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Copper-based semiconductor nanocrystals for optical applications

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Copper-based semiconductor nanocrystals (SNCs) are a broad class of nanomaterials that include pnictogenides, chalcogenides, and halides. The breadth of compositions, crystal structures, and electronic properties displayed by these SNCs is reflected in their rich photophysics, which can underpin strong photon absorption capabilities, efficient (photo)luminescence, and plasmonic behavior. Because of this optical versatility, copper-based SNCs have been proposed for applications in fields such as photovoltaics, sensing, bioimaging, and photocatalysis, in some cases reaching the level of commercialization. In this review, we provide an up-to-date overview of the research on these nanomaterials with optical properties of outmost technological relevance. After a brief introduction on light-(nano)matter interaction, we individually discuss pnictogenides, chalcogenides, and halides, pinpointing structure–property relationships, identifying the most common synthesis approaches, and highlighting cases of application we consider particularly relevant or novel. Lastly, we outline outstanding challenges, with the hope of spurring the ingenuity and curiosity of researchers towards the next discoveries in this multicoloured and multifaceted field.

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1. Introduction

Copper (Cu) is the 29th element of the periodic table and has key technological relevance. Classified as a transition metal, it gets its name from the island of Cyprus, where it was copiously mined in ancient times. Copper as a metal has played a key role in the development of early civilizations, so much so that an almost 2000-year long period of human history (from approximately mid-5th to early 3rd millennium BC) is referred to as “The Copper Age”.¹ During that prehistoric period, tools and weapons were mainly made of this rather malleable metal with a melting point of 1085 °C and capable of retaining sharp edges despite repeated use. Copper continued to play a central role also in the next era, “The Bronze Age”, since bronze is a Cu–Sn alloy. Even when iron took over the military and industrial scenes, copper remained a valuable metal used to produce copperware, jewelry, and mirrors. Moreover, it started

to become a valuable coinage metal. Fast forward to more recent history: in the 1820s, copper wires began being used in electrical wiring, *e.g.*, for telegraph lines and later for telephone lines.² Copper also features antimicrobial activity. Indeed, copper is the first solid antimicrobial material registered at the U.S. Environmental Protection Agency.³

But it is not only in its metallic form that copper has been used by humanity over the course of history: Ionic and non-ionic copper compounds have also found their place in human technology. CuSO₄ was already used already by ancient Egyptians, Greeks, and Romans for medicinal purposes, in inks, and as a mordant in dyeing. Records of these uses also remain also in the Middle Ages, and in the 19th century, it started to be used in electroplating, agriculture,⁴ and pest control. Cu-acetate has found similar application scenarios, while Cu(NO₃)₂ has been employed for textile dyeing, firework production, and in laboratory settings as an oxidizing agent and catalyst for organic reactions.⁵ CuCN is instead mainly used since the 18th century in electroplating, and CuCl₂ is a key catalyst in organic reactions.¹

While the story and use of metallic copper and copper compounds spans several centuries and millennia, it is only in the 20th century that Cu-based semiconductor materials began to be studied and applied. Indeed, Cu₂O is one of the earliest-known semiconductor materials,⁶ with the first patent on its use as current rectifier being filed in 1925. Cu₂S (chalcocite) and CuS (covellite) have been studied in the 1930s–1950s for their photoelectric properties. Since the 1970s, CuInSe₂ has

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been heavily investigated for photovoltaic applications, becoming a central material in so-called CIGS solar cells.

Moving to the nanoscale – even though Cu-doped CdS and ZnS colloidal nanocrystals have been reported since the mid-1980s⁷ – the history of Cu-based semiconductor nanocrystals (SNCs) starts in the late 1990s, early 2000s. These nanomaterials feature intriguing optical properties that somehow remained elusive until recently. To that end, a first example of photoluminescent CuInSe₂ nanocrystals was reported by Revaprasadu and co-workers,⁸ while the first report by the Burda group on the plasmonic properties of CuS nanocrystals was published in 2009.⁹ Since then, a plethora of studies on different copper-based SNCs have been published, reaching a tremendous level of understanding and control over the properties of these systems. The interest around Cu-based SNCs arises from the control that can be exerted on their optical pro-

perties through tuning of the composition, stoichiometry, structure, and copper's oxidation states.

Even though research on some of these nanomaterials peaked a few years ago, there is still a steady stream of publications on their fundamental properties and uses across a variety of fields, including biomedicine, energy conversion, catalysis, and lighting. Moreover, the recent push in the research on perovskite nanocrystals has led to the preparation of several Cu-based halide nanomaterials with unique properties.

The persistent interest in Cu-based SNCs is highlighted by excellent reviews published on the subject over the past few years. These reviews have varying degrees of comprehensiveness: some have a focus on chalcogenides,^{10–12} others on halides,^{13,14} while others even narrow the discussion down to a specific subclass (e.g., CuInS₂ QDs¹⁵). Some reviews describe

Table 1 Main semiconductors discussed in this Review, their respective crystal structure, bulk bandgap (when possible, the direct or indirect nature of the bandgap is reported), optical properties as SNCs, and main application fields

		Composition	Crystal structure	Bandgap (bulk) / eV	luminescence	plasmonic	Main applications	
Pnictogenides	Phosphides	Cu ₃ P	P6 ₃ cm	0-1.0			photonics, biomedicine photocatalysis	
	Nitrides	Cu ₃ N	Pm3m	0.8-1.9 (indir)				
Chalcogenides	Binary	Sulfides	Cu ₂ S	P6 ₃ /mmc (high chal.), P2 ₁ /c (low chal.)	1.1-2.0 (dir)	[shaded]	[shaded]	photovoltaics, photocatalysis, biomedicine <i>niche</i> : UC enhancement, SERS sensing
			Cu _{1.97} S	P2 ₁ /n (djurleite)				
			Cu _{1.81} S	P1 (roxbyite)				
			Cu _{1.86} S	R3m (digenite)				
			Cu _{1.75} S	Pnma (anilite)				
	CuS	P6 ₃ /mmc (covellite)						
	Selenides	Cu ₂ Se	Fm3m (berzelianite), P4 ₂ /n (bellidoite)	-2.0-2.3 (dir) and 1.1-1.5 (indir)	[shaded]	[shaded]		
		Cu _{1.5} Se	P4 ₂ /m (umangite)					
		CuSe	P6 ₃ /mmc (klockmannite)					
	Tellurides	Cu ₂ Te	P3m1 (weissite)	-1.1-1.5	[shaded]	[shaded]		
Cu ₃ Te		P4/nmm (rickardite)						
CuTe		Pmnm (vulcanite)						
Ternary	Cu-III-E	CuInS ₂ , CuInSe ₂ , CuInTe ₂	Mainly I42d (CP-type - CuInS ₂ : roquesite, CuGaS ₂ : gallite)	1.53 (dir), 1.07 (dir), 1.06 (dir)	[shaded]	[shaded]		
		CuGaS ₂ , CuGaSe ₂ , CuGaTe ₂	CuGaS ₂ : gallite	2.44, 1.68, 1.2-1.9				
		CuAlS ₂ , CuAlSe ₂	P6 ₃ /mc (WZ-type), F43m (ZB-type)	3.5, 2.7				
	Cu-IV-E	Cu ₂ SnS ₃ , Cu ₂ SnSe ₃	Cc (S: mohite, Se: okruginite), F43m (ZB-type), P6 ₃ /mc (WZ-type)	1.1-1.3 (dir), 0.4-0.9 (dir)				
		Cu ₂ SnS ₄	I42m (kuramite)	1.2-1.6 (dir)				
		Cu ₂ GeS ₃ , Cu ₂ GeSe ₃	Cc, P6 ₃ /mc (WZ-type), F43m (ZB-type)	1.5-1.6 (dir), 0.7-1.23 (dir)				
	Cu-V-E	Cu ₄ GeSe ₆	Pmn2 ₁ (argyrodite-type)	N/A				
		Cu ₂ BiS ₃ , Cu ₂ BiSe ₃	P2 ₁ 2 ₁ 2 (wittichenite), -	1.1-1.9, 1.1-1.9				
		Cu ₂ AsS ₄	I42m (luzonite)	1.3-1.4				
		Cu ₁₇ As ₄ S ₁₃	I43m (tennantite)	1.7				
		CuSbS ₂	Pnma (chalcostibite)	1.4-1.9 (dir)				
		Cu ₂ SbS ₄ , Cu ₂ SbSe ₄	I43m (famatinite), I42m (permingeaitite)	1.0-1.2				
	Cu-TM-E	Cu ₂ Sb ₂ S ₃	I43m (tetrahedrite)	1.8-2.0				
		Cu ₂ Sb ₃ S ₁₃	P2 ₁ /c (skinnerite)	1.4				
		CuFeS ₂ , CuFeSe ₂	I42d (chalcopyrite), P42c (eskebornite)	0.6, 0.16-0.60				
		Cu ₂ FeS ₄	Pbca (bornite)	0.8-1.3 [defect dependent]				
		Cu ₂ MoS ₄ , Cu ₂ CoS ₄ , Cu ₂ VSe ₄ , CuCrS ₂ , CuCrSe ₂ , CuCr ₂ Se ₄ , Cu ₂ MnS ₂	P42m, Fd3m, P43m, R3m, R3m, Fd3m, Fd3m					
Quaternary	Cu ₂ ZnSnS ₄	I4 (stannite), I43m (kesterite)	1.0-1.5 (dir), 0.8-1.4 (dir)					
[Cu ₁₀ CoSb ₂ S ₁₃ , Cu ₁₀ ZnSb ₂ S ₁₃ , Cu ₁₀ Ni ₂ Sb ₂ S ₁₃ , Cu ₂ ZnGeS ₄ , Cu ₂ CdGeS ₄ , Cu ₂ CoGeS ₄ , Cu ₂ ZnGeSe ₄]								
Halides	CsCuCl ₃ , CsCuBr ₃	P6 ₃ 22 or P6 ₃ 22 (1D), C222 ₁ (0D)	1.9-2.6 (indir), 1.9-2.6 (dir)					
	CsCu ₂ Cl ₃ , CsCu ₂ Br ₃ , CsCu ₂ I ₃	Cmcm (0D)	3.5-3.75 (dir)					
	Cs ₂ CuCl ₄ , Cs ₂ CuBr ₄	Pnma (0D), I4/mmm (2D)	2.6, 1.75					
	Cs ₂ Cu ₂ Cl ₃ , Cs ₂ Cu ₂ Br ₃ , Cs ₂ Cu ₂ I ₃	Pnma (0D) [only for Cl: Cmcm (1D)]	3.9-4.3					

DR = dielectric resonance, IB = interband absorption, LSPR = localized surface plasmon resonance, UC = upconversion, SERS = surface enhanced Raman spectroscopy, CP = chalcopyrite, WZ = wurtzite, ZB = zinc blende. The range of the spectral rectangles roughly indicates the variable extension of the bands, also considering size, morphology, and composition effects in reported SNCs. In binary chalcogenides, the darker black shading indicates to the position of the peak maximum shifting with varying material composition. Similarly, darker shading in the DR and LSPR rectangles for the Cu-Fe-E family highlights the peak maximum position.



selected application scenarios^{16,17} and others – while remaining cornerstone pieces – have been published more than five years ago.¹⁸ This review proposes a fresh and comprehensive look at a wide span of nanomaterial compositions, yet it is circumscribed to the analysis of the optical properties and the proposed applications based on them (Table 1). Herein, we highlight the structure–property relationships for each Cu-based nanomaterial, also identifying similarities between different nanomaterial families. To accomplish this goal, we first introduce basic concepts of light–matter interaction in nanomaterials (section 2). We then move to sections identified by the nature of the anion in the nanomaterial composition (Fig. 1): pnictogenides (section 3), chalcogenides (section 4),

and all-inorganic halides (section 5). Each section is further divided into subsections where selected families of Cu-based semiconductors are reviewed, touching upon their crystal structure, optical properties, synthesis methods, and applications. Whenever possible and relevant, parallels are drawn between different nanomaterial families. Lastly, we put the state-of-the-art of Cu-based nanomaterials in a broader context, highlighting future challenges and opportunities in this field (section 6).

2. Brief phenomenological background on light–matter interaction at the nanoscale

Before delving into Cu-based nanocrystals, it is useful to recall some basic concepts of light–matter interaction that guide the discussion on the optical properties of the nanomaterials presented in this Review (Fig. 2). Broadly speaking, the means of interaction of an optical nanomaterial with photons are twofold: absorption and scattering. The specifics of the underlying physical mechanisms vary depending on whether the nanomaterial displays plasmonic properties or not. Thus, we divide the discussion into plasmonic and non-plasmonic semiconductor nanomaterials. Moreover, we focus on the interaction with UV-Vis-NIR photons – *i.e.*, from approximately 250 to 2500 nm. We should also highlight that we do not aim for a rigorous mathematical description of the phenomena.

2.1. Non-plasmonic materials

2.1.1. Absorption. In non-plasmonic materials, absorption events occur when the energy of the impinging photon

The figure shows a periodic table with several elements highlighted. Copper (Cu) is highlighted in blue. Elements in pink (pnictogens) include N, P, As, Sb, Bi, and I. Elements in yellow (chalcogens) include S, Se, and Te. Elements in green (halogens) include F, Cl, Br, and I. Other elements are in white. Labels 'TM', 'halogens', 'chalcogens', and 'pnictogens' are placed above the periodic table with arrows pointing to the respective groups.

Fig. 1 Overview of the elements of interest in the context of copper-based SNCs. Cu is highlighted in blue, alongside the elements belonging to the pnictogen (pink), chalcogen (yellow), and halogen (green) groups that define the families of SNCs discussed in this review. The other elements in white background are found in multinary and/or alloyed copper-based SNCs.

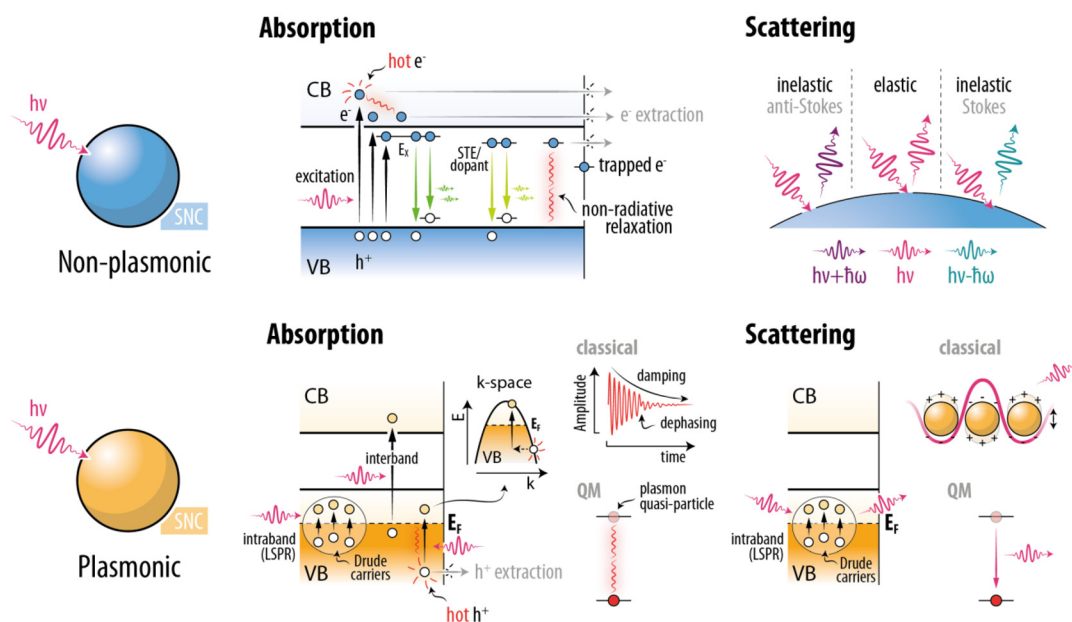


Fig. 2 Summary of the modalities of light–matter interaction in non-plasmonic and plasmonic SNCs.



matches the energy gap between two real electronic/energy states: One occupied and the other unoccupied. In SNCs, these states generally belong to the conduction and valence bands (CB and VB) or are introduced by dopants. Photoexcitation generally leads to the formation of free excitons (E_x): electron-hole pairs whose energy is stabilized *via* Coulomb interaction to a value (optical bandgap) lower than the energy difference between the CB bottom and the VB top (electronic bandgap). When the excited electron is promoted to an energy state above the lowest CB energy state, it is referred to as “hot electron”, since it has excess energy compared to thermal equilibrium. A hot electron can subsequently get trapped at a defect state related to vacancies, dangling bonds, or surface ligands. Alternatively, it can reach the bottom of the CB *via* thermalization, *i.e.*, transfer of momentum to the crystal lattice in the form of lattice vibrations (phonons) to reach the band edge.¹⁵ This process constitutes the generation of heat. From the bottom of the CB, the electron can occupy a free exciton state or be further stabilized in localized states due to, *e.g.*, electron-phonon interaction – self-trapped exciton (STE) – or presence of dopant-related states. From one of these states, the electron can radiatively recombine with a delocalized or localized (*i.e.*, occupying an energy state just above the VB) hole: A process referred to as photoluminescence. Note that even in the case of free-exciton luminescence, there is an energy difference between the absorption and emission photon (Stokes shift). The origin of this shift – whose explanation is beyond the scope of this review – include the presence of dark and bright exciton states (which are respectively inactive and active in optical transitions), energy transfer among different-sized nanocrystals, and electron-phonon coupling.¹⁶ Alternatively, the radiative transition can occur with the involvement of interbandgap states, *e.g.*, introduced by dopants. Excited electrons (in exciton or trap states) can also undergo non-radiative relaxation processes. Lastly, hot electrons, electrons at excitonic states, and in trap states can all be extracted from the semiconductor and transferred to another species, generally molecules. This process is harnessed in photocatalysis. Note that similar considerations made for hot electrons can be made for hot holes, with the caveat that the processes take place in the VB instead (see Fig. 2).

2.1.2. Scattering. When photons impinge on a nanomaterial, their electromagnetic field can be perturbed by the encounter with a system having a refractive index different from the one of the medium in which the photons are propagating. The photons are therefore scattered. Most photon scattering events are elastic in nature, meaning that the energy of the photon is conserved in the process. Within the elastic regime, scattering can be described more accurately using one of two formalisms depending on the ratio between the characteristic nanomaterial size (r) and the wavelength of the impinging photon (λ): Mie ($r/\lambda \gtrsim 1$) and Rayleigh ($r/\lambda \ll 1$) scattering. Mie scattering is characterized by preferential forward scattering and geometry-dependent resonance modes, while Rayleigh scattering is more isotropic and scales with the inverse fourth power of the wavelength of incident photons.

Inelastic scattering events involve energy exchange with the crystal structure and, thus, are non-conservative in nature. The resulting photons have either lower or higher energy than the incident one: these events are referred to as Stokes and anti-Stokes scattering, respectively. Raman spectroscopy analyses these photons to obtain information about the material under study.

2.2. Plasmonic materials

In plasmonic nanomaterials, a complete description of the light-matter interaction should consider both optical electronic transitions (see above) and the activation of confined, resonant, collective charge carrier oscillations referred to as localized surface plasmons resonance (LSPRs). The charge carriers participating in these oscillations are either quasi-free electrons or quasi-free holes close to the Fermi state, which falls within an electronic band rather than within the bandgap. From a classical perspective, the external oscillating field drives currents within the nanostructure, constituted by charge carriers in a partially filled band. The mobility and density of these carriers, together with the geometry of the structure, determine the natural oscillation frequencies that can be excited resonantly. This description is consistent with the framework of the Drude model and can be seen as the collective excitation of carriers to states near the Fermi level through intraband transitions. These collective excitation events constitute the excited currents and the LSPR. An alternative perspective, treating both light and the collective response in the material within a quantum mechanical framework, describes the plasmon as a quasiparticle (a boson) that represents a quantized, collective oscillation mode. A photon can therefore interact with a plasmonic nanomaterial when its energy matches the energy of the plasmon (quantum interpretation), which is directly related to the collective charge carrier oscillation frequency (classical interpretation).

Because of the topic of this review, it is useful to briefly highlight the unique features of plasmonic SNCs. Note that an in-depth discussion on this topic was published by Fauchaux *et al.*,¹⁹ and the interested reader is redirected to that publication. While in metal nanoparticles the main charge carriers supporting the LSPR are electrons (negative charges), in SNCs holes (positive charges) in the VB are the main responsible for the plasmonic behavior. In other terms, as the conductive phenomenon occurs close to the high-energy edge of the VB, the band curvature and population – which gives rise to the properties of the hole – also determine the conductivity and plasmonic properties of the system. A key differentiator between metal nanocrystals and SNCs is the charge density (N), which is orders of magnitude lower in SNCs (10^{16} – 10^{21} cm⁻³) compared to their metal counterpart (approx. 10^{23} cm⁻³).¹⁹ The bulk plasma frequency of a material scales with the charge carrier density according to:

$$\omega_p = \frac{Ne^2}{m^* \epsilon_0} \quad (1)$$



where e is the elementary charge, m^* is the charge carrier effective mass, and ϵ_0 is the vacuum permittivity.

Thus, a lower charge density translates to a lower plasma frequency (energy), which corresponds to longer wavelengths. Indeed, plasmonic SNCs show LSPR modes centered in the NIR-midinfrared (MIR) range, while metallic plasmonic nanocrystals are mainly active in the UV-visible. This extension of the range where plasmonic materials can operate using SNCs comes at the cost of an overall lower photon extinction cross section (*vide infra*). Importantly, in SNCs the carrier density can be tuned and dynamically controlled by intentional – and often reversible – introduction of vacancies, as well as by redox or electrochemical means. This charge density tuning offers an additional “knob” to modulate the optical properties of SNCs compared to metallic nanocrystals, where LSPR energy can be mainly modulated through size, morphology, and composition.

2.2.1. Absorption. Photon absorption events in plasmonic nanomaterials can occur due to the activation of LSPRs (supported by intraband electronic transitions) or can prompt interband electronic transitions. In the former case, the induced charge carrier oscillation undergoes a damping/dephasing. In the quantum mechanical interpretation, this corresponds to the plasmon quasiparticle decaying. In so doing, the energy can be dissipated through scattering events between the charge carriers and phonons or defects. This “frictional” decay occurs for excited charge carriers (*i.e.*, quasi-free carriers in the band of the majority carriers) interacting with the nuclei in the material structure, and it leads to heat generation by transferring their kinetic energy to phononic modes. More rarely, the energy of the decaying plasmon can be transferred to charge carriers close to the nanoparticle border, where scattering events can support the generation of hot carriers.²⁰ Hot carriers are also generated by direct optical interband excitation events. Like in the case of non-plasmonic materials, these hot carriers can either thermalize – thus generating heat – or be transferred to external moieties (*e.g.*, molecules).

2.2.2. Scattering. When the activated LSPR mode decays radiatively, the nanoparticle is said to have scattered the photon exciting it. This is an elastic scattering event in that the re-emitted photon has the same frequency as the exciting one. Scattering is dominant in larger plasmonic nanoparticles, scaling in the quasistatic approximation as the square of their polarizability rather than linearly like the absorption; but they also depend on their aspect ratio and overall geometry.²¹ Plasmonic nanoparticles with large scattering cross sections function as effective nanoantennas operating at visible and IR wavelengths, with their excitation dipoles coupling strongly with the far field.

2.3. Some important metrics

2.3.1. Extinction. The combined probability of photon scattering and absorption is referred to as photon extinction. The scattering, absorption, and extinction efficiency of a system are expressed by the respective cross sections (σ_{sc} , σ_{abs} , σ_{ext}). Molar

absorption and extinction coefficients are also used (ϵ_{abs} , ϵ_{ext}). One should note that in single-crystal optical materials (the background field of many researchers also working on optical nanomaterials), extinction and absorption are often used interchangeably; and rightfully so, since scattering is generally negligible in those systems. However, in nanomaterials, sometimes the scattering and absorption cross sections can be of similar order of magnitude, or scattering can even dominate. This is true, for example, in larger plasmonic nanomaterials or when nanoparticles aggregate to yield larger individual units. Careful distinction of scattering and absorption is therefore pivotal in the study of the optical properties of nanomaterials. For example, depending on the application, one would want to maximize solely the absorption or the scattering (rather than the total extinction!) of a nanomaterial.

2.3.2. Heat generation. Non-radiative recombination events lead to localized heat. The figure of merit for this photon-to-heat process is the heat conversion efficiency (HCE), which is defined as the fraction of energy converted to heat over the total amount of energy extinguished by a system.^{22,23}

$$\text{HCE} = \frac{Q_{\text{ext}} - Q_0}{I(1 - 10^{-A_\lambda})} \quad (2)$$

here Q_{ext} is the external heat flux, Q_0 is the heat generated by the dispersant (in the case of a SNC dispersion), I is the laser power, and A_λ is the absorbance at the laser wavelength.

The “external” version of this figure of merit (eHCE) has also been proposed, which considers also the mass of photon absorbing material.²⁴

$$\text{eHCE} = \text{HCE} \cdot a_\lambda = \text{HCE} \cdot \frac{A_\lambda}{\rho L} \quad (3)$$

where a_λ is the mass absorption coefficient of the species under study at the excitation wavelength, ρ is the mass concentration (in mg mL⁻¹) and L is the optical path (in cm) of the setup used to experimentally measure these metrics.

