



Cite this: *React. Chem. Eng.*, 2020, 5, 342

## Continuous manufacturing of silver nanoparticles between 5 and 80 nm with rapid online optical size and shape evaluation†

Bruno Pinho and Laura Torrente-Murciano \*

The physical and chemical properties of metal nanoparticles are strongly dependent on their size and shape. In this work, we present a flexible manufacturing approach for the synthesis of spherical silver nanoparticles with tuneable sizes between 5 to 80 nm. This unique size flexibility is enabled by rapid online characterisation coupling spectroscopy and a mathematical Mie theory-based algorithm for size and shape evaluation. While it is conventionally believed that narrow size distributions require a fast nucleation step, herein, we demonstrate that fast and controllable growth is also required. To achieve this, a combination of chemical and engineering approaches is presented to limit thermodynamically driven size focus, coalescence and secondary nucleation. We show that an optimum reducing agent to silver precursor to seeds ratio and pH range need to be maintained throughout the growth stage. Such demanding conditions can be achieved by accurate control of the feed points and fluid dynamics across a series of microfluidic helical reactors leading to low mixing times. In this way, particle sizes with narrow size distributions and spherical shapes can be easily tuned by just varying the reducing agent-to-precursor concentration in the growth stage in an approach directly applicable to other metal nanoparticles.

Received 25th November 2019,  
Accepted 6th December 2019

DOI: 10.1039/c9re00452a

rsc.li/reaction-engineering

### I. Introduction

Metal nanoparticles have a unique set of optical, electrical and magnetic properties that drives a wide range of applications in the fields of medicine,<sup>1–4</sup> chemistry,<sup>5</sup> agriculture<sup>6</sup> and biology<sup>7–9</sup> to name a few. The final properties required for each application strongly depends on the particle size, distribution,<sup>10–12</sup> composition,<sup>13–15</sup> surface modification (*i.e.* ligands, surfactants, and coordinating solvents)<sup>16,17</sup> and morphology (*i.e.* crystal structure and shape).<sup>18</sup> Thereby, the ability to fine-tune the colloidal properties of nanoparticles *via* wet synthesis<sup>19</sup> have attracted attention over the past decades as a way of developing a manufacturing technology.<sup>20,21</sup> As a result, a vast number of capping ligands,<sup>12,22,23</sup> surfactants<sup>24,25</sup> and reducing agents<sup>26,27</sup> have triggered the development of a large number of protocols both in batch and in continuous systems.<sup>28–34</sup> Regularly, these protocols rely on a presence of ligands and/or high molecular weight reducing agents (*e.g.* tannic acid), which adds an obstructive layer to the nanoparticle surface, being difficult to displace without disturbing the colloidal properties. Despite potential interferences in the final

application,<sup>12</sup> they facilitate the synthesis of particles with narrow size distributions.<sup>35,36</sup> Silver nanoparticles are particularly interesting due to their thermal/physical/optical properties,<sup>37,38</sup> low price (~1/90 of the price of gold),<sup>39</sup> and applications, such as nanotoxicology,<sup>40</sup> catalysis<sup>12,41</sup> and plasmonic devices.<sup>42</sup> The looked-for size ranges from a few nanometres to 100 nm depending on the application. To produce small-sized particles (up to 10 nm), a strong reducing agent is required to trigger rapid nucleation (*e.g.* sodium borohydride, NaBH<sub>4</sub>).<sup>19,40</sup> However, to grow particles over a wide range, a mild reducing agent is desired (*e.g.* as trisodium citrate, Na<sub>3</sub>C<sub>6</sub>H<sub>5</sub>O<sub>7</sub>).<sup>40</sup> Two synthesis routes promote controllable growth by the sequential additions of silver ions, being one discrete (stepwise growth)<sup>35,43</sup> and another continuous (dropwise growth).<sup>40</sup> As an example of a stepwise synthesis, Neus G. Bastús *et al.*<sup>35</sup> synthesized monodisperse spherical silver nanoparticles (standard deviation of 10%) with sizes ranging from 10 to 200 nm. This synthesis involved the reduction of silver nitrate in a batch reactor by a complex formed by trisodium citrate and tannic acid acting as both reducing and stabilising agent by attaching to the silver surface.<sup>44</sup> Shekhar Agnihotri *et al.*<sup>40</sup> applied a dropwise approach with steps with adjusted concentrations of silver nitrate, sodium borohydride, and citrate, producing particles from 5 to 100 nm. Nonetheless, both approaches were carried out in batches, requiring a large number of additions (more than 10) to achieve particle sizes of 100 nm, meticulous

Department of Chemical Engineering and Biotechnology, University of Cambridge, Philippa Fawcett Drive, CB3 0AS, Cambridge, UK. E-mail: lt416@cam.ac.uk

† Electronic supplementary information (ESI) available. See DOI: 10.1039/c9re00452a

preparation and/or post-synthesis purification. Besides, these dropwise methods normally present batch to batch variability and they are difficult to scale-up to continuous unsupervised manufacturing.

The literature is scarce for the continuous synthesis of nanoparticles with size tunability over a wide range. Recently, J. B. Deshpande *et al.*<sup>45</sup> synthesized silver nanoparticles using a CSTR (continuous stirred tank reactor) at pH 7.8, using AgNO<sub>3</sub> and trisodium citrate (Na<sub>3</sub>C<sub>6</sub>H<sub>5</sub>O<sub>7</sub>). Their synthesis strategy leads to high throughput production of particles but did not provide tight control over growth, producing broad particle size distributions with average sizes ranging from ~65 to 120 nm.

Despite the technology advancements, the controllable continuous synthesis of monodisperse particles from 5 to 80 nm is still a challenge. We have recently developed a novel approach for the continuous synthesis of small particles (<5 nm) in the absence of organic ligands<sup>19</sup> by using helical reactor for the promotion of Dean vortexes within the laminar flow (Reynolds <50) enabling the synthesis of small particles with narrow size distributions. We have also demonstrated the separation of the nucleation and growth-stages by rapid changes in chemical (*e.g.* reducing agents) and physical (*e.g.* temperatures) environments using microfluidic reactors connected in series.<sup>46</sup> Despite the promising results, this approach fails on the synthesis of NPs above 10 nm with narrow distribution due to the presence of secondary nucleation.

Taking advantage of the modular nature and hydrodynamic control of microfluidic reactors, this work presents a single step-up for the synthesis of silver particles with tuneable sizes between 5 to 80 nm. We show a distributed feed strategy along the continuous reactor to carefully control the growth rate supported by a novel online rapid (milliseconds time analysis) optical size and shape evaluation. This was possible through the development of a MATLAB mathematical algorithm to convert UV-vis spectra into particle size distribution based on Mie theory, enabling the assessment of the nanoparticle spherical quality (“spherical degree”). Success is only possible thanks to the simultaneous development of the synthetic understanding, reactor design and fast online characterisation to inspect time-resolved nanoparticle population in a flow stream.

## II. Materials and methods

Silver nitrate solution (AgNO<sub>3</sub>, 0.1 M), and sodium hydroxide solution (NaOH, 0.1 M) were purchased from Fisher Scientific. Trisodium citrate dihydrate (Na<sub>3</sub>C<sub>6</sub>H<sub>5</sub>O<sub>7</sub>, ≥99%, FG), citric acid (C<sub>6</sub>H<sub>8</sub>O<sub>7</sub>), bovine serum albumin (BSA), nitric acid (HNO<sub>3</sub>, 65%) and sodium borohydride solution (NaBH<sub>4</sub>, 12 wt% in NaOH 14 M) were purchased from Sigma-Aldrich. Ultra-pure water was obtained from a Merck Milli-Q IQ 7005 water purification system (18.2 MΩ cm resistivity). All chemicals were used without further purification. All the solutions were freshly prepared at the beginning of each experiment. The mixture of trisodium citrate and citric acid was defined as just citrate. The

concentrations for flow synthesis are defined as “concentrations before mixing” unless otherwise stated.

