



Table 1 Typical energy scales found in chemistry

Chemical event	Typical enthalpy value (kJ mol <sup>-1</sup> )
Reduction of lithium ions (Li <sup>+</sup> + e <sup>-</sup> → Li)	299
Hydration of lithium ions (Li <sup>+</sup> + 4H <sub>2</sub> O → Li <sup>+</sup> ·4H <sub>2</sub> O)	-508
Hydration of ammonium ions	-293
Hydration of magnesium ions	-1908
Solvation of lithium ions with propylene carbonate	-50
Hydrogen bridges	12–50
van der Waals forces	< 25
C–C bond	350
C–N bond	295

illustrative example, the oxidation of lithium metal to lithium cations corresponds to an enthalpy of  $-299 \text{ kJ mol}^{-1}$ , while the standard electrode potential of this redox process is  $-3.1 \text{ V}$  versus SHE.

To the surprise of most chemists, hydration and solvation enthalpies are rather high, and, in most cases, fully reversible. The hydration of lithium ions with 4 water molecules for instance in fact produces more energy than the oxidation of bare Li atoms, that is, one could hypothetically build a lithium “dehydration” device, with a voltage higher than that of a lithium battery. The numbers are of course smaller for coordination with organic solvents and also decrease upon an increase in the size of the ions, but in principle still lie in a range that is relevant for energy storage, and the scientific literature has already identified it as being relevant. For example, it has been observed that there is an anomalous increase in the energy storage capacity of electric double-layer capacitors if the size of the pores of the carbon (the material that is usually used as an electrode material) approaches the size of the electrolyte ions, thus resulting in desolvation. It is not surprising that the energy stored at each electrode (measured independently from one another) is about  $150 \text{ kJ mol}^{-1}$ , which is close to the solvation energy of the applied electrolyte (in this case tetraethylammonium tetrafluoroborate in acetonitrile). The energy stored by the cations is interestingly slightly lower.<sup>22</sup> The latter can be assigned to their lower solvation enthalpy caused by a larger diameter. Secondary bonding schemes, such as hydrogen bridges, are in the order of  $10\text{--}50 \text{ kJ mol}^{-1}$  each (and thus correspond to  $0.1\text{--}0.5 \text{ V}$  on the voltage scale) but can also contribute in a cooperative fashion to energy storage. The problem with all of these storage mechanisms is that the possible voltage to utilize them should be within the electrochemical stability range of the solvent applied, and this range is especially small for water.

Given the square dependence of the energy density on the applied voltage windows for EDLCs, the energy storage capacity can be enhanced by implementing electrolytes with higher electrochemical windows and stabilities. Ionic liquids (ILs), also known as low temperature molten salts (with melting points below  $100 \text{ }^\circ\text{C}$ ), are composed solely of ions and thus are promising alternatives to established solvent-based systems due to their wide operation voltage (as high as  $4 \text{ V}$ ), negligible volatility, non-flammability, as well as a wide potential operating temperature range.<sup>24–26</sup> In terms of the structures of ILs, the ion coordination numbers have been modelled, and indeed, as

exemplified for ethylammonium nitrate, a variety of coordination structures with varying coordination numbers of 2–5 can be imagined and are in an accessible energy range.<sup>27,28</sup> This means that the ion coordination structure can indeed play a role in charge storage, with uncompensated partial charges near the electrodes being taken up by electrons and holes with the appropriate voltage related to the energy of the structure change.

