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Robust large-area molecular junctions of self-assembled monolayers of a model helical paddlewheel complex†

Isabel Coloma,^a Thierry Buffeteau,^b Gilles Pecastaings,^c Santiago Herrero,^{a,d} Elizabeth Hillard,^e Patrick Rosa,^e Miguel Cortijo^{*a} and Mathieu Gonidec^{*e}

We report the preparation of a helical complex and its study in molecular junctions. We show that the SAMs of this racemic compound present electrically robust behaviour which will pave the way for future studies on the CISS effect with analogous enantiopure compounds.

Introduction

In recent years, there has been growing interest in the study of charge transport in chiral media. In particular, it has been shown that the current density across molecular junctions embedding chiral compounds and ferromagnetic electrodes can exhibit a great spin selectivity, which translates into strongly asymmetric J/V characteristics that critically depend on both the direction of magnetization of the electrode and the handedness of the chiral medium.^{1,2} This effect is just one of many manifestations of a phenomenon discovered rather recently, termed the chirality-induced spin selectivity (CISS) effect. Since its discovery in photoelectron scattering experiments in 1999,³ it has been the focus of intense experimental and theoretical efforts^{4–6} and has now been observed in a great variety of physical systems, most typically using self-assembled monolayers of chiral molecular systems.⁷ To this day, however, it remains highly challenging to predict the magnitude or even the occurrence of the CISS effect in new chiral molecular

systems, which is due both to the limited availability of comparable experimental data sets and to the poor correlation between theoretical predictions and actual experimental results, often disagreeing by many orders of magnitude. With this in mind, it is crucial to support theoretical work by developing new platforms for the study of the CISS effect in molecular systems that are versatile and can be tuned rather easily. In this regard, we firmly believe that coordination chemistry provides interesting avenues for the study of the CISS effect since it allows for an unparalleled ad-hoc tuning of the electronic properties of the compounds – by changing the core metal ions – while maintaining the organic ligands, and thus, the overall geometry of the compounds is almost unchanged. While there are some examples of coordination compounds and metalloproteins that have been studied for the CISS effect,^{8,9} the family of linear coordination clusters called the extended metal atom chains (EMACs)¹⁰ would constitute a particularly appealing target for CISS studies. They include chains of up to 11 metal ions¹¹ and possess an optical activity that rivals that of the best helices.^{12,13} Their chemistry, however, is somewhat challenging and the analogous, simpler, paddlewheel metal–metal bonded binuclear ruthenium compounds represent an interesting starting point, owing to their chemical stability and to their rather high spin moment. They have been considerably studied in the last few decades because of their potential applications in fields such as magnetism,^{14–16} molecular electronics,^{17–19} biomedicine,^{20–28} and catalysis.^{29–32} Most diruthenium complexes contain a Ru_2^{5+} core supported by four monoanionic equatorial bridging ligands and one or two axial ligands. The most common electronic configuration of these compounds is $\sigma^2\pi^4\delta^2(\pi^*\delta^*)^3$, with three unpaired electrons in the nearly degenerate π^* and δ^* molecular orbitals and a metal–metal bond order of 2.5.³³ Among diruthenium compounds, $[Ru_2(\mu\text{-ap})_4Cl]$ ($RuapCl$; $ap = 2\text{-anilinopyridinate}$), with four unsymmetrical ap ligands in the same orientation (4,0 isomer), was first reported a few decades ago.³⁴ This compound and its related derivatives have been studied mainly because of their isomerism and the effects of different axial

^aMatMoPol Research Group, Department of Inorganic Chemistry, Faculty of Chemical Sciences, Complutense University of Madrid, Avda. Complutense s/n, 28040 Madrid, Spain. E-mail: miguelcortijomontes@ucm.es

^bUniv. Bordeaux, CNRS, Bordeaux INP, ISM, UMR 5255, F-33405 Talence, France

^cUniv. Bordeaux, CNRS, CRPP, UMR 5031, F-33600 Pessac, France

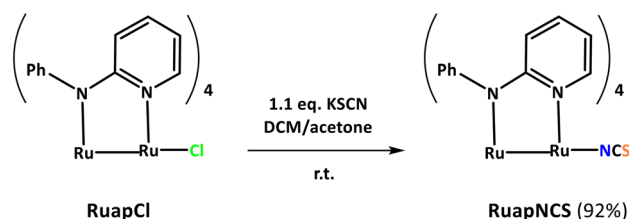
^dKnowledge Technology Institute, Complutense University of Madrid, Campus de Somosaguas, 28223 Pozuelo de Alarcón, Madrid, Spain

^eUniv. Bordeaux, CNRS, Bordeaux INP, ICMCB, UMR 5026, F-33600 Pessac, France. E-mail: mathieu.gonidec@icmcb.cnrs.fr

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Infrared spectroscopy also confirms the isolation of a monodentate isothiocyanate derivative, since a strong --N=C=S asymmetric stretching band is observed at 2034 cm^{-1} (Fig. S2†).⁴⁸ This relatively low value also suggests a strong Ru–N bond due to π back-bonding. The electronic structure of the new complex was studied by UV/vis-NIR spectroscopy. Fig. S3† shows the electronic spectrum in which four bands can be distinguished. Regarding the ultraviolet region, two absorption bands are observed at ~ 275 and $\sim 325\text{ nm}$, which can be assigned to an intra-ligand $\pi \rightarrow \pi^*$ transition and a $\pi(\text{N})/\delta^*(\text{Ru}_2) \rightarrow \pi^*(\text{aryl})$ transition, respectively. In the visible



Scheme 1 One-step synthesis of RuapNCS from the chloride analogue RuapCl.

