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# Maximized axial helicity in a Pd<sub>2</sub>L<sub>4</sub> cage: inverse guest size-dependent compression and mesocate isomerism†‡

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Helicity is an archetypal structural motif of many biological systems and provides a basis for molecular recognition in DNA. Whilst artificial supramolecular hosts are often helical, the relationship between helicity and guest encapsulation is not well understood. We report a detailed study on a significantly coiled-up Pd<sub>2</sub>L<sub>4</sub> metallohelicite with an unusually wide azimuthal angle (~176°). Through a combination of NMR spectroscopy, single-crystal X-ray diffraction, trapped ion mobility mass spectrometry and isothermal titration calorimetry we show that the coiled-up cage exhibits extremely tight anion binding (*K* of up to 10<sup>6</sup> M<sup>-1</sup>) by virtue of a pronounced oblate/prolate cavity expansion, whereby the Pd–Pd separation decreases for mono-anionic guests of increasing size. Electronic structure calculations point toward strong dispersion forces contributing to these host–guest interactions. In the absence of a suitable guest, the helical cage exists in equilibrium with a well-defined mesocate isomer that possesses a distinct cavity environment afforded by a doubled Pd–Pd separation distance.

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## Introduction

The prevalence of structural helicity in biological systems exemplifies its fundamental importance for complex functions. For example, the helical structure of DNA provides a chiral surface for recognition and facilitates tight packing of genetic information.<sup>1</sup> Coiled motifs in peptide chains, such as  $\alpha$ -helices, experience strong structural reinforcement through cooperative non-covalent interactions, giving rise to a stiff secondary structure. Clearly, helicity is important for defining the tertiary structure and biological function of proteins and other biomacromolecules.<sup>2</sup>

These findings have inspired chemists to prepare artificial assemblies with structural helicity as the blueprint, such as

helical molecules,<sup>3,4</sup> foldamers<sup>5,6</sup> and polymer aggregates.<sup>7</sup> Many of these helical assemblies find strong applications in various areas, including asymmetric catalysis<sup>8</sup> and biomedicine.<sup>9</sup> For example, dinuclear metallohelicites composed of two or three twisted ligand strands have been shown to inhibit malignant cell growth by binding the major groove of DNA.<sup>10–12</sup>

The family of quadruple-stranded metallohelicites intersects with the compound class of M<sub>2</sub>L<sub>4</sub> coordination cages, the former defined by a helical twist, the latter by an accessible cavity.<sup>13,14</sup> The degree of axial helicity for M<sub>2</sub>L<sub>4</sub> lantern-shaped cages can be defined by the azimuthal angle,  $\alpha$ , which is the angular projection of the bridging ligand respective to the metal–metal axis (Fig. 1a).<sup>15</sup> Many non-helical M<sub>2</sub>L<sub>4</sub> cages (Fig. 1b) have shown promising functions for molecular separation,<sup>16,17</sup> anion sequestration,<sup>18,19</sup> catalysis<sup>20,21</sup> and drug delivery.<sup>22–24</sup> Despite helicity being a functional motif in nature, efforts to augment the host–guest chemistry of M<sub>2</sub>L<sub>4</sub> cages have mainly focused on increasing their structural complexity.<sup>25</sup> The incorporation of functional ligands (chiral, photo-switchable, redox-active, etc.) into such cage architectures offers the prospect of multi-functionalisation through heteroleptic, heterometallic or other low-symmetry assembly approaches.<sup>25–34</sup> Moreover, structural motifs such as skewed geometry and double-cage interpenetration have been shown to give rise to phenomena such as shape recognition, allosteric binding and stimuli responsiveness.<sup>35–38</sup>

The relationship between axial helicity and the guest binding properties of M<sub>2</sub>L<sub>4</sub> coordination cages is somewhat ill-defined.<sup>19</sup>

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† Dedicated to Professor Peter Steel on the occasion of his 68<sup>th</sup> birthday.

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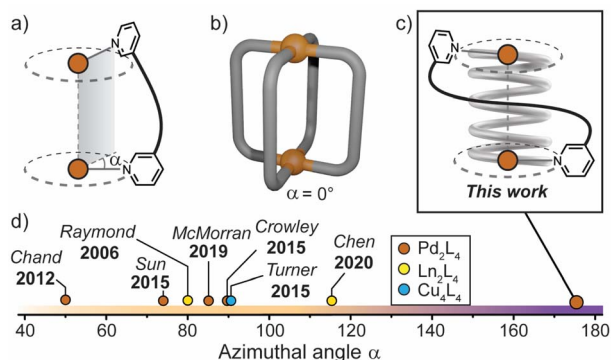


Fig. 1 (a) Azimuthal angle ( $\alpha$ ) for a  $M_2L_4$  helicate; (b) general structure of a non-helical  $M_2L_4$  cage, where  $\alpha = 0^\circ$ ; (c) this work: a tightly-wrapped  $Pd_2L_4$  helicate that behaves as a supramolecular spring; (d) azimuthal angles (extracted from X-ray diffraction data) of selected quadruple-stranded metallohelicates, shown for comparison.<sup>15,39,42–46</sup>

McMorran and Steel demonstrated that a helical  $Pd_2L_4$  cage assembled from flexible 4-bis(3-pyridyloxy)benzene ligands undergoes cavity adaptation by increasing its Pd–Pd separation and helical pitch upon binding anionic guests of increasing size.<sup>13,39</sup> Jung and Lee demonstrated that a non-helical  $Pd_2L_4$  cage assembled from alkyl-derived ligands behaves in a similar fashion, whereby an increase in the Pd–Pd separation is observed when the cage binds anions of increasing size.<sup>40</sup> It is therefore unclear whether axial helicity or ligand flexibility is primarily responsible for mediating cavity expansion in  $M_2L_4$  coordination cages.

