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Performance of the double-Wien filter of the Neoma MC-ICPMS/MS with an application to copper stable isotope compositions

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# Performance of the double-Wien filter of the Neoma MC-ICPMS/MS with an application to copper stable isotope compositions†

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The new Neoma MC-ICPMS/MS is equipped with a prefiltering system consisting of a double-Wien filter and a collision/reaction cell whose performances are challenged using different combinations of magnetic and electrostatic field values and adjustable slit apertures. The results show an asymmetrical attenuation of transmission relative to the chosen axial mass-to-charge value, with higher efficiency at removing low masses than high masses, even when magnetic induction is minimal. The resulting asymmetry of the bandpass window is fully predictable by theoretical calculations of ion trajectories in a Wien filter, either as a function of the magnetic field value or that of the aperture of the adjustable slit. With an axial mass-to-charge value set at  $^{120}\text{Sn}$ , the vertical deviation for a magnetic field value at 100% will be approximately 3 mm and 4 mm for a variation of mass-to-charge of  $\pm 20\%$ , respectively. In these conditions, Ar was already barely detectable with the lowest magnetic field value (10%) and a fully open adjustable slit, while Pb is quantitatively transmitted. We then use the prefiltering system to remove on-line the  $^{40}\text{Ar}^{23}\text{Na}^+$  compound that produces an isobaric interference with  $^{63}\text{Cu}$ , hampering high-precision measurement of Cu stable isotope composition ( $^{65}\text{Cu}/^{63}\text{Cu}$ ). While Na is not transmitted thanks to the double-Wien filter,  $^{40}\text{Ar}^{23}\text{Na}^+$  still interferes with  $^{63}\text{Cu}$ , demonstrating that this argide is produced in the plasma source but not in the reaction cell. Helium in the collision/reaction cell is necessary to remove the  $^{40}\text{Ar}^{23}\text{Na}^+$  interference. The MS/MS technology of the Neoma allows for the correction of the  $^{40}\text{Ar}^{23}\text{Na}^+$  interference up to a Na/Cu ratio of 10, where other classic MC-ICPMS already show an offset of the Cu stable isotope composition for a Na/Cu ratio of 1. The inability to correct the Cu stable isotope composition with a Na/Cu ratio higher than 10 suggests that the Na-based interference is no longer spectral and becomes linked to the matrix. We next measure the Cu stable isotope composition in eight certified reference materials prepared with a simple automated single step ion-chromatography procedure to purify Cu. The results show a very good agreement with previously reported values. The overall results suggest that the MS/MS technology of the Neoma MC-ICPMS allows efficient on-line isolation of analytes, therefore reducing potential spectral and matrix interferences to permit much better resolved and controlled subsequent effects in the collision/reaction cell.

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## Introduction

Copper has only two stable isotopes ( $^{63}\text{Cu}$  and  $^{65}\text{Cu}$ ), whose relative abundances ( $^{63}\text{Cu} = 69.17\%$  and  $^{65}\text{Cu} = 30.83\%$ ) have been determined in the 1960's by thermal ionization mass spectrometry (TIMS).<sup>1</sup> Isotopic variations were associated with poor analytical uncertainty ( $2\%$  per amu) because ionization yields of Cu were too low with the TIMS technique. It was not until the late 1990's with the arrival of the first commercialized multi-collector inductively coupled plasma mass spectrometry

(MC-ICPMS), at the time the Plasma 54 (VG Elemental, now ThermoFisher Scientific, Bremen), that it was possible to achieve precise ( $0.02\%$  per amu) Cu isotope ratio measurements.<sup>2</sup> The pioneering work of Maréchal *et al.*<sup>2</sup> showed that the instrumental mass bias of MC-ICPMS can be corrected by a combination of standard/sample bracketing and elemental doping with Zn, which was further modified by using Ni<sup>3,4</sup> or Ga.<sup>5,6</sup> The overall consistency of the measurements of Cu isotope compositions by MC-ICPMS leads to the rapid development of applications in various fields such as cosmochemistry,<sup>7,8</sup> igneous,<sup>9,10</sup> ore,<sup>11,12</sup> sediment<sup>13,14</sup> and river<sup>15,16</sup> geochemistry, oceanography,<sup>17,18</sup> tracing atmospheric<sup>19–21</sup> and soil<sup>22–24</sup> pollution, assessing the diagnosis and prognosis of metabolic<sup>25–27</sup> and neurodegenerative<sup>28–30</sup> diseases, cancer,<sup>31–33</sup> but also in palaeoanthropology<sup>34–36</sup> and archaeology.<sup>37–39</sup>

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Prior to isotopic analysis, Cu needs to be separated from the matrix and further purified using ion-exchange chromatography. Many protocols exist,<sup>3,40–42</sup> generally based on the use of the strongly basic AG MP-1 anion exchange resin (100–200 mesh, chloride form). These protocols are usually effective for the purification of Cu, but extreme care must be taken for the removal of Na, because an argide  $^{40}\text{Ar}^{23}\text{Na}^+$  compound interferes with the  $^{63}\text{Cu}$  isobar. Removing Na can be done quite easily for silicate materials, which generally exhibit a Na/Cu ratio ranging from  $10^2$  to  $10^3$ , but can be extremely challenging for biological fluids (urine, Na/Cu  $\sim 8 \times 10^4$ ; plasma, Na/Cu  $\sim 3 \times 10^6$ ) and above all, (seawater, Na/Cu  $\sim 5 \times 10^7$ ). For the extreme case of seawater, large volumes of eluent for Na separation and Cu purification are necessary for the handling of which automatic procedures have been recently developed.<sup>40,42</sup> Contrasting with these off-line procedures, an alternative on-line procedure now exists thanks to the MS/MS technology of the Neoma MC-ICPMS. The MS/MS technology consists of a precell mass filter and a collision cell.<sup>43</sup> Initially, the precell mass filter was a bespoke quadrupole, which allowed a prototype Proteus to measure several isotope systems.<sup>44–46</sup> Lack of sensitivity and non-reproducible mass bias prevented the Proteus from being widely adopted by the MC-ICPMS community, and the low-energy quadrupole was replaced by a high-energy double Wien filter on the Vienna prototype.<sup>47,48</sup>

The double-focusing electrostatic analyzer and magnetic sector geometry is standard and inspired from that of the Neptune MC-ICPMS, while the collection system has been totally redesigned as described in Dauphas *et al.*<sup>49</sup> and Télouk *et al.*<sup>50</sup>

In this study, we first explore the performance of the double-Wien filter to on-line prefilter ions with different  $m/z$  using various combinations of  $B$  and  $S$  values. Second, we compare the performances of the Neoma MC-ICPMS/MS to those of standard MC-ICPMS and third, we measure the  $^{65}\text{Cu}/^{63}\text{Cu}$  ratio in certified reference materials that have experienced a single purification step to validate the overall procedure.

