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# Yolk–shell nanostructures: synthesis, photocatalysis and interfacial charge dynamics

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Solar energy has long been regarded as a promising alternative and sustainable energy source. In this regard, photocatalysts emerge as a versatile paradigm that can practically transform solar energy into chemical energy. At present, unsatisfactory conversion efficiency is a major obstacle to the widespread deployment of photocatalysis technology. Many structural engineering strategies have been proposed to address the issue of insufficient activity for semiconductor photocatalysts. Among them, creation of yolk–shell nanostructures which possess many beneficial features, such as large surface area, efficient light harvesting, homogeneous catalytic environment and enhanced molecular diffusion kinetics, has attracted particular attention. This review summarizes the developments that have been made for the preparation and photocatalytic applications of yolk–shell nanostructures. Additional focus is placed on the realization of interfacial charge dynamics and the possibility of achieving spatial separation of charge carriers for this unique nanoarchitecture as charge transfer is the most critical factor determining the overall photocatalytic efficiency. A future perspective that can facilitate the advancement of using yolk–shell nanostructures in sophisticated photocatalytic systems is also presented.

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## 1. Introduction

The past decade has witnessed an urgent demand for alternative energy, which has prompted academia and industry to devote effort to green energy technology. The sun is one of the essential elements for human life, and it can be regarded as an almost inexhaustible clean, renewable energy source. The difficulty is how to utilize this natural source efficiently. Photocatalysts are media which can convert solar energy into chemical energy. This feature makes photocatalysis competitive among different energy conversion technologies. H<sub>2</sub>, O<sub>2</sub> and



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hydrocarbons, which can be obtained by photocatalytic water splitting<sup>1,2</sup> and CO<sub>2</sub> reduction,<sup>3,4</sup> respectively, are considered as potential alternative fuels to replace fossil fuels. No matter how these alternative fuels are consumed, the products after usage are clean and eco-friendly, complying with the core concept of sustainable development. In addition to producing green fuels, photocatalysis can also be operated to conduct wastewater treatment<sup>5,6</sup> and environmental remediation.<sup>7,8</sup> Using photocatalysts to decompose organic substances (*e.g.* antibiotics, dyes, bacteria) and inorganic toxicants (*e.g.* heavy metal ions, salt anions) is particularly attractive as it provides a feasible solution to the headache of the ever-rising universal environmental change. Moreover, photocatalysis can mediate the high costs and complicated processing in the manufacture of value-added chemicals.<sup>9,10</sup> Continuous effort has been devoted to the production of high value-added organic compounds by photocatalytic reactions.

The mechanism of photocatalysis can be divided into three primary steps. The first is to apply light irradiation with sufficient energy to semiconductor photocatalysts; the second is the generation of electron-hole pairs and the subsequent charge

carrier transfer; the third comprises the occurrence of surface redox reactions mediated by photoexcited electrons and holes. From a thermodynamics point of view, the photoexcited electrons must have a sufficiently high reducing power in order to manage the desired reduction reactions. In other words, the conduction band (CB) level of the selected photocatalyst must be located at a potential more cathodic than the potential required for the reduction reactions. Similarly, the valence band (VB) level of semiconductor should lie in an adequately anodic potential to endow photoexcited holes with enough oxidizing power capable of conducting oxidation reactions. During the photocatalytic process, there however exist many handicaps that can result in low photocatalytic activity. These disservices include poor light absorption, rapid electron-hole recombination, and sluggish charge transfer kinetics. As dragged by such feet of clay, only a limited amount of charge carriers can actually participate in surface redox reactions. This dilemma could be ameliorated by designing composite photocatalysts with well-defined structural and compositional heterogeneity, in which the attributes of individual components can be combined. Much significantly, integrating multiple components to create heterojunction can modify electronic band structure, rectifying charge transfer behavior to increase carrier utilization efficiency. Till now, various composite nanostructures have been proposed as the photocatalyst paradigm to construct superior photocatalytic systems. Representative examples are doped/alloyed nanocrystals,<sup>11–16</sup> particle-decorated nanostructures,<sup>17–40</sup> heterodimers,<sup>41</sup> core-shell nanocrystals,<sup>42–47</sup> and yolk-shell nanostructures.<sup>48,49</sup> Here, yolk-shell nanostructures stand for a newly emerging photocatalyst platform, which are composed of a movable core surrounded by a permeable shell with void space. Such a peculiar architecture possesses several noteworthy advantages, including confined space which can facilitate the diffusion of reacting species, and large surface area which can provide abundant active sites. Besides, the incident light may undergo multiple reflections inside the void space and thus conduce to photon harvesting. Most



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importantly, the electronic band alignment between yolk and shell components can induce effective separation of photoexcited electrons and holes, which is an indispensable feature for superior photocatalytic efficiency.

Considerable effort in recent years has been made to devise feasible synthetic approaches to yolk–shell nanostructures and demonstrate their practice in photocatalytic applications. As illustrated in Scheme 1, this review gives a comprehensive retrospection on the preparation and photocatalytic applications of yolk–shell nanostructures. Additional focus is placed on the realization of interfacial charge dynamics and the possibility of achieving spatial separation of charge carriers for this unique nanoarchitecture as charge transfer is the most critical factor determining the overall photocatalytic efficiency. The review ends with a summary and future perspective that can facilitate the advancement of using yolk–shell nanostructures in sophisticated photocatalytic systems. For readers who are interested in acquiring the development of photocatalysts technology based on other types of composite nanostructures, we recommend some recent review articles.<sup>50–58</sup>

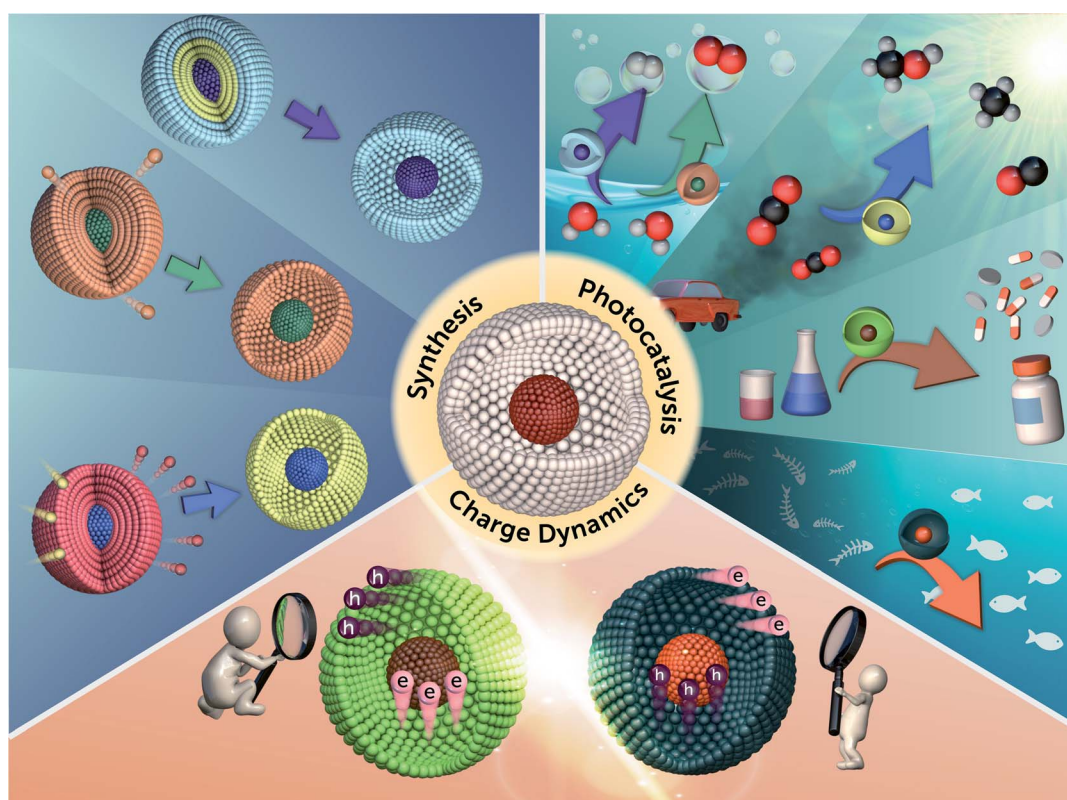
## 2. Synthetic approach

Base on the mechanism of void space formation, the synthetic approach to yolk–shell nanostructures can be divided into five categories. Kirkendall effect and galvanic replacement are both relevant to ion exchange, but the former is derived from different solubility products and the latter is caused by distinct

redox potentials. Ostwald ripening is a crystal growth process involving the dissolution of particles with smaller size for deposition on larger particles. This process can be exploited to produce void space as the dissolved particles contract. Chemical etching creates void space at the expense of sacrificial templates by selective etching. Hollow structure can also be generated through physical volatilization or oxidative decomposition by means of thermal treatment. The following section further describes the details of these synthetic systems.