2.3.3. Photoluminescence. Photoluminescence quantum yield (PLQY; the ratio between emitted and absorbed photons) is widely regarded as the key figure of merit for photoluminescence.

$$\text{PLQY} = \frac{n_{\text{em}}}{n_{\text{abs}}} \quad (4)$$

However, brightness (obtained as the product of PLQY and absorption cross section) is a more complete metric that considers both photon absorption and emission efficiency.²⁵

$$B = \sigma_{\text{abs}} \cdot \text{PLQY} \quad (5)$$

2.3.4. Photocatalysis. Alongside creating photogenerated currents in photodetection devices, a key application of extracted electrons is in photocatalysis/photochemistry. While there are several metrics for quantifying the efficiency of photocatalytic processes, internal and external quantum efficiencies are the most used one. The former considers the number of absorbed photons, while the latter the total number of excitation photons. The external quantum yield, also



referred to as apparent quantum efficiency (AQE), is therefore defined as the ratio between reacted electrons and absorbed photons, or – equivalently – between reaction rate (R) and photon flux (ϕ).²⁶

$$\text{AQE} = \frac{R}{\phi} \quad (6)$$

2.3.5. Quantum confinement. An important parameter in the context of SNCs is the exciton Bohr radius.²⁷ This magnitude (a_{B}) is the characteristic size of the quasiparticle exciton (an electron–hole pair interacting through Coulomb interactions), which is a function of the effective mass of the electrons and holes in a specific material.

$$a_{\text{B}} = \frac{\hbar^2 \epsilon}{e^2} \left(\frac{1}{m_{\text{e}}} + \frac{1}{m_{\text{h}}} \right) \quad (7)$$

where \hbar is the reduced Planck's constant, ϵ is the permittivity of the material, m_{e} and m_{h} are the electron and hole masses. When the size of the SNC approaches and gets smaller than the exciton Bohr diameter ($2a_{\text{B}}$), the SNC is said to enter the quantum confinement regime. In this regime, the energy difference between VB and CB increases (*i.e.*, the bandgap widens) following the particle-in-a-box model – hence the absorption and, if present, emission properties of the SNC can be controlled by changing the SNC size. In this Review, we decided to always use the more general term SNCs rather than the possibly more common and now mainstream term of quantum dots (QDs). That is because, strictly speaking, QDs are only those SNCs in quantum confinement regime. SNCs is thus a more generally applicable term.

After this introduction on the basics of light-nanomatter interaction, we focus our attention on the different families of Cu-based SNCs.

3. Pnictogenides

3.1. Copper phosphides

3.1.1. General considerations, structure, and optical properties. Together with copper chalcogenides (Cu–E; see section 4.1), copper phosphide (Cu_{3–x}P) nanocrystals belong to the family of self-doped doped SNCs featuring plasmonic properties. As the generic formula used in this review suggests, Cu_{3–x}P nanocrystals can be prepared in a variety of compositions with varying relative copper content. A theoretical study by Harper *et al.* identified several compositions (CuP₁₀, Cu₂P₇, CuP₂, Cu₂P, and Cu₃P) each with a unique crystal phase.²⁸ The nanocrystals reported in the literature crystallize in the $P6_3cm$ Cu₃P hexagonal phase (Fig. 3a).^{29,30} Correspondingly, the SNCs have a hexagonal nanoplate morphology (Fig. 3b).³¹ Cu_{3–x}P in the $P6_3cm$ phase has been recently suggested to be a semimetal³² – and hence it does not display a bandgap in the classical sense. Yet some studies show well-defined bandgaps with varying values up to >1.0 eV.^{33,34} Cu_{3–x}P SNCs are generally sub-stoichiometric in copper content, featuring free – here synonym of unbound-charge carriers (holes) in the VB and

hence a p-type behavior. This characteristic stems from the fact that the introduction of copper vacancies is thermodynamically favorable. The presence of copper vacancies in the structure also facilitates ion diffusion, thus enabling cation exchange processes.³⁵ In fact, Cu_{3–x}P nanocrystals have been used as sacrificial species for the preparation of, *e.g.*, InP nanocrystals.³⁶

As far as the optical properties are concerned, the positive charge carriers support LSPR modes centered in the NIR-to-MIR range (Fig. 3c).²⁹ The position of the LSPR peaks can be tuned by modulating the charge carrier density (N_{H}) in the nanocrystal. This strategy has been leveraged by Rachkov *et al.*, who developed Cu-coupled reversible redox strategies to controllably modify N_{H} – which lies in the 10^{21} – 10^{22} cm^{–3} range for Cu_{3–x}P nanocrystals.³⁵ Specifically, Cu_{3–x}P → Cu_{3–x–δ}P oxidation was promoted in the presence of I₂ or metal halides and TOP. This process could be reversed by exposing the nanocrystals to Cu⁺ in the form of tetrakis(acetonitrile)copper(i) hexafluorophosphate ([Cu(MeCN)₄]PF₆). The position of the LSPR peak center could be tuned with this strategy by 200 meV in the 700–900 meV range (Fig. 3d). It should be noted that this controlled redox approach draws inspiration from the protocols developed for copper chalcogenide nanocrystals, which have an analogous plasmonic behavior (see below).

3.1.2. Synthesis. In one of the first reports on this nanomaterial, Manna *et al.* prepared Cu_{3–x}P SNCs by bubbling PH₃ as phosphorous source – obtained *ex situ* from the reaction between Ca₃P₂ and HCl – into a solution of CuCl in a mixture of trioctyl phosphine (TOP), trioctyl phosphine oxide (TOPO), and oleylamine (OAm).³⁴ Therein, TOP acted as SNC size tuning agent by determining the nucleation rate and hence the total number of prepared nanocrystals. The central role of TOP in the preparation of Cu_{3–x}P SNCs was also highlighted in a report by De Trizio *et al.*³⁰ The authors developed a one-pot protocol whereby a CuCl solution in an alkylamine mixture is injected into a TOP and TOPO mixture at high temperature (370 °C). Depending on the amount of TOP in the reaction mixture, a final nanoplate morphology was obtained *via* homogeneous Cu_{3–x}P nucleation and growth or a gradual phosphorization of Cu intermediate nanocrystals (Fig. 3e). In the latter case, the conversion proceeded through growth along two main epitaxial Cu/Cu_{3–x}P interfaces to yield Janus Cu–Cu_{3–x}P nanocrystals that ultimately produced Cu_{3–x}P nanoplates (Fig. 3c).³⁰ Liu *et al.* proposed the preparation of Cu_{3–x}P SNCs *via* hot injection of a tris(trimethylsilyl)phosphine (TMS)₃P solution in 1-octadecene (ODE) into a CuCl solution in TOP/OAm kept at up to 180 °C. Intriguingly, a “wet annealing” process carried out at 300 °C led to a reshaping of the hexagonal nanoplates into disks, likely owing to the OAm-driven etching.²⁹ More recently, Rachkov *et al.* reported on a one-pot synthesis approach using copper halides, tris(diethylamino)phosphine [P(NEt₂)₃], OAm, and trioctylamine.³⁹ Therein, the authors placed the nucleation event at approximately 230 °C and identified the P/Cu ratio in the reaction mixture as the key parameter to control the nanocrystal size. Thorough investigation of the reaction intermediates *via* nuclear magnetic resonance (NMR) shed light on the





Fig. 3 Copper phosphide SNCs. (a) Unit cell and simulated X-ray powder diffraction (XRPD) pattern of the Cu₃P phase generally reported in SNCs. (b) Transmission electron microscopy (TEM) image of Cu₃P SNCs showing the characteristic hexagonal nanoplate morphology. Adapted with permission from ref. 29, Copyright 2016, John Wiley and Sons. (c) Extinction spectra of Cu₃P SNCs synthesized with different Cu:TOP feeding ratios. The LSPR peak shifts at longer wavelengths with increasing TOP content. Adapted with permission from ref. 29, Copyright 2016, John Wiley and Sons. (d) Scheme outlining the reversible redox process that allows controlling the carrier density and hence the LSPR peak position and intensity in Cu_{3-x}P SNCs. Adapted from ref. 37, Copyright 2024, American Chemical Society. (e) Two possible Cu₃P SNC growth mechanisms: homogeneous nucleation and growth (top) or gradual phosphorization of metallic Cu nanoparticles (bottom). Adapted with permission from ref. 30, Copyright 2012, American Chemical Society. (f) An example of Cu₃P SNC use in nanomedicine: thermal images of a tumour model during Cu₃P SNC-mediated photothermal therapy (top) and comparison of the effect in control models (bottom). The reported H&E-stained tumour sections and the photos of mice at the end of the treatment show a marked tumour reduction. Adapted with permission from ref. 38, Copyright 2021, American Chemical Society.

nature of the Cu–P species involved in the nanocrystal nucleation and growth. The synthesis of these SNCs was also reported *via* single-source precursor [Cu(H)(PPh₃)₆] by Downes *et al.*, who also identified the presence of a Cu nanocrystal intermediate gradually converted to Cu_{3-x}P.⁴⁰

3.1.3. Applications. Cu_{3-x}P SNCs have been proposed in fields including electrocatalysis, photonics, and nanomedicine. For example, Mondal *et al.* used Fe-doped Cu₃P SNCs for water splitting, while Downes *et al.* showcased the potential of these nanomaterials for CO₂ reduction.⁴¹ Liu *et al.* proposed instead the use of Cu₃P SNCs for laser applications in the telecom band. The carefully-tuned LSPR band at 1.5 μm made them suitable as nonlinear absorbers in a high-energy Q-switched laser.²⁹ Cu₃P SNCs have also been shown to be effective photothermal (PT) agents and photosensitizers (PS) in photothermal and photodynamic therapy, respectively (Fig. 3f).^{31,38} The LSPR properties of Cu₃P SNCs were also harnessed by Sun *et al.* to enhance the performance of graphene-based photodetectors, identifying the nature of the capping ligands as a key parameter in enabling good performance.⁴⁰

3.2. Copper nitrides

3.2.1. General considerations, structure, and optical properties. Cu₃N crystallizes in a cubic anti-ReO₃ structure (space

group *Pm* $\bar{3}$ *m*), where the Cu atoms occupy the corners and face centers of the cube, and nitrogen atoms sit at the body center, forming a 3D open framework (Fig. 4a).⁴² The morphology of the SNCs in this phase is also cubic (Fig. 4b). Moreover, the crystal structure is amenable to copper replacement with other metal ions. In fact, similarly to other Cu-based nanomaterials (see below), Cu₃N SNCs can act as cation exchange platforms for the synthesis of SNCs, such as Ni₄N and CoN.⁴³

In its bulk form, copper nitride (Cu₃N), is a semiconductor with a calculated indirect bandgap of 1.0 eV, while experimental reports show values in the 0.8–1.9 eV range – depending on the preparation method and actual stoichiometry. Aside from high values of bulk absorption coefficient (10⁴–10⁵ cm⁻¹),^{45–47} Cu₃N in the form of SNCs has been shown to exhibit electro- and photoluminescence in the visible-NIR region (Fig. 4c).^{44,48,49}

3.2.2. Synthesis. The synthesis of Cu₃N SNCs typically involves the reaction of copper salts with nitrogen precursors under reducing or inert conditions. A common approach proceeds through the thermal decomposition of a copper compound in the presence of an organic amine, which serves as both solvent and nitrogen source. The reaction is typically carried out at 220–260 °C under inert atmosphere, producing colloiddally stable nanocubes or quasi-spherical nanocrystals.



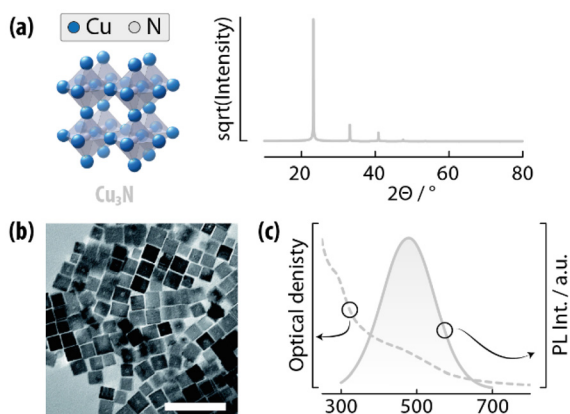


Fig. 4 Copper nitride SNCs. (a) Unit cell and simulated XRPD pattern of Cu₃N. (b) TEM image of Cu₃N SNCs showing the characteristic cubic morphology. Scale bar is 100 nm. (c) Extinction and emission spectrum of Cu₃N SNCs. Adapted with permission from ref. 44, Copyright 2020, Royal Society of Chemistry.

Explored copper precursors are Cu(NO₃)₂, copper methoxide (Cu(OMe)₂), and copper acetate. Typically employed amines are OAm, octadecylamine, and benzylamine. Alternatively, urea and ammonia were also explored as nitrogen sources, conducting the reaction in high-boiling point alcohols such as 1-nonanol and decreasing the length of the aliphatic chain all the way down to 1-pentanol.⁴⁸ The formation mechanism of Cu₃N SNCs entails an initial nucleation of small, quasi-spherical nanoparticles, followed by their ripening into nanocubes. Prolonging the reaction time beyond the formation of the nanocubes led to SNC degradation, with the appearance of metallic copper. A recent study by Zheng *et al.* showed the possibility to grow Cu₃N nanorods using Cu-acetate and OAm in a heat up approach, given the preferential adsorption of the ligands on the SNC so to favor growth along the [100] direction.⁵⁰ The nature of the amine was critical in the morphology and quality of the final crystals, albeit the authors did not provide a definitive explanation for this effect. The presence of O₂ in the reaction environment was shown to be pivotal in obtaining homogeneous morphology and size, hence the authors suggested an active role of this molecule as oxidizer of aminated species during the SNC synthesis.

Among the publications on the subject, the manuscript by the group of De Roo⁵¹ stands out as a pleasant example of a didactical discussion on the key role that the purification step plays in the preparation of high-quality, stable colloidal SNCs.

3.2.3. Applications. Reported examples of uses of Cu₃N SNCs involving light include mainly their use in photocatalysis.^{52,53} The most relevant is the report by Barman *et al.*, who showed that coupling of these SNCs with Au nanoparticles affords photocatalysts for organic dye degradation, thanks to the favorable band alignment between Au and the semiconductor.⁵⁴ Dasgupta *et al.* recently reported on Cu₃N-Ag SNCs for preparing a nanostructured electrode to be used in food spoilage detection systems.⁵⁵ The NH₃ sensor works

through a chemiresistive working principle, which is enhanced under UV exposure. This effect stems from the presence of Cu₃N-Ag SNCs, which promote electron-hole separation. Cu₃N SNCs have also been suggested as candidates for absorbing materials in photovoltaic devices, but we could not find reports on this use.

4. Chalcogenides

4.1. Binary compounds

4.1.1. General considerations, structure, and optical properties. Binary copper chalcogenides (Cu_{2-x}E, E = S, Se, Te) SNCs are the prototypical example of plasmonic semiconductor nanomaterials, with the first report on Cu_{2-x}S SNCs by Luther *et al.* in 2011.⁵⁶ Cu_{2-x}E SNCs are self-doped p-type semiconductors, with copper vacancies introducing free/unbound holes that result in tunable localized surface plasmon resonance (LSPR) in the NIR-to-MIR range. A special case is represented by CuS (covellite) SNCs, where the delocalized holes in the VB do not arise from copper vacancies (whose formation is energetically unfavored), but rather from the intrinsic band structure of the nanomaterial.⁵⁷ Carrier densities reaching ~10²¹–10²² cm⁻³ were reported – a value that approaches metallic regimes.

Copper sulfides. Materials of the Cu_{2-x}S family have direct bandgaps in the 1.1–2.0 eV range, depending on the actual composition of the material.¹¹ Indeed, Cu_{2-x}S (1 < x < 2) exhibits multiple crystalline forms, including chalcocite (Cu₂S; high chalcocite, *P*6₃/*mmc*; low chalcocite, *P*2₁/*c*), djurleite (Cu_{1.97}S; *P*2₁/*n*), roxbyite (Cu_{1.81}S; *P*1̄), digenite (Cu_{1.8}S; *R*3̄*m*), anilite (Cu_{1.75}S; *P**nma*), and covellite (CuS; *P*6₃/*mmc*) (Fig. 5a). In all these materials, copper is found almost exclusively in its Cu⁺ oxidation state, while sulfur can take up a valency of –1 (CuS) or –2 (Cu₂S).⁵⁸ Note that all the above-mentioned structures have well-defined stoichiometries. Yet, some of them tolerate introduction of copper vacancies, which induces deviations from the theoretical compositions. Importantly, the more the stoichiometry deviates from Cu₂S, the higher the value of N_H. That is to say, the higher the value of x in Cu_{2-x}S, the more charge carriers (holes) are present in the VB and, hence, the stronger the LSPR (Fig. 5b). The first report by Luther *et al.* showed a NIR LSPR peak for djurleite SNCs whose position corresponded to a hole density of 10²¹ cm⁻³ and, hence, a composition Cu_{1.93}S. Yet, CuS supports the strongest LSPR modes (Fig. 5c and d). Its SNCs generally grow along preferential crystallographic directions <010>, <100>, and <110>, often yielding a triangular or, more often, hexagonal nanoplate morphology (Fig. 5e and f). The nanoplate morphology of these SNCs results in both in-plane and out-of-plane LSPR modes.^{57,59} Moreover, very recently *via* single-particle study using electron energy loss (EEL) spectroscopy, Elibol *et al.* showed that CuS SNCs in the form of thin nanoplates feature a plethora of plasmonic modes.⁶⁰ The same authors also observed several visible and NIR emission signals under electron beam exposure.



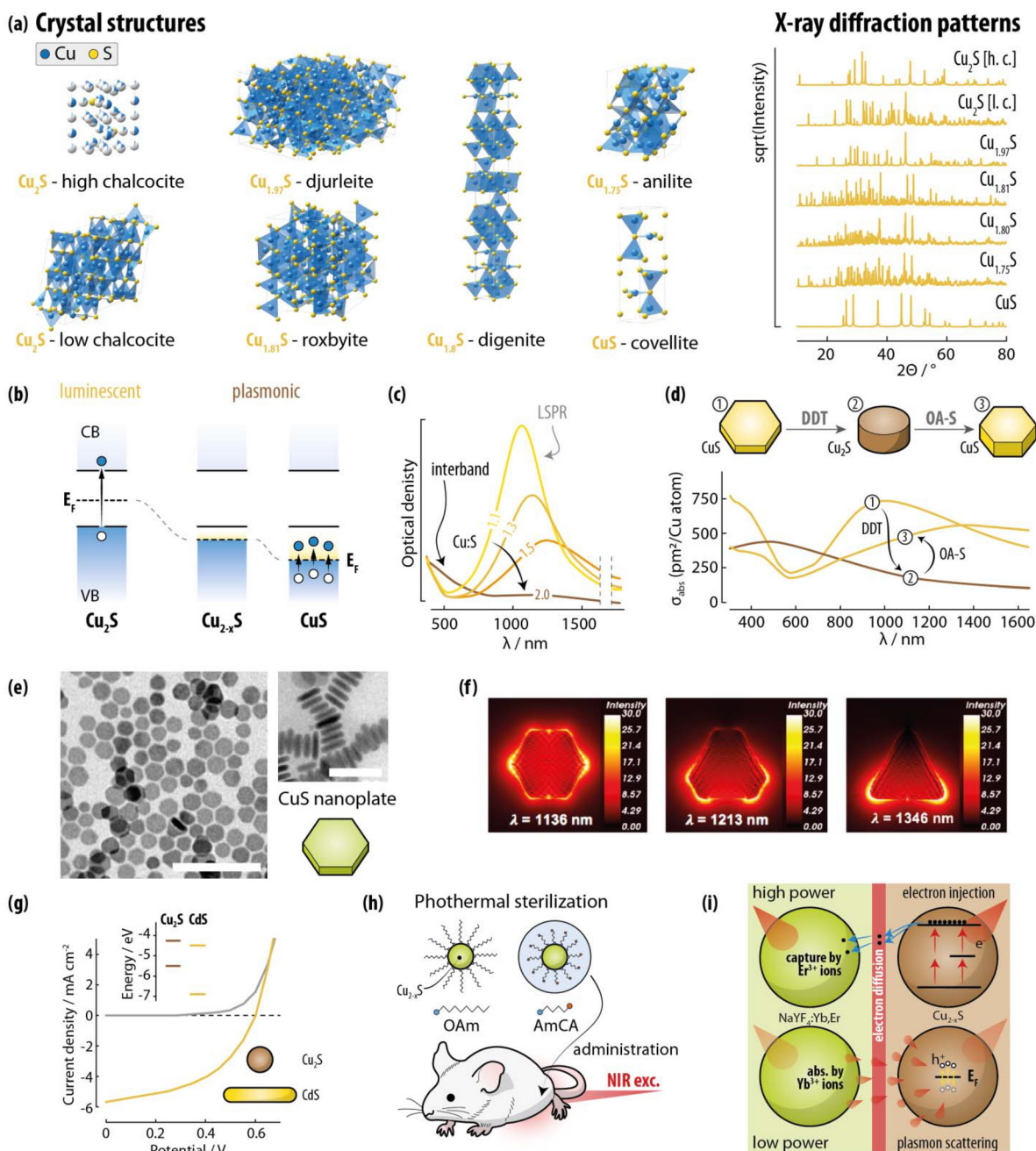


Fig. 5 Copper sulphide SNCs as representatives of the binary copper chalcogenide family. (a) Unit cell and simulated XRPD patterns of the most reported Cu_{2-x}S compositions. (b) While Cu_2S displays the electron structure of a typical SNC (with the Fermi level in the bandgap), Cu_{2-x}S compositions with $x > 0$ have excess holes in the VB and hence the Fermi level falls within this band. As such, LSPR-supporting intraband transitions can take place in SNCs with these compositions. (c) Extinction spectrum of Cu_{2-x}S SNCs with different Cu : S values, showing a gradual decrease in the LSPR mode intensity with increasing Cu^+ content. Adapted with permission from ref. 58, Copyright 2013, American Chemical Society. (d) Reversible transformation between CuS and Cu_2S SNCs through the use of DDT and OA-S, alongside the corresponding change in the extinction spectra. Adapted with permission from ref. 61, Copyright 2017, American Chemical Society. (e) TEM images of CuS SNCs showing the typical nanoplate morphology of this class of SNCs. Scale bars are 100 (left) and 50 (right) nm. Reprinted with permission from ref. 57, Copyright 2013, American Chemical Society. (f) Near-field intensity maps calculated at the resonant condition for CuS SNCs with different morphologies for in-plane illumination. Reprinted with permission from ref. 59, Copyright 2019, American Chemical Society. (g) Current density–voltage characteristics of a photovoltaic device implementing Cu_2S – CdS nanocrystals under no illumination (black) and under standard illumination (red). Adapted with permission from ref. 62, Copyright 2008, American Chemical Society. (h) Use of Cu_{2-x}S SNCs as light-driven sterilants. ref. 63. (i) Two mechanisms of upconversion emission enhancement in $\text{NaYF}_4:\text{Yb},\text{Er}$ supported by Cu_{2-x}S SNCs. Electron injection to Er^{3+} ions occurs at high excitation powers, while plasmonic scattering dominates at low excitation powers. Adapted with permission from ref. 64, Copyright 2016, American Chemical Society.



detailed study, the authors identified the crucial role played by the amine in coordinating Cu^+ ions, proposing a $\text{Cu}_x(\text{OAM})_y(\text{ODE})_z\text{Cl}_x$ complex as intermediate for the nucleation and growth of the nanoplates, with oleyl ammonium oleate acting as a growth controlling species. Single-source precursors are also used in either heat-up or hot-injection thermal decomposition methods. These precursors include copper dithiocarbamates, xanthates, thiobenzoate, thiocyanate, and selenocyanate.^{87,88} Intriguingly, already in 2003, Larsen *et al.* reported on a solventless thermal decomposition of Cu-DDT to yield Cu_2S nanorods and SNCs, even though no optical characterization was reported.⁸⁹ Hsu *et al.* followed in these footsteps by reporting on nanodisks with varying Cu_{2-x}S stoichiometry using both a solventless thermal decomposition approach and a heat-up, solvent-based technique, which allowed controlling the stoichiometry by tuning the Cu:S feeding ratio.⁹⁰ Saldanha *et al.* also proposed a generalized heat-up strategy for the preparation of a library of Cu_{2-x}S , $\text{Cu}_{2-x-y}\text{S}_x\text{Se}_y$, and $\text{Cu}_{2-x-y}\text{S}_x\text{Te}_y$ SNCs, where copper acetylacetonate and chalcogen powders (S, Se, Te) are dispersed in a DDT/OA mixture, followed by heating to 200–220 °C.⁸³ A one-pot, heat-up approach was proposed by Xiao *et al.* too, who used different Se sources (selenurea, SeO , Se) to obtain Cu_{2-x}Se nanomaterials with different morphology (spherical SNCs, nanorods, nanoplates).⁹¹ Similarly, Lord *et al.* obtained Cu_{2-x}Se SNCs starting from copper acetylacetonate and diphenyl diselenide in OAm and ODE.⁷³

The synthesis of pure Cu_{2-x}Te SNCs is relatively less explored, and the quality of the SNCs in terms of size homogeneity is generally poorer. An approach that yields uniform cubic Cu_{2-x}Te SNCs has been proposed by Li *et al.*, who used OAm, TOP, and TOPO as growth directing agents, while the Te precursor was prepared mixing TOP-Te and lithium bis(trimethylsilyl) amide in ODE.⁹² Nanorods and nanoplates were also obtained controlling the reaction parameters. Willhammar *et al.* later adopted this synthesis method and modified it into a one-pot heat-up approach, obtaining similar results in terms of SNCs quality.⁷⁹ An intriguing alternative was proposed by Kriegel *et al.*, who prepared sacrificial CdTe SNCs to then convert them into Cu_{2-x}Te SNCs.⁸⁰ This approach leverages decades-worth of knowledge on the synthesis of quantum dots (QDs) with controlled size and morphology (*e.g.*, spheres, rods, tetrapods), bypassing the difficulties in adjusting these features directly during the synthesis of Cu_{2-x}Te SNCs. While heat-up and hot-injection are arguably the most explored approaches for the preparation of Cu_{2-x}E SNCs, sonochemical approaches⁹³ and microwave-driven heating methods^{94,95} have also been proposed.

