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Mixed Tb/Dy coordination ladders based on tetra(carboxymethyl)thiacalix[4]arene: a new avenue towards luminescent molecular nanomagnets†

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The macrocyclic ligand calix[4]arene (L1) and its sulphur-containing analogue thia[4]calixarene (L2) are promising precursors for functional molecular materials as they offer rational functionalization with various organic groups. Here, we present the first example of lanthanide-based coordination polymers built from the macrocyclic thiacalix[4]arene backbone bearing four carboxylic moieties, namely, ligand H₄L3. The combination of H₄L3 with the Tb³⁺ and Dy³⁺ cations led to the formation of 1D ladder-type coordination polymers with the formula [Ln^{III}HL3DMF₃](DMF) (where DMF = dimethylformamide and Ln = Tb or Dy, denoted as HL3–Tb and HL3–Dy), which resulted from the coordination of the lanthanide cations with the partially deprotonated ligand HL3^{3–} that behaved as a T-shape connector. The coordination sphere around the metal was completed by the coordinated DMF solvent molecules. By combining both Tb³⁺ and Dy³⁺ cations, isostructural heterobimetallic solid solutions HL3–Tb_{1–x}Dy_x were also prepared. HL3–Tb and HL3–Dy showed visible light photoluminescence originating from the f–f electronic transitions of pale green emissive Tb³⁺ and pale yellow emissive Dy³⁺ with efficient sensitization by the functionalized thia[4]calixarene ligand HL3. In the HL3–Tb_{1–x}Dy_x solid solutions, the Tb/Dy ratio governed both the emission colour as well as the emission quantum yield, which reached even 28% at room temperature for HL3–Tb. Moreover, HL3–Dy exhibited a slow magnetic relaxation effect related to the magnetic anisotropy of the dodecahedral Dy³⁺ complexes, which were well isolated in the crystal lattice by expanded organic spacers.

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Introduction

Coordination polymers (CPs),^{1,2} MOFs³ and molecular networks^{4,5} have molecular architectures that offer properties such as gas storage,^{6,7} catalysis,^{8,9} magnetism^{10,11} and luminescence.^{12–14} The

generation of such molecular species is based on the formation of coordination bonds between organic coordinating ligands and metal centres or complexes.¹⁵ A large number of coordination polymers that display different connectivity patterns have been generated using different types of polytopic ligands and metals.

Concerning the luminescence properties⁷ of CPs, special attention has been devoted to lanthanide-based CPs,^{16–18} owing to their outstanding photoluminescence originating from the triplet excited states of f–f transitions. Lanthanide-based CPs may find applications in luminescence sensing, light emission together with photonics,^{19–21} and as ligands presenting a strong antenna effect,²² allowing energy transfer from an organic ligand to the metal. The use of macrocyclic ligands for the formation of such assemblies has been documented;²³ however, most of the reports are mainly related to porphyrin-based lanthanide CPs.²⁴

Concerning the magnetic properties of lanthanide-based CPs, examples of molecular nanomagnets have been reported earlier.^{25–27} In addition, LnCPs may combine both SMM behaviour and tunable photoluminescence, as shown recently.^{28–30}

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In this study, we report the structure of a series of isostructural new 1D lanthanide-based coordination networks,

The ligand fluorescence and Dy³⁺ and Tb³⁺ phosphorescence were thoroughly studied in the solid-state in the corresponding (Dy/Tb) solid solutions. In addition, the magnetic properties of the anisotropic Dy³⁺ compound were investigated in detail.

Synthesis and structural details

The crystal also contained a non-coordinated DMF molecule occupying the interstices between antiparallel 1D coordination networks without forming any specific interactions with them. The terbium atoms acted as T-shaped connectors, connecting three carboxylate groups belonging to three adjacent **HL3**³⁻ ligands, leading to a double chain ladder-like structure (Fig. 2a and b and schematically represented in d).

