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## Direct synthesis of metastable phases of 2D transition metal dichalcogenides

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The different polymorphic phases of transition metal dichalcogenides (TMDs) have attracted enormous interest in the last decade. The metastable metallic and small band gap phases of group VI TMDs displayed leading performance for electrocatalytic hydrogen evolution, high volumetric capacitance and some of them exhibit large gap quantum spin Hall (QSH) insulating behaviour. Metastable 1T(1T') phases require higher formation energy, as compared to the thermodynamically stable 2H phase, thus in standard chemical vapour deposition and vapour transport processes the materials normally grow in the 2H phases. Only destabilization of their 2H phase *via* external means, such as charge transfer or high electric field, allows the conversion of the crystal structure into the 1T(1T') phase. Bottom-up synthesis of materials in the 1T(1T') phases in measurable quantities would broaden their prospective applications and practical utilization. There is an emerging evidence that some of these 1T(1T') phases can be directly synthesized *via* bottom-up vapour- and liquid-phase methods. This review will provide an overview of the synthesis strategies which have been designed to achieve the crystal phase control in TMDs, and the chemical mechanisms that can drive the synthesis of metastable phases. We will provide a critical comparison between growth pathways in vapour- and liquid-phase synthesis techniques. Morphological and chemical characteristics of synthesized materials will be described along with their ability to act as electrocatalysts for the hydrogen evolution reaction from water. Phase stability and reversibility will be discussed and new potential applications will be introduced. This review aims at providing insights into the fundamental understanding of the favourable synthetic conditions for the stabilization of metastable TMD crystals and at stimulating future advancements in the field of large-scale synthesis of materials with crystal phase control.

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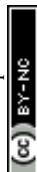
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## Introduction

The isolation of graphene has sparked widespread interest in layered materials. A large family of those materials is represented by the transition metal dichalcogenides which comprise about 40 compounds with a chemical formula of  $ME_2$  (M, transition metal; E, chalcogen).<sup>1</sup> While they have been studied for over half a century in their bulk form<sup>2,3</sup> for their photoelectrocatalytic activity,<sup>4–6</sup> lubricity<sup>7</sup> to superconductivity,<sup>8,9</sup> it was not until recently that they have been isolated in their elementary building blocks, which are triatomic layers, and several of their properties started to emerge and their potential identified. These emerging 2D materials offer new opportunities for the realization of novel technologies and for the exploration of novel condensed matter phenomena.

Interestingly, compounds of group IV, V, VI, VII and X are nearly isostructural with three possible crystallographic orders in the monolayer form,<sup>10</sup> and they present several analogies. Group VI TMDs have been studied more extensively than others and in particular the Se and S compounds ( $MoS_2$ ,  $MoSe_2$ ,  $WS_2$  and  $WSe_2$ ) owing to the richness and the uniqueness of the properties concerning each of their polymorphs. Additionally, the relative abundance of transition metal compounds, the good chemical stability in air compared to the dichalcogenides of other groups metals, make them more accessible. Group VI TMDs started to attract large interest a decade ago owing to the ability to display an indirect-to-direct bandgap crossover with bright photoluminescence (PL) in the monolayer form at room temperature<sup>11</sup> and for the valley degree of freedom associated to the non-centrosymmetric semiconducting phase.<sup>12</sup> Additional unique properties to this class of materials include mechanical flexibility, electrostatic coupling, large carrier mobility, gate-tunability, piezoelectricity in monolayers and odd number of layers, as discussed in recent reviews.<sup>13–16</sup>

It is now emerging that there is a variety of other properties associated to different polymorphs of group VI TMDs; some of those arising from the different crystallographic phases that they can acquire. In addition to the thermodynamically stable 2H polymorph with the trigonal-prismatic coordination of transition metal atoms, there are various polymorphs that are broadly defined as 1T phases (perfect 1T and distorted 1T', 1T'' and 1T''') and where the metal atoms present an octahedral (or distorted octahedral) coordination. These phases are generally characterized by a metal- (1T) or semimetal-like (1T') behaviour.

Indeed, the electronic band structure is heavily affected by the coordination of transition metal atoms which leads to a different splitting of the d-orbitals. The key opportunity arising from the TMD polymorphism is to influence the electronic structure of a given material and so to achieve a high density of electronic states at the Fermi level, which can enable efficient electrocatalytic hydrogen evolution from water. This ability was initially attributed to the localized metallic states on molybdenum-terminated edges of 2H  $MoS_2$  clusters.<sup>17</sup> Soon after, it became clear that the basal planes of the 1T phase  $MoS_2$  or  $WS_2$  could enable accessibility to a large number of electrocatalytically active sites thus approaching Pt performance.<sup>18</sup>

Several 1T' polymorphs have been demonstrated since. The 1T' phases of single layer  $WSe_2$  and  $WTe_2$  have been reported to be large gap quantum spin Hall (QSH) insulators suitable for application in spintronic devices and in the case of  $WSe_2$  also operable at ambient temperature.<sup>19–21</sup> Their chemical stability makes them more advantageous of the currently known large gap QSH insulators such as stanene<sup>22</sup> and two-dimensional In–Sb compounds which can be used in an inert atmosphere only.<sup>23</sup> On the other hand, it is predicted that a spontaneous symmetry breaking in the undistorted metallic 1T phase would open up a band gap and lead to the emergence of robust ferroelectricity in the novel 1T''' phase.<sup>24</sup>

Despite the rapid progress in investigating the properties of 1T(1T') phases of group VI TMDs over a small scale (about hundreds of nanometres over the basal plane), it is still quite arduous to obtain those phases as continuous over large areas. The 1T(1T') phases are metastable, thus their direct synthesis is very challenging because the formation energy of the 1T' phase is higher than that of the 2H phase;<sup>25</sup> therefore, the 2H  $ME_2$  is the preferable phase obtained.

The most common way, and historically the first one to be used, for achieving those 1T(1T') phases is *via* conversion of the 2H phase. This is generally achieved by increasing the electron density *via* exposure to alkali metals.<sup>26,27</sup> While the design of direct synthesis methods for the 1T(1T') phases is still at the early stages of investigation. Here, we provide a comprehensive overview of topochemical methods and direct synthesis methods designed so far to obtain the 1T(1T') phases of group VI TMDs. We will discuss the chemical prerequisites for the formation of the 1T(1T') phases in vapour phase and liquid phase. We will present the possible reaction pathways that can take place in liquid phase which lead to the nucleation and growth of crystalline metastable phases.

We will then review the intrinsic properties of 1T(1T') phases and we highlight the differences in the materials morphology and properties between 1T(1T') phases obtained *via* destabilization of 2H phase *versus* direct synthetic approaches. Finally, we will introduce newly emerging properties which are unique to the 1T(1T') phases as they arise from the lower symmetry. We will conclude the review with an outlook for the future of the different synthetic approaches to metastable phases in view of their applications and possible large-scale production of these materials.

### 1. Structural polymorphs of group VI TMDs

TMDs exhibit a variety of structural polymorphs. The polymorphs are determined by the very different chalcogen coordination geometry that transition metal atoms can accommodate, while the polytypes are determined by the stacking order of two or more individual triatomic layers with the same symmetry.

**1.1. 1H-type phases of group VI TMDs.** The first polymorph of group VI TMDs to be considered is characterised by a trigonal-prismatic coordination of transition metal atoms (point group  $D_{3h}$ ). In a single layer form, it is denoted as the 1H phase and its crystallographic projections along both the *b* (side view) and *c* (top view) axes are shown in Fig. 1. Depending on





**Fig. 1** Bulk polymorphs of group VI TMDs. In the schematics, the arrows indicate the pathways for possible structural transformation (i) from the 1H to the 1H' phase, (ii) from the 1H to the 1T phase, and (iii) from the 1T to the 1T' phase *via* either (i) metal plane gliding, (ii) chalcogen plane gliding, or (iii) spontaneous distortion in the metal sublattice. The crystallographic projections of the 1H, 1T, and 1T' monolayers along the *b* (side view) and *c* (top view) axes are shown. The known polytypes (stacking polymorphs) of the 1H, 1T, and 1T' structures are presented in the corresponding fields. Bottom left panel illustrates two additional clustering patterns – tetramerization (1T'') and trimerization (1T''') – proposed in the distorted 1T TMDs.

the stacking sequence of individual 1H triatomic layers, the 2H<sub>a</sub>, 2H<sub>c</sub> and 3R polytypes are distinguished, where H or R stands for the hexagonal or rhombohedral lattice system and the integer denotes the number of triatomic layers in a unit cell. The *AbA BaB* and *AcA BcB* stacking sequences result in the 2H<sub>a</sub> and 2H<sub>c</sub> polytypes, respectively, while the 3R phase displays the *AbA BcB CaC* sequence of staggered 1H layers sharing the same crystallographic orientation.<sup>1,28</sup> These structural polytypes are presented in Fig. 1. Under ambient conditions, the thermodynamically stable polymorph of the majority of group VI TMDs is the semiconducting 2H<sub>c</sub> phase,<sup>25,29–31</sup> whereas the 2H<sub>a</sub> phase is common in group V TMDs, such as NbSe<sub>2</sub>, and is metastable for group VI TMDs.<sup>32,33</sup> It has been predicted that the 2H<sub>c</sub>-to-2H<sub>a</sub> phase transition can be induced in bulk MoS<sub>2</sub> under pressure (13 GPa); however, such a transition is not expected in MoSe<sub>2</sub> and MoTe<sub>2</sub> since in these two cases the energy barrier of the structural transformation increases with pressure.<sup>33</sup> The experiments have indeed demonstrated that the 2H<sub>c</sub>-to-2H<sub>a</sub> phase transition in MoS<sub>2</sub> occurs under applied pressure of 19–40

GPa and is accompanied by a collapse in the *c* parameter leading to an overlap of the valence and conduction bands and to a metallic character of the 2H<sub>a</sub> phase.<sup>34,35</sup> In the 2H structure, the second layer is rotated by 60° around the *c* axis with respect to the first layer, whereas in the 3R structure all layers have the same orientation but are shifted with respect to one another, and this might be the reason why the 2H-to-3R transitions were not reported so far. Similar to the 1H single layers of group VI TMDs demonstrating a robust spin-valley coupled physics,<sup>12,36–38</sup> the semiconducting 3R polymorph lacks the inversion symmetry and therefore provides an ideal platform for realisation of spin- and valleytronic devices as well as of ultrathin nonlinear optical devices.<sup>39</sup> Direct synthesis of the phase-pure 3R polymorph is challenging and has only been reported *via* chemical vapour transport (CVT) approach.<sup>39–41</sup>

Two types of structural transitions in the 1H-type TMD monolayers associated with the atomic planes gliding within one triatomic layer are assumed to give rise to two distinct polymorphs: 1H' and 1T'<sup>42</sup> (Fig. 1, paths i and ii). The transversal





displacement of the transition metal plane could produce the distorted 2H phase, which is referred to as the 2H' phase and is considered as a 60° rotational phase of the 2H structure<sup>42</sup> (Fig. 1, path i). Experimentally, the formation of a new distorted 2H phase, 2H<sub>d</sub>, was observed in the few-layered 2H MoTe<sub>2</sub> and W<sub>x</sub>Mo<sub>1-x</sub>Te<sub>2</sub> when an external electric field was applied.<sup>43</sup> The phase was found to be unstable and is considered to be a transitional structural state between the semiconducting 2H<sub>c</sub> and the semimetallic 1T' (or Td) phases discussed below. The structure of the mentioned 2H<sub>d</sub> state was not inferred at this stage.

**1.2. 1T-type phases of group VI TMDs.** On the other hand, intralayer gliding of one of the chalcogen planes, as shown in Fig. 1, path ii, produces the 1T structure which can be visualised as a 1H three-atom thick E–M–E slab distorted along the [210] crystallographic direction. As a result of this, the metal centres adopt a perfect octahedral coordination (point group *O<sub>h</sub>*) shown in Fig. 1. In contrast to the rich polytypism of the H phases, only one trigonal 1T polymorph (space group *P3̄m1*), with the metal centres of two adjacent layers located above each other, is reported in literature.<sup>44</sup>

The perfect octahedral 1T structure of group VI disulphides and diselenides is typically higher in energy by 700–900 meV per formula unit than the trigonal-prismatic 1H structure,<sup>45</sup> which renders this phase metastable.<sup>25,46</sup> It is prone to spontaneous structural distortions, caused by charge density waves (CDW) which lead to clustering of transition metal atoms and buckling of chalcogen planes to accommodate these metal atom displacements<sup>47,48</sup> (Fig. 1, path iii). Different distorted 1T structures of chemically exfoliated MoS<sub>2</sub> and WS<sub>2</sub> have been proposed in literature back in 1990s.<sup>47</sup> The most common superstructures include the zigzag chain ( $2a \times a$ ) formation, tetramerization ( $2a \times 2a$ ), and trimerization ( $a\sqrt{3} \times a\sqrt{3}$ ) of transition metal atoms in the *ab* plane as shown in Fig. 1. The former one is widely referred to as the 1T' polymorph and was established by scanning tunnelling microscopy (STM) studies of chemically exfoliated MoS<sub>2</sub> and WS<sub>2</sub> single layers<sup>47,49–51</sup> (Fig. 1). The three-dimensional arrangement of the 1T' single layers can lead to the formation of two polytypes: the orthorhombic Td phase and the monoclinic 1T' phase (Fig. 1). Additionally, tetramerization and trimerization within the transition metal sublattice would result in the 1T'' and 1T''' phases, respectively.<sup>52</sup> According to first-principles calculations, the semimetallic 1T' phase of group VI TMDs monolayers generally possesses the lowest formation energy among the competing distorted octahedral phases.<sup>45,53</sup> However, a spontaneous trimerization of transition metal atoms, involving only a slight displacement from the equilibrium positions, is found to be favourable in 1T MoS<sub>2</sub> monolayers, leading to the formation of the semiconducting 1T''' phase.<sup>45,54</sup> Such distinct distortion patterns in 1T group VI TMD monolayers are in good agreement with the early predictions that weak distortions (trimerization) are to be expected in the systems with short M–E bonds; whereas, strong distortions (zigzag chain formation) are characteristic to the systems with long M–E bonds.<sup>55</sup> It has also been predicted that the spontaneous symmetry breaking in

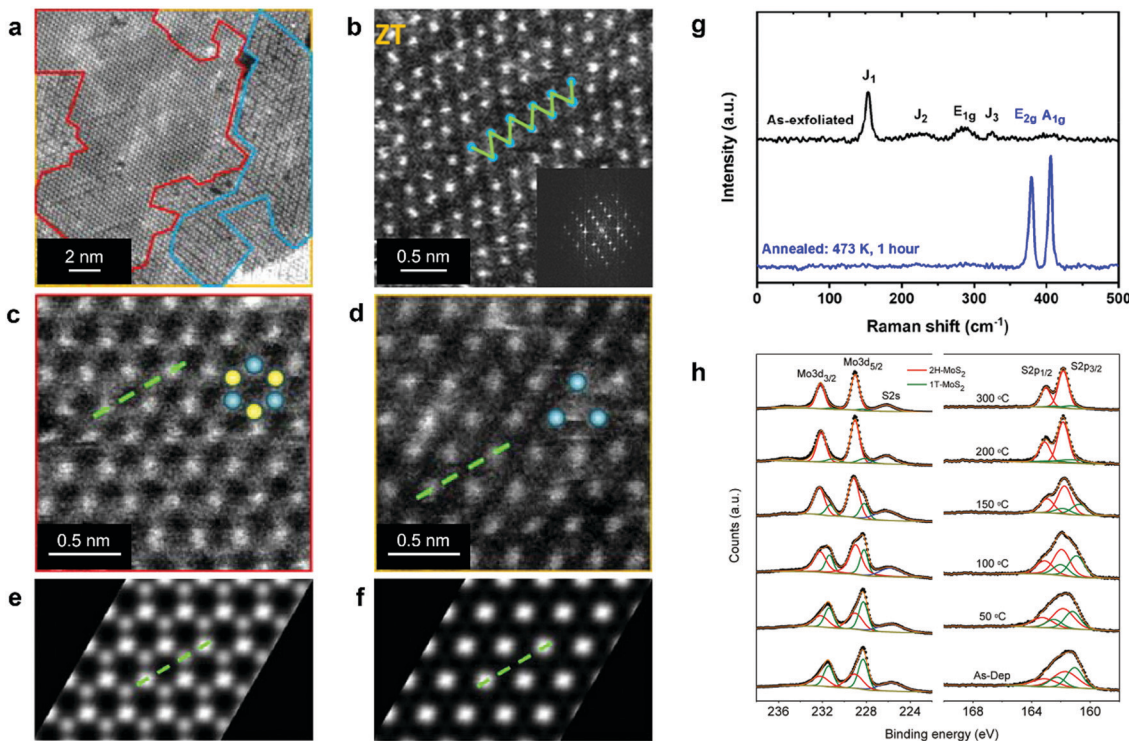
1T group VI TMD monolayers would lead to the emergence of robust ferroelectricity in the 1T''' phase, with the electrical dipoles perpendicular to the *ab* plane.<sup>24,54</sup> Recently, the preparation of 1T''' MoS<sub>2</sub> single crystals (space group *P3<sub>1</sub>m*) by deintercalation of KMoS<sub>2</sub> crystals has been reported.<sup>56</sup> Unlike the metallic 1T and semimetallic 1T' phases, bulk 1T''' MoS<sub>2</sub> was reported to be a semiconductor with an indirect band gap of 0.65 eV. Although the 1T''' phase was predicted to be formed only in MoS<sub>2</sub> monolayers, the small energy difference between the 1T' and 1T''' monolayers of WS<sub>2</sub> and MoSe<sub>2</sub> (36.5 and 57.8 meV per formula unit, respectively) suggests that these metastable phases might be obtained under favourable synthesis conditions.<sup>54</sup>