Given the widespread investigation on the use of Cu_{2-x}E SNCs for biomedical applications (see below), both water-transfer methods and direct synthesis in aqueous media have been developed – in both cases, drawing inspiration from the protocols developed for QDs. For instance, van Oversteeg *et al.* proposed ligand exchange of DDT or TOPO molecules on the Cu_{2-x}S SNCs surface for hydrophilic mercapto-acids (*i.e.*, 1-mercaptoundecanoic acid, 1-mercaptopropyl acid) or

sulfide ions (S^{2-}) from Na_2S .⁹⁶ Similarly to the approach followed by Saldanha *et al.*,⁸³ exchange with thiolated polyethylene glycol (PEG) was also demonstrated by Marin *et al.* for CuS nanoplates.⁸⁷ The same authors also demonstrated the straightforward synthesis of CuS SNCs in an aqueous solution of methoxy-terminated PEG molecules as growth controlling ligands.⁹⁷ The so-obtained CuS SNCs display an easily modifiable surface chemistry, which makes them amenable to applications in different scenarios. As shown by Gan *et al.*, Cu_{2-x}Se SNCs can also be prepared in aqueous environments using polyvinyl pyrrolidone (PVP) as growth directing agent at room temperature.⁹⁸ Cu_{2-x}S SNCs were also prepared directly in aqueous phase using citrate ions as growth controlling species and Na_2S as the precipitating agent.⁹⁹

4.1.3. Applications. The use of Cu_{2-x}E SNCs include energy conversion and storage, optoelectronics, catalysis, and photonics; yet, their use has been particularly investigated in the biomedical field.

Photovoltaics. The interest in Cu_{2-x}S SNCs for photovoltaic applications¹⁰⁰ is rooted in the discovery of the CdS/ Cu_2S heterojunction back in 1954.¹⁰¹ To that end, a mix of Cu_2S SNCs and CdS nanorods were used to prepare a solar cell with 1.6% power conversion efficiency (PCE) (Fig. 5g).⁶² Cu_{2-x}S SNCs featuring NIR LSPR too were used in photovoltaics, specifically to prepare conductive films later implemented in the architecture of a solar cell – again, together with CdS nanorods – to achieve PCE of 0.24%.¹⁰² CuS SNCs have also been used as hole-transporting-material in perovskite solar cells improving the cell's performance, but being limited in applicability by the large energy barrier between the VBs of perovskite and CuS.¹⁰³ On the other hand, Cu_2S SNCs have potential for the preparation of solution-processed electrodes in QD-sensitized solar cells (QDSSCs). The review by Burda and Zhao specifically tackles this application scenario of Cu_{2-x}S SNCs, alongside their use as electrodes for Li ion batteries.¹⁰⁰

Photocatalysis. The tunable optical properties of Cu_{2-x}E SNCs make these nanomaterials suitable for photocatalysis too. Yet, their stability during operation is an issue for sustainable use of these materials in photocatalysis. Lie *et al.* showed that Cu_{2-x}Se SNCs with controllable copper deficiency boost catalytic performance by >500× in luminol- H_2O_2 reactions, due to elevated hole density and ionic mobility. These features facilitate electron transfer processes that sustain H_2O_2 decomposition.¹⁰⁴ Sun *et al.* showed that Au- Cu_{2-x}S heterostructures act as effective photocatalyst through the study of rhodamine B as model pollutant,¹⁰⁵ while Hoàng Ly *et al.* used CuS hollow nanospheres for photocatalytic reduction of CO_2 into higher-value chemical species.¹⁰⁶ A recent review by Li *et al.* thoroughly discusses the latest advancements and challenges in Cu_{2-x}S photocatalysts specifically.¹⁰⁷

Biomedicine. In biomedicine, Cu_{2-x}E SNCs have been extensively studied as effective light-to-heat converters for photothermal applications. The possibility of obtaining LSPR in the NIR even at sizes well below 10 nm is an important differentiator compared to noble metal-based plasmonic nanomaterials – where NIR LSPR modes are obtained generally with bigger par-



ticles. The advantages of keeping the nanoparticle size small are two: (i) for some biomedical applications smaller particles are preferred due to their ability to elude certain biological barriers¹⁰⁸ and (ii) smaller plasmonic nanoparticles feature dominant absorption *vs.* scattering, which is pivotal for many therapeutic and imaging applications. For instance, CuS SNCs were used for NIR-driven male mice sterilization, thanks to combined heating and reactive oxygen species (ROS) production (Fig. 5h).⁶³ Similar materials and approach have been used for bacterial sterilization¹⁰⁹ and to promote osteogenesis in periodontal therapy¹¹⁰—in the latter case, upon incorporation of the SNCs in a polymeric hydrogel. In the same vein, incorporation of CuS SNCs – alongside deferoxamine – in a biomimetic hydrogel yielded patches for photothermal therapy of infected diabetic wounds.¹¹¹ Modification of CuS SNC surface with TAT and RGD peptides allowed Li *et al.* to achieve cell nucleus targeting at the tumor site, where NIR excitation resulted in localized photothermal therapy that prevented cancer recurrence.¹¹² Neutrophil-erythrocyte membrane-coated hollow CuS SNCs were employed by Xue *et al.* to treat osteoarthritis *via* NIR-triggered local heating of joints.¹¹³ Copper selenide and telluride SNCs have been employed too in similar scenarios.^{114–116} Cu_{2–x}Se SNCs and Au-Cu_{2–x}Se heterostructures have also been used as contrast agents for *in vivo* photoacoustic imaging by Liu *et al.*^{117,118} The groups of Zhang reported on ultrathin CuS SNCs for photoacoustic imaging too,¹¹⁹ while Marin *et al.* showcased the first example of negative optical coherence tomography contrast agents in the form of CuS nanoplates.⁸⁷

Others. A less common application scenario is the enhancement of upconversion emission *via* deposition of layers of luminescent NaYF₄:Yb,Er nanoparticles and Cu_{2–x}S SNCs separated by a buffer layer of MoO₃ of controlled thickness (8 nm) (Fig. 5i). The >200-value reached by the enhancement factor was explained considering the increased scattering of the 980 nm excitation photons by the plasmonic film (at low power), as well as electron transfer from Cu_{2–x}S SNCs to the upconverting nanoparticles upon two-photon excitation of the sulfide species (at high power).⁶⁴ Worth highlighting is also the work by Zhang *et al.*, who demonstrated the potential of these materials for sensing, showing that aggregates of Cu_{2–x}S SNC can be effectively used in surface-enhanced Raman scattering (SERS).¹²⁰

4.2. Ternary compounds and beyond

4.2.1. General considerations, structure, and optical properties. Ternary copper chalcogenide SNCs have general formula Cu_xM_yE_z, where M = In, Ga, Al, Fe, Bi, Sb, Sn, Ge, Cr, Zn. These materials can usually retain their crystal structure even when strong deviations from the theoretical stoichiometry occur. They can be prepared in alloyed forms, such as CuIn_{1–x}Ga_xS₂ or Cu₃(Sb_{1–x}As_x)S₄ or they can be alloyed with other binary semiconductors, such as CuInS₂-ZnS (generally indicated as Cu–In–Zn–S). These compositions are often indicated as quaternary – or even higher-order – compounds, and they often have intermediate properties compared to the indi-

vidual components. Indeed, this type of alloying is usually a strategy employed to fine tune the absorption and emission of photoluminescent SNCs between the two “extreme” properties featured by the two components.

Below, we provide a discussion on structure–property relationships in multinary Cu-based SNCs broken down by compound families: Cu-III-E, Cu-IV-E, and Cu-V-E (where III, IV, and V indicate the groups the metals belong to), Cu-TM-E (TM = transition metal), and quaternary materials.

Cu-III-E compounds. CuInE₂ and CuGaE₂. CuInE₂ and CuGaE₂ are families of p-type semiconductors. Together with CuGaE₂, AgInE₂, and AgGaE₂ are materials that have been investigated and/or are currently successfully employed for the preparation of photovoltaic devices including thin-film solar cells and QDSSCs (see below). Their properties are dominated by a marked tolerance against deviations from both cation and anion stoichiometries.

CuInS₂ is arguably the most studied ternary copper chalcogenide for the preparation of SNCs. This material crystallizes in three polymorphs: roquesite (*I42d*; the thermodynamically stable phase akin chalcopyrite), wurtzite (*P6₃mc*), and zinc blende (*F43m*) (Fig. 7a).¹⁵ Chalcopyrite and wurtzite are the polymorphs usually stabilized at the nanoscale (ch-CIS and wz-CIS). Most publications on CuInS₂ SNCs deal with chalcopyrite phase, although chalcopyrite can hardly be discerned from zinc blende simply *via* X-ray powder diffraction (XRD) – the prime technique employed for phase attribution for these materials. In this phase, Cu⁺ and In³⁺ ions occupy the crystal sites in an ordered fashion, while in the other two polymorphs the cations are randomly distributed over the cation sublattice. Another difference between phases lies in the morphology of the corresponding SNCs: ch-CIS takes up a tetrahedral habitus, while wz-CIS is generally assumed quasi-spherical (although actually having more complex morphologies such as truncated octahedron; Fig. 7b). Note that the chalcopyrite crystal phase is recurring in the chemistry of multinary Cu-chalcogenides, and it can be seen as a derivation of the zinc blende crystal structure (see below).

CuInS₂ has a direct bandgap of approximately 1.5 eV, absorption coefficients in the 10⁵ cm⁻¹ range, and a Bohr exciton radius of 4.1 nm. The interest around CuInS₂ SNCs stems from the visible and NIR photoluminescence that these nanomaterials display, with PLQY values at times higher than 90%.¹²⁵ Moreover, CuInS₂ SNCs feature an absorption cross section of 1.8 × 10⁻¹⁶ a³ cm² (where *a* is the SNC radius): On par with typical values of for CdSe QDs.¹²⁶ As such, CuInS₂ SNCs have brightness values comparable to those of CdSe QDs, while lacking more toxic Cd²⁺ in their composition. However, the optical properties of CuInS₂ SNCs are drastically different from the ones of binary QDs: (i) the excitonic absorption peak is poorly defined in the absorption spectrum, (ii) the emission peak is broad (fwhm > 300 meV) and (iii) characterized by a large Stokes shift (300–500 meV), (iv) the characteristic photoluminescence decay time is long (100–500 ns) (Fig. 7c). To explain these characteristics, several models have been proposed, which invoke electronic transitions of different





Fig. 7 Copper indium sulphide SNCs as representatives of the Cu-III-E family. (a) Unit cell and simulated XRPD patterns of the three CuInS_2 polymorphs. (b) TEM images and corresponding typical morphology of CuInS_2 in the chalcopyrite (ch- CuInS_2) and wurtzite (wz- CuInS_2) crystal structure. TEM images reprinted with permission from ref. 121 (Copyright 2010, American Chemical Society) and ref. 122 (Copyright 2017, American Chemical Society). (c) Typical extinction and emission spectrum (left), photoluminescence decay curve (centre), and photo under UV illumination (right) of a CuInS_2 SNCs dispersion. Adapted with permission from ref. 123 Copyright 2019, American Chemical Society. (d) Energy level schemes and corresponding transitions that underpin the optical properties of luminescent and plasmonic CuInS_2 SNCs. (e) Extinction spectrum of plasmonic CuInS_2 SNCs. Adapted with permission from ref. 124, Copyright 2012, American Chemical Society.

types: donor-acceptor-pair (DAP), band-to-band, free-to-bound, and self-trapped excitons. In a recent review by the Wang group, the debate over the origin of the luminescence in these nanomaterials is described in detail.¹⁵ Albeit no consensus has been reached yet, the community lean towards the free-to-bound model, where the CuInS_2 emission stems from the recombination between a free/unbound electron in the CB and a hole localized at a copper-related trap state just above the VB (Fig. 7d). This attribution is also supported by the striking similarities observed with the emission of Cu^+ -doped ZnS and CdS QDs.¹²⁷ Intriguingly, the emission of CuInS_2 SNCs can be tuned by adjusting the Cu:In ratio.¹²⁸ Specifically, within a certain range of off-stoichiometry, Cu-deficient SNCs display blue-shifted emission and an increased PLQY. The origin of the effect has not been clearly identified, with some authors claiming the involvement of copper vacancy-related states (V_{Cu}) in the radiative recombination process.¹²⁹ However, ensuring consistent particle size when varying the Cu:In ratio is difficult, thus complicating the study of off-stoichiometry effects. To that end, other approaches for tuning and improving the optical properties of CuInS_2 SNCs entail:

(i) Control of the SNC size to harness the quantum confinement effects.

(ii) The growth of core/shell structures with type-I band alignment using ZnS or CdS as shell material. In the case of ZnS, the reactivity of the ZnS precursors¹²⁵ and the surface chemistry of the core SNCs can be tuned to induce partial Cu^+ - Zn^{2+} cation exchange (hence forming a gradient shell) or exclusively ZnS overgrowth.¹³⁰

(iii) Alloying with other elements such as Ag^+ to substitute Cu^+ and Ga^{3+} or Al^{3+} (albeit with more difficulty due to the

ionic radius mismatch) for In^{3+} .^{131–135} This strategy enables spanning a broad spectral range, since the obtained SNCs can show emission spanning from the blue all the way to the NIR (approx. 450–1000 nm). Alloying with Zn^{2+} is also an explored avenue.^{136,137} This strategy exploits the miscibility over the whole composition range of CuInS_2 and ZnS, due to the structural similarities between the two materials.

Besides, photoluminescence, CuInS_2 SNCs can also sustain LSPR modes (Fig. 7e). The Rosenthal group was the first to report on plasmonic chalcopyrite CuInS_2 SNCs.¹²⁴ They were synthesized with highly reactive bis(trimethylsilyl) sulfide, which induced the incorporation of metal vacancies – due to the rapid SNC formation – that support LSPR. An $N_{\text{H}} = 2.6 \times 10^{20} \text{ cm}^{-2}$ was found analyzing the extinction spectrum according to the Drude formalism. Subsequently, the groups of Swihart¹³⁸ and Ghosh¹³⁹ showed that plasmonic CuInS_2 SNCs are obtained in the presence of excess Cu^+ in the structure. This photoluminescence-plasmonic duality is an intriguing feature of CuInS_2 SNCs, which highlights the versatility of this type of nanomaterial.

CuInSe_2 is similar in structure and properties to CuInS_2 .^{140,141} CuInTe_2 is instead mainly found in the chalcopyrite polymorph. In bulk, CuInSe_2 and CuInTe_2 have bandgaps of 1.04^{140,142,143} and 1.06 eV,¹⁴⁴ with absorption coefficients in the 10^5 cm^{-1} range. As their bandgap, the emission of these materials is shifted more towards the NIR range than CuInS_2 . The means of controlling the optical properties of these SNCs are also the same as the ones described for CuInS_2 SNCs.^{142,145–147} The emission of CuInSe_2 SNCs, for instance, can be effectively tuned from approx. 800 to 1400 nm by means of alloying with Ag^+ and Zn^{2+} , followed by overgrowth



of a ZnS shell for maximized PLQY.¹⁴⁸ Like in the case of binary Cu_{2-x}E SNCs, the telluride of the CuInE₂ family are underexplored compared to sulfides and selenides, and are often reported in alloys with lighter chalcogenides.¹⁴⁶

Similar considerations can be made for the CuGaE₂ SNC family. CuGaS₂, CuGaSe₂, and CuGaTe₂ all crystallize in the chalcopyrite crystal phase and have direct bulk bandgaps of 2.43, 1.68, and 1.2–1.9 eV, respectively. The optical properties of SNCs made of these materials (as well as the strategies proposed to tune them) resemble those reported for CuInE₂ SNCs, yet they are more shifted at higher energy (*i.e.*, shorter wavelengths).^{149–152}

CuAlE₂. Closely related to two families of semiconductors mentioned above are CuAlS₂ and CuAlSe₂. Both these semiconductors crystallize in the chalcopyrite phase and have wide bandgaps of approximately 3.5 and 2.7 eV, respectively.^{153–156} Although core-only CuInS₂ SNCs do not appear to feature appreciable photoluminescence, the Pandey group showed that preparation of type-II core/shell SNCs using CdS as shelling material results in tunable emission in the visible-NIR range.¹⁵⁷ The observed emission was explained in light of delocalization of the exciton in the shell, also given the dependence of its energy on the shell thickness. The same group also prepared CuAlS₂-ZnS heterostructures, observing a similar charge delocalization, thus achieving bandgap values in the 1.5–2.0 eV range despite the fact that both core and shell materials have bandgaps in the UV-blue region.^{158,159} Cu-Al-S/ZnS SNCs were later prepared by Hansen *et al.*, obtaining blue-emitting species with PLQY as high as 18%. The study highlighted the presence of a complex, defect-rich electronic structure of the SNCs, which was induced also by deliberate synthesis of highly Cu-defective nanomaterials. Combination of theoretical calculations and experimental observations suggested that the SNC cores were actually made of an ordered defect compound with composition CuAl₅S₈. Intriguingly, the optical properties of CuAlS₂ SNCs were also optimized *via* a combination of design-of-experiments and density functional theory (DFT) calculations by Baum *et al.*¹⁶⁰ In that study, also NIR absorption features were observed in some samples, suggesting the existence of resonances in CuAlS₂—which are currently underexplored.

Cu-IV-E compounds. Cu-Sn-E. In this family of semiconductors, the most common compounds reported in the form of SNCs are Cu₂SnS₃ (mohite, *Cc* – although other polymorphs are also known)¹⁶¹ and Cu₃SnS₄ (kuramite, *I42m*) (Fig. 8a). A metastable wurtzite structure is also reported in the literature only for SNCs, and Cu-Sn-S SNCs in the kesterite phase were obtained too.^{162,163} Cu₂SnSe₃ crystallizes in the cubic *F43m* (zinc blende)¹⁶⁴ or monoclinic *Cc* (okruginite)¹⁶⁵ structures. For an overview on Cu-Sn-Se SNCs, the reader is directed to ref. 164 and 166. Such richness of phases obtainable at the nanoscale is testament of the versatility of nanomaterials, where ligand and precursor nature and concentration act synergistically to determine the properties of the SNCs. SNCs of Cu-Sn-E display strong light absorption capabilities as well as LSPR modes supported by metal vacancies in

the structure. This observation was corroborated by a study by the Swihart group, who showed that the strength of the resonance decreased with increasing Sn content in the SNCs (Fig. 8b and c).¹⁶²

Cu-Ge-E. Among this family of materials, Cu₂GeS₃ and Cu₂GeSe₃ are the most studied as far as SNC production is concerned. Cu₂GeS₃ (bulk direct bandgap of 1.5–1.6 eV) is reported to crystallize in bulk and in SNC form in the *Cc* space group.¹⁶⁵ Yet, in SNCs it crystallizes also in cubic zinc blende-like (*F43m*) or metastable hexagonal wurtzite-like (*P6₃/mc*) phases.^{171,172} An orthorhombic (*Pmn2₁*) phase has been observed for Cu₈GeS₆ SNCs.^{173,174} Also for these nanomaterials, LSPR features have been observed, besides interband electronic transitions.¹⁷¹

Cu-V-E compounds. Cu₃BiS₃. Although some alternative compositions of Cu-Bi-based chalcogenide nanomaterials are reported,¹⁷⁵ Cu₃BiS₃ is by far the most prevalent.¹⁷⁶ Cu₃BiS₃ is a p-type semiconductor with bulk bandgap of 1.5 eV that crystallizes in the wittichenite phase (orthorhombic, *P2₁2₁2₁*). SNCs of this material are generally of interest due to their strong light absorption capabilities in the visible. Yan *et al.* also observed the presence of a NIR-centered absorption feature, which was ascribed to an LSPR mode.¹⁷⁷ Similar features in the NIR are observed in the study by Chakraborty *et al.* and Ye *et al.*,¹⁷⁸ while Paul *et al.* even recorded both transversal and longitudinal LSPR modes in alloyed Cu₃BiS_{3-x}Se_x nanorods obtained through cation exchange from Bi₂S_{3-x}Se_x.¹⁷⁹ Note that Cu₃BiSe₃ can also be directly obtained, liked showed by Du *et al.*, who synthesized hydrophilic SNCs of this material in ethylene glycol and observed LSPR modes centered above 1200 nm.¹⁸⁰

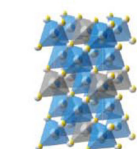
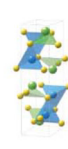
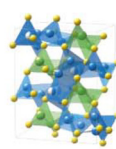
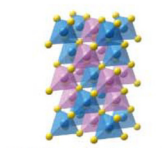
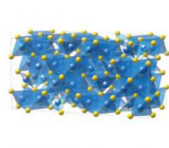
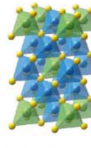
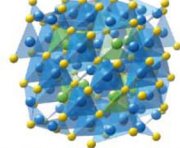
Cu-As-S and Cu-Sb-S. These families of materials include CuSbS₂ (chalcostibite, *Pnma*), Cu₃AsS₄ and Cu₃SbS₄ (luzonite and famatinite, *I42m*), Cu₁₂As₄S₁₃ and Cu₁₂Sb₄S₁₃ (tennantite and tetrahedrite, *I43m*).^{181–185} Also Cu₃SbS₃ (skinnerite, *P2₁/c*) has been reported (Fig. 8a).¹⁸⁶ These materials have various bandgaps from 1.0 to 2.0 eV and absorption coefficients in the 10⁵ cm⁻¹ in the visible range.^{185,187,188} Ternary chalcogenides of As and Sb are sometimes prepared in alloyed forms, given their good miscibility.¹⁸¹ Note that also Cu-As-Se and Cu-Sb-Se are reported, but we decided to omit them from this review given the limited number of reports on SNCs made of these materials.^{189,190} Aside from typical photo-induced interband electron transitions, it is worth mentioning the observation in Cu-Sb-S SNCs of a possible NIR LSPR by van Embden *et al.* (Fig. 8d and e)¹⁶⁷ and non-linear optical absorption, accompanied by ultrafast carrier dynamics, by Zhang *et al.*¹⁸⁸

Cu-TM-E. Cu-Fe-E. In these family of semiconductors,¹⁹¹ CuFeS₂ (Fig. 8f) and CuFeSe₂ are among the most studied. CuFeS₂ is often reported as an n-type semiconductor that crystallizes in the chalcopyrite phase with a bulk bandgap of 0.6 eV. The Cu-Fe-S system (often indicated as Cu_xFe_yS_z) is very rich in composition, and includes the orthorhombic Cu₅FeS₄ structure (bornite, *Pbca*) (Fig. 8a). CuFeSe₂ has a different, yet still tetragonal, crystal structure (eskebornite; *P42c*) with a bulk



(a) Crystal structures

● Cu ● Sn ● Sb ● Fe ● S

Cu₂SnS₃ - mohiteCu₃SnS₄ - kuramiteCuSbS₂ - chalcostibiteCu₃SbS₃ - skinneriteCuFeS₂ - chalcopyriteCu₃FeS₄ - borniteCu₃SbS₄ - famatiniteCu₁₂Sb₄S₁₃ - tetrahedrite