Pioneering experimental and theoretical studies have been devoted to exploring the underlying physics for IL charge storage under equilibrium conditions in carbon nanopores, and molecular dynamics simulations<sup>29</sup> and nuclear magnetic resonance (NMR) spectroscopic measurements<sup>30</sup> have provided information on the charge storage at the molecular level. The dependence of the capacitance on the pore size for IL electrolytes has been investigated in detail, especially the anomalous increase in the capacitance when the pore size decreases to the ion size.<sup>24,31,32</sup> Obviously, single cations or anions can penetrate ion-sized micropores, while the countercharge in the electrode and wall-ion interactions replace the interactions of the ions in the IL bulk. Pulling an entity from its own bulk environment into an otherwise foreign matrix is always accompanied with a loss of energy, and this energy corresponds to the voltage needed to drive this process. Using molecular dynamics simulations, the capacitance was also found to depend on the applied voltage, where the capacitance reaches a maximum when all of the co-ions are expelled from the pores.<sup>29</sup> Very recently, scattering experiments have been used to resolve the structure of IL molecules confined into nanoporous carbons.<sup>33</sup> The results indeed showed that the Coulombic ordering reduced when the pores could accommodate only a single layer of ions. This non-Coulombic ordering was even further enhanced in the presence of an external electric potential and surely represents the first example of a non-traditional energy term that contributes to charge storage in IL-based supercapacitors.

Up to now, the mechanisms for charge storage revealed in these studies have been based on processes occurring in the narrowest, “single ion” carbon micropores and have neglected possible conformation variations and phase changes in the bulk of the rather large IL ions during charging/discharging. The simulation studies have only employed coarse-grained, rigid models of IL ions, and the experimental studies have not taken into account the influence of structural changes, although they are very likely to occur as a function of the electric potential. Some simulation studies have revealed that ion-wall interactions can significantly affect the molecular conformation of IL ions.<sup>34</sup> Such structural (ordering) transitions are known to be potential-driven and are not limited to very narrow micropores.<sup>35</sup> In mesopores, the contribution of such transitions to the energy storage is significantly higher due to the greater number of involved solvent molecules and the related changes in the conformation and coordination energy.<sup>36</sup>

In a recent paper, we analyzed the capacitive charge storage in a variety of porous carbon nanomaterials and their hybrids using 1-ethyl-3-methylimidazolium tetrafluoroborate (EMImBF<sub>4</sub>) as a model ionic liquid between  $0$  and  $3.5 \text{ V}$ .<sup>21</sup> Fig. 1 presents a selected case involving a nanoporous carbon where the



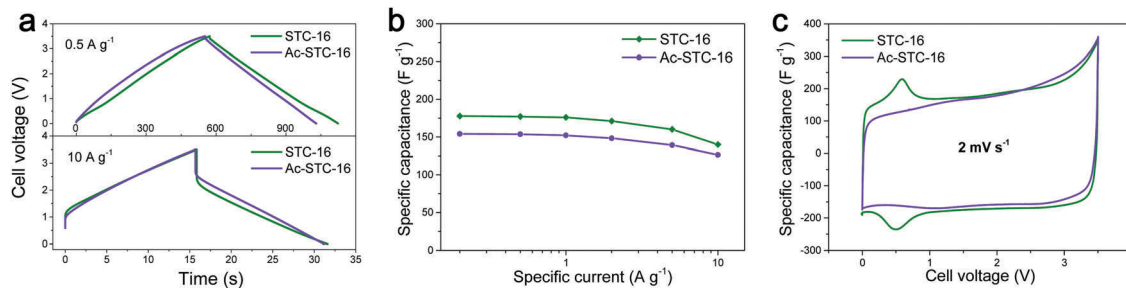


Fig. 1 EDLC performance comparison of STC-16 and Ac-STC-16 tested in EMImBF<sub>4</sub> ionic liquid using a two-electrode configuration: (a) the galvanostatic charge/discharge profiles at different specific currents, (b) the capacitance retention with specific current increase, and (c) the cyclic voltammograms at scan rates of 2 mV s<sup>-1</sup>. The difference is due to appropriately sized nanopores which, when enlarged by etching, lose their ability to tightly insert ions. Figures reproduced with permission from ref. 21.

