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In situ analysis of gas dependent redistribution kinetics in bimetallic Au-Pd nanoparticles†

Marta Perxés Perich,^a Christopher R. O'Connor,^b Koen M. Draijer,^a Nienke L. Visser,^a Nongnuch Artrith,^a Christian Reece,^b Petra E. de Jongh^a and Jessi E. S. van der Hoeven^{a*}

The catalytic and plasmonic properties of bimetallic gold–palladium (Au-Pd) nanoparticles (NPs) critically depend on the distribution of the Au and Pd atoms inside the nanoparticle bulk and at the surface. Under operating conditions, the atomic distribution is highly dynamic. Analyzing gas induced redistribution kinetics at operating temperatures is therefore key in designing and understanding the behavior of Au-Pd nanoparticles for applications in thermal and light-driven catalysis, but requires advanced *in situ* characterization strategies. In this work, we achieve the *in situ* analysis of the gas dependent alloying kinetics in bimetallic Au-Pd nanoparticles at elevated temperatures through a combination of CO-DRIFTS and gas-phase *in situ* transmission electron microscopy (TEM), providing direct insight in both the surface- and nanoparticle bulk redistribution dynamics. Specifically, we employ a well-defined model system consisting of colloidal Au-core Pd-shell NPs, monodisperse in size and uniform in composition, and quantify the alloying dynamics of these NPs in H₂ and O₂ under isothermal conditions. By extracting the alloying kinetics from *in situ* TEM measurements, we show that the alloying behavior in Au-Pd NPs can be described by a numerical diffusion model based on Fick's second law. Overall, our results indicate that exposure to reactive gasses strongly affects the surface composition and surface alloying kinetics, but has a smaller effect on the alloying dynamics of the nanoparticle bulk. Both our *in situ* methodology as well as the quantitative insights on restructuring phenomena can be extended to a wider range of bimetallic nanoparticle systems and are relevant in understanding the behavior of nanoparticle catalysts under operating conditions.

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Introduction

Bimetallic Au-Pd nanoparticles have extensively been explored in the fields of catalysis and plasmonics due to their favorable combination of optical and catalytic properties.^{1–3} Typically, Au offers high chemoselectivity in catalysis and a strong plasmonic response, whereas Pd provides high catalytic activity.^{3–6} Au-Pd NPs therefore present a valuable class of materials for thermal, photothermal and plasmon enhanced catalysis applications.^{1,2,4,7} In thermal catalysis, Au-Pd nanoparticles are valuable materials for selective hydrogenation and oxidation catalysis, where Pd facilitates fast H₂ and O₂ dissociation while Au steers the selectivity towards the desired partially hydrogenated or oxidated products.^{3,7–10} In both thermal and light enhanced catalytic applications, the metal distribution of the NP surface and NP

bulk are critical. The surface composition and nature of the active sites directly affect the catalytic activity and selectivity of the Au-Pd NPs,^{11–13} whereas the metal distribution of the NP bulk also impacts the plasmonic response of the metal nanoparticles.^{14–16} For instance, only by combining Au and Pd in a core-shell structure, enhanced catalytic performances compared to Pd could be achieved, while retaining most of the favorable plasmonic properties of Au.¹⁷ Hence, achieving a fine control over the metal distribution is critical in tuning the plasmonic and catalytic properties of Au-Pd NP systems.

The metal distribution within bimetallic nanoparticles is dynamic and can change under reaction conditions.^{18–21} Meta-stable structures such as Au-Pd core-shell NPs can exhibit strongly enhanced catalytic activity and selectivity compared to their alloyed counterparts,¹⁷ but are prone to restructure under reactive gasses and at elevated temperatures. Therefore, it is relevant to gain a detailed understanding of gas-dependent restructuring. Thermodynamically, Au should dominate the surface of Au-Pd NPs, because Au has a lower surface energy than Pd under vacuum.²² Indeed, Au surface segregation has been observed when Au-Pd NPs were heated under vacuum or H₂.^{23–26} However, under oxidizing conditions Pd can segregate to

^aMaterials Chemistry and Catalysis, Debye Institute for Nanomaterials Science, Utrecht University, 3584 CG Utrecht, The Netherlands. E-mail: j.e.s.vanderhoeven@uu.nl

^bRowland Institute at Harvard, Harvard University, Cambridge, Massachusetts 02142, USA

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the surface through formation of an oxidized palladium layer, lowering the surface energy of the Pd atoms.^{27,28} Theoretical studies based on density functional theory (DFT) have indeed confirmed Pd stabilization on the surface in the presence of O₂.^{27,29} However, most experimental and theoretical studies have focused on the final surface state after thermal treatment, whereas the dynamics of metal redistribution in reactive gasses remain unexplored for Au-Pd nanoparticle systems.

Direct observation of metal redistribution in bimetallic nanoparticles requires *in situ* techniques that can operate at high temperature and under gas atmospheres. Useful techniques include CO-DRIFTS, *in situ* EXAFS, and *in situ* TEM, all providing complementary information. CO adsorption measured by DRIFTS (diffuse reflectance infrared Fourier transform spectroscopy) is a surface specific, ensemble averaged technique that provides information both on the surface composition and nature of CO-adsorption sites averaged over a large number of Au-Pd NPs. Previous CO-DRIFTS studies on dilute Pd-in-Au nanoparticle alloys confirmed Au surface segregation upon annealing under vacuum²⁵ and Pd segregation under O₂ atmosphere,³⁰ matching with theory.²⁷ Additionally, CO-DRIFTS has been used to identify the Pd ensemble size at the nanoparticle surface, where the presence of Pd monomers, dimers or trimers in the NP surface can result in different catalytic properties.³¹ On the other hand, X-ray spectroscopy techniques such as EXAFS (extended X-ray absorption fine structure) have been used to analyze the metal distribution in dilute Pd-in-Au alloys after thermal treatment. The results revealed increased Pd-Pd bonds following O₂ pretreatment, indicating Pd-segregation.³¹ Ensemble techniques such as CO-DRIFTS and EXAFS provide averaged information about the sample, but have difficulties in assessing differences on a single nanoparticle level.

In situ gas-phase transmission electron microscopy (TEM) allows direct visualization of changes in individual metal

nanoparticles at elevated temperatures and/or upon exposure to reactive gasses.^{32–34} So far, several TEM studies on alloying dynamics of core-shell nanoparticles under high vacuum have been reported, using EDX (energy-dispersive X-ray) mapping,³⁵ fast tomography,^{36,37} and atomic resolution imaging.^{38,39} However, these methods are difficult to utilize in the presence of a gas atmosphere because (1) EDX is time consuming and requires a high electron dose rate, which might lead to undesired beam effects, (2) the fast tomography approach is not possible with the current *in situ* TEM holder technologies, which do not allow tilting at the needed range, and (3) high resolution studies are often limited to a few particles only, neglecting effects caused by polydispersity within a sample. Assessing the effect of a gas atmosphere on the alloying dynamics of NPs has therefore remained unexplored and requires new methodologies to monitor and extract the alloying kinetics from *in situ* gas-phase TEM data.

In this work, we combine time-lapse CO-DRIFTS and *in situ* gas-phase TEM to directly observe the effect of a reducing and oxidizing gas atmosphere on the surface and bulk alloying kinetics of Au-Pd nanoparticles. Using these techniques, the alloying process of a well-defined model system consisting of colloiddally prepared Au-core Pd-shell nanoparticles was assessed at the NP surface (CO-DRIFTS) and in the NP bulk (*in situ* TEM). Specifically, we focus on core-shell NPs with a 20 nm Au-core and a 1.7 nm-Pd shell, which are directly relevant for selective hydrogenation^{17,40} and oxidation^{40,41} catalysis. We developed a methodology to quantify the alloying dynamics of core-shell NPs from high angular annular dark field scanning transmission electron microscopy (HAADF-STEM) images acquired at atmospheric pressure without the need of EDX, fast tomography or atomic resolution imaging. Using this approach, we assess for the first time the NP bulk alloying dynamics from the *in situ* TEM data and show that the alloying dynamics can be predicted using a simple diffusion model based on Fick's second law.

