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Synthesis of a series of rare-earth-based multi-anion chalcogenide iodides $\text{RE}_3\text{Si}_2\text{Se}_x\text{S}_{8-x}\text{I}$ (RE = La, Ce, Pr, and Nd) using the flux-assisted boron–chalcogen mixture method†

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Single crystals and polycrystalline powders of rare earth mixed chalcogenide iodides $\text{La}_3\text{Si}_2\text{Se}_{1.21}\text{S}_{6.79}\text{I}$, $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$, $\text{Pr}_3\text{Si}_2\text{Se}_{1.22}\text{S}_{6.78}\text{I}$, and $\text{Nd}_3\text{Si}_2\text{Se}_{1.18}\text{S}_{6.82}\text{I}$ were prepared using the reactive flux-assisted boron–chalcogen mixture (BCM) method at 850 °C. All compounds crystallized in the monoclinic crystal system, space group $C2/c$ (space group number 15). The series adopts the $\text{La}_3\text{Si}_2\text{O}_8\text{Cl}$ structure type, containing isolated SiQ_4 tetrahedra connected by REQ_8 (RE = La, Ce, Pr and Nd) polyhedra; this arrangement creates tunnels that are filled by I atoms. The partial substitution of S by Se was carried out to modulate the optical properties. Phase pure samples and uniform solid solutions were obtained for all compositions as determined using powder X-ray diffraction patterns. Polycrystalline powders were used for physical property measurements, including magnetic susceptibility and UV-Vis diffuse reflectance. The solid-state UV-Vis data for the polycrystalline $\text{La}_3\text{Si}_2\text{Se}_{1.21}\text{S}_{6.79}\text{I}$, $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$, and $\text{Pr}_3\text{Si}_2\text{Se}_{1.22}\text{S}_{6.78}\text{I}$ samples revealed band gaps of $E_g = 2.5(1)$, $2.2(1)$, and $2.3(1)$ eV, typical of semiconductors. Magnetic measurements indicated that $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ and $\text{Nd}_3\text{Si}_2\text{Se}_{1.18}\text{S}_{6.82}\text{I}$ exhibit paramagnetic behavior with slightly negative Weiss constants $\theta = -25$ and -38 . The photoluminescence spectrum of $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ exhibits a broad emission band around ~ 493 nm.

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Introduction

Rare-earth chalcogenides exhibit diverse structural chemistry and display a broad range of physical properties, primarily due to their 4f electrons and, at times, due to their stability in higher coordination environments. This combination of factors leads to a variety of intriguing physical behaviors, making rare-earth chalcogenides promising candidates for different applications, including use as phosphors, magnetic materials, magneto-optical devices, and optical materials for various applications.^{1–5} The physical properties of rare-earth chalcogenides can be further altered by modifying their compositions, particularly through the incorporation of mixed cations or mixed anions.

The introduction of more than one type of anion into metal chalcogenides can significantly impact both the structure and the properties of the resulting materials. This enhancement often leads to novel characteristics that are unattainable in materials with only a single anion or in mixed-cation compounds,⁶ where examples include transparent conductors, inorganic pigments, luminescent materials, cathode materials in batteries, thermoelectric materials, and even superconductors.^{7–14} To date, however, mixed-anion compounds that incorporate heavier chalcogenides alongside halides remain relatively rare. In the quaternary system RE-Tt-Q-X (where RE = rare-earth metals, Tt = group 14 elements Si, Ge, or Sn, Q = chalcogenides S, Se, or Te, and X = halides Cl, Br, or I), only a few quaternary phases have been reported.^{15–20} Similarly, compounds that specifically explore mixed chalcogenide–halide compositions are rare. One notable system where this has been studied is $\text{Ce}_3\text{Si}_2(\text{S}_{1-y}\text{Se}_y)_8\text{I}$,²¹ in which sulfur (S) is partially or fully substituted by selenium (Se), resulting in three different compositions. This appears to be the only system in which this has been reported. The three distinct compositions each exhibit a striking change in the luminescence emission as sulfur is progressively replaced by selenium. This color tunability, driven by anion substitution,

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flux-assisted boron–chalcogen mixture (BCM) method. The synthesis process involved the use of RE_2O_3 (La or Nd), CeO_2 and Pr_6O_{11} , and SiO_2 oxide reagents as rare earth and silicon sources, respectively. Boron played a crucial role in this reaction by acting as a reducing agent to remove oxygen from the oxide reagents, which has been explained extensively in earlier work.^{1,22,32,33} This step was essential to ensure that the rare earth and silicon elements were available in their reduced forms, enabling them to react with elemental sulfur and selenium to form the desired chalcogenide compounds. In our experiments, an excess of NaI flux was employed to promote the formation of both single crystals and polycrystalline materials and also as a source of I.

