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# Lifting the discrepancy between experimental results and the theoretical predictions for the catalytic activity of RuO<sub>2</sub>(110) towards oxygen evolution reaction

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**Developing new efficient catalyst materials for the oxygen evolution reaction (OER) is essential for widespread proton exchange membrane water electrolyzer use. Both RuO<sub>2</sub>(110) and IrO<sub>2</sub>(110) have been shown to be highly active OER catalysts, however DFT predictions have been unable to explain the high activity of RuO<sub>2</sub>. We propose that this discrepancy is due to RuO<sub>2</sub> utilizing a different reaction pathway, as compared to the conventional IrO<sub>2</sub> pathway. This hypothesis is supported by comparisons between experimental data, DFT data and the proposed reaction model.**

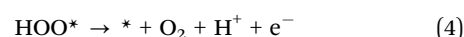
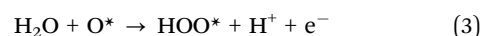
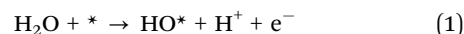
## Introduction

As part of a transition to a future sustainable economy, there is a need for sustainable fuel and energy storage. Hydrogen gas is an ideal candidate for such a fuel and storage compound, as it can be readily produced by the electrolysis of water.<sup>1</sup> Furthermore, the hydrogen is completely sustainable if the source of electricity is renewable.<sup>2</sup> Currently, the major challenge facing widespread electrolyzer use is the sluggish kinetics of the Oxygen Evolution Reaction (OER) at the anode,<sup>3</sup> fundamentally limited by the universal scaling relations.<sup>4,5</sup> Further development of the water electrolyzer thus requires finding efficient and practical catalysts to facilitate the OER. There are three types of water electrolyzers: alkaline water electrolyzers, proton exchange membrane (PEM) water electrolyzers and solid oxide water electrolyzers.<sup>6</sup> Of these three the alkaline water electrolyzer is the most mature and commercialized. Yet PEM technology has many advantages compared to the alkaline electrolyzer. Some examples include a much higher current density, purer gas, a smaller size for the same power and even the ability to operate at high pressure.<sup>6,7</sup> Currently the best candidates for PEM electrolyzer anode material are IrO<sub>2</sub> and RuO<sub>2</sub>, as these are both stable and active.<sup>8–13</sup> Both iridium and

ruthenium are however scarce materials and thus expensive.<sup>14</sup> It is therefore unrealistic to expect that these catalysts can be used on an industrial scale that would have an impact on society.<sup>1,15</sup> Due to the high catalytic performance of IrO<sub>2</sub> and RuO<sub>2</sub>,<sup>9</sup> a deeper understanding of how these catalysts interact with water could greatly improve the search for and development of new efficient OER catalysts.<sup>16</sup> Current theory and computational DFT models form a relatively accurate description of the experimental behaviour of IrO<sub>2</sub>.<sup>9,17,18</sup> However, for the case of RuO<sub>2</sub> there is a glaring discrepancy between theoretical predictions and experimental results.<sup>18–20</sup> This suggests that either the current reaction model is wrong in the case of RuO<sub>2</sub> or the under-evaluation is due to a computational artefact. In this work we propose that the discrepancy is due to RuO<sub>2</sub> utilizing an alternate reaction pathway for oxygen evolution as compared to IrO<sub>2</sub>. We therefore argue that it is not due to a computational artefact.

## Results-discussion

The conventional pathway describing the interaction between water molecules and the surface of an electrocatalyst was suggested in 2004.<sup>21</sup> During this reaction pathway three intermediates are produced *via* four electron–proton pair exchanges between the anode and the electrolyte, presented in the following reactions:



where \* indicates an active site of the surface and HO\*, O\*, HOO\* the adsorbed intermediates on that particular site. The above reaction path describes the oxygen evolution reaction taking place in an acidic environment but it can also be used

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Fig. 1 Illustration of two oxygen evolution reaction pathways. Each step is accompanied with the exchange of a proton–electron pair between the electrolyte and the catalyst surface. (a) The conventional OER path. (b) The proposed reaction mechanism for OER on  $\text{RuO}_2$  surfaces. (c) Illustration of the relevant  $\text{RuO}_2(110)$  surface sites.

for the thermodynamic description of the procedure happening in alkaline environment.<sup>21,22</sup>