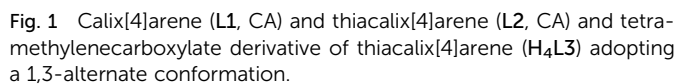


Table 1 Crystallographic parameters for HL3–Dy and HL3–Tb recorded at 173 K

Formula	C ₄₈ H ₅₃ O ₁₂ S ₄ Dy(C ₃ H ₇ NO) ₃ , C ₃ H ₇ NO, C ₆₀ H ₈₁ DyN ₄ O ₁₆ S ₄ , HL3–Dy	C ₄₈ H ₅₃ O ₁₂ S ₄ Tb(C ₃ H ₇ NO) ₃ , C ₃ H ₇ NO, C ₆₀ H ₈₁ TbN ₄ O ₁₆ S ₄ , HL3–Tb
Molecular weight	1405.02	1401.44
Crystal system	Triclinic	Triclinic
Space group	<i>P</i> $\bar{1}$	<i>P</i> $\bar{1}$
<i>a</i> (Å)	15.0091(17)	15.0160(17)
<i>b</i> (Å)	15.9798(18)	15.9630(14)
<i>c</i> (Å)	16.7401(18)	16.7660(16)
α (deg)	62.217(5)	62.226(3)
β (deg)	67.626(5)	67.513(3)
γ (deg)	79.181(6)	79.132(4)
<i>V</i> (Å ³)	3284.5(7)	3285.2(6)
<i>Z</i>	2	2
Colour	Colourless	Colourless
Crystal dim (mm ³)	0.090 × 0.100 × 0.100	0.090 × 0.100 × 0.100
<i>D</i> _{calc} (g cm ^{−3})	1.421	1.417
<i>F</i> (000)	1454	1452
μ (mm ^{−1})	1.332	1.270
Wavelength (Å)	0.71073	0.71073
Number of data meas	83 469	36 284
Number of data with <i>I</i> > 2σ(<i>I</i>)	17 032 [<i>R</i> (int) = 0.0752]	17 822 [<i>R</i> (int) = 0.0558]
<i>R</i>	<i>R</i> ₁ = 0.0557, <i>wR</i> ₂ = 0.1391	<i>R</i> ₁ = 0.0494, <i>wR</i> ₂ = 0.1073
<i>R</i> _w	<i>R</i> ₁ = 0.0756, <i>wR</i> ₂ = 0.1559	<i>R</i> ₁ = 0.0708, <i>wR</i> ₂ = 0.1186
GOF	1.057	1.011
Largest peak in final difference (eÅ ^{−3})	1.740 and −1.406	1.411 and −1.140

Consequently, the fourth carboxylic group remained protonated without interactions with the other components of the crystal.

As expected, different C–O distances were observed for the carboxylic moieties of **HL3**^{3−}, revealing the presence of one non-coordinated carboxylic group and three coordinated carboxylate moieties (see Table 2).

Within the 1-D arrays, the shortest Tb–Tb distance was equal to 11.879(3) Å, and the packing of the 1D networks in the *yOz* plane (Fig. 2b, packing along the *a*-axis) led to a shorter Tb–Tb distance of 7.836(4) Å between two Tb atoms belonging to two different chains.

Since both compounds **HL3–Dy** and **HL3–Tb** were isostructural (which is not the case for Gd and Eu analogues, for example), solid solutions of formula **HL3–Tb**_{1−*x*}**Dy**_{*x*} were

prepared using the same synthetic procedure as the one used for **HL3–Dy** and **HL3–Tb** (see Experimental section). Meanwhile, in our case, the related Eu compound obtained with **H₄L3** was not isostructural with **HL3–Tb** or **HL3–Dy**.

Starting with different experimental Dy/Tb ratios, microcrystalline compounds of the following formula were obtained: **HL3–Tb**_{0.07}**Dy**_{0.93}, **HL3–Tb**_{0.20}**Dy**_{0.80}, **HL3–Tb**_{0.42}**Dy**_{0.58}, **HL3–Tb**_{0.67}**Dy**_{0.33} and **HL3–Tb**_{0.95}**Dy**_{0.05}.

The structure of the solid solution was investigated by XRPD, as shown in Fig. 3 (for the determined cell parameters of the corresponding single crystals, see Table S1, ESI†). XRPD clearly demonstrated the structural homogeneity of the phase of the formed coordination compounds. All the **HL3–Tb**_{1−*x*}**Dy**_{*x*} solid

Table 2 Selected bond distances (Å) for HL3–Dy and HL3–Tb

	HL3–Dy	HL3–Tb
C–O (carboxyl-ate/-ic)	1.232(5) and 1.273(5) 1.239(5) and 1.268(4) 1.254(5) and 1.261(5) 1.189(5) and 1.323(5) (carboxylic)	1.215(4) and 1.269(4) 1.253(4) and 1.258(4) 1.243(4) and 1.259(4) 1.194(5) and 1.320(5) (carboxylic)
Ln–O (DMF)	2.354(3) 2.366(3) 2.380(3)	2.355(3) 2.369(2) 2.395(3)
Ln–O (carboxylate)	2.256(3) 2.380(3) 2.393(3) 2.404(3) 2.438(3)	2.277(2) 2.393(3) 2.411(3) 2.418(3) 2.440(2)