After synthesis, the reaction products were thoroughly washed with methanol to dissolve the residual NaI flux and eliminate any impurities. To verify the composition, we conducted energy dispersive spectroscopy (EDS) analyses. These analyses confirmed the presence of all expected elements in the crystal (Table S3†) and no extraneous elements were detected. The phase-purity of the samples was confirmed by PXRD analysis (Fig. 2 and Fig. S1–S3†), and these compounds were used for further characterization. Additionally, our observations revealed that all synthesized phases remained stable under ambient conditions. The thermal stability of the $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ polycrystalline powder samples was investigated through thermogravimetric analysis (TGA) conducted under a nitrogen atmosphere, with the temperature increasing progressively up to 500 °C (Fig. S4†). The TGA results indicated minimal weight loss over the examined temperature range, suggesting the stability of the material under an N_2 atmosphere. To verify the findings and assess any potential structural changes induced by heating, powder X-ray diffraction

(PXRD) analysis was performed on the samples after the TGA measurements. The PXRD patterns of the post-TGA samples revealed no observable changes, confirming that the phase is quite stable, and no oxidation or decomposition happened during the TGA experiment (Fig. S5†).

Crystal structure

The compounds $\text{RE}_3\text{Si}_2\text{Se}_x\text{S}_{8-x}\text{I}$ crystallize in the monoclinic $C2/c$ space group, sharing an isostructural relationship with the previously reported $\text{La}_3\text{Si}_2\text{O}_8\text{Cl}$ compound. The monoclinic crystal structure of the $\text{RE}_3\text{Si}_2\text{Se}_x\text{S}_{8-x}\text{I}$ compounds is also consistent with other known compounds such as $\text{RE}_3\text{Si}_2\text{O}_8\text{Cl}$, $\text{RE}_3\text{Si}_2\text{O}_8\text{X}$ ($\text{X} = \text{Cl}$, Br , and I), $\text{RE}_3\text{Ge}_2\text{S}_8\text{I}$, and $\text{RE}_3(\text{Ge}_{1-x}\text{Si}_x)_2\text{S}_8\text{I}$.^{16,18,34,35} The unit cell contains eight independent atomic sites, with four mixed chalcogenide sites (8f) being occupied by both S and Se. For the mixed-chalcogen $\text{RE}_3\text{Si}_2\text{Se}_x\text{S}_{8-x}\text{I}$ compounds, structural refinements provided evidence for site preferences of S and Se. Only one of the four chalcogen sites (labeled Q1 in Fig. 3d) contained a significant amount of Se (35%), while the remaining three sites are predominantly occupied by S with only ~10% Se for sites Q2 and Q3 and less than ~5% for site Q4. The details of the site occupancies of S and Se in $\text{RE}_3\text{Si}_2\text{Se}_x\text{S}_{8-x}\text{I}$ are provided in Table S1.† The Q atoms, specifically Q1, Q2, and Q3, which exhibit a higher concentration of selenium (Se), are strategically positioned around the channel within the unit cell structure where they can interact with the iodine in the center of the channel, as illustrated in Fig. 5. It is likely that the softer and more covalent selenium prefers interactions with the soft and covalent iodine, resulting in higher selenium occupancies in Q1, Q2, and Q3, vs. Q4.

The three-dimensional tunnel structure of $\text{RE}_3\text{Si}_2\text{Se}_x\text{S}_{8-x}\text{I}$ along the c -direction is shown in Fig. 3a. It consists of isolated SiQ_4 tetrahedra that are connected by REQ_8 ($\text{RE} = \text{La}$, Ce , Pr and Nd) polyhedra; this arrangement creates tunnels that are filled by I atoms (Fig. 3b, d and e). Both RE1 and RE2 form bicapped trigonal prisms with eight neighboring chalcogens, which extend to a tricapped trigonal prismatic structure once the additional long bond to I is included (Fig. 3c and f). The complete 3D structure contains two different layers ${}^2_{\infty}[\text{RE1SiQ}_4]$ and ${}^1_{\infty}[\text{RE2Q}_6]$ connected *via* edge sharing (Fig. 4c and d). The ${}^2_{\infty}[\text{RE1SiQ}_4]$ layers are constructed by zigzagging ribbons of face-sharing ${}^1_{\infty}[\text{RE1Q}_5]$ polyhedra that extend along the b -axis (Fig. 4a and b); SiQ_4 tetrahedra are located above and below the ${}^1_{\infty}[\text{RE1Q}_5]$ layers. The resulting 3D structure consisting of these two layers displays tunnels along the c -axis that are filled by iodine (Fig. 3a).³⁵