In these and related  $M_2L_4$  cages, the equilibrium constant for anion binding in polar solvents is rather low ( $<10^3 M^{-1}$ ) leading to the necessity of utilising non-polar solvents to achieve stronger host–guest interactions.<sup>41</sup> Moreover, the azimuthal angle ( $\alpha$ ) of so far reported helical  $M_2L_4$  cages ( $M = Pd^{II}$ ,  $Pt^{II}$  or lanthanides) has been largely limited to  $\leq 90^\circ$  (Fig. 1d); increasing  $\alpha$  beyond this limit represents uncharted territory for  $M_2L_4$  cages. We postulated that maximising the axial helicity of  $M_2L_4$  cages will lead to enhanced guest binding properties (compared to cages where  $\alpha \leq 90^\circ$ ) as the tightly-wrapped ligands should provide a large internal surface area and dis-favour guest unbinding.

Herein, we report a  $Pd_2L_4$  helical cage (**1**) with azimuthal angles ranging between  $171$ – $176^\circ$  (Fig. 1c and d), representing the most axially-twisted  $M_2L_4$  coordination cage reported so far. We closely examine the cage's unique structural changes associated with the binding of mono-anions through extensive X-ray crystallographic, NMR, MS and ITC analyses, as well as electronic structure calculations. Together, these experiments provide intimate structural details on the pronounced host–guest adaptability of **1**, facilitated by its significantly twisted, spring-like architecture. In addition, we probe the anion-dependent equilibrium between the coiled-up metallohelicate and its elongated mesocate isomer, which is a unique feature of this system. Together, these findings provide new detailed insights into the relationship between axial helicity, structural

adaptability and guest-binding phenomena of artificial host assemblies.

## Results and discussion

### Ligand design

To prepare a coordination cage with an azimuthal angle of  $>90^\circ$ , we designed a bis-monodentate ligand with a carbazole backbone and two inward-pointing 8-substituted isoquinoline donors (**L**).<sup>35</sup> This pincer-shaped ligand has a theoretical ligand bite angle of  $15^\circ$  (Fig. 2a) as defined by the intersection of a vertical vector along the ligand backbone and one running parallel to the N-donor.<sup>47</sup> Given the rigidity of **L** and the necessity of adopting a bite angle of  $\sim 90^\circ$  for lantern-shaped  $M_2L_4$  coordination cages (particularly with square-planar  $Pd^{II}$ ), we anticipated that the isoquinoline donors of **L** would need to undergo severe twisting to bridge two  $Pd^{II}$  ions. **L** was synthesized through a Sonogashira cross-coupling between 3,6-dibromo-9-hexyl-9H-carbazole and 8-ethynylisoquinoline in 53% yield (ESI†).

### Synthesis and solution-state analysis of **1**·X

Heating a 2 : 1 mixture of **L** with  $[Pd(CH_3CN)_4](BF_4)_2$  at  $70^\circ C$  in  $CD_3CN$  resulted in the assembly of the coordination cage  $[BF_4@Pd_2L_4](BF_4)_3$  (herein denoted as **1**· $BF_4$ ) as revealed by  $^1H$  NMR spectroscopy and electrospray ionisation mass spectrometry (ESI-MS). In the  $^1H$  NMR spectrum, the proton

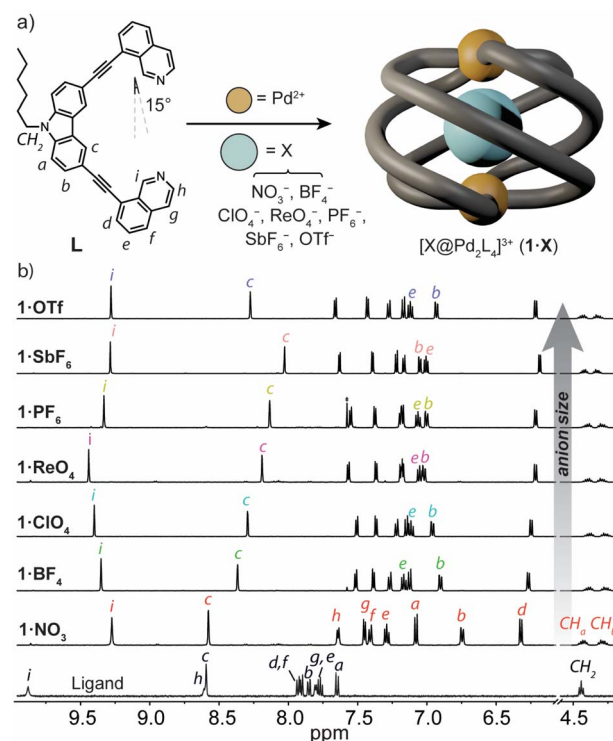


Fig. 2 (a) A scheme showing the structure and native bite angle of **L** and its self-assembly with  $Pd^{II}$  to give the  $Pd_2L_4$  helicate **1**·X, where X = different anionic guests; (b)  $^1H$  NMR spectra (500 MHz/ $CD_3CN$ ) of ligand **L** (bottom) and **1**·X (top).

resonances of the cage are significantly upfield-shifted relative to **L** (Fig. 2b), consistent with the presence of  $\pi$ -stacking interactions between neighbouring ligands in the cage assembly. The characteristic splitting of the diastereotopic *N*-hexyl CH<sub>2</sub> resonances (CH<sub>a</sub>, CH<sub>b</sub>, 4.18–4.48 ppm at 298 K) provides direct evidence for axial helicity, and the *P* and *M* enantiomers of the Pd<sub>2</sub>L<sub>4</sub> cage are in slow exchange relative to the <sup>1</sup>H NMR time scale. The encapsulation of a BF<sub>4</sub><sup>−</sup> anion within **1** was confirmed by <sup>19</sup>F NMR spectroscopy, which revealed two resonances at −151.8 ppm and −149.4 ppm (Fig. S10†). These are assigned to the free and tightly encapsulated BF<sub>4</sub><sup>−</sup> anion, respectively, and <sup>19</sup>F exchange spectroscopy (EXSY) confirmed their exchange at 25 °C (Fig. S15†).