## Experimental

### Reagents and materials

All experiments were carried out in laminar flow hoods in a clean laboratory at the LGL-TPE (Ecole Normale Supérieure de Lyon). Acids ( $\text{HNO}_3$ ,  $\text{HCl}$ , and  $\text{HF}$ ) were double distilled to reduce blank contaminations. Ultrapure water (resistivity  $>18.2$  M $\Omega$  cm) was obtained from a Milli-Q Element water purification system (Merck Millipore, Bedford, MA, USA). The performances of the double-Wien prefiltering was assessed using the multi-element standard solution SCP33MS containing 33 elements (SCP Sciences, Québec, Canada) diluted to 200 ng mL<sup>-1</sup>. Synthetic solutions were prepared by diluting Na and Zn Specpure solutions (Alfa Aesar, Karlsruhe, Germany) and Cu SRM-976 solution (National Institute of Standards and Technology, Gaithersburg, MD, USA) to reach a final concentration of 200 ng mL<sup>-1</sup> for Cu and Zn and various Na/Cu ratios. Eight geological certified reference materials (CRM) were obtained from the United States Geological Survey (USGS) and include the BHVO-1

Hawaiian basalt,<sup>51</sup> Glass Mountain rhyolite RGM-1,<sup>51</sup> the Columbia River basalt BCR-1,<sup>52</sup> the AGV-2 Guano Valley andesite,<sup>52</sup> the BIR-1 Icelandic basalt,<sup>53</sup> the DNC-1 North Carolina dolerite,<sup>53</sup> the Centreville diabase W2a,<sup>53</sup> and the Japanese alkali basalt JB1-a<sup>54</sup> from the Geological Survey of Japan (GSJ). A minimum sample size of 100 mg was weighed for CRM to avoid measurement uncertainties due to the heterogeneity of the reference material powder. Certified reference materials were digested with a mixture of 5 mL of 27 M distilled HF and 2.5 mL of 15 M distilled  $\text{HNO}_3$  at 120 °C for 12 hours and evaporated to dryness. Fluorides were redissolved using 2 mL of 6 M HCl and heated on a hotplate at 100 °C for 12 hours and then evaporated to dryness. The chemical separation of Cu was achieved using the automated chromatography system prepFAST-MC (Elemental Scientific, Omaha, USA) loaded with 500  $\mu\text{L}$  of Cu resin (Triskem, Rennes, France) following the procedure of Enge *et al.*<sup>55</sup>

### Instrumentation

The MS/MS equipment consists of a double-Wien filter and a collision/reaction cell.<sup>47</sup> The Wien filter is an important device in charged particle optics because it is a static field mass filter that deflects charged particles according to their velocity,<sup>56</sup> or equivalently their mass-to-charge ratio ( $m/z$ ). While the Wien filter properties can also be used in electron microscopy,<sup>57</sup> we will consider here the case where charged particles are ions. The Wien filter consists of a uniform electrostatic ( $\vec{E}$ ) field orthogonal to a magnetic ( $\vec{B}$ ) field. Charged particles passing through this arrangement of fields are subject to an electrostatic force ( $\vec{F}_E = q\vec{E}$ ) and a magnetic force ( $\vec{F}_B = q\vec{v} \times \vec{B}$ ), where  $q$  and  $\vec{v}$  are the charge and the velocity, respectively, of the ion. For a collimated ion beam perpendicular to both fields, the axial transmission of an ion  $m_0$  with a velocity  $v_0$  will happen when  $v_0 = |\vec{E}/\vec{B}|$  (or  $v_0 = E/B$ ). For ions with the same kinetic energy and  $m < m_0$ ,  $v > E/B$  because lighter ions travel faster, this ion beam will be deflected in the opposite direction of the  $\vec{E}$  field (Fig. 1). Inversely, for ions with  $m > m_0$ , the ion beam will be deflected in the direction of the  $\vec{E}$  field because  $v < E/B$  (Fig. 1). The mass prefiltering system of the Neoma MC-ICPMS/MS is a double-Wien filter, with the first filter deflecting ions away from the axial trajectory and the second filter refocusing ions back towards the axial trajectory, in both cases according to their ion velocity. Both Wien filters are equipped with an array of upper and lower baffles to avoid reflection on the side walls of ions with a too great angle of deflection and to avoid charging of the electrodes.<sup>47</sup> After being vertically deflected apart from the axial trajectory, those ions with a too large  $m/z$  difference relative to  $m_0/z$  are not selected using an adjustable slit located between the two Wien filters. A combination of  $E$ ,  $B$ , and adjustable slit aperture ( $S$ ) values thus permits the transmission of a certain mass range through the inversion lens (Fig. 1). The inversion lens is an Einzel lens that inverts the divergent trajectories of ion beams with  $m/z$  different from  $m_0/z$  towards the axial position, in inverse proportion to the degree of deviation introduced by the first Wien filter.<sup>46</sup> The inversion happens without altering energy, and the different ions enter the second Wien filter as an

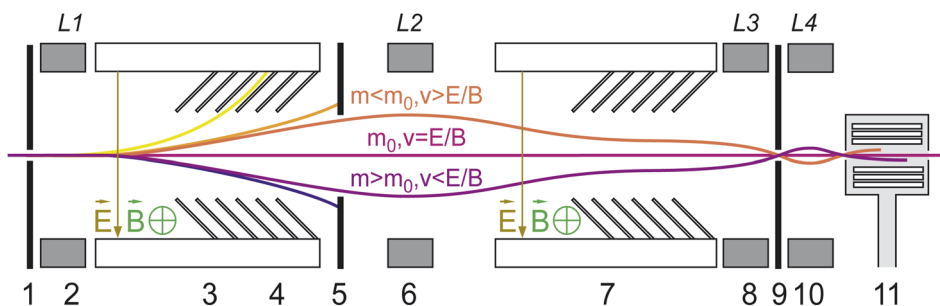


Fig. 1 Schematic overview of the ion path in the Neoma prefiltering MS/MS system. (1) Entrance aperture; (2), first lens (L1) to focus; (3), first Wien filter; (4), baffles; (5), adjustable slit; (6) second lens (L2) to inverse; (7), second Wien filter; (8), third lens (L3) to focus; (9), exit aperture; (10), fourth lens (L4) to focus; (11), collision/reaction cell.