The RE–Q bonds within the bicapped trigonal prisms of the $\text{RE}_3\text{Si}_2\text{Se}_x\text{S}_{8-x}\text{I}$ compositions exhibit bond lengths ranging from 2.9420(5) Å to 3.167(1) Å and the Si–Q bonds within these tetrahedra range from 2.1157(8) Å to 2.259(1) Å. These bond lengths are consistent with those found in related chalcogenide compounds, suggesting that the RE–Q and Si–Q bonding interactions in $\text{RE}_3\text{Si}_2\text{Se}_x\text{S}_{8-x}\text{I}$ are typical of this class of materials.^{36–41} The minimum RE...RE distance decreases from La ($d_{\text{La-La}} = 4.5034(5)$ Å) to Nd ($d_{\text{Nd-Nd}} = 4.4065(5)$ Å), presum-

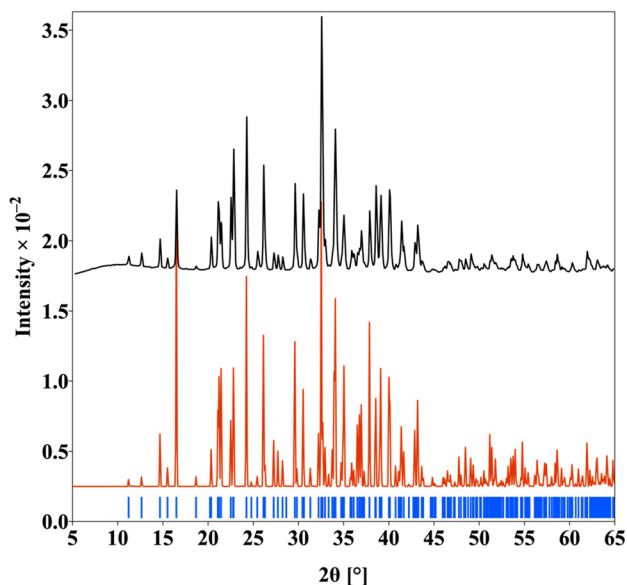


Fig. 2 The PXRD pattern of the polycrystalline $\text{Nd}_3\text{Si}_2\text{Se}_{1.18}\text{S}_{6.82}\text{I}$ compound (simulated pattern (red), experimental pattern (black) and Bragg position (blue)).