Results

Preparation of the Au-core Pd-shell nanoparticles

Au-Pd core-shell nanoparticles (NPs) monodisperse in size and composition were successfully synthesized *via* colloidal synthesis. The high angular annular dark field scanning transmission electron microscopy (HAADF-STEM) image (Fig. 1a) and size distribution (Fig. 1b) show that the Au-core Pd-shell particles have an overall diameter of 23.8 ± 1.3 nm, consisting of a 19.7 ± 1.7 nm Au-core and a 1.7 ± 0.4 nm Pd-shell. The 1.7 nm Pd-shell is clearly visible in the high resolution HAADF-STEM image (Fig. 1c), which also reveals the epitaxial growth of Pd onto the Au-surface and the overall penta-twinned symmetry, as expected from the synthesis protocol.^{5,13,42} The presence of the Pd-shell and the uniformity of the metal composition among the different particles was further confirmed by EDX mapping (Fig. 1d and e). The corresponding line scan in Fig. 1f clearly demonstrates a core-shell structure, where the Pd signal extends further than the Au signal. EDX quantification also revealed that the average atomic percentage of palladium was $36 \pm 5\%$ among nearly all particles, which is close to the 31.5



Jessi E. S. van der Hoeven

Dr. Jessi E. S. van der Hoeven is an assistant professor in the Debye Institute for Nanomaterials Science at Utrecht University. Her team focusses on advanced electron microscopy characterization of catalytic materials and on colloidal catalyst design, with a particular interest in bi- and tri-metallic core-shell systems. Using these core-shell catalysts, Jessi aims to contribute an impactful research line towards cleaner chemical conversions.

Previously, Jessi received her PhD in Chemistry at Utrecht University on the topic of gold nanorod synthesis for sensing and catalysis. She then moved to Harvard University to work as a postdoctoral fellow on inverse opal catalysts for selective hydrogenation catalysis.



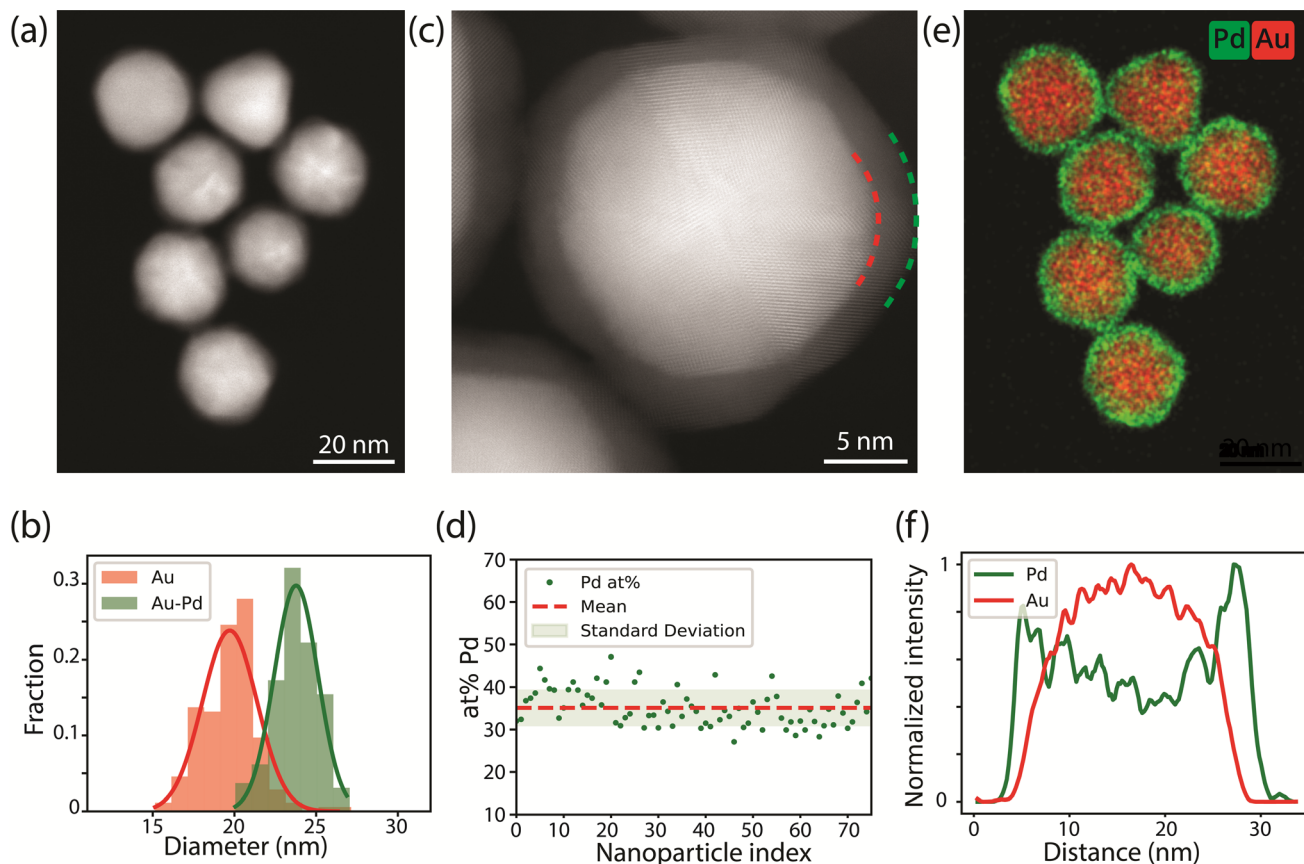


Fig. 1 Electron microscopy characterization of the as-synthesized Au-Pd core-shell nanoparticles. (a) Representative HAADF-STEM image. (b) Size distribution of the Au (red) and Au-Pd (green) NPs, obtained by measuring more than 200 NPs. (c) High resolution HAADF-STEM image of the core-shell NPs, clearly showing the 1.7 nm Pd-shell surrounding the 19.7 nm Au-core, and the overall twinned symmetry of the Au-core Pd-shell particle. The red and green dashed lines highlight the interface between the Au-core and the Pd-shell and the end of the Pd-shell respectively. (d) Pd content (in atom percent, at%) of 75 individual NPs, derived from EDX analysis. (e) Corresponding EDX map of (a) showing Au and Pd in red and green, respectively. (f) Line scan of the NP in the bottom of the EDX map in frame (e).

atomic% Pd obtained from bulk analysis with inductive coupled plasma (ICP). Prior to all following TEM experiments, the polyvinylpyrrolidone ligands present on the surface of the NPs were removed using an activated carbon method.⁴³

The Au-Pd core-shell NPs were supported on a silica support for the CO-DRIFTS measurements. TEM images show that the Au-Pd/SiO₂ were homogeneously distributed over the support and that the core-shell structure was maintained after calcination at 300 °C for 3 h (Fig. S1†). This thermal treatment was sufficient to remove the PVP ligands present on the NPs' surface, as confirmed by thermogravimetric analysis (Fig. S2†). Note that PVP is more easily removed from the metal surface and that any remainders are likely to be on the silica support only.⁴⁴ Thus, we expect that the Au-core Pd-shell NPs had a ligand free metal surface, both during CO-DRIFTS and TEM analysis.

