ChemComm

Check for updates

Induced chirality at the surface: fixation of a dynamic *M/P* invertible helical Co₃ complex on SiO₂†

Satoshi Muratsugu, [ID] *ab Kana Sawaguchi,ac Takafumi Shiraogawa, [ID] d Shunsuke Chiba,e Yoko Sakata, [ID] ef Sora Shirai,a Hiroshi Baba,a Masahiro Ehara, [ID] *d Shigehisa Akine [ID] *ef and Mizuki Tada [ID] *abc

Dynamic *M/P* invertible helicity was successfully induced at a SiO₂ surface immobilized with a dynamic helical trinuclear cobalt complex, [LCo₃(NHMe₂)₆](OTf)₃, using chiral ((R) or (S))-1-phenylethylamine. Solid-state CD spectra and theoretical calculations suggested that the fixation of the *M/P* helical complex on the surface *via* coordination interactions was the key factor of the induced chirality at the surface.

The helix is one of the chiral structures derived from spiral or propeller-like structures, and left- and right-handed orientations of spiral structures produce helicity.[1,2] Helical molecules have been applied to molecular machines,[3,4] molecular recognition,[5] optical probes,[6] and so on. In particular, helical molecules with a cage structure have attracted much attention, because these molecules could be used as a chiral host for enantioselective guest binding, a chiral nanoreactor, a chiral molecular switch triggered by guest binding, *etc.* The tris(saloph) organic cage (denoted as H₆L)[7–11] is one of the useful helical molecular cages, because metal complexes LM₃ were found to show unique dynamic structural conversions in response to guest binding.[9–11]

Among them, trinuclear cobalt(III) helical cage complexes, denoted as [LCo₃] unit, have great advantages for post-synthetic

a Department of Chemistry, Graduate School of Science, Nagoya University, Furo-cho, Chikusa-ku, Nagoya, 464-8602 Aichi, Japan. E-mail: smuratsugu@chem.nagoya-u.ac.jp, tada.mizuki.u6@f.mail.nagoya-u.ac.jp

b Integrated Research Consortium on Chemical Science (IRCCS), Nagoya University, Furo-cho, Chikusa-ku, Nagoya, 464-8602 Aichi, Japan

c Research Center for Materials Science (RCMS), Nagoya University, Furo-cho, Chikusa-ku, Nagoya, 464-8602 Aichi, Japan

d Institute for Molecular Science/School of Physical Sciences, Graduate University for Advanced Studies, Myodaiji, Okazaki, 444-8585 Aichi, Japan. E-mail: ehara@ims.ac.jp

e Graduate School of Natural Science and Technology, Kanazawa University, Kakuma-machi, Kanazawa, 920-1192 Ishikawa, Japan. E-mail: akine@se.kanazawa-u.ac.jp

f Nano Life Science Institute (WPI-NanoLSI), Kanazawa University, Kakuma-machi, Kanazawa, 920-1192 Ishikawa, Japan

† Electronic supplementary information (ESI) available: Whole image of Scheme 1, experimental details, XANES and CD spectra, computational details, and calculated UV-vis and CD spectra. See DOI: https://doi.org/10.1039/d3cc05534e

functionalization thanks to the additional coordination to the axial sites of the octahedral Co centres. When the complex has an achiral ligand such as NHMe₂, the left-handed (*M*)-[LCo₃(NHMe₂)₆]³⁺ and right-handed (*P*)-[LCo₃(NHMe₂)₆]³⁺ are enantiomeric and their ratio is 50:50 at equilibrium (Scheme 1). This achiral ligand can be replaced with a chiral ligand X, which makes the *M* and *P* isomers diastereomeric and biased *M/P* equilibrium. Indeed, this ligand exchange reaction smoothly proceeds upon the addition of chiral ligands such as ((R) or (S))-1-phenylethylamine (denoted as **A1**) and ((R) or (S))-3-amino-1,2-propanediol (denoted as **A2**), to give the *M/P* diastereomeric pair of the complexes, which exhibits intense circular-dichroism (CD) signals due to the shifted *M/P* equilibrium in a solution.[9] We expected that this additional coordination on the Co centre could be utilized both for the fixation of the helical structures of the [LCo₃X₆] complexes and for the equilibrium shift caused by the ligand exchange.[9]

A solid surface has been used for the support of functional compounds, and the attachment of metal complexes on a solid surface has been widely applied for functional materials such as heterogeneous catalysts and sensors.[12–15] The chemical attachment of a metal complex to a solid surface is designed by the chemical reaction of a metal complex precursor and functional groups that originally existed on the surface (*e.g.* silanol (Si–OH)) or functional groups that are grafted on the surface.[16,17] The attachment of metal complexes on solid surfaces often produces unique metal coordination structures, which are different from precursor complexes. This process is also expected to have a great influence on controlling the conformation of the helical structures of metal complexes.