### 2.1 Kirkendall effect

Diffusion couple has been widely exploited to synthesize hollow structure. On account of the different diffusion rates, voids can be generated as two ions encounter, which is known as Kirkendall effect. At present, plenty of works have adopted this principle to prepare for yolk–shell nanostructures with various compositions, e.g. Se–SnO<sub>2</sub>,<sup>59</sup> Sn/Sn<sub>2</sub>Co<sub>3</sub>–CoSnO<sub>3</sub>/Co<sub>3</sub>O<sub>4</sub>,<sup>60</sup> Pt–CoO,<sup>61</sup> Au–Fe<sub>3</sub>O<sub>4</sub>,<sup>62</sup> and Au–CdS.<sup>63</sup> Chiu *et al.* successfully demonstrated the growth of Au–CdS yolk–shell nanostructures by Kirkendall effect.<sup>63</sup> The experimental procedure is shown in Fig. 1A. The citrate-protected Au nanoparticles (NPs) were first subjected to coordination of Cu<sup>2+</sup> ions *via* the surface carboxyl groups. A uniform layer of Cu<sub>2</sub>O can then be coated on each of the Au NPs to produce Au–Cu<sub>2</sub>O core–shell NP template. Next, Na<sub>2</sub>S was added to the Au–Cu<sub>2</sub>O suspension, transforming Cu<sub>2</sub>O shell into Cu<sub>7</sub>S<sub>4</sub> as a result of the lower solubility product of Cu<sub>7</sub>S<sub>4</sub> than Cu<sub>2</sub>O. In addition to the compositional change,



Scheme 1 Schematic illustration for the synthesis, photocatalysis and charge dynamics of yolk–shell nanostructures.





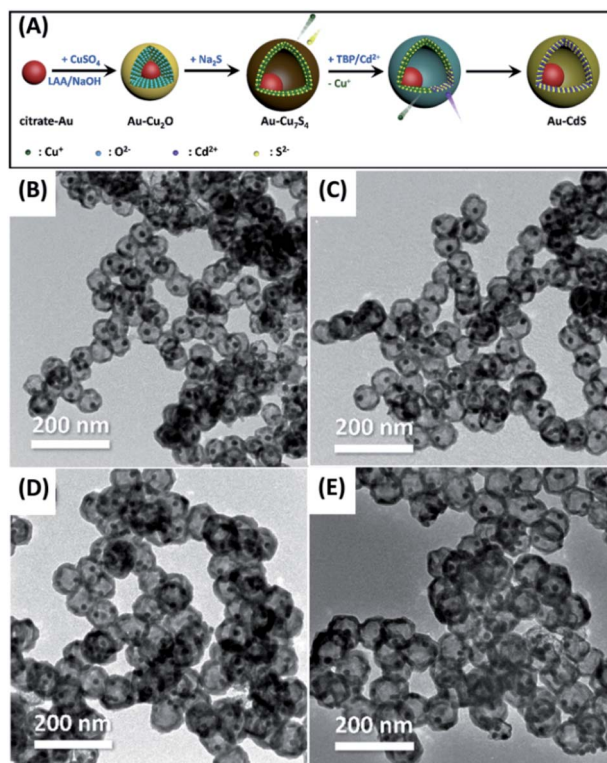


Fig. 1 (A) Synthetic approach to Au-CdS yolk-shell nanostructures by Kirkendall effect. (B-E) TEM images of Au-CdS with increasing void sizes. Reproduced with permission from ref. 63. Copyright 2019, Elsevier.

Kirkendall voids were also generated at the shell layer by virtue of the faster outward diffusion of  $\text{Cu}^{2+}$  relative to the inward diffusion of  $\text{S}^{2-}$ . Lastly, cation exchange was performed to convert  $\text{Cu}_7\text{S}_4$  to CdS, producing Au-CdS yolk-shell nanostructures. Notably, changing the Au NP/ $\text{Cu}^{2+}$  ratio during the synthesis of Au-Cu<sub>2</sub>O template can produce Au-CdS with controllable void size. Fig. 1B-E shows the TEM images of Au-CdS with increasing void sizes from  $40.2 \pm 3.3$  nm to  $64.0 \pm 6.3$  nm. These yolk-shell nanostructures all possessed a constant shell thickness of around  $11.5 \pm 0.4$  nm. Overall, Kirkendall effect represents a potentially universal strategy for the synthesis of yolk-shell nanostructures. As long as the difference in solubility product is deliberately considered, yolk-shell nanostructures with desired compositions can be obtained.

## 2.2 Galvanic replacement

By choosing two elements with suitable oxidation and reduction potentials, the replacement reaction can occur spontaneously. This principle is generally used to fabricate yolk-shell nanostructures with a wide range of compositions, including Ni-SnO<sub>2</sub>/Ni<sub>3</sub>Sn<sub>2</sub>,<sup>64</sup> Co-Au,<sup>65</sup> Au-AgPd,<sup>66</sup> Sb-C<sup>67</sup> and Pd-M<sub>x</sub>Cu<sub>1-x</sub> (M = Au, Pd, and Pt).<sup>68</sup> Liu *et al.* ingeniously designed hollow Sb-C yolk-shell spheres from Sn-C yolk-shell structures *via* galvanic replacement.<sup>67</sup> Fig. 2A displays the synthetic process. Hollow SnO<sub>2</sub> nanospheres were first obtained in the hydrothermal

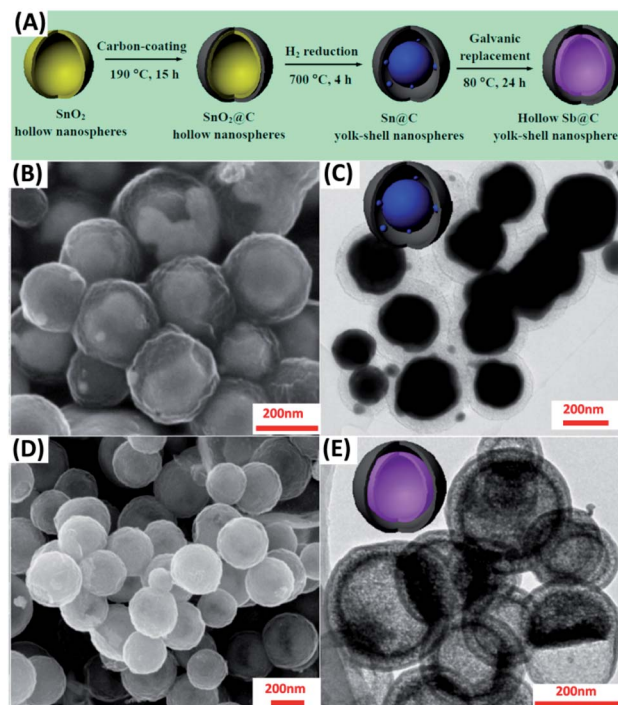


Fig. 2 (A) Synthetic approach to Sb-C yolk-shell nanospheres by galvanic replacement. (B and C) SEM and TEM images for Sn-C. (D and E) SEM and TEM images for Sb-C. Reproduced with permission from ref. 67. Copyright 2017, American Chemical Society.

reaction, followed by the uniform coating of a carbon shell. Further annealing under H<sub>2</sub> can convert hollow SnO<sub>2</sub> to solid Sn particles with empty space left between carbon shell and Sn core. Because Sb has a larger reducing potential than Sn ( $E_{\text{Sb}^{3+}/\text{Sb}} = 0.241$  V *vs.* SHE;  $E_{\text{Sn}^{4+}/\text{Sn}} = 0.0129$  V *vs.* SHE), adding Sb<sup>3+</sup> can further transform Sn-C to Sb-C. In this galvanic replacement process, Sn was oxidized to Sn<sup>4+</sup>, while Sb<sup>3+</sup> was reduced to Sb; meanwhile, structural transformation from solid Sn to hollow Sb occurred owing to the non-equilibrium of ion diffusion between outward Sn<sup>4+</sup> (faster) and inward Sb<sup>3+</sup> (slower). The microstructure of Sn-C and Sb-C was explored with scanning electron microscopy (SEM) and transmission electron microscopy (TEM). As Fig. 2B-E show, Sb-C inherited the size and morphology from the parent Sn-C. The solid Sn core in Sn-C and hollow Sb yolk in Sb-C can also be clearly identified. Analogous to Kirkendall effect, galvanic replacement is a general approach to the synthesis of yolk-shell nanostructures. By pairing up suitable redox potentials, yolk-shell nanostructures with a wide range of materials combination can be readily obtained.