Ex situ TEM analysis of the alloying process

To study the effect of the gas atmosphere on the alloying process in Au-Pd NPs, *ex situ* analysis after treatment under isothermal conditions (375 °C) in H₂ and O₂ were conducted.

The results in Fig. 2 show that depending on the gas atmosphere, the metal distribution of Au-Pd core-shell NPs differed after a 1 hour exposure to 375 °C. EDX mapping after heating under H₂ (10% balanced in Ar) shows uniform distribution of Au and Pd within each particle (Fig. 2a). The EDX line scans show overlapping Pd and Au signals, which confirms that the alloying process was nearly completed. In contrast, the core-shell structures persisted after 1 hour exposure to 375 °C under O₂ (10% balanced in N₂). Both the EDX maps (Fig. 2a) and corresponding line scans illustrate the remaining Pd surface enrichment with the Pd line profile clearly extending outside the Au profile. This suggests that the O₂ atmosphere partially stabilized the core-shell structure and retarded the alloying process.

In fact, the NPs heated under O₂ were less alloyed than those heated under H₂ over a range of heating times (5 to 90 min) and eventually completely alloyed after 12 h of heating. Fig. 2b shows the evolution of the degree of core-shell structure upon heating at different times. A value of 1 represents a core-shell structure and a value of 0 a full alloy. A value smaller than 0 would indicate Au enrichment in the NP surface. The degree of core-shell structure was derived by taking line scans from

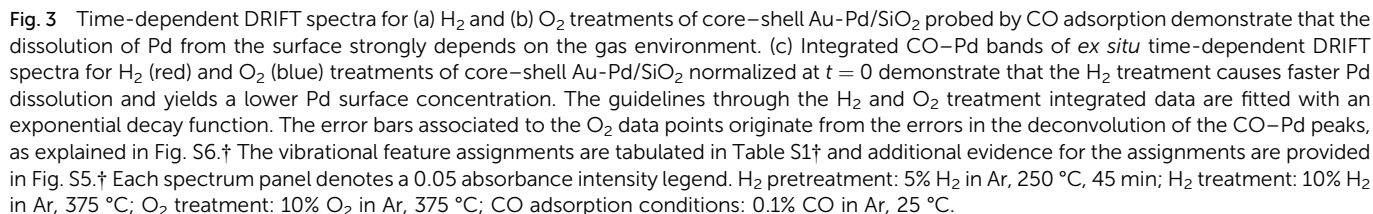




multiple particles in the EDX maps after heating for different amount of time and dividing the width of the Pd profile and the width of the Au profile. In the core-shell particle, the width of the Pd signal is larger than that of Au. Upon alloying the Pd and Au profile start to overlap. This approach is based on a previously validated methodology³⁵ and is further described in the experimental methods. Examples of EDX maps and line scans after different heating times can be found in Fig. S3.† Upon heating, the NPs evolve from a core-shell structure to an alloy, with the degree of core-shell structure decreasing from 1 to 0. The degree of core-shell structure was consistently higher under O₂ than under H₂, indicating that the alloying kinetics were slower under O₂. Due to the limited number of data points and the variation in alloying speed between the individual nanoparticles, the error bars are relatively large, preventing further quantification of the alloying kinetics based on these *ex situ* data. Hence, our *ex situ* results indicate that (1) the gas atmosphere influences the dynamics of the alloying process and (2) the AuPd alloy, which is the thermodynamically stable composition in the whole temperature range,⁴⁵ is eventually obtained in the NP bulk regardless of gas atmosphere.

To understand the impact of the gas atmosphere on the surface alloying of the core-shell Au-Pd/SiO₂ nanoparticles, we employed diffuse reflectance infrared Fourier transform spectroscopy measurements using CO adsorption (CO-DRIFTS). First, the samples were pretreated with H₂ (5% balanced in Ar, 250 °C, 45 min). Then, a series of H₂ or O₂ treatments (10% balanced in Ar, 375 °C) followed by a CO treatment (0.1% balanced in Ar, 25 °C) were performed to determine the relative population of adsorbed CO at distinct surface sites as a function of gas atmosphere and time. The surface structure was not

Sequential thermal treatments in H₂ at 375 °C over the course of 60 min show rapid Pd dissolution (Fig. 3a). A one-minute H₂ treatment induced a shift of CO-Pd⁰, CO-2Pd⁰, and CO-3Pd⁰ to lower frequency and a decrease in intensity of all features. A shift to lower frequency of CO bound to metallic Pd is consistent with previous observations attributed to alloying of the metallic Pd surface with Au.^{24,47} An incremental decrease in intensity of CO-Pd⁰, CO-2Pd⁰, and CO-3Pd⁰ was observed up to a 10 minutes H₂ treatment after which the relatively low intensity of the peaks is unchanging up to 60 minutes. Our observation of an exponential decay in total CO-Pd absorbance is consistent with first-order kinetics for surface Pd dissolution into bulk Au which has been observed on single crystal Pd/Au(111).²⁴ The limited ability of CO to adsorb to metallic Au under our CO adsorption conditions⁴⁹ cannot provide direct evidence of an Au-rich surface. However, the decrease in total CO-Pd absorbance from 100 to 15% is consistent with a transition from a Pd- to Au-rich surface and occurs within the first 10 minutes of thermal treatment (Fig. 3c).



We find clear differences in the CO-DRIFT spectra obtained upon thermal treatment in H₂ and O₂, providing insight in the

The effect of the gas-phase (H_2 or O_2) on the final surface termination of the Au-Pd NPs was further investigated using DFT calculations. In the calculations, the subsurface layers consisted of 36 atomic% Pd randomly distributed in Au, similar to our experimental system, and only the composition of the surface layer was varied. We considered four different surface layers: full Au (Au_{16}), full Pd (Pd_{16}), 6.25% Pd substituted in Au ($\text{Au}_{15}\text{Pd}_1$), and 12.5% Pd substituted in Au ($\text{Au}_{14}\text{Pd}_2$) (Fig. S7a†).

On these surfaces either a single H_2 or O_2 molecule was adsorbed. For the Pd-containing surfaces, adsorption of H_2 and O_2 always occurred on the Pd-site. On the Au_{16} surface adsorption took place on a Au-site. The full set of adsorption configurations are depicted in Fig. S8†. Note that the effect of subsurface Pd on the adsorption energies was negligible in comparison with the adsorption energies on the surface Pd (Fig. S9†).

The adsorption energies $E_{\text{ADS}}^{\text{H}_2}$ and $E_{\text{ADS}}^{\text{O}_2}$ for the four different surface compositions are shown in (Fig. S7b†) and become more negative with an increasing Pd surface fraction, indicating stronger binding on a Pd-enriched surface. The minimum Pd chemical potential μ_{Pd} was found for the $\text{Au}_{15}\text{Pd}_1$ -surface in O_2 (Fig. S7c†). O_2 adsorption on this Pd-monomer surface results in a μ_{Pd} that is 0.12 eV lower than on the $\text{Au}_{14}\text{Pd}_2$ -surface containing a Pd-dimer. Interestingly, the difference between the $\text{Au}_{15}\text{Pd}_1$ and $\text{Au}_{14}\text{Pd}_2$ -surface is much smaller in H_2 , only 0.02 eV. Thus, although $\text{Au}_{15}\text{Pd}_1$ -surface has a minimal μ_{Pd} in both O_2 and H_2 , the driving force for Pd surface segregation is larger in O_2 . Additionally, the DFT results indicate that Pd monomer formation is energetically favored. Both findings are consistent with the CO-DRIFTS results in Fig. 3, indicating a 15% Pd monomer concentration upon thermal treatment in O_2 . The DFT results are

also in line with previous theory studies indicating that Pd oxide formation stabilizes Pd at the nanoparticle surface under O_2 .²⁷

Direct visualization of nanoparticle bulk alloying with *in situ* gas-phase STEM

To directly visualize the alloying process in the Au-Pd NPs, *in situ* STEM experiments under atmospheric pressure of 10% H_2 or 10% O_2 in Ar were performed. The NPs were heated to 375 °C for a total time of 60 minutes and were cooled to room temperature at 18 different time steps to image. Cooling to room temperature for imaging allowed us to pause the alloying process and image several regions in the sample. This approach has previously been used for *in situ* heating TEM of Au-Ag NPs in vacuum.^{36,37,55} Here, we followed the alloying TEM process of 25 NPs. Since fast alloying (<10 minutes) was expected based on the CO-DRIFTS results, short heating steps were chosen in the first 15 minutes of the experiment. Fig. S10† depicts the temperature profile and Fig. S11† shows field of view of the regions followed through the experiment under H_2 . The NPs were stable and did not move over the SiN window. Only the few NPs that were very close together sintered during the pretreatment.