In this paper, we investigated the fixation of the dynamic *M/P* invertible helicity of [LCo₃(NHMe₂)₆](OTf)₃ on a SiO₂ surface on which ethylenediamine was functionalized (**en-SiO₂**). The addition of the chiral amine ligands ((R) or (S))-**A1** successfully induced the deviation of the dynamic helicity of the Co₃ complex at the surface. We demonstrated that the *M/P* equilibration was almost suppressed on the solid surface even in the

Fig. 2 (A) Calculated (a) UV-vis and (b) CD spectra of (M)-$[LCo_3(NHMe_2)_6]^{3+}$ (black) and (M)-$[LCo_3((R)$-**A1**$)_6]^{3+}$ (red) (in acetonitrile). (B) Calculated (a) UV-vis and (b) CD spectra of (M)-$[LCo_3(R)$-**A1**$)_6]^{3+}$ (black) and (M)-$[LCo_3en((R)$-**A1**$)_5]^{3+}$ (red) (in vacuum). (C) Calculated (a) UV-vis and (b) CD spectra of the modelled structure of **LCo₃((R)-A1)/en-SiO₂_calc** ((M) configuration) (black) and **LCo₃((R)-A1)/en-SiO₂_calc** ((P) configuration) (red) (in vacuum). (D) The summed simulated CD spectrum of the modelled structure of **LCo₃((R)-A1)/en-SiO₂_calc** ((M) configuration) and **LCo₃((R)-A1)/en-SiO₂_calc** ((P) configuration) with the ratio of 55/45. Abbreviated schematic structures (see ESI† and Fig. S2 for details) of the optimized trinuclear cobalt complexes are also presented. The torsion angles of each saloph moiety of the optimized structures under vacuum are also shown as ω.

Then, the addition of chiral (R)- or (S)-**A1** was conducted on **LCo₃/en-SiO₂**. 10 eq. (to [LCo₃] unit) of (R)- or (S)-**A1** was impregnated with **LCo₃/en-SiO₂** in methanol, and the DR UV-vis, Co K-edge XANES, and DR CD spectra of the obtained samples (denoted as **LCo₃((R)** or **(S)-A1)/en-SiO₂**) were measured. The DR UV-vis and Co K-edge XANES spectra of **LCo₃((R)** or **(S)-A1)/en-SiO₂** were almost similar to those of **LCo₃/en-SiO₂** as shown in Fig. 1(A)(c) and Fig. S1 (ESI†), indicating that the [LCo₃] unit was maintained after the coordination of the chiral **A1**.

On **LCo₃/en-SiO₂** before the addition of the chiral **A1**, its DR CD spectrum was silent, suggesting that the amounts of the P and M conformations were comparable on the SiO₂ surface (Scheme 1). In contrast, significant CD peaks at 645 nm, which were the positive/negative mirror peaks, were observed after the

addition of (R)-/(S)-**A1**, as shown in Fig. 1(B). The CD peak top position of **LCo₃((R)-A1)/en-SiO₂** was different from the peak top position observed in its UV-vis spectrum (491 nm) (Fig. 1(A and B)). The biased M/P equilibrium of [LCo₃] unit was suggested but the CD spectra of **LCo₃((R)-A1)/en-SiO₂** were remarkably different from those of the $[LCo_3((R)$-**A1**$)_6](OTf)_3$ in methanol (Fig. S4(E), ESI†).[9]