Next, we examined the self-assembly of **1** in the presence of anionic guests of different size and shape. We prepared a series of cage analogues (**1**·**X**), either from the respective Pd<sup>II</sup> salt (X = NO<sub>3</sub><sup>−</sup>, ClO<sub>4</sub><sup>−</sup>, PF<sub>6</sub><sup>−</sup>, SbF<sub>6</sub><sup>−</sup> and OTf<sup>−</sup>), or by anion exchange (X = ReO<sub>4</sub><sup>−</sup>). ESI-MS confirmed the Pd<sub>2</sub>L<sub>4</sub> assembly for each counterion by prominent signals corresponding to [Pd<sub>2</sub>L<sub>4</sub> + *n*X]<sup>(4−*n*)<sup>+</sup> cations; *n* = 1 and 2 (ESI†). When comparing their <sup>1</sup>H NMR spectra, the hydrogens that point directly into the cage cavity appear at distinct chemical shifts (see resonances *c* and *i*, Fig. 2b). As for **1**·BF<sub>4</sub>, also for **1**·PF<sub>6</sub> and **1**·OTf, <sup>19</sup>F NMR spectroscopy confirmed anion encapsulation, indicating that the cage cavity can accommodate anions of greatly different size and shape. In contrast to **1**·BF<sub>4</sub>, <sup>19</sup>F EXSY revealed that neither the larger PF<sub>6</sub><sup>−</sup> nor OTf<sup>−</sup> encapsulated anion exchanges with solvated anions under the conditions of the NMR measurement (Fig. S27 and S34†). In addition to the resonances that are directly sensitive to the bound guest (*c* and *i*), those corresponding to hydrogens that point outside of the cage cavity (*e* and *b*) also undergo notable changes in chemical shift across the host–guest complexes of **1**·**X** (Fig. 2b). Based on the data we discuss below, we interpret this as an anion-dependent ‘uncoiling/coiling-up’ behaviour.<sup>39</sup> This is further supported by the tendency of non-encapsulated anions to influence the chemical shift of proton resonances directly adjacent to the metal centre (in this case, *h* and *g*).<sup>19,48</sup></sup>

In order to gain further insight into the solution-state structure of **1**·**X**, we examined the host–guest complexes by <sup>1</sup>H-<sup>1</sup>H NOESY. These analyses revealed several intramolecular contacts diagnostic of axial helicity, including a strong interligand correlation between protons *f* and *a* (e.g., Fig. S20†). Cage complexes of larger anions (e.g., SbF<sub>6</sub><sup>−</sup>) display an additional NOE contact between protons *i* and *f* of neighbouring ligands (Fig. S39†). The former contact (*f*·*a*), in particular, confirms that the isoquinoline moiety of **L** is closely stacked on top of a neighbouring carbazole ligand backbone – a conformation that could only be afforded by a significantly twisted helical architecture.

The helical structure of **1**·**X** (X = NO<sub>3</sub><sup>−</sup>, ClO<sub>4</sub><sup>−</sup>, ReO<sub>4</sub><sup>−</sup>, PF<sub>6</sub><sup>−</sup>, SbF<sub>6</sub><sup>−</sup> and OTf<sup>−</sup>) was also examined by Ion-Mobility Mass Spectrometry (IM-MS). The experimental collisional cross section (CCS) values for the [X@Pd<sub>2</sub>L<sub>4</sub>]<sup>3+</sup> species were found to range between 560 ± 10 (**1**·BF<sub>4</sub>) and 563 ± 17 Å<sup>2</sup> (**1**·OTf) which correlates well to the calculated CCS obtained from their X-ray structures (Table S5.1† and *vide infra*). These CCS values

indicate that the cage host maintains its compact helical structure in the solution and gas phase, even when encapsulating anions of notably different size and shape.

### Solid-state analysis of **1**·**X**

With MeCN solutions of **1**·**X** in hand, we grew X-ray quality crystals of each sample (with the exception of **1**·NO<sub>3</sub>) by slow-vapour diffusion with diisopropyl ether as the antisolvent. **1**·BF<sub>4</sub>, **1**·ClO<sub>4</sub> and **1**·ReO<sub>4</sub> (anions with *T<sub>d</sub>* symmetry) crystallized in the tetragonal space group *P4<sub>2</sub>bc*, whilst **1**·PF<sub>6</sub>, **1**·SbF<sub>6</sub> and **1**·OTf (*O<sub>h</sub>* and *C<sub>3v</sub>* anion symmetry, respectively) crystallized in the monoclinic space group *P2<sub>1</sub>/n*. For each structure, the asymmetric unit was occupied by one full molecule of the Pd<sub>2</sub>L<sub>4</sub> cage with its respective counter ions and a variable number of solvent molecules. All forms of **1** are stabilised by strong  $\pi$ -stacking interactions between neighbouring ligands and these account for the upfield chemical shift of ligand resonances in the <sup>1</sup>H NMR spectra of **1**·**X** (Fig. 2b). Further inspection of the structures revealed that the azimuthal angle of the cage increases for larger anion guests (e.g., 171.6° and 176.2° for **1**·BF<sub>4</sub> and **1**·SbF<sub>6</sub>, respectively – Table ESI 4.3†). This represents the largest axial twist for a M<sub>2</sub>L<sub>4</sub> quadruple-stranded metallohelicite reported to date.