Example images of a Au-core Pd-shell NP at $t = 0$ min and of the corresponding alloyed Au-Pd NP at $t = 60$ min are shown in

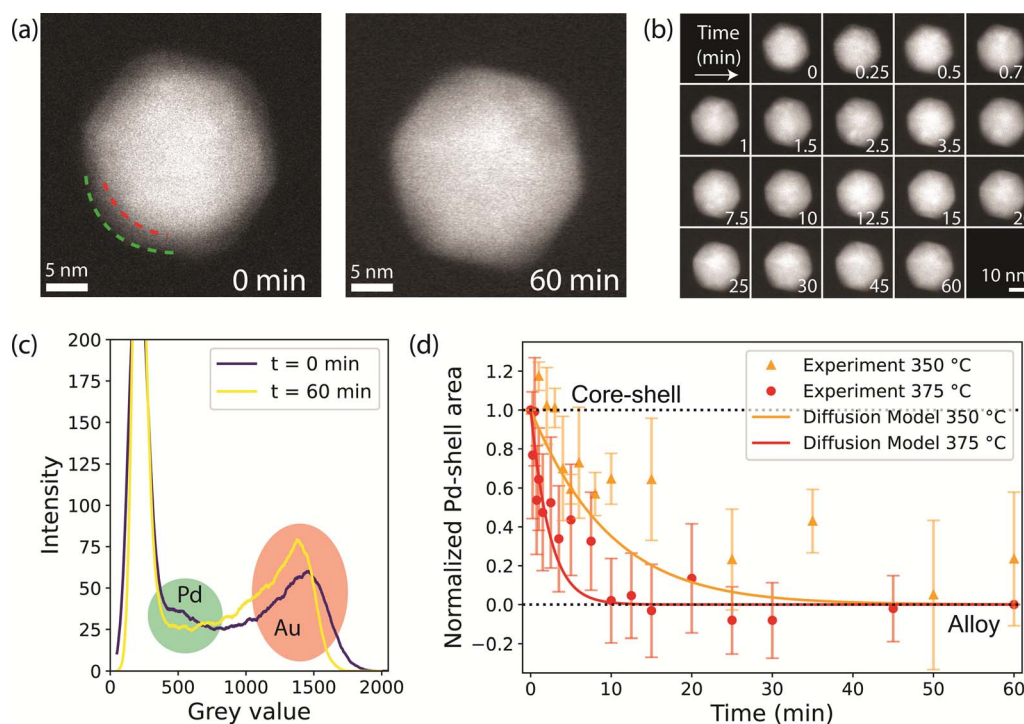


Fig. 4 *In situ* STEM visualizing the alloying of Au-Pd core-shell NPs in H_2 (10% H_2 in Ar, 375 °C). (a) Example HAADF-STEM image of a NP at $t = 0$ min (left), with the red and green dashed lines highlighting the interface between the Au-core and the Pd-shell and the nanoparticle surface, respectively, and at $t = 60$ min (right), with the alloying process completed. (b) Images of the same NP obtained at different total time (min) at 375 °C. (c) Grey-value histograms of the NP in (a) at $t = 0$ (purple) and $t = 60$ minutes (yellow), with the Pd and Au-related grey-values marked in green and red, respectively. (d) Normalized Pd-shell area decrease in a 10% H_2 in Ar atmosphere of 1 bar at 350 °C (orange circles) and 375 °C (red circles), indicating a shift from a core-shell structure to an alloy. The error bars represent the standards deviation of the Pd shell related area across measurements of multiple nanoparticles. The solid lines show the alloying kinetics over time calculated from Fick's diffusion law model at 350 °C (orange) and 375 °C (red). The diffusion coefficients were calculated from literature values (see Methods) and at $t = 0$ they are $3.18 \times 10^{-21} \text{ m}^2 \text{ s}^{-1}$ and $1.30 \times 10^{-20} \text{ m}^2 \text{ s}^{-1}$ for the diffusion of Pd into Au, and $1.73 \times 10^{-26} \text{ m}^2 \text{ s}^{-1}$ and $1.11 \times 10^{-25} \text{ m}^2 \text{ s}^{-1}$ for the diffusion of Au into Pd, at 350 and 375 °C, respectively.



Fig. 4a. The full set of images for this particle is shown in Fig. 4b, revealing the evolution of the NP during the different time steps in an H_2 atmosphere at 375 °C. The alloying dynamics were monitored by quantifying the disappearance of the Pd-shell in the HAADF-STEM images. Specifically, this was done by determining the loss of the grey value intensities corresponding to the Pd-shell. A core-shell structure, as shown in Fig. 4a, taken at $t = 0$ min, has three regions with distinct grey values: the Au-core in white, the Pd-shell in grey, and the background in black. This contrast predominantly originates from the atomic weight difference between Au and Pd. To illustrate this, the grey-value histograms for each region are shown in Fig. S12.† In the histogram of a NP, the Pd related grey values show as a shoulder between the darker pixels belonging to the background and the brighter pixels belonging to Au (Fig. 4c).

Upon alloying, the Pd-shell disappears, which leads to a decrease in the number of pixels corresponding to Pd in the histogram. An example set of histograms for one NP is shown in Fig. S13,† where the highlighted region indicates the grey values associated with the Pd-shell. It is important to note that apart from Z-contrast between Au and Pd, there is also thickness contrast, that intrinsically leads to darker pixels at the NP edges even when the metal distribution is uniform. To account for this, we use the histogram at $t = 60$ min as a baseline and compute the area between the histogram lines at each time point and the final state ($t = 60$ min) for the grey values that correspond to the Pd-shell. A detailed explanation of the method is found in the ESI† (Fig. S12–S15†). To validate our data analysis approach, we performed an *in situ* heating experiment under vacuum following the alloying process both by EDX (as previously performed in ref. 35) and HAADF-STEM imaging. This allowed us to benchmark the results obtained from the new HAADF-STEM data analysis methodology with an already established method.³⁵ The degree of alloying *vs.* temperature data in Fig. S16† confirm that the results were very similar and that our HAADF-STEM analysis method can be used to reliably quantify the alloying dynamics.

Fig. 4d shows the decrease in Pd-shell area averaged over all 25 particles as a function of time during thermal treatment in H_2 at 375 °C. The data indicate that the alloying was complete within 10 minutes, closely matching our CO-DRIFTS results. We verified that the alloying kinetics derived from the *in situ* TEM data were not impacted by electron beam irradiation. To do so, a total of 20 NPs were imaged less frequently during the same experiment, thereby accounting for less than half of the electron dose. The Pd-shell grey value area evolved in the same manner as for the more irradiated NPs (Fig. S17†). Furthermore, EDX showed fully alloyed NPs in both irradiated and non-irradiated areas (Fig. S18†). These results therefore proof that the electron beam did not substantially influence the alloying dynamics.