## 2.3 Ostwald ripening

Ostwald ripening can account for the growth of yolk-shell nanostructures possessing single composition (homogeneous) or multiple compositions (heterogeneous). The formation of homogeneous yolk-shell structures involves the dissolution of a core particle and re-deposition of a shell layer on the gradually dissolved core. Heterogeneous yolk-shell structures on the













demonstrated remarkable photocatalytic activity toward CO<sub>2</sub> reduction, including Au-TiO<sub>2</sub>,<sup>94</sup> Au-g-C<sub>3</sub>N<sub>4</sub>/SnS,<sup>95</sup> Ni-SiO<sub>2</sub>,<sup>96</sup> TiO<sub>2</sub>,<sup>97</sup> and TiO<sub>2-x</sub>.<sup>98</sup>

Ziarati *et al.* obtained a yolk-shell nanostructure of TiO<sub>2</sub> (Y@S-TiO<sub>2</sub>) based on the Ostwald ripening approach in a solvothermal reaction.<sup>98</sup> After hydrogenation treatment, the TiO<sub>2-x</sub> yolk-shell nanostructures (Y@S-TiO<sub>2-x</sub>) were decorated with Co-Al layered double hydroxides (LDHs), producing Y@S-TiO<sub>2-x</sub>/LDH core/shell nanospheres. Here, hydrogenation can create oxygen vacancies for TiO<sub>2</sub>, which not only reduced bandgap of TiO<sub>2</sub>, but also served as adsorption sites for CO<sub>2</sub>. The decoration of Co-Al LDHs on the other hand can improve charge separation of TiO<sub>2-x</sub>. The porous features of LDHs can also promote CO<sub>2</sub> adsorption. These attributes were all beneficial for increasing photocatalytic activity of CO<sub>2</sub> reduction. Fig. 8A and B clearly reveal the yolk-shell structural features for TiO<sub>2-x</sub> and the successful decoration with Al-Co LDHs. Fig. 8C compares the results of photocatalytic CO<sub>2</sub> reduction among different samples in terms of CH<sub>3</sub>OH production. Several important points can be observed. First, P25 were almost inactive due to the poor visible light absorption. Second, P25/LDH showed better photocatalytic activity than P25 and LDH did, highlighting the synergy of LDHs with TiO<sub>2</sub> for promoting the efficiency of CO<sub>2</sub> reduction. Third, hydrogenated TiO<sub>2</sub> showed much higher activity than non-treated TiO<sub>2</sub> (H-P25/LDH *vs.* P25/LDH and Y@S-TiO<sub>2-x</sub>/LDH *vs.* Y@S-TiO<sub>2</sub>/LDH). This revealed the importance of hydrogenation treatment for promoting CO<sub>2</sub> reduction activity of TiO<sub>2</sub>. Compared with P25 and LDH, a 20-fold increase in CH<sub>3</sub>OH production can be achieved on Y@S-TiO<sub>2-x</sub>/LDH. Such a superior performance can be realized by the attributes associated with structural engineering and charge transfer as illustrated in Fig. 8D. It is interesting to note that CH<sub>3</sub>OH was the main product at the first

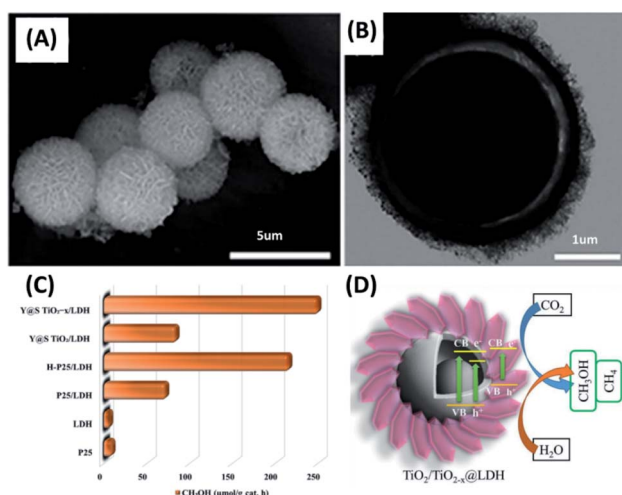


Fig. 8 (A and B) SEM and TEM images of Y@S-TiO<sub>2-x</sub>/LDH. (C) Comparative results of CO<sub>2</sub> reduction in terms of CH<sub>3</sub>OH production under visible light illumination. (D) Plausible charge transfer mechanism for superior photocatalytic efficiency of Y@S-TiO<sub>2-x</sub>/LDH toward CO<sub>2</sub> reduction. Reproduced with permission from ref. 98. Copyright 2019, American Chemical Society.

2 h of photocatalytic reaction. Further extending the reaction time led to the gradual production of CH<sub>4</sub> in addition to CH<sub>3</sub>OH. This finding illustrated that the selectivity of CO<sub>2</sub> reduction on Y@S-TiO<sub>2-x</sub>/LDH could be adjusted by controlling the reaction time. It should be noted that photocatalytic CO<sub>2</sub> reduction on yolk-shell nanostructures is at the infant stage, making it difficult to assess the actual performance. Quantitative indexes, *e.g.* apparent quantum yield (AQY) and solar-to-fuel conversion efficiency, enabling the reliable evaluation and global comparison of CO<sub>2</sub> reduction activity for different yolk-shell nanostructure systems are indispensable.

### 3.4 H<sub>2</sub> production

H<sub>2</sub> production on semiconductor photocatalysts from water splitting represents a viable, green approach to creation of alternative energy. Ever since Fujishima and Honda demonstrated water photolysis on TiO<sub>2</sub> photoelectrodes,<sup>99</sup> photocatalytic H<sub>2</sub> production has been an ultimate goal of solar energy conversion in the search for clean energy. There are many strategies developed to enhance the efficiency of solar H<sub>2</sub> production. Employing unique nanostructures to better harvest solar energy is particularly important because light absorption is the leading process determining how many charge carriers can be produced. In this regard, yolk-shell nanostructures offer an ideal structural platform since light harvesting efficiency can be enhanced as a result of multi-reflection of incident light within the void space. Till now, a variety of yolk-shell nanostructures have demonstrated remarkable activity of photocatalytic H<sub>2</sub> production, such as Au-TiO<sub>2</sub><sup>100,101</sup> and Au-CdS.<sup>63,102</sup>

Shi *et al.* prepared for Au-TiO<sub>2</sub> yolk-shell nanospheres with a chemical etching method.<sup>100</sup> These nanospheres were assembled into three-dimensional arrays (3D-array) to study the multi-scattering effect on the photocatalytic activity of H<sub>2</sub> production. Fig. 9A depicts the charge transfer mechanism for Au-TiO<sub>2</sub>. The plasmonic Au yolk can be excited by visible light, generating hot electrons for enhancing the photoactivity of TiO<sub>2</sub>. As Fig. 9B shows, the as-prepared Au-TiO<sub>2</sub> was highly

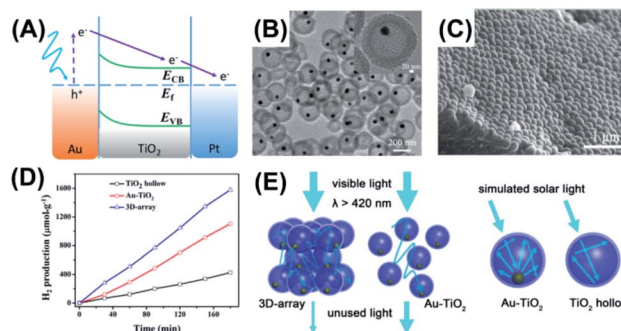


Fig. 9 (A) Plausible charge transfer mechanism for superior photocatalytic efficiency of Au-TiO<sub>2</sub> toward H<sub>2</sub> production, (B) TEM image of Au-TiO<sub>2</sub>, (C) SEM image of 3D-array, (D) comparative results of H<sub>2</sub> production under simulated daylight illumination, (E) schematic illustration of multi-scattering effect and increased light reflection for 3D-array and Au-TiO<sub>2</sub>. Reproduced with permission from ref. 100. Copyright 2016, ACS publications.





mono-dispersed with a uniform shell thickness of 46 nm and inner cavity size of 176 nm. The SEM image in Fig. 9C clearly reveals the orderly packed, hexagonal structure for 3D-array. These arrays comprised several layers of Au–TiO<sub>2</sub>, which may facilitate light scattering between the individual Au–TiO<sub>2</sub> spheres. Fig. 9D compares the photocatalytic activity of H<sub>2</sub> production among three relevant samples. Compared with hollow TiO<sub>2</sub>, both 3D-array and non-assembled Au–TiO<sub>2</sub> showed much higher activity of H<sub>2</sub> production, manifesting the beneficial function of the plasmonic Au. Significantly, 3D-array showed 1.4 folds higher activity than non-assembled Au–TiO<sub>2</sub>. As Fig. 9E illustrates, the superior activity of 3D-array can be attributed to the improved photon harvesting resulting from the multi-scattering of incident light between the individual Au–TiO<sub>2</sub>. On the other hand, photon harvesting can also be augmented in each of the yolk–shell particles because incident light can be reflected many times within the void space. As a result of the increased light reflection inside the TiO<sub>2</sub> shell, the plasmonic Au yolk can generate more hot electrons for improving the photocatalytic activity. The superior photocatalytic performance of yolk–shell nanostructures toward H<sub>2</sub> production can be reflected from the unprecedented AQY value. In Chiu's work,<sup>63</sup> Au–CdS yolk–shell nanostructures achieved an AQY value of 4.22% at 420 nm without the use of Pt co-catalysts. This value was substantially larger than that of Pt-deposited Au–CdS core–shell nanostructures (1.9%),<sup>103</sup> highlighting the great promise of yolk–shell nanostructures for photocatalytic H<sub>2</sub> production.