X-ray diffraction patterns

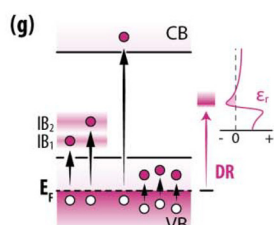
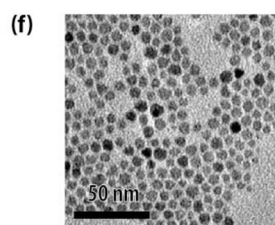
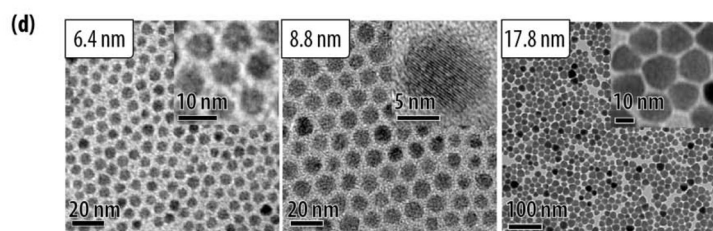
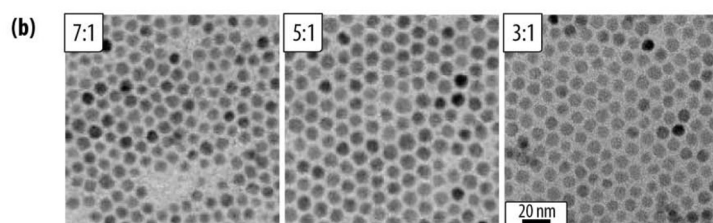
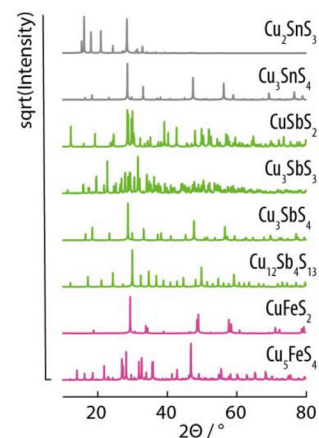


Fig. 8 Copper tin sulphide, copper antimony sulphide and copper iron sulphide SNCs as representatives of other ternary copper chalcogenides. (a) Unit cell and simulated XRPD patterns of the most common compositions reported for Cu–Sn–S, Cu–Sb–S, and Cu–Fe–S SNCs. (b) Representative TEM images and (c) extinction spectra of Cu₂SnS₃ SNCs prepared with different Cu:Sn ratios. Reproduced with permission from ref. 162, Copyright 2015, American Chemical Society. (d) Representative TEM images and (e) extinction spectra of Cu₁₂Sb₄S₁₃ SNCs of different sizes. Adapted with permission from ref. 167, Copyright 2013, American Chemical Society. (f) TEM image of CuFeS₂ SNCs. Reprinted with permission from ref. 168, Copyright 2021, American Chemical Society. (g) Electronic structure typical of Cu–Fe–S SNCs including intrabandgap states conducive of dielectric resonances. (h) Effect of oxidation and Fe removal on the optical properties of CuFeS₂ SNCs. Adapted with permission from ref. 169, Copyright 2021, American Chemical Society. (i) Photoluminescence of CuFeS₂/CdS core/shell SNCs of different sizes. Adapted with permission from ref. 170 Copyright 2016, American Chemical Society.

bandgap varying considerably from one report to another in the 0.16–0.6 eV range.^{192–194}

The presence of iron in the composition of these materials results in a unique electronic structure, since the empty 3d

orbitals of this metal introduce intermediate band (IB) states, which partake in electronic transitions (Fig. 8g). Moreover, this electronic structure leads to negative permittivity values within specific wavelength ranges that support so-called quasi-static



dielectric resonances (DRs).¹⁹⁵ In fact, various recent studies observed both DR and more common LSPR contributions in the extinction spectra of both chalcopyrite and bornite SNCs.^{168,196–198} These two resonant contributions are fundamentally different in nature (Fig. 8h) and their relative contributions to overall extinction can be modulated *via* controlled oxidation. In a follow-up study employing transient absorption spectroscopy, Kuszynski *et al.* found that the amount of Fe in the bornite SNCs correlates with the position of the IB, also identifying its central role in the relaxation dynamics of both hot holes and exciton carriers.¹⁹⁹

Remarkably, CuFeS₂ SNCs have been used also as cores for photoluminescent core/shell nanomaterials (Fig. 8i). Bhattacharyya and Pandey showed that, similarly to CuAlS₂, core-only CuFeS₂ SNCs are not luminescent, but the growth of a gradient CdS shell activates their luminescence.¹⁷⁰ The emission could be tuned in the visible range controlling the size of the core in the core/shell structure; however, the actual emission mechanism was not definitively interpreted. Vlasovets *et al.* also showed photoluminescence in Cu–Fe–S SNCs incorporated into a polymeric film,²⁰⁰ while CuAl_xFe_{1-x}S₂/ZnS were observed to have tunable emission in the visible-NIR range with large Stokes shifts.²⁰¹

Others. Other less-reported SNCs are made of materials including Cu₂MoS₄ (*P42m*),²⁰² CuCo₂S₄ (*Fd3m*),²⁰³ Cu₃VS₄ (*P43m*),²⁰⁴ CuCrS₂ (*R3m*), CuCrSe₂ (*R3m*),²⁰⁵ CuCr₂Se₄ (*Fd3m*),²⁰⁶ Cu₂MnS₂ (*Fd3m*). The optical properties of these SNCs are dominated by interband photon absorption.

Quaternary compounds. The definition of quaternary Cu-chalcogenide compounds can be blurry. In fact, many claimed quaternary materials can be considered alloys of other (ternary and/or binary compounds), with their electronic and optical properties falling in between the ones of the constituent materials. Moreover, the higher the number of the combined elements, the higher the chance of introducing disorder in the structure. Here, we consider as quaternary compounds those with well-defined stoichiometries. There are many possible such materials, and here we select a few that are well studied. Among those, quaternary semiconductors that include Zn²⁺ and/or Sn⁴⁺ are the most prevalent.

With its 1.0–1.5 eV direct bandgap,²⁰⁷ good stability, and high absorption coefficient,²⁰⁸ Cu₂ZnSnS₄ is used for the preparation of solar cells given the fact that it is solely made of Earth-abundant elements. In its bulk form, Cu₂ZnSnS₄ crystallizes in the kesterite (*I4*) and stannite (*I42m*) crystal phases (Fig. 9a), both closely related to chalcopyrite. Indeed, they are virtually distinguishable *via* XRD characterization alone, especially at the nanoscale. Kesterite is the most common Cu₂ZnSnS₄ crystal structure in SNCs. To that end, Collord and Hillhouse reported on kesterite SNCs with different compositions and sizes (Fig. 9b–e).²⁰⁹ Intriguingly, the authors found that the optical properties of larger SNCs are better described by a direct electronic transition, while small SNCs show an indirect bandgap with and sustain LSPR modes (Fig. 9c). A wurtzite phase has been observed too in SNCs. Initially reported by Mainz *et al.* as a transitional phase towards the for-



Fig. 9 Copper zinc tin sulphide SNCs as representatives of quaternary copper chalcogenides. (a) Unit cell and simulated XRPD patterns of the two most common compositions reported for Cu₂ZnSnS₄ SNCs. (b) TEM images of kesterite SNCs at different reaction times, alongside (c) the extinction spectra of SNCs obtained by precipitating smaller and larger SNC populations. Scale bars are 50 nm. Adapted with permission from ref. 209, Copyright 2015, American Chemical Society. (d) Band alignment of a direct (top) and indirect (bottom) bandgap semiconductor.

mation of stable kesterite SNCs,²¹⁰ some authors later reported on stable wurtzite SNCs likely obtained through the initial formation of Cu_{2-x}S and/or Cu₂SnS₃ seeds.^{211–213} The interest around wurtzite phase Cu₂ZnSnS₄ stems from the stoichiometric flexibility of this material – and hence high tunability of its optical properties – given the random distribution of the metals over the cation sublattice.

The sustainability of Cu₂ZnSnS₄ was further boosted in a recent study by Jones *et al.*, where the authors demonstrated scalability of the optimized synthesis of these SNCs, and gauged a >30% environmental impact reduction through life-cycle assessments of the materials.²¹⁴

Other quaternary SNCs include, among others, those made of Cu₁₀Co₂Sb₄S₁₃, Cu₁₀Zn₂Sb₄S₁₃, Cu₁₀Ni_{1.5}Sb₄S₁₃, Cu₂ZnGeS₄, Cu₂CdGeS₄, Cu₂CoGeS₄, and Cu₂ZnGeSe₄.^{215–220}

4.2.2. Synthesis. The preparation of multinary Cu-based chalcogenide SNCs has been explored through a multitude of approaches. Like in the case of binary Cu_{2-x}E SNCs, hot-injection and heat-up approaches are the most commonly employed ones, alongside solvo- and hydrothermal, sonochemical, microwave-driven, as well as microfluidic-based methods. In addition, to these “direct” methods, multi-step protocols have been developed, which are based on cation exchange processes. Notably, also most direct methods occur through the formation of SNCs of a specific composition fol-



lowed by slower cation exchange or cation incorporation processes to achieve the final material. However, for the sake of the discussion, these two approaches are presented independently (Fig. 10), depending on whether the cation exchange is kept by design clearly separated in the synthesis.

Direct approaches. The fundamental challenge in the synthesis of multinary compounds in direct approaches lies in the different nature of the elements that compose the material. In that regard, let us consider the case of CuInS_2 . The metal ions in this material are Cu^+ and In^{3+} , which display soft and hard Lewis acid behavior, respectively (Fig. 6). As such, they interact differently with different molecules, them being amines, carboxylate, thiolate, phosphines, or phosphine oxides. Hence, different ligands – at carefully controlled concentrations – should be introduced in the reaction environment to form *in situ* metal precursors with balanced reactivity. This is necessary to achieve the targeted multinary SNC rather than two populations of, in this case, Cu_{2-x}S and In_2S_3 SNCs.

An alternative is the use of single-source precursors that contain both metals in the right stoichiometric proportion (here, equimolar Cu and In amounts) alongside the chalcogen. The thermal decomposition of the precursor results in the simultaneous release of the elements in the right proportion, thus favoring the formation of the multinary compound. This approach seems promising on paper and thus attracted most of the attention in the early days of research on these SNCs. However, the time-consuming preparation of the precursors, their often air-sensitive nature, and our advanced understanding of wet chemistry approaches for nanoparticle synthesis have favored other, more user-friendly methods.

In hot-injection methods, metal precursors (salts and/or complexes) are mixed with a solution containing one or more species of ligands, the slurry is then heated to a target temperature generally $>100\text{ }^\circ\text{C}$ under vacuum to promote the precursor dissolution and *in situ* formation of metal complexes (e.g., metal carboxylates, metal thiolates). The temperature is then increased to a target value and a solution containing the chalcogen (in the form of elemental S, Se, or Te or a thermally labile chalcogen-containing molecule) is rapidly injected into

the metal solution. CuFeSe_2 nanocubes have been prepared by Wang *et al.* following this type of approach and using diphenyl selenide as chalcogen precursor.²²¹ The Dennis group reported the synthesis of bornite SNCs through a similar approach with OAm-S as the chalcogen solution, which was gradually injected over the course of 5 min at $180\text{ }^\circ\text{C}$.¹⁶⁹ The authors also observed that fast injection of the chalcogen solution in a glovebox led to a better controlled SNC size. Even quaternary SNCs can be prepared with this approach, as shown by Zhu *et al.*, who synthesized $\text{Cu}_2\text{CoGeS}_4$ SNCs, also demonstrating the possibility of replacing effortlessly Co for other TM such as Fe, Ni, Cd, and Mn.²¹⁹ Similar approaches are reported for most of the SNC families described above. Very recently, in an elegant study Chen *et al.* reported on a generalized strategy for the preparation of highly monodispersed quaternary wurtzite Cu-based selenide SNCs *via* injection of a diphenyl diselenide solution in a mixture of OAM, DDT, and the metal precursors.²²²

In a heat up approach, all the species are mixed together and the mixture is gradually heated to a target temperature. Lower-reactivity chalcogen species are generally used in this approach, to avoid uncontrolled SNC growth already at low temperature. In the case of CuInS_2 SNCs, the heat up approach has been perfected over the years: The preferred method entails the use of DDT in large excess – sometimes together with some OA and diluted with ODE – in the presence of CuI and In-acetate. Upon increasing the temperature, yellow metal-thiolates are formed above $100\text{ }^\circ\text{C}$, which start decomposing around $180\text{ }^\circ\text{C}$ to give ch-CIS QDs. Note that various Cu^+ precursors can be used, including Cu^{2+} salts. Use of divalent copper is also possible, despite the +1 oxidation state of copper in practically all multinary Cu-based semiconductors, thanks to the reducing capabilities of some molecules such as DDT or amines (e.g., OAm).²¹⁵ Moreover, the copper counterion can partake in the SNC growth, remaining incorporated in the crystal structure too.²²³ Quaternary, $\text{Cu}_2\text{ZnGeS}_4$ SNCs were prepared *via* a heat up approach by Fan *et al.*²¹⁷ Yet, as mentioned previously, the growth mechanism was shown to be step-wise, with an initial formation of $\text{Cu}_{1.75}\text{S}$ SNCs, followed by rapid incorporation of Zn^{2+} to yield a $\text{Cu}_x\text{Zn}_y\text{S}$ wurtzite phase, followed by gradual Ge^{4+} incorporation to give $\text{Cu}_2\text{ZnGeS}_4$ SNCs in an orthorhombic crystal phase.

Cation exchange approaches. The advantages of using cation exchange approaches are several: (i) a single batch of sacrificial (often Cu_{2-x}E) SNCs can be used to obtain wildly different compositions, (ii) crystal structures otherwise unachievable *via* direct methods can be stabilized for a given composition depending on the structure of parent SNCs, and (iii) upon controlling the reactivity of the exchanging ion solution, unique heterostructures can be unlocked.^{224,225} The clear drawback of this approach is its time-consuming and less-straightforward nature compared to, e.g., heat-up approaches. However, in the case of Cu-chalcogenide SNCs, cation exchange methods are possibly unrivalled in terms of flexibility and degree of control on the optical properties of the produced SNCs.²²⁶



Fig. 10 Scheme of direct synthesis and cation exchange approaches.



The kinetics of cation exchange processes are tuned by adjusting the composition of the cation exchange solution. Specifically, the ligand should promote extraction of the ion in the sacrificial SNC rather than too markedly stabilizing in solution the ion to be pushed in the structure. Because Cu_{2-x}E SNCs are most often the sacrificial species, thiols and phosphines (soft Lewis bases) are used to promote Cu^+ extraction. For example, Akkerman *et al.* showed that low chalcocite Cu_{2-x}S SNCs prepared *via* a thiourea-mediated heat up approach can be used as sacrificial species for preparing CuInS_2 SNCs.²²⁷ Incorporation of In^{3+} was achieved in the presence of TOP at 120 °C, achieving wurtzite CuInS_2 SNCs. An additional cation exchange process with Zn^{2+} , OAm, and octylamine (OcAm) led to incorporation of this divalent cation and a consequent enhancement of the photoluminescence intensity of $\text{CuInS}_2/\text{ZnS}$ with a gradient ZnS shell.²²⁷ Shamraienko *et al.* treated instead CuSe nanoplates with a Zn^{2+} and Sn^{2+} solution in OAm/OcAm in the presence of excess TOP, to obtain $\text{Cu}_2\text{ZnSnSe}_4$ nanoplates.⁷⁸ One should also consider the oxidation state of the cation to be exchanged, thus promoting ligand-free incorporation. To that end, Liu *et al.* showed that Sn^{2+} can be effectively incorporated in CuS SNCs to prepare Cu_3SnS_4 SNCs, while Sn^{4+} could only be incorporated to yield Cu_2SnS_3 SNCs if DDT was also added in the dispersion.²²⁸ The authors explained this effect considering the S–S reducing capabilities of Sn^{2+} , which promoted conversion of the crystal structure from covellite to kuramite.

4.2.3. Applications. Similarly to Cu_{2-x}E SNCs, the main application fields for multinary Cu-based SNCs are photovoltaics and optoelectronics, photocatalysis, and biomedicine.

Photovoltaics & optoelectronics. Photovoltaic cells with $\text{Cu}(\text{In}, \text{Ga})(\text{S}, \text{Se})_2$ absorbing layer are commercially available. These are known as CIGS solar cells and display some of the highest absorption coefficients among photovoltaic devices. Therefore, SNCs with such compositions are used to prepare inks for the fabrication of printable photovoltaic devices (Fig. 11a).^{229–231} Intriguingly, plasmonic CuInS_2 SNCs were shown to be more effective than their non-plasmonic counterpart for this application. This is because of the absorption enhancement ensured by the partial overlap between the interband absorption and LSPR bands (Fig. 7e).¹²⁴ Along the same lines, inks of $\text{Cu}_2\text{ZnSnS}_4$ (CZTS) and $\text{Cu}_2\text{ZnSnSe}_4$ (CZTSe) SNCs have been proposed to prepare printable devices,^{232–235} and other compositions have been explored too. As a matter of fact, the first reported synthesis of CZTS and CZTSe SNCs was reported by Guo *et al.*, who demonstrated the suitability of those SNCs dispersion for preparing solar cells.²³² Indeed, the interest around these materials stems from a theoretical 32% PCE value of thin-film solar cells based on kesterite.²³⁶ The film preparation techniques that guarantee better film homogeneity and thickness control include (multistep) spin coating, spray deposition, blade coating, and slot die coating followed by an annealing process to promote grain growth and reduce phenomena of carrier recombination and scattering at the SNC boundaries.^{214,233,237,238} Like in the case of classically grown CZTS films, selenization is sometimes performed to

increase the performance of SNC-derived CZTS films too.²³⁹ The (partial) conversion to CZTSe shifts the absorption to longer wavelengths (given that CZTS and CZTSe have bandgaps of 1.6 and 1.0 eV, respectively) and is reported to improve carrier mobility and concentration.²⁴⁰ Note that also SNC-based cells have been reported that do not require post-deposition annealing. To that end, Korala *et al.* highlighted the central role played by SNC surface chemistry to achieve effective, uniform passivation and hence increase the values of open-circuit potential (V_{oc}) by hundreds of mV.²⁴¹ Cu-based multinary SNCs are also used in QDSSCs as the photoanode material.^{242,243} Already in 2014 the Kamat group reported on the size-dependent efficiency of CuInS_2 SNC based QDSSCs, identifying a threshold size above which the PCE value dropped drastically (Fig. 11b).²⁴⁴ This cutoff was interpreted in light of a decreased charge separation achieved with narrowing bandgap. In addition, luminescent SNCs can be used to prepare light-converting layers in luminescent solar concentrators (LSCs). CuInS_2 -based SNCs are particularly suited for this type of device, since a large Stokes shift (Fig. 7c) is required for minimizing self-reabsorption over long distances (Fig. 11c).^{245–247} Core/shell QDs, gradient shells, and doping of the SNCs are all strategies that have been pursued to increase the efficiency of LSCs based on these SNCs. Moreover, Wu *et al.* recently showed that the Stokes shift can be further enhanced by dispersing the SNCs in thiol-ene polymers, reaching record PCE values of 1.36% for a $29 \times 29 \text{ cm}^2$ LSC device.²⁴⁸

Other optoelectronics uses of multinary Cu-based SNCs include photodetectors^{256–259} and light-emitting devices (LEDs).^{125,260,261} Regarding the latter application, the tunability of the photoluminescence of, *e.g.*, CuInS_2 -based SNCs makes them suitable for visible as well as NIR-emitting LEDs. Moreover, the SNCs can be used both as electroluminescent species²⁶¹ or as light-converting layer in blue-emitting LEDs to achieve white light emission.²⁶² However, one limitation in the use of these species in LED-based applications is the relatively broad emission spectrum, which limits the color purity achievable with these SNCs.

In a recent publication, Sachdeva *et al.* demonstrated that plasmonic CuInS_2 SNCs are also of interest in optoelectronics due to their ability of generating hot holes under sub-bandgap (NIR) excitation.¹³⁹ Intriguingly, when a CdS shell is grown on plasmonic CuInS_2 SNCs, the hot holes are delocalized to the shell material, where they can be extracted for further use.²⁶³

Photocatalysis. Multinary Cu-based chalcogenides have been used for photocatalysis above all in hydrogen evolution, pollutant degradation, and CO_2 reduction. Often, the SNCs are supported on a substrate that participates in the catalytic reaction forming a heterojunction. TiO_2 and $\text{g-C}_3\text{N}_4$ are often the substrate of choice, with the SNCs acting as sensitizers for the substrate material. Like in the case of biomedical applications (see below), systems based on $\text{CuIn}(\text{S}, \text{Se})_2$ SNCs are the most explored ones. This is mainly because they feature a favorable band alignment to achieve an efficient S-scheme (Step-scheme) heterojunction and can extend the absorption capa-





Fig. 11 Applications and case studies for multinary Cu-based SNCs. (a) Photo of a CuInSe_2 SNC ink and its deposition onto glass substrates (left), alongside a representative current-vs-potential curve under dark conditions (black) or light exposure (yellow) of a corresponding photovoltaic device (centre) and its ICPE spectrum (right). Adapted with permission from ref. 229, Copyright 2008, American Chemical Society. (b) Photo of ITO/ CuInS_2 SNCs photoanodes (top) and power conversion efficiency of the corresponding solar cells (bottom) for different SNC sizes. Adapted with permission from ref. 244, Copyright 2014, American Chemical Society. (c) Scheme showing the working principle of a LSC (top) alongside a photograph of a $30 \times 30 \times 0.7$ cm LSC waveguide of CuInS_2 SNC in PMMA. Reprinted with permission from ref. 245, Copyright 2019, John Wiley and Sons. (d) Scheme showing the approximate band alignment and electron transfer processes (black lines) in the graded $\text{CuInSe}_2/(\text{CuInSe}_2\text{S}_{1-x})/\text{CuInS}_2$ SNCs and TiO_2 photoanode. Adapted with permission from ref. 249, Copyright 2021, Elsevier. (e) TEM images of kesterite-wurtzite (red-green) $\text{Cu}_2\text{ZnSnS}_4$ single (top) and double (bottom) homojunctions, alongside a high-resolution image showing the atomic arrangement in the two polymorphs. Scale bars are 5 nm. Adapted with permission from ref. 250, Copyright 2022, Springer. (f) Scheme showing the energy levels of p- and n-type $\text{Cu}_2\text{ZnSnS}_4$ SNCs and their use in CO_2 photoreduction. Adapted with permission from ref. 251, Copyright 2022, American Chemical Society. (g) Use of $\text{CuInSe}_2/\text{ZnS}$ SNCs coated with phospholipids and antibodies for *in vivo* tumour imaging. From left to right: The emission spectra, functionalization process, study of cell interaction, and NIR images in nude mice are reported. Adapted with permission from ref. 252, Copyright 2020, Elsevier. (h) Study of the cytotoxicity profile of CuInS_2 -based SNCs. From left to right: photos of CuInS_2 and $\text{CuInS}_2/\text{ZnS}$ SNCs in simulated biological fluid at different days, degradation of CuInS_2 (CIS), Cu-In-Zn-S (CISZ), and $\text{CuInS}_2/\text{ZnS}$ (CIS/ZnS) SNCs monitored via optical absorption spectroscopy, organ-specific distribution (as % of In^{3+} initial dose) and summed In^{3+} content found in the body for CIS, CISZ, and CIS/ZnS. Adapted with permission from ref. 253, Copyright 2020, American Chemical Society. (i) Scheme representing the effect of Cd^{2+} occupying surface traps in CuInS_2 SNCs responsible for photoluminescence quenching. Adapted with permission from ref. 254, Copyright 2025, American Chemical Society. (j) Emission spectrum of CuInS_2 SNCs during a loading/unloading cycle in a diamond anvil cell, with the corresponding calibration dataset using peak maximum position as the manometric readout feature. Adapted with permission from ref. 255, Copyright 2023, John Wiley and Sons.



bilities of the substrate in the visible/NIR range. To that end, Li *et al.* reported on CuInSe₂/(CuInSe_xS_{1-x})₅/CuInS₂ core/shell/shell SNCs for photoelectrochemical hydrogen generation.²⁴⁹ The SNCs were prepared *via* cation exchange starting from CdSe/(CdSe_xS_{1-x})₅/CdS and then introducing in the lattice first Cu⁺, followed by In³⁺ (Fig. 11d). The resulting 1200-nm emitting SNCs were used for photo-electrochemical generation of H₂ in the presence of Na₂SO₃ with photocurrent density as high as ~4.5 mA cm⁻². A similar concept was recently proposed by Lee *et al.* in a study aimed at understanding the formation mechanisms of CuInS₂ SNCs *via in situ* small-angle X-ray scattering and *ex situ* XRPD and X-ray absorption.²⁶⁴ After investigation of the growth dynamics of the SNCs, the authors used core-only CuInS₂ SNCs to prepare a photoelectrochemical device that showed a current density >8 mA cm⁻². CuInS₂ SNCs were also employed by Zhang *et al.* together with g-C₃N₄ to prepare an S-scheme heterojunction for both H₂ evolution and pollutant (tetracycline) degradation, showing the versatility of this type of composite for photocatalysis applications.²⁶⁵ In a recent study, Yu and co-workers showcased the suitability of quaternary SNCs too for solar-to-hydrogen conversion.²⁵⁰ In the study, the authors developed a synthesis method to grow polytypic SNCs where a kesterite phase is epitaxially grown onto a wurtzite structure. Single- and double-homojunctions can be created at the single-SNC level (Fig. 11e), with double-homojunction systems featuring a better photocatalytic performance. Several other examples of similar S-scheme systems for H₂ evolution can be found in the literature that are based on other multinary copper chalcogenides, including, *e.g.*, CuCo₂S₄,^{266,267} Cu-Zn-S,²⁶⁸ Cu-Sb-S.²⁶⁹ Clearly the band alignment is pivotal to ensure effective photocatalytic performance. To that end, Chai *et al.* demonstrated that the amount of Sn²⁺ and Cu²⁺ doping in Cu₂ZnSnS₄ nanosheets could be controlled through the synthesis temperature.²⁵¹ An increased amount of this type of doping induced a change from p- to n-type semiconductor, and a change in the band structure and Fermi energy position. As a result, faster CO₂ photoreduction was achieved, with reported values of 48.14 and 25.04 μmol g⁻¹ h⁻¹ of CO and CH₄, respectively (Fig. 11f). While this is but a small collection of studies on the subject, the interested reader could find additional material on the subject in a dedicated review.²⁷⁰

Biomedicine. Multinary Cu-based SNCs are generally used in biomedicine because of their capabilities of acting as luminescent probe or for their ability to convert light into heat for photothermal applications.