A schematic representation of the conventional OER reaction path is depicted in Fig. 1a, while the  $\text{RuO}_2$  surface sites are illustrated in Fig. 1c. First, water approaches the surface and the first intermediate  $\text{HO}^*$  is created. Second, the oxygen forms another bond to the surface losing the remaining proton as shown in eqn (2). Third, a water molecule binds to the surface-bound oxygen, dropping a proton in the process. Simultaneously the double bond of  $\text{O}^*$  breaks, creating the third intermediate  $\text{HOO}^*$ . The fourth and last step, happens while the oxygen atoms bond with each other, breaking their bonds with the surface and the hydrogen respectively. This reaction mechanism succeeds to accurately describe the trends of catalytic activity for the majority of the metal oxides. For example the per site DFT activity of  $\text{IrO}_2$  places  $\text{IrO}_2$  among the best catalysts, in agreement with the experimental observations concerning the performance of the material either in the nanoparticle or the single crystal form.<sup>4,9,17,23</sup>

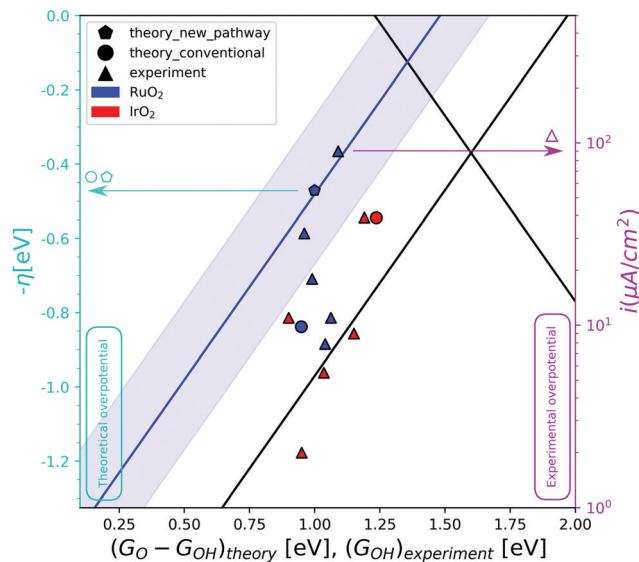
DFT calculations following this model have however been unable to explain the high experimental activity of  $\text{RuO}_2$ .<sup>18–20</sup> Previous studies have shown that the usage of a dopant at the bridge site of  $\text{RuO}_2$  makes the bridge oxygen more eager to interact with the proton of the intermediates.<sup>24,25</sup> We propose that  $\text{RuO}_2$  follows such a mechanism without the usage of a dopant, as illustrated in the schematic representation Fig. 1b. This  $\text{RuO}_2$  pathway is very reminiscent of the conventional pathway, differing only in steps one and three. Here, instead of the proton being directly bound to the adsorbed intermediate, the proton has migrated to the bridge site oxygen. This complex is stabilized by a weak interaction between this bridge-bound proton and the intermediate (indicated by a yellow dotted line). This hypothesis is supported by the experimental works of Rao *et al.*, who identified a  $-\text{OO}$  species at high potentials.<sup>17,26</sup> This  $-\text{OO}$  species is the experimental equivalent of the third intermediate of the  $\text{RuO}_2$  pathway depicted in Fig. 1b.

The different way that the first and third intermediates are adsorbed on the surface, has an effect on their binding