We conducted DFT calculations to look for the possible origin of the unique CD spectra of **LCo₃((R)** or **(S)-A1)/en-SiO₂**. We modelled the plausible coordination structures of the SiO₂-supported [LCo₃] unit and simulated their UV-vis and CD spectra (see ESI†). First, the two helical structures of (M) and (P)-$[LCo_3(NHMe_2)_6]^{3+}$ were optimized, and the optimized structures agreed well with those obtained from single crystal X-ray

diffraction (Fig. S3, ESI†). The TD-DFT calculations of the excited states of the complexes provided the simulated UV-vis and CD spectra of $[LCo_3(NHMe_2)_6]^{3+}$ (Fig. 2(A) and Fig. S3, ESI†), which suitably reproduced the experimentally observed UV-vis spectra of (M) and (P)-$[LCo_3(NHMe_2)_6]^{3+}$. $[LCo_3]$ unit at which all NHMe_2 ligands were substituted to (R)-A1, (M)-$[LCo_3((R)$-A1$)_6](OTf)_3$ and (P)-$[LCo_3((R)$-A1$)_6](OTf)_3$ (Fig. S4, ESI†), were also optimized, and the calculated UV-vis and CD spectra of (M)-$[LCo_3((R)$-A1$)_6](OTf)_3$,[9] showed similar trends to those experimentally obtained (Fig. 2(B)).

Then, ligand substitution from A1 to ethylenediamine at $[LCo_3]$ unit was investigated. (M)-$[LCo_3en((R)$-A1$)_5]^{3+}$ was designed by using $H_2N(CH_2)_2NH_2$ as the model of the ethylenediamine site on SiO_2 and coordinating it to one of the Co sites in $[LCo_3]$ unit (Fig. 2). The theoretical UV-vis and CD spectra of (M)-$[LCo_3((R)$-A1$)_6]^{3+}$, (M)-$[LCo_3en((R)$-A1$)_5]^{3+}$, and (M)-$[LCo_3(NHMe_2)_6]^{3+}$ were found to be almost similar to each other (Fig. 2(B)(a and b)). The ligand exchange to (R)-A1 and the ethylenediamine unit on SiO_2 was suggested not to give significant changes in their UV-vis and CD spectra.

Finally, we investigated the interaction from the surface Si–OH group on $LCo_3((R)$-A1$)$/en-SiO_2. The structural models of $LCo_3((R)$-A1$)$/en-SiO_2 with the coordination of ethylenediamine and Si–O$^-$ (H$^+$ was removed from Si–OH.) were calculated (denoted as $LCo_3((R)$-A1$)$/en-SiO_2_calc). By comparison between the calculated UV-vis spectra of (M)-$[LCo_3((R)$-A1$)_6]^{3+}$, (M)-$[LCo_3en((R)$-A1$)_5]^{3+}$, and (M)-$LCo_3((R)$-A1$)$/en-SiO_2_calc (Fig. 2(C)(a)), the peak intensity at around 480 nm of (M)-$[LCo_3((R)$-A1$)_6]^{3+}$ and (M)-$[LCo_3en((R)$-A1$)_5]^{3+}$ became small on that of (M)-$LCo_3((R)$-A1$)$/en-SiO_2_calc. In the CD spectra of (M)- and (P)-$LCo_3((R)$-A1$)$/en-SiO_2_calc (Fig. 2(C)(b)), it should be noted that the CD spectral intensities at 500–600 nm became large compared with those at 350–450 nm.

The summed CD spectra of (M)- and (P)-$LCo_3((R)$-A1$)$/en-SiO_2_calc were simulated by changing the molar ratio of (M) and (P) diastereomers (see ESI† for detailed calculation methods). It should be noted that the simulated CD spectra show a dominating peak at 500–600 nm and negligible peak at 350–450 nm and the observed CD intensities are well reproduced at the diastereomeric ratio of M/P = 55/45 (Fig. 2(D) and Fig. S5, ESI†). The estimated M/P ratio was less biased than the equilibrated ratio of M/P = 88:12 in a methanol solution (c = 0.2 mmol L^{-1}), which was experimentally observed for $[LCo_3((R)$-A1$)_6]^{3+}$.[9] Thus, the emergence of this characteristic CD spectra suggested that the M/P equilibrium shift of $[LCo_3((R)$-A1$)_6]^{3+}$ would be suppressed by the attachment to the en-SiO_2 surface in the current system. Due to this equilibrium suppression, the differences in the peak top positions of the UV-vis and CD spectra of $LCo_3((R)$-A1$)$/en-SiO_2 emerged, which were supported by the current calculations.