It is important to note that for group VI disulphides and diselenides, the 1T' polymorph is metastable under ambient conditions and can only be obtained under specific experimental conditions. On the contrary, bulk WTe<sub>2</sub> forms the orthorhombic Td polymorph (space group *Pnm2<sub>1</sub>*).<sup>57</sup> Bulk MoTe<sub>2</sub> is normally obtained in the 2H polymorph ( $\alpha$ -phase); however, the monoclinic 1T' MoTe<sub>2</sub> polymorph (space group *P2<sub>1</sub>/m*) exists at temperatures above 900 °C ( $\beta$ -phase).<sup>57,58</sup> Due to a very small energy difference between the 2H and 1T' polymorphs of MoTe<sub>2</sub> (40 meV per formula unit<sup>25</sup>), both polymorphs are reported almost equally probable depending on the synthesis conditions.<sup>59,60</sup> The orthorhombic Td polymorph ( $\gamma$ -phase) of MoTe<sub>2</sub> can be obtained by cooling the monoclinic 1T' ( $\beta$ -) phase down to cryogenic temperatures<sup>61–64</sup> or it can be stabilized at room temperature in the presence of substitutional doping with tungsten.<sup>65</sup> On the other hand, the Td polymorph of WTe<sub>2</sub> can be transformed into the lower symmetry 1T' phase at pressure over 15.5 GPa.<sup>66,67</sup> Due to a large energy difference between the 2H and 1T' phases, the metastable group VI diselenides and disulphides are challenging to obtain in a pure form and the reports on their successful synthesis are limited.<sup>21,68–71</sup> The crystal structure of the monoclinic 1T' phase of group VI diselenides and disulphides is still debated with the possible space groups being either *P2<sub>1</sub>/m* or *C2/m*; the latter has only been reported recently and is referred to as the 2M phase, where M stands for the monoclinic lattice system.<sup>72</sup> Unlike the semiconducting 2H counterparts, the corresponding 1T' and Td phases of group VI TMDs are semimetallic.<sup>61</sup> Furthermore, the non-centrosymmetric Td polymorph exhibits topologically nontrivial surface electronic states,<sup>73,74</sup> and bulk Td phases of WTe<sub>2</sub> and MoTe<sub>2</sub> have been extensively investigated as potential candidates for the observation of 3D Weyl fermions.<sup>64,75–77</sup> Recently, single layers of group VI TMDs with the 1T' symmetry have been predicted to host topologically protected edge states,<sup>46,78</sup> and large QSH gaps have been verified experimentally in monolayer 1T' WTe<sub>2</sub>,<sup>19</sup> 1T' WSe<sub>2</sub>,<sup>20,21</sup> and 1T' MoS<sub>2</sub>.<sup>79</sup>

**1.3. Distinguishing the metastable group VI TMD polymorphs.** Reliably distinguishing the metastable polymorphs of group VI TMDs is of paramount importance for the experimental investigation of electronic properties of the respective phases. In recent literature, the difference between 1T and 1T' structural polymorphs is often neglected and the 1T' phase is frequently reported as 1T<sup>80–83</sup> or is more generally referred to







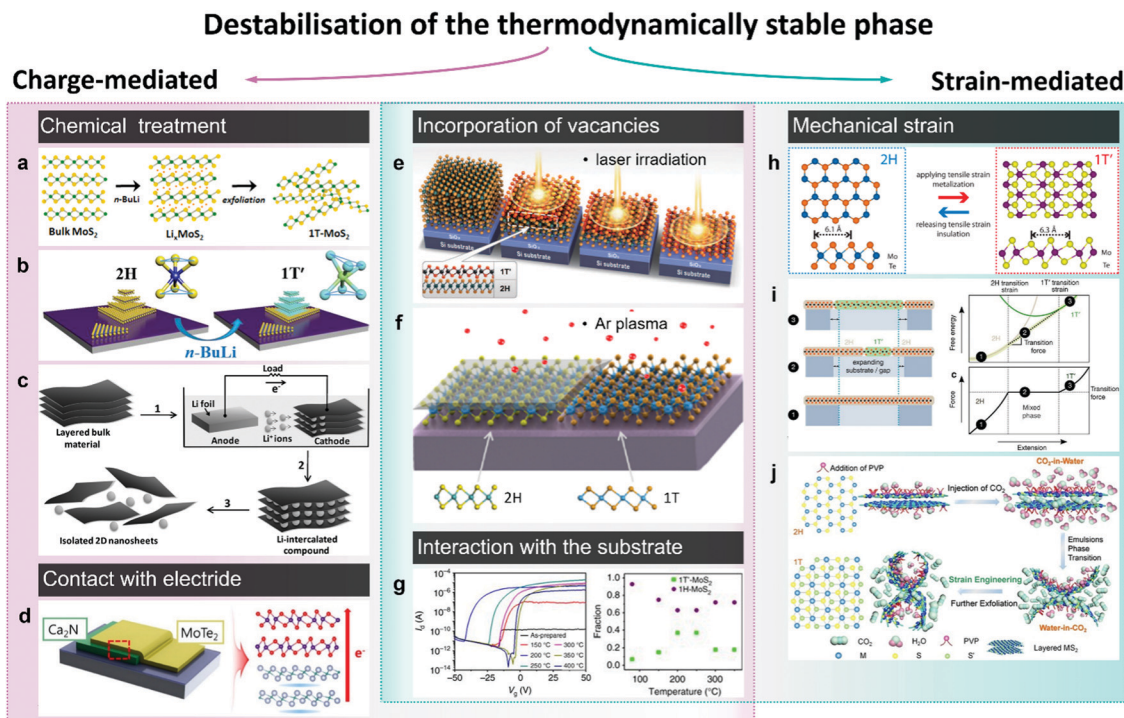
**Fig. 2** Distinguishing the 1H, 1T, and 1T' polymorphs. (a) STEM image of chemically exfoliated MoS<sub>2</sub> displaying co-existing 1H, 1T, and 1T' phases; the corresponding areas are enclosed by red, yellow, and blue curves, respectively. (b–d) Representative HAADF STEM images of the 1T', 1H, and 1T MoS<sub>2</sub> lattice, respectively. (e and f) Simulated HAADF STEM images of the ideal 1H and 1T MoS<sub>2</sub> monolayers, respectively. (g) Raman spectra of chemically exfoliated 1T' (top) and annealed 2H (bottom) MoS<sub>2</sub> flakes, illustrating the characteristic vibrational modes of these two phases. (h) XPS spectra of the Mo 3d, S 2s and S 2p core level electrons of mixed-phase 1T/2H MoS<sub>2</sub> nanosheets, demonstrating the gradual recovery of the 2H phase during thermal annealing. The 1T and 2H phase contributions are given in green and red, respectively. Separate panels are reproduced with permission: (a and c–f) from ref. 85, copyright 2012, American Chemical Society; (b) from ref. 51, copyright 2016, American Chemical Society; (g) from ref. 263, copyright 2019, American Chemical Society; and (h) from ref. 264, copyright 2011, American Chemical Society.

as the metallic phase.<sup>84</sup> The 1T and 1T' domains can be unambiguously distinguished in high resolution transmission electron microscopy (HR TEM) and scanning transmission electron microscopy (HR STEM) images of atomically-thin nanosheets where the 1T' phase exhibits characteristic zigzag chains of transition metal atoms<sup>85,86</sup> (Fig. 2a and b). However, the 1T and 2H phases, displaying a three-fold symmetry in the *ab* plane, are virtually undistinguishable in HR TEM images. It has recently been established that the 1T and 2H domains in group VI TMD monolayers can be distinguished in high-angle annular dark-field (HAADF STEM) images by the intensity of chalcogen columns as shown in Fig. 2c and d.<sup>85</sup> The discernible variations in signal intensity arise since in the trigonal-prismatic 2H structure, two chalcogen atoms overlapping in the direction of the incident beam make the contrast almost comparable to that of a metal site producing a honeycomb lattice in STEM images (Fig. 2e), whereas in the octahedral 1T structure, a strong contrast between transition metal and chalcogen sites results in a hexagonal lattice (Fig. 2f). Moreover, the stacking sequence of 1H monolayers in bi- and trilayer nanosheets can be deduced. However, extreme caution must be taken when interpreting the microscopy data and especially Fourier filtered images since residual aberrations may give rise to quasi atomic columns.<sup>28</sup> Additionally, it is suggested that the

2H and 3R polytypes are possible to distinguish by the intensity profiles of selected area electron diffraction (SAED) patterns taken along the [001] zone axis.<sup>41</sup> The {100} and {110} diffraction spots display nearly equal intensities in the 2H phase, while in the 3R phase, their intensities differ by almost three times.<sup>41</sup>

Spectroscopic techniques, such as Raman spectroscopy and X-ray photoemission spectroscopy (XPS), offer a relatively accurate way of detecting co-existing polymorphs in TMD nanostructures over micron-sized areas. In the original work by J. Sandoval *et al.*, it has been reported that the octahedral 1T phase of chemically exfoliated MoS<sub>2</sub> displayed two Raman-active modes, E<sub>g</sub> and A<sub>1g</sub>, at 287 and 408 cm<sup>-1</sup>, respectively, while the E<sub>2g</sub><sup>1</sup> mode of the 2H phase (at 383 cm<sup>-1</sup>) was absent.<sup>87</sup> In addition to those, three previously unassigned modes, J<sub>1</sub>, J<sub>2</sub>, and J<sub>3</sub>, at 156, 226, and 333 cm<sup>-1</sup>, respectively, were found (Fig. 2g). These modes were attributed to the 2a × a superstructure in the octahedral phase. The wider set of Raman-active modes of the 1T' and Td phases is indicative of their lower symmetry as compared to the 2H phase and is usually reported for WTe<sub>2</sub> and MoTe<sub>2</sub>.<sup>62,88–90</sup> The 1T' signature J<sub>1</sub>, J<sub>2</sub>, and J<sub>3</sub>, modes were also observed in the 1T' MoS<sub>2</sub>, WS<sub>2</sub>, MoSe<sub>2</sub>, and WSe<sub>2</sub> single crystals<sup>68,91</sup> and few-layered nanosheets.<sup>71,86,92,93</sup> It should be noted that the J<sub>1</sub> mode consistently exhibits the highest intensity in Raman spectra of nearly phase-pure 1T' group





**Fig. 3** Accessing the metastable 1T(1T') phase via destabilisation of the thermodynamically stable (2H) phase. The approaches to induce the 2H-to-1T' transition in group VI TMDs are divided in two categories: charge-mediated, as chemical treatment (panels a–c) and contact with an electrode (panel d); and strain-mediated, such as tensile strain-induced transition between the 2H and 1T'(1T) polymorphs (panels h–j). Panels e–g represent the phase transitions caused by the combination of charge and strain as in the case of chalcogen vacancies (panels e and f) or due to the interaction with the substrate (panel g). Separate panels are adapted with permission: (a) from ref. 100, copyright 2017, American Chemical Society; (b) from ref. 114, copyright 2018, American Chemical Society; (c) from ref. 108, copyright 2011, John Wiley & Sons; (d) from ref. 123, copyright 2017, American Chemical Society; (e) from ref. 155, copyright 2015, American Association for the Advancement of Science; (f) from ref. 138, copyright 2017, American Chemical Society; (g) from ref. 128, copyright 2017, Springer Nature; (h) from ref. 130, copyright 2016, American Chemical Society; (i) from ref. 25, copyright 2014, Springer Nature; and (j) from ref. 134, copyright 2017, John Wiley & Sons.

VI disulphides.<sup>86,92,93</sup> Further, the overall Raman signal intensity is significantly lower in the 1T(1T') phase TMDs than in the 2H samples of equivalent thickness as has been experimentally demonstrated in the case of MoS<sub>2</sub><sup>94–96</sup> and WSe<sub>2</sub> nanosheets.<sup>71,97</sup>

In XPS spectra, both the 1T and 1T' phases are reported to present peaks systematically shifted to lower binding energies as compared to the 2H phase, however, XPS is not sensitive to the differences between the 1T and 1T' phases<sup>98</sup> (Fig. 2h). In chemically exfoliated TMDs, the shift to lower binding energies is typically ascribed to partial reduction of transition metal centres during alkali metal intercalation.<sup>91</sup>

## 2. Access to the metastable crystal phases of group VI TMDs

The metastable crystal phases can be either accessed through destabilisation of the thermodynamically stable polymorphs, for instance, due to electron transfer, or synthesised directly, when the growth conditions are optimised in such a way that the metastable product is kinetically preferred.

### 2.1. 2H-to-1T(1T') phase transformation

**2.1.1. Chemical treatment.** The first reports on the metallic 1T(1T') group VI TMD polymorphs date back to the pioneering works by R. Frindt *et al.* on chemical exfoliation of these materials<sup>27,50</sup> (Fig. 3a). Chemical exfoliation of layered materials typically proceeds in two steps.<sup>47,50,99</sup> At first, bulk TMD powders

are intercalated with alkali metals. The most commonly used intercalants include *n*-BuLi,<sup>83,100,101</sup> *t*-BuLi,<sup>102</sup> and LiBH<sub>4</sub>.<sup>50,103</sup> The subsequent hydrolysis of the intercalated alkali metal cations leads to the generation of hydrogen gas bubbles that split the individual layers apart. The exfoliation can be facilitated by gentle shaking or sonication. It has been demonstrated that upon intercalation of alkali metal species between the layers of bulk crystals of group VI TMDs, the crystal phase would change from the 2H to 1T. It has been suggested that the mechanism of this structural transition should be imputed to the electron transfer from the intercalated alkali metal species leading to an increased electron count on the transition metal d-orbitals.<sup>47</sup> Specifically, for the case of MoS<sub>2</sub> monolayers, it has been shown that the injection of 4 electrons per unit cell allows lowering the 2H-to-1T transition barrier from 1.59 to only 0.27 eV.<sup>48</sup> Thus, upon this electron transfer (chemical reduction), the octahedral coordination of transition metal becomes favoured over the trigonal-prismatic and thus the 2H-to-1T structural phase transition occurs.<sup>104</sup> In the absence of intercalated species, the 2H-to-1T phase transition due to accumulation of negative charges was also observed in Re doped MoS<sub>2</sub> monolayers under continuous electron beam scanning during STEM imaging.<sup>42</sup> The transition occurs through the initial formation of numerous non-parallel molybdenum zigzag chains which is followed by chalcogen plane