energy and thereby on the overall activity. The energy interdependency of the  $\text{HO}^*$  and  $\text{HOO}^*$  intermediates thus changes from  $\approx 3.2$  eV, as dictated by the universal scaling relations,<sup>4,5</sup> to  $\approx 2.7$  eV, a value that is closer to the ideal value of 2.46 eV. This relation is depicted in the activity volcano of Fig. 2 by the blue trend line. The activity volcano supports our hypothesis, as the DFT data point corresponding to  $\text{RuO}_2$  following pathway Fig. 1b (blue pentagon), holds a lower overpotential compared to the data point corresponding to the conventional pathway (blue circle). This places it right on top of the blue trend line. An observation that is strengthening our analysis, is that the blue trend line is followed by experimental data produced in the work of Suntivich *et al.*<sup>27</sup> In their experiments,  $\text{RuO}_2(110)$  surfaces were synthesized and their electrochemical response in different pH is recorded. Furthermore, they assign the first and second pre-oxidation peaks observed at the cyclic voltammeteries, as the  $\text{HO}^*$  and  $\text{O}^*$  intermediates respectively. The experimental  $\text{HO}^*$  energies serve the role of the descriptor for the experimental data at the activity volcano in the diagram (Fig. 2). The red triangles corresponding to  $\text{IrO}_2$ , reproduced from another work of Suntivich *et al.*,<sup>28</sup> tend as an ensemble to be placed closer towards the strong binding side of the conventional activity volcano. This is an indication that  $\text{IrO}_2$  follows the conventional reaction pathway. The  $\text{RuO}_2$  experimental data points are however spread. The two most active data points, corresponding to pH 13 (the left-most point) and pH 1 (the right-most point), are placed right on top of the blue trend-line together with the theoretical prediction for pathway Fig. 1b. In contrast, those corresponding to neutral, weakly acidic and weakly alkaline electrolytes are placed closer to the conventional activity volcano. This V-shaped activity trend has been attributed to a dependence of the  $\text{HOO}^*$  formation energy on pH in tandem with the fact that there is a clear pathway to  $\text{HOO}^*$  formation (from  $\text{O}^*$ ) in both high and low pH.<sup>27</sup> Furthermore, recent studies have suggested that the activity at high pH is mediated by cations.<sup>29</sup>

In Fig. 3 the scaling relations between  $\text{HO}^*$  and  $\text{O}^*$  intermediates on the cus site are depicted for both experimental and

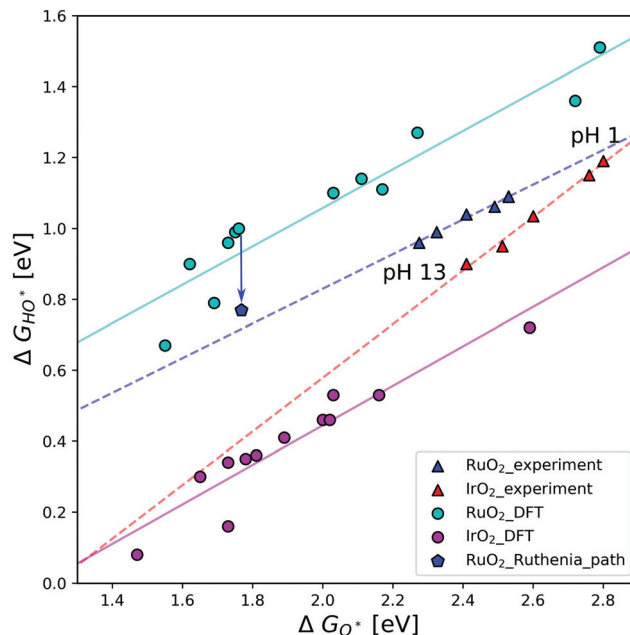




**Fig. 2** The OER activity volcano. The red data points correspond to IrO<sub>2</sub> and the blue ones to RuO<sub>2</sub>. The triangles represent the experimental data in varying pH conditions,<sup>27,28</sup> while the circle and the pentagon are theoretical data for the conventional and the RuO<sub>2</sub> pathway respectively. The blue trend line corresponds to the strong binding side of the volcano if the RuO<sub>2</sub> pathway is followed. The blue shaded area around the blue trend line, is the DFT error of  $\pm 0.2$  eV. The left y-axis (cyan) presents the theoretical overpotential and the right y-axis (magenta) presents the experimental overpotential, as the logarithm on the current is what you would expect from the Butler–Volmer equation to compare with potential. As we can only hope to compare trends, and not absolute numbers, the theoretical and experimental overpotentials are calibrated by overlapping the theoretical and experimental overpotential for IrO<sub>2</sub>(110). We cannot expect the scale of changes to be the same for calculations and experiments as the predicted differences are often larger than the measured. This figure is therefore two different figures overlapped for the comparison of trends. This also means that the two y-axes can be scaled and translated relative to each other. Two different descriptors with the same scale are used on the x-axis.  $G_O - G_{HO}$  and  $G_{HO}$  for the theoretical and the experimental data respectively.

theoretical results. The calculated data subtracted from the work of Federico Calle-Vallejo *et al.*,<sup>19</sup> were produced by DFT calculations on (110) surfaces of IrO<sub>2</sub> and RuO<sub>2</sub>, using different implementations of DFT. In this figure, the experimental points relating to the aforementioned V-shaped activity pattern constitute a clear trend-line. This hints at them following the same pathway, even though their activity dips at neutral pH, and suggests that the activity dip might be mechanistic in nature.