### Continuous synthesis of silver nanoparticles

The synthesis of Ag nanoparticles was carried out inflow using microreactors connected in series as depicted in Fig. 1. In a typical synthesis of particles >10 nm, the nucleation step took place at 90 °C in the first reactor (R0, seed reactor) by mixing 62.5 μL min<sup>-1</sup> of 3 μM NaBH<sub>4</sub> and 20 μM Na<sub>3</sub>C<sub>6</sub>H<sub>5</sub>O<sub>7</sub> and 62.5 μL min<sup>-1</sup> of 1 μM AgNO<sub>3</sub> in a T shape microchannel (ETFE, 0.02 inches through-hole, IDEX Health & Science LLC) using programmable syringe pumps from Chemyx (Fusion 200, accuracy ±0.35%). The temperature was controlled using two types of aluminium heating modules, one for the T-mixer and another for the reactor. Special care should be paid to avoid the degradation of the solutions and thus, the silver solution was kept under darkness and the NaBH<sub>4</sub> and Na<sub>3</sub>C<sub>6</sub>H<sub>5</sub>O<sub>7</sub> in ice-cold to avoid the hydrolysis of NaBH<sub>4</sub>. The silver solution was introduced perpendicular to the continuous flow. Growth took place at 90 °C in reactors R1 to R5 connected in series. At the entrance of each growth reactor, 25 μL min<sup>-1</sup> of an aqueous solution of AgNO<sub>3</sub>, Na<sub>3</sub>C<sub>6</sub>H<sub>5</sub>O<sub>7</sub> and C<sub>6</sub>H<sub>8</sub>O<sub>7</sub> was added with citrate/AgNO<sub>3</sub> concentrations range from 3.5/1 mM to 14/4 mM using T shape microchannels, similar to the one in the seed-stage, and a rack syringe pump. The reaction pH was adjusted in the growth-stage by adding 25 μL min<sup>-1</sup> of a 60 mM aqueous NaOH solution (unless otherwise specified; pH<sub>IN</sub> ~ 12) at the entrance of R1. All the reactors (R0 to R5) were 3 m long PFA tubing (perfluoroalkoxy, IDEX Health & Science LLC) with an outer and inner diameter of 1/16 and 0.03 inches, respectively. The tubing was coiled on an inverted flow helical 3D-printed frame (Fig. 1) (made with a FormLabs Form 1+ stereolithographic printer) with a helix and pitch diameter of 1 cm and 1/π cm, respectively. The outlet of the system was introduced into an ice-cold vial to reduce Brownian motion and, consequently, reduce agglomeration. These flow conditions have a Re between 15 and 23, and De between 4 and 6. The formation of silver oxide was avoided during the preparation of the stock solutions. A cleaning protocol was carried out at the beginning of each experimental day using nitric acid, water, silicone oil and NaOH solutions.

In a typical synthesis of particles between 5 to 10 nm, the system consisted on a seed (R0) and a single growth reactor (R1).<sup>46</sup> The nucleation step was carried out in the seed reactor (R0), by mixing 125 μL min<sup>-1</sup> of 0.05 mM AgNO<sub>3</sub> with 0.025 mM NaBH<sub>4</sub> and 0.35 mM Na<sub>3</sub>C<sub>6</sub>H<sub>5</sub>O<sub>7</sub> at 60 °C. This seed-stage produces particles of 5.1 ± 1.1 nm (~6.6 × 10<sup>12</sup> particles per mL). In the growth reactor (R1), Na<sub>3</sub>C<sub>6</sub>H<sub>5</sub>O<sub>7</sub> and AgNO<sub>3</sub> were pumped at 25 μL min<sup>-1</sup>, and the temperature was set at 90 °C. The final particle size can be fine-tuned by changing the Na<sub>3</sub>C<sub>6</sub>H<sub>5</sub>O<sub>7</sub> to AgNO<sub>3</sub> ratio (6.2 to 9.3 nm using 3.5/0.5 to 14/2 mM/mM, respectively).<sup>46</sup> No NaOH was required to correct the pH.



**Fig. 1** Schematic of the experimental set-up for the continuous synthesis of Ag nanoparticles. The set-up is divided into two stages (R0), seed-stage and growth-stage (R1–R5). In the seed-stage,  $\sim 5$  nm nanoparticles are synthesized using a strong reducing agent,  $\text{NaBH}_4$ . In the growth-stage, citrate is used as a mild reducing agent. At the entrance of each growth reactor, periodic additions of  $\text{AgNO}_3$  + citrate solutions (trisodium citrate + citric acid) are continuously injected. The pH is corrected at the entrance of R1 by the addition of a  $\text{NaOH}$  solution. The inverted coiled reactors used have an internal volume of 1.36 mL each.

### Characterisation of silver nanoparticles

Silver nanoparticles were characterised by transmission electron microscopy (TEM, FEI Tecnai 20), dynamic light scattering (DLS, Malvern Zetasizer Nano S90) and ultraviolet-visible spectroscopy (Agilent Cary 60), using a flow cell. Specimens for TEM analysis were prepared by mixing  $0.3 \text{ g ml}^{-1}$  of the BSA protein with the silver nanoparticle solution in a 1/1 v/v ratio to avoid their agglomeration during sample preparation.<sup>19</sup> The solution was kept at  $5^\circ \text{C}$  for at least 2 hours before preparing the TEM grid. The TEM images were analysed using the image processing software, ImageJ. Ultraviolet-visible spectroscopy measurements were performed on an Agilent Cary 60 UV-vis spectrophotometer in the 300 to 800 nm wavelength range with a resolution of 1 nm. For the DLS analysis, only particles above 10 nm were considered (intensity-based). Particles were analysed without any intermediate purification or separation steps.

## III. Mathematical algorithm for online nanoparticle size and shape evaluation

UV-vis spectroscopy is a widely used analytical technique to characterise the optical extinction of plasmonic metal nanoparticles due to its close dependence on particle size and size distribution (PSD).<sup>47,48</sup> There are unique prospects and superiority over other characterization techniques (DLS, TEM and ICP-MS) due to its millisecond detection, non-disturbance, precision and wide applicability range.<sup>49</sup> In addition, it can be easily connected (online) to the outlet of flow reactors, allowing quick characterisation of colloidal nanoparticles. However, despite its wide application, UV-vis spectroscopy does not provide a direct measurement of particle size distribution and shape. A simple peak position correlation to the particle size is not reliable as it does not account/report: i) peak shape, ii) peak shift due to changes in the chemical environment, iii) presence

of bimodal distributions, and iv) non-sphericity. To provide this information, in this work we have developed a mathematical algorithm based on the Maxwell's equations. Mie theory can accurately predict light scattering and absorption cross-sections for monodisperse spherical particles from a few nanometers up to several thousands of micrometers.<sup>47,50,51</sup> This novel approach enables a fast (*i.e.* millisecond) inspection of particle size, distribution and quality for spherical nanoparticles. This algorithm can be used as a screening tool but it has the potential to be used as exploration tool.

The algorithm is depicted in Fig. 2 and it consists of three parts. The first part corrects the Mie theory optical extinction (*e.g.* size of the particles) by adjusting the dielectric function assuming a mono-size distribution. The second part defines an objective function to account for particle size distribution. The third part establishes a protocol to assure that the optical size-dependent response follows the experimental standards (correlation of UV-vis spectra and TEM microscopy obtained during calibration). Similar approaches for the first and second parts have been previously used with success for gold and silver from  $\sim 5$  to  $\sim 10$  nm.<sup>47,52</sup> The third part is a novel approach to account for the overall complexity of the colloidal system and the lack of reliable dielectric models,<sup>53</sup> mainly due to changes in the chemical environment by the presence of a multitude of species (*e.g.* citrate and  $\text{NaBH}_4$ ) and uncertainties on the complex metal-dielectric function. This algorithm is made openly available to the community in the University of Cambridge open data repository. The key input parameters considered in the algorithm are: 1) spectral wavelengths analysed from 390 to 600 nm with 1 nm resolution, 2) particle size evaluation (radius) from 2 to 80 nm with 0.26 nm resolution, 3) Johnson and Christy dielectric function,<sup>54</sup> 4) Daimon and Masumura water refraction index wavelength dependency.<sup>55</sup> A constrained local minimization approach based on the simplex search method by Lagarias *et al.*<sup>56</sup> was adopted to speed up the algorithm performance, taking about 36 seconds on an i5-350U 1.70GHz computer.



**Fig. 2** A. Schematic of the Mie theory-based calibration algorithm used to estimate particle size and distributions. B. Size-dependent optical response before and after calibration on part III of the algorithm.