When an excess amount (100 eq. to $[LCo_3]$ unit) of (S) or (R)-A2 was added to $LCo_3((R)$ or (S)-A1$)$/en-SiO_2, respectively, the DR CD spectra completely inverted as shown in Fig. S6 (ESI†), and the peak tops that appeared at 564 nm and 467 nm were close to the CD spectrum of $[LCo_3((S)$ or (R)-A2$)_6](OTf)_3$ (538 nm and 467 nm) in methanol (c = 0.2 mmol L^{-1}).[9] These results

suggested that all axial A1 ligands of the $[LCo_3]$ unit were exchanged to (S) or (R)-A2, and the helicity of the $[LCo_3]$ unit was totally inverted by the ligand exchange. These results also supported that the origin of the unique induced CD spectra would be derived from the dual coordination of the ethylenediamine unit and the surface Si–O$^-$ group on the SiO_2 surface.

In summary, unique helicity induction was observed on a SiO_2 surface by the attachment of the $[LCo_3(NHMe_2)_6](OTf)_3$ complex and the addition of chiral (R) or (S)-phenylethylamine A1. Solid-state CD spectra and theoretical calculations suggested that the coordination of both the ethylenediamine site and Si–OH would be key for the immobilized helical structures, and the M/P ratio of $[LCo_3]$ unit on en-SiO_2 was found to be different from the corresponding ratio observed in a solution. The current results demonstrate that a solid surface acts as the field of fixation of dynamic helicity, and that oxide surfaces are promising to work as a tuner of the M/P ratio of the dynamic helical complexes.

Conflicts of interest

There are no conflicts to declare.

Notes and references

1 T. Mori, *Chem. Rev.*, 2021, **121**, 2373.
2 A. Tsurusaki and K. Kamikawa, *Chem. Lett.*, 2021, **50**, 1913.
3 N. Ousaka and E. Yashima, *Chem. Lett.*, 2021, **50**, 320.
4 E. Yashima and K. Maeda, *Bull. Chem. Soc. Jpn.*, 2021, **94**, 2637.
5 H. Gholami, D. Chakraborty, J. Zhang and B. Borhan, *Acc. Chem. Res.*, 2021, **54**, 654.
6 R. Carr, N. H. Evans and D. Parker, *Chem. Soc. Rev.*, 2012, **41**, 7673.
7 S. Akine, M. Miyashita and T. Nabeshima, *J. Am. Chem. Soc.*, 2017, **139**, 4631.
8 S. Akine, M. Miyashita and T. Nabeshima, *Chem. – Eur. J.*, 2019, **25**, 1432.
9 Y. Sakata, S. Chiba, M. Miyashita, T. Nabeshima and S. Akine, *Chem. – Eur. J.*, 2019, **25**, 2962.
10 S. Akine, M. Miyashita and T. Nabeshima, *Inorg. Chem.*, 2021, **60**, 12961.
11 Y. Sakata, S. Chiba and S. Akine, *Proc. Natl. Acad. Sci. U. S. A.*, 2022, **119**, e2113237119.
12 J. M. Thomas, R. Raja and D. W. Lewis, *Angew. Chem., Int. Ed.*, 2005, **44**, 6456.
13 J. M. Notestein and A. Katz, *Chem. – Eur. J.*, 2006, **12**, 3954.
14 M. Tada and Y. Iwasawa, *Coord. Chem. Rev.*, 2007, **251**, 2702.
15 C. Copéret, A. Comas-Vives, M. P. Conley, D. P. Estes, A. Fedorov, V. Mougel, H. Nagae, F. Núñez-Zarur and P. A. Zhizhko, *Chem. Rev.*, 2016, **116**, 323.
16 S. Muratsugu, Z. Weng, H. Nakai, K. Isobe, Y. Kushida, T. Sasaki and M. Tada, *Phys. Chem. Chem. Phys.*, 2012, **14**, 16023.
17 S. Muratsugu, M. H. Lim, T. Itoh, W. Thumrongpatanaraks, M. Kondo, S. Masaoka, T. S. A. Hor and M. Tada, *Dalton Trans.*, 2013, **42**, 12611.