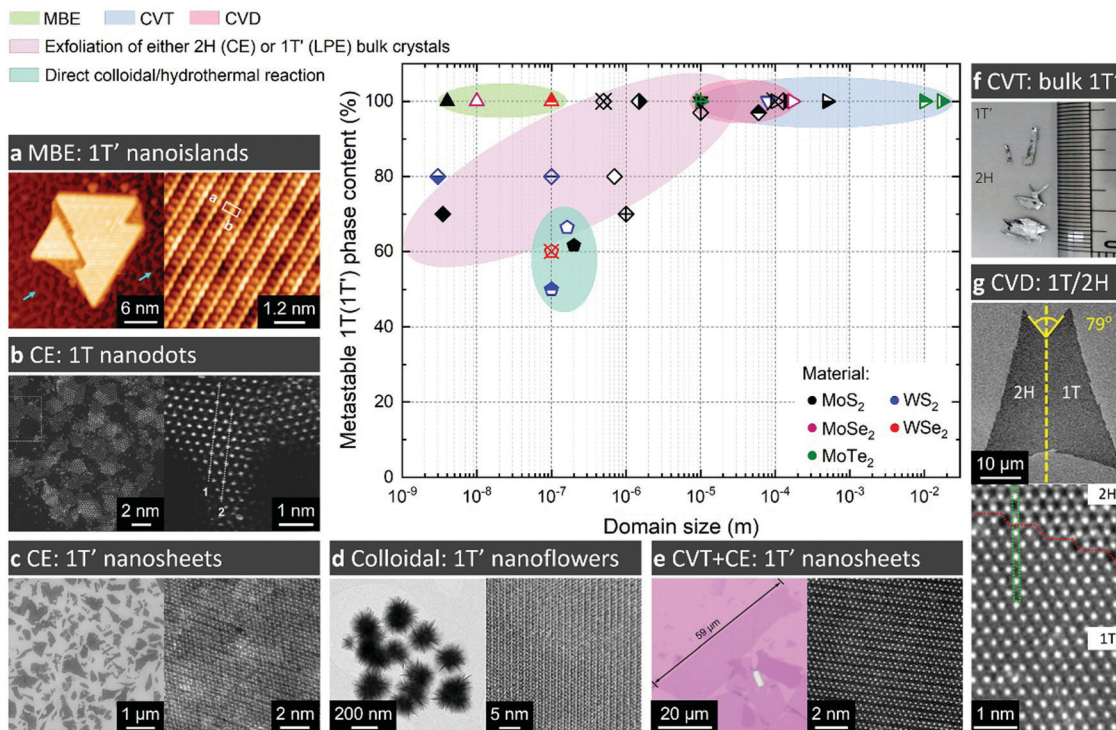


Fig. 4 Comparison of various experimental approaches to produce the metastable 1T(1T') phases of group VI TMDs. The approaches are grouped into the following categories: molecular beam epitaxy (MBE); chemical vapour transport (CVT) growth; chemical vapour deposition (CVD); chemical (CE) or liquid-phase (LPE) exfoliation of bulk crystals; and direct synthesis approaches, such as colloidal and hydrothermal growth. The experimental data are taken from the following articles: ref. 21, 56, 58, 68–71, 79–81, 83, 93, 101, 103, 112, 144, 147, 148, 153, 162 and 164. Separate panels are adapted with permission: (a) from ref. 70, copyright 2019, American Chemical Society; (b) from ref. 112, copyright 2018, John Wiley & Sons; (c) from ref. 103, copyright 2013, American Chemical Society; (d) from ref. 71, copyright 2019, Springer Nature; (e) from ref. 69, copyright 2019, John Wiley & Sons; (f) from ref. 153, copyright 2015, Springer Nature; and (g) from ref. 164, copyright 2018, American Chemical Society.

gliding, reducing mechanical strain and forming a domain of the trigonal 1T phase.<sup>42</sup>

Generally, the 1T(1T') phases produced *via* chemical exfoliation approach are attained as small crystalline (within the range 1–10 nm, Fig. 4) domains often embedded into the 2H matrix<sup>85</sup> (Fig. 2a and 4c). Patches of both 1T and 1T' phases are found in chemically exfoliated MoS<sub>2</sub> nanosheets.<sup>49,85,101</sup> Whereas a characteristic distortion in the *ab* plane of chemically exfoliated WS<sub>2</sub> indicates the crystal phase to be 1T',<sup>47,51</sup> although the resulting crystal phase is commonly denoted as 1T.<sup>83</sup> The ratio between the converted 1T(1T') and pristine 2H phases strongly depends on the type and concentration of reducing agent (alkali metal) and the TMD material.<sup>102</sup> The maximum amount of 1T(1T') phases achievable *via* chemical exfoliation approach is reported to be ~70–80% of the overall material<sup>92,101</sup> (typically, estimated from by XPS data, Fig. 4). The metastable 1T(1T') phases obtained by chemical exfoliation are reported to gradually decay into the 2H polymorph over time; this process is considerably promoted in the restacked films once the stabilising ions have been removed during drying.<sup>49,87,105</sup> A full conversion to the 2H phase occurs if the material is subjected to high temperatures, such as ~100 °C for MoS<sub>2</sub><sup>27</sup> and ~200 °C for WS<sub>2</sub>.<sup>50</sup>

Recently, it has been demonstrated that the lifetime of the 1T(1T') phases in intercalated few-layered MoS<sub>2</sub> flakes can be

considerably extended *via* the hydrogenation of intercalated lithium species.<sup>86</sup> Results of DFT calculations suggest that LiH is a good electron donor that stabilises the 1T' phase in multi-layered flakes for over 3 months. Although LiH, encapsulated between the TMD layers, is stable, the surface bound lithium hydride can easily convert into Li<sub>2</sub>CO<sub>3</sub> and Li<sub>3</sub>N, what makes this approach unsuitable for stabilisation of the 1T' phase in chemically exfoliated single layers.

A morphological aspect should be also taken into consideration as the non-uniform distribution of intercalated lithium hinders the control over the dimensions and distribution of the converted 1T(1T') patches. The intercalation kinetics affects the conversion efficiency and the lateral dimensions of the patches of metastable phase.<sup>106,107</sup> A uniform lithiation has been reported by an electrochemical intercalation process where Li foil serves as an anode and bulk TMD material as a cathode<sup>108</sup> (Fig. 3c). After lithium insertion, lithiated powders are exfoliated by mild sonication in water producing predominantly monolayered nanosheets with lateral dimensions of up to 1 μm. Moreover, it has been demonstrated that the lithium migration within the van der Waals (vdW) gaps of lithiated few-layered material can be controlled by an external electric field so that a reversible 2H-to-1T' phase transformation can be uniformly extended over micron-sized areas.<sup>109</sup> Such controllable phase switching has been demonstrated in lithiated MoS<sub>2</sub> thin films,





which have been used as neuromorphic memory devices,<sup>109</sup> and in ultra-thin ( $\sim 2\text{--}5$  nm) MoS<sub>2</sub> nanosheets upon a reversible electrochemical lithiation.<sup>110</sup> Another method to achieve a uniform intercalation in group VI TMDs has been recently reported by J. Zheng *et al.*<sup>111</sup> Performed prior to intercalation, the expansion of TMD powders by reacting with hydrazine under hydrothermal conditions has facilitated the alkali metal naphthalenide insertion. This allowed producing single layer (80%) MoS<sub>2</sub> sheets of up to 400  $\mu\text{m}^2$  in size; the metastable 1T(1T') phase content, however, was not specified.<sup>111</sup>

It is interesting to note that combined with ball-milling of 2H-phase bulk crystals, chemical (organolithium) treatment and exfoliation approach has recently been shown to produce MoS<sub>2</sub>, MoSe<sub>2</sub>, WS<sub>2</sub>, Mo<sub>0.5</sub>W<sub>0.5</sub>S<sub>2</sub> and MoSSe nanodots with lateral sizes of 2–5 nm and a high (70%) content of the metastable 1T phase<sup>112</sup> (Fig. 4b).

Phase conversion can also be achieved *via* a post-growth treatment of CVD grown 2H TMD few-layered flakes and films with solutions of *n*-BuLi in nonpolar solvents<sup>113,114</sup> (Fig. 3b). Considering that CVD grown TMD flakes present much larger (tens of microns) lateral sizes compared to the ones of chemically exfoliated flakes (submicron for monolayers), this could allow for the production of larger continuous domains of the metastable 1T(1T') phases. The *n*-BuLi treatment of polycrystalline films of 2H–MoS<sub>2</sub> and 2H–MoSe<sub>2</sub> grown on Si/SiO<sub>2</sub> wafers led to a partial conversion into the metallic 1T phase, with an overall conversion efficiency of  $\sim 50\%$ .<sup>113</sup> Further, the lithium-induced structural transformations in the CVD grown 2H–MoS<sub>2</sub> flakes appear to be thickness-dependent.<sup>114</sup> The increased energy barrier of the 2H-to-1T phase transition in MoS<sub>2</sub> monolayers implies much longer treatment time in order to achieve the phase conversion compared to the multi-layered counterparts.<sup>114</sup> This, therefore, paves the way to the 2D metal–semiconductor junctions in CVD grown TMD devices by controlling the exposure to lithiating agent (*n*-BuLi). In the case of 2H-WSe<sub>2</sub> monolayers, a complete conversion into the metallic 1T phase has been reported and the deposition of a PMMA mask has allowed for the patterning of 2H/1T heterojunctions.<sup>97</sup> While a complete conversion to the 1T(1T') phases *via* lithium intercalation of bulk MoSe<sub>2</sub> and WSe<sub>2</sub> powders have not been achieved.<sup>102,115</sup> Finally, the 2H-to-1T phase conversion has been demonstrated to be induced also by exposure to the vapours of electron donating compounds, such as butylamine and trimethylamine, as was shown in the cases of MoS<sub>2</sub> and MoSe<sub>2</sub>.<sup>116</sup> It has been proposed that this approach can be potentially employed to fabricate new passive sensing devices based on detecting the abrupt changes in material conductivity in the presence of strong donor analytes, such as nerve gas.

**2.1.2. Electrostatic gating.** The dynamical control of the structural phase transitions in TMD monolayers by an applied electric field can be potentially interesting for application in light-weight flexible electronic devices, including non-volatile memory devices.<sup>117</sup> Bias-voltage driven T-to-H crystal phase transformations have been previously demonstrated in TaS<sub>2</sub> crystals.<sup>118</sup> Y. Li *et al.* have proposed that the semiconductor-to-

semimetal phase transition can be achieved in some group VI TMDs using an electrostatic gating.<sup>119</sup> It has been found that the transition between the 2H and 1T' phases can be driven by excess electric charge in the monolayer. Thus, the excess charge densities less than  $-0.04e$  or greater than  $0.09e$  per formula unit would suffice to induce the 2H-to-1T' transformation in an undoped MoTe<sub>2</sub> monolayer; however, to achieve this transition in lighter TMD monolayers, such as MoS<sub>2</sub> and WS<sub>2</sub>, much greater excess charge densities are needed. Alternatively, a constant voltage of  $-1.6$  or  $4.4$  V is required to induce the semiconductor-to-semimetal transition in monolayer MoTe<sub>2</sub> on HfO<sub>2</sub> in the electrostatic gating device configuration.<sup>119</sup> A reversible 2H-to-1T' phase transition was achieved in MoTe<sub>2</sub> by ionic liquid gating in a field-effect transistor configuration.<sup>120</sup> The required gate voltages could be substantially reduced by varying the chemical composition, such as by alloying<sup>119,121</sup> or forming Janus monolayers.<sup>122</sup>

**2.1.3. Electron doping by contacting with an electron reservoir.** Long-range crystal lattice modulation can be achieved by contacting the 2H TMD (MoTe<sub>2</sub>) with a single crystalline 2D electride as has been demonstrated in ref. 123 (Fig. 3d). Electrides are a class of inorganic compounds in which electrons serve as anions.<sup>124</sup> Layered electrides containing two-dimensionally confined anionic electrons are generally characterised by exceptionally high charge carrier concentrations ( $1.4 \times 10^{22} \text{ cm}^{-3}$ ) and high electron mobilities at room temperature ( $160 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ ).<sup>125</sup> Similarly to van der Waals solids, layered electrides can be exfoliated into 2D sheets.<sup>126</sup> Unlike the chemical treatment discussed earlier, electron transfer from the electride allows for achieving higher levels of doping (about  $10^{14} \text{ cm}^{-2}$ ). Such high doping densities can also be achieved by ionic modulation, however, contacting with electride is effective over larger distances from the contact interface. For instance, forming a vertical heterostructure between  $[\text{Ca}_2\text{N}]^+ \cdot e^-$  and few-layered MoTe<sub>2</sub> has led to a uniform 2H-to-1T' phase transformation over a micron sized area even as far away as 100 nm from the electride.<sup>123</sup>

A reversible 2H-to-1T phase transition has been demonstrated in MoS<sub>2</sub> monolayers by an injection of hot electrons generated by plasmon excitation in Au nanospheres.<sup>127</sup> Beyond that, electron injection from gold substrate was reported to cause the 1H-to-1T' phase transition in MoS<sub>2</sub> single layers during thermal annealing leading to the formation of co-existing domains of 1H and 1T' phases; this electron-driven transition was facilitated by the interfacial strain<sup>128</sup> (Fig. 3g). The 1T' phase content reached its maximum (37%) at the annealing temperature about 200–250 °C.

**2.1.4. Applying mechanical strain.** According to density functional calculations, group VI TMDs are predicted to undergo a transition from the 2H to the 1T' phase under tensile strain.<sup>25</sup> It is identified that a 0.3–3% uniaxial strain applied along the *b* (armchair) direction of 2H MoTe<sub>2</sub> monolayers allows for the transition into the 1T' phase at room temperature<sup>25</sup> (Fig. 3i). The strain-induced 2H-to-1T' phase transition in MoTe<sub>2</sub> monolayers is presumed to occur due to simultaneous transition



metal and tellurium atoms displacement and does not proceed through an intermediate 1T state.<sup>129</sup> The process involves a distortion of the transition metal sublattice accompanied by a simultaneous gliding and rippling of chalcogen plane.<sup>129</sup> Experimentally, a reversible strain-induced 2H-to-1T' phase transition has been demonstrated in 20 nm thick MoTe<sub>2</sub> films by applying a tensile strain using an AFM tip<sup>130</sup> (Fig. 3h). Notably, such a phase transition was realised even at room temperature if a small tensile strain of 0.2% was applied.

Among the emerging approaches to induce the 2H-to-1T' phase transformation in group VI TMDs, we would like also to mention the microwave plasma treatment<sup>131,132</sup> and the supercritical CO<sub>2</sub> treatment.<sup>133,134</sup> These less conventional approaches allow for a controllable phase engineering also in the liquid phase. Recently, X. Tong *et al.* have discovered that during the supercritical CO<sub>2</sub> treatment, aqueous suspensions of TMD nanosheets transform into water-in-CO<sub>2</sub> emulsions, causing a strain-induced 2H-to-1T reconstruction in WS<sub>2</sub> and MoS<sub>2</sub> flakes confined at the convex water/CO<sub>2</sub> interface<sup>134</sup> (Fig. 3j). The metastable 1T phase content achieved by this approach is reported to vary from 70%<sup>133</sup> to 90%.<sup>134</sup> Similar to supercritical CO<sub>2</sub> treatment, mild microwave plasma treatment has been reported to introduce the 1T phase domains in the exfoliated 2H MoS<sub>2</sub> nanosheets.<sup>132</sup> C. Sharma *et al.* have used forming gas microwave plasma to convert exfoliated 2H MoS<sub>2</sub> flakes into the 1T phase with a 70% efficiency. The phase transition is assumed to proceed due to the momentum transfer from the plasma ions to the MoS<sub>2</sub> lattice, causing the chalcogen plane gliding.<sup>131</sup>

**2.1.5. The effect of chalcogen vacancies.** Another type of electron donors are the chalcogen vacancies formed in TMD nanosheets. The presence of chalcogen vacancies plays a key role in promoting the 2H-to-1T' phase transition by a dual mechanism: generation of extra charges and weakening of the metal-chalcogen bonds. Both facilitate the structural transformation.<sup>135</sup> Specifically, in MoS<sub>2</sub>, the difference in the formation energy between the 2H and 1T phase decreases almost to zero when the vacancy concentration reaches 8 at%.<sup>136</sup> Acting as point defects, the vacancies also serve as nucleation sites of the 1T' phase being formed in the 2H matrix.

Local 2H-to-1T' transition occurring in mechanically exfoliated few-layered MoTe<sub>2</sub> flakes under laser irradiation is commonly attributed to the formation of Te vacancies due to local heating<sup>137</sup> (Fig. 3e). This transition is irreversible; however, the 2H phase can be restored by high temperature annealing in tellurium atmosphere.<sup>137</sup>

In the case of MoS<sub>2</sub>, it has been shown that the chalcogen vacancies can be introduced into the material by gentle Ar plasma treatment<sup>138</sup> (Fig. 3f). The phase transformation was demonstrated to occur in the vicinity of single point defects and the lateral sizes of the 1T' MoS<sub>2</sub> domains were typically about a few nm.<sup>138</sup> Combined with shadow mask technique, lateral 2H/1T' homojunctions, including 1D conducting 1T' channels, were patterned by Ar plasma treatment in MoS<sub>2</sub> monolayers.<sup>139</sup>

A controlled laser beam 2H/1T' patterning has also been performed in few-layered MoS<sub>2</sub> flakes.<sup>140</sup> However, Ar plasma treatment might be preferable for the phase patterning on TMD monolayers since this method is considered less disruptive compared to laser beam irradiation.<sup>139</sup> Other methods of introducing chalcogen vacancies in MoS<sub>2</sub>, as for example harsh electrochemical etching, to cause the 2H-to-1T structural transformation are just emerging.<sup>135</sup> In their recent work, X. Gan *et al.* have demonstrated that the 1T phase content in multi-layered MoS<sub>2</sub> can be varied by either increasing the etching potential or the reaction time, although the spatial distribution of the metastable phase is not possible to control.<sup>135</sup>

**2.2. Direct growth of the metastable 1T' phase of group VI TMDs.** Direct growth techniques, such as vapour-phase deposition and solution-phase reactions, are based on chemical reactions between the constituent precursors and commonly enable a good level of control of the synthetic conditions. The methods to obtain the metastable 1T(1T') phase of group VI TMDs presented in literature can be divided into two categories: the first rely on the formation of the metastable crystal phase due to the interaction with charged reaction by-products or with the substrate; while the second rely on optimising the reaction kinetics in such a way that the kinetically-preferred, yet thermodynamically metastable, product is formed.