It is the relative strong binding of the oxygen intermediate, which makes the calculated activity of RuO<sub>2</sub> smaller than IrO<sub>2</sub>. Whereas the binding of HO\* and HOO\* on the RuO<sub>2</sub> cus site is similar, the O\* binding is much stronger than that on IrO<sub>2</sub>. This could be an artefact of the DFT calculations, however, it is seen to hold across DFT implementations. The binding energies vary between the different methods, but the difference between HO\* and O\* binding is close to constant. Previous experimental studies also show that RuO<sub>2</sub> binds oxygen stronger than IrO<sub>2</sub> for the same HO\* binding,<sup>27</sup> even if the difference is smaller than that found in the DFT data. In contrast to the DFT



**Fig. 3** Scaling relation of HO\* binding energies against O\* on the cus site for IrO<sub>2</sub> and RuO<sub>2</sub>(110) surfaces. The circles correspond to DFT data following different DFT implementations as reproduced from Federico Calle-Vallejo *et al.*<sup>19</sup> In contrast, the triangles correspond to experimental data in varying electrolyte conditions.<sup>27,28</sup> Each of these data sets have their own corresponding trend-line. The blue pentagon represents RuO<sub>2</sub> DFT data following pathway Fig. 1b, and the blue arrow represents the difference in  $\Delta G_{HO^*}$  as a result of following this pathway. This indicated difference is the same size as the difference between the experimental and DFT trend-lines for RuO<sub>2</sub>.

data, differences in experimental data is due to varying electrolyte pH. As previously mentioned, the experiments measure the potentials for the first and second oxidation peaks, those are normally assumed to be related to the HO\* and O\* intermediates on the cus site. However, it could just as well be related to the reaction path Fig. 1b. The stronger the relative O\* binding on the cus site, the more likely reaction (1) becomes, as the then relatively unstable HO\* intermediate is avoided. Doing a DFT calculation following the RuO<sub>2</sub> pathway creates the point represented by a blue pentagon. This point is placed towards a stronger binding of HO\*, and thus comes very close to the experimental trend-line. As points corresponding to different DFT implementations could lie anywhere on the teal trend-line, it is fair to assume that the calculated pentagon could lie anywhere on the experimental trend-line, as dependent on the specific DFT implementation. It therefore seems that pathway Fig. 1b accurately describes the experimental trend.

## Computational methods

Density functional theory calculations were done using Grid-based Projector Augment Wave (GPAW),<sup>30,31</sup> assisted by the Atomic Simulation Environment (ASE)<sup>32</sup> interface. Using the generalized gradient approximation (GGA), the BEEF-vdW functional<sup>33</sup> expressed exchange and correlation. This functional was chosen



specifically, as it gives a better description of the long range interactions, compared to the initially chosen RPBE.<sup>34</sup> RPBE provides trustworthy results for strong interacting chemical complexes (chemi-adsorption of a molecule on a surface), but it is not as efficient as BEEF-vdW in describing long range interactions. In this particular case the usage of RPBE, placed the blue trend-line of Fig. 2 much closer to the strong side of the conventional volcano. This result differentiates the interpretation of the phenomena and thus the conclusions of the study. The IrO<sub>2</sub> and RuO<sub>2</sub>(110) surfaces consist of four atomic layers where the two bottom layers were fixed in their initial position, as to mimic the bulk of the corresponding material. The top layers were free to converge to their minimum electronic energy positions. In the *x* and *y* directions the structures were replicated by 1 and 3 times respectively. The sampling of the Brillouin zone was done with a *k*-point mesh of (3, 2, 1) and the calculations were conducted with a grid spacing of 0.18 Å. Above and below the structures a vacuum of 15 Å was introduced to avoid unintended interactions between the slab and itself. The structures were relaxed until the total forces in the system were below 0.05 eV Å<sup>-1</sup>.