micropore size matches the size of the IL ions (denoted as STC-16). Galvanostatic charging/discharging curves with potential limitation are close to the symmetric triangle at a low specific current of 0.5 A g<sup>-1</sup>, typically characteristic of double-layer capacitor behavior (Fig. 1a and b). However, all of the cyclic voltammetry (CV) curves exhibited small peaks at a cell voltage of ~0.6 V, which in addition are fully reversible. An even more pronounced peak starts to emerge at higher cell voltages (Fig. 1c). The interesting dimension of these experiments is that we also CO<sub>2</sub>-etched this specific carbon sample (denoted as Ac-STC-16) to create a similar structure, but with an even greater specific surface area (2833 m<sup>2</sup> g<sup>-1</sup> instead of 2324 m<sup>2</sup> g<sup>-1</sup>) and enlarged micropores. Obviously, both the specific storage peak at 0.6 V and some of the charge storage capacity (178 F g<sup>-1</sup> resp. 7.5 μF cm<sup>-2</sup> for STC and 154 F g<sup>-1</sup> resp. 5.4 μF cm<sup>-2</sup> for the etched sample) vanish, that is, increasing the surface area while otherwise preserving most other structural parameters counter-intuitively reduces the energy storage performance. This was a model experiment to show that double-layer compression only partially describes the energy storage, and that the energy lost is one of the special ion conformations in the appropriately sized pores.

The origin of the peaks in the CVs of the STCs at a cell voltage of ~0.6 V is in a strict sense still unproven, but from the rate dependence measurements, we would like to rule out a Faradic origin. In addition, we only see such peaks at this potential in materials where the micropore sizes are close to the size of BF<sub>4</sub><sup>-</sup> ions. On an enthalpy scale, 0.6 V translates into 60 kJ mol<sup>-1</sup>, which is in the range of strong inter-molecular forces and Coulomb isolation. In our opinion, the origin of the peak is likely due to the insertion of single ions into the micropores and the reversible formation and release of inter-molecular forces. For example, the entry of single cations into the wall has been predicted by numerical simulations. Here, a single ion loses all its neighbouring Coulombs, while some of the energy is returned through van der Waals interactions with the carbon and of course localized electrons in the wall to compensate for the ion charge.

In addition, a rise at a very high voltage was observed in all of the CVs of the mesopore-containing carbons, and, according to our experiments, this is not related to material degradation. For the purposes of this present discussion paper, we ran

additional CVs well into the region of metastability, and the energy uptake was found to peak significantly and even turned out to be partially reversible (Fig. 2a). This peak was observed for carbon materials with different pore sizes as well as for different ILs (Fig. 2b). The peak-like CV shape was found to be present only if the ethyl group in the cation of the EMImBF<sub>4</sub> was changed to a butyl group and also if the anion was changed from BF<sub>4</sub><sup>-</sup> to TFSI<sup>-</sup>. After the addition of adiponitrile solvent to the IL electrolyte, the peak disappeared, and the CV exhibited a rather typical EDLC-like appearance (Fig. 2c).

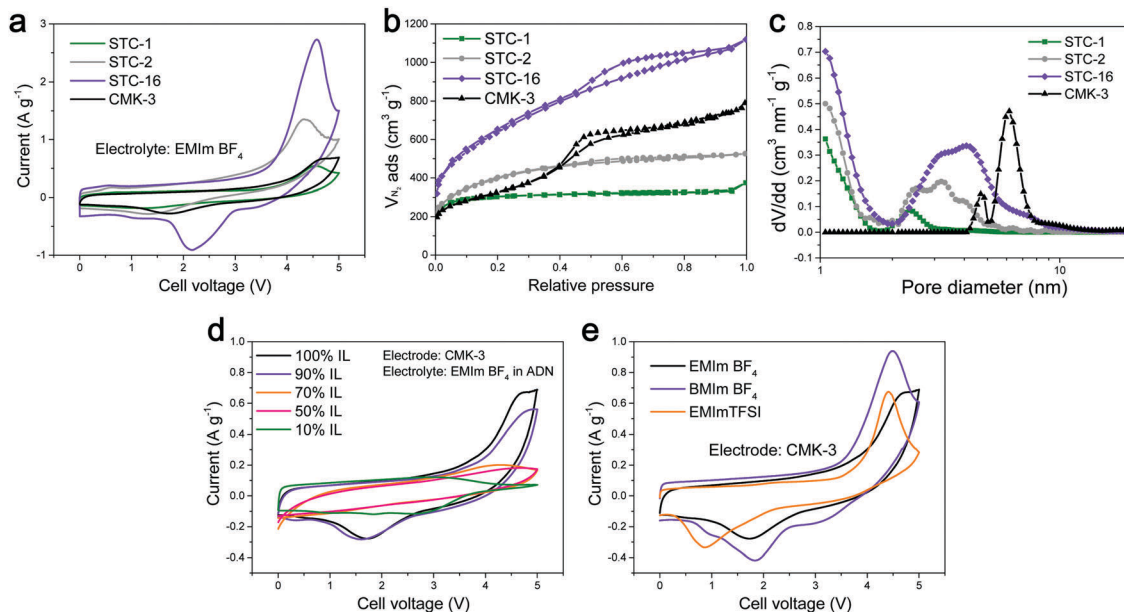
This voltage peak translates to a much higher transition enthalpy of above 350–450 kJ mol<sup>-1</sup>, which is higher than that of covalent bonds, but in the range of the melting of salts and other collective properties and structural transitions involving serious changes in the ion coordination numbers (Fig. 3).