**2.2.1. The role of a substrate in a crystal phase-selective growth.** The use of engineered substrates has been found beneficial for growing the metastable phases of materials with large energy difference between the stable and metastable polymorphs (550 and 330 meV per formula unit in MoS<sub>2</sub> and MoSe<sub>2</sub>, respectively<sup>25</sup>). Nanometre-sized 1T' MoSe<sub>2</sub> islands were selectively formed on Au(111) substrates which were treated with selenium prior to the deposition, whereas only the thermodynamically stable 1H polymorph was obtained if the layer of selenium was not pre-deposited<sup>70</sup> (Fig. 4a). Such a difference in the resulting crystal phase was attributed to the formation of an electron donating Mo layer between the Au surface and the pre-sputtered Se overlayer that facilitate the growth of the 1T' phase. MBE deposition of 1T' MoS<sub>2</sub> nanoclusters (30–50 atoms) on Au(111) substrates has also been reported.<sup>79</sup> It has been also hypothesized that the gold substrate itself serves as an electron reservoir that stabilised the nucleating 1T' phase.

Thermal expansion coefficients of various substrates might be another important growth parameter to consider when the formation of a metastable phase is targeted. Typically, the formation of metastable polymorphs of MoS<sub>2</sub> does not occur in a direct chemical vapour deposition (CVD) reaction due to the huge energy barrier for the 1H-to-1T phase transition. Recently, Z. Wang *et al.* reported the direct growth of the 1T-1H bi-phase TMD monolayers, in which the formation of the metastable phase is associated with the in-plane thermal strain imposed by the substrate during the cooling step.<sup>141</sup>

**2.2.2. Synthesis from electron donating precursors.** High phase-purity mm-sized 1T and 1T' phase MoS<sub>2</sub> single crystals have been demonstrated through a delicate deintercalation of Li or K



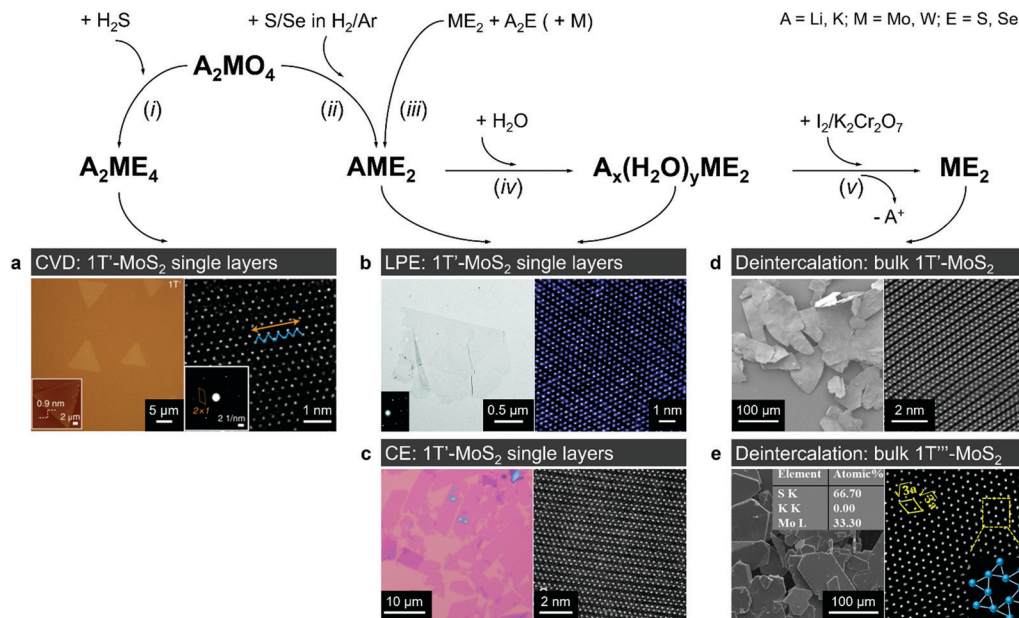
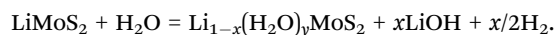
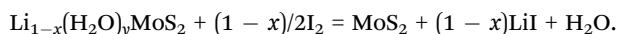


Fig. 5 Direct vapour-phase (CVT, CVD) growth of the 1T'(1T') polymorphs from alkali metal containing precursors. Bulk 1T' crystals can be produced by gentle deintercalation of bulk intercalation compounds  $A_x(H_2O)_yME_2$  (panels d and e); 1T' monolayers can be obtained by chemical (CE, panel c) or liquid-phase (LPE, panel b) exfoliation of bulk intercalation compounds  $A_x(H_2O)_yME_2$  and  $AME_2$ ; high-quality 1T' monolayers can be grown by CVD method from  $A_2ME_4$  precursor (panel a). Separate panels are adapted with permission: (a) from ref. 147, copyright 2018, Springer Nature; (b) from ref. 144, copyright 2017, The Royal Society of Chemistry; (c) from ref. 69, copyright 2019, John Wiley & Sons; (d) from ref. 68, copyright 2018, Springer Nature; and (e) from ref. 56, copyright 2019, American Chemical Society.

ions from  $LiMoS_2$  and  $KMoS_2$  compounds.<sup>27,69,142,143</sup> Ternary  $KMoS_2$  is typically produced by reacting the corresponding alkali metal molybdate,  $K_2MoO_4$ , with dry  $H_2S$  that is then followed by a mild reduction in  $H_2/N_2$  atmosphere<sup>27,142</sup> (Fig. 5, reactions i and ii). Alternatively,  $LiMoS_2$  and  $KMoS_2$  can be obtained by fusing  $MoS_2$  with either  $Li_2S$  or  $K_2S$  under vacuum in the presence of metallic  $Mo$ <sup>56,143,144</sup> (Fig. 5, reaction iii). The deintercalation of these alkali metal thiomolybdate precursors is reported to proceed in two stages. At first, the ternary compounds are exposed to water to hydrate the alkali metal ions intercalated between  $MoS_2$  layers<sup>143</sup> (Fig. 5, reaction iv):



This leads to the lattice expansion in the  $c$  direction. It is assumed that the increased interlayer distance facilitates the alkali metal ions removal owing to the reduced interaction with the chalcogen atoms of basal planes.<sup>143</sup> The alkali metal is removed by a treatment with a mild oxidiser, such as  $I_2$ <sup>27,69</sup> (Fig. 5, reaction v):



Using lithiated compounds might be preferable since smaller ionic radius of lithium leads to its more effective hydration and complete removal from the interlayer gaps, while the deintercalation of larger potassium ions is believed to cause numerous lattice defects.<sup>142,145</sup> It has been reported that the structural disorder in the intercalation compounds of  $MoS_2$  is externally

stabilised by hydrated ions trapped between the layers.<sup>146</sup> Thus,  $K_{0.7}MoS_2$ , similarly to  $Li_xMoS_2$  ( $x \leq 1$ ), typically demonstrates the  $2a \times 2a$  superstructure,<sup>145</sup> whereas the structure of the hydrated  $K_x(H_2O)_yMoS_2$ , refined by diffraction methods (XRD, SAED) and STM imaging, is assigned to the monoclinic symmetry, similar to the one reported for  $WTe_2$  ( $2a \times a$ ).<sup>142,146</sup> Depending on the starting composition of the intercalation compound, different distorted phases of  $MoS_2$  were demonstrated.<sup>146</sup> Gentle deintercalation of  $Li_x(H_2O)_yMoS_2$  with iodine solution was shown to result the high-quality 1T  $MoS_2$  single crystals.<sup>143</sup> Distorted  $MoS_2$  structures produced by oxidation of  $K_x(H_2O)_yMoS_2$  strongly depend on the amount of intercalated potassium.<sup>146</sup> An incomplete oxidation of  $K_x(H_2O)_yMoS_2$  ( $x \approx 0.3$ ) leads to the formation of  $MoS_2$  crystals with the  $2a \times a$  (1T') structure, whereas the final product of a prolonged (complete) oxidation displays the  $a\sqrt{3} \times a\sqrt{3}$  (1T'') superstructure in the  $ab$  plane, where, according to STM, Mo atoms are slightly displaced from their ideal positions forming triangles (trimerization).<sup>58,146</sup> A similar  $2a \times a$  (1T') structure was obtained by oxidation of  $K_x(H_2O)_yMoS_2$  ( $x < 0.3$ )<sup>146</sup> implying a composition-dependant difference in the symmetry of these intercalation compounds. Additionally, the strength of the oxidiser is also reported to affect the crystal symmetry of deintercalated metastable  $MoS_2$ . By using a mild oxidiser, like  $I_2$ , to treat hydrated  $KMoS_2$  crystals, the 1T' phase  $MoS_2$  crystals were synthesised.<sup>68</sup> While by using strong oxidisers,  $K_2Cr_2O_7$  and  $Br_2$ , it was possible to attain the 1T'' polymorph of  $MoS_2$ <sup>27,56,146</sup> (Fig. 5e). It is also reported that intercalation compounds





containing other alkali metals, such as sodium and rubidium, can be used as precursors in the synthesis of 1T' MoS<sub>2</sub> single crystals.<sup>56</sup>

Bulk crystals of the 1T' MoSe<sub>2</sub> with lateral size ~100 μm have been demonstrated through a similar deintercalation (with I<sub>2</sub> solution) of a precursor obtained through a chemical reaction between alkali metal containing molybdenum precursors, such as potassium or sodium molybdate, and elemental selenium annealed in an inert atmosphere in an ampoule<sup>69</sup> or in an open system in a reducing H<sub>2</sub>/Ar atmosphere<sup>68</sup> (Fig. 5d). The role of alkali metal cations (Na, K) is similar to that of lithium in chemical exfoliation process; namely, being electron donors, these cations presumably stabilise the 1T' phase over the 2H phase during the solid-state reaction of the precursor formation. To the best of our knowledge, a direct growth of bulk 1T' WS<sub>2</sub> and 1T' WSe<sub>2</sub> crystals has not been reported so far.

High-purity 1T' monolayers can then be isolated from the parent bulk crystals either mechanically<sup>68</sup> or *via* chemical exfoliation in solution.<sup>69</sup> In the case of chemical exfoliation, an additional lithium intercalation stage of bulk 1T' crystals is introduced followed by exfoliation in water. With such approach, the 1T' phase content in the final product reaches 97%<sup>69</sup> (Fig. 4e and 5c). A successful exfoliation of 1T' MoS<sub>2</sub> monolayers by gentle sonication and even shaking of the hydrated Na<sub>x</sub>MoS<sub>2</sub> or LiMoS<sub>2</sub> crystals immersed in water has also been demonstrated<sup>69,93,144</sup> (Fig. 5b). As shown by C. Guo *et al.*, the alkali metal cations can then be removed from the colloids of MoS<sub>2</sub> monolayers by repeated dialysis.<sup>144</sup> The reported sizes of thus exfoliated 1T' MoS<sub>2</sub> monolayers vary from a few μm<sup>144</sup> to a few tens of μm<sup>69</sup> in the case of agitation-assisted exfoliation and chemical exfoliation, respectively. Direct exfoliation of bulk 1T' TMD crystals can effectively avoid an incomplete phase transition that is typically manifested as co-existing 2H/1T/1T' phases in chemically exfoliated material.<sup>85</sup> Moreover, this approach also allows protecting the in-plane framework of the 1T' phase in TMD monolayers from fragmentation due to lithium intercalation thus potentially producing 1T' domains with larger lateral dimensions.

A phase-selective CVD growth of MoS<sub>2</sub> monolayers on mica based on using a similar type of precursor – potassium thiomolybdate, K<sub>2</sub>MoS<sub>4</sub> – has been proposed recently<sup>147</sup> (Fig. 5a). The authors found out that the crystallisation of the 1T' phase of MoS<sub>2</sub> becomes favourable in potassium-enriched CVD growth due to the stabilisation of the 1T' phase in the intermediate K<sub>x</sub>MoS<sub>2</sub> compound. The crystal phase control was achieved by changing the atmosphere from inert to reducing by adding H<sub>2</sub> into the carrier gas. The MoS<sub>2</sub> flakes produced by decomposition of K<sub>2</sub>MoS<sub>4</sub> precursor in an inert (Ar) atmosphere display the 2H phase, whereas high-purity 1T' MoS<sub>2</sub> flakes were produced if H<sub>2</sub> gas was introduced. Thus, 1T' MoS<sub>2</sub> monolayers with over 90% phase purity can be selectively grown on mica in the temperature range 650–750 °C and H<sub>2</sub> concentration of approximately 5–12% in the carrier gas mixture (Fig. 4). The choice of a suitable substrate may also play an important role since mica may assist in trapping potassium cations produced

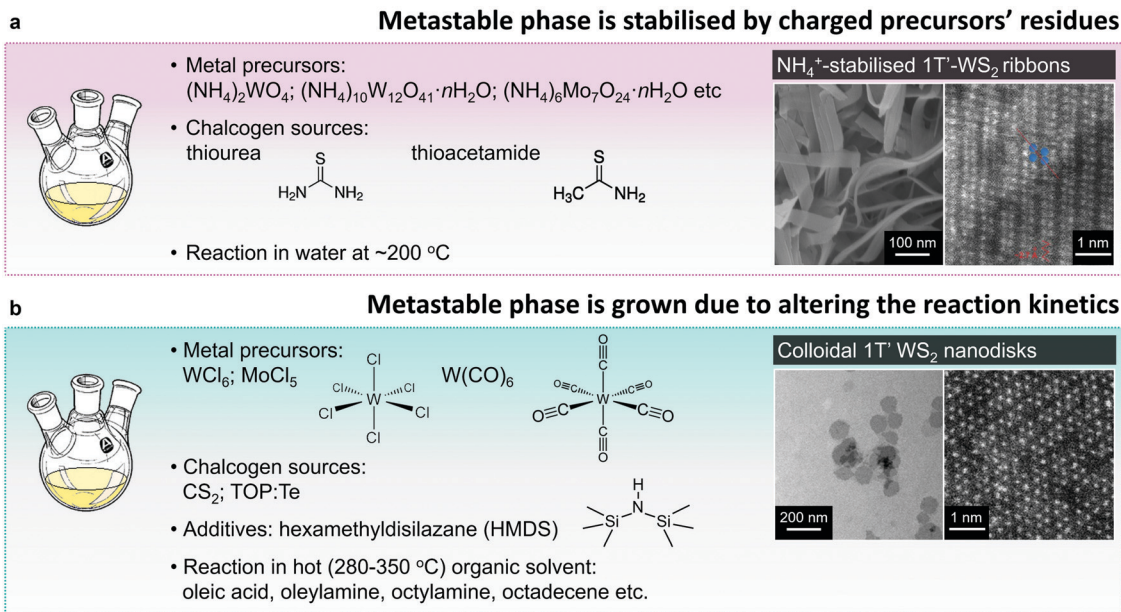
during the precursor decomposition facilitating the formation of the 1T' phase. Interestingly, 1T'/2H heterobilayers could be selectively produced by simply changing pure Ar to Ar/H<sub>2</sub> mixture during the crystal growth. The authors have also demonstrated that Cs<sub>2</sub>MoS<sub>4</sub> and Rb<sub>2</sub>MoS<sub>4</sub> could be suitable for a phase-selective synthesis of 1T' MoS<sub>2</sub>.<sup>147</sup> Additionally, the growth of 1T' WS<sub>2</sub> under the same conditions was achieved if K<sub>2</sub>WS<sub>4</sub> was used as an alkali metal containing precursor.<sup>147</sup>