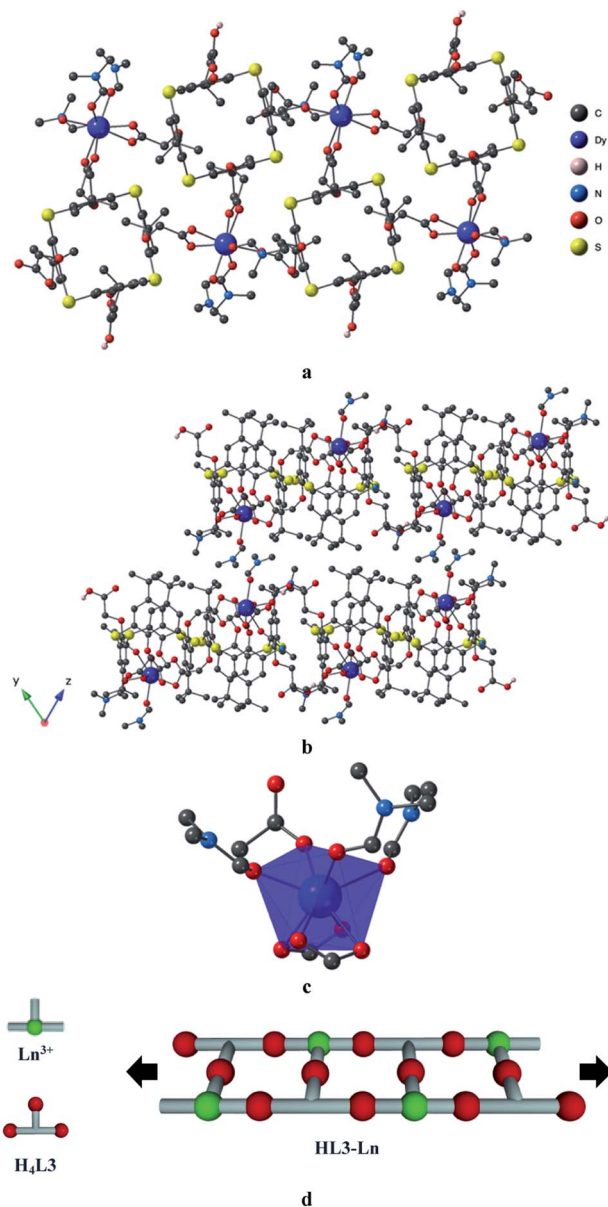


Fig. 2 (a) A portion of the crystal structure of HL3-Tb, (b) crystal packing of HL3-Tb, (c) the coordination environment of Ln^{3+} and (d) the schematic representation of connectivity in HL3-Tb.

solutions were isostructural and displayed the same structure as those of HL3-Tb or HL3-Dy. Since there was only one type of crystallographically independent Ln ion in the unit cell, owing to similar size of both lanthanide cations, one may assume a statistical distribution of the lanthanide cations in the crystal.

The composition of the HL3-Tb_{1-x}Dy_x solid solutions was also checked using energy-dispersive X-ray spectroscopy (EDS), as shown in ESI.†

The experimentally determined Tb/Dy ratios measured were in good agreement with the relative quantities of the lanthanide salts used for synthesis. The slight differences were probably due to weighing inaccuracy and the overestimation of Tb³⁺ content in the EDS measurements (Table 3).

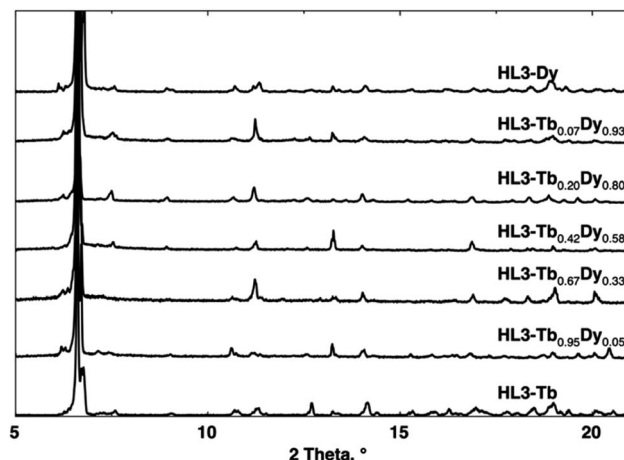


Fig. 3 Comparison of the simulated PXRD patterns for HL3-Tb (bottom) and HL3-Dy (top) and the recorded PXRD patterns for HL3-Tb_{1-x}Dy_x solid solutions. Discrepancies in intensity between the observed and simulated patterns are due to the preferential orientations of the microcrystalline powders.