### 3.5 O<sub>2</sub> production

Compared to H<sub>2</sub> production, there are relatively few kinds of photocatalysts that exhibit valid activity of O<sub>2</sub> production from water splitting. The cause is the higher kinetics difficulty of water oxidation, which requires four equivalent holes in order to produce an O<sub>2</sub> molecule. Despite the difficulty, O<sub>2</sub> production is still an important reaction with regard to the realization of overall water splitting. The use of yolk–shell nanostructures for conducting photocatalytic O<sub>2</sub> production is quite limited. The successful example was demonstrated on Fe<sub>2</sub>TiO<sub>5</sub>–TiO<sub>2</sub>.<sup>104</sup>

Among the different semiconductor photocatalysts, TiO<sub>2</sub> and  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> are particularly suited for conducting photocatalytic O<sub>2</sub> production by virtue of the sufficiently anodic VB level ( $E_{\text{VB}} = +2.91$  eV vs. NHE for TiO<sub>2</sub>;  $E_{\text{VB}} = +2.48$  eV vs. NHE for  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>). Waqas *et al.* prepared for Fe<sub>2</sub>TiO<sub>5</sub>–TiO<sub>2</sub> yolk–shell hollow spheres (FTYS-HS) with a thermal treatment method and demonstrated effective photocatalytic O<sub>2</sub> production.<sup>104</sup> Here, Fe<sub>2</sub>TiO<sub>5</sub> has a band structure straddling that of TiO<sub>2</sub> and  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>. By combining Fe<sub>2</sub>TiO<sub>5</sub> with TiO<sub>2</sub>, the photocatalytic O<sub>2</sub> production efficiency can be enhanced as a result of the improved charge transfer.<sup>105</sup> Fig. 10A shows the high-angle annular dark-field (HAADF) scanning TEM image and the corresponding energy-dispersive spectrometry (EDS) mapping data for FTYS-HS, which revealed uniform distribution of Ti and Fe at the yolk and shell regions. Fig. 10B compares the photocatalytic activity of O<sub>2</sub> production among four relevant samples. As expected, FTYS-HS achieved a much higher activity of O<sub>2</sub>

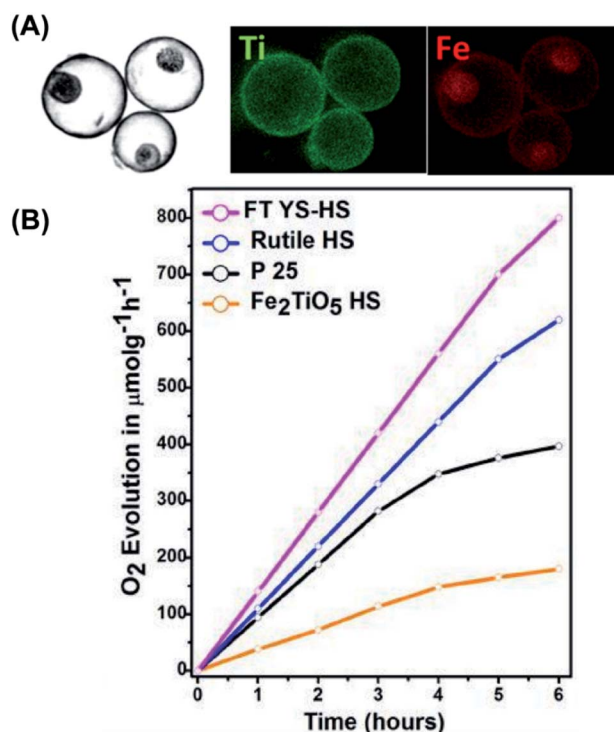


Fig. 10 (A) HAADF scanning TEM image and the corresponding EDS data of FTYS-HS, (B) comparative results of O<sub>2</sub> production under white light illumination. Reproduced with permission from ref. 104. Copyright 2017, Elsevier.

production than P25, rutile TiO<sub>2</sub> hollow spheres (rutile HS), and Fe<sub>2</sub>TiO<sub>5</sub> hollow spheres (Fe<sub>2</sub>TiO<sub>5</sub> HS). The superiority of FTYS-HS can be partly assigned to Fe<sub>2</sub>TiO<sub>5</sub>/TiO<sub>2</sub> heterojunction which improved charge carrier separation.<sup>106</sup> The multi-reflection of incident light inside the yolk–shell structures contributed to the enhanced photocatalytic activity of FTYS-HS as well. Similar to the case of CO<sub>2</sub> reduction, only a limited number of examples have demonstrated the practice of yolk–shell nanostructures in photocatalytic O<sub>2</sub> production. Future effort can be focused on the survey of more yolk–shell nanostructures candidates to construct photocatalytic systems for efficient O<sub>2</sub> production.

### 3.6 Fine chemicals synthesis

Fine chemicals synthesis by means of photocatalytic redox reactions is a prosperous topic attracting interdisciplinary scientists. Traditionally, fine chemicals are manufactured under harsh conditions, such as high temperature and highly oxidative environment. Although high yield with a good selectivity can be achieved, the traditional approach is costly and harmful, and therefore, is not viable from the economic and environmental points of view. As an alternate, using photocatalysts to synthesize fine chemicals has received increasing attention. Examples include the production of aldehydes by photocatalytic oxidation<sup>107–110</sup> and the generation of amines by photocatalytic reduction.<sup>111–114</sup> Yolk–shell nanostructures have



also demonstrated remarkable photocatalytic activity for synthesis of fine chemicals.<sup>115–120</sup>

In a study by Ziarati *et al.*, yolk-shell TiO<sub>2</sub> structures (Y@S-TiO<sub>2</sub>) were prepared in a solvothermal reaction *via* Ostwald ripening.<sup>118</sup> The obtained Y@S-TiO<sub>2</sub> was further modified with di-amines to wrap graphene sheets in a high level. The resultant graphene highly wrapped yolk-shell TiO<sub>2</sub> (G-HW-Y@S-TiO<sub>2</sub>) architectures showed good selectivity of producing aldehydes in the photocatalytic oxidation of aromatic alcohols. Fig. 11A–C show the SEM and TEM images for G-HW-Y@S-TiO<sub>2</sub>. Successful wrapping of 5 nm-thick graphene sheets can be recognized. In Fig. 11D, the photocatalytic efficiency of G-HW-Y@S-TiO<sub>2</sub> toward selective oxidation of benzyl alcohol was investigated and compared with other counterpart samples, including graphene wrapped TiO<sub>2</sub> (G-W-TiO<sub>2</sub>), graphene wrapped hollow TiO<sub>2</sub> (G-W-H-TiO<sub>2</sub>) and graphene lightly wrapped yolk-shell TiO<sub>2</sub> (G-W-Y@S-TiO<sub>2</sub>). G-W-TiO<sub>2</sub> showed mediocre activity with a moderate conversion rate (47%) and benzaldehyde selectivity (61%). For G-W-H-TiO<sub>2</sub>, the conversion rate (78%) and