### Part I – estimation of optical extinction based on mono-size spherical dispersion

The first part of the algorithm is set to correct the Mie theory spectral response by fitting a size-dependent polynomial function,  $A(r)$ . This function acts on the complex metal-dielectric function (free electrons, intraband transition) by adding a size-dependent contribution, giving a more representative signature of the extinction response by Mie theory ( $\sigma_{\text{ext}}$  or  $Q_{\text{ext}}$ ).<sup>50,53</sup> As it is a fitted empirical function, it can account for some uncertainties on the complex dielectric constants ( $\varepsilon_{\text{Re}}$  and  $\varepsilon_{\text{Im}}$ ), and other factors

affecting the surface plasmon resonance (SPR) (position and width), such as chemical interface dumping (CID) promoted by molecules bonded to the metal surface. This adjustment is key when covering a large set of particle size ranges, from a radius of 5 to 100 nm. If a short range of particle sizes are evaluated, the calibration of  $A(r)$  is straightforward, and a simple regression between the particle size obtained by Mie and TEM is adequate.<sup>19,47</sup> By extending the Mie theory to a broad range of particle sizes, the calibration strategy changes to account for all the size intrinsic effects. Here  $A(r)$  is determined assuming that the commercial samples are monosized. For this, the Mie theory model is used based on the multipole expansion of the electric and magnetic fields of the nanoparticle suspension. The cross-section extinction response ( $\sigma_{\text{ext}} = Q_{\text{ext}} \pi r^2$  [m<sup>2</sup>]) for a spherical particle with  $r$  radius depends on the metal-dielectric function and the surrounding dielectric environment (*i.e.*, chemical environment), and it can be expressed as:

$$\sigma_{\text{ext}} = \frac{2\pi}{k^2} \sum_{L=1}^{\infty} (2L+1) \text{Re}(a_{L^*} + b_{L^*}) \quad (1)$$

where  $Q_{\text{ext}}$  is the extinction efficiency,  $r$  [m] is the particle radius,  $a_{L^*}$  and  $b_{L^*}$  are the Mie coefficients in term of the spherical Riccati–Bessel functions,<sup>57</sup>  $k$  is the wave vector ( $k = 2\pi\varepsilon_{\text{m}}/\lambda$ ),  $\varepsilon_{\text{m}}$  is the real refractive index of the surrounding medium,  $\lambda$  is the evaluated wavelength, and  $L^*$  is the multipolar order, which has been set to 3. The terms  $a_{L^*}$  and  $b_{L^*}$  depends on the ratio between refractive index of the metal and the medium,  $m = \varepsilon(\lambda, r)/\varepsilon_{\text{m}}$ . The  $\varepsilon(\lambda, r)$  is the complex refractive index (or dielectric function) of a metal sphere with  $r$  radius, and it can be defined as:

$$\varepsilon(\lambda, r) = \varepsilon_{\text{Re}}(\lambda) + \varepsilon_{\text{Im}}(\lambda)i + \Delta\varepsilon(\lambda, r) \quad (2)$$

where  $\varepsilon_{\text{Re}}$  and  $\varepsilon_{\text{Im}}$  are the real and imaginary refractive index relative to the metal nanoparticle, respectively, being set according to Johnson and Christy.<sup>54</sup> To account for surface effects, the dielectric function is written to account for the particle size dependency, which follows the Drude-based model relationship. It considers that electrons are free and can be forced to oscillate by an electromagnetic wave with a frequency  $\omega$ .<sup>50</sup>

$$\Delta\varepsilon(\lambda, r) = \frac{\omega_p^2}{\omega} \left( \frac{1}{\omega + i\omega\Gamma_{\infty}} - \frac{1}{\omega + i\Gamma(r)} \right) \quad (3)$$

$$\Gamma(r) = \Gamma_{\infty} + A(r)v_{\text{F}}/r \quad (4)$$

where  $\omega = 2\pi c/\lambda$  [rad s<sup>-1</sup>] is the bulk angular frequency with  $c$  [m s<sup>-1</sup>] as the speed of light, and  $\omega_p$  [rad s<sup>-1</sup>] is the resonance frequency (also named plasma frequency, 9.05 e/h for silver),  $\Gamma_{\infty}$  [s<sup>-1</sup>] is the bulk metal value frequency (0.0205 e/h for silver),  $\Gamma(r)$  [s<sup>-1</sup>] is the size-dependent

relaxation frequency,  $A(r)$  is a polynomial empirical function, and  $v_F = 1.39 \times 10^6$  [m s<sup>-1</sup>] is the Fermi velocity. The frequency ( $\omega$ ) is often reported in eV units as  $\omega\hbar/e$  [eV], where  $\hbar$  is the reduced Planck's constant,  $1.054 \times 10^{-34}$  [J s], and  $e$  is electronvolt to Joules conversion,  $1.6 \times 10^{-19}$  [J eV<sup>-1</sup>].

The experimental data  $[\ln(I_0/I)]$  is fitted by the cross-section extinction response from the Mie theory  $[\sigma_{\text{ext}}(\lambda, r, \varepsilon)]$  through an objective function. On its classical and simplest form, the least-square method can be defined for a monodisperse spherical sample as:

$$f_{\text{obj},1} = \min_r \sum_{\lambda_0}^{\lambda_r} \left( \frac{\ln(I_0/I)_{\lambda}}{\text{Exp.data}} - \frac{\sigma_{\text{ext}}(\lambda, r, \varepsilon) LC_p}{\text{Lambert Beer law}} \right)^2 \quad (5)$$

where  $(I_0/I)_{\lambda}$  is the experimental extinction response of monochromatic light at wavelength  $\lambda$ ,  $L$  [m] is the optical length and  $C_p$  is the total number of particles per volume of solution [particles per m<sub>sol</sub><sup>3</sup>].

### Part II – definition of an objective function to account for particle size distribution

The objective function  $f_{\text{obj},1}$  does not account for particle size distribution (PSD) and resonance shifts by changes in the chemical environment. To account these, the following equations were derived (refer to the ESI† for more details).

$$\text{RSS} = \sum_{\lambda_0}^{\lambda_r} \left( \frac{\ln(I_0/I)_{\lambda+\lambda_{\text{shift}}}}{\text{Exp. data}} - \underbrace{\frac{LN_p}{4} \int_{r_0}^{r_r} \frac{Q_{\text{ext}}(\lambda, r, \varepsilon)}{r} f(r) d(r)}_{\text{Calculated}} \right) \quad (6)$$

$$f_{\text{obj},2} = \min_{f(r)\lambda_r} \left( \beta \frac{\text{RSS} + \alpha \|f(r)\|_2^2}{\sum_{\lambda_0}^{\lambda_r} \ln(I_0/I)_{\lambda+\lambda_{\text{shift}}}} + (1-\beta) |\lambda_{\text{shift}}| \right) / \lambda^{\max(I/I_0)} \quad (7)$$

$$\lambda_{\text{shift}} = \lambda^{\max(I/I_0)} - \lambda^{\max(\text{Calculated})} \quad (8)$$

where  $N_p$  [particles] is the total number of particles,  $\lambda_{\text{shift}}$  [nm] is the peak location difference between the experimental and calculated data. The RSS (residual sum of squares) corresponds to a direct fit between experimental and calculated optical responses, and it takes into account the contribution of PSD in a form of volume frequency distribution [m<sub>particle</sub><sup>3</sup> m<sub>sol</sub><sup>-3</sup>],  $0 \leq f(r) \leq 1$ , here defined as a normal distribution. The other terms ( $\alpha$ : Tikhonov regularization term,  $\beta$ : weight between RSS and peak location) present in the eqn (7) assure an accurate mathematical convergence. This equation differs from the literature<sup>58</sup> as it accounts for  $\lambda_{\text{shift}}$  and  $\beta$ . Mie calculated curves were allowed to shift ( $\lambda_{\text{shift}}$ ) to better match the

position of the experimental spectrum.<sup>47</sup> To control the maximum shift acceptable,  $\beta$  was empirically set to 0.1.

### Part III – determination of the particle size by correction of optical responses

Spectral responses (frequency and size dependency) for metal colloids are strongly dependent on the chemical environments (e.g. presence of a multitude of species such as citrate and NaBH<sub>4</sub>) and the uncertainties on the complex metal-dielectric function. Typical dielectric corrections account for the contribution of the free (e.g. Drude-Sommerfeld) and bound electrons.<sup>50,53</sup>

The ideal way of determining  $A(r)$  is by directly solving the  $f_{\text{obj},2}$  with a set of commercial standards, using their PSD (mean and standard deviation) as an input. Unfortunately, this methodology is computationally intense and thus, a trade-off between computational time and accuracy was selected, and  $A(r)$  was calculated in part I assuming the standards were mono-size particles. Since the correct dielectric function does represent the whole particle size distribution (PSD), a linear regression between the calculated and the real PSD (TEM-based) is created [ $r_{\text{cal}} = g(r_{\text{TEM}})$ ]. The result obtained is used to correct the final Mie theory response, in a form of  $\sigma_{\text{ext}}(\lambda, g(r), \varepsilon)$  to finally calculate the actual particle size and distribution.

### Validation of the Mie theory-based algorithm – optical size distribution evaluation

To validate the Mie theory-based algorithm, the UV-vis spectra of a range of commercial standards were used to estimate their particle size and distribution. The results were compared to the actual values determined by TEM microscopy as shown in Table 1. The calculated radius and PSD present a relative deviation of 6.9% and 50.4%, respectively. While the calculated particle sizes are accurate and reliable, the PSD is underestimated, due to the trade-off between computational time and accuracy in the calibration process and/or the choice of the optimization method (local minimum to reduce running time). This high deviation on the particle size distribution is due to a mathematical inverse problem, which is stated as an ill-posed problem (Fredholm integral equation of first kind). This means that a system does not have a unique solution, but a multitude of solutions. Narrow the number of solutions can be achieved by a regularization parameter. Herein, we used a classical Tikhonov methodology coupled with a local minimum optimization algorithm. However, one should notice that the commercial standards normally present very narrow size distributions (<15–20%), where small deviations in the estimated PSD (i.e. 1.0 nm instead to 0.8 nm) lead to large deviations (i.e. 22.5%).