The luminescence and magnetic properties of this HL3-Dy series of compounds are presented below.

Photoluminescence properties

For all the investigated compounds, the measurements were carried out in the solid-state using polycrystalline powders.

Among the bands of the excitation spectra of HL3-Tb and HL3-Dy at RT in the 350–500 nm region with emission at $\lambda_{\text{em}} = 534$ nm and 576 nm, respectively (see Fig. S1, ESI†), those attributed to the typical f–f transitions of the Tb³⁺ and Dy³⁺ cations presented lower intensities compared with the band attributed to the $\pi \rightarrow \pi^*$ transitions of H₄L3 (300–350 nm). These observations clearly indicated the occurrence of an energy transfer from the excited level of ligand HL3³⁻ to the emitting state of the Tb³⁺ (or Dy³⁺) cations, as seen in the case of antenna effect.⁷⁷ The capability of calixarene carboxylate derivatives to act as antennae for Tb³⁺ complexes has already been documented.⁵³

The room-temperature emission spectra of the HL3-Tb and HL3-Dy coordination networks, when excited at $\lambda_{\text{ex}} = 350$ nm, are depicted in Fig. 4. For HL3-Tb, the emission spectrum is composed of the characteristic five electronic transitions, i.e., $^5\text{D}_4 \rightarrow ^7\text{F}_6$, $^5\text{D}_4 \rightarrow ^7\text{F}_5$ (strongest), $^5\text{D}_4 \rightarrow ^7\text{F}_4$, $^5\text{D}_4 \rightarrow ^7\text{F}_3$ and $^5\text{D}_4$

Table 3 Comparison of the observed and experimental Tb/Dy ratios in HL3-Tb_{1-x}Dy_x coordination polymers

	Tb/Dy experimental ratio	Tb/Dy observed ratio (EDS)
HL3-Tb _{0.95} Dy _{0.05}	95 : 05	95 : 05
HL3-Tb _{0.67} Dy _{0.33}	60 : 40	67 : 33
HL3-Tb _{0.42} Dy _{0.58}	40 : 60	42 : 58
HL3-Tb _{0.20} Dy _{0.80}	20 : 80	20 : 80
HL3-Tb _{0.07} Dy _{0.93}	10 : 90	07 : 93

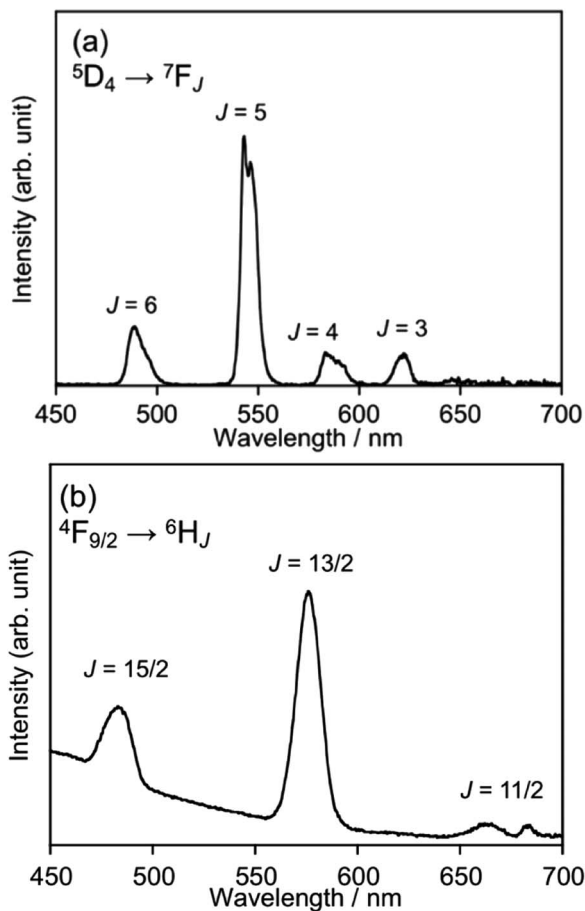


Fig. 4 Emission spectra ($\lambda_{\text{ex}} = 350$ nm) of HL3-Tb (a) and HL3-Dy (b) at RT.