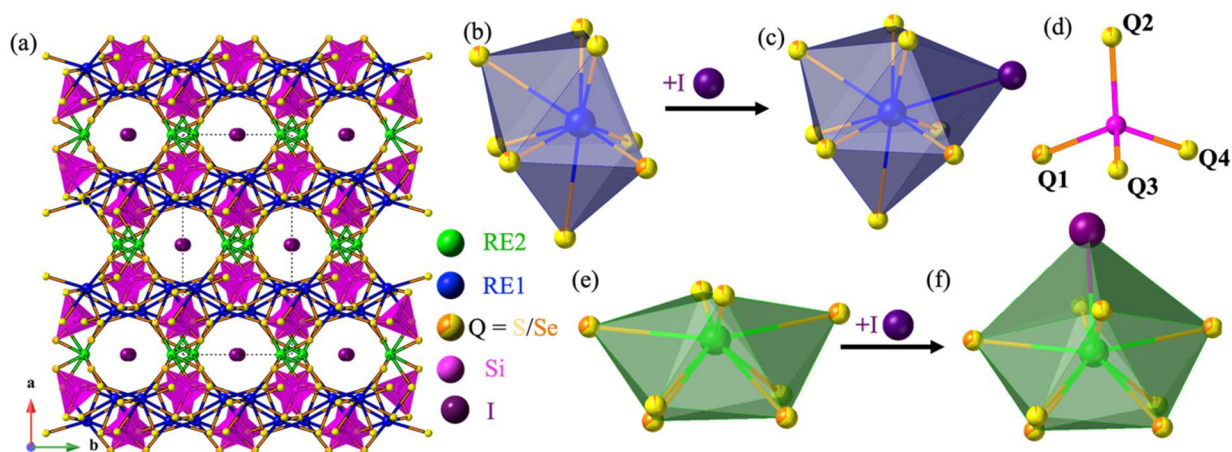


Fig. 3 The (a) unit cell view of $\text{RE}_3\text{Si}_2\text{Se}_{8-x}\text{I}$ compounds, (b) RE1Q_8 bicapped trigonal prism, (c) RE2Q_8 bicapped trigonal prism, (d) SiQ_4 tetrahedron, (e) RE2Q_8 bicapped trigonal prism, and (f) $\text{RE1Q}_8\text{I}$ tricapped trigonal prism.