For all solid-state structures of **1**·**X**, one disordered anion was located inside the cage cavity (Fig. 3c), and additional anions sit in close proximity to the Pd<sup>II</sup> centres. The encapsulated anions are stabilised by multiple weak hydrogen bonds (C–H···X = 2.2–2.5 Å) arising from the carbazole hydrogens that point directly inside the cavity of **1**. For **1**·PF<sub>6</sub> and **1**·SbF<sub>6</sub>, the bound anion is tilted by *ca.* 50° with respect to the longest anion axis and Pd(isoquinoline)<sub>4</sub> plane, which contrasts with the vertical orientation of PF<sub>6</sub><sup>−</sup> in other Pd<sub>2</sub>L<sub>4</sub> cages<sup>38</sup> as well as the bound OTf<sup>−</sup> in **1**·OTf. The tilted orientation of PF<sub>6</sub><sup>−</sup> and SbF<sub>6</sub><sup>−</sup> presumably maximises host–guest interactions by orienting the fluoride substituents toward the surrounding hydrogen atoms of the host (also the case for the vertically oriented OTf<sup>−</sup> anion).

M<sub>2</sub>L<sub>4</sub> cages typically accommodate guests of increasing size by increasing the separation between the metal centres.<sup>39,40,49</sup> In analysing the solid-state structures of **1**·**X**, we found the opposite trend for the Pd···Pd separation (*d*<sub>1</sub>) with respect to the size and shape of the encapsulated anion. Fig. 3a and b depict how **1** adapts to anions of increasing size. Counter-intuitively, the cage compresses along the Pd–Pd axis for larger anions; *d*<sub>1</sub> (Pd···Pd distance) = 9.29 Å for **1**·BF<sub>4</sub> and 8.99 Å for **1**·ReO<sub>4</sub>. A further decrease of *d*<sub>1</sub> is observed for **1**·PF<sub>6</sub> (*d*<sub>1</sub> = 8.71 Å) followed by a small increase for **1**·SbF<sub>6</sub> and a much larger increase for the non-globular OTf<sup>−</sup> guest (**1**·OTf *d*<sub>1</sub> = 9.40 Å). Simultaneously, *d*<sub>2</sub> (average distance between the opposing carbazole N atoms) increases incrementally from **1**·BF<sub>4</sub> (*d*<sub>2</sub> = 14.56 Å) to **1**·SbF<sub>6</sub> (*d*<sub>2</sub> = 15.12 Fig. 3b). It is worth noting that despite the similar molecular volumes of SbF<sub>6</sub><sup>−</sup> and OTf<sup>−</sup> (84.9 and 85.4 Å<sup>3</sup> respectively; RB3LYP/6-31g(d), LANL2DZ) **1** closely adapts to their different shape. This is reflected by large differences in *d*<sub>1</sub> and *d*<sub>2</sub> when comparing **1**·OTf to **1**·SbF<sub>6</sub> (+0.59 Å and −0.27 Å, for *d*<sub>1</sub> and *d*<sub>2</sub>, respectively; Fig. 3b). Combined, this shows that the wide azimuthal angle (or strong helical pitch) of **1** affords



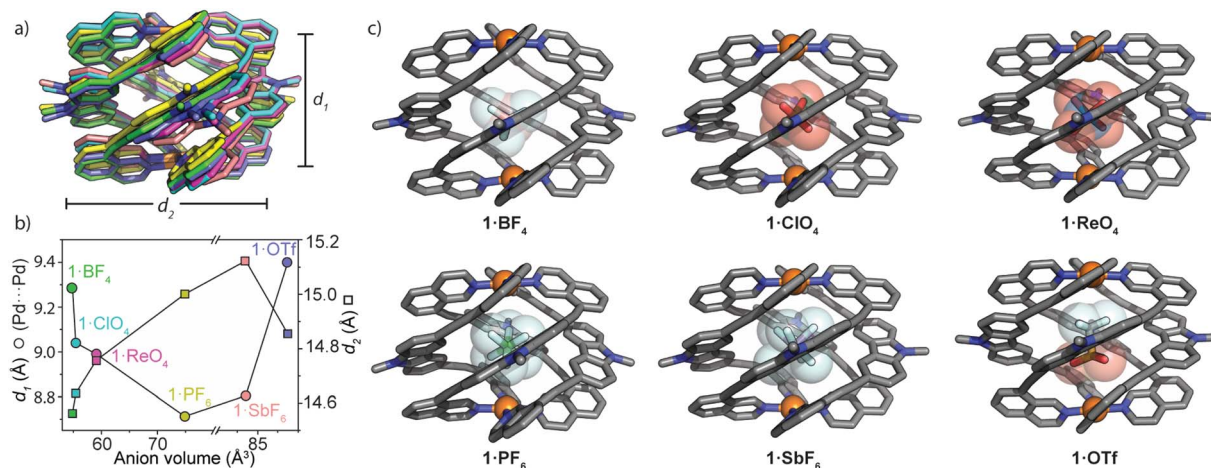


Fig. 3 Solid-state comparison of  $1 \cdot X$ ; (a) an overlay of the six X-ray structures of  $1$  (the structures are centred on Pd1, and counter-ions and hexyl chains are removed for clarity).  $d_1$  = Pd···Pd separation,  $d_2$  = average distance between carbazole-N atoms.; (b) a plot of the solid-state structural parameters obtained from single-crystal X-ray analysis. Anion volumes were calculated by DFT geometry optimizations (RB3LYP; B, O, F, Cl, P, S: 6-31g(d); Re, Sb: LANL2DZ); (c) X-ray crystal structures of  $1 \cdot X$ . Hexyl chains, hydrogen atoms and unbound anions are omitted for clarity.

size and shape adaptability in guest binding, where the cavity trends from a prolate to an oblate shape; as  $d_1$  decreases,  $d_2$  increases. This behaviour enables  $1$  to alter its cavity volume by up to 18%; for example, the cavity volumes of  $1 \cdot \text{BF}_4$  and  $1 \cdot \text{SbF}_6$  are  $119 \text{ \AA}^3$  and  $140 \text{ \AA}^3$  respectively, as determined by VOIDOO<sup>50</sup> calculations. These adaptations enable host-guest complexes of  $1$  to maintain a favourable packing coefficient of  $0.55 \pm 0.09$ .<sup>51</sup>