When the same experiment was performed under 10% O_2 , the normalized Pd-shell related area in the histograms decreased in slightly faster than in our H_2 experiments (Fig. 5a and S19†), and the alloying process seemed to be completed after 5 minutes. After 60 minutes at 375 °C, the NPs were fully alloyed, both in irradiated and non-irradiated regions

(Fig. S20†). The alloying under O_2 *in situ* occurred faster than expected based on our *ex situ* experiments. We investigated if this discrepancy could be attributed to electron beam irradiation. Consistent with the H_2 experiment, regions imaged with less than half of the dose evolved in a similar way compared to NPs that received a higher electron dose (Fig. S19b†). To further investigate the possibility of electron beam effects, the experiment was repeated without beam irradiation following the same heating profile as used in the *in situ* TEM measurements. EDX after thermal treatment confirmed that all NPs were fully alloyed (Fig. S21†), and hence no indications for the beam effects altering the alloying process. However, the alloying kinetics did depend on the oxygen pressure used. When repeating our *in situ* TEM measurement with a higher partial pressure of oxygen (100%) the alloying process slowed down, and the nanoparticles alloyed after 30 minutes, coinciding with our CO-DRIFTS results (Fig. 5). Hence, a higher partial pressure of O_2 was needed in the *in situ* TEM setup to observe the effect of the O_2 seen in the CO-DRIFTS and *ex situ* TEM experiments.

We also investigated the temperature effect on the alloying dynamics by performing an additional experiment in H_2 at 350 °C. The Pd-shell related area decreased significantly slower compared to the experiment at 375 °C (Fig. 4d), indicating substantially slower alloying. This result demonstrates that (1) the temperature significantly affects the alloying dynamics of the NP bulk and that (2) our data analysis methodology is capable of detecting significant speed differences in the alloying process. Interestingly, our data can be described well by a numerical model based on Fick's second law. In short, the model is used to calculate the amount of Au and Pd atoms passing through a Au-Pd interface at each time point, and has previously been applied to describe the alloying of Au-Ag nanorods.³⁵ The diffusion coefficients were calculated with bulk data of the frequency factor (D_0) and activation energy (Q),⁵⁶ corrected for nanoparticle finite-size effects, which account for a ~4% decrease in Q .^{35,57} The model relies on iterating the calculation of the quantity of atoms that diffuse at each time step, followed by updating of the composition of the interface. Using this model, we find that the diffusion

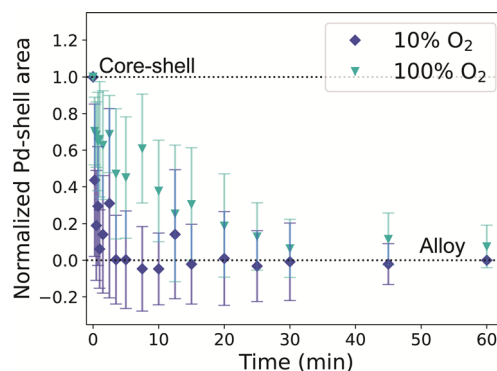


Fig. 5 *In situ* STEM visualizing the alloying of Au-Pd core-shell NPs in O_2 at 375 °C. The plot shows the normalized Pd-shell area as a function of time for a 10% O_2 in Ar (dark blue) and 100% O_2 (light blue) gas atmosphere. The error bars represent the standard deviation of the Pd shell related area across measurements in multiple nanoparticles.



coefficients at $t = 0$ are $3.18 \times 10^{-21} \text{ m}^2 \text{ s}^{-1}$ and $1.30 \times 10^{-20} \text{ m}^2 \text{ s}^{-1}$ for the diffusion of Pd into Au, and $1.73 \times 10^{-26} \text{ m}^2 \text{ s}^{-1}$ and $1.11 \times 10^{-25} \text{ m}^2 \text{ s}^{-1}$ for the diffusion of Au into Pd, at 350 and 375 °C, respectively. The modeled diffusion behavior at 350 and 375 °C is shown in Fig. 4d, and it shows a good agreement with datapoints derived from the *in situ* TEM data both at 350 and 375 °C. This validates our *in situ* TEM approach, and it also shows that the modeled diffusion is accurate for nanoparticle systems.

Discussion

CO-DRIFTS showed that the final surface composition of the Au-Pd nanoparticles is highly affected by the gas atmosphere, and that Pd is, in part, stabilized at the surface after thermal treatment in an O₂ atmosphere. Under H₂, after 1 h at 375 °C, the Au-Pd NPs surface was Au-rich and 85% of the Pd had diffused to the bulk of the NPs. Contrarily, under O₂, the surface still consisted of a mixture of Au and oxidized Pd. While Pd segregation into the bulk has also been observed under O₂, it was less prominent than under H₂, with a loss of 77% of the original surface Pd. These findings are consistent with literature reporting on the gas dependent segregation effects in Au-Pd alloy nanoparticles. Similar to our work, Pd-enriched surfaces have been observed after thermal treatment in O₂ (ref. 30) whereas Au-rich surfaces have been reported for thermal treatment in H₂ (ref. 25) or vacuum (ref. 23). Experimental EXAFS and XPS results have shown that Pd is oxidized under O₂, and catalytic testing confirmed that O₂ treated Au-Pd NPs are more active in hydrogenation and oxidation reactions than H₂ treated Au-Pd NPs, possibly due to Pd surface enrichment upon O₂ treatment.^{27,58} Thus, our findings demonstrating the difference in Au-Pd surface composition upon H₂ and O₂ treatment are in good agreement with previous studies on dilute Au-Pd nanoparticle systems (<10 atomic% Pd) and confirm that the same trends hold for higher Pd-content Au-Pd systems.

Our time-lapse CO-DRIFTS measurements (Fig. 3) allowed direct assessment of the Pd surface segregation kinetics in O₂ and H₂. Under H₂, Pd quickly segregated to the NP bulk and the surface composition stabilized after 10 minutes at 375 °C. Under O₂, the Pd segregation kinetics were two times slower, and a stable composition only reached after 30 minutes at 375 °C. Thus, Pd dissolution was directly affected by the gas atmosphere, where Pd stabilization at the surface affects its dissolution kinetics into the nanoparticle bulk. Our *ex situ* TEM experiments (Fig. 2) validate the CO-DRIFTS results and confirm that the alloying happens slower under O₂ than under H₂. The differences in time scales for surface Pd dissolution may be attributed to the higher surface sensitivity of CO-DRIFTS, which only analyzes the composition of the outermost atomic layer. To the best of our knowledge, only one study has previously provided dynamic measurements on a Pd/Au (111) single crystal model systems,²⁴ but this had never been extended to Au-Pd nanoparticle systems. Our observation of an exponential decay in total CO-Pd absorbance is consistent with first-order kinetics for surface Pd dissolution into bulk Au observed on single crystals,²⁴ suggesting a similar Pd-dissolution kinetic behavior

between single crystals and nanoparticles. Our DFT calculations correctly predict incomplete Pd dissolution from the surface in both gas atmospheres, with more Pd stabilization under O₂, in line with our CO-DRIFTS results and previous DFT studies.²⁷ Moreover, the chemical potential is the lowest for Pd monomers, indicating their stabilization at the surface. This stabilization of the Pd monomers is indeed observed in the CO-DRIFTS, where the CO-Pd peak disappears slower when compared to the CO-2Pd or CO-3Pd peak. The stabilization of surface Pd after O₂ treatment likely leads to enhanced catalytic activity due to Pd surface enrichment.²⁷