Almost all the literature on photoluminescent multinary Cu-based SNCs for biomedical applications deals with CuInS₂ and in some cases CuInSe₂. Since the first demonstration of the use of CuInS₂/ZnS SNCs for *in vivo* lymph node imaging in 2010 by Pons *et al.*,²⁷¹ a plethora of publications on the biomedical use of these SNCs have seen the light. For instance, chitosan-coated CuInS₂/ZnS SNCs were used for both cell and *in vivo* imaging by Deng *et al.*²⁷² Lv *et al.* reported on pegylated phospholipid-coated CuInS₂/ZnS SNCs as a theranostic platform capable of performing as fluorescence and photoacoustic

contrast agent, while acting as photothermal and photodynamic therapy agent.²⁷³ CuInSe₂/ZnS SNCs were also used by Lian *et al.* as NIR-II (1000–1350 nm) emitting nanoparticles for developing bioassays targeting circulating tumor cells (Fig. 11g).²⁵² Moving from In-based and luminescent systems, CuFeS₂-Au nanohybrids were employed by Wen *et al.* as SERS sensors for the detection of lung cancer cells and the relative biomarkers.²⁷⁴

Crucially, the toxicity of this type of SNCs started being explored too, with early reports on both cells and small animal models (*C. elegans* and mice) showing negligible toxicity due to the high chemical stability of the SNCs.²⁷⁵ More recent studies by the Dennis group introduced a more critical view on the matter, finding that the presence of a ZnS shell is critical in imparting sufficient stability to minimize – yet not completely eliminate – CuInS₂ dissolution and hence release of toxic copper ions (Fig. 11h).²⁵³

Virtually any poorly- or non-luminescent, Cu-based multinary SNC can be used in light-to-heat conversion applications, given the large absorption coefficient displayed by these materials extending towards the NIR region. Indeed, SNCs made of CuFeS₂,^{276–278} Cu₃BiS₃,²⁷⁹ Cu-Sb-S,¹⁸⁵ Cu₃SnS₄,²⁸⁰ and Cu-Co-S²⁸¹ are just some of the nanocrystals explored in this context. Often, the light-to-heat conversion capabilities are harnessed for preparing a theranostic nanosystem. To that end, Hou *et al.* showed how poly(vinylpyrrolidone)-coated Cu-Sb-S SNCs can be used simultaneously for photothermal and photodynamic therapy, while acting as photoacoustic imaging contrast agents.¹⁸⁵ Yuan *et al.* also showed that CuCo₂S₄ SNCs have further potential as magnetic resonance contrast agents – due to the presence of Co in their composition – besides acting as effective photothermal agents, with HCE values >70%.²⁸¹ The presence of Bi in Tween-20-coated Cu₃BiS₃ SNCs was instead leveraged by Liu *et al.* for computed tomography as an additional imaging modality besides photoacoustic.²⁷⁹ Of course, plasmonic SNCs are better suited for these applications, since their NIR LSPR modes can be harnessed to use optical excitation falling at different wavelengths. By the same token, Cu-Fe-S SNCs – with their NIR absorption supported by IB transitions and LSPR modes – are particularly attractive in this context.

Others. Other applications of these SNCs include sensing of metal species, as well as luminescence thermometry and manometry. Alayeto *et al.* recently showcased the use of bright CuInS₂-based SNCs for selective Cd²⁺ sensing harnessing the photoluminescence enhancement induced by incorporation of the heavy metal ion in the SNC structure (Fig. 11i).²⁵⁴ Marin *et al.* prepared CuInS₂ SNC-polymer composites capable of acting as luminescent thermometers,¹²³ while Duda *et al.* applied a multimodal approach to analyse the luminescence signal of CuInS₂/ZnS SNCs for enhanced thermal sensing performance.²⁸² Aldaz-Caballero *et al.* also proposed CuInS₂ SNCs as luminescent manometers, showing tuning the SNC size and Cu-content impacted the sensing performance of the probes (Fig. 11j).²⁵⁵ It should also be highlighted that the company UbiQD is on the market with CuInS₂-



based solutions for optimization of the solar spectrum in greenhouses. The technology relies on light absorption and conversion in specific spectral windows to maximize crop yield.²⁸³

5. Halides

5.1. General considerations, structure, and optical properties

The research on ternary copper halides SNCs has mainly stemmed from the research on halide perovskites, in pursuit of less toxic and more stable nanomaterials that are equally performing in optical and optoelectronic applications. This family of materials is the most recent development in the research of Cu-based SNCs, with the seminal paper on Cs₃Cu₂I₅ and CsCu₂I₃ SNCs published by Cheng *et al.* in 2019.²⁸⁴ The most common stoichiometries encountered for these materials are ACuX₃, ACu₂X₃, A₂CuX₄, A₃Cu₂X₅ – where A = K, Rb, Cs and X = Cl, Br, I.¹⁴ These materials crystallize mainly into 0D or 1D structures,²⁸⁵ where Cu-containing clusters are individually distributed (0D) or arranged along ribbons (1D) interspersed by A⁺ ions. Given that the SNCs that can be found in the literature are mainly Cs-based, here we report on those materials specifically (Fig. 12).

In the CsCuX₃ family, CsCuCl₃ crystallizes in a chiral structure, in either *P*₆₂₂ (right-handed) or *P*₆₅₂ (left-handed) space group with octahedral [CuCl₆] units.^{291,292} Intriguingly, also CsCuBr₃ takes up a crystal structure with *P*₆₅₂ space group above 420 K, while the thermodynamically stable phase is orthorhombic (*C*222₁), with dimeric units of face-sharing [CuBr₆] octahedra. The reported bandgap values of these materials vary substantially between reports. CsCuCl₃ has an indirect bandgap (1.9–2.6 eV), while CsCuBr₃ has a somewhat wider, direct bandgap (2.63 eV). There is currently no report on CsCuI₃.

The three ternary halides of the CsCu₂X₃ family all crystallize in the *Cmcm* (or possibly *Pbnm*) space group, featuring tetrahedral [CuX₄] units sharing edges to form 1D chains along the *c*-axis.^{293,294} At room temperature, reported bandgaps are in the 3.0–3.75 eV range and direct in nature.^{295,296} Cs₂CuCl₄ and Cs₂CuBr₄ are found to crystallize mainly in the orthorhombic *Pnma* space group where Cu²⁺ resides in distorted tetrahedra, although mixed phases with Cs₂Cu(Cl_{1-x}Br_x)₄ at specific growth conditions can take up a tetragonal *I4/mmm* structure.²⁹⁷ No instances of Cs₂CuI₄ are reported in the literature.

The Cs₃Cu₂X₅ family (with bulk bandgaps in the 3.9–4.3 eV range) is of particular interest for the preparation of SNCs,²⁹⁸ since these semiconductors generally feature strong emission intensity (reaching PLQY ≈ 1), coupled with efficient light absorption capabilities (albeit values of absorption coefficient are not reported). This family of materials tends to crystallize in the orthorhombic *Pnma* space group where [Cu₂X₅]³⁻ dimers are separated by A⁺ ions to give a 0D structure. Cs₃Cu₂Cl₅ has been shown to take up also the *Cmcm* space group with two non-equivalent Cu⁺ ions disordered over two

tetrahedral sites to form 1D structures.^{288,299} Intriguingly, doping can be used to enhance the optical properties of these materials. To that end, Chen *et al.* demonstrated that Cs₃Cu₂I₅ SNCs can be doped with Rb⁺ – albeit only with an efficiency close to 20%. This process yielded wider band gap and higher PLQY (reaching almost unity), owing to the distortion induced by Rb⁺ substitution at one of the Cs⁺ sites.³⁰⁰ An overview of the structures, dimensionality, and cluster geometry of Cs–Cu–X – which are the most studied materials of this type – is reported in Fig. 12a.

The luminescence of these materials is characterized by the presence of strongly localized self-trapped excitons (STEs). STEs arise when there is strong electron–phonon coupling, and the electron–hole pair is strongly confined due to lattice distortions. Indeed, photoexcitation in these materials also results in a marked Jahn–Teller effect due to the change from Cu(I) 3d¹⁰ to Cu(II) 3d⁹ electronic configuration in tetrahedral coordination,²⁸⁷ which stabilizes the STE (Fig. 12c and d). Note that depending on the dimensionality of the structure, the excitons are localized over individual Cu-containing clusters (0D), or along ribbons (1D) and planes (2D) of clusters separated by A⁺ ions. The presence of STEs yields in these materials broad emission peaks,²⁸⁶ lifetimes up to 10s of microseconds, and marked Stokes shifts – at times larger than 1.5 eV (Fig. 12b).²⁹³

Much alike halide perovskites, a strategy employed for fine-tuning the optical properties of these SNCs is mixing different halides to obtain alloyed SNCs (Fig. 12e).^{288,301,302} This is a particularly successful strategy for Cs₃Cu₂X₅, where all the halides crystallize in the same structure. Another important feature of these materials is the relatively low charge mobility, owing to rather flat electronic bands (resulting in high effective mass).

The most common compositions synthesized at the nanoscale include CsCu₂X₃ and Cs₃Cu₂X₅, where copper is in its +1 oxidation state.

5.2. Synthesis

Synthesis protocols for the preparation of SNCs made of these materials take clear inspiration from the methods developed for halide perovskites. Specifically, the synthesis approach generally entails the rapid injection of Cs⁺ precursors in a Cu + X solution to induce controlled precipitation of the SNCs in the presence of ligands such as fatty acids and amines. In the first report on this family of materials, selective access to SNCs of CsCu₂I₃ or Cs₃Cu₂I₅ was achieved by controlling the injection temperature of Cs-oleate in a CuI, ODE, OA, and OAm mixture (Fig. 12f).²⁸⁴ Later, Chatterjee *et al.* reported on the preparation of both quasi-spherical SNCs and rod-shaped Cs₃Cu₂I₅ microcrystals using respectively a controlled precipitation method at low temperature (25 °C) and a hot injection protocol (130 °C).³⁰³ In both cases, the injection of the Cs⁺ precursor is followed by rapid quenching in ice to prevent further crystal growth. Park *et al.* demonstrated that in this type of reaction the presence of excess Cu-oleate is pivotal to obtain well-formed and passivated cubic SNCs, and identified MnI₂ as a



(a) Crystal structures



X-ray diffraction patterns

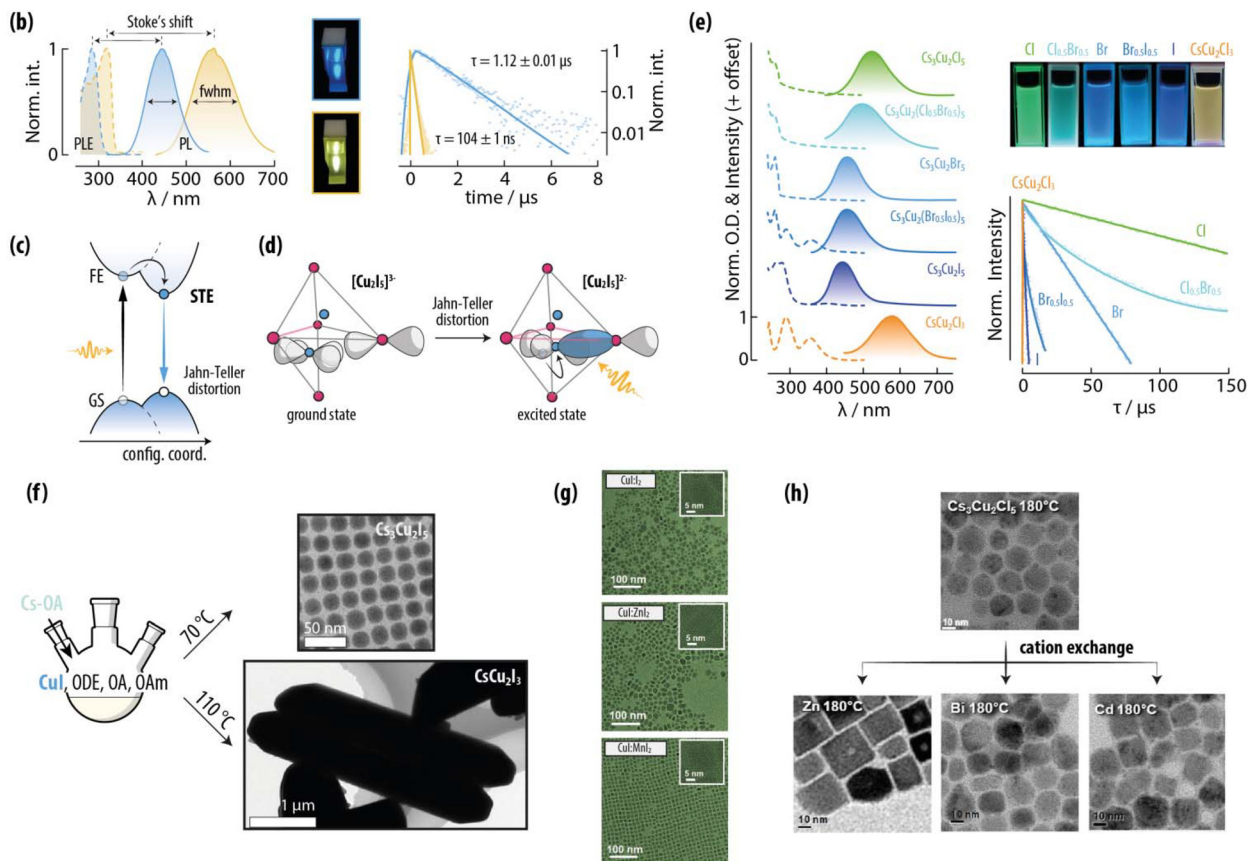


Fig. 12 Caesium copper halides as representative of copper halide SNCs. (a) Unit cell, copper cluster geometry, and simulated XRPD patterns of the most common compositions reported for Cs–Cu–X (where X = Cl, Br, I). (b) Absorption (dashed lines) and emission (solid lines) spectra of $\text{Cs}_3\text{Cu}_2\text{I}_5$ (blue) and CsCu_2I_3 (yellow), alongside a photo of the dispersion under UV excitation, and the corresponding photoluminescence decay curves. Adapted with permission from ref. 286, Copyright 2019, American Chemical Society. (c) The emission stems from a self-trapped exciton. (d) Representation of the Jahn–Teller distortion that occurs after excitation in $\text{Cs}_3\text{Cu}_2\text{I}_5$. Adapted with permission from ref. 287, Copyright 2025, Springer. (e) Absorption and emission spectra, photoluminescence decay curves, and photo under UV excitation of a series of alloyed $\text{Cs}_3\text{Cu}_2\text{X}_5$ SNCs, alongside CsCu_2Cl_3 SNCs. Adapted with permission from ref. 288, Copyright 2020, American Chemical Society. (f) Scheme of the synthesis protocol for the synthesis of $\text{Cs}_3\text{Cu}_2\text{I}_5$ SNCs or CsCu_2I_3 microcrystals, alongside the TEM images of the respective materials. Adapted with permission from ref. 284, Copyright 2019, John Wiley and Sons. (g) TEM images of $\text{Cs}_3\text{Cu}_2\text{I}_5$ SNCs obtained using different halide additives, showing optimized morphology and narrow size distribution in the presence of MnI_2 . Adapted with permission from ref. 289, Copyright 2022, American Chemical Society. (h) TEM images of parent $\text{Cs}_3\text{Cu}_2\text{Cl}_5$ synthesized at 180 °C and used as sacrificial species for cation exchange with Zn, Bi, and Cd. The TEM images of the SNCs obtained from the cation exchange procedures are reported in the bottom row. Adapted with permission from ref. 290, Copyright 2023, American Chemical Society.



better iodide additive compared to I_2 or ZnI_2 (Fig. 12g).²⁸⁹ The authors attributed this observation to the capability of Mn^{2+} to favor Cu-oleate from CuI in the presence of OA. This is the result of the catalytic capabilities of manganese, driven by the ion's multiple possible oxidation states. Moreover, Le *et al.* identified the length of the alkyl chain of the organic ligands (acid-amine pair) as a key parameter in controlling the size and morphology of $Cs_3Cu_2Br_5$ SNCs prepared *via* hot injection.³⁰⁴ In the same study, alloying with different halides (Cl and I) was also achieved to fine-tune the SNC optical properties.

Zhang *et al.* showed that the reverse approach can also be used, and they achieved the preparation of $Cs_3Cu_2I_5$ SNCs by injecting a CuI solution in a mixture of Cs_2CO_3 , ODE, OA, and OAm.³⁰⁵ These types of reactions can also be streamlined into continuous flow systems – as demonstrated by Urban and co-workers³⁰⁶ – for enhanced reproducibility and tunability of the SNC size and composition.

Intriguingly, Kwon *et al.* proposed the use of $Cs_3Cu_2Cl_5$ SNCs as sacrificial seeds for conducting cation exchange processes to access Cs_2ZnCl_4 , Cs_3BiCl_6 , and $CsCdCl_3$ SNCs (Fig. 12h).²⁹⁰ The structural changes (also accompanied by morphological variations) that cation exchange induced are such that further exchange with Cu^+ leads to SNCs of a material different from the starting one: $CsCu_2Cl_3$. Very recently, Rogach and co-workers also demonstrated that $CsPbBr_3$ - $Cs_3Cu_2Br_5$ nano-heterostructures can be prepared *via* epitaxial growth. These composite nanosystems show a pseudo-type II band alignment and offer an attractive platform to prepare heterojunctions with fine-tuned band architectures.³⁰⁷

5.3. Applications

Given the more recent history of SNCs with these compositions, examples of their applications are relatively limited. Yet, there are promising indications of the suitability of Cu-based halide SNCs in some specific fields. The attractiveness of these SNCs compared to, *e.g.*, Pb-based perovskite SNCs is generally attributed to the higher environment friendliness coupled with better stability – often imparted through careful doping strategies.

5.3.1. X-ray imaging. The scintillation capabilities of $CsCu_2X_3$ and $Cs_3Cu_2X_5$ SNCs have been fully embraced by the research community over the past five years. Examples of X-ray stimulated luminescence (*i.e.*, radioluminescence) have been reported for SNCs of $CsCu_2I_3$,³⁰⁸ $Cs_3Cu_2Cl_5$,^{309,310} and $Cs_3Cu_2I_5$.^{300,311,312} In fact, the X-ray absorption coefficient of these materials is on par with CsI:Tl (Fig. 13a) – a staple scintillator material. The high radioluminescence efficiency is supported by high PLQY values at times close to 90%. Very recently Chen *et al.* showed that doping with Rb^+ is an effective strategy to increase both the PLQY and the stability under continuous X-ray irradiation of $Cs_3Cu_2I_5$ SNCs.³⁰⁰ *Via* simulations, the authors identified doping of Rb at a Cs site as the most probable doping mechanism, which resulted in reduced electron–phonon coupling and overall suppressed

non-radiative recombination pathways. In the various studies mentioned above, the preparation of X-ray detection screens was achieved by depositing the SNCs on various surfaces, including paper, or incorporating them into a polymeric matrix to prepare standalone films. The attractiveness of these strategies resides in the easy processability of SNC-based films and the lack of highly toxic thallium in the scintillator.

5.3.2. Light emitting devices. $Cs_3Cu_2X_5$ SNCs have been used in LEDs both as electroluminescent material³¹⁴ and to prepare light conversion layers. To that end, in one of the first reports on this use of $Cs_3Cu_2X_5$ SNCs, Wang *et al.* prepared a blue LED with half-lifetime >100 h and an external quantum efficiency of 1.12% (Fig. 13b).²⁹⁸ Therein, the SNC layer acted as the electroluminescent element, showing good resistance to heating to >80 °C and storage in air for >30 days. The groups of Park³¹⁵ and Ye,³¹⁶ instead used respectively Mn-doped and Mg-doped $Cs_3Cu_2X_5$ SNCs as light converting layers in LEDs based on UV-emitting chips. Mn^{2+} introduced additional radiative relaxation pathways (through d–d transitions) that extended the color gamut of the SNCs. The resulting LED showed optical output powers on par with commercial white LEDs, yet a better spectral overlap (88.8 vs. 32.4%) with the solar spectrum. Mg^+ doping was instead leveraged to increase the PLQY of the SNCs, while also improving their stability when exposed to high temperatures and high-humidity environments.

5.3.3. Others. Other proposed applications for these nanomaterials include photovoltaic devices, photodetectors, photocatalysis, luminescent inks, and sensing. As mentioned above, the mobility of the charge carriers in these materials is relatively low, as evinced by the flat electronic band structure (Fig. 13a). Therefore, the performance of SNCs with these compositions is limited in applications that require fast charge separation. Still, some examples are reported in the literature. In that vein, Daliran *et al.* developed a mesoporous metal organic framework (MOF) incorporating $CsCu_2I_3$ SNCs as a composite photocatalyst.³¹⁷ The *in situ* growth within the MOF pores granted control over the SNC size and imparted stability to the halides. The authors demonstrated the suitability of the system to catalyze cycloaddition and photo-oxidation/Knoevenagel condensation cascade reactions under visible LED illumination (1 W, 32 000 lux), identifying the presence of multiple Lewis acid sites (Fe^{3+} , Zr^{4+} , Cu^+ , and Cs^+) as a key element.

$Cs_3Cu_2Cl_5$ ³¹³ and $Cs_3Cu_2I_5$ ³¹⁸ SNCs were also explored in photovoltaic devices for their UV harvesting capabilities. Albeit a large portion of UV light with wavelength shorter than 300 nm does not reach Earth surface due to ozone absorption, the interest around photovoltaic devices working in that spectral range is underpinned by their possible use in space power stations (Fig. 13c). To that end, Wang *et al.* showed that $Cs_3Cu_2Cl_5$ SNCs are suitable for LSC and as luminescent downshifting (LDS) layers³¹³ when incorporated in a siloxane polymer (PDMS), owing to the large Stokes shift – limiting reabsorption – and PLQY > 80%. A similar concept was behind the work of Liu *et al.*, who boosted the efficiency of Si photodetectors in the UV through the implementation of a downshifting $Cs_3Cu_2Cl_5$ SNC polymeric layer.³¹⁹



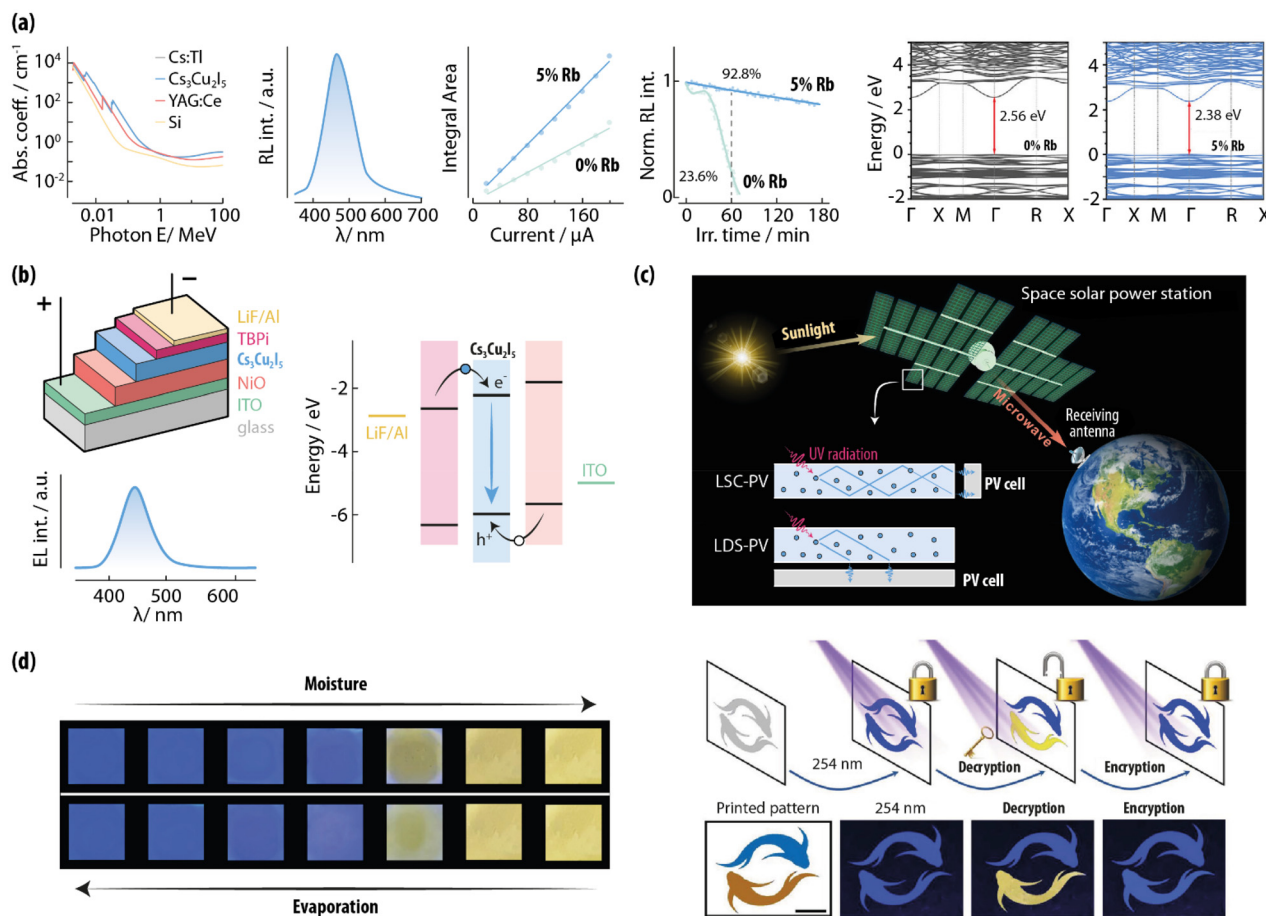


Fig. 13 Caesium copper halides as representative of copper halide SNCs. (a) Potential of Cs₃Cu₂I₅ SNCs for the preparation of scintillation devices. From left to right: X-ray absorption coefficient for Cs₃Cu₂I₅ SNCs compared to other materials used for scintillation, emission spectrum of Cs₃Cu₂I₅:Rb(5%) SNCs under X-ray excitation, the integrated area of undoped and 5% Rb doped SNCs as a function of driving voltage of the X-ray source, stability study under continuous X-ray irradiation, and band structure of the undoped and doped SNCs. Adapted with Permission from ref. 300, Copyright 2025, John Wiley and Sons. (b) Scheme of the optoelectronic device used as blue LED with a Cs₃Cu₂I₅ SNC emissive layer, alongside the electroluminescence spectrum of the LED, and the approximate band alignment of the various layers. Adapted with Permission from ref. 298, Copyright 2020, American Chemical Society. (c) Scheme showing the envisioned applicability in space power stations of luminescent solar concentrators containing Cs₃Cu₂I₅ SNCs. Adapted with permission from ref. 313, Copyright 2023, American Chemical Society. (d) Change of colour in an anticounterfeiting ink driven by the moisture-driven conversion between Cs₃Cu₂I₅ and CsCu₂I₃ SNCs. An example of the use of this ink for data encryption is also shown. Adapted with permission from ref. 305, Copyright 2021, John Wiley and Sons.