### Anion exchange experiments

The versatile anion encapsulation properties of  $1$  motivated us to investigate anion exchange experiments in  $\text{CD}_3\text{CN}$  solution. A fresh sample of  $1 \cdot \text{NO}_3$  was titrated with tetrabutylammonium salts of  $\text{BF}_4^-$ ,  $\text{ClO}_4^-$ ,  $\text{ReO}_4^-$ ,  $\text{PF}_6^-$ ,  $\text{SbF}_6^-$  and  $\text{OTf}^-$ . In the  $^1\text{H}$  NMR spectrum of these samples, slow exchange relative to the NMR timescale was observed in each case. Initial calculations revealed that for  $\text{BF}_4^-$ ,  $K \approx 10^3 \text{ M}^{-1}$  whilst for the remaining anions,  $K$  was estimated to be  $>10^5 \text{ M}^{-1}$ , thus approaching the accuracy limit of  $K$  determination by  $^1\text{H}$  NMR spectroscopy.<sup>§</sup>

To examine anion exchange more accurately, we turned to Isothermal Titration Calorimetry (ITC). Titrating  $1 \cdot \text{NO}_3$  with tetrabutyl ammonium salts of  $\text{BF}_4^-$ ,  $\text{ClO}_4^-$ , and  $\text{ReO}_4^-$  yielded titration curves consistent with the expected 1:1 host-guest binding model (Fig. S52†). The equilibrium constants for anion exchange followed an anion size-dependent trend:  $K = 1.9 \times 10^3 \text{ M}^{-1}$  ( $\text{BF}_4^-$ ),  $9.2 \times 10^4 \text{ M}^{-1}$  ( $\text{ClO}_4^-$ ) and  $1.1 \times 10^6 \text{ M}^{-1}$  ( $\text{ReO}_4^-$ ) (Fig. 4a and b and Table S3.1†). The thermodynamic data revealed that anion exchange of  $1 \cdot \text{NO}_3$  with  $\text{BF}_4^-$  is entropically driven, whilst exchange with  $\text{ClO}_4^-$  and  $\text{ReO}_4^-$  is increasingly enthalpically driven (Fig. 4b). Given this, the encapsulated  $\text{BF}_4^-$  guest may be bound with more degrees of freedom compared to larger  $\text{ClO}_4^-$  and  $\text{ReO}_4^-$ . The high enthalpic contributions observed in the exchange of  $\text{ClO}_4^-$  and  $\text{ReO}_4^-$  with  $\text{NO}_3^-$  in  $1$  point towards a more optimized fit, which presumably maximizes the electrostatic and dispersive interactions in these host-guest complexes.

Monitoring anion exchange of  $1 \cdot \text{NO}_3$  with  $\text{PF}_6^-$ ,  $\text{SbF}_6^-$  and  $\text{OTf}^-$  by ITC was uninformative due to the slow kinetics for the bimolecular exchange processes of these particular anions (e.g., complete exchange with  $\text{SbF}_6^-$  requires 1.5 h). Time-resolved  $^1\text{H}$  NMR array experiments revealed the rates of anion exchange to be  $\text{OTf}^- > \text{PF}_6^- > \text{SbF}_6^-$ ;  $3.0 \times 10^{-3} \text{ s}^{-1}$ ,  $2.3 \times 10^{-3} \text{ s}^{-1}$ ,  $1.2 \times 10^{-3} \text{ s}^{-1}$  respectively (Fig. S53†). In contrast, ITC analysis indicated that anion exchange of  $1 \cdot \text{NO}_3$  with  $\text{BF}_4^-$ ,  $\text{ClO}_4^-$  or  $\text{ReO}_4^-$  is complete within seconds. It is worth noting that the  $K$  values for the exchange of nitrate in  $1$  with polyatomic anions are considerably high ( $\geq 10^5 \text{ M}^{-1}$  for all

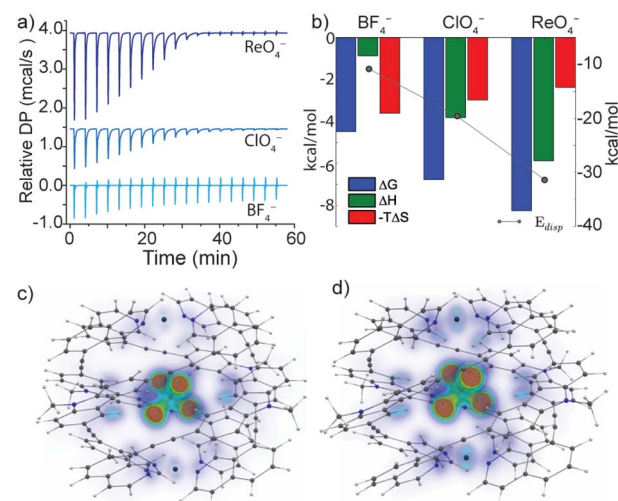


Fig. 4 (a) Relative heat flow produced upon titrating a MeCN solution of  $1 \cdot \text{NO}_3$  with aliquots of tetrabutyl ammonium  $\text{BF}_4^-$ ,  $\text{ClO}_4^-$  or  $\text{ReO}_4^-$  (as determined by ITC analysis); (b) graphical representation of the thermodynamic parameters for guest exchange with a theoretical  $E_{disp}$  obtained from electronic structure calculations; calculated dispersion interaction density (VOXEL-DID) of (c)  $1 \cdot \text{ClO}_4$  gas-phase structure (d)  $1 \cdot \text{ClO}_4$  solid-state structure.

anions except  $\text{BF}_4^-$ ) compared to other  $\text{M}_2\text{L}_4$  cages and related cage structures reported in the literature.<sup>19,52,53</sup> For  $\text{M}_2\text{L}_4$  cages ( $\text{M} = \text{Pd}^{\text{II}}$  or  $\text{Pt}^{\text{II}}$ ), equilibrium exchange constants for mono-anionic guests in polar solvent are usually in the order of  $<10^3 \text{ M}^{-1}$ ,<sup>39,46</sup> and non-polar solvents, large non-coordinating counter-ions or endohedral binding sites can be employed to further maximize host-guest interactions.<sup>41,54,55</sup> Here, the exceptionally strong anion encapsulation properties of **1** can be credited to the large internal surface area of the tightly-wrapped cavity, which is afforded by the unusually wide azimuthal angles ( $171$ – $176^\circ$ ).