uncollimated beam. The two Wien filters are of similar geometry, and their  $E$  and  $B$  fields are identical and controlled by a unique power supply. Thus, in the second Wien filter, ions with  $m < m_0$  will be deflected in the opposite direction of the  $\vec{E}$  field (and inversely for ion  $m > m_0$ ), such that the ion beam becomes collimated again before the exit aperture. The mass prefilter contains four lenses that need to be tuned accordingly (Fig. 1). The first lens (L1) focuses the ion beams through the adjustable slit, L2 (the inversion Einzel lens) bends the ion beams back toward the second Wien filter, the third lens (L3) focuses the ion beam into the exit aperture and the fourth lens (L4) focuses the ion beams to the collision/reaction cell (CRC), where they are kept focused using an hexapole despite scattering due to collisions with gas (Fig. 1). The CRC allows the introduction of four different reactive ( $O_2$ ,  $NH_3$ ,  $H_2$ ) and non-reactive (He) gases. The interfering species, which are polyatomic or molecular ions, collide with the cell gas more frequently than analyte ions because they have a larger collisional cross section. The mechanism for the elimination of the pre-existing polyatomic ions in collision with a non-reactive gas is the selective energy loss displayed by polyatomic ions compared with monoatomic ions. By applying a small, fixed bias voltage at the cell exit, it is possible to keep the polyatomic ions and newly formed species in the CRC while monoatomic ions have sufficient kinetic energy to exit.<sup>58</sup>

The Cu isotopic compositions were measured using the Nu Plasma MC-ICPMS (Nu Instruments, Wrexham, UK) following the original procedure described by Maréchal *et al.*<sup>2</sup> or using the Neoma MC-ICPMS/MS. The instrument parameters are summarized in Table 1. On the day of analysis, Cu-purified solutions were diluted in a Zn-doped solution (Zn JMC 3-0749L, Johnson Matthey Royston, UK) to match the concentration of the standard bracketing solution ( $200 \text{ ng mL}^{-1}$ ). Instrumental mass bias and temporal drift were corrected with an exponential law using Zn as an internal standard, combined with sample-standard bracketing, as recommended by Maréchal *et al.*<sup>2</sup> All the results of isotopic measurements are given in the delta notation (expressed in ‰) and reported relative to the international isotopic standard solutions SRM-976 (National Institute of Standards and Technology, Gaithersburg, MD, USA) using:

$$\delta^{65}\text{Cu} = \left[ \frac{(^{65}\text{Cu}/^{63}\text{Cu})_{\text{sample}}}{(^{65}\text{Cu}/^{63}\text{Cu})_{\text{standard}}} - 1 \right] \times 1000$$

We use the ‘33 Elements’ solution at  $200 \text{ ng mL}^{-1}$  and mass scans ranging between 20 to 250 atomic mass units (amu) with an axial value set at the middle of the mass range, *i.e.*, at  $^{120}\text{Sn}$  and then at  $^{63}\text{Cu}$ . Closing the slit necessitates to tune the  $E$  value

Table 1 Instrument settings and data acquisition parameters for MC-ICP-MS(/MS) analyses

	Neoma MS/MS	Nu Plasma
RF power (W)	1200	1350
Plasma condition	Wet, quartz cyclonic/Scott double spray chamber	Wet, cyclonic spray chamber
Nebulizer Ar flow ( $\text{L min}^{-1}$ )	0.9	1
Mass resolution	2250	400
Sampling cone	standard	Ni wet
Skimmer cone	H	Ni wet
Cup configuration	H5: $^{71}\text{Ga}$ ; H4: $^{70}\text{Zn}$ ; H3: $^{69}\text{Ga}$ ; H2: $^{68}\text{Zn}$ ; H1: $^{67}\text{Zn}$ ; Ax: $^{66}\text{Zn}$ ; L1: $^{65}\text{Cu}$ ; L2: $^{64}\text{Zn}$ ; L3: $^{63}\text{Cu}$ ; L4: $^{62}\text{Ni}$ ; L5: $^{61}\text{Ni}$	H4: $^{68}\text{Zn}$ ; H3: 67.5; H2: $^{67}\text{Zn}$ ; Ax: $^{66}\text{Zn}$ ; L1: 65.5; L2: $^{65}\text{Cu}$ ; L3: $^{64}\text{Zn}$ ; L4: $^{63}\text{Cu}$ ; L5 $^{62}\text{Ni}$
Sensitivity ( $\text{V ppm}^{-1}$ )	$\sim 65$	$\sim 20$
Blank signal (2% $\text{HNO}_3$ )	$^{63}\text{Cu} \sim 12 \text{ mV}$	$^{63}\text{Cu} \sim 5 \text{ mV}$
Integration time (s)	4.194	10
Cycles	40	30

by few volts. All statistical analyses was performed using the R software.<sup>59</sup>