selectivity (94%) can be largely increased. Significantly, G-W-Y@S-TiO<sub>2</sub> and G-HW-Y@S-TiO<sub>2</sub> showed nearly 100% of benzaldehyde selectivity, with G-HW-Y@S-TiO<sub>2</sub> exhibiting a little bit higher conversion. The superior performance of G-HW-Y@S-TiO<sub>2</sub> for selective oxidation of benzyl alcohol can be attributed to the synergy among several factors. Note that graphene had a work function (4.42 eV) larger than TiO<sub>2</sub> (4.21 eV).<sup>121</sup> The photoexcited electrons of TiO<sub>2</sub> were then transferred to graphene, leaving photoexcited holes at TiO<sub>2</sub>. Therefore, the photocatalytic oxidation reaction mostly occurred at the TiO<sub>2</sub> surface of G-HW-Y@S-TiO<sub>2</sub>. As illustrated in Fig. 11E, the hydrophilic alcohol reactant first diffused through the mesoporous shell, reacted with the photoexcited holes at TiO<sub>2</sub> yolk, and produced benzyl alcohol cation radicals. At the same time, the photoexcited electrons at graphene reacted with the dissolved O<sub>2</sub> and generated <sup>•</sup>O<sub>2</sub><sup>-</sup>. The cation radicals further reacted with <sup>•</sup>O<sub>2</sub><sup>-</sup> and produced benzaldehydes. Such a partially oxidized product was hydrophobic, which can be repulsed by the hydrophilic surface of TiO<sub>2</sub> and would thus diffuse outward prior to further oxidation. The wrapped graphene sheets on the other hand can prevent the effluent benzaldehydes from direct contact with TiO<sub>2</sub>, prohibiting them from being oxidized by the photoexcited holes at TiO<sub>2</sub>. This explained why such a high selectivity can be achieved on G-HW-Y@S-TiO<sub>2</sub>.

Since the pioneering work published by Akira Suzuki in 1979,<sup>122</sup> Suzuki–Miyaura coupling reaction has been widely utilized to synthesize a variety of technologically important organic molecules, such as styrenes and biphenyls. Typical Suzuki–Miyaura coupling reactions use Pd as catalysts. Photocatalysts containing Pd are also capable of conducting efficient coupling reactions.<sup>119,123–125</sup> In the study by Rohani *et al.*,<sup>128</sup> yolk-shell structures comprising hydrogenated urchin-like TiO<sub>2</sub> shell and TiO<sub>2</sub> yolk (HUY@S-TOH) were prepared in an Ostwald ripening-based dissolution–recrystallization process. These yolk-shell structures were decorated with Au–Pd core-shell particles (HUY@S-TOH/AuPd), which were employed as photocatalysts to carry out Suzuki–Miyaura coupling reaction. In Fig. 12A and B, the TEM images clearly revealed the yolk-shell feature of HUY@S-TOH and the successful decoration of 5 nm-sized Au/Pd particles. The photocatalytic efficiency was studied by carrying out coupling reaction of 4-iodotoluene with phenylboronic acid under visible light illumination. The expected product was biphenyls. Four relevant samples, including P25 TiO<sub>2</sub>, P25/AuPd, Y@S-TOH/AuPd, HUY@S-TOH/Au, HUY@S-TOH/Pd and HUY@S-TOH/AuPd, were tested and compared. Notably, P25 TiO<sub>2</sub> showed negligible activity, while P25/AuPd was significantly active. This finding pointed out the necessity of using Au/Pd particles to conduct Suzuki–Miyaura coupling reaction. On the other hand, because of the higher surface area and enhanced photon harvesting inside the yolk-shell structures, Y@S-TOH/AuPd showed much higher activity than P25/AuPd did. Among the different samples, HUY@S-TOH/AuPd showed the best activity in terms of biphenyls yield with a turnover frequency (TOF) of 7095 h<sup>-1</sup>. As summarized in Fig. 12C, this TOF value was considerably higher than the performance of the state-of-the-art photocatalysts ever reported. Fig. 12D further depicts the charge transfer mechanism

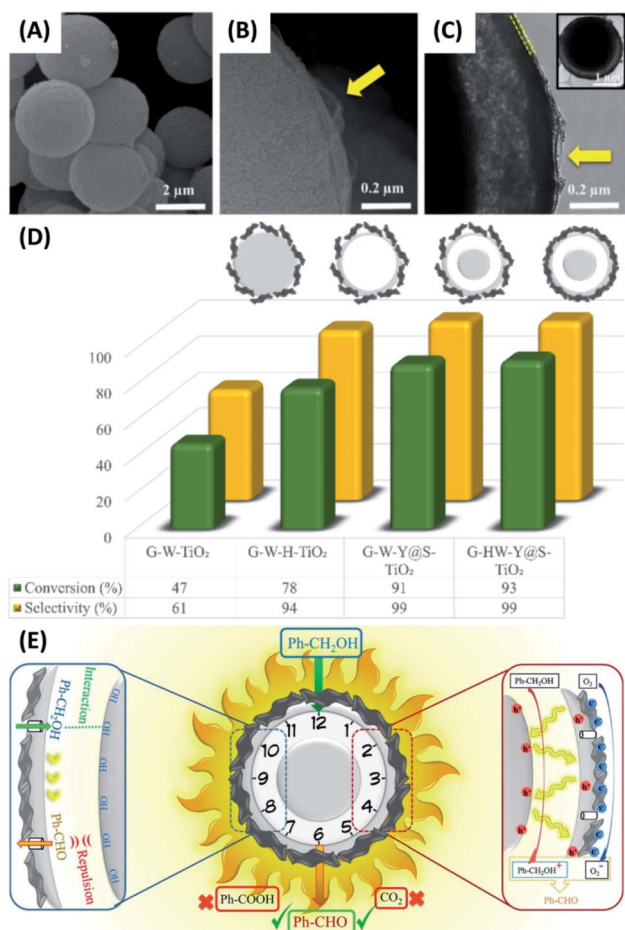


Fig. 11 (A–C) SEM and TEM images of G-HW-Y@S-TiO<sub>2</sub>. (D) Comparative results of selective benzyl alcohol oxidation under visible light illumination. (E) Plausible charge transfer mechanism for superior photocatalytic activity of G-HW-Y@S-TiO<sub>2</sub> toward selective benzyl alcohol oxidation. Reproduced with permission from ref. 118. Copyright 2018, Elsevier.





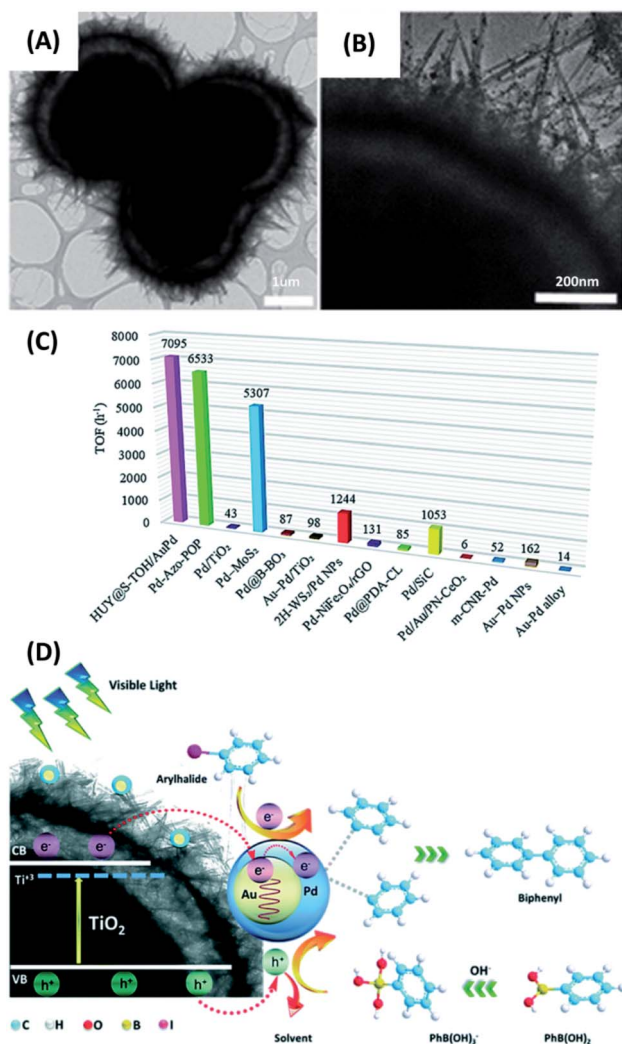


Fig. 12 (A and B) TEM images of HUY@S-TOH/AuPd. (C) Comparison of TOF value with those achieved by state-of-the-art photocatalysts reported in the literature. (D) Plausible charge transfer mechanism for superior photocatalytic activity of HUY@S-TOH/AuPd toward Suzuki–Miyaura coupling reaction. Reproduced with permission from ref. 119. Copyright 2019, Royal Society of Chemistry.

accountable for the superior activity of HUY@S-TOH/AuPd. Hydrogenated TiO<sub>2</sub> was capable of harvesting visible photons to equip TiO<sub>2</sub> with visible photoactivity. Upon visible irradiation, the photoexcited electrons of hydrogenated TiO<sub>2</sub> may transfer to Au/Pd; meanwhile, hot electrons induced by plasmonic absorption of Au were injected to Pd. The C–X bonding of aryl halides can then be activated by the electron-enriched Pd, facilitating oxidative addition process of Suzuki reaction.<sup>126</sup> In the meantime, phenylboronic acid could be oxidized by photoexcited holes. Cross-coupling reaction then took place as the oxidized phenylboronic acids encountered the activated aryl halides, producing the expected product.