Similarly to vapour deposition approaches, the reported protocols of hydro- or solvothermal syntheses of the metastable 1T(1T') group VI TMD polymorphs utilise a reaction involving ammonium or alkali metal containing precursors. Typically, 1T(1T') WS<sub>2</sub> and MoS<sub>2</sub> nanostructures are obtained by reacting ammonium containing salts, like (NH<sub>4</sub>)<sub>2</sub>WO<sub>4</sub>,<sup>148</sup> (NH<sub>4</sub>)<sub>10</sub>W<sub>12</sub>O<sub>41</sub>·H<sub>2</sub>O,<sup>82</sup> and (NH<sub>4</sub>)<sub>6</sub>Mo<sub>7</sub>O<sub>24</sub>,<sup>81</sup> with the chalcogen precursor, such as thiourea or thioacetamide, or by thermal decomposition of ammonium thiomolybdate<sup>149</sup> in solution phase (Fig. 6a). Thus produced material exhibits the morphology of ultra-thin nanosheets or nanoribbons composed of nm-sized co-existing 2H and 1T(1T') single-crystalline domains. The 1T(1T') phase content in these nanostructures reaches ~61.6%<sup>81</sup> (Fig. 4). The metastable phases are likely to be stabilised by charged precursor residues trapped between the layers of few-layered nanosheets and nanoribbons. Specifically, presence of the intercalated ammonium uniformly distributed across the nanosheets has been repeatedly verified by elemental mapping and XPS analysis.<sup>81,82,150</sup> Apart from cationic species, it is also assumed that the 1T' phase in MoS<sub>2</sub> produced by a solvothermal reaction from (HN<sub>4</sub>)<sub>2</sub>MoS<sub>4</sub> can be stabilised by intercalated anionic thiomolybdate moieties.<sup>149</sup> To verify that, the complete removal of salt residues was achieved by prolonged dialysis and the absence of intercalated ammonium in the synthesised material has been demonstrated by XPS analysis. The solvothermal strategy has been further adapted to form arrays of the 1T'-rich (60%) nanosheets (MoS<sub>2</sub>) on conductive support (single wall carbon nanotubes) for the application in electrochemical devices.<sup>150</sup>

An interesting chemical approach to produce the 1T phase of TMDs that rely on the conversion of non-van der Waals solids, such as MAX and MXene phases, has been reported recently.<sup>151,152</sup> In this approach, the transition metal atom coordination might be templated by the symmetry of a parent MAX phase. J. Tang *et al.* have demonstrated that high quality TiS<sub>2</sub> sheets can be synthesised by a high temperature chalcogenation of Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> in H<sub>2</sub>S atmosphere.<sup>151</sup> This approach has been further extended to produce group IV and VI disulphides and diselenides, and 1T phases of MoS<sub>2</sub> and WS<sub>2</sub>, in particular.<sup>152</sup> Interestingly, the metastable 1T phases have only been obtained in the presence of phosphorus vapour in the chalcogen atmosphere, leading to the formation of P-doped 1T/2H mixed-phase material.<sup>152</sup>

**2.2.3. Reaction kinetics-driven direct synthesis of the metastable phase.** Phase-selective synthesis through a fine tuning of the growth conditions is achievable in TMDs with a small energy difference between the thermodynamically stable and metastable polymorphs. MoTe<sub>2</sub> displays the difference of only





**Fig. 6** Direct solution-phase growth of the 1T' polymorphs. (a) Hydro-/solvothetmal reactions between the ammonium containing precursors produce the metastable 1T' phase stabilised by precursors' residues. (b) Colloidal synthesis of the metastable 1T' phase achieved due to the fine control over the reaction kinetics. Figures in panels (a) and (b) are adapted with permission from ref. 82, copyright 2015, John Wiley & Sons, and ref. 80, copyright 2014, American Chemical Society.

about 35–40 meV per formula unit between the 2H and 1T' polymorphs,<sup>25,153</sup> which is markedly different from the rest of group VI TMDs, and therefore is the most suitable candidate for developing the phase-selective growth strategies. The second candidate could be  $\text{WTe}_2$  with the difference between stable 1T' and metastable 2H being only 90 meV per formula unit.<sup>25</sup> However, the synthesis of the metastable 2H  $\text{WTe}_2$  has not been demonstrated so far.

In vapour phase deposition approaches, the crystal phase tunability has been demonstrated by finely controlling the reaction parameters. For instance, islands of 1T'  $\text{WSe}_2$  extended over  $\sim 100$  nm have been selectively synthesised on bilayer graphene on 6H-SiC(0001) substrates by molecular beam epitaxy (MBE) by keeping the deposition temperature not higher than  $280^\circ\text{C}$ <sup>20,21</sup> (Fig. 4). In CVD approach, the crystal phase selectivity is often achieved by controlling the precursor flux. Thus, polycrystalline films of 1T'  $\text{MoTe}_2$  with single crystalline domains of  $\sim 10$  nm are commonly reported through tellurisation of Mo films or Mo islands in a tellurium-deficient atmosphere.<sup>60,154</sup> In this case, the formation of the metastable phase is attributed to tellurium deficiency in the produced  $\text{MoTe}_x$  material that lowers the free energy of 1T' polymorph.<sup>154</sup> In particular, Te deficiency greater than 2 at% suffices the stabilisation of the 1T' phase over the 2H in  $\text{MoTe}_2$ .<sup>25,155</sup> The lateral dimensions of such phase-uniform polycrystalline 1T'  $\text{MoTe}_2$  films are reported to reach a few  $\text{cm}^2$ .<sup>59</sup> Interestingly, prolonging the tellurisation time has facilitated a complete solid-to-solid 1T'-to-2H phase transformation; moreover, the formation of single crystalline 2H  $\text{MoTe}_2$  domains as large as 2.34 mm in diameter has been achieved by annealing in a tellurium-enriched atmosphere.<sup>154</sup> Additionally, the formation of the metastable 1T'

phase by tellurisation of Mo films was attributed to a mechanical strain accumulated between the tellurised and pristine Mo regions.<sup>59</sup> In such a case, during a prolonged tellurisation, gradual relaxation of the strain triggers structural transformation into the 2H phase.<sup>59</sup> It has also been reported that either Mo or  $\text{MoO}_3$  films can be used as a precursor for the synthesis of few-layered 1T'  $\text{MoTe}_2$ .<sup>156</sup> Regardless of the precursor choice, the 1T' polymorph of  $\text{MoTe}_2$  was obtained if Te supply was insufficient; however, using  $\text{MoO}_3$  as precursor led to the formation of morphologically uniform,  $\sim 3$  nm thick films (film roughness  $\approx 1$  nm), whereas the roughness of the films produced from Mo precursor under the same conditions was  $\approx 5$  nm.<sup>156</sup> The choice of molybdenum precursor may, however, plays an important role in the CVD synthesis under high tellurium flux. Thus, the same research group have found that  $\text{MoO}_3$  reacts more easily with Te vapour forming the 2H phase, while Mo or  $\text{MoO}_{3-x}$  produce the 1T' phase under the same conditions.<sup>157</sup> Recently, L. Zhou *et al.* have demonstrated that planar 1T'/2H  $\text{MoTe}_2$  heterostructures can be produced by a one-step tellurisation of  $\text{MoO}_3$  films using a partially overlapped geometry.<sup>158</sup> The formation of 2H phase is attributed to the trapping of Te gas in the gap between the substrate and the cover and therefore to higher concentration of Te residing over the surface of  $\text{MoO}_3$  film that led to more complete tellurisation; whereas uniform 1T' phase is obtained in the exposed areas of the substrate. L. Yang *et al.* have demonstrated that the rate of  $\text{MoO}_3$  film tellurisation can be controlled by adjusting both the carrier gas flow rate and the reaction temperature; this has allowed the authors to identify the conditions for the selective growth of 1T' and 2H  $\text{MoTe}_2$  few-layered films.<sup>159</sup> Alternatively, lateral 1T'/2H homojunctions can be formed in few-layered  $\text{MoTe}_2$  films by optimising the precursor



flux in such a way that the incomplete 1T'-to-2H transition is arrested leading to isolated 2H domains being formed in continuous 1T' matrix.<sup>60,160</sup> However attractive as a lever to direct the crystallisation of MoTe<sub>2</sub> in the metastable crystal phase, tellurium deficiency of the 1T' phase produced *via* metal film tellurisation can significantly hamper its potential use in electronics where high structural perfection is required.<sup>161</sup> High-quality 1T' MoTe<sub>2</sub> single crystals can be grown directly from elemental Mo and Te by chemical vapour transport (CVT) reaction.<sup>58,153,162</sup> In this method, a stoichiometric mixture of elemental precursors is sealed in an evacuated ampoule and then maintained at elevated temperature, typically 800–1000 °C, to produce polycrystalline MoTe<sub>2</sub> product, which is further recrystallised by CVT method using either TeCl<sub>4</sub>,<sup>63,77</sup> TeBr<sub>4</sub>,<sup>162</sup> or I<sub>2</sub><sup>58</sup> as transport agent. The monoclinic 1T' phase of MoTe<sub>2</sub> is thermodynamically stable at the temperatures reported for the CVT processes (900–1100 °C),<sup>27,153</sup> and can be preserved in the final product if rapidly quenched down to room temperature.<sup>58,77</sup> The dimensions of thus produced 1T' MoTe<sub>2</sub> single crystals are reported to be as large as 17 × 10 × 2 mm<sup>3</sup> (162) (Fig. 4f). Moreover, D. Keum *et al.* have shown that stoichiometric 1T' MoTe<sub>2</sub> single crystals can be grown from Mo and Te powders in a molten salt (NaCl).<sup>153</sup> The use of salt was recognised essential since this allows dissolving Te and thus prevents its sublimation and formation of Te deficient product. Similar to ref. 58, 1T' MoTe<sub>2</sub> single crystals were attained if the product was quenched, whereas the 2H phase is a product of slow cooling to room temperature. Atomically-thin 1T' MoTe<sub>2</sub> flakes with the lateral size exceeding 10 μm and exhibiting high structural perfection can then be mechanically exfoliated from the bulk crystals onto supports (Si/SiO<sub>2</sub>) for further examination.<sup>58,153</sup>

By optimising the quenching regime, all three, 1T, 1T', and 2H, phases of MoTe<sub>2</sub> can be selectively produced starting from MoO<sub>3</sub> and Te powder precursors.<sup>163</sup> The fact that all three phases were obtained through the same growth procedure with the only variation being the cooling rate led the authors to conclude that during the direct CVD reaction, MoTe<sub>2</sub> nucleates in the 1T phase. The 1T phase is obtained if after the growth stage the final product is rapidly cooled in the presence of CO<sub>2</sub> and H<sub>2</sub> gases. Trapping of the metastable 1T phase in this case might be caused by the interplay of CO<sub>2</sub> and H<sub>2</sub> adsorption and interaction with the substrate. If the temperature is reduced more gradually then the 1T' phase is found to be dominating in the final product. Finally, natural cooling allows enough time for the crystal lattice to transform into the 2H phase.<sup>119</sup>

A different approach to obtain the metastable phases on monolayered TMDs *via* CVD has been demonstrated for the synthesis of the 1T polymorph of WS<sub>2</sub><sup>164</sup> and it is based on the use of a catalyst. A mixture of Fe<sub>3</sub>O<sub>4</sub> and NaCl, serving as a catalyst and a growth promoter, was added into the CVD-tubular furnace (WO<sub>3</sub> and elemental S as material precursors) leading to the formation of monolayered 1T/2H WS<sub>2</sub> flakes with a butterfly-like shapes, where one of the laterally stitched wings adopted the 1T and the other – the 2H phase. Single crystalline domains in these butterfly-like WS<sub>2</sub> monolayers reached 80 μm (Fig. 4g). It has also been shown that these 1T/2H WS<sub>2</sub> structures

could serve as seeds for the laterally expanded 1T/1T and 2H/2H WS<sub>2</sub>/MoS<sub>2</sub> planar heterostructures.<sup>164</sup>

Owing to the non-equilibrium nature of liquid-phase crystallisation, the reaction pathway in the systems with diverse polymorphism greatly depends on subtle variations in growth conditions and may often lead to the formation of metastable reaction products.<sup>165–167</sup> First, the crystal growth conditions may be optimised in such a way that the metastable reaction product is kinetically preferred over the equilibrium product.<sup>168</sup> In fact, it is known that in the systems with numerous polymorphs present, the crystallisation often proceeds through the formation of thermodynamically unstable polymorphs that successively transform into the equilibrium phase; this is known as Ostwald's step rule (Fig. 7b).<sup>169</sup> The formation of the metastable phase can dominate the crystallisation kinetics owing to a lower nucleation barrier than that of the equilibrium phase<sup>170</sup> (Fig. 7a) and so one can assume that the metastable polymorph can be isolated if the crystallisation is arrested before the next, more stable, polymorph nucleates and grows. It is assumed that high supersaturations might be prerequisite for such multi-step crystallisation.<sup>171,172</sup> Also, the free energy of different polymorphs may be altered due to the interaction of the growing crystal with its environment, such as solvent and, especially in the case of colloidal growth, capping ligands, thus lowering the energy barrier for the preferential formation of metastable polymorph and potentially preventing its transformation into the thermodynamically stable one.<sup>80</sup> In TMD systems, the effects of nanocrystal nucleation and growth kinetics on the preferential formation of one polymorph over the other are still at very early stages of investigation.

Y. Sun *et al.* have reported that the metastable 1T' polymorph of MoTe<sub>2</sub> is directly accessible in colloidal synthesis.<sup>173</sup> In their work, flower-like 1T' MoTe<sub>2</sub> nanostructures having diameters of 100–150 nm were grown in solution phase by continuously injecting molybdenum precursor into a hot mixture (300 °C) containing trioctylphosphine telluride (TOP:Te) in

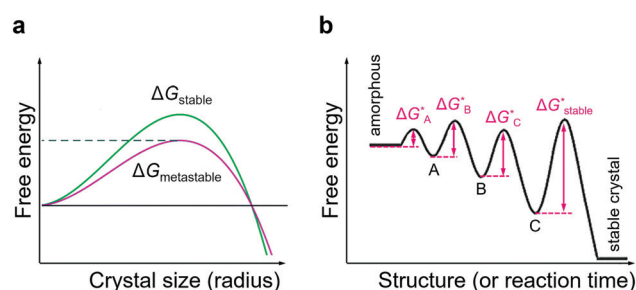


Fig. 7 Homogenous nucleation of the metastable phase in a system with diverse polymorphism. (a) Nucleation barriers of the metastable (purple) and the stable (green) phase in a supersaturated solution. (b) Schematic diagram illustrating the crystallisation pathway of the stable phase proceeding through a series of intermediate metastable phases A, B, and C;  $\Delta G^*$  denotes the energy barriers of the respective phase transitions. Separate panels are adapted with permission: (a) from ref. 170, copyright 2017, The Royal Society of Chemistry; (b) from ref. [Nature Physics volume 5, pages 68–73 (2009); 10.1038/nphys1148], copyright 2008, Springer Nature.





trioctylphosphine (TOP) and oleylamine (OlAm). Molybdenum precursor was prepared by dissolving  $\text{MoCl}_5$  in oleic acid (OlAc). Adding hexamethyldisilazane (HMDS) into tellurium precursor prior to the molybdenum complex injection is required for the formation of crystalline, uniform  $\text{MoTe}_2$  nanostructures. The authors hypothesised that the interaction with organic ligands can effectively modify the surface energy of  $\text{MoTe}_2$  nanocrystals and so further lower the energy barrier for forming the metastable structure. Once formed, the metastable  $1\text{T}'$  phase is kinetically trapped due to the pinning at grain boundaries in the petals of  $\text{MoTe}_2$  nanoflowers. Following a similar synthesis protocol, Y. Sun *et al.* have obtained high-quality  $\text{WTe}_2$  nanoflowers using  $\text{WCl}_6$  complex in OlAc as transition metal precursor.<sup>174</sup> Interestingly, colloidal  $\text{WTe}_2$  nanoflowers are attained in the orthorhombic Td phase. Additionally, the authors have demonstrated that by using mixtures of  $\text{MoCl}_5$  and  $\text{WCl}_6$  as transition metal source, nanostructured  $\text{W}_x\text{Mo}_{1-x}\text{Te}_2$  solid solutions with composition-dependant Td/ $1\text{T}'$  crystal phase heterogeneity spanning the Td  $\text{WTe}_2$  and  $1\text{T}'$   $\text{MoTe}_2$  polymorphs can be accessed.<sup>174</sup>