The mathematical algorithm was further validated using a large set of samples (>100) synthesised in flow reactors during this study. The samples were analysed by DLS

**Table 1** Comparison between actual size (measured by TEM) and calculated sizes (using Mie theory-based algorithm) of a range of commercial standards

Sample supplier	TEM		Calculated using algorithm		Deviations	
	Average radius (nm)	Radius std (nm)	Average radius (nm)	Radius std (nm)	Average radius relative deviation (%)	Radius std relative deviation (%)
A	4.1	0.8	4.9	1.0	19.5	22.5
B	5.0	1.0	5.0	1.2	1.1	21.7
C	11.0	2.5	12.7	1.0	15.3	59.2
B	14.5	1.5	13.5	0.9	7.2	36.9
B	20.5	1.5	21.5	1.0	5.0	31.4
C	27.3	3.4	24.8	1.3	9.3	61.1
B	36.5	4.0	35.9	1.4	1.5	66.0
C	37.9	5.6	40.5	1.4	6.6	75.1
B	39.5	3.5	38.4	1.3	2.9	62.9
B	47.5	4.5	47.9	1.5	0.7	67.2

Sample (A) was synthesised according to literature.<sup>46</sup> (B) are samples from Nanocomposix and (C) from BBi.

(intensity-based) after dilution and without any further purification and/or separation. As the DLS usually overestimates the particle size, its value was corrected using a TEM-DLS correlation of commercial standards from BBi (see Fig. 3A where  $d_{DLS}$  is shown as a function of  $d_{TEM}$ ). Calibrated DLS particle sizes show a good agreement with the Mie theory calculated particle sizes as shown in Fig. 3B. It is worth noting that DLS has a considerably higher relative deviation ( $\sim 18.7\%$ ) compared to the Mie theory calculated size deviation of 6.9% (Table 1).

### Optical nanoparticle shape evaluation – spherical degree

Mie theory cannot predict the optical response of non-spherical particles, such as spheroids, cylinders, rods, prisms and disks. Fig. 4 (and S1†) shows a characteristic UV-vis spectra, size distribution, shape and aspect ratio sample [ $A_R$ ] defined as the ratio between the shortest ( $r_A$ ) and the highest radius ( $r_B$ )] of spherical and non-spherical nanoparticles

respectively. In the spherical samples (Fig. 4A), there is a good agreement between experimental and calculated (using the developed algorithm) spectral signatures (width and resonance position). However, when analysing a sample with a majority of non-spherical particles (Fig. 4B), there is a clear difference between both spectral signatures. In this case, the experimental absorption spectrum is non-Gaussian, being composed by at least two absorption bands, blue- and red-shifted. Other methodologies/tools can be used to account for non-spherical shapes using finite element approaches, such as discrete dipole approximation (DDA) and boundary element method (BEM). In this work, we use this property of the Mie theory to add shape evaluation capabilities to Mie theory through our new methodology. In this context, spherical particles are stated as a group of quasi-spherical particles with the underlying symmetries of face centre cubic (fcc) lattice.

In this sense, several statistical parameters were considered to evaluate the “spherical degree” of the silver nanoparticles by comparing the experimental spectra of five citrate-based synthesized samples with spherical and non-spherical particles against the spherical standards, including RSS (residual sum of squares) (eqn (6)), the objective function in the algorithm (eqn (7)), FWHM of the UV-vis spectra (full width at half maximum), skewness<sub>EXP</sub>/skewness<sub>Mie</sub> and kurtosis<sub>EXP</sub>/kurtosis<sub>Mie</sub> (refer to ESI† for further details). We found that the RSS is the most suitable parameter as “spherical degree” (refer to Fig. S2†) because it better delimits the region between a sample with spherical and non-spherical particles. To define the RSS threshold region, individual samples with spherical (5.4 and 79.8 nm,  $\sim$ yellow/red colloidal) and non-spherical (majority of triangular prisms, blue colloidal) particles were mixed at different optical density ratios, Fig. 5A. There is a clear difference between the spectral shape and peak location for spheres and triangular prisms NPs. Silver spherical NPs show a blue-shift quasi-Gaussian peak ( $\sim 389$  to  $460$  nm<sup>-1</sup>), however triangular prisms NPs show a red-shift non-Gaussian peak ( $\sim 630$  nm<sup>-1</sup>). Using different spherical/non-spherical



**Fig. 3** A. Validation of the Mie theory-based algorithm using  $>100$  synthesized Ag nanoparticles samples characterised by DLS. Data plotted in function of the spherical degree (RSS). B. Calibration of the DLS intensity-based particle size by TEM of commercial particles from BBi.



Fig. 4 Characterization of two synthesized samples with different particle shapes and sizes, A – quasi-spherical (spherical and spheroidal) particles and B – non-spherical particles, by spectral and TEM analysis. The experimental optical response (red) was compared to the estimated one using the Mie theory-based algorithm (blue).

mixtures, a RSS threshold to  $\leq 0.4$  (grey zone in Fig. 5B) is determined which represents a spherical to triangular optical ratios up to 0.15. In other words, a blue-shifted spectrum with  $1 \text{ a.u. cm}^{-1}$  can have up to  $0.15 \text{ a.u. cm}^{-1}$  of non-

spherical particles. If a higher spherical content is required, the value of RSS threshold can be decreased. In summary, a spherical degree is determined for the accurate use of the Mie theory-based algorithm where particles with a  $\text{RSS} \leq 0.4$  are considered spherical/spheroidal. This zone depends on the algorithm properties, such as spectral resolution, wavelength range, particle size range analysed, *etc.* and thus, it should be reconsidered if different equipment or experimental protocol is used.

#### IV. Synthesis of silver nanoparticles with tuneable sizes

Silver nanoparticles with tuneable sizes between 10 and 80 nm were synthesised using a series of microfluidic reactors connected in series as shown in Fig. 1 where the outlet nanoparticles were optically characterised to evaluate the size, size distribution and shape using the developed Mie-theory algorithm. It is important to highlight that based on its principle, the algorithm is also able to detect any phenomena affecting the spectral response of the particles, not only growth but also aggregation or dissolution. Microreaction technology enables precise control over the hydrodynamics in the system as we have recently demonstrated by the use of inverted flow helical reactors to promote the formation of Dean vortices under laminar regime. The presence of Dean vortices improves the radial mixing and narrow residence time distribution, both leading to narrow particle size distribution and minimum agglomeration even in the absence of steric capping ligands.<sup>19</sup> Electronic stabilisation by citrate anions is used to avoid the agglomeration of the particles. However, such stabilisation does not limit particle



Fig. 5 A. Spectral response of mixtures of spherical and non-spherical (majority of triangular prisms) silver particles. B. RSS as a function of the optical density ratio of triangular to spherical particles. The RSS was calculated based on the Mie theory algorithm applied to the spectra present on A. The grey rectangular limits the zone where spherical/spheroids particles are in the majority of the sample ( $\text{RSS} \leq 0.4$ ).

growth by steric effects. To control growth nanoparticles, an optimised reactor configuration with  $Re \approx 23$  and  $De \approx 6$  (at the end of last reactor) was used.

To achieve narrow size distributions, careful control of the nucleation and growth stages is essential which can be exceptionally achieved by rapid changes in the chemical environment of the different reactors connected in series.<sup>59</sup> In this case, a seed-mediated growth method was used for the synthesis of the silver particles where nucleation takes place in the first reactor (seed stage, R0, Fig. 1) followed by controlled growth in reactors (growth stage, R1 to R5, Fig. 1), depending on the desired size of the particles. The resulting nanoparticles were optically characterised online to evaluate the size, size distribution and shape using the developed Mie-theory algorithm.

The nucleation step is carried out using a strong reducing agent ( $\text{NaBH}_4$ ) to trigger rapid nucleation which leads to small particles ( $\sim 5$  nm) and narrow particle size distribution ( $<20\%$ ). A relatively long residence time ( $\sim 8$  min) is used to ensure full conversion of the silver ions as well as full degradation of the  $\text{NaBH}_4$ .<sup>19</sup> Both conditions are essential to avoid secondary nucleation in the consecutive growth reactors. Particle agglomeration is kept at a minimum by the presence of citrate ( $20 \mu\text{M}$ ) to electrostatically stabilise the particles. Once that the seeds are formed, controllable growth takes place in reactors R1 to R5 (Fig. 1) by adding silver ions at the entrance of each growth reactor. A mild reducing agent (citrate) is also added to promote the autocatalytic reduction of silver ions and limit secondary nucleation to a minimum.<sup>35,60,61</sup> The sequential addition of silver precursor and reducing agent enables a controllable and slow growth in a reliable and versatile way. Similarly to the seed-stage, a relatively long residence time ( $\sim 8$  min) is used to ensure full conversion of the silver ions in each of the growth reactors.