$\rightarrow {}^7\text{F}_2$ (weakest), whereas for **HL3-Dy**, the ${}^4\text{F}_{9/2} \rightarrow {}^6\text{H}_{15/2}$, ${}^4\text{F}_{9/2} \rightarrow {}^6\text{H}_{13/2}$ (strongest) and ${}^4\text{F}_{9/2} \rightarrow {}^6\text{H}_{11/2}$ (weakest) electronic transitions are observed, as shown in Fig. 4a and b, respectively. These were typical f-f transitions with emission in the green region for Tb^{3+} and yellow-white region for Dy^{3+} . The corresponding emission spectra at 77 K are presented in Fig. S2a and b, ESI†

The photophysical behaviour of the mixed **HL3-Tb_{1-x}Dy_x** solid solutions was analysed at RT and 77 K. The luminescence spectra of the **HL3-Tb_{1-x}Dy_x** solid solutions observed at RT and 77 K are shown in Fig. 5 and S3.†

At RT, the **HL3-Tb_{1-x}Dy_x** solid solutions exhibited the typical luminescence bands of the f-f transitions of Tb^{3+} at 545 nm and those of Dy^{3+} at 574 nm. The band intensity at 545 nm decreased when x was decreased, and that at 574 nm increased when x was decreased, as shown in Fig. 5a and b. The introduction of Dy^{3+} in the **HL3-Tb_{1-x}Dy_x** solid solutions allowed slight tuning of the emission colour ranging from pale green to pale yellow, as shown in Fig. 5a and b.

The absolute quantum yields of Tb^{3+} in the series of **HL3-Tb_{1-x}Dy_x** solid solutions were observed (see Fig. S3a and b, ESI†), and the values at 77 K were found to be 58.9% (27.9% at RT) for **HL3-Tb**, 32.6% for **HL3-Tb_{0.95}Dy_{0.05}**, and 15.1% for

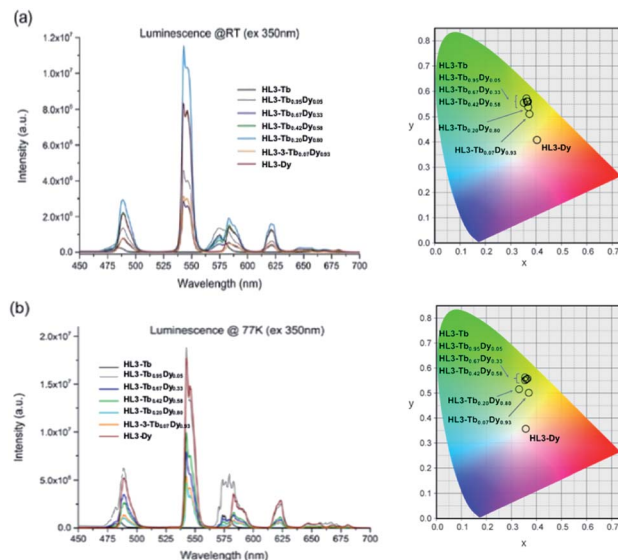


Fig. 5 Emission spectra ($\lambda_{\text{ex}} = 350$ nm) of HL3-Dy (black) and HL3-Tb (brown), and the solid solutions **HL3-Tb_{0.07}Dy_{0.93}** (grey), **HL3-Tb_{0.20}Dy_{0.80}** (blue), **HL3-Tb_{0.42}Dy_{0.58}** (green), **HL3-Tb_{0.67}Dy_{0.33}** (light blue) and **HL3-Tb_{0.95}Dy_{0.05}** (orange) at (a) room temperature and (b) 77 K and the corresponding emission colours presented on the CIE 1931 chromaticity diagrams.

HL3-Tb_{0.20}Dy_{0.80}. Therefore, the existence of Dy^{3+} ion in the solid solution played a role in decreasing the luminescence quantum yields of the **HL3-Tb_{1-x}Dy_x** solid solutions.

The luminescence decay profile and fitted results are summarized in Tables S2 and S3 and Fig. S3–S5 in ESI†. As described above, the luminescence quantum yields of Dy^{3+} in the neat and solid solutions with Tb^{3+} were quite low, and it was not possible to obtain the decay profiles at both RT and 77 K for **HL3-Dy**. The lifetimes of **HL3-Tb** at RT and 77 K were estimated as a single component, and the τ values were 0.866 and 0.916 ms, respectively. After the addition of Dy^{3+} in the **HL3-Tb_{1-x}Dy_x** solid solutions, the τ values decreased, and the number of luminescence components of Tb^{3+} increased to two or three (see Tables S2 and S3 in ESI†). The luminescence decay curve of **HL3-Tb_{0.2}Dy_{0.8}** could be divided into two components, namely 0.378 and 0.932 ms ($\lambda = 543$ nm), at RT. The latter value corresponded to the one observed for **HL3-Tb**. The second and third emission lifetimes were related to the increasing number of Dy^{3+} centers surrounding the Tb^{3+} centers.