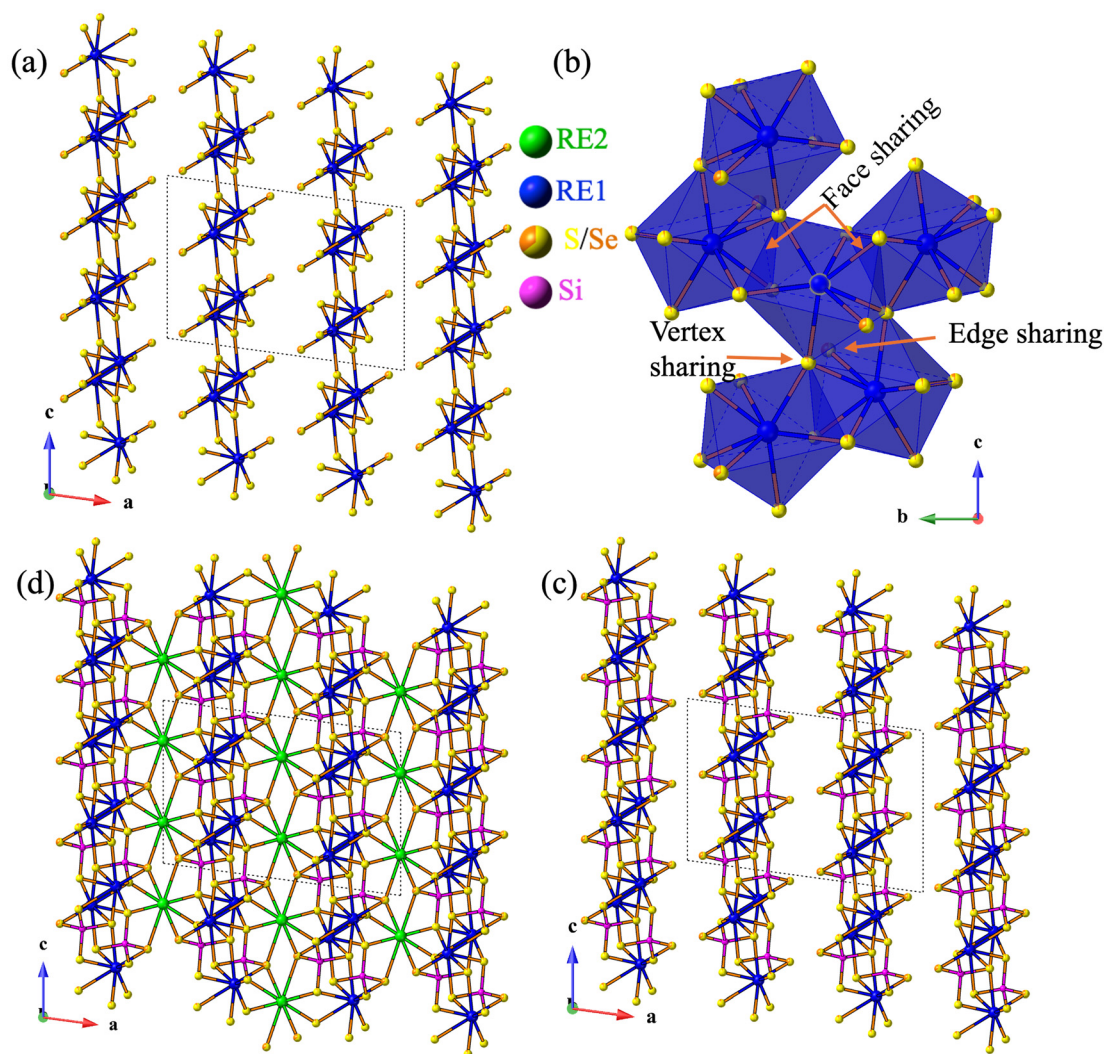


Fig. 4 (a) $2\text{D } \frac{1}{\infty} [\text{RE1Q}_5]$ layers along the b -axis, (b) $2\text{D } \frac{1}{\infty} [\text{RE1Q}_5]$ ribbons along the a -axis, (c) $\frac{2}{\infty} [\text{RE1SiQ}_4]$ layers viewed along the b -direction, and (d) total 3D view of $[\text{RE}_3\text{Si}_2\text{Se}_{8-x}\text{I}]^{1+}$, with RE2Q_8 bicapped trigonal prisms bridging between $\frac{2}{\infty} [\text{RE1SiQ}_4]^{1-}$ layers.



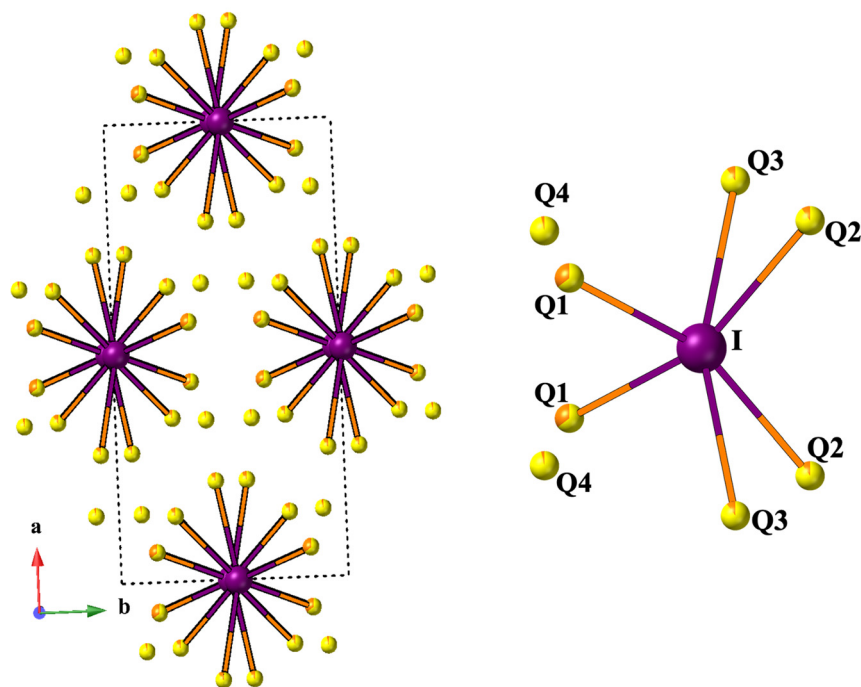


Fig. 5 The coordination environment of I with Q atoms (Q = S/Se). Iodine is purple and Q is S/Se yellow/orange.

ably due to the lanthanide contraction. The lanthanide contraction also may cause the shrinkage of the cavity size where the minimum I...I distance also decreases from 5.4966(4) Å to 5.4067(4) Å. The RE–I bonds range from RE1–I of 3.4773(5) Å to 3.4488(5) Å and RE2–I of 3.3333(5) Å to 3.2496(5) Å.

Magnetic properties

The temperature-dependent magnetization of polycrystalline samples of $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ and $\text{Nd}_3\text{Si}_2\text{Se}_{1.18}\text{S}_{6.82}\text{I}$ was measured over the temperature range of 2 K–300 K. The magnetic susceptibility (χ) was found to follow the Curie law upon cooling down to 2 K with no evidence of long-range magnetic ordering for $\text{Nd}_3\text{Si}_2\text{Se}_{1.18}\text{S}_{6.82}\text{I}$; however, an antiferromagnetic transition ($T_N = 6$ K) was observed for the $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ compound. The magnetic susceptibility vs. temperature plots of $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ and $\text{Nd}_3\text{Si}_2\text{Se}_{1.18}\text{S}_{6.82}\text{I}$ are shown in Fig. 6a and b. The M vs. H plots of $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ and $\text{Nd}_3\text{Si}_2\text{Se}_{1.18}\text{S}_{6.82}\text{I}$ are shown in Fig. 6c and d which are collected over an applied magnetic field of 7 T to –7 and 5 T to –5 T. The inverse susceptibility data were fitted to the Curie law using the high temperature region (200–300 K) to calculate the Curie (C) and Weiss (θ) constants (see the inset of Fig. 6a and b).