Moreover, we directly visualized and quantified the alloying dynamics in the bulk of core-shell bimetallic NPs in gas-phase through *in situ* TEM for the first time by following the changes in the grey value histograms. Previous studies have mainly used aberration corrected environmental TEM to visualize the loss²⁰ or the formation²¹ of a core-shell structure. Studies on the dynamics of the alloying process upon heating have been performed under vacuum, using EDX,³⁵ fast electron tomography,^{36,37} or atomic resolution imaging.^{38,39} These approaches offer a more precise quantification method but are difficult to perform under atmospheric pressure of gas. The main limitation of our data analysis method is the large deviation in the Pd-shell related area among nanoparticles (see error bars in Fig. 4 and 5), which restricts drawing conclusions for small differences in alloying dynamics among nanoparticles. The element-specific Z-contrast in HAADF-STEM mode is complicated by the additional thickness- and diffraction-based contrast, which makes it more difficult to quantify the loss in Pd-shell contrast from our 2D STEM images. The thickness dependent contrast causes the edges of the NP to look darker than the center. On the other hand, the diffraction contrast broadens and displaces the white peak in the histogram. While we correct for the thickness contrast, the Pd-shell related area does not always follow the expected decreasing behavior at individual NPs, resulting in relatively large error bars. We validated our methodology in vacuum and show that the results obtained from the HAADF-STEM data analysis closely match the results obtained from EDX when imaging under vacuum (Fig. S16†). Furthermore, our methodology does not only allow acquiring images at atmospheric pressure, but also drastically reducing the electron dose and acquisition time, especially when compared to EDX mapping. Reducing the imaging time allows imaging at more time points or at different regions of the sample, allowing the *in situ* monitoring of many different particles within the same experiment.

The effect of the gas atmosphere (H₂ or O₂) on the Au-Pd alloying dynamics observed in our *in situ* TEM experiments was smaller than the effect of different temperatures (350 vs. 375 °C). Under 10% H₂, the alloying process was completed after 10 minutes, matching our CO-DRIFTS results (Fig. 4). However, the expected slowdown of the alloying process did not occur under 10% O₂, but only happened when using higher oxygen pressures (100%). The alloying process was significantly slowed down under 100% O₂, and the process was completed after 30 minutes, coinciding with our CO-DRIFTS results (Fig. 5). This likely indicates that a higher partial pressure is needed during the *in situ* TEM measurements to allow the formation of the Pd oxide to



stabilize Pd on the surface of the NPs, and hence retard the alloying. Indeed, the formation of a Pd oxide has been observed *in situ* in Au-Pd nanoparticles under 1 atm O₂.⁵⁹

The alloying kinetics derived from our *in situ* TEM results closely match our model based on Fick's diffusion model and show that a decrease in temperature of 25 °C leads to a pronounced slowdown in alloying dynamics. This diffusion model has previously been validated for Au-Ag core-shell nanorods of ~70 nm length.³⁵ Furthermore, for Au-Ag NPs at 450 °C, the diffusion simulation yields diffusion coefficients with the same order of magnitude to those previously observed experimentally.^{36,37} With this *in situ* study, we show that the alloying behavior of Au-Pd NPs is also described well with this model, when considering the physical properties of the metals and a correction for NP size and shape. We therefore expect that this diffusion model is able to predict alloying dynamics in a whole range of metals, temperatures, and nanoparticle sizes. Given the superior catalytic activity of Au-Pd core-shell nanoparticles with respect to their alloyed counterparts,¹⁷ this model can help in designing the catalytic testing conditions such that alloying and the associated loss in catalytic activity do not occur during the reaction.

Conclusion

Our combined CO-DRIFTS and *in situ* TEM study highlights the crucial role of the gas atmosphere on the surface alloying of Au-Pd NPs, and its more limited impact on the alloying dynamics of the NP bulk. Exposure to H₂ at 375 °C led to faster Pd dissolution from the surface than under O₂ showing that Pd was partially stabilized at the NP surface in an O₂-rich atmosphere. The DFT-calculations indicate that this can be ascribed to a lower chemical potential for Pd-monomers at the surface under an O₂ atmosphere. *In situ* TEM allowed the direct visualization of the alloying dynamics in the nanoparticle bulk and showed that the diffusion in the NP bulk is mostly impacted by the heating temperature, and to a certain extent also by the partial pressure of O₂. Through our data analysis methodology, we were able to extract the alloying kinetics from the HAADF-STEM images allowing for short imaging times, low electron doses, and for a large number of nanoparticles followed at multiple time steps. The experimentally attained diffusion kinetics were captured well by a numerical model based on Fick's second law of diffusion, and we expect that this model can be translated to other nanoparticle systems as well. Altogether, our work introduces a new approach to quantitatively assess the alloying dynamics in bimetallic nanoparticle systems through *in situ* TEM in reactive gasses and presents new insights into gas dependent segregation phenomena on surface and bulk alloying of multi-metallic nanoparticle systems.

Experimental section

Chemicals

Trisodium citrate dihydrate (SC, ≥99.0%), tannic acid (TA), potassium carbonate (K₂CO₃, ≥99.0%), chloroauric acid (HAuCl₄, 99.9%), sodium tetrachloropalladate (II) (Na₂PdCl₄, 98%) and polyvinylpyrrolidone (PVP, MW 55 000) were obtained from Merck.

Aerosil OX 50 was purchased from Degussa. Correl-like activated carbon was purchased from Norit. Ultrapure water with a resistivity of 18.2 mΩ cm⁻¹ (Millipore Milli-Q grade) was used. All chemicals were used without further purification. The glassware used for AuNP synthesis and Pd overgrowth was cleaned overnight with aqua regia (HCl/HNO₃ mixture in a 3 : 1 ratio by volume), rinsed thoroughly with Milli-Q water and dried in an oven.

Au NP synthesis

Gold nanoparticles of around 20 nm in diameter were synthesized following the procedure of Piella *et al.*⁴² Briefly, a solution of 450 mL of sodium citrate 2.20 mM, 0.30 mL 2.50 mM tannic acid and 3.00 mL 150 mM potassium carbonate were heated to 70 °C while stirring vigorously. Then, 3.00 mL 25.0 mM HAuCl₄ were added, and the temperature was maintained for 10 minutes. Within 3–5 min, the color changed from pale yellow to red, indicating the formation of small gold nanoparticles, and the pH shifted from 10 to 8. Then, 165 mL of this seed solution were extracted, and 165 mL of 2.20 mM sodium citrate solution were added. When the temperature reached 70 °C again, 1.50 mL 25.0 mM HAuCl₄ were injected twice with a 10 minutes interval. This process was repeated for an additional 8 times, until a total of 30.0 mL 25.0 mM HAuCl₄ had been added. After the last step, the reaction mixture was allowed to cool down and stored in the dark, and the UV-vis spectra of each growth cycle was measured.

Pd overgrowth

For the Pd overgrowth, 1.50 mL polyvinylpyrrolidone solution (1 g/10 mL H₂O, Mw = 55 000 g mol⁻¹) were added to 150 mL of the as-synthesized sodium-citrate-capped gold nanoparticles and stirred overnight. The pH was adjusted to 4 with 0.10 M HCl (2.50 mL) to ensure a slow enough reaction rate during the Pd overgrowth.⁵ Next, 7.80 mL Na₂PdCl₄ 10 mM were added to the solution and it was stirred for 5 minutes before rapidly adding 7.80 mL of 40 mM ascorbic acid under vigorous stirring to ensure homogeneous reduction of the Pd precursor onto the Au-cores. The solution was left stirring overnight at 400 rpm. Afterwards it was centrifuged at 12 000 rcf for 1 h and kept in half the volume of water for storage in the fridge.