The τ_{obs} values (Tables S2 and S3 in ESI†) were in accordance with those reported for other calixarene-based coordination compounds in the solid-state.^{53,78}

Magnetic measurements

As **HL3-Dy** contains the magnetically anisotropic Dy^{3+} cations, which may lead to a potential single-molecule magnet character, its magnetic properties were investigated. The magnetic behaviour of **HL3-Tb** is not presented here because of the lack of significant anisotropy for Tb. As shown in Fig. 6, at 300 K, the $\chi_{\text{M}}T$ product reached $14.1 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$, which is the exact value



expected for free-ion approximation. A typical decrease in the $\chi_M T$ product was observed upon cooling, which could be related to thermal depopulation within the m_J levels of the ground $^6H_{15/2}$ multiplet of Dy^{3+} . Besides a slow decrease in the signal, the $\chi_M T$ versus T plot did not display abrupt changes, which suggested the lack of magnetic interactions down to 1.8 K as a result of the efficient magnetic isolation of lanthanide centres in the crystal lattice. Magnetization measured at $T = 1.8$ K increased quickly with increasing magnetic field to 10 kOe and much slower at higher fields, reaching $5.4 \mu_B$ at 60 kOe without saturation, as expected for anisotropic Dy^{3+} (inset in Fig. 6).⁷⁹

To investigate the slow magnetic relaxation effects in **HL3-Dy**, the alternate-current (ac) magnetic characteristics were gathered at $H_{dc} = 0$ (Fig. 7a). The maxima of the frequency dependence of out-of-phase magnetic susceptibility, χ_M'' , appeared on the edge of the measurement range, even at 1.8 K, quickly moving towards higher frequencies while the temperature increased. This indicated the presence of slow magnetic relaxation, a characteristic effect of the single-molecule magnet (SMM) behaviour. To slow down the relaxation by reducing the presumably disturbing quantum tunnelling of the magnetization (QTM) process, an external dc magnetic field was applied (see Fig. S7, ESI†). The field-variable ac magnetic characteristics were analysed using the generalized Debye model for a single relaxation.⁸⁰ The extracted relaxation times were plotted against the magnetic field, and the resulting dependence was analysed by taking into account the QTM effect (Quantum Tunnelling of the Magnetisation), which was dominant at low fields, and a field-induced direct process responsible for the acceleration of the relaxation process at higher fields (see ESI, Fig. S7†).⁸¹ An equilibrium between these processes appeared at $H = 1000$ Oe, which was selected as the optimal dc field to investigate magnetic relaxation in **HL3-Dy**. The dc magnetic characteristics were gathered in the 1.8–5 K temperature range (Fig. 7b and see Fig. S8, ESI†). They were analysed using the generalized Debye model for a single relaxation, and the resulting relaxation times are presented as the function of temperature (Fig. 7c).

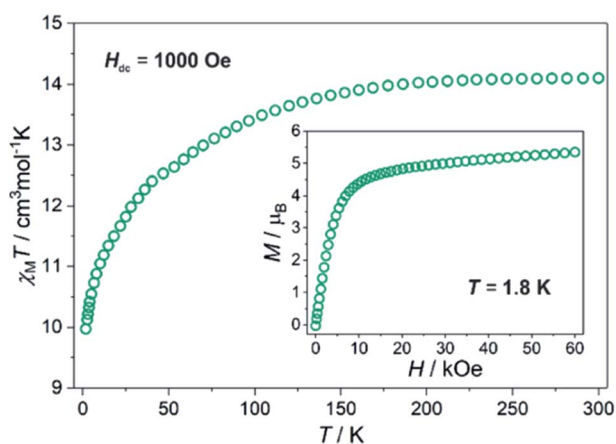


Fig. 6 The temperature dependence of the $\chi_M T$ product of **HL3-Dy** at $H_{dc} = 1000$ Oe, and the field dependence of molar magnetization at $T = 1.8$ K (inset).