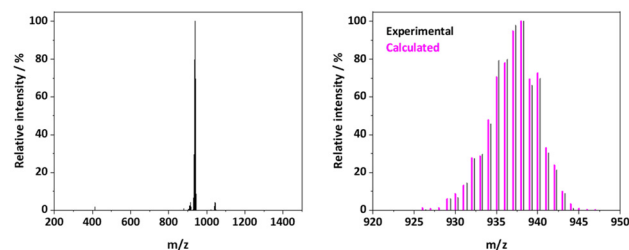


Fig. 1 ESI⁺ spectra (left) and enlargement of the [M]⁺ base peak (right) of RuapNCS.

The morphology of the Au^{TS}-**RuapNCS** substrate after 3.5 h of incubation was probed by AFM in tapping mode and was shown to be extremely smooth (see Fig. 2). The roughness of

ToF-SIMS was performed in both negative and positive modes for substrates incubated in **RuapNCS** and **RuapCl** solutions. The spectra recorded for the Au-**RuapNCS** sample show several Au clusters and three peaks related to the complex can clearly be identified (Fig. 3a and S6[†]) thanks to their rather unique isotopic distribution while the negative mode was found to be not very informative (Fig. S6[†]). In particular, in the positive mode, fragments corresponding to $[M - NCS]^+$ are observed at $m/z = 880.03$. Besides, a minor contribution from the molecular peak is observed at $m/z = 939.18$, which can be assigned to $[M + H]^+$. Altogether, these results suggest the formation of self-assembled monolayers of the complex **RuapNCS** over template-stripped gold (Au^{TS}-**RuapNCS**), where **RuapNCS** molecules are anchored to the substrate through the isothiocyanate axial ligand. Moreover, a two-dimensional mapping analysis of the main peak was performed and the resulting ToF-SIMS image is shown in Fig. 3b. These 2D ToF-SIMS results show good homogeneity of the samples with a rather constant signal across the substrate and no significant local concentration of species that could result from aggregates. The

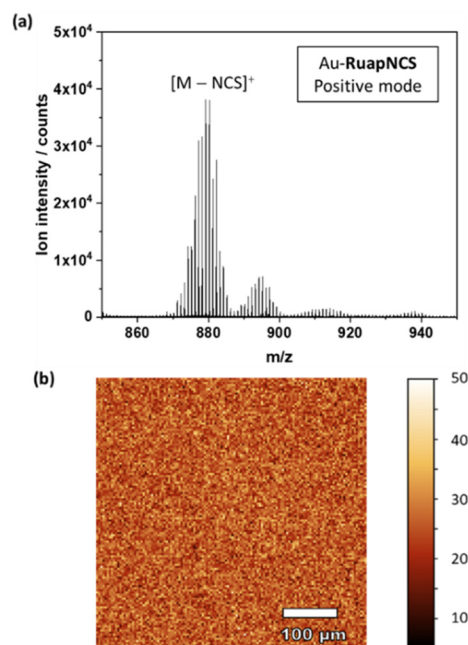


Fig. 2 Tapping-mode AFM topography image for Au^{TS}-RuapNCS.

ToF-SIMS positive spectrum for the control sample Au–RuapCl (see Fig. S7†) also showed a prominent $[M - X]^+$ fragment but its intensity was reduced by a factor of approximately 3 with respect to the Au–RuapNCS monolayer sample (see Fig. S8†). This result likely indicates the presence of some physisorbed RuapCl molecules on the gold surface, with the amount likely depending on the amount of solvent used during the rinsing procedure.

The PM-IRRAS spectra recorded for the Au–RuapNCS (Fig. 4) and Au–RuapCl (Fig. S9†) samples mainly show the same bands as observed in the IR spectrum of the bulk compounds. Importantly, for Au–RuapNCS, a clear shift in the $-N=C=S$ asymmetric stretching band can be observed, which is due to the anchoring of the complex to the substrate through the sulfur atom. The band is shifted from 2034 cm^{-1} to 2083 cm^{-1} , a variation that is consistent with an NCS bridging ligand (Au–SCN–Ru).⁴⁸ Moreover, as expected based on the selection rules of PM-IRRAS, there is a clear effect of the orientation of the molecules on the surface in the PM-IRRAS spectrum with respect to the isotropic ATR spectrum. Some bands are essentially suppressed while others are enhanced. These bands were assigned (Fig. S10†) on the basis of DFT calculations on the isolated molecule at the B3LYP/LANL2DZ level. The $-N=C=S$ asymmetric stretching band is clearly intense, supporting the attachment of the molecule *via* the sulfur and aligned with the $S=C=N$ –Ru–Ru direction normal to the substrate.