Given the brightness of these SNCs, it is unsurprising that preparation of luminescent inks has also been reported.³²⁰ Some of the anticounterfeiting strategies are multi-level in nature, such as the moisture-sensitive inks developed by Zhang *et al.* (Fig. 13d),³⁰⁵ Wang *et al.*,³²¹ and Feng *et al.*³²²

Sensing is yet another application field explored for these nanomaterials, with a report on chemical sensing of tebuconazole (TEB, a fungicide).³²³ The sensing mechanism is based on photoluminescence quenching of 1,10-phenanthroline-passivated Cs₃Cu₂I₅ SNCs, which is induced by TEB-induced aggregation in a concentration-dependent fashion.

6. Conclusions

We have reviewed the recent advances in the field of Cu-based semiconductor nanocrystals (SNCs) and their optical pro-

perties. Specifically, these nanomaterials exhibit photoluminescence and plasmonic properties, making them attractive for applications such as photovoltaics, lighting, sensing, photocatalysis, and X-ray imaging.

Of the three families discussed herein, pnictogenides and chalcogenides display stark similarities, featuring robustness against stoichiometry deviations and tunability of the optical properties depending on the composition of the material. Specifically, SNCs with compositions characterized by partially-filled valence bands display quasi-free holes that can support localized plasmonic resonances. A unique case is represented by CuFeE₂ SNCs, which support both localized plasmonic resonances and dielectric resonance modes arising from Fe 3d intrabandgap states. Photoluminescence in these families of SNCs stems instead from different mechanisms. For CuInE₂ and CuGaE₂ SNCs – the more explored compo-



sitions in pursuit of photoluminescence – the consensus in the community points to a free-electron-to-bound-hole as the most probable radiative recombination mechanism.

The third family, Cu-based halide SNCs, has markedly different properties. The main interest in them stems from an efficient photoluminescence, which makes these SNCs an attractive, less toxic alternative to halide perovskites. The crystal structure of Cu-based halide SNCs is characterized by Cu-halide clusters arranged in 0D, 1D, or 2D patterns where self-trapped excitons (STEs) get localized.

Despite the exciting optical features of Cu-based SNCs, our increased understanding of structure-properties relationships, and refinement of the synthesis protocols, several outstanding challenges and opportunities lie ahead in this field.

For starters, in the vast body of literature on Cu-based SNCs, some compositions are much less explored than others. This is particularly the case of nitrides and tellurides. Nitride SNCs are generally challenging to synthesize *via* colloidal chemistry, and Cu_3N is indeed the more synthetically-accessible composition in the family.³²⁴ The research on nitride SNCs is also being driven by the interest in the plasmonic properties of these materials.³²⁵ The identification of precursors that are safe to handle and with controlled reactivity, as well as alternative synthesis routes would certainly stir up interest for Cu-based nitride SNCs, allowing a more thorough study of their size- and morphology-dependent optical properties. The scarcity of tellurium-containing precursors and ligands is behind the reduced number of reports on Cu-based SNCs containing this chalcogen. In fact, tellurium chemistry is characterized by its malodorous and toxicity, as well as its instability and sensitivity to air and light, making this element less than appealing to work with. Yet tellurium is unique in that it has a diffuse electron density and shows increased spin-orbit coupling. One direct consequence is the reduced bandgap often featured by tellurides compared to sulfides and selenides, which results in more red-shifted optical absorption and photoluminescence. The electronic band structure is also modulated, resulting in tunable electronic properties, such as carrier mobility. As such, increased effort in the preparation of telluride Cu-based SNCs could lead to nanomaterials with uncommon, attractive optical properties.

Focusing on chalcogenides, phase transitions upon exposure to even mild temperatures and during storage are a concern in Cu-E SNCs. This intrinsic phase instability arises from the multiple polymorphs observed in Cu-E and the low diffusion barrier in these materials. Structural changes alter electronic/optical properties, hence reducing the credibility of technological solutions based on Cu-E SNCs. A solution to this issue might require a multipronged approach to obtain new, more complex Cu-E SNCs with minimized ion diffusion (*e.g.*, *via* doping or alloying) and resistance to surface oxidation (*e.g.*, through the growth of core/shell nanostructures). Future research aimed at understanding phase transitions in these nanomaterials at the atomic level and, thus, guiding the design of stable Cu-E SNCs could yield results with tremendous technological impact.

Cu-based halide SNCs are generally considered more stable than, *e.g.*, Pb-based halide perovskite SNCs; however, their long-term stability remains a concern. In these materials, exposure to air and/or moisture can lead to surface oxidation, mainly driven by the redox chemistry of the Cu(I)/Cu(II) system. It is conceivable that strategies recently developed to increase the surface stability of halide perovskites can be applied to these Cu-based SNCs, such as the use of multidentate ligands³²⁶ and shell formation with inorganic materials.³²⁷ Partial substitution of Cu(I) for other metals, *e.g.*, Ag(I), that are less readily oxidized can also improve stability – though inducing changes in the optical and electronic properties. Additionally, the soft crystal structure of halides (extensively discussed in halide perovskites³²⁸) leads to substantial anion migration within the lattice, which can adversely affect the stability of optoelectronic devices. Strategies used to limit ion migration in Pb-halide perovskites are expected to be similarly effective in Cu-based systems. These include defect engineering *via* aliovalent element doping and tuning of metal-halogen bond strength.³²⁹ The former could be achieved *via* interstitial doping with ions such as Ca^{2+} and Ba^{2+} to increase the activation energy for halide migration. The latter could be accomplished by partial cation substitution of Cu(I) with Au(I), which forms a stronger metal-halide bond.

The relatively low mobility of charge carriers in Cu-based halide SNCs poses instead major limitations to their use in applications that require efficient charge extraction, such as photocatalysis or photoelectronic devices. An exciting avenue for future research is the identification of approaches to modify the electron structure of the material to modulate the effective mass of the charge carriers, including doping and strain engineering.

The preparation of complex heterostructures is another approach expected to yield materials with a unique combination of optical and electronic properties. For example, plasmonic and luminescent properties can be achieved in a single heterostructure based completely on copper. In addition, more effective carrier extraction can be obtained by modulating the electronic band alignment in multiple materials in contact within the same heterostructure. Cation (and anion) exchange strategies, as well as controlled epitaxial growth, are ideal tools to achieve such structures. However, the ultimate success of these sophisticated structures remains contingent upon resolving the stability issues mentioned above.

To conclude, the technological relevance – either already validated and exploited or just postulated – of Cu-based SNCs is expected to drive further research on these nanomaterials. As such, we believe that the future of this field still has plenty to offer chemists, physicists, materials scientists, and engineers alike. We hope that more of the systems proposed as proof-of-concept would move towards large-scale implementation and on the market, where they can benefit society for a brighter future.



Author contributions

R. M. contributed to the conceptualization, funding acquisition, visualization, and writing of the original draft. L. V. B. and P. C. contributed to reviewing and editing. P. C. also contributed to data visualization.

Conflicts of interest

There are no conflicts to declare.

Data availability

No primary research results, software or code have been included and no new data were generated or analysed as part of this review.

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References

- <https://www.copper.org/education/history/> accessed on 31/10/2025.
- <https://sites.dartmouth.edu/toxmetal/more-metals/copper-an-ancient-metal/> accessed on 31/10/2025.
- G. Grass, C. Rensing and M. Solioz, *Appl. Environ. Microbiol.*, 2011, **77**, 1541–1547.
- G. Fiske Johnson, *Agric. Hist.*, 1935, **9**, 67–79.
- M. Gao, R. Ye, W. Shen and B. Xu, *Org. Biomol. Chem.*, 2018, **16**, 2602–2618.
- B. K. Meyer, A. Polity, D. Reppin, M. Becker, P. Hering, P. J. Klar, T. Sander, C. Reindl, J. Benz, M. Eickhoff, C. Heiliger, M. Heinemann, J. Bläsing, A. Krost, S. Shokovets, C. Müller and C. Ronning, *Phys. Status Solidi B*, 2012, **249**, 1487–1509.
- K. E. Knowles, K. H. Hartstein, T. B. Kilburn, A. Marchioro, H. D. Nelson, P. J. Whitham and D. R. Gamelin, *Chem. Rev.*, 2016, **116**, 10820–10851.
- M. A. Malik, P. O'Brien and N. Revaprasadu, *Adv. Mater.*, 1999, **11**, 1441–1444.
- Y. Zhao, H. Pan, Y. Lou, X. Qiu, J. Zhu and C. Burda, *J. Am. Chem. Soc.*, 2009, **131**, 4253–4261.
- S. Sun, P. Li, S. Liang and Z. Yang, *Nanoscale*, 2017, **9**, 11357–11404.
- W. van der Stam, A. C. Berends and C. de Mello Donega, *ChemPhysChem*, 2016, **17**, 559–581.
- L. Wu, Y. Li, G.-Q. Liu and S.-H. Yu, *Chem. Soc. Rev.*, 2024, **53**, 9832–9873.
- J. Qu, S. Xu, H. Shao, P. Xia, C. Lu, C. Wang and D. Ban, *J. Mater. Chem. C*, 2023, **11**, 6260–6275.
- Z. Ma, X. Ji, S. Lin, X. Chen, D. Wu, X. Li, Y. Zhang, C. Shan, Z. Shi and X. Fang, *Adv. Mater.*, 2023, **35**, 2300731.
- B. Chen, W. Zheng, F. Chun, X. Xu, Q. Zhao and F. Wang, *Chem. Soc. Rev.*, 2023, **52**, 8374–8409.
- J. Ma, N. Li, J. Wang, Z. Liu, Y. Han and Y. Zeng, *Exploration*, 2023, **3**, 20220161.
- B. Bhattacharyya, C. Balischiwski, C. Pacholski, A. Pandey, I. Bald and A. Taubert, *Adv. Opt. Mater.*, 2023, **11**, 2202411.
- C. Coughlan, M. Ibáñez, O. Dobrozhan, A. Singh, A. Cabot and K. M. Ryan, *Chem. Rev.*, 2017, **117**, 5865–6109.
- J. A. Faucheaux, A. L. D. Stanton and P. K. Jain, *J. Phys. Chem. Lett.*, 2014, **5**, 976–985.
- J. B. Khurgin, *Faraday Discuss.*, 2019, **214**, 35–58.
- M. Pelton and G. W. Bryant, *Introduction to metal-nanoparticle plasmonics*, John Wiley & Sons Inc, Hoboken, NJ, Online-Ausg., 2013.
- A. Paściak, A. Pilch-Wróbel, Ł. Marciniak, P. J. Schuck and A. Bednarkiewicz, *ACS Appl. Mater. Interfaces*, 2021, **13**, 44556–44567.
- D. K. Roper, W. Ahn and M. Hoepfner, *J. Phys. Chem. C*, 2007, **111**, 3636–3641.
- A. Paściak, R. Marin, L. Abiven, A. Pilch-Wróbel, M. Misiak, W. Xu, K. Prorok, O. Bezkrvnyi, Ł. Marciniak, C. Chanéac, F. Gazeau, R. Bazzi, S. Roux, B. Viana, V.-P. Lehto, D. Jaque and A. Bednarkiewicz, *ACS Appl. Mater. Interfaces*, 2022, **14**, 33555–33566.
- S. E. Braslavsky, *Pure Appl. Chem.*, 2007, **79**, 293–465.
- S. Ye, C. Ding and C. Li, *Advances in Inorganic Chemistry*, Elsevier, 2019, vol. 74, pp. 3–59.
- G. H. Wannier, *Phys. Rev.*, 1937, **52**, 191–197.
- A. F. Harper, M. L. Evans and A. J. Morris, *Chem. Mater.*, 2020, **32**, 6629–6639.
- Z. Liu, H. Mu, S. Xiao, R. Wang, Z. Wang, W. Wang, Y. Wang, X. Zhu, K. Lu, H. Zhang, S. Lee, Q. Bao and W. Ma, *Adv. Mater.*, 2016, **28**, 3535–3542.
- L. De Trizio, A. Figuerola, L. Manna, A. Genovese, C. George, R. Brescia, Z. Saghi, R. Simonutti, M. Van Huis and A. Falqui, *ACS Nano*, 2012, **6**, 32–41.
- Y. Liu, J. Wu, Y. Jin, W. Zhen, Y. Wang, J. Liu, L. Jin, S. Zhang, Y. Zhao, S. Song, Y. Yang and H. Zhang, *Adv. Funct. Mater.*, 2019, **29**, 1904678.
- A. Crovetto, T. Unold and A. Zakutayev, *Chem. Mater.*, 2023, **35**, 1259–1272.



- 33 X. Yue, S. Yi, R. Wang, Z. Zhang and S. Qiu, *Nanoscale*, 2016, **8**, 17516–17523.
- 34 G. Manna, R. Bose and N. Pradhan, *Angew. Chem., Int. Ed.*, 2013, **52**, 6762–6766.
- 35 D. Bandyopadhyay, A. Upcher, L. Houben, V. Ezersky, S. Ganguly, S. Ghosh and M. Bar Sadan, *Small Struct.*, 2025, **6**, 2500338.
- 36 D. Stone, X. Li, T. Naor, J. Dai, S. Remennik and U. Banin, *Chem. Mater.*, 2023, **35**, 10594–10605.
- 37 A. G. Rachkov, K. Chalek, H. Yin, M. Xu, G. P. Holland and A. M. Schimpf, *ACS Nano*, 2024, **18**, 5282–5296.
- 38 F. Qi, Y. Chang, R. Zheng, X. Wu, Y. Wu, B. Li, T. Sun, P. Wang, H. Zhang and H. Zhang, *ACS Biomater. Sci. Eng.*, 2021, **7**, 2745–2754.
- 39 A. G. Rachkov and A. M. Schimpf, *Chem. Mater.*, 2021, **33**, 1394–1406.
- 40 T. Sun, Y. Wang, W. Yu, Y. Wang, Z. Dai, Z. Liu, B. N. Shivananju, Y. Zhang, K. Fu, B. Shabbir, W. Ma, S. Li and Q. Bao, *Small*, 2017, **13**, 1701881.
- 41 C. A. Downes, N. J. Libretto, A. E. Harman-Ware, R. M. Happs, D. A. Ruddy, F. G. Baddour, J. R. Ferrell III, S. E. Habas and J. A. Schaidle, *ACS Appl. Energy Mater.*, 2020, **3**, 10435–10446.
- 42 J. G. Zhao, S. J. You, L. X. Yang and C. Q. Jin, *Solid State Commun.*, 2010, **150**, 1521–1524.
- 43 L. Yang, L. Zhang, Y. Li, B.-H. Lee, J. Kim, H. S. Lee, J. Bok, Y. Ma, W. Zhou, D. Yuan, A.-L. Wang, M. S. Bootharaju, H. Zhang, T. Hyeon and J. Chen, *J. Am. Chem. Soc.*, 2024, **146**, 12556–12564.
- 44 R. Kadzutu-Sithole, L. F. E. Machogo-Phao, T. Kolokoto, M. Zimuwandeyi, S. S. Gqoba, K. P. Mubiayi, M. J. Moloto, J. Van Wyk and N. Moloto, *RSC Adv.*, 2020, **10**, 34231–34246.
- 45 K. Matsuzaki, T. Okazaki, Y.-S. Lee, H. Hosono and T. Susaki, *Appl. Phys. Lett.*, 2014, **105**, 222102.
- 46 S. Majumder, M. Ohishi, K. Saito, Q. Guo, M. A. M. Patwary and T. Tanaka, *Mater. Sci. Semicond. Process.*, 2024, **182**, 108702.
- 47 C. Navío, M. J. Capitán, J. Álvarez, F. Yndurain and R. Miranda, *Phys. Rev. B:Condens. Matter Mater. Phys.*, 2007, **76**, 085105.
- 48 T. Nakamura, H. Hayashi and T. Ebina, *J. Nanopart. Res.*, 2014, **16**, 2699.
- 49 A. Stróżecka, J. Li, R. Schürmann, G. Schulze, M. Corso, F. Schulz, C. Lotze, S. Sadewasser, K. J. Franke and J. I. Pascual, *Phys. Rev. B:Condens. Matter Mater. Phys.*, 2014, **90**, 195420.
- 50 D. Zheng, Y.-Q. He, H. Ou-Yang, J.-D. Zhang, G. Zhang and S.-K. Han, *Nano Lett.*, 2024, **24**, 8887–8893.
- 51 M. Parvizian, A. Duràn Balsa, R. Pokratath, C. Kalha, S. Lee, D. Van den Eynden, M. Ibáñez, A. Regoutz and J. De Roo, *Angew. Chem., Int. Ed.*, 2022, **61**, e202207013.
- 52 R. K. Sithole, L. F. E. Machogo, M. J. Moloto, S. S. Gqoba, K. P. Mubiayi, J. Van Wyk and N. Moloto, *J. Photochem. Photobiol., A*, 2020, **397**, 112577.
- 53 Q. Guo, F. Liang, Z. Sun, Y. Wang, X.-B. Li, S.-G. Xia, Z. C. Zhang, L. Huang and L.-Z. Wu, *J. Mater. Chem. A*, 2020, **8**, 22601–22606.
- 54 D. Barman, S. Paul, S. Ghosh and S. K. De, *ACS Appl. Nano Mater.*, 2019, **2**, 5009–5019.
- 55 U. Dasgupta, M. Ghosh, R. Gangopadhyay, D. D. Klyamer, P. Venkatesan, T. Basova, R.-A. Doong and A. D. Chowdhury, *Chem. Eng. J.*, 2025, **521**, 166394.
- 56 J. M. Luther, P. K. Jain, T. Ewers and A. P. Alivisatos, *Nat. Mater.*, 2011, **10**, 361–366.
- 57 Y. Xie, L. Carbone, C. Nobile, V. Grillo, S. D'Agostino, F. Della Sala, C. Giannini, D. Altamura, C. Oelsner, C. Krysch and P. D. Cozzoli, *ACS Nano*, 2013, **7**, 7352–7369.
- 58 Y. Xie, A. Riedinger, M. Prato, A. Casu, A. Genovese, P. Guardia, S. Sottini, C. Sangregorio, K. Miszta, S. Ghosh, T. Pellegrino and L. Manna, *J. Am. Chem. Soc.*, 2013, **135**, 17630–17637.
- 59 R. M. Córdova-Castro, M. Casavola, M. Van Schilfgaarde, A. V. Krasavin, M. A. Green, D. Richards and A. V. Zayats, *ACS Nano*, 2019, **13**, 6550–6560.
- 60 K. Elibol, F. Davoodi, U. Parekh, M. Taleb, S. Scheel, M. Burghard, P. A. Van Aken, C. Klinke, N. Talebi and R. Lesyuk, *Adv. Opt. Mater.*, 2025, **13**, 2402965.
- 61 Y. Liu, M. Liu and M. T. Swihart, *Chem. Mater.*, 2017, **29**, 4783–4791.
- 62 Y. Wu, C. Wadia, W. Ma, B. Sadtler and A. P. Alivisatos, *Nano Lett.*, 2008, **8**, 2551–2555.
- 63 Z. Liu, X. Liu, Y. Du, J. Ren and X. Qu, *ACS Nano*, 2015, **9**, 10335–10346.
- 64 D. Zhou, D. Liu, W. Xu, Z. Yin, X. Chen, P. Zhou, S. Cui, Z. Chen and H. Song, *ACS Nano*, 2016, **10**, 5169–5179.
- 65 I. Kriegel, C. Jiang, J. Rodríguez-Fernández, R. D. Schaller, D. V. Talapin, E. da Como and J. Feldmann, *J. Am. Chem. Soc.*, 2012, **134**, 1583–1590.
- 66 K. H. Hartstein, C. K. Brozek, S. O. M. Hinterding and D. R. Gamelin, *J. Am. Chem. Soc.*, 2018, **140**, 3434–3442.
- 67 R. Alam, M. Labine, C. J. Karwacki and P. V. Kamat, *ACS Nano*, 2016, **10**, 2880–2886.
- 68 D. Dorfs, T. Härtling, K. Miszta, N. C. Bigall, M. R. Kim, A. Genovese, A. Falqui, M. Povia and L. Manna, *J. Am. Chem. Soc.*, 2011, **133**, 11175–11180.
- 69 V. B. Llorente, V. M. Dzhagan, N. Gaponik, R. A. Iglesias, D. R. T. Zahn and V. Lesnyak, *J. Phys. Chem. C*, 2017, **121**, 18244–18253.
- 70 W. van der Stam, S. Gudjonsdottir, W. H. Evers and A. J. Houtepen, *J. Am. Chem. Soc.*, 2017, **139**, 13208–13217.
- 71 R. D. Heyding and R. M. Murray, *Can. J. Chem.*, 1976, **54**, 841–848.
- 72 A. L. N. Stevels and F. Jellinek, *Recl. Trav. Chim. Pays-Bas*, 1971, **90**, 273–283.
- 73 R. W. Lord, J. Fanghanel, C. F. Holder, I. Dabo and R. E. Schaak, *Chem. Mater.*, 2020, **32**, 10227–10234.
- 74 G. Gariano, V. Lesnyak, R. Brescia, G. Bertoni, Z. Dang, R. Gaspari, L. De Trizio and L. Manna, *J. Am. Chem. Soc.*, 2017, **139**, 9583–9590.