The effective magnetic moments determined from the linear fit are $4.61\mu_B$ and $6.57\mu_B$ per formula unit, for the Ce and Nd compositions, which are in good agreement with the theoretical moments expected for Ce^{3+} and Nd^{3+} ions, with values of $4.39\mu_B$ and $6.27\mu_B$, respectively. The negative Weiss constants ($\theta = -25$ and -38) observed for the $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ and $\text{Nd}_3\text{Si}_2\text{Se}_{1.18}\text{S}_{6.82}\text{I}$ compounds indicate antiferromagnetic

interactions between f-electrons on neighboring rare-earth ions, likely due to the intervening sulfur and selenium ions. The minimum distance between neighboring Nd atoms observed is 4.4065(5) Å. No difference between the field-cooled (FC) and zero field-cooled (ZFC) data was observed; the FC data are shown in Fig. 6 and the ZFC data are shown in Fig. S6.† The M vs. H plot measured at 2 K begins to saturate at higher magnetic fields as expected, Fig. 6c and d.

UV-visible diffuse reflectance spectroscopy

The optical properties of the quaternary compounds $\text{La}_3\text{Si}_2\text{Se}_{1.21}\text{S}_{6.79}\text{I}$, $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$, and $\text{Pr}_3\text{Si}_2\text{Se}_{1.22}\text{S}_{6.78}\text{I}$ were examined using polycrystalline powder samples. The Kubelka–Munk equation $\alpha/S = (1 - R)^2/2R$ was used (where α is the absorption coefficient, R is the reflectance, and S is the scattering coefficient) to analyze the optical absorption data, as shown in Fig. 7. The absorption data reveal band gaps of 2.5(1) eV, 2.2(1) eV and 2.3(1) eV for the polycrystalline $\text{La}_3\text{Si}_2\text{Se}_{1.21}\text{S}_{6.79}\text{I}$, $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$, and $\text{Pr}_3\text{Si}_2\text{Se}_{1.22}\text{S}_{6.78}\text{I}$ samples, consistent with the semiconducting nature of the composition. This band gap value of $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ is lower than that observed for the related compound $\text{Ce}_3\text{Si}_2\text{S}_8\text{I}$, and $\text{Ce}_3\text{Si}_2\text{S}_6\text{SeI}$ which have reported band gaps of 2.78 eV and 2.39 eV, respectively, and higher than that of $\text{Ce}_3\text{Si}_2\text{S}_8\text{I}$ (1.92 eV).²¹ The observed reduction in the band gap may be attributed to the substitution of selenium (Se) for sulfur (S) in the lattice. Selenium, being less electronegative and having a higher atomic radius than sulfur, can contribute to a decrease in the overall band gap by enhancing the covalency in the bonding. Consequently, the introduction of Se into