On the other hand, colloidal synthesis of the metastable polymorphs of group VI disulphides and diselenides remains majorly unexplored. Generally, the metastable polymorphs of these materials are considerably less accessible due to the greater energy barriers of their formation.<sup>46</sup> Consequently, the solution-phase growth of group VI disulphides and diselenides, whether by a thermolytic decomposition of a single source precursor, like  $[\text{Mo}_2\text{O}_2\text{S}_2(\text{S}_2\text{COEt})_2]$ ,<sup>175</sup> or by a direct reaction between molecular precursors, such as  $\text{MoCl}_5$  and  $\text{CS}_2$ ,<sup>176</sup> commonly results in the thermodynamically stable 2H polymorph. To date, the reported protocols include the use of tungsten and molybdenum hexacarbonyls<sup>177</sup> and chlorides,<sup>178–180</sup> molybdenyl acetylacetonate,<sup>181</sup> sodium molybdate<sup>179,182</sup> and sodium tungstate,<sup>183</sup> as a source of transition metal. Chalcogen precursors include carbon disulphide,<sup>176,180</sup> diphenyl diselenide,<sup>177,180,181</sup> complexes of elemental sulphur<sup>178</sup> or selenium<sup>179,182,183</sup> in organic solvents, like oleylamine or octadecene. Despite such a variety of precursors available for a direct solution-phase growth of group VI TMDs, only the semiconducting 2H phase was synthesised. However, a slow sulphidation of colloidal  $\text{W}_{18}\text{O}_{49}$  nanorods, serving as sacrificial seeds, by *in situ* generated  $\text{H}_2\text{S}$  has led to the formation of irregularly-shaped ultrathin  $\text{WS}_2$  disks demonstrating co-existing 2H and  $1\text{T}'$  crystal phases.<sup>184</sup>

B. Mahler *et al.* have reported on a phase-selective colloidal synthesis of  $\text{WS}_2$  nanosheets<sup>80</sup> (Fig. 6b). Disk-like  $\text{WS}_2$  nanostructures with the distorted octahedral  $1\text{T}'$  crystal structure were synthesised through a reaction between  $\text{WCl}_6$  complex in oleic acid and  $\text{CS}_2$  in hot (320 °C) oleylamine, whereas the 2H  $\text{WS}_2$  nanosheets were grown from the same molecular precursors in the presence of an additive, hexamethyldisilazane (HMDS). HMDS is an essential chemical in colloidal synthesis of another class of layered materials: narrow-gap  $\text{A}^{\text{IV}}\text{-B}^{\text{VI}}$  semiconductors, such as GeS and SnS.<sup>185,186</sup> It has been noted that the growth from the respective transition metal halides in hot oleylamine in the absence of HMDS leads to amorphous

product, whereas in the presence of HMDS, highly-uniform, single crystalline plates of SnS and GeS are being formed.<sup>185–187</sup> In the case group VI TMDs, Y. Sun *et al.* have reported that the presence of HMDS is required for the formation of highly crystalline  $\text{MoTe}_2$ ,<sup>173</sup>  $\text{WTe}_2$ <sup>174</sup> and  $\text{W}_x\text{Mo}_{1-x}\text{Te}_2$ <sup>174</sup> colloidal nanosheets. M. Kobylinski *et al.* have suggested that in general, the role of HMDS in colloidal synthesis from transition metal chlorides is two-fold: first, HMDS may react with the transition metal source, forming the corresponding (trimethylsilyl)amino complex; second, HMDS may bind chloride ions and remove them from the reaction mixture as (trimethylsilyl) chloride. While B. Mahler *et al.* hypothesised that in the specific case of colloidal synthesis of  $\text{WS}_2$ , HMDS may effectively tune the reactivity of tungsten precursor enabling the phase-selective growth.<sup>80</sup> Specifically, they presume that HMDS may partially substitute the oleate ligands complexing tungsten ions, leading to the formation of a more active form of tungsten precursor. Their alternative hypothesis suggests that HMDS may act as a capping ligand during the nanocrystal growth, stabilising the 2H polymorph over the  $1\text{T}'$ . In order to verify which of the two possible mechanism is more likely to drive the reaction, the authors replace oleic acid (OlAc) with another strongly binding ligand – dodecanethiol (DDT). The authors suggest that unlike OlAc, DDT does not react with HMDS thus the formation of an activated tungsten (trimethylsilyl)amino complex is not possible in the DDT protocol. Indeed, the material grown from  $\text{WCl}_6$ -DDT complex in the presence of HMDS exhibits the  $1\text{T}'$  phase. These findings suggest that the formation of the thermodynamically stable 2H phase of  $\text{WS}_2$  can be mainly attributed to the activation of tungsten complex by reacting with HMDS, whereas less reactive tungsten oleate may lead to the direct growth of the  $1\text{T}'$  phase. It should be noted that a detailed study of the suggested intermediates in the  $\text{WCl}_6$ -OlAc and  $\text{WCl}_6$ -DDT systems, their activity and possible reactions with HMDS is not provided in the cited work.<sup>80</sup>

In addition to  $1\text{T}'$   $\text{WS}_2$ , the direct growth of the metastable  $1\text{T}'$  polymorph of  $\text{WSe}_2$  has been reported recently<sup>71,188</sup> (Fig. 4d). It is worth mentioning that the  $1\text{T}'$  polymorphs of  $\text{WSe}_2$  and  $\text{MoSe}_2$  were not previously demonstrated by Li intercalation and were only achieved *via* direct synthesis.<sup>68,71</sup> Uniform  $1\text{T}'$   $\text{WSe}_2$  nanoflowers with an average diameter of 200 nm and composed of highly-crystalline atomically-thin petals were grown through a reaction between  $\text{W}(\text{CO})_6$  and trioctylphosphine selenide (TOP:Se) in hot oleic acid (300 °C).<sup>71</sup> Prior to that work, only a few experimental protocols on colloidal synthesis of  $\text{WSe}_2$  nanostructures were present, but only the formation of thermodynamically stable 2H polymorph was reported.<sup>177,180</sup> In order to elucidate the reaction pathway, M. Sokolikova *et al.* have performed a parametric study of growth conditions. The activity of both precursors was systematically varied by replacing  $\text{W}(\text{CO})_6$  with  $\text{WCl}_6$ , TOP:Se with elemental Se; additionally, the activity of tungsten precursor was modified by complexing with two additives: hexamethyldisilazane (HMDS) and tetradecylphosphonic acid (TDPA). Interestingly, the metastable  $1\text{T}'$  crystal phase was consistently grown under these conditions; only negligible changes in the



morphology were reported. Based on these observations, the authors have proposed that under the described experimental conditions, the formation of metastable 1T' polymorph might be attributed to the crystallisation in kinetically-driven regime; however, switching between the supposedly kinetically-driven and thermodynamically-controlled regimes has not been unambiguously demonstrated.

Later, J. Geisenhoff *et al.* have demonstrated a phase-selective synthesis of WSe<sub>2</sub> nanoflowers that was achieved in a biligand system of trioctylphosphine oxide (TOPO) and oleic acid (OAc).<sup>188</sup> The metastable 1T' phase of WSe<sub>2</sub> was obtained when the fraction of OAc in the ligand mixture was high, while keeping the oleic acid fraction to minimum has led to the growth of the thermodynamically stable 2H phase. The authors indicate that by coordinating with oleic acid, a less reactive form of tungsten precursor is formed; this leads to nucleation of fewer seeds and to larger nanocrystals compared to the synthesis in TOPO-rich mixtures. However, the effect of ligand mixture composition on the preferential nucleation of one of the other crystal phases remains unclear. Regardless of the ligand ratio, colloidal WSe<sub>2</sub> was found to nucleate in the metastable 1T' phase; however, in the systems, containing low fraction of OAc, it converts into the 2H phase over time.<sup>188</sup> Whereas in the case of ligand mixtures with high OAc content or of syntheses in OAc, the final product seemingly retains the metastable 1T' phase.<sup>188</sup>

Recently, it has been proposed that the energy barrier for the homogeneous nucleation of the metastable 1T phase could be lowered in a magneto-hydrothermal process.<sup>136</sup> The preferential nucleation of the 1T polymorph in high magnetic fields is facilitated by the higher magnetic susceptibility of the 1T structure, compared to the 2H phase. Using this approach, W. Ding *et al.* have selectively synthesised 1T MoS<sub>2</sub> and WS<sub>2</sub> nanosheets through a hydrothermal reaction between either ammonium heptamolybdate ((NH<sub>4</sub>)<sub>6</sub>Mo<sub>7</sub>O<sub>24</sub>·4H<sub>2</sub>O) and thiourea or tungsten hexachloride (WCl<sub>6</sub>), and thioacetamide under an external magnetic field of ~9 T.<sup>136</sup> The nucleation of the 2H phase from the same precursors in the absence of magnetic field (control) has also been demonstrated.

Considering the present evidence, we suggest that the understanding of crystallisation pathways leading to the formation of metastable polymorphs of group VI TMDs in colloidal solution is at early stages. A thorough investigation of the nucleation mechanism, including the nature of reaction intermediates formed from molecular precursors, possible prenucleation clusters, energetics of the nanocrystal–solution interface, and surface dynamics of capping ligands on the nanocrystal facets, is needed in order to gain control over the phase-selective growth of various polymorphs of group VI TMDs.

It is interesting to note that unlike other van der Waals solids, such as colloidal SnS<sup>186</sup> and SnSe,<sup>189</sup> GeS and GeSe,<sup>185</sup> InSe,<sup>190</sup> and group IV<sup>191,192</sup> and V<sup>192</sup> disulphides and diselenides, that are acquired in a form of two-dimensional platelets and disks, group VI TMDs obtained through a solution-phase (hydro-/solvothermal<sup>81,193</sup> and colloidal<sup>71,173,174,179,180</sup>) reaction between similar type of precursors typically exhibit branched

morphology of flower-like structures composed of ultra-thin petals. The overall anisotropic morphology of van der Waals compounds is expected since the lateral extension in the *ab* plane is energetically preferred due to surface energy difference between the undercoordinated edge facets and the basal planes, where dangling bonds are absent.<sup>192,194</sup> In some cases, like layered SnSe<sup>189</sup> and Bi<sub>2</sub>Te<sub>3</sub>,<sup>195</sup> and non-layered PbS<sup>196</sup> systems, the growth of large regularly-shaped single-crystalline nanostructures does not proceed through the extension of primary nucleated seeds but rather occurs *via* the two-dimensional oriented attachment of the seeds. The self-assembly process is driven by the minimisation of surface energy associated with the unsaturated bonds, for example, in the lateral surface of SnSe seeds, and can be effectively controlled by varying the precursors' concentrations.<sup>189</sup> However, neither of these crystallisation mechanisms can explain the formation of branched flower-like structures of group VI TMDs produced *via* a solution-phase reaction. The formation of 3D flower-like nanostructures composed of individual 2D nanosheets might be caused by a nucleation from a dense amorphous state formed in the supersaturated solution as proposed for TiS<sub>2</sub> nanoflowers.<sup>197</sup> S. Prabakar *et al.* have found out that TiS<sub>2</sub> flower-like structures evolve in colloidal solution starting with the formation of amorphous globules and proceeding through the appearance of randomly oriented crystalline platelets on the surface of amorphous globules to the growth of these petals originating from the central core so that the amorphous core is no longer apparent.<sup>197</sup> Similar formation mechanism was later described by O. Meiron *et al.*<sup>198</sup> and by C. Zhang *et al.*<sup>182</sup> for colloidal MoS<sub>2-x</sub>Se<sub>2-2x</sub> and MoSe<sub>2</sub> nanoflowers, respectively. O Meiron *et al.* have identified amorphous blobs that are formed at the very early stages of the reaction; these later serve as material reservoirs for the curled and tangled nanosheets that crystallise inside the amorphous clusters.<sup>198</sup> Over time, amorphous material is fully consumed and uniform nanoflowers are formed. Despite a considerable progress in controlled synthesis of TMDs, we highlight that a deeper understanding of solution-phase nucleation and crystallisation processes of group VI TMDs is required in order to fully employ their potential in practical applications.

### 3. Comparison of various techniques to access the metastable 1T(1T') phases

Fig. 4 provides a comprehensive overview of the reported approaches to access the metastable 1T(1T') phases of group VI TMDs. The data points were taken from the references, where the lateral dimensions of continuous domains (not necessarily single crystalline) of the metastable 1T(1T') phase were clearly stated and the metastable phase content was determined from the results of either XPS or Raman spectroscopy. These are listed in Table 1.

The most commonly reported approach to synthesise the metastable phase in measurable quantities is chemical exfoliation from bulk TMD powders. In chemically exfoliated MoS<sub>2</sub> and WS<sub>2</sub>, the 1T(1T') phase is found as nm-scale domains embedded in the 2H matrix, although the total 1T(1T') phase content reaches ~70–80%. It should be noted that there is little evidence that this approach is effective in the case of group VI diselenides.



**Table 1** Comparison of various synthesis techniques to access the metastable 1T(1T') phases of group VI TMDs

Material	Lateral size	% of 1T(1T')	Ref.
<b>Molecular beam epitaxy</b>			
MoS <sub>2</sub>	4 nm	100	79
MoSe <sub>2</sub>	10 nm	100	70
WSe <sub>2</sub>	100 nm	100	21
<b>Chemical vapour transport</b>			
MoS <sub>2</sub>	100 μm	100	56
	150 μm	100	68
	500 μm	100	69
MoSe <sub>2</sub>	170 μm	100	68
MoTe <sub>2</sub>	10 μm	100	58
	1 cm	100	153
	1.7 cm	100	162
<b>Chemical vapour deposition</b>			
MoS <sub>2</sub>	10 μm	100	147
WS <sub>2</sub>	80 μm	100	164
<b>Chemical exfoliation (both 2H and 1T' crystals)</b>			
MoS <sub>2</sub>	3.5 nm	70	112
WS <sub>2</sub>	3 nm	80	112
	100 nm	80	83
	1 μm	70	101
MoS <sub>2</sub>	700 nm	80	103
	10 μm	97	69
	60 μm	97	69
	500 nm	100	93
	1.5 μm	100	265
	1.5 μm	100	265
<b>Direct colloidal/hydrothermal synthesis</b>			
MoS <sub>2</sub>	200 nm	61.6	81
WS <sub>2</sub>	100 nm	50	80
	162 nm	66.4	148
WSe <sub>2</sub>	100 nm	60	71

On the other hand, CVT growth, that up until recently was employed to synthesise cm-sized single crystals of 1T' MoTe<sub>2</sub> only, has been extended to the case of molybdenum disulphide and diselenide. The growth of the metastable phase is achieved with the assistance of alkali metal containing precursors and 1T' MoS<sub>2</sub> and 1T' MoSe<sub>2</sub> crystals with lateral dimensions of a few hundreds of μm are produced by gentle deintercalation of the intermediate products. Furthermore, chemical exfoliation from bulk 1T' MoS<sub>2</sub> and 1T MoSe<sub>2</sub> crystals results in unprecedentedly large nanosheets with the metastable phase content of nearly 100%. CVD growth, that typically leads to the formation of the thermodynamically stable polymorph, can be optimised in such a way that the metastable phase is selectively produced. Thus, in the presence of a catalyst or alkali metal containing precursors, the growth of single crystalline domains of 1T WS<sub>2</sub> and 1T' MoS<sub>2</sub> with the sizes of a few tens of μm has been reported. Another deposition approach, molecular beam epitaxy, allows for the growth of islands of 1T' phase with lateral dimensions varying from a few nm as in the case of MoS<sub>2</sub> and MoSe<sub>2</sub> to nearly reaching 100 nm in the case of WSe<sub>2</sub>. It should be noted that the selectivity of MBE process is rather low at present as the formation of separate domains of the 2H phase is often reported.

Solution-phase approaches, such as hydrothermal and colloidal reactions, allow for the synthesis of measurable quantities of the 1T' phase containing material, in some cases

reaching a few grams. The lateral size of 1T' domains are reported within the range from a few tens to a few hundreds of nm, and the metastable phase content is about 50–60%.