## 4. Interfacial charge dynamics

The efficiency of semiconductor photocatalysts is dependent on intrinsic properties, including light absorption, carrier transfer and recombination, and kinetics of surface redox reactions. The capability of light absorption inherits from the optical properties of the chosen photocatalysts, which can be augmented by means of dopant introduction and materials combination. Charge dynamics at interface dictates charge transfer and carrier utilization, the two determinant processes involved in the photocatalytic reactions. Under light illumination, the photoexcited electrons and holes are transported from the bulk of the photocatalysts to the surface region for conducting redox reactions. These charge carriers are however inclined to be trapped and recombined before they can eventually reach surface region. This would jeopardize carrier utilization efficiency and decrease the photocatalytic performance, making interfacial charge dynamics the most critical factor determining the overall photocatalytic efficiency. An in-depth understanding of how photoexcited charge carriers can be transported, trapped and/or recombined will enable performance optimization of semiconductor photocatalysts. In particular, establishing correlations between interfacial charge dynamics and photocatalytic activity can offer empirical yet practicable knowledge for devising versatile nanostructures for superior photocatalytic applications.

Compared to the extensive studies of interfacial charge dynamics on composite nanostructures like particle-decorated nanostructures<sup>19–21,23–29,31–40</sup> and core–shell nanocrystals,<sup>43,45–47</sup> examples are relatively scarce for the topic of studying charge dynamics on yolk–shell nanostructures.<sup>63,81,101,102,127,128</sup> In this section, we introduce two representative studies that explored the correlations of interfacial charge dynamics with the photocatalytic performance of H<sub>2</sub> production on yolk–shell nanostructures (Au–CdS<sup>63</sup> and Au–TiO<sub>2</sub><sup>101</sup>). In addition to realization of charge dynamics, manipulation of charge transfer behaviors is also imperative with regard to the regulation of photocatalytic reactions. For yolk–shell nanostructures, spatial separation of photoexcited electrons and holes at the yolk and shell region is possible since the yolk is inherently separated from the shell. The spatial separation of electrons and holes not only guarantees long-lived charge separation but also inhibits backward reactions, both of which are crucial for reinforcing the utility in photocatalytic reactions. This concept has not been realized on yolk–shell nanostructures, but it has been widely demonstrated on hollow nanostructures, a structurally analogous example to yolk–shell model. Additional focus on the strategies of achieving spatial separation of charge carriers for hollow nanostructures will be placed in this section with the aim to provide applicable guidelines to yolks-shell nanostructures.

### 4.1 Charge dynamics of Au–CdS

In Chiu's study,<sup>63</sup> transient absorption (TA) spectroscopy was utilized to explore charge transfer dynamics of Au–CdS nanostructures. A definite correlation of interfacial charge dynamics and photocatalytic efficiency with the void size of Au–CdS was



established. Note that TA is a powerful ultrafast laser spectroscopic technique used to monitor the electronic transition processes of photoexcited charge carriers within semiconductor nanostructures. By identifying the photophysics beyond the recorded TA spectra, the interfacial charge dynamics of the samples can be clearly interpreted.

Fig. 13A and B show the TA spectra for pure hollow CdS and Au–CdS-4 recorded by increasing delay times from –2 to 100 ps at the probe wavelength between 450 to 750 nm. For pure CdS, a narrow negative bleaching was noticed at 475 nm, which can be ascribed to the typical ground state bleaching of charge carriers from the excited CdS. This transient bleaching was however red-shifted to 515 nm for Au–CdS-4, reflecting pronounced electronic interaction between Au yolk and CdS shell. Most importantly, Au–CdS-4 showed much quicker decay of transient bleaching than pure CdS, suggesting the prevalence of additional electron transportation pathway from CdS to Au. Fig. 13C further compares the decay kinetics of TA bleaching for pure CdS and the four CdS–Au samples. Here, a counterpart sample by physically mixing pure CdS and Au NPs (denoted as CdS + Au) was also examined in order to highlight the superior feature of yolk–shell nanostructures. These TA profiles were fitted with a biexponential function to obtain the quantitative information of charge dynamics. Two lifetime components ( $\tau_1$  and  $\tau_2$ ) can be generated from the fitting, which was respectively associated with exciton decay processes induced by shallow and deep traps of CdS. To offer a global comparison, an intensity-weighted average lifetime,  $\langle\tau\rangle$ , was presented, respectively giving 27.44 ps, 24.74 ps, 14.66 ps, 13.79 ps, 12.34 ps, and 11.71 ps for pure CdS, CdS + Au, Au–CdS-1, Au–CdS-2, Au–CdS-3, and Au–CdS-4. By considering that electron transportation from CdS to Au caused the reduced lifetime of Au–CdS, the rate constant of interfacial charge transfer ( $k_{ct}$ ) can be obtained using the equation

$$k_{ct}(\text{CdS} \rightarrow \text{Au}) = 1/\langle\tau\rangle(\text{Au–CdS}) - 1/\langle\tau\rangle(\text{pure CdS})$$

The calculated  $k_{ct}$  value was respectively  $0.40 \times 10^{10} \text{ s}^{-1}$ ,  $3.18 \times 10^{10} \text{ s}^{-1}$ ,  $3.61 \times 10^{10} \text{ s}^{-1}$ ,  $4.46 \times 10^{10} \text{ s}^{-1}$ , and  $4.90 \times 10^{10} \text{ s}^{-1}$

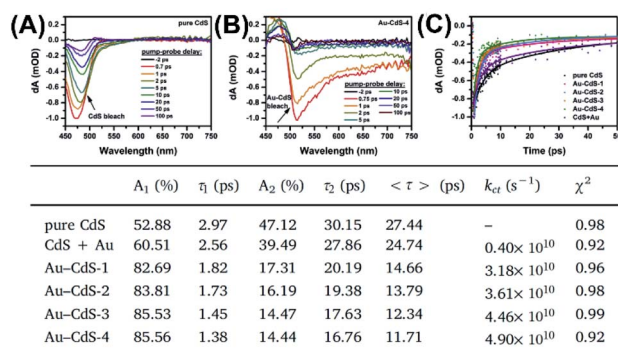


Fig. 13 (A and B) TA profiles of pure CdS and Au–CdS-4. (C) Kinetics of TA bleaching for relevant samples. Inserted table summarizes the fitting results. Reproduced with permission from ref. 63. Copyright 2019, Elsevier.

for CdS + Au, Au–CdS-1, Au–CdS-2, Au–CdS-3, and Au–CdS-4. There were two significant points to highlight. Firstly, compared to CdS + Au, the four Au–CdS samples all exhibited larger  $k_{ct}$ , signifying more efficient charge dynamics within the yolk–shell structure. Secondly, the  $k_{ct}$  of Au–CdS was enhanced as void size increased, suggesting that void size was a vital cause influencing charge transfer from CdS to Au. The implications of these two features in photocatalytic reactions were further explored. The comparative results of photocatalytic H<sub>2</sub> production showed consistency with the variation of  $k_{ct}$ , in which the four Au–CdS surpassed CdS + Au and CdS–Au-4 had the highest activity. This work highlighted the importance of understanding charge dynamics of yolk–shell nanostructures toward the optimization of photocatalytic performance.