## Results and discussion

### Prefiltering performances

We first obtain a scan with the adjustable slit fully open ( $S = 100\%$ ) and minimal magnetic field ( $B = 10\%$ , Fig. 2A) for a  $m_0$  set at 120 amu ( $^{120}\text{Sn}$ ). This configuration represents the basic instrumental conditions that the MS/MS will alter either with varying magnetic field, or closing the adjustable slit, or a combination thereof. It is noteworthy that these basic instrumental conditions already influence the transmission of low masses at about 100 amu away from  $m_0$ . Indeed, the minimal  $B$  value totally suppresses the signal intensities of elements such as Na, Mg, and S (Fig. 2A), and, of interest, also Ar, whose signal is only  $\sim 0.3$  V for  $^{40}\text{Ar}$  and  $\sim 1$  V for  $^{40}\text{ArH}$ . The prefiltering effects are asymmetrical, as high masses are transmitted efficiently (Fig. 2A). We next explore the effect of augmenting the magnetic field with an increment of ten percent of  $B$ , while keeping a constant  $E/B$  ratio and the adjustable slit fully open. Fig. 2B shows the difference in signal intensities between two consecutive  $B$  increments of ten percent. The transmission of the elements in the range of 40 to 70 amu is completely eliminated when  $B$  is increased to 20%, and the transmission of the elements in the range of 70 to 90 amu is completely eliminated when  $B$  is increased to 30% (Fig. 2B). Lead is no longer transmitted for  $B = 40\%$ . Slightly increasing  $B$  to 20% or 30% not only decreases the transmission of elements with low and high masses, but also increases that of elements with intermediate masses, *i.e.*, close to  $m_0 = 120$  amu (Fig. 2B). This can be explained by an

enhanced transmission of elements with intermediate masses that is made easier by the removal of  $\text{Ar}^+$  ions that produce significant space charge with a defocusing effect. Increasing the  $B$  value up to 40% reduces the signal intensities by about 80% for elements being ten amu lighter than the 120  $m/z$ . However, when  $B$  is set  $\geq 70\%$ , the transmission likely becomes unstable as some masses slightly higher than the 120  $m/z$  exhibit both increased and decreased voltages (Fig. 2B). A close-up from 90 to 150  $m/z$  is given in Fig. S1.† Observed enhanced signal intensities can be significant ( $\sim 2$  V) but remain lower than generalized reduced signal intensities ( $\sim -8$  V). Measuring both enhanced and reduced transmission for a given mass is obviously an artifact. A high measurement rate during scan acquisitions (about 20 measurements by amu) produces a lot of data, which can yield positive or negative differences when subtracted from two consecutive scans. Overall, this indicates that the instrument is suffering from unstable conditions at high  $B$  values. The performance of the MS/MS prefilter at high  $B$  values requires further investigation, but we anticipate that setting such high  $B$  values to remove elements will instead be preferentially achieved by closing the adjustable slit.

We next study the prefiltering effects for  $B$  values set at 10%, 30%, or 50%,  $E/B$  being constant, and with the adjustable slit fully open, or 50% or 70% closed ( $S = 100\%$ , 50% and 30%, respectively). The resulting scans are given in Fig. 3. Closing the adjustable slit when  $B$  is minimal produces very asymmetrical effects, with high masses not being affected until the slit is 50% closed, while low masses are filtered efficiently. The magnetic field must be set at half its maximum value, with a fully open adjustable slit to obtain a more symmetrical filtering. This combination ( $B = 50\%$ ,  $S = 100\%$ ) produces a shape of the

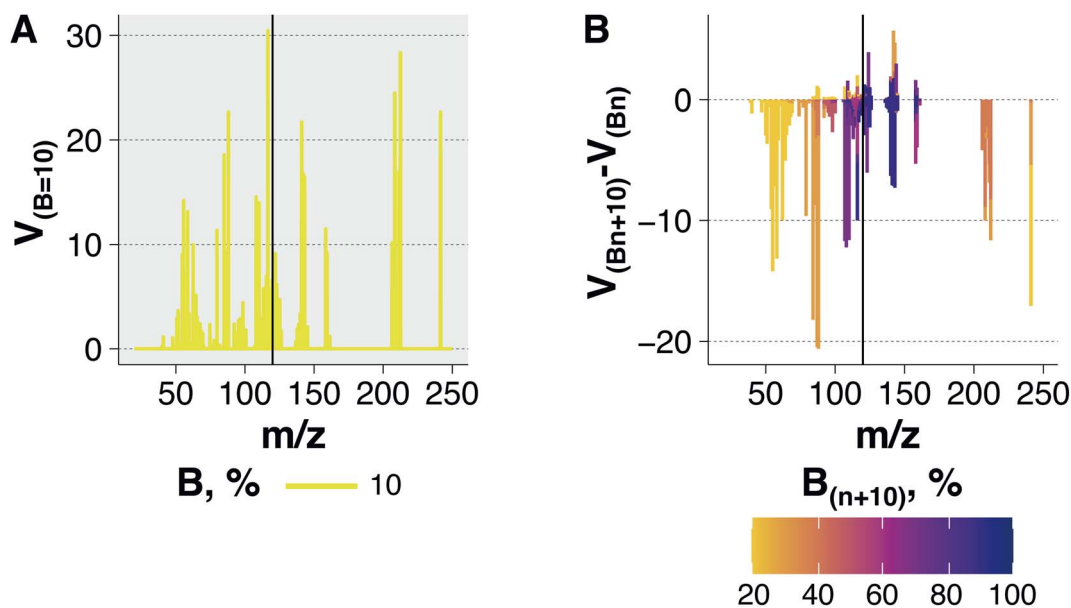


Fig. 2 A) Overall signal intensities for a mass scan with  $m_0 = 120$  from 20 to 250  $m/z$  with induction set at  $B = 10\%$ . (B) Difference of signal intensities between two successive mass scans with induction set at  $B(n + 10)$  and  $B(n)$ , with  $n$  increasing by a 10% increment. Amplifier of the axial collector was set at  $10^{11} \Omega$ .