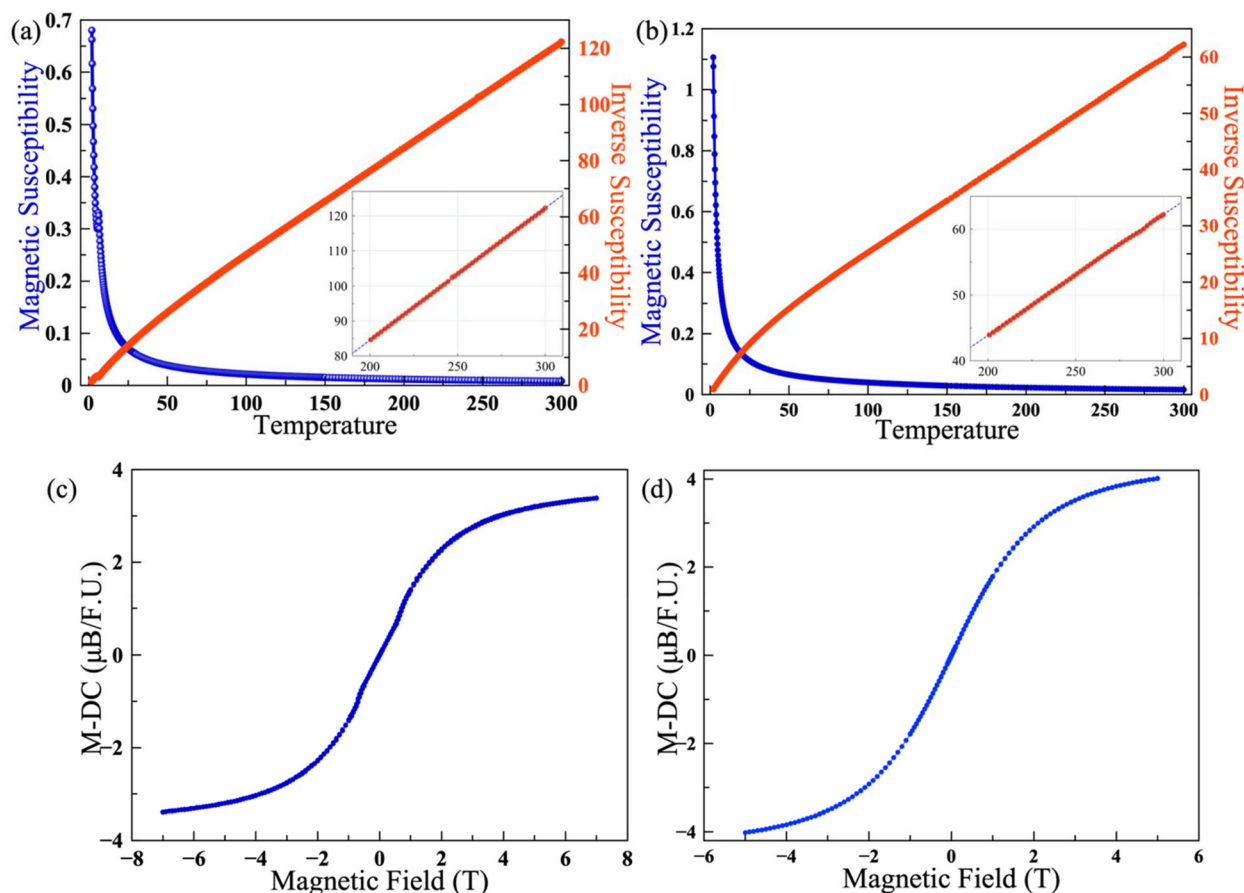


Fig. 6 Molar susceptibility and inverse molar susceptibility vs. temperature plots for polycrystalline (a) $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ and (b) $\text{Nd}_3\text{Si}_2\text{Se}_{1.18}\text{S}_{6.82}\text{I}$ measured in an applied field of 1000 Oe; magnetization as a function of applied field plots for polycrystalline (c) $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ and (d) $\text{Nd}_3\text{Si}_2\text{Se}_{1.18}\text{S}_{6.82}\text{I}$ measured at 2 K.

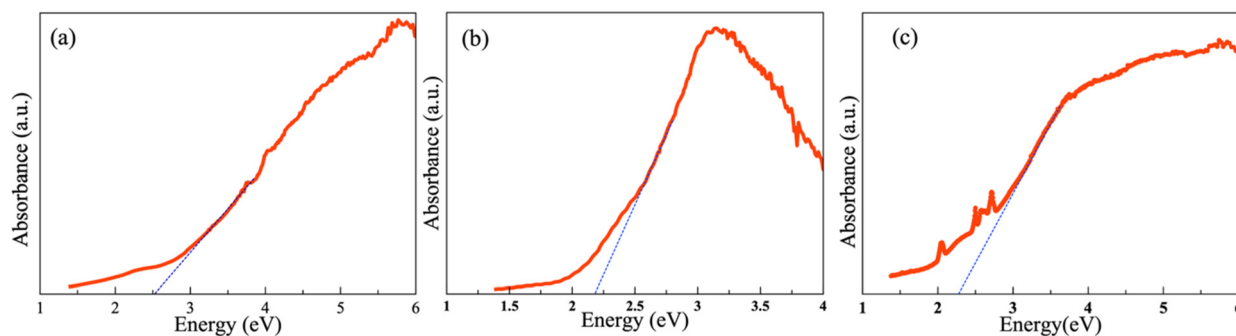


Fig. 7 Optical absorption plot of polycrystalline (a) $\text{La}_3\text{Si}_2\text{Se}_{1.21}\text{S}_{6.79}\text{I}$, (b) $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$, and (c) $\text{Pr}_3\text{Si}_2\text{Se}_{1.22}\text{S}_{6.78}\text{I}$.

$\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ is the likely reason for the narrower band gap in $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ relative to $\text{Ce}_3\text{Si}_2\text{S}_8\text{I}$.