#### 4. How stable the metastable 1T(1T') phases are?

Although being thermodynamically metastable under ambient conditions, the 1T(1T') polymorph is kinetically stable over the course of a few days to a few months, depending on the material morphology and environment. Thus, chemically exfoliated 1T' and 1T'' MoS<sub>2</sub> monolayers gradually convert back to the 2H phase over 3–5 days in air as demonstrated by Raman<sup>87</sup> and STM<sup>146</sup> studies. The conversion is attributed to the restacking of exfoliated material in suspension that leads to the van der Waals interactions between the layers being recovered. The transformation into the 2H phase is significantly promoted when the exfoliated 1T' TMDs are deposited to form thin films owing to the removal of water from the interlayer space of exfoliated sheets during drying.<sup>87</sup> Synthesised through a direct hydrothermal reaction, the 1T-phase MoS<sub>2</sub> was reported to slowly convert into the 2H phase over 90 days in water. This was accompanied by a gradual loss of the catalytic activity.<sup>199</sup>

#### 5. Property-related applications of the metastable 1T(1T') phases

The metastable 1T(1T') polymorphs of group VI TMDs are found to be promising for various applications, including electrocatalysis,<sup>83,103</sup> electrochemical energy storage,<sup>101,200</sup> sensors,<sup>201,202</sup> transparent electrodes for nanoscaled optoelectronics<sup>110</sup> and photovoltaics,<sup>203</sup> transport layers for solar cells,<sup>204</sup> and even microwave absorbers for application in electromagnetic shielding and stealth camouflage.<sup>205</sup> The incorporation of domains of the 1T phase in the 2H matrix induces a robust ferromagnetism in the intrinsically non-magnetic semiconducting TMDs, further expanding the range of potential applications to spintronic devices.<sup>206,207</sup> The difficulty in accessing the input of the metastable phase is commonly associated with the fact that in many cases, the 1T(1T') phase co-exists with the thermodynamically stable 2H phase in these produced materials. Here, we review the state-of-the-art research articles related to the understanding of how the intrinsic properties of the 1T(1T') phases can enable novel applications.

##### 5.1. Basal plane activation for efficient electrocatalysis.

Group VI TMDs have been actively studied as catalysts for the hydrogen evolution reaction (HER).<sup>83,115,208–215</sup> Electrocatalytic hydrogen generation is practically achieved by water electrolysis, typically performed in acidic medium. The initial reduction of solvated protons by electrons from the external circuit leads to the formation of adsorbed neutral hydrogen atom on the catalyst surface; this step is known as the Volmer reaction. Then, the formation of molecular hydrogen proceeds through either the Tafel reaction or the Heyrovsky reaction. The former one involves a reaction between two adsorbed surface hydrogen atoms. In the latter, the adsorbed hydrogen reacts with the solvated proton followed by reduction by electrons from the external circuit.<sup>216</sup> Depending on the electrocatalyst nature, the rate limiting step of the HER can be either the initial proton reduction or the following formation of molecular hydrogen.







**Fig. 8** Applications of the metastable 1T(1T') phases of group VI TMDs in electrochemical energy storage and conversion. (a) SEM image (left) of micron-sized chemically exfoliated 1T' MoS<sub>2</sub> nanosheets and the polarisation curves (right) of 1T' and 2H MoS<sub>2</sub> demonstrating the basal plane activation in the 1T' polymorph. (b) SEM image (left) of 1T' WSe<sub>2</sub> nanosheets grown directly on carbon fibres demonstrating lower HER overpotentials and lower Tafel slopes (right) compared to the corresponding 2H phase. (c) Polarisation curves (left) and the corresponding Tafel slopes (right) of the CVD grown 1T' MoS<sub>2</sub> flakes. (d) SEM image (left) of micron-sized chemically exfoliated MoS<sub>2</sub> nanosheets and CV curves (right) of a symmetric supercapacitor device in an organic electrolyte acquired at various scan rates. (e) A schematic (top) illustrating the domain size reduction during a continuous cycling of lithiated MoS<sub>2</sub> nanosheets and charge-discharge plots (bottom) at different rates for the lithiated MoS<sub>2</sub>-based batteries. Separate panels are adapted with permission: (a) from ref. 103, copyright 2013, American Chemical Society; (b) from ref. 71, copyright 2019, Springer Nature; (c) from ref. 147, copyright 2018, Springer Nature; (d) from ref. 101, copyright 2015, Springer Nature; and (e) from ref. 200, copyright 2016, American Chemical Society.

The widely accepted descriptor for the HER rate is the Gibbs free energy  $\Delta G_{\text{H}}$  which characterises the adsorption of hydrogen on the active site.<sup>217</sup> A good catalyst for the HER binds hydrogen neither weakly nor strongly ( $\Delta G_{\text{H}}$  is close to 0). First-principles calculations have identified MoS<sub>2</sub> as a novel type of catalyst with a potential to replace Pt-group materials in electrocatalytic hydrogen production.<sup>218</sup> Although bulk MoS<sub>2</sub> is a poor catalyst, the catalytic activity of nanostructured MoS<sub>2</sub> is attributed to the distinct metallic states localised at the Mo-terminated (10–10) edges.<sup>17</sup> These one-dimensional states were first predicted by M. Bollinger *et al.*<sup>219</sup> and later visualised using scanning tunnelling microscopy<sup>219</sup> and scanning transmission X-ray microscopy.<sup>220</sup> The thermochemistry predictions were verified by T. Jaramillo *et al.* by measuring the catalytic activity of UHV sputtered MoS<sub>2</sub> nanoislands of various sizes.<sup>17</sup> The observed correlation between the rate of hydrogen evolution and the fraction of edge sites of MoS<sub>2</sub> nanoparticles unambiguously suggests that in the case of semiconducting 2H MoS<sub>2</sub> nanostructures, the electrocatalytic HER predominantly takes place at the edge sites.<sup>17</sup> Among 2H group VI TMDs, MoSe<sub>2</sub> is expected to demonstrate the highest catalytic activity since both (10–10) metal- and (–1010) chalcogen-terminated edges have  $\Delta G_{\text{H}}$  close to thermoneutral

(0.02 and –0.05 eV, respectively).<sup>221</sup> MoSe<sub>2</sub> is closely followed by WS<sub>2</sub>, both edges of which are potentially active for the catalytic HER (–0.04 and –0.05 eV), whereas only the Se edge of WSe<sub>2</sub> (–0.05 eV) and the Mo edge of MoS<sub>2</sub> (0.06 eV) were identified as catalytically active.<sup>221</sup> The basal planes of 2H group VI TMDs are found to be catalytically inert with  $\Delta G_{\text{H}}$  close to 2 eV.<sup>221</sup> Theoretically, higher affinity of basal planes of metallic 1T(1T') TMDs to hydrogen can lead to a remarkable improvement of their catalytic performance as compared to the semiconducting 2H counterparts.<sup>18,221,222</sup> Experimentally, the activation of the basal plane of the metallic 1T phase has been demonstrated by partially oxidising chemically exfoliated MoS<sub>2</sub> nanosheets.<sup>103</sup> D. Voiry *et al.* have compared the performance of 1T and 2H MoS<sub>2</sub> nanosheets before and after oxidative treatment. They reported that the catalytic activity of edge oxidised 2H MoS<sub>2</sub> nanosheets decreased considerably, while the performance of 1T phase remained unaffected. The activation of basal planes and so the higher density of catalytically active sites in metallic polymorphs is considered as one of the reasons why the 1T(1T') polymorphs outperform their 2H counterparts in catalytic HER.<sup>83,103,223</sup> This activity enhancement has been demonstrated for both group VI TMD nanosheets<sup>103</sup> and quantum dots<sup>112</sup> (Fig. 8a).



The enhanced catalytic activity of the metastable polymorphs of group VI TMDs is also associated with higher conductivity of the metallic (semimetallic) 1T(1T') phases. Recently, it has been suggested that the characteristics of the interface between active material and functional substrate, as well as the electron transport across active material, affect the catalytic performance even more than the thermodynamics of catalytic reaction. Compared to the semiconducting 2H counterparts, the 1T(1T') phases of WS<sub>2</sub>, MoS<sub>2</sub> and MoSe<sub>2</sub> greatly benefit from faster charge transfer in metallic polymorphs.<sup>102,199,224</sup> Additionally, in the case of metallic polymorphs, the catalyst–substrate interface becomes essentially metal–metal, which leads to more effective electron injection into the active material.<sup>225</sup> It is worth mentioning that the 1T' phase of WSe<sub>2</sub> have shown the enhanced electrocatalytic activity compared to its 2H counterpart (Fig. 8b); however, in this case the material was obtained uniquely *via* direct synthesis.<sup>71</sup>

The catalytic activity of metastable 1T and 1T' polymorphs of group VI TMDs cannot be compared directly as both phases are often reported interchangeably.<sup>83</sup> Further, L. Wang *et al.* have suggested that during electrochemical cycling, the 1T phase of MoS<sub>2</sub> irreversibly transforms into the more catalytically active 1T' phase.<sup>226</sup> According to first principle calculations, the 1T phase readily transforms into the 1T' phase at room temperature and this transition is considerably promoted by hydrogen adsorption during the H<sub>2</sub> evolution. Moreover, the authors have experimentally proved that the 1T-to-1T' transition does occur in chemically exfoliated MoS<sub>2</sub> and they suggest that such a transition may explain the self-optimising catalytic behaviour reported for the group VI TMD catalysts.<sup>226</sup> Additionally, there is not yet a clear evidence of the effect that a fully extended 1T' phase and large single crystalline domains of 1T' phase can have on the HER performance. For instance, the enhanced catalytic activity in mixed-phase 1T'/2H, as well as in polycrystalline single phase, TMD systems is occasionally attributed to the strained crystalline domain boundaries in the basal plane, providing a higher density of catalytically active sites.<sup>227,228</sup> Certainly, the bottom-up synthesis approach paves the way to expand the portfolio of 1T' phases and to control the crystallinity, morphology and crystal phase in multi-layered form (Fig. 8c).

**5.2. Increased conductivity for application in energy storage and sensors.** Due to the increased electrical conductivity of the metastable 1T and 1T' polymorphs compared to the semiconducting 2H counterparts, these materials are being extensively explored for energy storage applications. Thus, combined high electrical conductivity and large accessible surface area of metallic TMD nanosheets have been employed in electrochemical double-layer capacitors.<sup>101,229</sup> M. Acerce *et al.* have demonstrated that the electrodes formed from restacked chemically exfoliated 1T MoS<sub>2</sub> nanosheets can intercalate ions, such as H<sup>+</sup>, Li<sup>+</sup>, Na<sup>+</sup>, and K<sup>+</sup>, and can achieve volumetric capacitance of up to 650 F cm<sup>-3</sup> (at 20 mV s<sup>-1</sup>) in aqueous electrolytes.<sup>101</sup> Moreover, thus prepared electrodes are suitable for operation in non-aqueous electrolyte (ENIM BF<sub>4</sub> in MeCN) and demonstrate capacitance as high as 250 F cm<sup>-3</sup> (at 5 mV s<sup>-1</sup>) and good stability over 5000 cycles<sup>101</sup> (Fig. 8d).

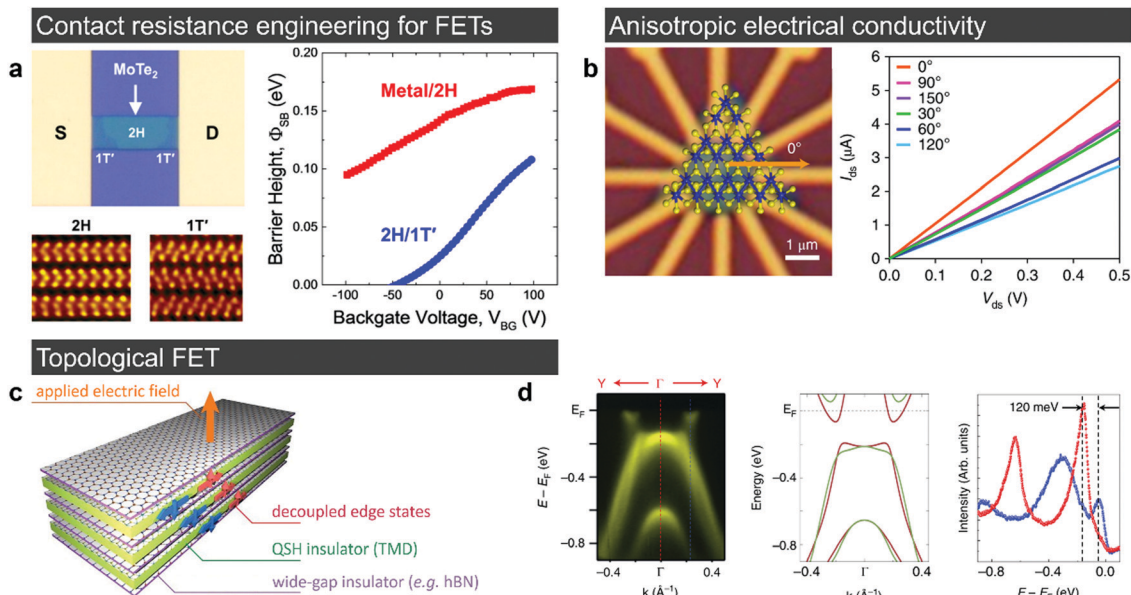
Grown by solution-phase approach material, exhibiting flower-like morphology and containing high fraction of metallic 1T phase, is gaining interest as a potential electrode material due to high specific surface area accessible to the electrolyte. R. Naz *et al.* have reported on highly defective MoS<sub>2</sub> nanosheets grown on reduced graphene oxide framework achieving high gravimetric capacitance of 442 F g<sup>-1</sup> (at 1 A g<sup>-1</sup>) in aqueous electrolyte (1 M H<sub>2</sub>SO<sub>4</sub>) with capacitance retention of over 84% during 3000 cycles.<sup>230</sup> Furthermore, 1T phase MoS<sub>2</sub> nanosheets grown through a magneto-hydrothermal reaction deliver a gravimetric capacitance of 379 F g<sup>-1</sup> (at 1 A g<sup>-1</sup>) in 1 M Li<sub>2</sub>SO<sub>4</sub> and show an exceptionally high stability (nearly 100%) over 10 000 cycles.<sup>231</sup>

Based on first-principle calculations, 1T MoS<sub>2</sub> monolayers have been identified as a suitable anode material for lithium-ion batteries due to their ability to absorb Li ions at high concentrations without aggregation.<sup>232</sup> T. Xiang *et al.* have realised hybrid 1T MoS<sub>2</sub>/graphene electrodes exhibiting a capacity of 666 mA h g<sup>-1</sup> at high current density (3500 mA g<sup>-1</sup>).<sup>233</sup> It is interesting to note that the arrays of vertically aligned few-layered MoS<sub>2</sub> nanosheets composing the electrodes were grown on graphene *via* a single-step hydrothermal reaction and the interlayer spacing in thus grown MoS<sub>2</sub> reaches 9.8 Å dictated by the graphene lattice; such lattice expansion facilitates a faster Li diffusion essential for a high performance lithium-ion battery. In bulk MoS<sub>2</sub> crystals, the repeating lithiation process is reported to form a mosaic-like nanocrystalline domains, which ensure high rate capability and excellent stability in rechargeable lithium-ion batteries.<sup>200</sup> This process is schematically shown in Fig. 8e. Such phase restructuring in pre-lithiated MoS<sub>2</sub> ensures a faster and more stable charge/discharge rate than in either chemically exfoliated or bulk material.<sup>200</sup> The 1T phases have seen recent utilization also in composites for batteries electrodes and we refer the reader to the following reviews for more details on this subject.<sup>234–236</sup> Covering these works is beyond the scope of the current review as here, we aim at discussing the applications that directly arise from the intrinsic characteristics of the metallic 1T phases of group VI TMDs.

The metastable 1T(1T') phases of group VI TMDs have been also proposed to be used in miniaturised biosensors. In particular, chemically exfoliated metallic 1T WS<sub>2</sub> has been employed as a conductive platform for an indirect enzymatic detection of fenitrothion; the biosensor expressed good linearity in the a broad concentration range (1–1000 nM) with the detection limit as low as 2.86 nM of fenitrothion.<sup>201</sup> On the other hand, the semimetallic 1T' MoTe<sub>2</sub> and WTe<sub>2</sub> polymorphs are reported to be ultra-sensitive surface enhanced Raman scattering (SERS) platforms for molecular detection.<sup>237</sup> In this case, the enhancement of Raman signal is achieved due to the polarisation of analyte molecules (rhodamine 6G) owing to the coupling with the 1T' phase SERS platform. The reported detection limits are as low as 40 and 400 fM of rhodamine 6G on pristine few-layered 1T' WTe<sub>2</sub> and 1T' MoTe<sub>2</sub> flakes, respectively; however, these can be further improved when the sensor is integrated on a Bragg reflector.