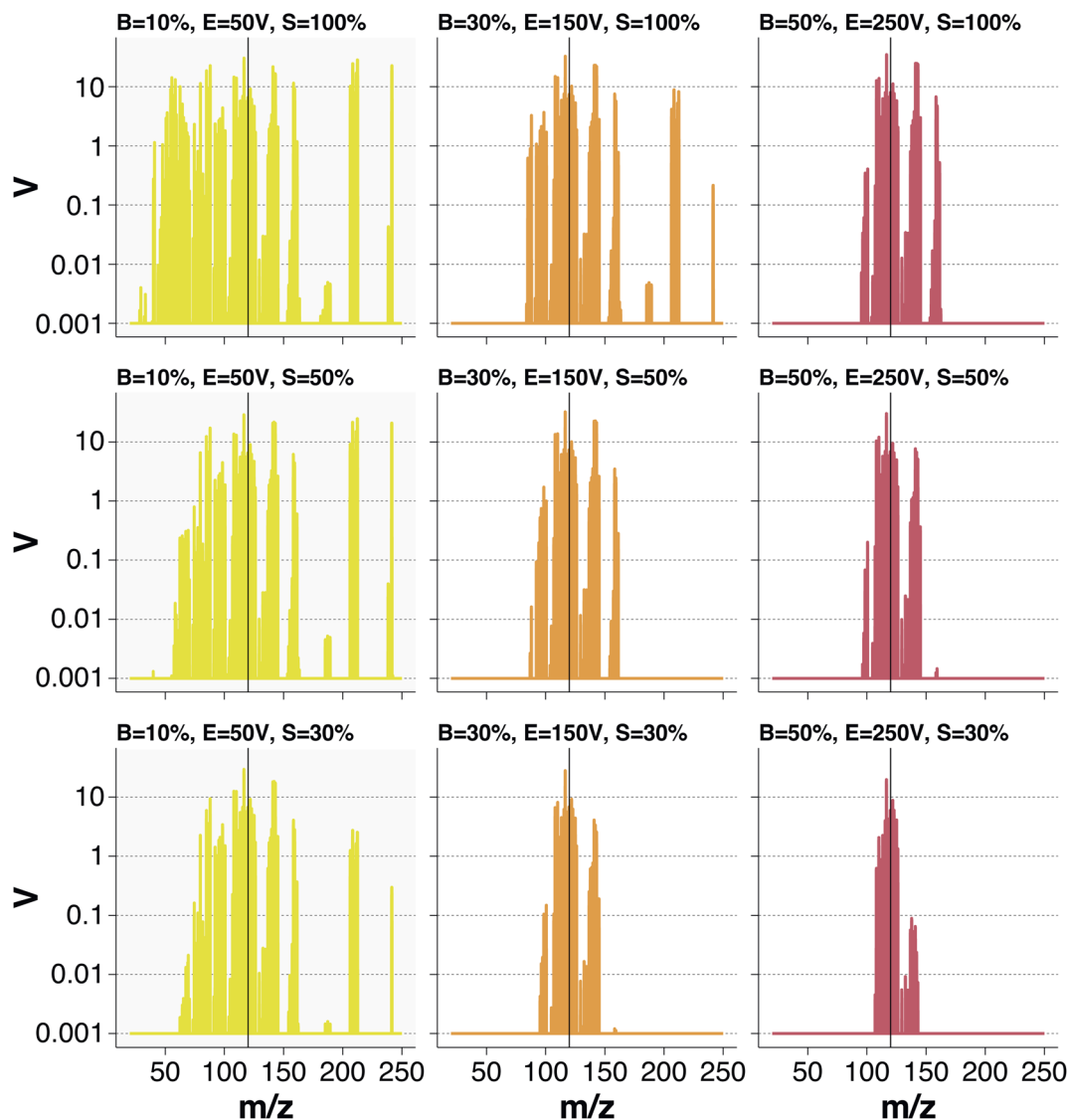


Fig. 3 Overall signal intensities for a mass scan with  $m_0 = 120$  from 20 to 250  $m/z$  for various values of induction and adjustable slit aperture. Amplifier of the axial collector was set at  $10^{11} \Omega$ .

bandpass window equivalent to that with  $B = 30\%$  and  $S = 50\%$ , for which masses are sharply filtered at about 40 amu away from the axial 120  $m/z$ . Decreasing the width of the bandpass window to 30 amu from either side of the axial 120  $m/z$  is obtained using combinations of  $B = 50\%$  and  $S = 50\%$ , or  $B = 30\%$  and  $S = 30\%$ . A half width of 20 amu of the bandpass window can be obtained with a combination of  $B = 50\%$ ,  $S = 30\%$ . Thus, the behavior of the MS/MS prefilter is in agreement, but for low masses only, with the expectations of Craig *et al.*<sup>48</sup> for the Vienna prototype, *e.g.*, that closing the slit will decrease the overall size of the bandpass window without altering the sides steepness. However, further increasing  $B$  by 20 percent (to 50%) with a full open adjustable slit efficiently steepens the bandpass window on the low mass side but does not greatly affect high masses. Again, this behavior of the MS/MS prefilter is in agreement, but for low masses only, with that anticipated from the Vienna prototype.<sup>48</sup>

We further analyze the behavior of the MS/MS prefilter for a  $m_0$  lower than  $^{120}\text{Sn}$  set at 63 amu ( $^{63}\text{Cu}$ ) with the adjustable slit fully open, or 50% or 70% closed ( $S = 100\%$ , 50% and 30%, respectively) and  $B$  values set at 10%, 30%, or 50%, while keeping  $E/B$  constant. The resulting scans are given in Fig. 4. Closing the adjustable slit when  $B$  is minimal produces an asymmetrical bandpass window, but with an asymmetry less pronounced than when  $m_0$  was set at 120 amu (Fig. 3). When  $B$  is minimal (10%) and with a slit 70% closed, or when  $B$  is set at 30% with a slit 100% open, it now produces a more symmetric bandpass window, which was not the case when  $m_0$  was set at 120 amu. Thus, the capability of the MS/MS prefilter to remove low and high masses varies as a function of  $m_0$ .

To investigate this issue, we calculated the theoretical total deviation in the  $X$  direction (*i.e.*, the direction of the electric field in the Wien filter) of a beam as a function of its mass, for masses above and below  $m_0$  as described in ref. 47 (see details in