Due to the onset of material decomposition, electrochemical measurements can only be ran at 5 V for analytical purposes with the present IL. Indeed, below the peak at high voltage, the constant increase in the current may indicate that irreversible electrolysis processes occur in parallel to the ordering transitions, and the peak intensity weakened upon measuring a greater number of cycles. However, in the standard, stable range we were able to tap into the footing of this process, *i.e.* we observed higher capacities due to this process and expected to observe even further significant increases in the storage capacities, for appropriate solvents systems and optimized solvent-wall interactions.

The amount of specific energy stored in the present systems was found to be one order of magnitude higher than that of commercial EDLCs, without any loss in the excellent power capability.<sup>21</sup>

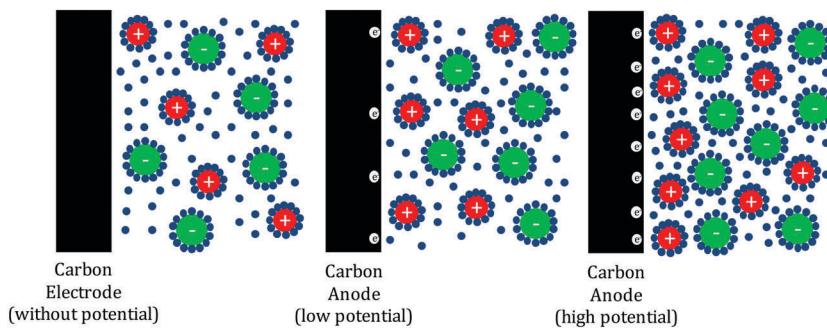
The sole contribution of double-layer formation and compression can – by the sheer amount of stored energy – be safely excluded. In our previous paper,<sup>21</sup> we translated the stored charge per surface area into the pro forma area per molecule at the electrochemical interface, and up to 2 ions (and the same number of electrons in the carbon phase) per nm<sup>2</sup> were calculated, which can be approximated to a dense layer of cations on the carbon surface. In contrast to the aqueous and organic electrolytes, the overall density of cations and anions in ionic liquids is expected to remain close to the bulk value and the energy cannot be stored in ion compression by removal of