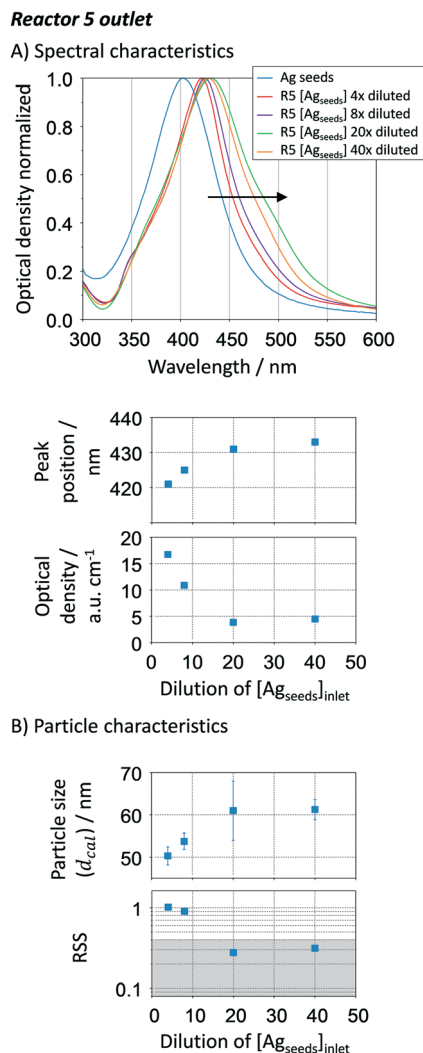
### Effect of Ag seed-to-growth solution ratio

Seed-mediated growth relies on the autocatalytic reduction of metal salts in the metal nanocrystal surface.<sup>46</sup> The activation energy barrier is much lower in this process when comparing to spontaneous nucleation.<sup>62</sup> Defining the seed concentration is key to control growth and consequently the resulting size. Depending on the seed concentration, two different growth mechanisms (diffusion- or kinetic-controlled) can dominate this stage,<sup>63</sup> the former believed to lead to relatively narrow size distribution.<sup>61</sup>

To optimise the initial seed to silver precursor concentration added in the growth-stage (Fig. 1), a seed concentration of  $\sim 2 \times 10^{10}$  particles per mL with an estimated size of  $34 \pm 7$  nm and RSS of 0.16 was systematically varied after the seed reactor (R0) by dilution. The dilution degree was defined as in the volumetric form. In the growth stage (R1 to R5), the sequential additions contained a  $\text{AgNO}_3/\text{citrate}$  ratio of 3 mM/14 mM.

Increasing the Ag seed dilution from 4 to 40 (*i.e.* decreasing the Ag seed concentration) leads to a red-shift of the UV-vis

spectra of the resulting silver NPs after reactor R5, at the same time that the optical density decreases (Fig. 6A). The plasmon resonance shift (421 nm to 433 nm) is characteristic of a gradual increase in particle size, which is confirmed by the Mie theory-based algorithm (Fig. 6B),  $\sim 50$  to  $\sim 60$  nm. To increase particle size beyond 60 nm, a considerably higher amount of silver precursor in the growth stage is needed. Peak location and optical density indicate that when increasing the seed dilution, a lower number of particles with higher particle sizes are being formed. This is an expected result as the total number of seeds decreases, leading to a subsequent decrease



**Fig. 6** Effect of the initial Ag seed to Ag growth concentration ratio. A. UV-vis spectra spectral peak location and optical density (peak absorbance/optical length) for the Ag particles after R5; B. particle size distribution (average and standard deviation) and spherical degree (RSS) calculated using the Mie theory-based algorithm. Conditions: nucleation reactor  $\rightarrow 60$  °C, 0.025 mM  $\text{NaBH}_4$ , 0.35 mM  $\text{Na}_3\text{C}_6\text{H}_5\text{O}_7$  and 1.25 mM  $\text{AgNO}_3$ , total flow rate:  $250 \mu\text{L min}^{-1}$ , residence time: 5.4 min (seed concentration of  $2 \times 10^{10}$  particles per mL). 5 growth reactors  $\rightarrow 90$  °C, sequential addition of  $25 \mu\text{L min}^{-1}$  of  $\text{AgNO}_3/\text{citrate}$ : 3 mM/14 mM, total residence time: 31.2 min, initial  $\text{pH}_{\text{IN}} \sim 12$  (NaOH added at  $25 \mu\text{L min}^{-1}$ ). All the concentrations are present before mixing.

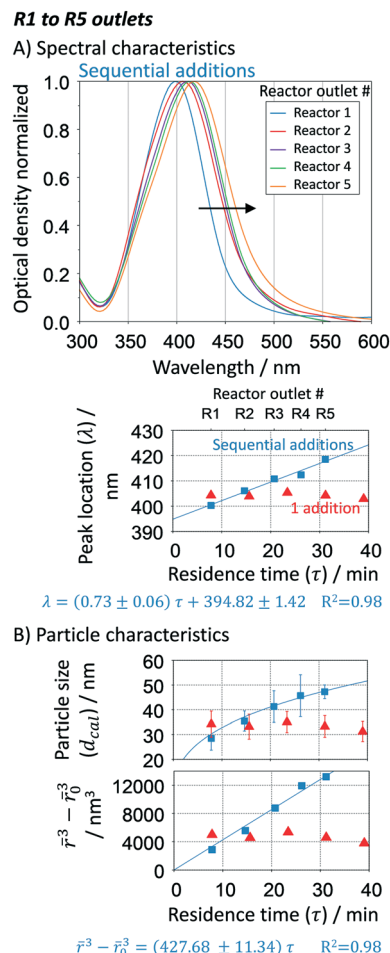


of the solid surface area. The slowdown is beneficial towards particle growth as it adds more control over the kinetics, being closer to a diffusion-controlled regime. The particle size distribution, represented by the error bars in Fig. 6B, does not seem to be affected by the Ag seed-to-growth solution ratio within the studied range. On the other hand, the RSS value decreases as the dilution factor increases, ranging from 1 to 0.31 (Fig. 6B), which suggests that the Ag seed concentration plays a crucial role not only on particle size but also on the shape. To guaranty spherical particles with a narrow size distribution, an optimum Ag seed to Ag growth concentration of  $\sim 5 \times 10^8$  NP mL<sup>-1</sup> (in this particular case requiring a dilution factor of 40) is determined which corresponds to  $\sim 1.8 \times 10^{-6}$  m<sup>2</sup> Ag per mL.

### Effect of number of growth reactors

In order to synthesise monodispersed particles and control the growth stage, the effect of the number of reactors was evaluated. Keeping the metal ions concentration during the growth stage<sup>63,64</sup> below the critical supersaturation and far above the solubility level is of the utmost importance to limit secondary nucleation and thermodynamically driven size focusing, respectively, which are responsible for particle polydispersity.<sup>61,65</sup> In other words, it is possible to growth particles in a kinetic regime (surface reduction). In batch synthesis, these optimal conditions are achieved by a dropwise addition or stepwise method which potentially leads to concentration profiles across the reactors. In our study, we overcome this issue by distributing the feed at the inlet of several growth reactors (R1 to R5). To control a size focus growth rate, it is necessary to have a fine balance between the number, concentration and flow rate of the additions of silver ions and reducing agent.

A single addition (*i.e.* no further additions in the growth reactors) shows that particles only grow in R1 (up to 34 nm) and its size, size distribution and optical density ( $\sim 1.4$  a.u. cm<sup>-1</sup>) does not change by increasing the residence time (*e.g.* increasing number of reactors from R1 to R5 without additional additions, Fig. 7B red symbols). Only after R3 ( $\sim 24$  min), the intra-particle thermodynamically driven size focusing starts to slightly impact the particle size. Such size focusing phenomena can be easily reduced by adding Ag<sup>+</sup> content at the entrance of R2. If the concentration of AgNO<sub>3</sub> is increased in the single addition, the solution stability is compromised due to the precipitation of silver oxide, leading to uncontrollable growth as well as secondary nucleation according. When performing continuous additions (R1 to R5), the absorption spectra show a clear shift towards higher wavelengths, Fig. 7A. This spectra behaviour is characteristic from a gradual increase in the particle size, which is confirmed by our Mie theory-based algorithm (Fig. 7B), going from  $\sim 28$  nm to  $\sim 50$  nm. For all the results analysed, a high spherical degree is observed (RSS  $\leq 0.4$ ). A closer look to the particle size distributions reveals that the particle cubic radius increases linearly as a function of residence time [ $\bar{r}^3 -$



**Fig. 7** Effect of the number of growth reactors. A. UV-vis spectra and peak location at the outlet of each reactor; B. particle size distribution (average and standard deviation) calculated using Mie theory-based algorithm and regression of the particle cubic radius as a function of the residence time. Conditions: nucleation reactor  $\rightarrow 90$  °C, 3  $\mu$ M NaBH<sub>4</sub>, 20  $\mu$ M Na<sub>3</sub>C<sub>6</sub>H<sub>5</sub>O<sub>7</sub> and 1  $\mu$ M AgNO<sub>3</sub>, total flow rate: 125  $\mu$ L min<sup>-1</sup>, residence time: 10.9 min. Growth reactors  $\rightarrow$  refer to Fig. 6. Single addition  $\rightarrow$  only 1 addition of AgNO<sub>3</sub>/citrate (3 mM/14 mM) at R1 instead of 5.

