
174 C. Wang, D. Du, M. Song, Y. Wang and F. Li, *Adv. Energy Mater.*, 2019, **9**, 1900022.
175 G. Yan, S. Mariyappan, G. Rousse, Q. Jacquet, M. Deschamps, R. David, B. Mirvaux, J. W. Freeland and J.-M. Tarascon, *Nat. Commun.*, 2019, **10**, 585.
176 B. Zhang, R. Dugas, G. Rousse, P. Rozier, A. M. Abakumov and J.-M. Tarascon, *Nat. Commun.*, 2016, **7**, 10308.
177 J. Chen, L. Wei, A. Mahmood, Z. Pei, Z. Zhou, X. Chen and Y. Chen, *Energy Storage Mater.*, 2020, **25**, 585–612.
178 J. Qian, C. Wu, Y. Cao, Z. Ma, Y. Huang, X. Ai and H. Yang, *Adv. Energy Mater.*, 2018, **8**, 1870079.
179 R. Rehman, J. Peng, H. Yi, Y. Shen, J. Yin, C. Li, C. Fang, Q. Li and J. Han, *RSC Adv.*, 2020, **10**, 27033–27041.
180 B. Wang, Y. Han, X. Wang, N. Bahlawane, H. Pan, M. Yan and Y. Jiang, *iScience*, 2018, **3**, 110–133.
181 Y. Ma, Y. Ma, S. L. Dreyer, Q. Wang, K. Wang, D. Goonetilleke, A. Omar, D. Mikhailova, H. Hahn, B. Breitung and T. Brezesinski, *Adv. Mater.*, 2021, **33**, 2101342.
182 S.-M. Xu, X. Liang, X.-Y. Wu, S.-L. Zhao, J. Chen, K.-X. Wang and J.-S. Chen, *Nat. Commun.*, 2019, **10**, 5810.

183 S.-H. Chung and A. Manthiram, *Adv. Mater.*, 2019, **31**, 1901125.
184 Y. Wang, R. Zhang, J. Chen, H. Wu, S. Lu, K. Wang, H. Li, C. J. Harris, K. Xi, R. V. Kumar and S. Ding, *Adv. Energy Mater.*, 2019, **9**, 1900953.
185 Z.-Y. Wang, D.-D. Han, S. Liu, G.-R. Li, T.-Y. Yan and X.-P. Gao, *Electrochim. Acta*, 2020, **337**, 135772.
186 R.-Z. Zhang and M. J. Reece, *J. Mater. Chem. A*, 2019, **7**, 22148–22162.
187 M. J. R. Haché, C. Cheng and Y. Zou, *J. Mater. Res.*, 2020, **35**, 1051–1075.
188 H. Raza, J. Cheng, C. Lin, S. Majumder, G. Zheng and G. Chen, *EcoMat*, 2023, **5**, e12324.
189 L. Tian, Z. Zhang, S. Liu, G. Li and X. Gao, *Nano Energy*, 2023, **106**, 108037.
190 L. Cao, D. Li, T. Pollard, T. Deng, B. Zhang, C. Yang, L. Chen, J. Vatamanu, E. Hu, M. J. Hourwitz, L. Ma, M. Ding, Q. Li, S. Hou, K. Gaskell, J. T. Fourkas, X.-Q. Yang, K. Xu, O. Borodin and C. Wang, *Nat. Nanotechnol.*, 2021, **16**, 902–910.
191 W. Gao, J. Michalička and M. Pumera, *Small*, 2022, **18**, 2105572.
192 J. Xing, Y. Zhang, Y. Jin and Q. Jin, *Nano Res.*, 2023, **16**, 2486–2494.
193 A.-S. Etman, J. Zhou and J. Rosen, *Electrochem. Commun.*, 2022, **137**, 107264.
194 X. Xu, Y. Du, C. Wang, Y. Guo, J. Zou, K. Zhou, Z. Zeng, Y. Liu and L. Li, *J. Alloys Compd.*, 2020, **822**, 153642.
195 E. Shen, X. Song, Q. Chen, M. Zheng, J. Bian and H. Liu, *ChemElectroChem*, 2021, **8**, 260–269.
196 W. Jiang, T. Wang, H. Chen, X. Suo, J. Liang, W. Zhu, H. Li and S. Dai, *Nano Energy*, 2021, **79**, 105464.
197 W. Wang, X. Li, Y. Cheng, M. Zhang, K. Zhao and Y. Liu, *J. Taiwan Inst. Chem. Eng.*, 2023, **143**, 104714.
198 S.-J. Zhao, N. Li, L.-P. Sun, Q. Li, L.-H. Huo and H. Zhao, *J. Alloys Compd.*, 2022, **895**, 162548.
199 Z. Lin, B. Ma, Z.-H. Chen and Y.-K. Zhou, *Ceram. Int.*, 2023, **49**, 23057.
200 M. A. Buckingham, J. M. Skelton and D. J. Lewis, *Cryst. Growth Des.*, 2023, DOI: **10.1021/acs.cgd.3c00712**.
201 C.-M.-D. Silva, H. Amara, F. Fossard, A. Girard, A. Loiseau and V. Huc, *Nanoscale*, 2022, **14**, 9832–9841.
202 G.-R. Dey, C.-R. McCormick, S.-S. Soliman, A.-J. Darling and R.-E. Schaak, *ACS Nano*, 2023, **17**, 5943–5955.
203 N.-L.-N. Broge, M. Bondesgaard, F. Søndergaard-Pedersen, M. Roelsgaard and B.-B. Iversen, *Angew. Chem., Int. Ed.*, 2020, **132**, 22104.