#### 4.2 Charge dynamics of Au–TiO<sub>2</sub>

In Dillon's work,<sup>101</sup> Au–TiO<sub>2</sub> yolk–shell nanostructures with gradually improved shell crystallinity were prepared. The influence of shell crystallinity on the photophysics and its correlation with the resultant photocatalytic performance was investigated. In this study, time-resolved photoluminescence (PL) was employed to follow the dynamics of the photoexcited electrons and holes. Compared to TA, time-resolved PL offers a more intuitive viewpoint to explore the crystallinity effect on charge dynamics because PL emissions are highly sensitive to the crystallographic structure.<sup>129–131</sup>

Fig. 14A and B displays the time-resolved PL spectra recorded at 500 nm after 266 nm laser excitation at two time regimes (1 ns and 20 ns). Four Au–TiO<sub>2</sub> samples were compared, including the least crystalline sample A and the most crystalline sample D. Here, the emissions at 500 nm can be associated with charge recombination involving self-trapped excitons of TiO<sub>2</sub> or surface defects within TiO<sub>2</sub> lattice. The dynamics of this emission can thus reflect the degree of crystallinity of TiO<sub>2</sub>. At both time regimes, sample D showed considerably slower PL decay

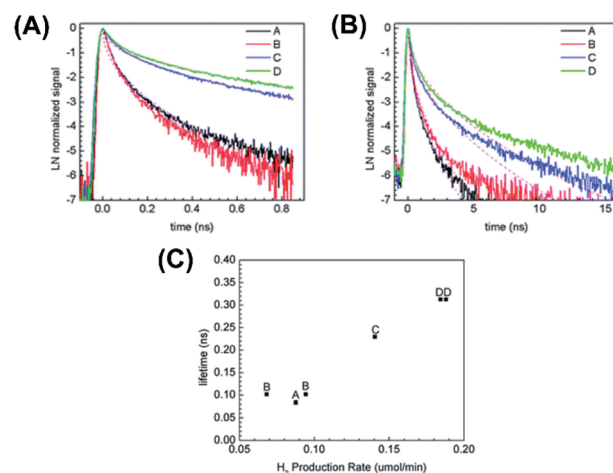


Fig. 14 (A and B) Time-resolved PL spectra of four Au–TiO<sub>2</sub> samples recorded at two time regimes. (C) Correlation of PL lifetime and H<sub>2</sub> production rate with the crystallinity of Au–TiO<sub>2</sub>. Reproduced with permission from ref. 101. Copyright 2013, Royal Society of Chemistry.





kinetics than sample A did, suggesting that better crystallinity may lead to slower charge recombination as a result of the smaller number of defect sites. These kinetics profiles can be analyzed using a stretched exponential model, from which a characteristic lifetime was received. At 20 ns time regime, the computed lifetime of sample A, B, C, D was respectively 0.084 ns, 0.10 ns, 0.23 ns, 0.31 ns. The correspondence between charge dynamics and crystallinity degree was further correlated with the photocatalytic activity of H<sub>2</sub> production. In Fig. 14C, a clear correlation of PL lifetime with H<sub>2</sub> production rate was found. This correlation demonstrated that high degree of crystallinity was conducive to an enhanced photocatalytic activity for Au-TiO<sub>2</sub> as the number of defect sites may be reduced. This work also illustrated the promise of employing time-resolved PL to evaluate the photocatalytic performance of yolk-shell nanostructures.

### 4.3 Spatial separation of charge carriers

The concept of spatially separating photoexcited electrons and holes stems from the introduction of co-catalysts that have been widely employed to reduce the overpotentials of photocatalytic reactions. There are two types of co-catalysts used in photocatalysis; one functions as an electron trap site due to its large work function, such as Pt, Ag, Au particles (also known as reduction co-catalysts),<sup>86,120,132,133</sup> and the other one serves as a hole trap site, such as IrO<sub>2</sub>, PbO<sub>x</sub>, PbS, MnO<sub>x</sub>, CoO<sub>x</sub> (also known as oxidation co-catalysts).<sup>120,133–138</sup> The introduction of co-catalysts can anchor electrons or holes to guarantee long-lived charge separation.<sup>120,123,125,133</sup> Significantly, simultaneous loading of dual co-catalysts makes possible the spatial separation of electrons from holes, which can be utilized to achieve the concurrent occurrence of redox reactions at different sites. However, random deposition of dual co-catalysts may lead to the interference of redox reactions, arousing charge recombination and backward reactions to deteriorate the overall photocatalytic efficiency. Therefore, introducing dual co-catalysts at separate sites of photocatalysts is indispensable to achieving spatial separation of charge carriers.

The introduction of dual co-catalysts at separate sites has been realized on solid nanostructures with well-defined exposed facets, in which site-selective deposition of co-catalysts can be accomplished *via* photo-deposition.<sup>133</sup> For yolk-shell and hollow nanostructures which inherently have two separate sites, *i.e.* the inner and outer surfaces, selective deposition of dual co-catalysts at these surfaces can be more readily achieved. Domen's group demonstrated a successful practice of spatially selective deposition of dual co-catalysts on Ta<sub>3</sub>N<sub>5</sub> hollow spheres, in which reduction (Pt) and oxidation (IrO<sub>2</sub>) co-catalyst was respectively loaded at the inner and outer surface.<sup>134</sup> Upon light irradiation, the photoexcited electrons were attracted by Pt, while the photogenerated holes were trapped by IrO<sub>2</sub>. As a result of the spatial separation of electrons and holes, charge recombination was greatly restrained, leading to a much enhanced photocatalytic activity of overall water splitting. In Zheng's work, hollow C<sub>3</sub>N<sub>4</sub> structures were simultaneously deposited with Pt and Co<sub>3</sub>O<sub>4</sub> co-catalysts to

demonstrate the effectiveness of photocatalytic overall water splitting.<sup>138</sup> The spatial separation of Pt (at inner surface of C<sub>3</sub>N<sub>4</sub>) from Co<sub>3</sub>O<sub>4</sub> (at outer surface of C<sub>3</sub>N<sub>4</sub>) not only improved charge separation, but also inhibited the further reduction of O<sub>2</sub> at Pt. As Fig. 15 compares, a stoichiometric ratio of the produced H<sub>2</sub> to O<sub>2</sub> can be attained as Pt and Co<sub>3</sub>O<sub>4</sub> were separately loaded on C<sub>3</sub>N<sub>4</sub>. If Pt and Co<sub>3</sub>O<sub>4</sub> were both deposited at the outer surface of C<sub>3</sub>N<sub>4</sub>, a decrease in overall water splitting activity accompanied by reduced O<sub>2</sub> production otherwise resulted.

It should be noted that the dispersion and size of the loaded co-catalysts pose significant effect on the photocatalytic efficiency. Poor dispersion and uneven size distribution of co-catalysts can cause significant particle aggregation, reducing long-term stability to deteriorate the practice of photocatalysis. Devising a versatile synthetic procedure capable of depositing co-catalysts with good dispersion and uniform size is important to ensure the sustainable operation of photocatalysts. For yolk-shell and hollow nanostructures, it is difficult to control the dispersion and size of the deposited co-catalysts at inner surface, especially when the initial template or the resulting shell contains irregular porosity.<sup>139</sup> She *et al.*<sup>137</sup> and Sun *et al.*<sup>139</sup> employed zeolites with ordered micropores as a structurally confined template to allow for the loading of well-dispersed co-catalysts at the inner surface of hollow nanostructures. As displayed in Fig. 16, the nanoconfinement effect of zeolites enabled a fine tunability of dispersion, size and amount for the loaded MO<sub>x</sub> (M = Pd, Co, Ni or Cu).<sup>137</sup> Further external growth of TiO<sub>2</sub> followed by the removal of zeolites resulted in the formation of TiO<sub>2</sub> hollow nanostructures with MO<sub>x</sub> exactly deposited at the inner surface. The outer photo-deposited Pt can then be divided from the inner MnO<sub>x</sub> by TiO<sub>2</sub>, achieving great spatial separation for the dual co-catalysts to show remarkable

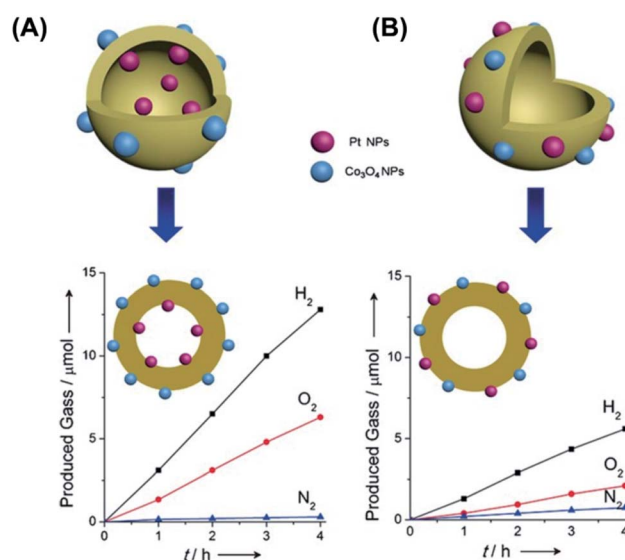


Fig. 15 Comparative results of photocatalytic water splitting under UV illumination on C<sub>3</sub>N<sub>4</sub> hollow structures with (A) spatially separate co-catalysts and (B) randomly distributed co-catalysts. Reproduced with permission from ref. 138. Copyright 2016, Wiley.