- 75 S. C. Riha, D. C. Johnson and A. L. Prieto, *J. Am. Chem. Soc.*, 2011, **133**, 1383–1390.
- 76 L. E. Marbella, X. Y. Gan, D. C. Kaseman and J. E. Millstone, *Nano Lett.*, 2017, **17**, 2414–2419.
- 77 V. Lesnyak, R. Brescia, G. C. Messina and L. Manna, *J. Am. Chem. Soc.*, 2015, **137**, 9315–9323.
- 78 V. Shamraienko, D. Spittel, R. Hübner, M. Samadi Khoshkhoo, N. Weiß, M. Georgi, K. B. L. Borchert, D. Schwarz, V. Lesnyak and A. Eychmüller, *J. Mater. Chem. C*, 2021, **9**, 16523–16535.
- 79 T. Willhammar, K. Sentosun, S. Mourdikoudis, B. Goris, M. Kurttepel, M. Bercx, D. Lamoen, B. Partoens, I. Pastoriza-Santos, J. Pérez-Juste, L. M. Liz-Marzán, S. Bals and G. Van Tendeloo, *Nat. Commun.*, 2017, **8**, 14925.
- 80 I. Kriegel, J. Rodríguez-Fernández, A. Wisnet, H. Zhang, C. Waurisch, A. Eychmüller, A. Dubavik, A. O. Govorov and J. Feldmann, *ACS Nano*, 2013, **7**, 4367–4377.
- 81 Y. Wu, F. Juan, B. Wang, S. Sun, J. Jia, H. Wei and B. Cao, *Optik*, 2022, **271**, 170229.
- 82 H. Li, R. Brescia, M. Povia, M. Prato, G. Bertoni, L. Manna and I. Moreels, *J. Am. Chem. Soc.*, 2013, **135**, 12270–12278.
- 83 P. L. Saldanha, R. Brescia, M. Prato, H. Li, M. Povia, L. Manna and V. Lesnyak, *Chem. Mater.*, 2014, **26**, 1442–1449.
- 84 Z. Xu, C. Shen, Y. Hou, H. Gao and S. Sun, *Chem. Mater.*, 2009, **21**, 1778–1780.
- 85 J. W. Thomson, K. Nagashima, P. M. Macdonald and G. A. Ozin, *J. Am. Chem. Soc.*, 2011, **133**, 5036–5041.
- 86 G. Fonder, F. Laffineur, J. Delhalle and Z. Mekhalif, *J. Colloid Interface Sci.*, 2008, **326**, 333–338.
- 87 R. Marin, J. Lifante, L. V. Besteiro, Z. Wang, A. O. Govorov, F. Rivero, F. Alfonso, F. Sanz-Rodríguez and D. Jaque, *Adv. Healthcare Mater.*, 2020, **9**, 1901627.
- 88 W. P. Lim, C. T. Wong, S. L. Ang, H. Y. Low and W. S. Chin, *Chem. Mater.*, 2006, **18**, 6170–6177.
- 89 T. H. Larsen, M. Sigman, A. Ghezlbash, R. C. Doty and B. A. Korgel, *J. Am. Chem. Soc.*, 2003, **125**, 5638–5639.
- 90 S.-W. Hsu, C. Ngo and A. R. Tao, *Nano Lett.*, 2014, **14**, 2372–2380.
- 91 G. Xiao, J. Ning, Z. Liu, Y. Sui, Y. Wang, Q. Dong, W. Tian, B. Liu, G. Zou and B. Zou, *CrystEngComm*, 2012, **14**, 2139–2144.
- 92 W. Li, R. Zamani, P. Rivera Gil, B. Pelaz, M. Ibáñez, D. Cadavid, A. Shavel, R. A. Alvarez-Puebla, W. J. Parak, J. Arbiol and A. Cabot, *J. Am. Chem. Soc.*, 2013, **135**, 7098–7101.
- 93 A. Dey, S. Noyel Victoria, M. S. Gaikwad and A. Sao, *Mater. Lett.*, 2024, **361**, 136093.
- 94 D. K. Becerra-Paniagua, E. B. Díaz-Cruz, A. Baray-Calderón, A. R. Garcia-Angelmo, E. Regalado-Pérez, M. Del Pilar Rodríguez-Torres and C. Martínez-Alonso, *J. Mater. Sci.:Mater. Electron.*, 2022, **33**, 22631–22667.
- 95 M. R. Kim, H. A. Hafez, X. Chai, L. V. Besteiro, L. Tan, T. Ozaki, A. O. Govorov, R. Izquierdo and D. Ma, *Nanoscale*, 2016, **8**, 12946–12957.
- 96 C. H. M. Van Oversteeg, F. E. Oropeza, J. P. Hofmann, E. J. M. Hensen, P. E. De Jongh and C. De Mello Donega, *Chem. Mater.*, 2019, **31**, 541–552.
- 97 R. Marin, A. Skripka, L. V. Besteiro, A. Benayas, Z. Wang, A. O. Govorov, P. Canton and F. Vetrone, *Small*, 2018, **14**, 1803282.
- 98 X. Y. Gan, S. E. Crawford, E. A. Eikey, R. Sen, J. R. Killinger and J. E. Millstone, *J. Phys. Chem. C*, 2020, **124**, 4747–4754.
- 99 W. Gao, Y. Sun, M. Cai, Y. Zhao, W. Cao, Z. Liu, G. Cui and B. Tang, *Nat. Commun.*, 2018, **9**, 231.
- 100 Y. Zhao and C. Burda, *Energy Environ. Sci.*, 2012, **5**, 5564–5576.
- 101 D. C. Reynolds, G. Leies, L. L. Antes and R. E. Marburger, *Phys. Rev.*, 1954, **96**, 533–534.
- 102 L. Liu, B. Zhou, L. Deng, W. Fu, J. Zhang, M. Wu, W. Zhang, B. Zou and H. Zhong, *J. Phys. Chem. C*, 2014, **118**, 26964–26972.
- 103 J. Tirado, C. Roldán-Carmona, F. A. Muñoz-Guerrero, G. Bonilla-Arboleda, M. Ralaiarisoa, G. Grancini, V. I. E. Queloz, N. Koch, M. K. Nazeeruddin and F. Jaramillo, *Appl. Surf. Sci.*, 2019, **478**, 607–614.
- 104 S. Q. Lie, D. M. Wang, M. X. Gao and C. Z. Huang, *Nanoscale*, 2014, **6**, 10289–10296.
- 105 M. Sun, X. Fu, K. Chen and H. Wang, *ACS Appl. Mater. Interfaces*, 2020, **12**, 46146–46161.
- 106 N. Hoàng Ly, V. Thi Huong, B. Van Duc, P.-D. Nguyen, T. M. Aminabhavi, Y. Vasseghian and S.-W. Joo, *Chem. Eng. J.*, 2023, **462**, 142135.
- 107 C. Li and J. Wang, *Small*, 2024, **20**, 2404798.
- 108 W. Poon, Y.-N. Zhang, B. Ouyang, B. R. Kingston, J. L. Y. Wu, S. Wilhelm and W. C. W. Chan, *ACS Nano*, 2019, **13**, 5785–5798.
- 109 Y. Li, M. Su, T. Yan, Z. Wang and J. Zhang, *ACS Appl. Bio Mater.*, 2023, **6**, 2468–2476.
- 110 Y. Yang, C. Xu, S. Xu, Y. Li, K. Chen, T. Yang, J. Bao, Y. Xu, J. Chen, C. Mao, L. Chen and W. Sun, *Biomater. Sci.*, 2025, **13**, 1434–1448.
- 111 H. Shen, C. Zhang, Y. Meng, Y. Qiao, Y. Ma, J. Chen, X. Wang and L. Pan, *Adv. Healthcare Mater.*, 2024, **13**, 2303000.
- 112 N. Li, Q. Sun, Z. Yu, X. Gao, W. Pan, X. Wan and B. Tang, *ACS Nano*, 2018, **12**, 5197–5206.
- 113 X. Xue, H. Liu, S. Wang, Y. Hu, B. Huang, M. Li, J. Gao, X. Wang and J. Su, *Composites, Part B*, 2022, **237**, 109855.
- 114 G. Zhou and M. Li, *Chem. Eng. J.*, 2022, **450**, 138348.
- 115 S. Zhang, G. Li, D. Deng, Y. Dai, Z. Liu and S. Wu, *Adv. Ther.*, 2021, **4**, 2100043.
- 116 X.-M. Wang, L. Huang, Y.-J. Wang, L. Xuan, W.-W. Li and L.-J. Tian, *Chem. Eng. J.*, 2021, **405**, 126711.
- 117 X. Liu, W. Law, M. Jeon, X. Wang, M. Liu, C. Kim, P. N. Prasad and M. T. Swihart, *Adv. Healthcare Mater.*, 2013, **2**, 952–957.
- 118 X. Liu, C. Lee, W.-C. Law, D. Zhu, M. Liu, M. Jeon, J. Kim, P. N. Prasad, C. Kim and M. T. Swihart, *Nano Lett.*, 2013, **13**, 4333–4339.



- 119 M. Su, Z. Wu, T. Yan, N. Li, X. Li, T. Hou, J. Liu, C. Zhang, C. Zhu, Z. Wang and J. Zhang, *Adv. Funct. Mater.*, 2024, **34**, 2409580.
- 120 J. Zhang, T. Xing, M. Zhang and Y. Zhou, *Chem. Eng. J.*, 2022, **434**, 134457.
- 121 H. Zhong, S. S. Lo, T. Mirkovic, Y. Li, Y. Ding, Y. Li and G. D. Scholes, *ACS Nano*, 2010, **4**, 5253–5262.
- 122 C. Xia, J. D. Meeldijk, H. C. Gerritsen and C. De Mello Donega, *Chem. Mater.*, 2017, **29**, 4940–4951.
- 123 R. Marin, A. Vivian, A. Skripka, A. Migliori, V. Morandi, F. Enrichi, F. Vetrone, P. Ceroni, C. Aprile and P. Canton, *ACS Appl. Nano Mater.*, 2019, **2**, 2426–2436.
- 124 J. S. Niezgodna, E. Yap, J. D. Keene, J. R. McBride and S. J. Rosenthal, *Nano Lett.*, 2014, **14**, 3262–3269.
- 125 Z. Liu, C. Hao, Y. Sun, J. Wang, L. Dube, M. Chen, W. Dang, J. Hu, X. Li and O. Chen, *Nano Lett.*, 2024, **24**, 5342–5350.
- 126 C. A. Leatherdale, W.-K. Woo, F. V. Mikulec and M. G. Bawendi, *J. Phys. Chem. B*, 2002, **106**, 7619–7622.
- 127 K. E. Knowles, K. H. Hartstein, T. B. Kilburn, A. Marchioro, H. D. Nelson, P. J. Whitham and D. R. Gamelin, *Chem. Rev.*, 2016, **116**, 10820–10851.
- 128 B. Chen, H. Zhong, W. Zhang, Z. Tan, Y. Li, C. Yu, T. Zhai, Y. Bando, S. Yang and B. Zou, *Adv. Funct. Mater.*, 2012, **22**, 2081–2088.
- 129 A. S. Fuhr, H. J. Yun, N. S. Makarov, H. Li, H. McDaniel and V. I. Klimov, *ACS Photonics*, 2017, **4**, 2425–2435.
- 130 A. C. Berends, W. Van Der Stam, J. P. Hofmann, E. Bladt, J. D. Meeldijk, S. Bals and C. De Mello Donega, *Chem. Mater.*, 2018, **30**, 2400–2413.
- 131 A. Farid, L. Qiu, X. Zhang, B. Hu, Z. Wu and G. Wang, *ACS Appl. Nano Mater.*, 2025, **8**, 14703–14712.
- 132 W.-S. Song, J.-H. Kim, J.-H. Lee, H.-S. Lee, Y. R. Do and H. Yang, *J. Mater. Chem.*, 2012, **22**, 21901.
- 133 H. Guo, P. Yang, J. Hu, A. Jiang, H. Chen, X. Niu and Y. Zhou, *ACS Omega*, 2022, **7**, 9642–9651.
- 134 J. Zhang, B. Wang, M. Tebyetekerwa, Y. Zhu, B. Liu, H. T. Nguyen, S. Tian, Y. Zhang and Y. Lu, *J. Mater. Chem. C*, 2019, **7**, 15074–15081.
- 135 W. Peng, J. Du, Z. Pan, N. Nakazawa, J. Sun, Z. Du, G. Shen, J. Yu, J.-S. Hu, Q. Shen and X. Zhong, *ACS Appl. Mater. Interfaces*, 2017, **9**, 5328–5336.
- 136 A. Bora, A. Prudnikau, N. Fu, R. Hübner, K. B. L. Borchert, D. Schwarz, N. Gaponik and V. Lesnyak, *Chem. Mater.*, 2022, **34**, 9251–9260.
- 137 Z. Guan, H. Ye, P. Lv, L. Wang, J. Zhang, B. Zou and A. Tang, *J. Mater. Chem. C*, 2021, **9**, 8537–8544.
- 138 X. Wang and M. T. Swihart, *Chem. Mater.*, 2015, **27**, 1786–1791.
- 139 M. Sachdeva, N. Kharbanda, N. Ghorai and H. N. Ghosh, *J. Phys. Chem. Lett.*, 2025, **16**, 3302–3308.
- 140 V. Sousa, B. F. Gonçalves, M. Franco, Y. Ziouani, N. González-Ballesteros, M. Fátima Cerqueira, V. Yannello, K. Kovnir, O. I. Lebedev and Y. V. Kolen'ko, *Chem. Mater.*, 2019, **31**, 260–267.
- 141 S. Bang, J. Liu, B. Wang, C. M. Perez, T.-R. Liu, K. D. Crans, Z. Sun, A. Strzelecki, O. V. Prezhdo, Y.-T. Shao, X. Guo and R. L. Brutchey, *Chem. Mater.*, 2025, **37**, 2611–2618.
- 142 H. Zhong, Z. Wang, E. Bovero, Z. Lu, F. C. J. M. Van Veggel and G. D. Scholes, *J. Phys. Chem. C*, 2011, **115**, 12396–12402.
- 143 A. Rockett and R. W. Birkmire, *J. Appl. Phys.*, 1991, **70**, R81–R97.
- 144 P. Janicek, V. Kucek, J. Kasparova, T. Plechacek, E. Cernoskova, L. Benes, M. Munzar and C. Drasar, *J. Electron. Mater.*, 2019, **48**, 2112–2119.
- 145 W. Lian, D. Tu, X. Weng, K. Yang, F. Li, D. Huang, H. Zhu, Z. Xie and X. Chen, *Adv. Mater.*, 2024, **36**, 2311011.
- 146 S. Kim, M. Kang, S. Kim, J.-H. Heo, J. H. Noh, S. H. Im, S. I. Seok and S.-W. Kim, *ACS Nano*, 2013, **7**, 4756–4763.
- 147 E. Yang, L. Yu, G. Li, W. Zhang, X. Wang, R. Chen, H. Kang, Z. Chen, E. Guo and T. Wang, *Chem. Mater.*, 2024, **36**, 8753–8762.
- 148 C.-J. Lin, Y. Chen, P. Yang and H.-S. Chen, *J. Alloys Compd.*, 2025, **1021**, 179647.
- 149 A. Lohar, M. Abdul Haque, Y. Jadhav, A. Sagdeo, S. Rane, S. Jadhkar, S. Sartale, Y. D. Kolekar and S. Mahamuni, *J. Phys. Chem. C*, 2025, **129**, 10539–10549.
- 150 R. Sun, J. Zhao, O. Lin, Y. Li, X. Xie, W. Niu, Z. Yin and A. Tang, *J. Mater. Chem. C*, 2024, **12**, 4593–4598.
- 151 M. A. Haque, A. Lohar, Y. Jadhav, R. Kumar, S. N. Jha, D. Bhattacharyya, S. Jadhkar, S. Sartale and S. Mahamuni, *J. Mater. Chem. A*, 2024, **12**, 10726–10736.
- 152 H. Zhi, X. Tong, Y. You, A. I. Channa, X. Li, J. Wu, G. S. Selopal and Z. M. Wang, *Sol. RRL*, 2023, **7**, 2300641.
- 153 C.-H. Ho and C.-C. Pan, *Opt. Mater. Express*, 2013, **3**, 480.
- 154 M. Bettini, *Solid State Commun.*, 1973, **13**, 599–602.
- 155 Y. Shim, K. Hasegawa, K. Wakita and N. Mamedov, *Thin Solid Films*, 2008, **517**, 1442–1444.
- 156 N. N. Syrbu, B. V. Korzun, A. A. Fadzeyeva, R. R. Mianzelen, V. V. Ursaki and I. Galbic, *Phys. B*, 2010, **405**, 3243–3247.
- 157 B. Bhattacharyya, T. Pandit, G. P. Rajasekar and A. Pandey, *J. Phys. Chem. Lett.*, 2018, **9**, 4451–4456.
- 158 B. Bhattacharyya, A. K. Simlandy, A. Chakraborty, G. P. Rajasekar, N. B. Aetukuri, S. Mukherjee and A. Pandey, *ACS Energy Lett.*, 2018, **3**, 1508–1514.
- 159 A. Mukherjee, P. Dutta, B. Bhattacharyya, G. P. Rajasekar, A. K. Simlandy and A. Pandey, *Nano Sel.*, 2021, **2**, 958–966.
- 160 F. Baum, T. Pretto, R. A. Gouvêa and M. J. L. Santos, *Cryst. Growth Des.*, 2022, **22**, 3669–3679.
- 161 Y. Liu, M. Liu, D. Yin, W. Wei, P. N. Prasad and M. T. Swihart, *Chem. Mater.*, 2017, **29**, 3555–3562.
- 162 X. Liu, X. Wang and M. T. Swihart, *Chem. Mater.*, 2015, **27**, 1342–1348.
- 163 U. V. Ghorpade, M. P. Suryawanshi, S. W. Shin, I. Kim, S. K. Ahn, J. H. Yun, C. Jeong, S. S. Kolekar and J. H. Kim, *Chem. Mater.*, 2016, **28**, 3308–3317.
- 164 M. Ibáñez, D. Cadavid, U. Anselmi-Tamburini, R. Zamani, S. Gorsse, W. Li, A. M. López, J. R. Morante, J. Arbiol and A. Cabot, *J. Mater. Chem. A*, 2013, **1**, 1421–1426.



- 165 M. Ahmadi, S. S. Pramana, S. K. Batabyal, C. Boothroyd, S. G. Mhaisalkar and Y. M. Lam, *Inorg. Chem.*, 2013, **52**, 1722–1728.
- 166 J.-J. Wang and K. M. Ryan, *CrystEngComm*, 2016, **18**, 3161–3169.
- 167 J. Van Embden, K. Latham, N. W. Duffy and Y. Tachibana, *J. Am. Chem. Soc.*, 2013, **135**, 11562–11571.
- 168 Y. Yao, A. Bhargava and R. D. Robinson, *Chem. Mater.*, 2021, **33**, 608–615.
- 169 J. C. Kays, C. R. Conti, A. Margaronis, J. E. Kuszynski, G. F. Strouse and A. M. Dennis, *Chem. Mater.*, 2021, **33**, 7408–7416.
- 170 B. Bhattacharyya and A. Pandey, *J. Am. Chem. Soc.*, 2016, **138**, 10207–10213.
- 171 C. Dong, D. Yao, J. Feng, T. Huang, X. Hu, Z. Wu, Y. Liu, B. Yang and H. Zhang, *Chem. Mater.*, 2016, **28**, 9139–9149.
- 172 P. Ramasamy and J. Kim, *Chem. – Asian J.*, 2015, **10**, 1468–1473.
- 173 Y. Tomm, S. Schorr and S. Fiechter, *J. Cryst. Growth*, 2008, **310**, 2215–2221.
- 174 M. D. Regulacio, S. Y. Tee, S. H. Lim, Z. Zhang and M.-Y. Han, *CrystEngComm*, 2018, **20**, 6803–6810.
- 175 H. Li, Q. Zhang, A. Pan, Y. Wang, B. Zou and H. J. Fan, *Chem. Mater.*, 2011, **23**, 1299–1305.
- 176 M. Chakraborty, R. Thangavel, P. Komninou, Z. Zhou and A. Gupta, *J. Alloys Compd.*, 2019, **776**, 142–148.
- 177 C. Yan, E. Gu, F. Liu, Y. Lai, J. Li and Y. Liu, *Nanoscale*, 2013, **5**, 1789.
- 178 B. Li, K. Ye, Y. Zhang, J. Qin, R. Zou, K. Xu, X. Huang, Z. Xiao, W. Zhang, X. Lu and J. Hu, *Adv. Mater.*, 2015, **27**, 1339–1345.
- 179 S. Paul, S. Ghosh, B. Dalal, P. Chal, B. Satpati and S. K. De, *Chem. Mater.*, 2018, **30**, 5020–5031.
- 180 J. Du, X. Wang, X. Dong, C. Zhang, L. Mei, Y. Zang, L. Yan, H. Zhang and Z. Gu, *Nanoscale*, 2019, **11**, 7157–7165.
- 181 R. B. Balow, C. K. Miskin, M. M. Abu-Omar and R. Agrawal, *Chem. Mater.*, 2017, **29**, 573–578.
- 182 R. B. Balow, E. J. Sheets, M. M. Abu-Omar and R. Agrawal, *Chem. Mater.*, 2015, **27**, 2290–2293.
- 183 A. Das, A. Shamirian and P. T. Snee, *Chem. Mater.*, 2016, **28**, 4058–4064.
- 184 S. A. McClary, R. B. Balow and R. Agrawal, *J. Mater. Chem. C*, 2018, **6**, 10538–10546.
- 185 M. Hou, C. Yan, Z. Chen, Q. Zhao, M. Yuan, Y. Xu and B. Zhao, *Acta Biomater.*, 2018, **74**, 334–343.
- 186 E. Dutková, P. Levinský, J. Hejtmánek, K. Knížek, L. Findoráková, M. Baláž, M. Fabián, K. Gáborová, V. Puchý and P. Baláž, *Crystals*, 2025, **15**, 511.
- 187 K. Chen, J. Zhou, W. Chen, P. Zhou, F. He and Y. Liu, *Part. Part. Syst. Charact.*, 2015, **32**, 999–1005.
- 188 F. Zhang, K. Chen, X. Jiang, Y. Wang, Y. Ge, L. Wu, S. Xu, Q. Bao and H. Zhang, *J. Mater. Chem. C*, 2018, **6**, 8977–8983.
- 189 A. Moser, O. Yarema, G. Garcia, M. Luisier, F. Longo, E. Billeter, A. Borgschulte, M. Yarema and V. Wood, *Chem. Mater.*, 2023, **35**, 6323–6331.
- 190 S. Wan, S. Xiao, M. Li, X. Wang, K. H. Lim, M. Hong, M. Ibáñez, A. Cabot and Y. Liu, *Small Methods*, 2024, **8**, 2301377.
- 191 B. Bhattacharyya, C. Balischewski, C. Pacholski, A. Pandey, I. Bald and A. Taubert, *Adv. Opt. Mater.*, 2023, **11**, 2202411.
- 192 H. Liu, C. Chen, H. Wen, R. Guo, N. A. Williams, B. Wang, F. Chen and L. Hu, *J. Mater. Chem. A*, 2018, **6**, 18839–18846.
- 193 Y.-K. Hsu, Y.-G. Lin and Y.-C. Chen, *Mater. Res. Bull.*, 2011, **46**, 2117–2119.
- 194 M. Ma, B. Ruan, M. Zhou, Y. Gu, Q. Dong, Q. Yang, Q. Wang, L. Chen, Y. Shi, J. Yi, G. Chen and Z. Ren, *J. Cryst. Growth*, 2023, **622**, 127398.
- 195 R. Gaspari, G. Della Valle, S. Ghosh, I. Kriegel, F. Scotognella, A. Cavalli and L. Manna, *Nano Lett.*, 2017, **17**, 7691–7695.
- 196 S. Ghosh, T. Avellini, A. Petrelli, I. Kriegel, R. Gaspari, G. Almeida, G. Bertoni, A. Cavalli, F. Scotognella, T. Pellegrino and L. Manna, *Chem. Mater.*, 2016, **28**, 4848–4858.
- 197 G. Gabka, P. Bujak, A. Ostrowski, W. Tomaszewski, W. Lisowski, J. W. Sobczak and A. Pron, *Inorg. Chem.*, 2016, **55**, 6660–6669.
- 198 Y. Kitagawa and H. Yao, *J. Phys. Chem. C*, 2023, **127**, 22149–22156.
- 199 J. E. Kuszynski, X. Zhong, S. A. McGill, A. M. Dennis and G. F. Strouse, *ACS Mater. Lett.*, 2024, **6**, 3367–3375.
- 200 R. Vlasovets, A. Iodchik, J. Tang and M. Artemyev, *ACS Appl. Nano Mater.*, 2025, **8**, 4279–4287.
- 201 R. Yadav, B. Bhattacharyya, S. K. Saha, P. Dutta, P. Roy, G. P. Rajasekar, A. Narayan and A. Pandey, *J. Phys. Chem. C*, 2021, **125**, 2511–2518.
- 202 K. Zhang, W. Chen, Y. Wang, J. Li, H. Chen, Z. Gong, S. Chang, F. Ye, T. Wang, W. Chu, C. Zou and L. Song, *AIP Adv.*, 2015, **5**, 077130.
- 203 X. Zhang, J. Liu, X. Yang, G. He, B. Li, J. Qin, P. R. Shearing, D. J. L. Brett, J. Hu and X. Lu, *Nanoscale*, 2019, **11**, 9733–9742.
- 204 Y. Dong, S. Dong, C. Yu, J. Liu, S. Gai, Y. Xie, Z. Zhao, X. Qin, L. Feng, P. Yang and Y. Zhao, *Sci. Adv.*, 2023, **9**, eadi9980.
- 205 Y. Wang, B. Wang, Z. Yan, Z. Song, H. Yang, F. Wang, X. Xu and Y. Liu, *J. Colloid Interface Sci.*, 2025, **699**, 138277.
- 206 S. Harvey, J. M. DeStefano, J.-H. Chu, D. R. Gamelin and B. M. Cossairt, *Chem. Mater.*, 2024, **36**, 10746–10757.
- 207 C. Steinhagen, M. G. Panthani, V. Akhavan, B. Goodfellow, B. Koo and B. A. Korgel, *J. Am. Chem. Soc.*, 2009, **131**, 12554–12555.
- 208 C. Ritchie, A. S. R. Chesman, M. Styles, J. J. Jasieniak and P. Mulvaney, *Langmuir*, 2018, **34**, 1655–1665.
- 209 A. D. Collord and H. W. Hillhouse, *Chem. Mater.*, 2015, **27**, 1855–1862.
- 210 R. Mainz, A. Singh, S. Levchenko, M. Klaus, C. Genzel, K. M. Ryan and T. Unold, *Nat. Commun.*, 2014, **5**, 3133.