Deposition on a silica support

200 mg of (Aerosil OX50) were mixed with 10 mL EtOH and sonicated for 15 minutes to ensure proper dispersion. Half of the synthesis volume of the Au-Pd NPs was centrifuged and resuspended in 20 mL EtOH and added to the silica suspension. After sonicating for approximately one hour, the purple dispersion was divided in two 50 mL centrifuge tubes, and 15 mL of toluene was added to each tube as an antisolvent, leading to homogeneous deposition of the AuPd NPs on the silica support.¹⁷ Then, it was centrifuged at 2000 rcf for 5 minutes and the clear supernatant was removed. The pellet was dried in an oil bath at 60 °C overnight. To remove the ligands, the supported AuPd/SiO₂ NPs were heated in a U-shaped reactor with a ramp of 2 °C per minute and a flow of 100 mL min⁻¹ of 10% O₂ balanced in N₂ and were kept at 300 °C for 3 hours. This resulted in a 4.7 wt% of metal on the Aerosil OX50 support.



Diffuse reflectance infrared Fourier transform spectroscopy (CO-DRIFTS) experiments. DRIFTS experiments were carried out in a low-dead volume reaction chamber (Harrick Scientific) equipped with ZnSe windows, mounted inside the sample compartment of a Bruker Invenio FT-IR spectrometer using a Praying Mantis diffuse reflectance accessory (Harrick Scientific). The gas flow through the DRIFTS reactor was controlled by a rapid switching gas handling system.⁶⁰ The catalyst sample was prepared by pressing and sieving between 304 stainless-steel meshes (50 × 50 mesh and 70 × 70 mesh). The DRIFTS reactor was loaded with approximately 135 mg of 46 grit SiC, an inert support with high thermal conductivity, followed by approximately 15 mg of the prepared AuPd/SiO₂ catalyst. Large temperature gradients can exist between the thermocouple contact in a DRIFTS reactor cell and the catalyst surface temperature exposed to the infrared beam.⁶⁴ Therefore, the temperature was calibrated by accounting for the thermal gradient between the thermocouple and the catalyst surface by using an optical pyrometer with an emissivity of 0.95. All DRIFTS experiments used a total volumetric flow rate of 50 mL min⁻¹. Each absorbance spectrum was obtained by averaging 200 background and sample scans at a resolution of 4 cm⁻¹ using a liquid-nitrogen-cooled HgCdTe (MCT) detector, while the Praying Mantis diffuse reflectance accessory and FT-IR spectrometer was purged with dry N₂. The AuPd/SiO₂ pretreatment involved reduction under H₂ (5% H₂ in Ar, 250 °C) for 45 min. The sequential H₂/O₂ treatments (10% H₂/O₂ in Ar, 375 °C) involved heating, holding for the denoted time and cooling in the H₂/O₂ environment. CO-DRIFTS measurements were taken after 0, 1, 5, 10, 15, 30 and 60 minutes (total time at 375 °C). The sample measurements were acquired at 25 °C once the saturation of adsorbed CO in a 0.1% CO in Ar was achieved. Quantitative analysis of DRIFTS spectra can be prone to errors because adsorbate concentration is truly neither proportional to absorbance intensity nor Kubelka–Munk (KM)^{62,63} intensity but can be nearly proportional to one intensity unit depending on the relative reflectivity of the DRIFTS features.⁶⁴ The high relative reflectivity of our DRIFTS features demonstrate the integration of absorbance intensity is expected to be nearly proportional to concentration while Kubelka–Munk intensity is less appropriate (Fig. S22†).⁶⁴

Around 5 mL of AuPd NPs were centrifuged and resuspended in EtOH and then drop-casted onto a carbon coated TEM grid (H₂ experiments) or a SiN TEM chip (O₂ experiments). Prior to heating, the PVP ligands were removed by washing the grid with

The degree of core-shell structure in Fig. 2c was calculated by a procedure similar to literature.³⁵ At each time point, a minimum of 8 line scans (30×5.5 nm) from different nanoparticles were taken from the EDX maps, to get some statistics on the distribution of Au and Pd within the nanoparticle. Note that after oven heating, particles that were close to each other aggregated. The line scans used for analysis were obtained from non-aggregated particles, or taken in the direction where the particles did not touch each other. The degree of core-shell structure was calculated by the division of the full width at half maximum of Au by that of Pd in the line scans. It was then normalized so that the degree of core-shell structure is 1 for the core-shell structure, and 0 when the alloying is complete. Note that it is the same procedure than the one in ref. 35 for the degree of alloying, but the normalization process is different.

Transmission electron microscopy (TEM) was performed on a Talos F200X (Thermo Fisher Scientific) operated at 200 kV. Particle size distributions were obtained by measuring 100 to 200 nanoparticles from bright field TEM images. First, a threshold was applied to mask the nanoparticles, and the size was analyzed using the “analyze particles function” in ImageJ.⁶⁵ For energy-dispersive X-ray (EDX) mapping, the microscope was operated in scanning-TEM (STEM) mode, with a screen current of 1.18 nA and a dwell time of 2–5 μ s per pixel, a pixel size of 0.3827 nm, 512×512 pixels per image and a camera length of 98 mm. The total time of an EDX map was of 15 minutes. The Super-XTM EDX detector present in the microscope was used to collect the EDX signal and it was quantified using the Velox software. The EDX maps were prefiltered by averaging 5 pixels before quantification. Both Au and Pd were quantified using the L-lines. High resolution electron microscopy was acquired with a double aberration corrected Spectra 300 (Thermo Fisher Scientific) with 2048×2048 pixels of 13.6 pm pixel size, a dwell time per pixel of 2.5 μ s, a screen current of 0.42 nA, camera length of 90 mm and convergence angle of 30 mrad.

The *in situ* gas-phase STEM measurements were performed using a Protochips Atmosphere system and STEM/EDX optimized chips on a Talos F200x. The chips were glow discharged for 90 s with 10 V of tension. Then, several droplets of Au-Pd NPs dispersed in ethanol were drop-casted on the top chip. The chip was washed with the activated carbon method described in the previous section and let dry before assembling the holder. Once in the microscope, the holder was pump-purged three times to 0.5 mbar to remove air from the system and leak checking was

performed. Then, the sample was heated to 250 °C in 5% H₂ in Ar for 45 minutes to ensure the removal of water from the system and reduction of potentially oxidized Pd species, and cooled down to room temperature. For the experiments under oxygen atmosphere, the holder was pump/purged with Ar 10 times to 20–200 mbar to ensure H₂ removal and to avoid its mixture with O₂. During the experiment, the gasses (10% H₂ or O₂ diluted in Ar or 100% O₂) were introduced at 1 bar with a 0.1 sccm flow rate. Then, the sample was heated to 375 °C for a total time of 1 h. The heating was interrupted after 0.25, 0.5, 0.75, 1, 1.5, 2.5, 3.5, 5, 7.5, 10, 12.5, 15, 20, 25, 30, 45 and 60 minutes, and the imaging was performed at room temperature at 4 or 5 different regions, accounting for 29 NPs under 10% H₂ and 31 NPs under 10% O₂ and 33 NPs under 100% O₂. To study the effect of the electron beam dose, some regions were only imaged after 1, 5, 15, 30, 45 and 60 minutes in the 10% H₂ and O₂ experiments. The temperature ramp was always 10 °C s^{−1} and the beam was always blanked at high temperature. The imaging was performed at room temperature in HAADF-STEM mode, camera length of 125 mm, 2048 × 2048 pixels with 0.09567 nm pixel size, a dwell time of 5 μs per pixel, a screen current of 25 pA, accounting for a dose rate of 41.3 e[−] (Å² s)^{−1}. Moreover, 20 images taken with 100 ns dwell time per pixel were taken and put together using a DCFI algorithm in VELOX. The total dose was 15 570 e[−] Å^{−2} for the regions with 18 heating steps and 6055 e[−] Å^{−2} for the regions with 7 heating steps. After the experiment, the chips were disassembled and kept for inspection and EDX mapping under vacuum. An extra experiment was performed at 350 °C, where after the pretreatment, the heating was interrupted after 1, 2, 3, 4, 5, 6, 8, 10, 25, 35, 50 and 60 minutes. A total of 13 NPs were followed through the experiment. To ensure complete alloying after the experiment, the NPs were heated to 450 °C for 5 minutes and imaged afterwards. The temperature plots of the experiments are shown in Fig. S10.†

***In situ* heating transmission electron microscopy (under vacuum)**

The *in situ* heating TEM experiments under vacuum were performed using a DENS solutions Wildfire heating holder, in a similar fashion as in literature.³⁵ The sample was drop-casted on the chip, and was heated step-wise from room temperature to 600 °C with steps of 50 °C. At each temperature, an EDX map and several HAADF-STEM images were taken.