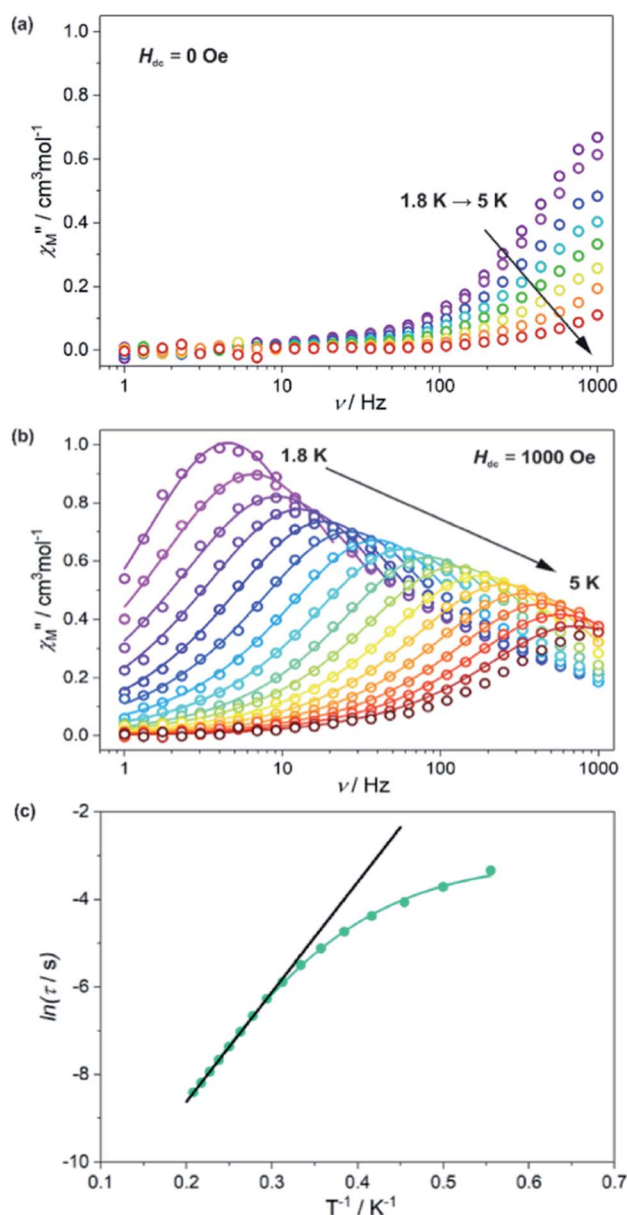


Fig. 7 For **HL3-Dy**: (a) Zero dc field $\chi_M''(\nu)$ curves in the range of 1.8–5 K, (b) dc field $\chi_M''(\nu)$ curves in the range 1.8–5 K under a dc field of 1 kOe, and (c) temperature dependence of the relaxation time, τ , at 1 kOe, fitted using the Arrhenius law in the limited 3.4–5 K range (extrapolated black line) and within the full 1.8–5 K range using the contributions from the Orbach, Raman, direct and QTM relaxation processes (green line; see text for details).

In the higher temperature range of 3.4–4.8 K, the temperature dependence of the relaxation times was close to linear, and thus, it could be analysed by a simple Arrhenius law, which is presented in eqn (1):

$$\ln(\tau) = \frac{U_{\text{eff}}}{k_B T} + \ln(\tau_0) \quad (1)$$

where U_{eff} represents the effective thermal energy barrier, while τ_0 is the relaxation attempt time for spin reversal at infinite temperature. The obtained parameters $U_{\text{eff}}/k_B = 25.08(9)$ K, $\tau_0 =$

$1.2(2) \times 10^{-6}$ s were within the characteristic limits of single-molecule magnets with moderate magnetic anisotropy.⁷⁹ More precisely, the full temperature dependence curve of the relaxation times was fitted using contributions from the Orbach and Raman relaxation processes together with direct and QTM processes taken from the previous analysis of field-dependent relaxation. Therefore, eqn (2) was used:

$$\tau^{-1} = \tau_1^{-1} \exp(-U_{\text{eff}}/k_{\text{B}}T) + B_{\text{Raman}}T^n + \frac{a(1 + c^2H^2)}{(1 + bH^2)} + ATH^4 \quad (2)$$

where the first term represents the Orbach process related to the Arrhenius law, the second introduces a two-phonon Raman relaxation pathway, the third denotes the QTM effect and the last one shows a direct process.^{82,83} To avoid over-parameterization, the two last terms describing the QTM effect and direct process taken from the field-dependence analysis and the related parameters were not optimized during the fitting procedure (Fig. S8†). The obtained values $U_{\text{eff}}/k_{\text{B}} = 31(3)$ K and $\tau_0 = 1.2(7) \times 10^{-6}$ s were in good agreement with those extracted from the Arrhenius law. The Raman process had to be included with $B_{\text{Raman}} = 0.27(1) \text{ s}^{-1} \text{ K}^{-n}$ and $n = 6.0(3)$. Thus, the last value expected for lanthanide-based SMMs would be located in the 2–9 range.³⁹ As several examples of SMMs lacking the Orbach process have been reported,⁸⁴ we tried to fit the temperature dependence of the relaxation times excluding Orbach process; however, it was not possible to obtain a satisfying fit.