The pyridine ring ($C=C$ and $C=N$) stretching band at 1540 cm^{-1} which has a transition dipole along the Ru–Ru direction (Fig. S11†) is also clearly enhanced. In contrast, the pyridine in-plane CH deformation and phenyl in-plane CH deformation bands at 1284 and 1215 cm^{-1} (Fig. S11†), respectively, that both correspond to transverse modes, are strongly suppressed in the PM-IRRAS spectrum. Altogether, this analysis confirms that the molecules are grafted on the surface *via* the sulfur atom with the Ru–Ru bond essentially normal to the surface.

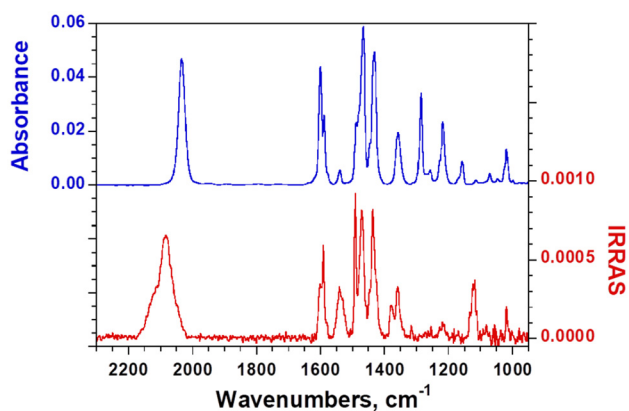


Fig. 4 Comparison between the bulk FTIR-ATR spectrum for RuapNCS and the PM-IRRAS spectrum for Au–RuapNCS.

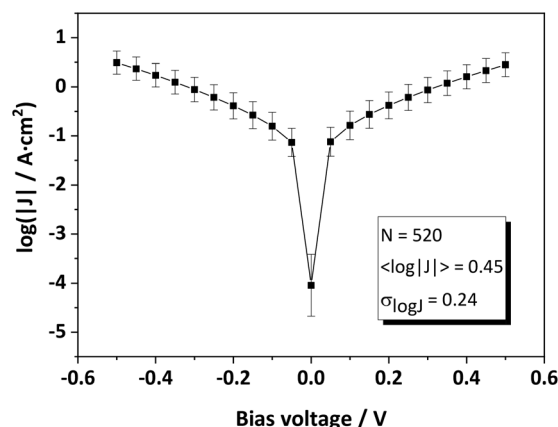


Fig. 5 Plot of $\langle \log(|J|) \rangle$ versus applied bias voltage (V) for Au^{TS}–RuapNCS//Ga₂O₃/EGaIn junctions. The error bars correspond to the standard deviation.

Au^{TS}–RuapNCS//Ga₂O₃/EGaIn molecular junctions

In order to evaluate the robustness of the monolayers with respect to charge transport measurements, we then studied the Au^{TS}–RuapNCS SAMs using eutectic gallium–indium (EGaIn) top electrodes by using the classical conical tip method at room temperature. The resulting Au^{TS}–SAM//Ga₂O₃/EGaIn large-area junctions were measured in the $\pm 0.5\text{ V}$ range. To obtain statistically robust data, we measured 14 junctions and accumulated 520 traces. The results are shown in Fig. 5. The junction yield was excellent; in fact, we did not observe a single short in the 14 junctions, and there were no significant outliers in the dataset.

The current density values were also found to be very robust and repeatable, with an average value of $\log(|J|/\text{A cm}^{-2}) = 0.45 \pm 0.24$ ($N = 520$) at 0.5 V ; the standard deviation is essentially as good as that seen for model SAMs of *n*-alkanethiols⁵² and the current density histogram presents a well-defined normal distribution (see Fig. S12†). There is no significant rectification in these junctions, which is not surprising given the molecular structure of the compounds, and this is consistent with the fact that the highest occupied molecular orbitals (as calculated by DFT, see Fig. S13†) are essentially delocalized over the whole complex. Moreover, the current density at 0.5 V ($\log|J| = 0.45$) is of the same order of magnitude as that observed for oligophenyl-methanethiol SAMs of similar length.⁵³ For the sake of comparison, we also attempted to measure molecular junctions of a control sample of Au^{TS}–RuapCl. However, as expected, the stability and/or quality of the physisorbed layer was too poor to allow measurement of its charge transport characteristics, and the EGaIn junctions invariably led to shorts.

Conclusions

We have shown here, with a model compound, that substituting linear coordination clusters with an axial isothiocyanate



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