Data analysis

Extended information of the methodology of the *in situ* STEM data analysis is presented in the ESI.†

Diffusion model

A numerical diffusion model based on the approximation of Fick's second law of diffusion was used to describe the alloying behavior of the Au-core Pd-shell nanoparticles. This model has previously been used and validated for the alloying in Au-core Ag-shell nanoparticles.³⁵ Briefly, the numerical calculations to compute the number of Pd and Au atoms diffusing through the

Au–Pd interface per time step at a specific temperature using the following formula (eqn (1)):

$$\frac{n_{\text{Pd}}}{\Delta t} = \left(\frac{A}{r}\right) \cdot D_0^{\text{Pd}} e^{-\frac{Q_{\text{Pd}}}{RT}} \cdot |C_{\text{core}}^{\text{Pd}} - C_{\text{shell}}^{\text{Pd}}| \quad (1)$$

where A is the interface area, r the radius of the NP, D_{Pd} the frequency factor, Q_{Pd} the activation energy, R the gas constant, T the temperature, $C_{\text{core}}^{\text{Pd}}$ the Pd concentration of the core and $C_{\text{shell}}^{\text{Pd}}$ the Pd concentration of the shell. The bulk frequency factor D_0^{Pd} and activation energy Q_{Pd} were taken from literature and were 0.2 cm² s^{−1} and 47 kcal mol^{−1}, respectively.⁵⁶ An analogous formula was used for the diffusion of Au through the interface, with frequency factor D_0^{Au} and activation energy Q_{Au} 0.12 cm² s^{−1} and 62 kcal mol^{−1}, respectively.⁵⁶ The bulk Q values were corrected for the NP size effects in the same manner as in literature (eqn (2)), and consist on the calculation of a shape factor (α_{shape}), described in eqn (3), where D is the nanoparticle diameter, γ_s and γ_l are the surface energies in the solid and liquid phases (obtained from²²), respectively, S the surface area of the NP, V the volume and $\Delta H_{\text{m},\infty}$ the bulk melt enthalpy.^{35,57} After this finite-size effects correction, Q_{Au} and Q_{Pd} were 59.6 and 45.2 kcal mol^{−1}, respectively.

$$\frac{Q_{\text{NP}}}{Q_{\text{bulk}}} = 1 - \frac{\alpha_{\text{shape}}}{D} \quad (2)$$

$$\alpha_{\text{shape}} = \frac{D(\gamma_s - \gamma_l)}{\Delta H_{\text{m},\infty}} \times \frac{S}{V} \quad (3)$$

The initial state is a pure Au-core and pure Pd-shell phase connected through an interface. For each time step of 1 s the number of Au and Pd atoms passing through the interface was calculated, and the composition of both sides of the interface was updated. The Pd atoms migrate from the shell to the core and the Au atoms from the core to the shell. These steps were iterated until the fully alloyed state was reached, in which the concentration of Pd in the core was equal to that in the shell. In our case, the calculations were run for $T = 350$ and 375 °C.

Density functional theory (DFT) calculations

All DFT calculations were performed using the FHI-aims software⁶⁶ and the Perdew–Burke–Ernzerhof (PBE)⁶⁷ generalized gradient approximation (GGA) exchange–correlation functional. Long-range van der Waals interactions were corrected with the Tkatchenko–Scheffler method.⁶⁸ The Pd_{0.36}Au_{0.64} composition (suggested by experimental results) was modelled as an 80-atom random alloy in the face-centered cubic (FCC) structure and a lattice constant of 4.076 Å. The (1 1 1) surface was modeled using a periodic slab model with 5 atomic layers and a 4 × 4 (11.54 × 11.54 Å) surface cell. A 16-atom monolayer of pure Au (Au₁₆), pure Pd (Pd₁₆), monomeric Pd in Au (Au₁₅Pd₁), or dimeric Pd in Au (Au₁₄Pd₂) was added on top of the random alloy surface slab (see Fig. S7a†). A vacuum region with a width of ~15 Å was introduced to avoid erroneous interactions. For each of the surface models, the bottom 5 atomic layers were frozen at the ideal bulk positions and the top layer was optimized until the forces were converged to a threshold of 0.01 eV



$\text{\AA}^{-1}$. Brillouin zone sampling was done with a Γ -centered $7 \times 7 \times 1$ k -point mesh. Geometry optimizations used the Broyden–Fletcher–Goldfarb–Shanno (BFGS) optimization algorithm.^{69–73} Geometry optimizations with FHI-aims employed the pre-defined ‘light’ settings (4th order expansion of the Hartree potential, radial integration grids with 302 points in the outer shell, and a tier 1 basis set). The structural energies were subsequently refined with the ‘tight’ settings (6th order expansion, 434 grid points, and a tier 2 basis set). Self-consistency calculations were considered converged when the volume-integrated root-mean square change of the charge density between the present and previous self-consistent field iteration was less than $10^{-6} \text{ eV \AA}^{-3}$ and the total energy changed less than 10^{-6} eV .

The adsorption energy E_{ADS} (eV) of O and H was calculated as (eqn (4))

$$E_{\text{ADS}}^{\text{ad}_2} = E_{\text{ad}_2@\text{surface}} - E_{\text{surface}} - E_{\text{ad}_2} \quad (4)$$

where ad_2 is either H_2 or O_2 , $E_{\text{ad}_2@\text{surface}}$ is the DFT energy of the surface with the adsorbed molecule, E_{surface} is the DFT energy of the clean surface, and E_{ad_2} is the DFT energy of the gas phase H_2 or O_2 molecule.

The chemical potential μ_{Pd} (eV) of the Pd surface atoms was calculated as (eqn (5))

$$\mu_{\text{Pd}}^{\text{surf}} = (E_{\text{ADS}}^{\text{ad}}\{\text{Au}_{N-n}\text{Pd}_n\} - E_{\text{ADS}}^{\text{ad}}\{\text{Au}_N\})/n \quad (5)$$

where $E_{\text{ADS}}^{\text{ad}}$ is the adsorption energy of adsorbate A in eV, N is the total number of surface atoms and n is the number of Pd atoms in the surface. Zero-point energy contributions to the chemical potential were neglected.

Data availability

The data from the Density Functional Theory (DFT) calculations can be obtained from the GitHub repository at <https://github.com/atomisticnet/2024-AuPd-DFT-data>.⁷⁴ The dataset contains atomic structures and energies in the formats used by the FHI-aims software.⁶⁶ The raw *in situ* TEM datasets and grey-value histograms can be obtained from the Dataverse repository.⁷⁵ The python code for the analysis of the *in situ* TEM data and Fick's diffusion model is available upon request to the corresponding author.

Conflicts of interest

There are no conflicts to declare.

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