**5.3. Patterned 2H/1T' homojunctions for nanoelectronics.** Owing to their suitable bandgaps, atomically smooth interfaces,





**Fig. 9** Applications of the metastable 1T(1T') phases of group VI TMDs in novel electronics. (a) Optical image (left) of an FET device with patterned 1T'/2H MoTe<sub>2</sub> contacts and the Schottky barrier height (right) in the 2H MoTe<sub>2</sub> only device and in the device with patterned 1T'/2H contacts. (b) Electrical characterisation of a CVD grown 1T' MoS<sub>2</sub> flake (left) demonstrating an anisotropic electrical transport (right). (c) A schematic of a topological FET device based on van der Waals heterostructures of alternating monolayer 1T' group VI TMD and single-/few-layer wide-gap insulators. (d) Electronic characterisation of 1T' WSe<sub>2</sub> monolayer. High-resolution ARPES band dispersion (left) and the corresponding calculated bands (middle) demonstrate the band inversion in monolayer 1T' WSe<sub>2</sub>; an indirect band gap in 1T' WSe<sub>2</sub> calculated from two energy dispersion curves (right) is characteristic to 2D TI. Separate panels are adapted with permission: (a) from ref. 160, copyright 2019, American Chemical Society; (b) from ref. 147, copyright 2018, Springer Nature; (c) from ref. 46, copyright 2014, American Association for the Advancement of Science; and (d) from ref. 20, copyright 2018, Springer Nature.

and good thermal stability, semiconducting group VI TMDs, and MoS<sub>2</sub> single layers in particular, are being investigated for the next generation electronics and optoelectronics.<sup>238–241</sup> Recently, a field-effect transistor (FET) has been experimentally demonstrated with a single mechanically exfoliated MoS<sub>2</sub> flake as a conductive channel.<sup>238</sup> The device demonstrated an on/off ratio of 10<sup>8</sup> at room temperature and high channel mobilities of at least 200 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> were achieved by dielectric screening with HfO<sub>2</sub> as a gate insulator.<sup>238</sup> The electronics based on atomically thin TMDs is majorly plagued by the large contact resistance between the semiconductor and the source/drain metal electrodes.<sup>242</sup> This can be solved by locally inducing the 2H-to-1T' transformation on TMD nanosheets to significantly lower the contact resistance between the monolayer and bulk metal contacts (Fig. 9a). Local phase patterning to realise ohmic contacts has been demonstrated in MoS<sub>2</sub> and MoTe<sub>2</sub> transistors.<sup>95,155</sup> The 1T phase electrodes were locally patterned on CVD grown 2H MoS<sub>2</sub> nanosheets by exposing the material to *n*-BuLi, while the semiconducting channel was preserved by protecting this area by a PMMA mask.<sup>95</sup> In 1–3 layered MoS<sub>2</sub> nanosheets, the heterophase 1T'/2H homojunction is characterised by a contact resistance of 200–300 Ω μm at zero gate bias. Thus fabricated MoS<sub>2</sub> FET exhibits the channel mobilities of ~50 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> and on/off ratios exceeding 10<sup>7</sup>.<sup>95</sup> Similarly, the 1T'/2H homojunctions induced in MoTe<sub>2</sub> by local laser irradiation have increased the channel mobility by a factor of 50, while preserving the high on/off ratios of ~10<sup>6</sup>.<sup>155</sup> Interestingly, in MoTe<sub>2</sub>, lateral 1T'/2H homojunctions can be obtained *via* direct CVD synthesis.<sup>160,243,244</sup> The contact barrier height

at the lateral 1T'/2H interface is reported to be as low as 25–30 meV,<sup>160,243</sup> and in the coplanar 1T'/2H contact MoTe<sub>2</sub> FET, the channel mobility reached 16.2 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> at room temperature.<sup>243</sup>

**5.4. Dynamic phase transitions for non-volatile memory devices.** Fast and reversible structural transitions between the crystal phases characterised by a large difference in the electrical resistance between them, lay a foundation for non-volatile memory devices and neuromorphic computing.<sup>245</sup> F. Zhang *et al.* have recently demonstrated that the dynamic switching between high resistive (2H) and low resistive (1T') states in few-layered MoTe<sub>2</sub> and W<sub>x</sub>Mo<sub>1-x</sub>Te<sub>2</sub> can be controlled by a vertical electric field.<sup>43</sup> Remarkably, the switching in such resistive random access memory (RRAM) device is achieved within 10 ns and for the Al<sub>2</sub>O<sub>3</sub>/MoTe<sub>2</sub> stacked configurations, the on/off current ratios are as high as 10<sup>6</sup>.<sup>43</sup> In MoS<sub>2</sub> films, a reversible 2H-to-1T' phase switching controlled by the in-plane Li ion migration has also been demonstrated.<sup>109</sup> A reversible local phase transformation caused by the difference in Li ion concentration intercalated between MoS<sub>2</sub> layers was effectively modulated by an applied in-plane electric field. Furthermore, it has been shown that the individual devices can be coupled through the local ionic exchange, illustrating their potential for neuromorphic computing. Finally, the material produced in solution phase, such as colloidal MoS<sub>2</sub> nanosheets, can be assembled *via* spray coating, enabling a wafer-scale fabrication of RRAM arrays on flexible substrates for integration in next generation mobile and wearable electronic devices.<sup>176</sup>





**5.5. Emerging properties and applications.** In the semi-conducting 2H polymorph, superconductivity can be induced by intercalation of alkali metals<sup>246</sup> as was demonstrated in the case of MoS<sub>2</sub> with the critical temperatures  $T_c$  varying from 1 to 6.5 K.<sup>8,247</sup> Typically, the fraction  $x$  of alkali metal (A = Li, Na, K, Rb) in these intercalation compounds A<sub>x</sub>MoS<sub>2</sub> was as high as ~0.4.<sup>247</sup> Later, the alkali metal intercalation has been shown to cause the transition into superconducting state in both the 1T and 1T' bulk polymorphs of MoS<sub>2</sub> with the  $T_c$  of 2.8 and 4.6 K, respectively.<sup>248</sup> The presence of alkali metal in the intercalated bulk powders was found crucial for the superconducting phase transition.<sup>248</sup> On the other hand, pristine 1T' MoTe<sub>2</sub> is found to undergo a superconducting phase transition at 0.1 K and the transition temperature can be dramatically increased to 8.2 K under applied pressure (11.8 GPa).<sup>63</sup> Similarly, pressure-driven dome-shaped superconductivity has been reported in bulk WTe<sub>2</sub> with the highest  $T_c$  (7 K) achieved at 16.8 GPa.<sup>249</sup> Further, dome-shaped superconducting behaviour was observed in pressurised 2M WSe<sub>2</sub> crystals with the resistance drop onset at 4.2 GPa with the  $T_c$  of 4.3 K and the  $T_c$  maximum (7.3 K) at 10.7 GPa.<sup>72</sup> The development of a facile synthesis route to produce bulk 1T(1T') phases of group VI TMDs *via* gentle deintercalation of A<sub>x</sub>MoS<sub>2</sub> intercalation compounds has presented an opportunity to study the superconducting phase transition in pure 1T TMD polymorph. Thus, 1T MoS<sub>2</sub> is reported to be an intrinsic superconductor at temperatures below 4 K.<sup>143</sup>

Due to the structural anisotropy in the *ab* plane, 1T' single layers have been predicted to exhibit anisotropic transport characteristics. Anisotropic electrical conductivity of the 1T' phase have been recently demonstrated in chemically exfoliated<sup>250</sup> and CVD grown 1T' MoS<sub>2</sub> flakes<sup>147</sup> (Fig. 9b). The conductance reaches its maximum along the *a* (zigzag) direction and the anisotropic ratio is as high as 2,<sup>147</sup> that exceeds the reported anisotropic conductance for black phosphorus.<sup>251</sup> This leads to envision application of the 1T' phases for neuromorphic devices where anisotropic transport is sought in different device components. Black phosphorous, a highly anisotropic 2D material, is being investigated in synaptic devices and higher devices complexities could be enabled by 2D materials with different anisotropic transport.

Bulk nonmagnetic Td WTe<sub>2</sub> demonstrates an extremely large magnetoresistance (XMR); the primary source of which is established to originate from the carrier compensation.<sup>252</sup> Thus, at 4.5 K, positive magnetoresistance of 452 700% was observed in Td WTe<sub>2</sub> in a magnetic field of 14.7 T; moreover, the magnetoresistance value reaches 13 000 000% at 0.53 K in a magnetic field of 60 T and no saturation of magnetoresistance is evidenced even in very high magnetic fields.<sup>253</sup> Later, the XMR behaviour was discovered in the Td phase of MoTe<sub>2</sub>.<sup>254,255</sup> Furthermore, the magnetoresistance observed in Td MoTe<sub>2</sub> in high magnetic fields can be tuned by an applied tensile strain (0.5%).<sup>256</sup> At low temperatures and in high magnetic fields, the magnetoresistance variation approaches 30% under uniaxial tensile strain perpendicular to the zigzag chain direction, while the strain along the zigzag chain direction effectively suppresses the magnetoresistance.<sup>256</sup>

## Conclusions and future outlook

In summary, in the past few years there have been exciting fast developments in the elucidation of the potential of the metastable phases of TMDs, ranging from electrocatalysis<sup>83,103</sup> to condensed matter phenomena such as 2D topological insulators.<sup>19</sup> Here, we have reviewed the concomitant significant step towards the production of these phases in measurable quantities and the key applications that these phases can uniquely enable. The most widely used method to synthesise metastable phases is based on the intercalation with alkali metals in organic solvent<sup>83,114</sup> or in a dry form.<sup>50,103</sup> This is a versatile route which can be applied to bulk powders, with an associated exfoliation of the flakes, or on individual flakes, often CVD grown, on a substrate. Thus, it provides with the ability to produce highly concentrate suspensions of the 1T(1T') phases in water as well as monolayers extended over large areas. The latter is compatible with the integration in planar devices such as field-effect transistors.<sup>95,155</sup> The first provides stable colloidal suspensions of the 1T phases in water which can be processed from solution *via* vacuum filtration<sup>103</sup> or dip-coating,<sup>257</sup> to generate uniform thin films for different applications from electrocatalysis<sup>103</sup> and photoassisted electrocatalysis,<sup>257</sup> to energy storage devices.<sup>101</sup>

In the form of thin films, the 1T phase can be converted into the 2H phase upon annealing.<sup>27,50</sup> Thus, the accessibility of the 2H phase from the 1T(1T') phase is advantageous as it provides an alternative approach to the commonly used exfoliation of the 2H phase in organic solvent by applying a shear force,<sup>258</sup> allowing for the processability of water-based suspensions which are compatible with large-scale production and deposition techniques, such as ink-jet printing,<sup>259</sup> 3D printing,<sup>260,261</sup> and spray coating,<sup>262</sup> to name a few.

While chemical treatment has intrinsic advantages and it has been very effective for initial studies of 1T phases, as the field progresses, the limitations are now emerging. The first one is represented by the use of Li-chemistry which is air sensitive, involving the use of a controlled atmosphere. Although the intercalation of MoS<sub>2</sub> and WS<sub>2</sub> bulk powders have demonstrated the 2H-to-1T(1T') phase transformation, the conversion of the corresponding diselenides (MoSe<sub>2</sub> and WSe<sub>2</sub>) have not been reported. While the 1T' phases of MoSe<sub>2</sub> and WSe<sub>2</sub> have been obtained *via* direct synthesis, demonstrating that these metastable phases can actually be achieved.<sup>68,71</sup>

Furthermore, the 1T' phase resulting from the Li intercalation is limited to the 75% of the material and it has a patchy distribution across the basal plane of individual flakes with a polycrystalline nature, where individual domains can be limited to very few nm.<sup>85</sup> The impact of the polycrystallinity on the electrical, mechanical and electrocatalytic properties is still not known yet. Thus, the search for bottom-up approaches to the synthesis arises from the need of fully controlling the phase morphology, the materials processability and the possibility to establish a wider spectrum of metastable phases and compounds composition.

Progress has been made in the direct synthesis of metastable phases *via* both vapour-based deposition techniques and



in liquid phase from molecular precursors which led to the first demonstrations of the 1T' phase of MoS<sub>2</sub><sup>147</sup> and WSe<sub>2</sub>.<sup>71</sup> Most of the reported vapour-phase syntheses are based on chemical vapour transport and have demonstrated the synthesis of bulk crystals of metastable 1T' MoSe<sub>2</sub><sup>68</sup> and 1T' MoS<sub>2</sub> phases.<sup>68,69</sup> The achievement of the 1T' phase monolayers *via* CVD approach has been limited to MoTe<sub>2</sub>, WTe<sub>2</sub>,<sup>161</sup> and MoS<sub>2</sub>.<sup>147</sup> These bottom-up approaches lead to nearly continuous 1T' phase across large areas (over tens of microns) of over 80% and exhibit a single crystalline nature across hundreds of nanometres. As the atomically thin nature of the flakes is not preserved with CVT synthesis, when bulk crystals are produced, they need to undergo a liquid phase exfoliation step to be used in their 2D form. CVD, which can enable scalability of the synthesis, have the advantage of producing anisotropic monolayers over large areas for application in catalysis, nano-electronics, neuromorphic systems, quantum technologies and advanced spectroscopies; however, altering the reaction energetics to selectively produce the metastable phase is still challenging.

MBE have been used to demonstrate single crystalline phase of 1T' WSe<sub>2</sub><sup>20,21</sup> on bilayer graphene terminated SiC substrates, and of 1T' MoSe<sub>2</sub><sup>70</sup> and 1T' MoS<sub>2</sub><sup>79</sup> on Au substrates. Although the 1T' domains are small (less than 100 nm), these are sufficiently extended to demonstrate new condensed-matter phenomena so far just predicted, thanks to the single crystalline nature and of the 1T' phase.

Direct synthesis in solution is the most recent synthesis method demonstrated and it has the advantage of involving low temperatures (<300 °C), high percentage of the 1T' phase in the final product and single crystalline nature across flakes sizes of 100 nm.<sup>71,80,173</sup> Synthesised through a direct colloidal reaction, the nanosheets can form stable dispersions in non-polar organic solvents. This synthesis route presents features compatible with a large-scale production and the possibility of translating the chemistry of the synthesis into a hydrothermal process.

Branched flower-like nanostructures with atomically thin petals are generally obtained *via* liquid phase synthesis.<sup>71,81,173,199</sup> Such morphology is very advantageous for catalysis and supercapacitors as it maximizes the exposure of active sites *versus* a 2D nanosheets geometry where flakes are planarly stacked, losing some of the accessible surface area. These nanoflowers can also be directly grown on the electrode materials where they can be strongly anchored enabling a good electrical contact. The possibility of obtaining individualized planar 2D nanosheets could, however, expand the applications of the 1T' phases produced in liquid phase.

The next key challenge lies in understanding the nucleation stage in solution phase to be able to reliably control the flake morphology (between branched sheets and individual sheets) and the crystal phase. As stated by the Ostwald's step rule, which has been based on empirical observations in zeolite materials, the crystallization from a solution occurs in steps in such a way that often thermodynamically unstable phases occur first, followed by a transformation in the

thermodynamically stable phase.<sup>169</sup> Thus, more theoretical understanding and experimental studies of the early stages of nucleation of TMDs are needed to predict the course of the growth and thus to design the ways of controlling morphology and to lock the crystal phase. The ability to produce individualized flakes in solution with micrometric lateral sizes can expand the potential applications competing with CVD methods. Accessing pure 1T phases of single crystalline nature across large areas of TMDs is necessary to study the intrinsic physical and chemical properties and so far there is a clear indication that only direct synthesis methods enable to approach 100% of metastable phase with large single crystals. Also, to fully understand the crystallographic structure of the most exotic metastable 1T' and 1T'' phases and their related properties is necessary to have high-quality single crystals.

For practical applications, the metastable phases should be stabilized in the absence of charged species, *i.e.* precursor residues or intercalated cations, to produce pristine, kinetically stable materials. To this goal, the Ostwald's step rule could be again useful to predict how to preserve a metastable phase beyond the nucleation point.

The recent evidence that the 1T TMD<sub>s</sub> can also be obtained *via* chemical transformation of non-van der Waals materials, further expands the potential routes of the synthetic approaches possibly making them numerous.<sup>152</sup> Thus, these new synthesis strategies should pave the way to a new class of atomically thin materials which present new physicochemical features promising in electrocatalytic and energy storage applications as well as new fundamental condensed matter phenomena which can enable new technologies. Recent demonstration of 2D topological insulating properties in the 1T' TMD monolayers could lead to the fabrication of topological field-effect transistors operating at room temperature<sup>46</sup> (Fig. 9c and d). The anisotropic transport in 1T' phases dictated by the anisotropic crystal structure could see applications in artificial synaptic devices to emulate biological systems.

The investigation on the exotic physical properties related to the most rarely observed or predicted metastable crystal phases of group VI TMDs (1T' and 1T'') is progressing very rapidly and there is an incredible potential to exploit these highly anisotropic but stable phases in future nanoelectronics, optoelectronic systems and quantum technology.

## Conflicts of interest

The authors declare no conflict of interest.

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