$\bar{r}_0^3 = K\tau$ , where  $K$  is the volumetric growth rate (nm<sup>3</sup> min<sup>-1</sup>),  $\bar{r}_0$  is the initial average particle radius, and  $\tau$  is the residence time as  $\tau = \frac{V}{Q} = \frac{\text{reactor vol.}}{\text{flow rate}}$ . The linear regression between changes in size and residence time is characteristic of diffusion-controlled growth.<sup>66</sup> This key observation indicates high control over the particle growth with a single growth rate across all the reactors (R1 to R5),  $K = 427.68 \pm 11.34$  nm<sup>3</sup> min<sup>-1</sup>. The quality of the linear regression is unexpectedly high, which excludes particle growth governed by stochastic collisions and the presence of secondary nucleation. The particle growth seems to be mostly governed by size focus, where the ion reduction happens at the surface of the silver precipitates (in a kinetic regime). Nevertheless, since this is a continuous process, thermodynamically driven size focusing may be present as the ion concentration decreases inside each growth reactor. To suppress it, an infinite number of

reactors/additions are necessary (*e.g.* mimicking the dropwise methodology).

These results demonstrate that continuous synthesis in microfluidics coupled with the sequential addition of precursors can provide exquisite control over the growth of spherical nanoparticles due to their low diffusion time. This unique approach enables a particle growth up to ~50 nm with just 5 reactors (34 min residence time) without user intervention. However, it is important to note that the growth rate is highly sensitive to external factors such as degradation of the stock solutions, solution preparation, pH variations, fouling in the reactor, flow fluctuations, *etc.*

### Effect of pH in the growth reactors

The pH is known to impact the growth of the silver seeds<sup>67,68</sup> as it affects not only the Ag<sup>+</sup> reduction rate but also the speciation of citrate<sup>69</sup> and silver species. In addition, the presence of different citrate (Cit) and dicarboxylic acid (DCA) species formed as a by-product of the decarboxylation of citrate during the silver reduction:

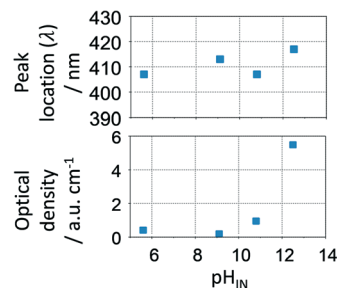


To study the pH effect on continuous particle growth, a feed of NaOH was added at the entrance of the first growth reactor, R1 (Fig. 1) to vary the initial pH (pH<sub>IN</sub> at the inlet of R1) from 5.6 to 12.5 (sample S1 to S4). Further acidic conditions (pH < 5) were not tested as they would lead to the dissolution of silver nanoparticles. The rest of the growth conditions were kept constant with 25 μL min<sup>-1</sup> sequential additions of AgNO<sub>3</sub> to citrate ratio of 3 mM/14 mM in each growth reactor. It is important to note that the seed reactor is kept at pH ~ 7, since speciation is not relevant when NaBH<sub>4</sub> is used as a reducing agent.

Fig. 8A shows the optical response at the outlet of R5 as a function of the inlet pH (pH<sub>IN</sub>). As the pH<sub>IN</sub> decreases, the absorbance peak becomes broader, non-Gaussian, its location shifts towards lower wavelengths and the optical density greatly decreases. Variations in optical density confirm changes in the Ag<sup>+</sup> reduction rate. Regarding the other peak characteristics, the Mie theory-based algorithm suggests that these variations have a small impact on the average particle size, ranging from ~40 to ~50 nm (Fig. 8B). However, higher standard deviations and RSS (above the threshold of 0.4) are obtained in all cases except when pH<sub>IN</sub> is ~12, indicating the formation of non-spherical particles. Although the reduction rate is slower when pH<sub>IN</sub> < 12 (confirmed by lower optical density), it cannot explain (by itself) the broader size distributions and the differences in RSS values. There are two plausible explanations for this behaviour. The first one is linked to the presence of secondary nucleation due to the accumulation of silver ions between reactors, which might lead to hyper saturation. However, secondary nucleation by itself cannot explain the presence of much broader absorption peaks (at the calculated sizes) when compared to

### Reactor 5 outlet

#### A) Spectral characteristics



#### B) Particle characteristics

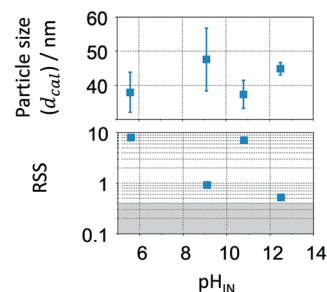


Fig. 8 Effect of initial pH (pH<sub>IN</sub>) in the growth-stage. A. UV-vis spectra, spectral peak location and optical density at the outlet of R5; B. particle size distribution (average and standard deviation) and spherical degree (RSS). Conditions: nucleation reactor → refer to Fig. 7. 5 growth reactors → same as Fig. 6 and 7 except for changes in the pH<sub>IN</sub>.

pH ~ 12.5. This is because it will lead to low growth and lower particle sizes. Thermodynamically driven size focusing is not likely to be present since the ion concentration is at the vicinity of the supersaturation critical level,<sup>65</sup> where it can be neglected. The second explanation is related to the speciation of citrate and dicarboxylic acid (DCA),<sup>70</sup> detailed in Fig. 9.

In the pH<sub>IN</sub> ranging from 5.6 to 10.8, the dicarboxylic acid is completely hydrolysed to DCA<sup>2-</sup> while a mixture of H<sub>2</sub>Cit<sup>-</sup> and HCit<sup>2-</sup> are present. For pH<sub>IN</sub> ~ 12, the resulting nanoparticles have a considerably narrower absorption peak (almost Gaussian) due to the presence of completely hydrolysed DCA<sup>2-</sup> and Cit<sup>3-</sup> species. These results demonstrate that the presence of H<sub>2</sub>Cit<sup>-</sup> and/or HCit<sup>2-</sup> seems to play a role in the non-spherical shape of the resulting silver nanoparticles, likely by selectively binding to the silver nanoparticles leading to anisotropic growth.<sup>43,71</sup> Besides, Cit<sup>3-</sup> species present at pH > 10.8 (*i.e.* S4) has a considerably



**Fig. 9** A. Mole fraction of the different hydrolysed citrate and dicarboxylic acid species as a function of pH. The  $pK_a$  values for citrate hydrolysis are 3.1, 4.76, 6.4, respectively; the  $pK_a$  values for dicarboxyaceton are 2.74 and 3.72, respectively. B. The pH of the samples S1 to S4 at the entrance of the first growth reactor (R1) and at the end of the system (R5, outlet). For S1, all growth reactors operate at similar pH levels due to the buffer properties of the citrate species. For S2 and S3, a significant decrease in pH is observed, caused by the dilution of NaOH, the sequential additions and the buffer properties of citrate species. For S4, the pH decrease is related to the dilution of NaOH and sequential additions.

higher reducing potential, leading to a faster reduction of the silver ions during the growth-stage, translated in a narrow size distribution due to homogeneous growth. It is important to note that small changes in the initial pH (even below 0.1) can lead to big changes in the speciation of the citrate species. During the continuous synthesis these species can have detrimental effects on the reproducibility of the synthesis if not carefully controlled.