To investigate the role of London dispersion interactions in the guest binding of **1**, we performed electronic structure calculations on the host-guest complexes examined by ITC (*i.e.*,  $\mathbf{1} \cdot \text{BF}_4^-$ ,  $\mathbf{1} \cdot \text{ClO}_4^-$  and  $\mathbf{1} \cdot \text{ReO}_4^-$ ; all as  $3+$  cations with hexyl chains truncated to methyl substituents). All host-guest systems were optimized using the BP-86<sup>56,57</sup> functional with a def2-SVP<sup>58–61</sup> basis and the D3-correction[D3] including three-body-terms.<sup>62,63</sup> The optimised gas-phase structures were subjected to DF-SCS-LMP2/cc-pVTZ electronic structure calculations and the results were compared to the energies obtained from the corresponding solid-state geometries at the same level of theory.

Indeed, a strong correlation between the measured guest binding enthalpies and calculated dispersion contributions<sup>64–66</sup> was found for both solid-state and gas-phase geometries, where **1** maintains a large contact surface to the bound guest (Fig. 4b–d, S71 and S72†). Additional calculations (DF-SCS-LMP2/VTZ level of theory) revealed that compression of the Pd⋯Pd axis of **1** is energetically demanding in the absence of a bound guest (Fig. S73†), which leads to the conclusion that dispersion interactions have a strong impact on the guest binding and corresponding structural adaptations of **1**. The trend of increasing dispersion contributions for larger anions is consistent between the solid-state and gas-phase-optimised geometries (Table S5.3†), suggesting that the influence of crystal packing on the solid-state structure of **1** is minimal. This is supported by the fact that  $\mathbf{1} \cdot \text{BF}_4^-$ ,  $\mathbf{1} \cdot \text{ClO}_4^-$  and  $\mathbf{1} \cdot \text{ReO}_4^-$  crystallise in the same tetragonal space group with almost identical unit cell dimensions ( $a, b = 30 \text{ \AA}$ ,  $c = 29 \text{ \AA}$ , Table S4.2†).

### Helicate/mesocate isomerism of **1**

Helicate/mesocate isomerism is prevalent for double and triple-stranded metallohelicates, and control over their selective formation and isolation has been the subject of intense study.<sup>67–69</sup> This in part is motivated by their respective properties; for example, triple-stranded di-ruthenium helicates exhibit selectivity toward certain cancer cells as compared to their mesocate counterparts.<sup>12</sup> For these types of metallohelicates, ligand design (*e.g.*, odd/even spacer length) or steric hindrance can lead to the formation of one isomer over the other.<sup>70,71</sup> In many cases, both helicate and mesocate isomers can form from the same ligand and metal ion and exist either as separable static structures<sup>72</sup> or in dynamic equilibrium;<sup>67,73</sup> the latter is usually the case for helicates based on non-chelated

coordination environments with high conformational flexibility (concerning the ligand–metal bonds). However, factors that drive the selective formation of helicates over mesocates remain ill-defined, especially for  $\text{M}_2\text{L}_4$  quadruple-stranded metallohelicates. This may relate to their acute azimuthal angles ( $\alpha < 90^\circ$ ), where only subtle structural distortions are required to invert helicity.<sup>74</sup> This means that nearly all so far postulated  $\text{M}_2\text{L}_4$  mesocates are transient species, rather than isolable ones.<sup>75,76</sup>

In our attempts to grow single crystals of the helical cage  $\mathbf{1} \cdot \text{NO}_3^-$ , we were surprised to obtain crystals of a markedly different  $[\text{Pd}_2\text{L}_4]^{4+}$  cage complex (here denoted as **1'**). X-ray analysis revealed that **1'** crystallizes in the orthorhombic space group *Pccn* with half of the  $\text{Pd}_2\text{L}_4$  cage in the asymmetric unit. In the structure of **1'**, the isoquinoline donors of **L** adopt a *syn* orientation and twist away from the carbazole backbone (giving individual ligands a  $C_s$  symmetry) – a conformation which gives rise to a mesocate (in contrast to the  $C_2$ -symmetric ligand conformation in the helicates). In this elongated isomer, the Pd⋯Pd separation is  $16.8 \text{ \AA}$  – which is almost double that of the most contracted form of helical  $\mathbf{1} \cdot \text{X}$ . In contrast to the helicate, neighbouring carbazole backbones in mesocate **1'** participate in mutual edge-to-face  $\pi$ -stacking (closest C–H⋯ $\pi$  contact between two ligand backbones =  $3.0 \text{ \AA}$ ) and form the faces of