Based on these results, we could clearly postulate that the Dy^{3+} complexes in **HL3-Dy** exhibit single-molecule magnet characteristics with a moderated thermal energy barrier originating from the Orbach relaxation process and a strong QTM effect, partially reduced by using the external dc magnetic field.

Conclusions

In this work, we have demonstrated the possibility of growing a series of lanthanide-based isostructural homometallic coordination polymers using thiacalix[4]arene bearing carboxylic moieties. Owing to their isostructural nature, a series of unprecedented heterometallic solid solutions based on the macrocyclic thiacalix[4]arene derivatives were generated. The single crystals of the two isostructural Tb^{3+} - and Dy^{3+} -based coordination polymers (**HL3-Tb** and **HL3-Dy**, formula $[\text{Ln}^{\text{III}}\text{HL3DMF}_3] \cdot (\text{DMF})$) were structurally characterized using X-ray diffraction, and their photophysical properties were investigated. The formation of the 1D ladder-type coordination compounds resulted in the specific binding of the Ln^{3+} cation to the partially deprotonated organic ligands that behaved as a “T-Shape” connector. In the solid solutions that combined Tb^{3+} and Dy^{3+} forming **HL3-Tb_{1-x}Dy_x**, the Ln^{3+} cations are randomly distributed within the crystals.

The photoluminescence properties of the coordination polymers were investigated. Ligand **HL3**³⁻ sensitized the luminescence of the Tb^{3+} or Dy^{3+} cations in **HL3-Tb** and **HL3-Dy** at RT, and they exhibited the characteristic f–f transition due to the presence of the Ln^{3+} cations. The introduction of Dy^{3+} in

HL3-Tb allowed the tuning of the emission colour of the solid solutions, which ranged between pale yellow and pale green.

In addition, the magnetic measurements of **HL3-Dy** revealed significant magnetic anisotropy in the Dy^{3+} complexes attached to the thiacalix[4]arene moieties, leading to SMM behaviour. Therefore, this series of new coordination polymers can be considered as a novel source of bifunctional magneto-luminescent molecule-based materials.⁸⁵ We present a novel and elegant pathway for the construction of luminescent molecular nanomagnets by the incorporation of lanthanide ions in coordination chains based on functionalized thiacalix[4]arene ligands.

Experimental methods

X-ray crystallography

Data were collected at 173(2) K on a Bruker Apex-II-CCD diffractometer equipped with an Oxford Cryosystem liquid N_2 device using graphite-monochromated $\text{Mo-K}\alpha$ ($\lambda = 0.71073$ Å) radiation. For all the structures, diffraction data were corrected for absorption. The structures were solved using SHELXS-97 and refined by full matrix least-squares on F^2 using SHELXL-97. Hydrogen atoms were introduced at calculated positions and refined using a riding model.⁸⁶ The structures can be obtained free-of-charge from the Cambridge Crystallographic Data Centre via www.ccdc.cam.ac.uk/datarequest/cif. CCDC: 1898825 (**HL3-Tb**) and 1898826 (**HL3-Dy**).

Powder diffraction studies (PXRD)

The diffraction diagrams were collected on a Bruker D8 diffractometer using monochromatic $\text{Cu-K}\alpha$ radiation in a scanning range of 4 and 40° using a scan step size of 8° min^{-1} .

Elemental analysis

Elemental analysis was performed by the Analysis Service of the Faculty of Chemistry of the University of Strasbourg.

EDS experiments

Elemental analyses using EDS were performed on a Bruker-QUANTAX System, where the EDS Detector was equipped with a GEMINI FE-SEM ULTRA55 (Carl Zeiss). The samples were mixed with colloidal graphite in isopropanol, dried and pressed into pellets. The EDS analysis was conducted at an accelerating voltage of 15 kV.

Photophysical properties

Luminescence and excitation spectra were recorded on a Fluorolog 3-22 (Horiba Jobin Yvon).

The absolute luminescence quantum yields and luminescence lifetimes were determined using an absolute luminescence quantum yield spectrometer C9920-02 (Hamamatsu Photonics K. K.) and a Quantaurus-Tau C11367-12 (Hamamatsu Photonics K. K.), respectively, with pulsed excitation light sources.



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