Photoluminescence

Photoluminescence spectra were recorded using a large (2.4 mm × 0.8 mm × 0.6 mm) yellow crystal of $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$. The spectra were recorded using a 375 nm ultraviolet diode laser as the excitation source. The spectra

revealed a broad emission band centered around ~493 nm, which is attributed to the 5d–4f electronic transition of the Ce^{3+} ions (Fig. 8). This emission, observed in the blue-green region, is similar to that reported for the related compound $\text{Ce}_3\text{Si}_2\text{S}_8\text{I}$ which shows strong luminescence in the blue region at ~462 nm. The slight redshift in emission observed for $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ can likely be attributed to differences in the local environment around the Ce atom, possibly due to the



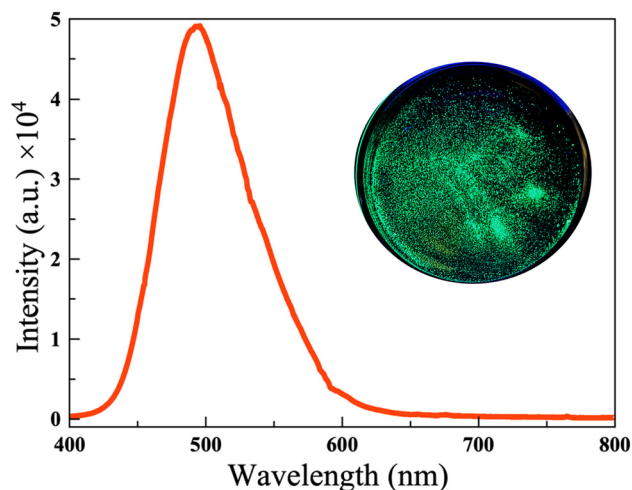


Fig. 8 Emission spectra of the $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ single crystal (the inset figure shows the green color of the sample under exposure to UV light).

influence of the selenium atoms. The photoluminescence in this compound thus appears to be sensitive to the local coordination environment around the Ce atom. The little asymmetric shape of the emission peak may arise from the presence of two distinct types of emissive centers, which is consistent with the existence of two rare-earth (RE) sites within the structure. These RE sites experience slightly different local environments due to the inherent disorder of the surrounding sulfur (S) and selenium (Se) atoms. Variations in the local coordination geometry and ligand field strength at each site can lead to differences in emission energy, resulting in overlapping of emission bands. This overlap can cause spectral broadening and asymmetry in the photoluminescence peak. We also investigated its scintillation ability by illuminating a powder sample of $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ with Cu X-rays. The green scintillation that results confirms the ability of this material to scintillate. (Fig. S7†). These observations suggest that selective substitution of chalcogen atoms in this family of compounds can be used to fine-tune the photoluminescence emission behavior.

Conclusion

The rare-earth mixed chalcogenide iodides $\text{La}_3\text{Si}_2\text{Se}_{1.21}\text{S}_{6.79}\text{I}$, $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$, $\text{Pr}_3\text{Si}_2\text{Se}_{1.22}\text{S}_{6.78}\text{I}$, and $\text{Nd}_3\text{Si}_2\text{Se}_{1.18}\text{S}_{6.82}\text{I}$ were successfully synthesized using a flux-assisted BCM method, yielding phase-pure samples that crystallize in the monoclinic $C2/c$ space group. These compounds exhibit a three-dimensional tunnel structure resembling the $\text{La}_3\text{Si}_2\text{O}_8\text{Cl}$ type, with isolated SiQ_4 tetrahedra connected to rare-earth polyhedra to generate a tunnel structure in which the tunnels are occupied by iodide ions. Partial substitution of S with Se impacts the optical properties of these materials, as demonstrated by UV-visible spectroscopy and photoluminescence studies. Specifically, $\text{Ce}_3\text{Si}_2\text{Se}_{1.39}\text{S}_{6.61}\text{I}$ has a band gap of 2.2(1) eV classifying it as a semiconductor. It also exhibits luminescence be-

havior with a broad emission band around 493 nm. The materials appear to be paramagnetic, as exemplified by $\text{Nd}_3\text{Si}_2\text{Se}_{1.18}\text{S}_{6.82}\text{I}$, which exhibits paramagnetic behavior with a negative Weiss constant. These findings highlight the potential of these mixed chalcogenide iodides for optical applications, where the optical properties can be influenced by the substitution of Se for S.

Data availability

The CCDC 2388413–2388416† entries encompass the supplementary crystallographic data associated with this paper. Raw data for other measurements are available upon request.

Conflicts of interest

The authors declare no competing financial interest.

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