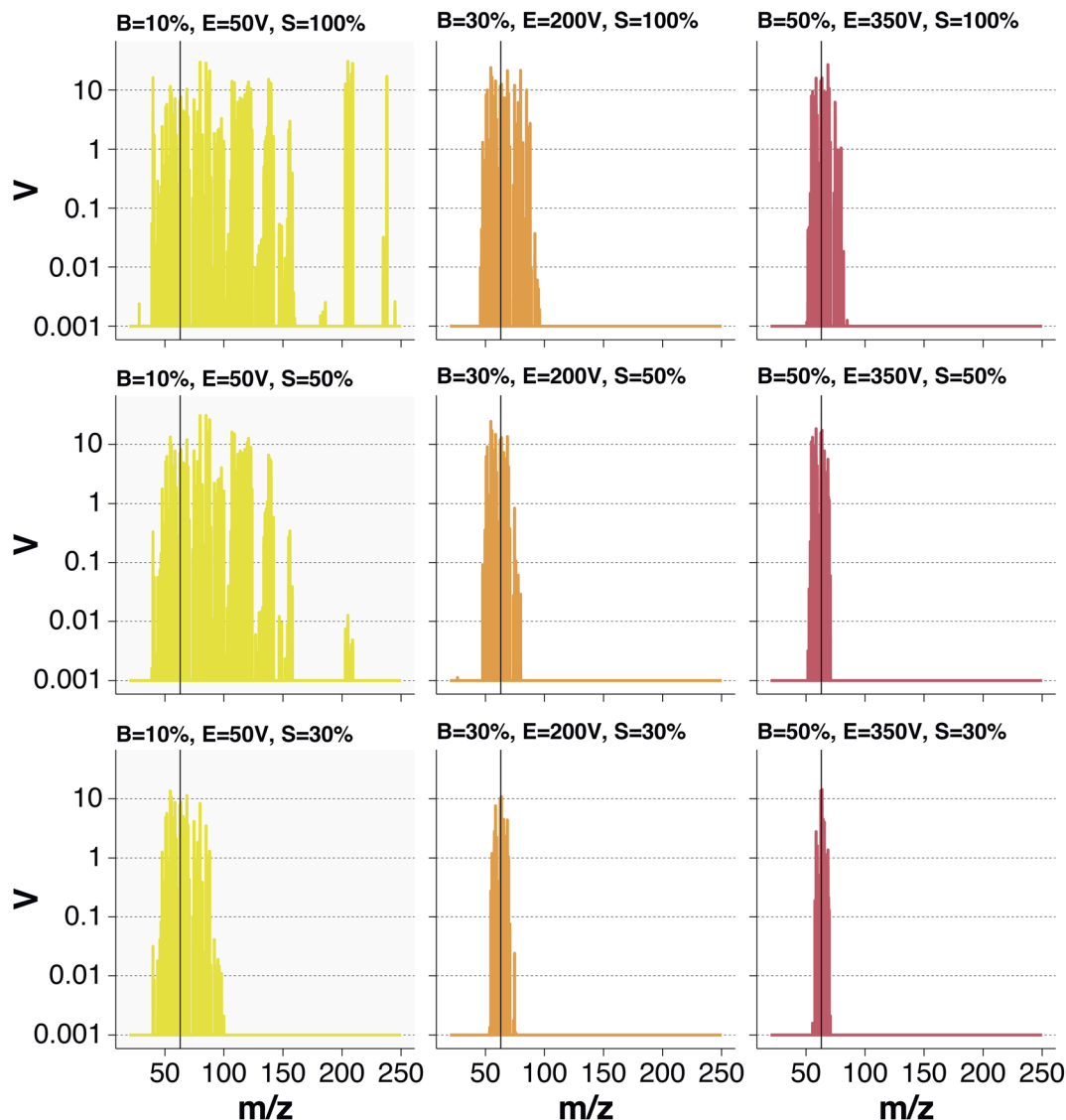


Fig. 4 Overall signal intensities for a mass scan with  $m_0 = 63$  from 20 to 250  $m/z$  for various values of induction and adjustable slit aperture. Amplifier of the axial collector was set at  $10^{11} \Omega$ .

ESI<sup>+</sup>). The value of the total deviation at the position of the slit ( $\Delta X_{\text{tot}}$ ) was shown to be equal to:

$$\Delta X_{\text{tot}} = \left( E - \sqrt{\frac{2E_{\text{kin}}}{m}} B \right) \frac{ql^2}{4E_{\text{kin}}} + l_2 \frac{\left( \frac{E}{B} - v_{\text{init}} \right) \sin\left( \frac{\omega l}{v_{\text{init}}} \right)}{\frac{E}{B} + \left( v_{\text{init}} - \frac{E}{B} \right) \cos\left( \frac{\omega l}{v_{\text{init}}} \right)}$$

where  $E_{\text{kin}}$  is the kinetic energy of ions,  $m$  their mass,  $v_{\text{init}}$  the initial velocity,  $\omega = qB/m$  and  $l$  the length of the Wien filter along its main axis. This non-linear equation giving  $\Delta X_{\text{tot}}$  as a function of the initial velocity was solved numerically to calculate the deviation in the  $x$ -direction for a given mass, as a function of slit opening and magnetic field in the first Wien filter. As shown in Fig. 5, the calculated deviation of the beam as a function of the magnetic field is not symmetric for high and low masses. For  $m_0 = 120$  amu, the vertical deviation for a  $B$  value of 100% will be 3 mm and 4 mm for a variation of mass-to-charge of

$\pm 20\%$ , respectively. Thus, by either increasing the magnetic field (or closing symmetrically the slit) our calculations show an asymmetric trimming of ion in the low mass and high mass regions, with low mass being more efficiently eliminated.

To evaluate the reliability of the modelling, we next compare the theoretical calculations with the results obtained at  $m_0$  set at 120 amu and 63 amu with different combinations of  $B$  and  $S$  values shown in Fig. 3 and 4. The variation of the asymmetry of the bandpass window as a function of  $B$  value is given for  $m_0$  set at 120 amu and 63 amu in Fig. 6A and B, respectively, and shows a good agreement between modelled and measured results. The modelled transmitted low and high masses adequately reproduce the observation that the symmetry of the bandpass window is improved for weak magnetic field value (30%) when  $m_0$  is low (Fig. 6B). We note, however, a small discrepancy for high masses between 150 and 200 amu that we attribute to an inaccurate measurement of the bandpass window due to a lack

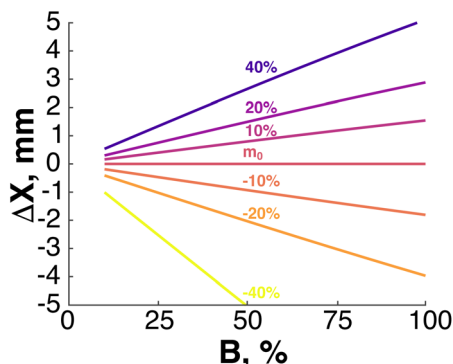


Fig. 5 Deviation in x-axis at the position of the slit located between the two Wien filters for various masses shifted by  $X\%$  as labeled on the curves, as a function of magnetic field in the first Wien filter given in %. The line with no deviation corresponds to the mass  $m_0 = 120$  amu. The length of the Wien is assumed to be 7 cm and the kinetic energy of the singly charged ions is 2 keV. Details for calculations are given in the ESI.†

of elements with those masses (PGE and REE) in the 33 elements solution. The variation of the asymmetry of the bandpass window as a function of the slit aperture is given for  $m_0$  set at 120 amu and 63 amu in Fig. 7A and B, respectively. Again, a good agreement is observed between modelled and measured results, which confirms that the adjustable slit needs to be closed to about 70% to produce a more symmetric bandpass.