- 211 U. V. Ghorpade, M. P. Suryawanshi, S. W. Shin, C. W. Hong, I. Kim, J. H. Moon, J. H. Yun, J. H. Kim and S. S. Kolekar, *Phys. Chem. Chem. Phys.*, 2015, **17**, 19777–19788.
- 212 B. Zhou, D. Xia and Y. Wang, *RSC Adv.*, 2015, **5**, 70117–70126.
- 213 Y. Gao, F. Long, J. Wang, J. Zhang, S. Mo and Z. Zou, *Mater. Design*, 2017, **123**, 24–31.
- 214 M. D. K. Jones, B. L. Willis, S. Campbell, G. Kartopu, P. Maiello, P. Punathil, W. M. Cheung, E. Woolley, L. C. R. Jones, O. Oklobia, A. Holland, V. Barrioz, G. Zoppi, N. S. Beattie and Y. Qu, *ACS Sustainable Chem. Eng.*, 2024, **12**, 11613–11627.
- 215 M. Zubair, V. A. Lebedev, M. Mishra, T. E. Adegoke, I. S. Amiin, Y. Zhang, A. Cabot, S. Singh and K. M. Ryan, *Chem. Mater.*, 2022, **34**, 10528–10537.
- 216 A. S. R. Chesman, J. Van Embden, E. Della Gaspera, N. W. Duffy, N. A. S. Webster and J. J. Jasieniak, *Chem. Mater.*, 2014, **26**, 5482–5491.
- 217 C.-M. Fan, M. D. Regulacio, C. Ye, S. H. Lim, S. K. Lua, Q.-H. Xu, Z. Dong, A.-W. Xu and M.-Y. Han, *Nanoscale*, 2015, **7**, 3247–3253.
- 218 S. Huang, J. Zai, D. Ma, Q. He, Y. Liu, Q. Qiao and X. Qian, *Chem. Commun.*, 2016, **52**, 10866–10869.
- 219 Y. Zhu, K. He, T. Tadesse Tsega, N. Ali, J. Zai, S. Huang, X. Qian and Z. Chen, *Sustainable Energy Fuels*, 2020, **4**, 2892–2899.
- 220 P. Ramasamy, M. Kim, H.-S. Ra, J. Kim and J.-S. Lee, *Nanoscale*, 2016, **8**, 7906–7913.
- 221 W. Wang, J. Jiang, T. Ding, C. Wang, J. Zuo and Q. Yang, *ACS Appl. Mater. Interfaces*, 2015, **7**, 2235–2241.
- 222 S. Chen, B. Zu, Q. Jin, X. Wu, Z. Xu and L. Wu, *ACS Appl. Mater. Interfaces*, 2025, **17**, 24382–24389.
- 223 R. Marin, A. Skripka, Y.-C. Huang, T. A. J. Loh, V. Mazeika, V. Karabanovas, D. H. C. Chua, C.-L. Dong, P. Canton and F. Vetrone, *Chem. Commun.*, 2020, **56**, 3341–3344.
- 224 L. Chen, Z. Kong, H. Tao, H. Hu, J. Gao and G. Li, *Nanoscale*, 2022, **14**, 3907–3916.
- 225 C. R. McCormick, R. R. Katzbaer, B. C. Steimle and R. E. Schaak, *Nat. Synth.*, 2023, **2**, 152–161.
- 226 F. Shabani, B. Canimkurbey, M. Ahmad, S. Delikanli, Z. Dikmen, F. Isik and H. V. Demir, *Chem. Mater.*, 2025, **37**, 6173–6181.
- 227 Q. A. Akkerman, A. Genovese, C. George, M. Prato, I. Moreels, A. Casu, S. Marras, A. Curcio, A. Scarpellini, T. Pellegrino, L. Manna and V. Lesnyak, *ACS Nano*, 2015, **9**, 521–531.
- 228 Y. Liu, M. Liu, D. Yin, L. Qiao, Z. Fu and M. T. Swihart, *ACS Nano*, 2018, **12**, 7803–7811.
- 229 M. G. Panthani, V. Akhavan, B. Goodfellow, J. P. Schmidtke, L. Dunn, A. Dodabalapur, P. F. Barbara and B. A. Korgel, *J. Am. Chem. Soc.*, 2008, **130**, 16770–16777.
- 230 I. Sharma, K. V. K. Adari, L. Biddika and K. Borse, *Discover Mater.*, 2024, **4**, 76.
- 231 W. van der Stam, A. C. Berends and C. de Mello Donega, *ChemPhysChem*, 2016, **17**, 559–581.
- 232 Q. Guo, H. W. Hillhouse and R. Agrawal, *J. Am. Chem. Soc.*, 2009, **131**, 11672–11673.
- 233 V. A. Akhavan, B. W. Goodfellow, M. G. Panthani, C. Steinhagen, T. B. Harvey, C. J. Stolle and B. A. Korgel, *J. Solid State Chem.*, 2012, **189**, 2–12.
- 234 Y. Ji and P. Mulvaney, *arXiv*, 2024, preprint, arXiv:2408.15795, DOI: [10.48550/arXiv.2408.15795](https://doi.org/10.48550/arXiv.2408.15795).
- 235 K. P. Stroh, M. Szablewski and D. P. Halliday, *Mater. Adv.*, 2024, **5**, 1045–1055.
- 236 S. Campbell, G. Zoppi, L. Bowen, P. Maiello, V. Barrioz, N. S. Beattie and Y. Qu, *ACS Appl. Energy Mater.*, 2023, **6**, 10883–10896.
- 237 D. C. Hayes, S. A. Langdon, R. M. Spilker and R. Agrawal, *ACS Appl. Energy Mater.*, 2024, **7**, 885–895.
- 238 H. Zhou, W.-C. Hsu, H.-S. Duan, B. Bob, W. Yang, T.-B. Song, C.-J. Hsu and Y. Yang, *Energy Environ. Sci.*, 2013, **6**, 2822.
- 239 M. Heidariramsheh, H. Jalalichamani, M. Shabanzadeh, S. M. Mahdavi and N. Taghavinia, *Mater. Res. Bull.*, 2025, **184**, 113243.
- 240 Z. Wang, W. He, M. Ma and Y. Sui, *Energy Rep.*, 2022, **8**, 37–44.
- 241 L. Korala, M. B. Braun, J. M. Kephart, Z. Tregillus and A. L. Prieto, *Chem. Mater.*, 2017, **29**, 6621–6629.
- 242 I. Hwang and K. Yong, *ChemElectroChem*, 2015, **2**, 634–653.
- 243 G. Wang, H. Wei, J. Shi, Y. Xu, H. Wu, Y. Luo, D. Li and Q. Meng, *Nano Energy*, 2017, **35**, 17–25.
- 244 D. H. Jara, S. J. Yoon, K. G. Stamplecoskie and P. V. Kamat, *Chem. Mater.*, 2014, **26**, 7221–7228.
- 245 A. Anand, M. L. Zaffalon, G. Gariano, A. Camellini, M. Gandini, R. Brescia, C. Capitani, F. Bruni, V. Pinchetti, M. Zavelani-Rossi, F. Meinardi, S. A. Crooker and S. Brovelli, *Adv. Funct. Mater.*, 2020, **30**, 1906629.
- 246 L. Jin, G. S. Selopal, X. Tong, D. F. Perepichka, Z. M. Wang and F. Rosei, *Adv. Mater.*, 2024, **36**, 2402912.
- 247 M. R. Bergren, N. S. Makarov, K. Ramasamy, A. Jackson, R. Guglielmetti and H. McDaniel, *ACS Energy Lett.*, 2018, **3**, 520–525.
- 248 Y. Wu, J. Huang, J. Zang, J. Zhou, C. Cheng, Z. Hu, D. Shan, W. Yang, I. Sychugov, L. Sun and B. Xu, *Energy Environ. Sci.*, 2024, **17**, 6338–6349.
- 249 F. Li, M. Zhang, D. Benetti, L. Shi, L. V. Besteiro, H. Zhang, J. Liu, G. S. Selopal, S. Sun, Z. Wang, Q. Wei and F. Rosei, *Appl. Catal., B*, 2021, **280**, 119402.
- 250 L. Wu, Q. Wang, T.-T. Zhuang, G.-Z. Zhang, Y. Li, H.-H. Li, F.-J. Fan and S.-H. Yu, *Nat. Commun.*, 2022, **13**, 5414.
- 251 Y. Chai, Y. Chen, B. Wang, J. Jiang, Y. Liu, J. Shen, X. Wang and Z. Zhang, *ACS Sustainable Chem. Eng.*, 2022, **10**, 8825–8834.
- 252 W. Lian, D. Tu, P. Hu, X. Song, Z. Gong, T. Chen, J. Song, Z. Chen and X. Chen, *Nano Today*, 2020, **35**, 100943.
- 253 J. C. Kays, A. M. Saeboe, R. Toufanian, D. E. Kurant and A. M. Dennis, *Nano Lett.*, 2020, **20**, 1980–1991.
- 254 I. Alayeto, P. Gutierrez-Gonzalez, Q. Dongmei, A. Pallarés Vilar, K. Hamraoui, C. Civera, J. P. Cascales Sandoval,



- S. Melle, O. G. Calderón, D. Méndez-González, J. Rubio-Retama and G. Villaverde, *ACS Appl. Nano Mater.*, 2025, **8**, 14461–14469.
- 255 L. Aldaz-Caballero, U. R. Rodríguez-Mendoza, V. Lavín, P. Canton, A. Benayas and R. Marin, *Adv. Sens. Res.*, 2023, **2**, 2300078.
- 256 Y. Liu, C. Zhao, J. Li, S. Zhao, X. Xu, H. Y. Fu, C. Yu, F. Kang and G. Wei, *ACS Appl. Electron. Mater.*, 2021, **3**, 1236–1243.
- 257 G. Liang, C. Li, J. Zhao, Y. Fu, Z. Yu, Z. Zheng, Z. Su, P. Fan, X. Zhang, J. Luo, L. Ding and S. Chen, *SusMat*, 2023, **3**, 682–696.
- 258 B. S. Shah, S. H. Chaki, S. V. Bhatt and M. P. Deshpande, *Mater. Sci. Semicond. Process.*, 2024, **179**, 108495.
- 259 B. Kumar, S. V. Singh, A. Chattopadhyay, S. Biring and B. N. Pal, *ACS Omega*, 2020, **5**, 25947–25953.
- 260 Z. Liu, C. Hao, Y. Liu, R. Wu, J. Zhang, Z. Chen, F. Wang, L. Guan, X. Li, A. Tang and O. Chen, *ACS Nano*, 2024, **18**, 21523–21533.
- 261 L. J. Lim, X. Zhao and Z. Tan, *Adv. Mater.*, 2023, **35**, 2301887.
- 262 H. Li, X. Jiang, A. Wang, X. Chu and Z. Du, *Front. Chem.*, 2020, **8**, 669.
- 263 N. Kharbanda, M. Sachdeva, N. Ghorai, A. Kaur, V. Kumar and H. N. Ghosh, *J. Phys. Chem. Lett.*, 2024, **15**, 5056–5062.
- 264 H. C. Lee, H. Kim, K. Kim, K. Lee, W. Chung, S. B. Ha, M. Kim, E. Ahn, S. Li, S. Ji, G. Lee, H. Ma, S. J. Lim, H. Choi, J. Kim, H. Ahn, S. In and J. Yang, *Adv. Sci.*, 2025, **12**, e00829.
- 265 J. Zhang, Y. Zhao, K. Qi and S. Liu, *J. Mater. Sci. Technol.*, 2024, **172**, 145–155.
- 266 M. Zhang, N. Fang, X. Song, Y. Chu, S. Shu and Y. Liu, *ACS Omega*, 2020, **5**, 32715–32723.
- 267 Y. Rao, M. Sun, B. Zhou, Z. Wang, T. Yan and Y. Shao, *ACS Appl. Energy Mater.*, 2024, **7**, 5457–5466.
- 268 T. Tran, T. P. Sheng, C. Huang, J. Li, L. Chen, L. Yuan, C. A. Grimes and Q. Cai, *Chem. Eng. J.*, 2012, **210**, 425–431.
- 269 U. V. Ghorpade, M. P. Suryawanshi, S. W. Shin, X. Wang, E. Jo, H. Bae, K. Park, J.-S. Ha, S. S. Kolekar and J. H. Kim, *J. Mater. Chem. A*, 2018, **6**, 19798–19809.
- 270 J. Wang, S. Lin, N. Tian, T. Ma, Y. Zhang and H. Huang, *Adv. Funct. Mater.*, 2021, **31**, 2008008.
- 271 T. Pons, E. Pic, N. Lequeux, E. Cassette, L. Bezdetnaya, F. Guillemin, F. Marchal and B. Dubertret, *ACS Nano*, 2010, **4**, 2531–2538.
- 272 D. Deng, Y. Chen, J. Cao, J. Tian, Z. Qian, S. Achilefu and Y. Gu, *Chem. Mater.*, 2012, **24**, 3029–3037.
- 273 G. Lv, W. Guo, W. Zhang, T. Zhang, S. Li, S. Chen, A. S. Eltahan, D. Wang, Y. Wang, J. Zhang, P. C. Wang, J. Chang and X.-J. Liang, *ACS Nano*, 2016, **10**, 9637–9645.
- 274 H. Wen, H. Wang, J. Hai, S. He, F. Chen and B. Wang, *ACS Sustainable Chem. Eng.*, 2019, **7**, 5200–5208.
- 275 C.-W. Chen, D.-Y. Wu, Y.-C. Chan, C. C. Lin, P.-H. Chung, M. Hsiao and R.-S. Liu, *J. Phys. Chem. C*, 2015, **119**, 2852–2860.
- 276 J. Yang, R. Xie, L. Feng, B. Liu, R. Lv, C. Li, S. Gai, F. He, P. Yang and J. Lin, *ACS Nano*, 2019, **13**, 13144–13160.
- 277 X. Jiang, Y. Han, H. Zhang, H. Liu, Q. Huang, T. Wang, Q. Sun and Z. Li, *ACS Appl. Mater. Interfaces*, 2018, **10**, 43396–43404.
- 278 X. Jiang, S. Zhang, F. Ren, L. Chen, J. Zeng, M. Zhu, Z. Cheng, M. Gao and Z. Li, *ACS Nano*, 2017, **11**, 5633–5645.
- 279 J. Liu, P. Wang, X. Zhang, L. Wang, D. Wang, Z. Gu, J. Tang, M. Guo, M. Cao, H. Zhou, Y. Liu and C. Chen, *ACS Nano*, 2016, **10**, 4587–4598.
- 280 Y. Yang, C. Wang, N. Wang, J. Li, Y. Zhu, J. Zai, J. Fu and Y. Hao, *J. Nanobiotechnol.*, 2022, **20**, 195.
- 281 B. Li, F. Yuan, G. He, X. Han, X. Wang, J. Qin, Z. X. Guo, X. Lu, Q. Wang, I. P. Parkin and C. Wu, *Adv. Funct. Mater.*, 2017, **27**, 1606218.
- 282 M. Duda, P. Joshi, A. Borodziuk, K. Sobczak, B. Sikora-Dobrowolska, S. Maćkowski, A. M. Dennis and Ł. Kłopotowski, *ACS Appl. Mater. Interfaces*, 2024, **16**, 60008–60017.
- 283 <https://www.ubiqd.com/> accessed on 31/10/2025.
- 284 P. Cheng, L. Sun, L. Feng, S. Yang, Y. Yang, D. Zheng, Y. Zhao, Y. Sang, R. Zhang, D. Wei, W. Deng and K. Han, *Angew. Chem., Int. Ed.*, 2019, **58**, 16087–16091.
- 285 Y. Nah, D. Solanki and D. H. Kim, *Cell Rep. Phys. Sci.*, 2022, **3**, 101171.
- 286 P. Vashishtha, G. V. Nutan, B. E. Griffith, Y. Fang, D. Giovanni, M. Jagadeeswararao, T. C. Sum, N. Mathews, S. G. Mhaisalkar, J. V. Hanna and T. White, *Chem. Mater.*, 2019, **31**, 9003–9011.
- 287 R. Huang, L. Yang, F. Yang, Y. Puttisong, Q. Hu, G. Li, J. Hu, Z. Hu, L. Li, J. Tang, W. Chen, Y. Han, J. Luo and F. Gao, *Nat. Commun.*, 2025, **16**, 7264.
- 288 Y. Li, P. Vashishtha, Z. Zhou, Z. Li, S. B. Shivarudraiah, C. Ma, J. Liu, K. S. Wong, H. Su and J. E. Halpert, *Chem. Mater.*, 2020, **32**, 5515–5524.
- 289 C.-X. Li, S.-B. Cho, D.-H. Kim and I.-K. Park, *Chem. Mater.*, 2022, **34**, 6921–6932.
- 290 H.-G. Kwon, S. M. Lee, J. Ryu, J. H. Park, S. K. Kwak and S.-W. Kim, *Chem. Mater.*, 2023, **35**, 1301–1309.
- 291 Z. Cao, J. He, C. Jiao, Z. Liu, L. Xu, C. Zheng, S. Peng and B. Chen, *J. Phys. Chem. Lett.*, 2023, **14**, 2533–2541.
- 292 S. Cui, Y. Chen, S. Tao, J. Cui, C. Yuan, N. Yu, H. Zhou, J. Yin and X. Zhang, *Eur. J. Inorg. Chem.*, 2020, **2020**, 2165–2169.
- 293 R. Lin, Q. Guo, Q. Zhu, Y. Zhu, W. Zheng and F. Huang, *Adv. Mater.*, 2019, **31**, 1905079.
- 294 S. Hull and P. Berastegui, *J. Solid State Chem.*, 2004, **177**, 3156–3173.
- 295 R. Roccanova, A. Yangui, G. Seo, T. D. Creason, Y. Wu, D. Y. Kim, M.-H. Du and B. Saporov, *ACS Mater. Lett.*, 2019, **1**, 459–465.
- 296 J. J. Van Blaaderen, L. A. Van Den Brekel, K. W. Krämer and P. Dorenbos, *Chem. Mater.*, 2023, **35**, 9623–9631.
- 297 N. Krüger, S. Belz, F. Schossau, A. A. Haghghirad, P. T. Cong, B. Wolf, S. Gottlieb-Schoenmeyer, F. Ritter and W. Assmus, *Cryst. Growth Des.*, 2010, **10**, 4456–4462.



- 298 L. Wang, Z. Shi, Z. Ma, D. Yang, F. Zhang, X. Ji, M. Wang, X. Chen, G. Na, S. Chen, D. Wu, Y. Zhang, X. Li, L. Zhang and C. Shan, *Nano Lett.*, 2020, **20**, 3568–3576.
- 299 R. Zhang, X. Mao, D. Zheng, Y. Yang, S. Yang and K. Han, *Laser Photonics Rev.*, 2020, **14**, 2000027.
- 300 J. Chen, J. Li, T. Yin, Z. Yin, Y. Li, C. Wang, L. Guan, J. Xiao, F. Teng and A. Tang, *Adv. Funct. Mater.*, 2025, **35**, 2506331.
- 301 L. Lian, M. Zheng, W. Zhang, L. Yin, X. Du, P. Zhang, X. Zhang, J. Gao, D. Zhang, L. Gao, G. Niu, H. Song, R. Chen, X. Lan, J. Tang and J. Zhang, *Adv. Sci.*, 2020, **7**, 2000195.
- 302 K. Watanabe, K. Imai and H. Nishinaka, *ACS Omega*, 2025, **10**, 10972–10978.
- 303 S. Chatterjee, P. Mukherjee, A. Sen and P. Sen, *J. Phys. Chem. Lett.*, 2024, **15**, 4191–4196.
- 304 T.-H. Le, S. Lee, H. Jo, M. Kim, J. Lee, M. Chang and H. Yoon, *ACS Appl. Nano Mater.*, 2021, **4**, 7621–7627.
- 305 F. Zhang, W. Liang, L. Wang, Z. Ma, X. Ji, M. Wang, Y. Wang, X. Chen, D. Wu, X. Li, Y. Zhang, C. Shan and Z. Shi, *Adv. Funct. Mater.*, 2021, **31**, 2105771.
- 306 K. Arslanova, P. Ganswindt, T. Lorenzen, E. Kostyurina, K. Karaghiosoff, B. Nickel, K. Müller-Caspary and A. S. Urban, *Small*, 2024, **20**, 2403572.
- 307 K. Vighnesh, M. S. Hassan, Z. Li, A. V. Sokolova, A. S. Portniagin, H. Liu, X. Zhao, Y. Wu, X. Zhong and A. L. Rogach, *Nano Lett.*, 2025, **25**, 12414–12421.
- 308 L. Li, Z. Fan, J. Zhang, D. Fan, X. Liu and Y. Wang, *Inorg. Chem.*, 2023, **62**, 19848–19855.
- 309 L. Han, J. Zhao, Y. Miao, W. Zhou, M. Yuan, J. Cui and Z. Ci, *Mater. Res. Bull.*, 2024, **179**, 112980.
- 310 L. Han, B. Sun, C. Guo, G. Peng, H. Chen, Z. Yang, N. Li, Z. Ci and Z. Jin, *Adv. Opt. Mater.*, 2022, **10**, 2102453.
- 311 B. Wang, P. Li, Y. Zhou, Z. Deng, X. Ouyang and Q. Xu, *ACS Appl. Nano Mater.*, 2022, **5**, 9792–9798.
- 312 L. Huang, H. Ye, W. Xiang, H. Fan and X. Liang, *J. Mater. Chem. C*, 2023, **11**, 8524–8532.
- 313 J. Wang, T. Cai and O. Chen, *Nano Lett.*, 2023, **23**, 4367–4374.
- 314 J. Yao, H. Liu, Z. Chen, B. Ma and L. Xu, *Angew. Chem., Int. Ed.*, 2025, **64**, e202423185.
- 315 C. Li, S. Cho, S. Sohn, D. Kwak and I. Park, *Small*, 2025, **21**, 2408468.
- 316 X. Zhang, Y. Tang, C. Lin, H. Xie, R. Fan, D. Yan, X. Wang, Y. Liu and W. Ye, *Opt. Lett.*, 2025, **50**, 3840.
- 317 S. Daliran, M. Khajeh, A. R. Oveisi, J. Albero and H. García, *ACS Appl. Mater. Interfaces*, 2022, **14**, 36515–36526.
- 318 X.-R. Yuan, X.-S. Zhang, X.-Y. Zhao, X.-K. Gong, L. Kong, B.-Z. Zhou, J.-P. Xu and L. Li, *ACS Appl. Nano Mater.*, 2024, **7**, 23214–23227.
- 319 C. Liu, L. Wang, F. Fang, Z. Zhao, J. Pan, J. Akram, S. B. Shafie, R. Z. Talaighil, Q. Li, Z. Zhao, J. Wu, Z. Zhu, W. Lei, X. Zhang and J. Chen, *Front. Mater.*, 2021, **8**, 682833.
- 320 F. Gao, X. Zhu, Q. Feng, W. Zhong, W. Liu, H. Xu and Y. Liu, *Nano Energy*, 2022, **98**, 107270.
- 321 Y. Wang, Y. Yan, D. Li, W. Zhao, S. Chen, Q. Zhong, J. Liu, F. Diarra, M. Cao and Q. Zhang, *Dalton Trans.*, 2021, **50**, 12826–12830.
- 322 J. Feng, J. Wang, D. Wang, M. Han, G. Qian, F. Wu, Q. Lin and Z. Hu, *ACS Appl. Electron. Mater.*, 2022, **4**, 225–232.
- 323 M. R. Patel, P. Chetti, T. J. Park and S. K. Kailasa, *ACS Appl. Nano Mater.*, 2025, **8**, 16186–16195.
- 324 M. Parvizian and J. De Roo, *Nanoscale*, 2021, **13**, 18865–18882.
- 325 M. Dasog, *Chem. Mater.*, 2022, **34**, 4249–4258.
- 326 R. Nie, P. Zhang, J. Gao, C. Wang, W. Chu, L. Li, K. Wang, D. Qian, F. Lin, X. Xia, Y. Wu, L. Chao, C. Miao, X. Zhao, W. Guo and Z. Zhang, *Nat. Commun.*, 2025, **16**, 6438.
- 327 G. H. Ahmed, J. Yin, O. M. Bakr and O. F. Mohammed, *ACS Energy Lett.*, 2021, **6**, 1340–1357.
- 328 K. Sakhatskyi, R. A. John, A. Guerrero, S. Tsarev, S. Sabisch, T. Das, G. J. Matt, S. Yakunin, I. Cherniukh, M. Kotyrba, Y. Berezovska, M. I. Bodnarchuk, S. Chakraborty, J. Bisquert and M. V. Kovalenko, *ACS Energy Lett.*, 2022, **7**, 3401–3414.
- 329 T. Hu, Y. Wang, K. Liu, J. Liu, H. Zhang, Q. U. Khan, S. Dai, W. Qian, R. Liu, Y. Wang, C. Li, Z. Zhang, M. Luo, X. Yue, C. Cong, Y. Yongbo, A. Yu, J. Zhang and Y. Zhan, *Nano-Micro Lett.*, 2025, **17**, 145.