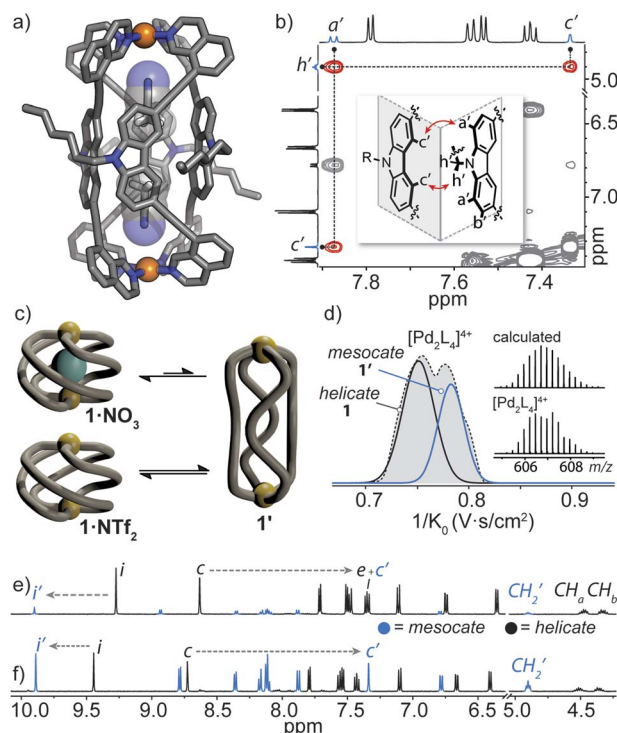


Fig. 5 (a) X-ray structure of mesocate **1'** showing the MeCN solvent guests; (b)  $^1\text{H}$ – $^1\text{H}$  NOESY spectrum of  $\mathbf{1} \cdot \text{NTf}_2$  ( $\text{R} = \text{C}_6\text{H}_{13}$ ), showing the important contacts in the inset; (c) a schematic of the anion-dependant equilibrium between the helicate **1** and mesocate **1'**; (d) Trapped Ion Mobility (TIMS) mobilogram alongside the measured and calculated isotope pattern of the  $4+$  species of  $\mathbf{1} \cdot \text{NTf}_2$ ; (e)  $^1\text{H}$  NMR spectra (500 MHz,  $\text{CD}_3\text{CN}$ ) of  $\mathbf{1} \cdot \text{NO}_3^-$  after 2 weeks at  $25^\circ\text{C}$  (signal  $c'$  arising from **1'** overlaps with signals of **1**); (f)  $\mathbf{1} \cdot \text{NTf}_2$  after two weeks at  $25^\circ\text{C}$ .



the tubular, almost cuboid Pd<sub>2</sub>L<sub>4</sub> structure. Instead of encapsulating a NO<sub>3</sub><sup>−</sup> counterion, the C<sub>4h</sub> symmetric mesocate hosts two MeCN solvent molecules (Fig. 5a) that orient their nitrogen atoms toward the opposing Pd<sup>II</sup> centres. This type of electrostatic interaction has been previously observed for other Pd<sup>II</sup>-based cages<sup>77</sup> and in the case of 1', the MeCN guests are further stabilised by CH-π interactions (C-H<sub>MeCN</sub>⋯C<sub>carbazole</sub> = 2.66 Å) inside the aromatic-rich cavity. Although we were not able to obtain an X-ray structure of the expected helical form of 1·NO<sub>3</sub>, the chemical shifts of inward-pointing protons *i* and *c* (Fig. 2a and 5e and f) are consistent with a helically-twisted cage encapsulating a NO<sub>3</sub><sup>−</sup> guest.

To understand the helicate/mesocate isomerism of 1, we further examined the self-assembly process between L and Pd(NO<sub>3</sub>)<sub>2</sub> in solution. As expected, the <sup>1</sup>H NMR spectrum obtained minutes after combining L with Pd(NO<sub>3</sub>)<sub>2</sub> clearly reflects the helical cage species, owing to the characteristic splitting pattern of the diastereotopic N-CH<sub>2</sub> protons (CH<sub>a</sub>/CH<sub>b</sub>). However, when the CD<sub>3</sub>CN solution of 1·NO<sub>3</sub> was allowed to equilibrate over two weeks at room temperature, another set of signals evolved which were assignable to a new Pd<sup>II</sup> complex of L (Fig. 5e, ratio = 1:0.15 in favour of 1·NO<sub>3</sub>). The chemical shifts of the newly evolved aromatic resonances appeared downfield compared to 1·NO<sub>3</sub>, with exception of *c* (denoted as *c'*), which appeared significantly upfield (−1.27 ppm). The significant shielding of proton *c'* is in agreement with the X-ray structure of mesocate 1' as this particular hydrogen points directly toward the π-system of an adjacent carbazole backbone in the mesocate (Fig. 5a and b).

The equilibrium between 1 and 1' was further investigated by synthesizing 1·NTf<sub>2</sub> (Fig. S42†). We chose NTf<sub>2</sub><sup>−</sup> as a counterion because its volumetric size (157 Å<sup>3</sup>) is larger than the maximal cavity size of 1 (≤140 Å<sup>3</sup>). Thus, an 'empty' form of 1 was prepared by combining L with [Pd(CH<sub>3</sub>CN)<sub>4</sub>](NTf<sub>2</sub>)<sub>2</sub> at room temperature. Indeed, the absence of encapsulated NTf<sub>2</sub><sup>−</sup> was confirmed by <sup>19</sup>F NMR spectroscopy, which showed one isolated signal corresponding to free NTf<sub>2</sub><sup>−</sup> present in solution (Fig. S43†). The <sup>1</sup>H NMR spectrum of 1·NTf<sub>2</sub> after 5 minutes of mixing the ligand and Pd<sup>II</sup> salt revealed two sets of signals in a 1:0.24 ratio, with the helical, empty form of the cage predominating. After 2 weeks at room temperature, the sample equilibrated to a 1:1 ratio (Fig. 5f) of 1 and 1'. Importantly, <sup>1</sup>H-<sup>1</sup>H NOESY unequivocally confirmed that the second set of signals corresponds to mesocate 1'. Two characteristic NOE contacts were observed: *a'*⋯*c'* and *c'*⋯*h'* (Fig. 5b and S46†). These contacts originate from the orthogonal relationship of neighbouring carbazole backbones (an arrangement only present in 1') and are not observed in the NOESY spectrum of 1 or L. The difference in the equilibrium ratios of 1 and 1' for 1·NO<sub>3</sub> and 1·NTf<sub>2</sub> reveals that the bound NO<sub>3</sub><sup>−</sup> anion stabilizes the helicate 1 but is not necessarily required to template it (Fig. 5c). Mesocate 1' was not observed to encapsulate anions, however, upon addition of a suitable guest (e.g., PF<sub>6</sub><sup>−</sup>), the equilibrium shifts in favour of the metallohelicate 1 (Fig. S47†).