In conclusion, the double Wien prefiltering is a highly efficient system for the on-line removal of elements that are 30 to 20 amu lower than a given analyte. It provides an ion beam that only contains the elements of interest to enter the CRC to allow better resolved reactions and collisions.

#### On-line removal of matrix effects

Copper isotopes were measured on the Nu Plasma and Neoma MS/MS according to instrument parameters given in Table 1. For the Neoma MS/MS, setting the axial mass at  $^{66}\text{Zn}$  and the

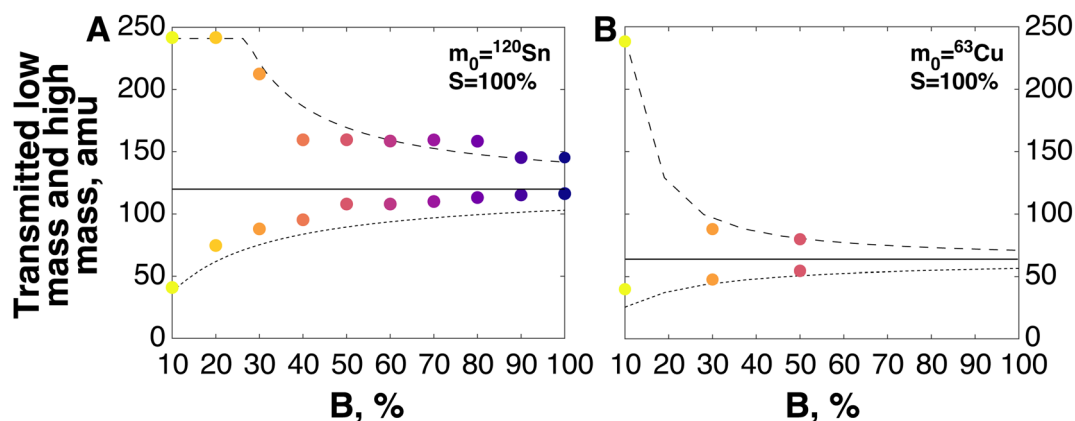


Fig. 6 Calculated low mass (dotted line) and high mass (dashed line) transmitted through the slit located between the two Wien filters as a function of the magnetic field in the first Wien filter (in %). Details for calculations are given in the ESI.† (A)  $m_0 = 120$  amu. Solid circles are from direct measurements as shown in Fig. 3 with a threshold value of 1 V. (B)  $m_0 = 63$  amu. Solid circles are from direct measurements as shown in Fig. 4 with a threshold value of 1 V.

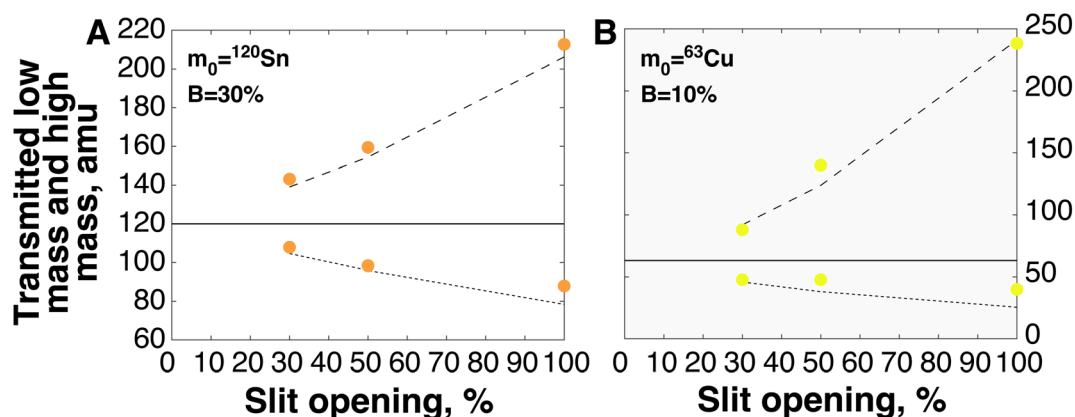


Fig. 7 Calculated low mass (dotted line) and high mass (dashed line) transmitted through the slit located between the two Wien filters as a function of the opening of the slit (in %). Details for calculations are given in the ESI.† (A)  $m_0 = 120$  amu. Solid circles are from direct measurements as shown in Fig. 3 with a threshold value of 1 V. (B)  $m_0 = 63$  amu. Solid circles are from direct measurements as shown in Fig. 4 with a threshold value of 1 V.



magnetic field at 30% with the adjustable slit 70% open efficiently removes any possible argide interferences (Fig. S2†). When Na is added in the Cu–Zn solution with a Na/Cu ratio of 10, the MS/MS prefiltering efficiently removes Na, which is 43 amu lighter than  $^{66}\text{Zn}$  (Fig. S3†). However, a peak scan in an acid blank solution containing only Na still shows the existence of the  $^{40}\text{ArNa}^+$  compound that interferes with the  $^{63}\text{Cu}^+$  isobar (Fig. S4†), demonstrating that the formation of the  $^{40}\text{ArNa}^+$  ion mainly occurs in the plasma source. The addition of He in the CRC is necessary to remove this argide (Fig. S5†), however we found the appearance of a new interference on  $^{68}\text{Zn}$  when He is introduced (Fig. S6†). We hypothesize that the interfering species might be formed of a O- and, in a lesser extent, N-based compound as these are known to form in wet conditions in the presence of He.<sup>60</sup>

The instrumental mass fractionation shows a linear relationship between  $\ln(^{65}\text{Cu}/^{63}\text{Cu})$  vs.  $\ln(^{66}\text{Zn}/^{64}\text{Zn})$  with a slope of  $0.998 \pm 0.049$  and an intercept of  $0.235 \pm 0.036$  ( $R^2 = 0.963$ ) and, despite the presence of the double-Wien filter and the use of He in the CRC, the fractionation factors  $f_{\text{Cu}}$  and  $f_{\text{Zn}}$  are calculated to be  $2.11 \pm 0.01$  ( $\pm 2$  SD,  $n = 18$ ) and  $2.13 \pm 0.01$  ( $\pm 2$  SD,  $n = 18$ ), respectively, thus similar to values reported in the literature (e.g., with the Plasma 54 instrument<sup>2</sup>).