The optimum pH value  $> 10.8$  in the growth-stage is not in agreement with previously reported studies in batch with dropwise/stepwise addition of NaOH, where an optimum pH value is between 7.7 and 10.5 was determined.<sup>40,72</sup> It is important to note that our approach at high pH produces large spherical particles with (RSS close to 0.4), high particle concentration (high optical density) and low residence time ( $\sim 34$  min). The differences can potentially be related to inefficient mixing in batch reactors which can lead to localised pH and concentration gradients. This gets worse when the ion reduction rate increases as the mixing time remains high. In contrast, the present microfluidic device enables unique control over the fluid mechanics and mass transfer. This is due to its inherent laminar flow regime, presence of Dean vortices,<sup>19,73</sup> low mixing time,<sup>30</sup> and high mass transfer.<sup>73</sup> The mixing time is  $\sim \frac{x^2}{D} \sim 3$  ms, where  $x$  is the channel width and  $D$  the diffusion coefficient, which for  $AgNO_3$  in water is  $1.54 \text{ cm}^2 \text{ s}^{-1}$ .<sup>74</sup>

### Effect of citrate/ $AgNO_3$ ratio in the additions in the growth-stage

The citrate/ $AgNO_3$  concentration ratio in the sequential additions (inlet, R1 to R5) has a strong effect on the final size, distribution and shape of the resulting silver nanoparticles. Its impact was evaluated by varying the citrate/

$AgNO_3$  ratio between 3.5/14 mM and 0.1/4 mM, respectively, while keeping constant the remaining conditions. Fig. 10A shows the spectral response of the nanoparticles at the outlet of R5. Increasing the citrate to  $AgNO_3$  ratio leads to a decrease in the absorption peak wavelength, which is translated into an expected decrease in particle size (Fig. 10B). Interestingly, there is a clear trend between growth and reducing-to-precursor concentration. If citrate to  $AgNO_3$  ratio is above  $\sim 40$ , size focusing is expected to dominate particle growth as the  $Ag^+$  concentration is far from the vicinity of the supersaturation concentration (here set as 0.1 mM of  $AgNO_3$ ). The growth is minimum leading to small



**Fig. 10** Effect of the [citrate]/ $AgNO_3$  ratio of the sequential additions in the growth-stage. A. UV-vis spectra, peak location and optical density of particles in R5. B. Calculated average particle size and spherical degree (RSS) using the Mie theory-based algorithm. For smaller particle sizes, the number of growth reactors should be reduced. Conditions: nucleation reactor  $\rightarrow$  refer to Fig. 5 growth reactors  $\rightarrow$  same as Fig. 6 and 7 except for changes in the citrate/ $AgNO_3$  ratio.

particle sizes, large size distributions (represented by the error bars in Fig. 10B), and large RSS values. The RSS confirms that such uncontrollable system is accompanied by the formation of non-spherical particles. Anisotropic structures are known to be formed when the citrate concentration is high.<sup>72</sup> On the other end, when the citrate to AgNO<sub>3</sub> ratio is below 0.8, it is expected to have the formation of large lumps by the coalescence of destabilized particles due to the poor electrostatic stabilisation under these conditions. Therefore, this range was not tested. Within these two limits, changing the citrate to AgNO<sub>3</sub> ratio between 0.8 and 40, provides a smart approach to tune up the particle size from 40 to 80 nm with narrow size distributions and high spherical degree. As previously discussed, keeping the Ag<sup>+</sup> concentration within a certain range by sequential additions, focusing growth mechanism dominates over size focusing, even if the Ag<sup>+</sup> concentration decreases along the reactor. For smaller particle sizes, one should decrease the number of growth reactors as shown above (*i.e.* collect the resulting nanoparticles in R1 instead of R5).

## V. Conclusions

In this work, we demonstrate the continuous synthesis of silver nanoparticles with tuneable sizes between 5 and 80 nm (biggest range demonstrated to date), narrow size distribution and high spherical degree. This flexible manufacturing technology allows not only the separation of nucleation and growth, but also an exquisite growth control by maintaining the concentration range in a number of reactors connected in series. Online size and shape evaluation facilitates systematic parametric studies through the development of a mathematical algorithm based on the Mie model solution of the Maxwell's equations. In this way, we demonstrate that size tunability can be easily achieved by changing the number of reactors and the citrate to silver precursor ratio added in the growth-stage, between 0.8 and 40. To ensure spherical shapes and narrow size distribution, the Ag seed to Ag growth surface area should be  $\sim 1.8 \times 10^{-6}$  m<sup>2</sup> Ag per mL and the growth should take place at pH levels above 10.8. This approach provides a unique control and understanding of particle growth in continuous, mostly thanks to the online optical evaluation of size and shape, which can apply to the synthesis of other metal nanoparticles.

## Abbreviation list

$A(r)$  Empirical polynomial function to adjust the metal-dielectric function  
 $a_{L^*}$  and  $b_{L^*}$  Mie coefficients in term of the spherical Riccati-Bessel functions  
 $c$  Speed of light  
 $C_p$  Total number of particles per volume of solution  
 CSTR Continuous stirred tank reactor  
 $d$  Particle size expressed in diameter

De Dean dimensionless number  
 DLS Dynamic light scattering  
 $e$  Electronvolt to Joules conversion  
 $g(r)$  Polynomial function to fit calculated and TEM particle radius  
 $\hbar$  Reduced Planck's constant  
 ICP-MS Inductively coupled plasma mass spectrometry  
 $(I_0/I)_\lambda$  Experimental extinction response of monochromatic light at wavelength  
 $k$  Wave vector  
 $K$  Growth rate  
 $L$  Optical length  
 $L^*$  Multipolar order  
 $N_p$  Total number of particles  
 Optical density Spectral absorbance divided by the optical length ( $L$ )  
 PSD Particle size distribution  
 $Q$  Volumetric flow rate  
 $Q_{\text{ext}}$  Extinction efficiency  
 $r$  Particle radius  
 $R\#$  Reactor number #  
 Re Reynolds dimensionless number  
 RSS Residual sum of squares, here also defined as spherical quality degree  
 SEM Scanning electron microscope  
 SPR Surface plasmon resonance  
 TEM Transmission electron microscopy  
 $V$  Reactor volume  
 $v_F$  Fermi velocity  
 $\Gamma_\infty$  Bulk metal value frequency  
 $\Gamma(r)$  Size dependent relaxation frequency  
 $\epsilon_m$  Real refraction index of the surrounding medium  
 $\epsilon_{\text{Re}}$  and  $\epsilon_{\text{Im}}$  Metal-dielectric constants (real and imaginary)  
 $\lambda$  Evaluated wavelength  
 $\lambda_{\text{shift}}$  Peak location difference between the experimental and calculated data  
 $\tau$  Residence time  
 $\omega$  Bulk angular frequency  
 $\omega_p$  Resonance frequency

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

The authors greatly acknowledge the financial support from UK Engineering and Physical Science and Research Council (grant number EP/L020443/2).

## References

- 1 S. Prabhu and E. K. Poulou, *Int. Nano Lett.*, 2012, 2, 32.
- 2 J. S. Kim, E. Kuk, K. N. Yu, J.-H. Kim, S. J. Park, H. J. Lee, S. H. Kim, Y. K. Park, Y. H. Park, C.-Y. Hwang, Y.-K. Kim, Y.-S. Lee, D. H. Jeong and M.-H. Cho, *Nanomed.: Nanotechnol., Biol. Med.*, 2007, 3, 95–101.