## Conclusion

In this work we are studying the discrepancy between DFT and experimental results, regarding the oxygen evolving reactivity of RuO<sub>2</sub>. We propose that the reaction pathway for electrochemical water oxidation on RuO<sub>2</sub>(110) surfaces, at least in acidic conditions, is slightly different from the reaction path on IrO<sub>2</sub>. In particular, the differences are located at the first and third intermediates, where the protons of HO\* and HOO\* are migrating towards the bridge oxygen surface. The energy interdependency of HO\* and HOO\* is 2.7 eV for the RuO<sub>2</sub> pathway, and is much closer to the ideal difference of 2.46 eV. As a consequence the DFT activity is much higher than the one produced by the conventional mechanism and thus the structure is placed closer to the apex of the activity volcano. Furthermore the new placement of RuO<sub>2</sub>(110) on the activity volcano, is at the same region of the RuO<sub>2</sub> experimental results for highly acidic and highly alkaline electrolytes. By using the conventional pathway we have a very weak interaction of HO\* with the surface's cus site. On the other hand, using the RuO<sub>2</sub> pathway widens the energy difference between HO\* and O\*, placing this DFT calculation closer to experimental trend-lines. This theoretical-experimental agreement, indicates that the RuO<sub>2</sub> mechanism is generally followed. It is to be expected that this trend extends to other active RuO<sub>2</sub> facets, as it is the strong binding of the O\* intermediate that makes the conventional path too difficult. Given that other facets share this strong oxygen binding, the alternate mechanism would likely apply.

## Conflicts of interest

There are no conflicts to declare.

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## References

- H. B. Gray, *Nat. Chem.*, 2009, **1**, 7.
- Z. Yan, J. L. Hitt, J. A. Turner and T. E. Mallouk, *Proc. Natl. Acad. Sci. U. S. A.*, 2020, **117**, 12558–12563.
- J. Durst, C. Simon, F. Hasché and H. A. Gasteiger, *J. Electrochem. Soc.*, 2014, **162**, F190–F203.
- I. C. Man, H.-Y. Su, F. Calle-Vallejo, H. A. Hansen, J. I. Martínez, N. G. Inoglu, J. Kitchin, T. F. Jaramillo, J. K. Nørskov and J. Rossmeisl, *ChemCatChem*, 2011, **3**, 1159–1165.
- S. Divanis, T. Kutlusoy, I. M. Ingmer Boye, I. C. Man and J. Rossmeisl, *Chem. Sci.*, 2020, **11**, 2943–2950.
- Q. Feng, X. Yuan, G. Liu, B. Wei, Z. Zhang, H. Li and H. Wang, *J. Power Sources*, 2017, **366**, 33–55.
- G. Matute, J. Yusta and L. Correas, *Int. J. Hydrogen Energy*, 2019, **44**, 17431–17442.
- S. Trasatti, *J. Electroanal. Chem. Interfacial Electrochem.*, 1980, **111**, 125–131.
- Y. Lee, J. Suntivich, K. J. May, E. E. Perry and Y. Shao-Horn, *J. Phys. Chem. Lett.*, 2012, **3**, 399–404.
- R. Frydendal, E. A. Paoli, B. P. Knudsen, B. Wickman, P. Malacrida, I. E. L. Stephens and I. Chorkendorff, *Chem-ElectroChem*, 2014, **1**, 2075–2081.
- O. Kasian, S. Geiger, T. Li, J.-P. Grote, K. Schweinar, S. Zhang, C. Scheu, D. Raabe, S. Cherevko, B. Gault and K. J. J. Mayrhofer, *Energy Environ. Sci.*, 2019, **12**, 3548–3555.
- M. Escudero-Escribano, A. F. Pedersen, E. A. Paoli, R. Frydendal, D. Friebel, P. Malacrida, J. Rossmeisl, I. E. L. Stephens and I. Chorkendorff, *J. Phys. Chem. B*, 2018, **122**, 947–955.
- M. Bernt, A. Siebel and H. A. Gasteiger, *J. Electrochem. Soc.*, 2018, **165**, F305–F314.
- D. Hageluen and V. den Broeck, *Precious Materials Handbook*, 2012, pp. 10–35.
- P. C. K. Vesborg and T. F. Jaramillo, *RSC Adv.*, 2012, **2**, 7933–7947.
- A. Grimaud, A. Demortière, M. Saubanère, W. Dachraoui, M. Duchamp, M.-L. Doublet and J.-M. Tarascon, *Nat. Energy*, 2016, **2**, 16189.
- R. Rao, M. Kolb and L. E. A. Giordano, *Nat Catal.*, 2020, **3**, 516–525.
- J. Rossmeisl, Z.-W. Qu, H. Zhu, G.-J. Kroes and J. Nørskov, *J. Electroanal. Chem.*, 2007, **607**, 83–89.
- L. G. V. Briquet, M. Sarwar, J. Mugo, G. Jones and F. Calle-Vallejo, *ChemCatChem*, 2017, **9**, 1261–1268.
- Y. Lee, J. Suntivich, K. J. May, E. E. Perry and Y. Shao-Horn, *J. Phys. Chem. Lett.*, 2012, **3**, 399–404.