Next, Trapped Ion Mobility (TIMS) ESI-MS analysis was carried out on an equilibrated sample of 1·NTf<sub>2</sub> (containing both helicate 1 and mesocate 1'), in order to gain insight into

the gas-phase behaviour of these isomeric structures. Based on previous studies examining structural variations in heteroleptic palladium cages,<sup>78</sup> we anticipated that the drift times of the two isomers of 1 should be resolvable. Indeed, we observed two distinct mobilities corresponding to 4+ ions at 606.7 *m/z* (Fig. 5d and S69†). Derivation of the experimental collisional cross sections (eCCS) for these mobilities revealed that the two 4+ species have an effective area of 605 Å<sup>2</sup> and 631 Å<sup>2</sup>, which correspond well to the relative difference between the theoretically determined tCCS values of 1 and 1' (568.6 Å<sup>2</sup> and 578.9 Å<sup>2</sup>, respectively). Thus, the TIMS data confirms the helicate/mesocate pairs (*i.e.*, isomers 1 and 1') in the gas phase, suggesting that the MeCN solvent does not play a significant role in templating the formation of 1'. To investigate this hypothesis further, we examined the formation of 1·BF<sub>4</sub> in DMSO by <sup>1</sup>H NMR spectroscopy. In contrast to the exclusive formation of the helicate in MeCN, the <sup>1</sup>H NMR spectrum of 1·BF<sub>4</sub> in DMSO revealed the presence of both the helicate and mesocate in a ~2:1 ratio, respectively. Thus, both anionic guests and solvent influence the equilibrium distribution of the helicate/mesocate pair.

## Conclusions

In summary, we have presented a comprehensive analysis of the solution state, solid state, and gaseous state structure and host-guest chemistry of the most axially twisted Pd<sub>2</sub>L<sub>4</sub> cage reported to date. Through careful ligand design, the helical twist of this Pd<sub>2</sub>L<sub>4</sub> cage (azimuthal angle of up to ~176°) results in exceptionally tight anion binding which is mediated by prolate/oblate cavity size and shape adaptability. Notably, a comparison of six X-ray crystal structures of 1·X (X = BF<sub>4</sub><sup>−</sup>, ClO<sub>4</sub><sup>−</sup>, ReO<sub>4</sub><sup>−</sup>, PF<sub>6</sub><sup>−</sup>, SbF<sub>6</sub><sup>−</sup> and OTf<sup>−</sup>) provided detailed insights on the compression/decompression of the Pd-Pd distance (*d*<sub>1</sub>), whereby *d*<sub>1</sub> decreases when larger anions are bound within the cage cavity (with the exception of the non-globular OTf<sup>−</sup> guest). In the binding of anions of increasing size, we also observed a subtle increase in the azimuthal angle of the cage host. <sup>1</sup>H NMR and ITC analysis revealed the benefit of large axial helicity for 1 by *K* values as high as 10<sup>6</sup> M<sup>−1</sup> (ReO<sub>4</sub><sup>−</sup>) for the bimolecular exchange of NO<sub>3</sub><sup>−</sup> with larger anions. This can be attributed to the tightly-wrapped helical cage, which provides a large contact surface that facilitates strong dispersion forces toward the bound anion. In the absence of a suitable guest, 2D NMR, X-ray diffraction and TIM-MS revealed that 1 exists in a well-defined helicate/mesocate equilibrium, and convergence to the helicate is observed in the presence of an appropriate guest.

Although factors such as cavity volume,<sup>79,80</sup> ligand-guest interactions,<sup>55</sup> host asymmetry<sup>35</sup> and metal-metal distance<sup>19</sup> all play a role in the host-guest chemistry of M<sub>2</sub>L<sub>4</sub> coordination cages, our work exemplifies how maximising their axial helicity correlates with an exceptionally tight and adaptable binding pocket with a strong affinity toward a range of anionic guests. We anticipate this work will inspire further research into guest recognition/separation by tightly-wrapped helical cage compounds.



## Data availability

The synthetic procedures, characterization, spectral data, and additional computational and crystallographic details are available in the ESI.† CCDC 1851098–1851104 contain the supplementary crystallographic data for this paper. Additional data is available from the authors upon request.

## Author contributions

W. M. B. and G. H. C. conceived the idea. W. M. B. conducted most of the synthesis, characterization, and analysis, and wrote the manuscript. S. H. and C. D. acquired and analysed the IM-MS data. J. J. H. and S. H. solved and refined the X-ray structures. R. A. M. and A. W. performed electronic structure calculations. W. H. performed 2D NMR measurements. G. H. C. managed the project, aided in experimental design, and reviewed the manuscript. All authors provided comments and approved the final version of the manuscript.

## Conflicts of interest

There are no conflicts to declare.

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## Notes and references

§ Furthermore, attempts to obtain reliable values of  $K$  for  $1 \cdot \text{X}$  by  $^1\text{H}$  NMR spectroscopy were convoluted by the presence of proton resonances of mesocate  $1'$  which arose in samples of  $1 \cdot \text{NO}_3$  and  $1 \cdot \text{BF}_4$  after long periods of equilibration.

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