- 3 T. Neuberger, B. Schöpf, H. Hofmann, M. Hofmann and B. von Rechenberg, *J. Magn. Magn. Mater.*, 2005, **293**, 483–496.
- 4 L. Zhang, F. Gu, J. Chan, A. Wang, R. Langer and O. Farokhzad, *Clin. Pharmacol. Ther.*, 2008, **83**, 761–769.
- 5 D. V. Talapin and E. V. Shevchenko, *Chem. Rev.*, 2016, **116**, 10343–10345.
- 6 R. Prasad, A. Bhattacharyya and Q. D. Nguyen, *Front. Microbiol.*, 2017, **8**, 1014.
- 7 O. V. Salata, *J. Nanobiotechnol.*, 2004, **2**, 3.
- 8 D. M. Kendziora, I. Ahmed and L. Fruk, *RSC Adv.*, 2014, **4**, 17980–17985.
- 9 W.-T. Liu, *J. Biosci. Bioeng.*, 2006, **102**, 1–7.
- 10 A. Moissala, A. G. Nasibulin and E. I. Kauppinen, *J. Phys.: Condens. Matter*, 2003, **15**, S3011–S3035.
- 11 K. An and G. A. Somorjai, *ChemCatChem*, 2012, **4**, 1512–1524.
- 12 C. Dong, C. Lian, S. Hu, Z. Deng, J. Gong, M. Li, H. Liu, M. Xing and J. Zhang, *Nat. Commun.*, 2018, **9**, 1252.
- 13 A. Zaleska-Medynska, M. Marchelek, M. Diak and E. Grabowska, *Adv. Colloid Interface Sci.*, 2016, **229**, 80–107.
- 14 X. Wang, Y. Zhu, A. Vasileff, Y. Jiao, S. Chen, L. Song, B. Zheng, Y. Zheng and S.-Z. Qiao, *ACS Energy Lett.*, 2018, **3**, 1198–1204.
- 15 G. Sharma, A. Kumar, S. Sharma, M. Naushad, R. Prakash Dwivedi, Z. A. Allothman and G. T. Mola, *J. King Saud Univ., Sci.*, 2019, **31**, 257–269.
- 16 I. J. Bruce and T. Sen, *Langmuir*, 2005, **21**, 7029–7035.
- 17 A. K. Gupta and M. Gupta, *Biomaterials*, 2005, **26**, 3995–4021.
- 18 X. S. Shen, G. Z. Wang, X. Hong and W. Zhu, *Phys. Chem. Chem. Phys.*, 2009, **11**, 7450–7454.
- 19 K.-J. Wu, G. M. De Varine Bohan and L. Torrente-Murciano, *React. Chem. Eng.*, 2017, **2**, 116–128.
- 20 P. K. Jain, X. Huang, I. H. El-Sayed and M. A. El-Sayed, *Plasmonics*, 2007, **2**, 107–118.
- 21 S. Chen, R. Yuan, Y. Chai and F. Hu, *Microchim. Acta*, 2013, **180**, 15–32.
- 22 S. Chandran, V. Ravichandran, S. Chandran, J. Chemmunda and B. Chandarshekar, *J. Appl. Res. Technol.*, 2016, **14**, 319–324.
- 23 M. Chastellain, A. Petri and H. Hofmann, *J. Colloid Interface Sci.*, 2004, **278**, 353–360.
- 24 M. Grzelczak, J. Pérez-Juste, P. Mulvaney and L. M. Liz-Marzán, *Chem. Soc. Rev.*, 2008, **37**, 1783–1791.
- 25 N. R. Jana, L. Gearheart and C. J. Murphy, *Adv. Mater.*, 2001, **13**, 1389–1393.
- 26 S. Panigrahi, S. Kundu, S. Ghosh, S. Nath and T. Pal, *J. Nanopart. Res.*, 2004, **6**, 411–414.
- 27 S. Iravani, *Green Chem.*, 2011, **13**, 2638–2650.
- 28 E. Shahbazali, V. Hessel, T. Noël and Q. Wang, *Nanotechnol. Rev.*, 2014, **3**, 65.
- 29 V. Sebastian Cabeza, S. Kuhn, A. A. Kulkarni and K. F. Jensen, *Langmuir*, 2012, **28**, 7007–7013.
- 30 S. Marre, Y. Roig and C. Aymonier, *J. Supercrit. Fluids*, 2012, **66**, 251–264.
- 31 M. S. Bakshi, *Cryst. Growth Des.*, 2016, **16**, 1104–1133.
- 32 D. Li, C. Wang, D. Tripkovic, S. Sun, N. M. Markovic and V. R. Stamenkovic, *ACS Catal.*, 2012, **2**, 1358–1362.
- 33 J. Xiao and L. Qi, *Nanoscale*, 2011, **3**, 1383–1396.
- 34 J. Boleininger, A. Kurz, V. Reuss and C. Sönnichsen, *Phys. Chem. Chem. Phys.*, 2006, **8**, 3824–3827.
- 35 N. G. Bastús, F. Merkoçi, J. Piella and V. Puentes, *Chem. Mater.*, 2014, **26**, 2836–2846.
- 36 Y. Yu, Q. Zhang, X. Lu and J. Y. Lee, *J. Phys. Chem. C*, 2010, **114**, 11119–11126.
- 37 D. Cambié, C. Bottecchia, N. J. W. Straathof, V. Hessel and T. Noël, *Chem. Rev.*, 2016, **116**, 10276–10341.
- 38 K. L. Kelly, E. Coronado, L. L. Zhao and G. C. Schatz, *J. Phys. Chem. B*, 2003, **107**, 668–677.
- 39 K. Schweikert, *J. Bank. Financ.*, 2018, **88**, 44–51.
- 40 S. Agnihotri, S. Mukherji and S. Mukherji, *RSC Adv.*, 2014, **4**, 3974–3983.
- 41 U. Aslam, S. Chavez and S. Linic, *Nat. Nanotechnol.*, 2017, **12**, 1000.
- 42 J. J. Baumberg, J. Aizpurua, M. H. Mikkelsen and D. R. Smith, *Nat. Mater.*, 2019, **18**, 668–678.
- 43 X. Dong, X. Ji, H. Wu, L. Zhao, J. Li and W. Yang, *J. Phys. Chem. C*, 2009, **113**, 6573–6576.
- 44 K. Ranoszek-Soliwoda, E. Tomaszewska, E. Socha, P. Krzyczmonik, A. Ignaczak, P. Orłowski, M. Krzyzowska, G. Celichowski and J. Grobelny, *J. Nanopart. Res.*, 2017, **19**, 273.
- 45 J. B. Deshpande and A. A. Kulkarni, *Chem. Eng. Technol.*, 2018, **41**, 157–167.
- 46 K.-J. Wu and L. Torrente-Murciano, *React. Chem. Eng.*, 2018, **3**, 267–276.
- 47 V. Amendola and M. Meneghetti, *J. Phys. Chem. C*, 2009, **113**, 4277–4285.
- 48 D. K. Bhui, H. Bar, P. Sarkar, G. P. Sahoo, S. P. De and A. Misra, *J. Mol. Liq.*, 2009, **145**, 33–37.
- 49 M. van Dijk, A. Tchebotareva, M. Orrit, M. Lippitz, S. Berciaud, D. Lasne, L. Cognet and B. Lounis, *Phys. Chem. Chem. Phys.*, 2006, **8**, 3486–3495.
- 50 L. J. M. Herrera, D. M. Arboleda, D. C. Schinca and L. B. Scaffardi, *J. Appl. Phys.*, 2014, **116**, 233105.
- 51 K. Kolwas and A. Derkachova, *J. Quant. Spectrosc. Radiat. Transfer*, 2013, **114**, 45–55.
- 52 P. K. Jain, K. S. Lee, I. H. El-Sayed and M. A. El-Sayed, *J. Phys. Chem. B*, 2006, **110**, 7238–7248.
- 53 A. Derkachova, K. Kolwas and I. Demchenko, *Plasmonics*, 2016, **11**, 941–951.
- 54 P. B. Johnson and R. W. Christy, *Phys. Rev. B: Solid State*, 1972, **6**, 4370–4379.
- 55 M. Daimon and A. Masumura, *Appl. Opt.*, 2007, **46**, 3811–3820.
- 56 J. C. Lagarias, J. A. Reeds, M. H. Wright and P. E. Wright, *SIAM J. Optim.*, 1998, **9**, 112–147.
- 57 E. Stefan Kooij and B. Poelsema, *Phys. Chem. Chem. Phys.*, 2006, **8**, 3349–3357.
- 58 L. Wang, X. Sun and J. Xing, *J. Mod. Opt.*, 2012, **59**, 1829–1840.
- 59 K. J. Wu, Y. H. Gao and L. Torrente-Murciano, *Faraday Discuss.*, 2018, **208**, 427–441.

- 60 T. M. Rahman and V. E. Rebrov, *Processes*, 2014, **2**, 466–493.
- 61 R. Zong, X. Wang, S. Shi and Y. Zhu, *Phys. Chem. Chem. Phys.*, 2014, **16**, 4236–4241.
- 62 X. C. Jiang, C. Y. Chen, W. M. Chen and A. B. Yu, *Langmuir*, 2010, **26**, 4400–4408.
- 63 N. T. K. Thanh, N. Maclean and S. Mahiddine, *Chem. Rev.*, 2014, **114**, 7610–7630.
- 64 V. K. LaMer and R. H. Dinegar, *J. Am. Chem. Soc.*, 1950, **72**, 4847–4854.
- 65 E. C. Vreeland, J. Watt, G. B. Schober, B. G. Hance, M. J. Austin, A. D. Price, B. D. Fellows, T. C. Monson, N. S. Hudak, L. Maldonado-Camargo, A. C. Bohorquez, C. Rinaldi and D. L. Huber, *Chem. Mater.*, 2015, **27**, 6059–6066.
- 66 I. M. Lifshitz and V. V. Slyozov, *J. Phys. Chem. Solids*, 1961, **19**, 35–50.
- 67 M. Darroudi, M. B. Ahmad, A. H. Abdullah, N. A. Ibrahim and K. Shameli, *Int. J. Mol. Sci.*, 2010, **11**, 3898–3905.
- 68 M. K. Alqadi, O. A. Abo Noqtah, F. Y. Alzoubi, J. Alzoubi and K. Aljarrah, *Mater. Sci.-Pol.*, 2014, **32**, 107–111.
- 69 J. Liu, C. Dai and Y. Hu, *Environ. Res.*, 2018, **161**, 49–60.
- 70 Y. Gao and L. Torrente-Murciano, *Nanoscale*, 2019 (under review).
- 71 Z. Chen, J. W. Chang, C. Balasanthiran, S. T. Milner and R. M. Rioux, *J. Am. Chem. Soc.*, 2019, **141**, 4328–4337.
- 72 X. Dong, X. Ji, J. Jing, M. Li, J. Li and W. Yang, *J. Phys. Chem. C*, 2010, **114**, 2070–2074.
- 73 J. H. Xu, J. Tan, S. W. Li and G. S. Luo, *Chem. Eng. J.*, 2008, **141**, 242–249.
- 74 Z. Koczorowski, *J. Electroanal. Chem. Interfacial Electrochem.*, 1981, **125**, 225–229.