- 21 J. K. Nørskov, J. Rossmeisl, A. Logadottir, L. Lindqvist, J. R. Kitchin, T. Bligaard and H. Jónsson, *J. Phys. Chem. B*, 2004, **108**, 17886–17892.
- 22 J. Rossmeisl, A. Logadottir and J. Nørskov, *Chem. Phys.*, 2005, **319**, 178–184.
- 23 K. A. Stoerzinger, O. Diaz-Morales, M. Kolb, R. R. Rao, R. Frydendal, L. Qiao, X. R. Wang, N. B. Halck, J. Rossmeisl, H. A. Hansen, T. Vegge, I. E. L. Stephens, M. T. M. Koper and Y. Shao-Horn, *ACS Energy Lett.*, 2017, **2**, 876–881.
- 24 N. B. Halck, V. Petrykin, P. Krtil and J. Rossmeisl, *Phys. Chem. Chem. Phys.*, 2014, **16**, 13682–13688.
- 25 M. Busch, N. Halck, U. Kramm, S. Siahrostami, P. Krtil and J. Rossmeisl, *Nano Energy*, 2016, **29**, 126–135.
- 26 R. R. Rao, M. J. Kolb, N. B. Halck, A. F. Pedersen, A. Mehta, H. You, K. A. Stoerzinger, Z. Feng, H. A. Hansen, H. Zhou, L. Giordano, J. Rossmeisl, T. Vegge, I. Chorkendorff, I. E. L. Stephens and Y. Shao-Horn, *Energy Environ. Sci.*, 2017, **10**, 2626–2637.
- 27 D.-Y. Kuo, H. Paik, J. Kloppenburg, B. Faeth, K. M. Shen, D. G. Schlom, G. Hautier and J. Suntivich, *J. Am. Chem. Soc.*, 2018, **140**, 17597–17605.
- 28 D.-Y. Kuo, J. K. Kawasaki, J. N. Nelson, J. Kloppenburg, G. Hautier, K. M. Shen, D. G. Schlom and J. Suntivich, *J. Am. Chem. Soc.*, 2017, **139**, 3473–3479.
- 29 R. R. Rao, B. Huang, Y. Katayama, J. Hwang, T. Kawaguchi, J. R. Lunger, J. Peng, Y. Zhang, A. Morinaga, H. Zhou, H. You and Y. Shao-Horn, *J. Phys. Chem. C*, 2021, **125**, 8195–8207.
- 30 J. Enkovaara, C. Rostgaard, J. J. Mortensen, J. Chen, M. Dułak, L. Ferrighi, J. Gavnholt, C. Glinsvad, V. Haikola and H. A. Hansen, *J. Phys.: Condens. Matter*, 2010, **22**, 253202.
- 31 J. J. Mortensen, L. B. Hansen and K. W. Jacobsen, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2005, **71**, 035109.
- 32 A. H. Larsen, J. J. Mortensen, J. Blomqvist, I. E. Castelli, R. Christensen, M. Dułak, J. Friis, M. N. Groves, B. Hammer and C. Hargus, *J. Phys.: Condens. Matter*, 2017, **29**, 273002.
- 33 J. Wellendorff, K. T. Lundgaard, A. Møgelhøj, V. Petzold, D. D. Landis, J. K. Nørskov, T. Bligaard and K. W. Jacobsen, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2012, **85**, 235149.
- 34 B. Hammer, L. B. Hansen and J. K. Nørskov, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1999, **59**, 7413–7421.