## Review

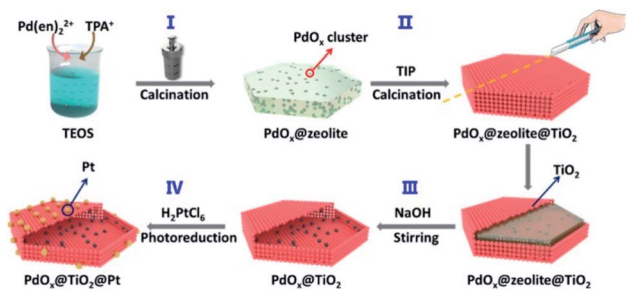


Fig. 16 Schematic illustration of using zeolites as structurally confined template to achieve spatial separation of dual co-catalysts on  $\text{TiO}_2$ . Reproduced with permission from ref. 137. Copyright 2020, American Chemical Society.

photocatalytic activity. On the other hand, atomic layer deposition (ALD) also offers feasibility for deposition of highly-dispersed co-catalysts at the selected positions. Zhang *et al.* employed ALD to sequentially deposit Pt nanoclusters and  $\text{TiO}_2$  layer on carbon nanocoil (CNC) template.<sup>135</sup> Upon a removal of CNC, porous  $\text{TiO}_2$  nanotubes with Pt deposited at the inner surface were obtained. Further ALD deposition of  $\text{CoO}_x$  produced  $\text{TiO}_2$  nanotubes with Pt and  $\text{CoO}_x$  respectively loaded at the inner and outer surfaces. The amount and location of the loaded Pt and  $\text{CoO}_x$  can be finely controlled by adjusting ALD conditions, proven to be crucial for interfacial charge dynamics and determinant to the resultant photocatalytic activity.

Despite extensive efforts, it has been pointed out that the spatially separated co-catalysts can only enhance charge separation at surface and are not efficient for promoting charge separation in bulk. To address this issue, Li *et al.* prepare for  $\text{TiO}_2/\text{In}_2\text{O}_3$  mesoporous hollow spheres and selectively deposited Pt and  $\text{MnO}_x$  at the inner and outer surface.<sup>140</sup> As illustrated in Fig. 17A and B, the idea was to combine the advantages of  $\text{TiO}_2/\text{In}_2\text{O}_3$  as thin heterojunction, and Pt and  $\text{MnO}_x$  as spatially separated co-catalysts. Here, the thin heterojunction (overall thickness = 20 nm) can promote effective charge separation in the whole bulk region as the nominal diffusion path of charge carriers is comparable to the heterojunction thickness. The separated charge carriers can then be transported in an opposite direction to the surface region, where the electrons and holes were respectively trapped by Pt and  $\text{MnO}_x$ . With this synergy, much enhanced photocatalytic activity of  $\text{O}_2$  production can be attained. Such a regulation of charge separation in the bulk can be evidenced by studying charge dynamics of the samples as they were placed in an environment containing electron and hole scavengers. Under this situation, charge dynamics of the samples can be dictated by bulk recombination. The resultant time-resolved PL data suggested that bulk recombination was indeed prohibited by the introduction of thin heterojunction. Likewise, Wang *et al.* proposed a thin heterojunction based on  $\text{TiO}_2/\text{CdS}$  hollow spheres (overall thickness = 23 nm) and loaded Pt and  $\text{Co}_3\text{O}_4$  at the inner and outer surface, respectively.<sup>136</sup> As Fig. 17C and D illustrates, the synergy of thin heterojunction and spatially separated co-

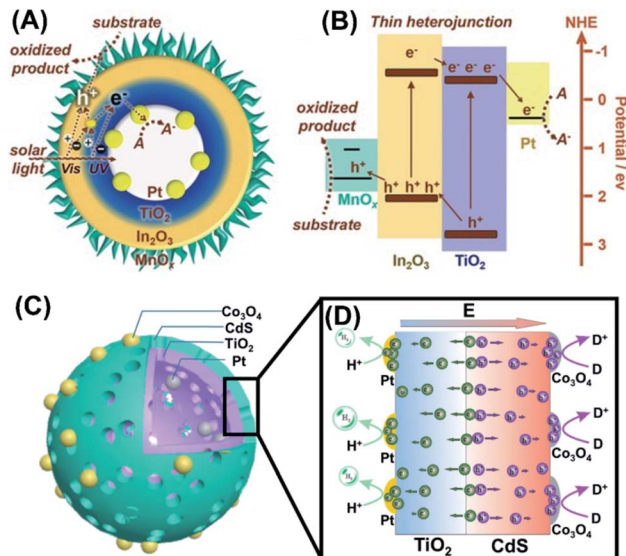


Fig. 17 (A) Microstructural model for  $\text{TiO}_2/\text{In}_2\text{O}_3$  hollow spheres with Pt and  $\text{MnO}_x$  respectively deposited at the inner and outer surface. (B) Plausible charge transfer mechanism for superior photocatalytic activity of  $\text{Pt}/\text{TiO}_2/\text{In}_2\text{O}_3/\text{MnO}_x$ . Reproduced from ref. 140. Copyright 2016, Wiley. (C) Microstructural model for  $\text{TiO}_2/\text{CdS}$  hollow spheres with Pt and  $\text{Co}_3\text{O}_4$  respectively deposited at the inner and outer surface. (D) Plausible charge transfer mechanism for superior photocatalytic activity of  $\text{Pt}/\text{TiO}_2/\text{In}_2\text{O}_3/\text{Co}_3\text{O}_4$ . Reproduced with permission from ref. 136. Copyright 2016, Wiley.

catalysts gave rise to a remarkable enhancement in the performance of solar  $\text{H}_2$  production.

## 5. Conclusions and future perspective

In this review, the synthetic schemes for yolk-shell nanostructures are systematically organized. The formation of hollow shell can be attributed to the difference in solubility product for Kirkendall effect and the difference in redox potential for galvanic replacement. By pairing up materials with appropriate solubility products and redox potentials, yolk-shell nanostructures composed of various components can be readily obtained. Ostwald ripening takes the advantage of the dissolution-re-deposition principle during the process of crystal growth. The shell layer grows as the core particles dissolve, or the shell layer (or core particle) itself dissolves and re-crystallizes, both of which generate empty space between the core and shell as a result of the contraction of the dissolved components. Both chemical etching and thermal treatment are based on the assistance of sacrificial templates like  $\text{SiO}_2$  and carbon. Upon a removal of the templates, cavity inside the shell can evolve. The necessity of removing templates by strong acids/bases or high-temperature calcination however poses significant concerns on whether or not the remaining components will be deteriorated in such a harsh condition. Definitely, yolk-shell nanostructures with tunable structural parameters are essential for further tailoring and optimizing the intrinsic





properties. Despite extensive synthetic efforts, it is difficult to achieve exact control over the size of void space and thickness of encircling shell because void formation and shell growth proceed simultaneously in most cases. Continuous efforts must be made to manipulate the growth of hollow shell and obtain yolk-shell nanostructures with controllable void size and shell thickness.

Compared to the structurally analogous counterpart, *i.e.* core-shell nanostructures, yolk-shell nanostructures have exceptional structural features that can be exploited to tailor the photocatalytic properties. For example, the movable yolk particles can stir the reaction solution inside the void space, creating a homogeneous reaction environment for accelerating mass transfer kinetics and thus enhancing the photocatalytic activity. An in-depth understanding of mass transfer kinetics across the shell can lay the foundation for constructing robust nano-reactors based on yolk-shell nanostructures. In terms of charge carrier transfer, a small number of studies are available on deriving the correlations of interfacial charge dynamics with the photocatalytic performance for yolk-shell nanostructures. The charge transfer between the yolk and the shell and its influence on the photocatalytic properties has seldom been examined. In particular, the movable yolk can cause ambiguous interpretation of charge transfer dynamics. Because the yolk particles are constantly moving during the photocatalytic process, complicated charge transfer and separation behaviors are expected. How to devise a realistic charge transfer mechanism by considering the influence of movable yolk is another focus of future direction. One future scenario of using yolk-shell nanostructures for advanced photocatalytic applications is spreading their powder into a polluted pond or river for environmental remediation. The empty space in the hollow shell enables yolk-shell nanostructures to float on water, providing high adaptability and recyclability for their participation in photocatalytic processes. On the other hand, spatial separation of photoexcited electrons and holes at the yolk and shell region is achievable for yolk-shell nanostructures. These separated charge carriers can be more easily and separately accessible by the reactive species, allowing for the concurrent yet separate occurrence of reduction and oxidation reactions. This achievement is particularly appealing to overall water splitting and CO<sub>2</sub> reduction, the scenario of which can better operate as the redox reactions occur at different sites. It is thus not surprising to witness the realization of artificial photosynthesis by utilizing yolk-shell nanostructures. These demonstrations are not yet fully comprehended, requiring further efforts to make significant breakthrough.

## Conflicts of interest

There are no conflicts to declare.

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