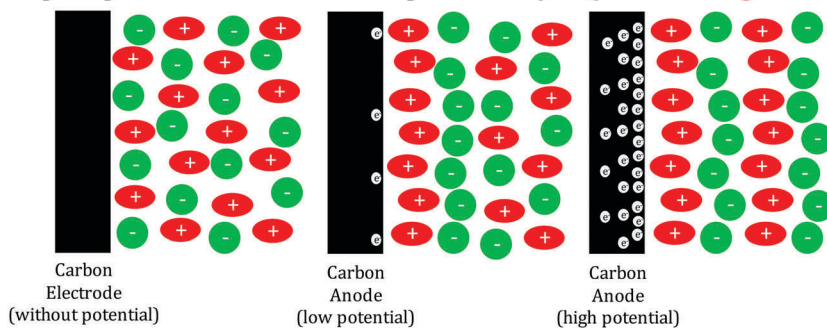


**Fig. 2** (a) Cyclic voltammograms (1st cycle) of microporous STC-1, micro- and mesoporous STC-2/STC-16, and ordered mesoporous CMK-3 carbon measured up to 5 V at  $2 \text{ mV s}^{-1}$ . (b) CV curves of CMK-3 in different electrolytes. (c) CV curves of CMK-3 in EMImBF<sub>4</sub> electrolyte at different concentrations in adiponitrile (ADN) solvent. It should be noted that the CV works on the absolute voltage scale, while the 2-electrode supercapacitors have floating potentials, *i.e.* there is a shift of  $\sim 0.5 \text{ V}$  when comparing the diagrams. (d) Nitrogen physisorption isotherms of the utilized carbon materials and (e) the corresponding pore size distributions indicating that the high voltage peak is related to processes enabled by the incorporation of mesopores.

### Supercapacitor anodes with solvent-based electrolytes (compression double-layer)



### Supercapacitor anodes with ionic liquid electrolytes (local ordering transitions)



**Fig. 3** Schematic of the (negative) electrode–electrolyte interface in EDLCs working with solvent-based (top) and IL (bottom) electrolytes. While energy storage in the former is based on double-layer compression and thus requires transport of electrolyte ions, local ordering transitions (*i.e.* changes in the ion coordination numbers) are responsible for the energy storage in the absence of a solvent.

non-charged solvent. On the contrary, ion rearrangement towards lower coordination numbers is an effect that can progress for nanometers into the bulk phase of the ionic liquid

and is therefore not controlled by the specific surface area but rather by pore volume, in agreement with the experimental findings (Fig. 3).



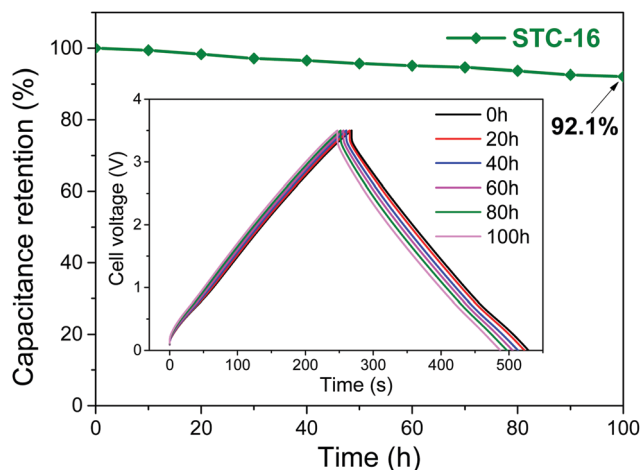


Fig. 4 Floating stability of STC-16 at a specific current of  $1 \text{ A g}^{-1}$  for 100 h of floating at 3.5 V. The inset shows the charge/discharge profiles every 20 h. Figures reproduced with permission from ref. 21.

This indeed opens up new possibilities for energy storage, as here the majority of energy is not stored in micropores (and cannot be optimized by the high specific surface area), but rather in appropriately designed mesopores, which support high energy Coulombic rearrangements of the IL in the bulk.

The involved processes are fully reversible, as indicated by the high stability of the devices and the nearly perfect symmetrical charge–discharge curves at low charge densities (Fig. 4), but are still limited by the electrochemical stability window of the IL electrolyte. Structural changes in liquids are in addition comparably fast (much faster than ion insertions or other battery processes), as they are not determined by transport and viscosity, but instead rely on local anion–cation exchange processes on the Angstrom scale. This causes capacitor-like rate performances, thus resulting in favorable power densities. This is mentioned in spite of the very high viscosities of ionic liquids, which are more restrictive for diffusion processes than for local structural rearrangements. Using these “new” effects, supercapacitors with all their strengths (cost, sustainability, rate, stability), even with their main weakness, the specific energy density, move into the range of batteries. The future is notoriously hard to predict, but we assume that going to smaller, higher coordinating ionic liquids, as well as optimizing wall–liquid interactions to promote the smoothness of structural changes, will allow the gap to be closed in current Li-ion battery technology.

## Conflicts of interest

There are no conflicts to declare.

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