In the present study, we have evaluated the non-spectral matrix effects of Na with a Cu-normalized ratio ranging from 1 to 25 using the Nu Plasma (Fig. 8A) and to 50 using the Neoma MS/MS (Fig. 8B). The results obtained with the Nu Plasma compare well with the literature values, regardless of the type of MC-ICPMS used, the Sapphire (Nu Instrument, Wrexham, UK)

in the conventional high energy pathway<sup>61</sup> or the Neptune.<sup>5,6,10,42</sup> The overall results show that the  $\delta^{65}\text{Cu}$  value begins to be significantly offset with a Na/Cu ratio of 1 (Fig. 8A). In comparison, the measured  $\delta^{65}\text{Cu}$  values with the Neoma MS/MS remain accurate up to a Na/Cu ratio of 10 (Fig. 8B and S7†). At this stage, the  $\delta^{65}\text{Cu}$  value is no longer corrected and becomes negatively biased, suggesting that the  $^{40}\text{Ar}^{23}\text{Na}^+$  interference is no longer spectral and becomes a matrix effect.

### Accuracy of the $\delta^{65}\text{Cu}$ values measurements

Blank signal measured in  $\text{HNO}_3$  2% is 12 mV on  $^{63}\text{Cu}$  for a sensitivity of about  $5 \text{ V ppm}^{-1}$  (Table 1). The repeatability (or short-term external precision) was evaluated by repeated analyses of the SRM-976 solution at  $0.2 \text{ mg L}^{-1}$ , which yields an initial and preliminary value of  $0.03\text{‰}$  ( $\pm 2$  SD,  $n = 13$ ), which will need to be confirmed by further analyses.

Eight geological CRM were purified for Cu using a single step preparation ion-exchange chromatography using the PrepFAST-MC system, and the measured  $\delta^{65}\text{Cu}$  values are given in Table S1† and shown in Fig. 9. The accuracy with already reported  $\delta^{65}\text{Cu}$  values is in good agreement, e.g.,  $-0.04\text{‰}$  for BHVO-1,  $0.02\text{‰}$  for W2a,  $0.00\text{‰}$  for BIR-1,  $0.03\text{‰}$  for AGV-2. We measured, however, a slightly higher  $\delta^{65}\text{Cu}$  value ( $0.22\text{‰}$ ) for BCR-1 than previously measured on different instruments ( $0.09 \pm 0.02\text{‰}$ , Table S1†), for which we do not have at the moment any explanation to propose. Three blood samples show similar  $\delta^{65}\text{Cu}$  values on the Nu Plasma and the Neoma MS/MS within  $0.05\text{‰}$  uncertainties.

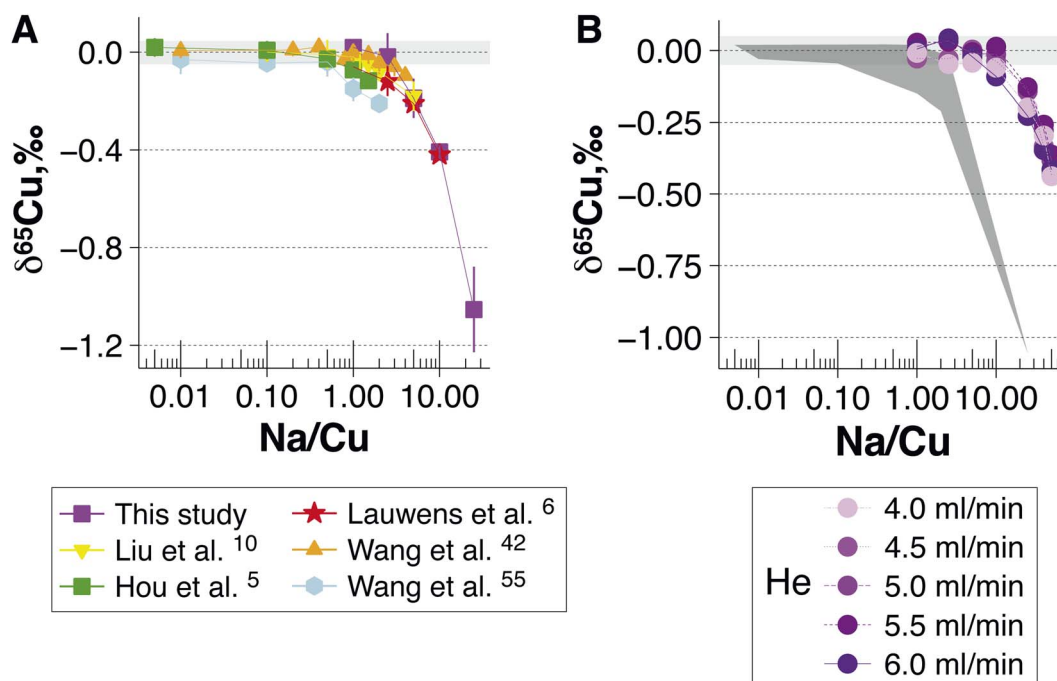


Fig. 8 (A) Matrix effects on the  $\delta^{65}\text{Cu}$  value of the SRM-976 solution as a function of the Na/Cu ratio measured with the Nu Plasma and compared to the literature. The light grey area represents  $\pm 0.05\text{‰}$  deviation. The error bars are  $\pm 2$  SD of the mean. (B) Matrix effects on the  $\delta^{65}\text{Cu}$  value of the SRM-976 solution as a function of the Na/Cu ratio measured with the Neoma MS/MS. Several He flow are tested. The light grey area represents  $\pm 0.05\text{‰}$  deviation and the dark grey area represents the hull of the panel A for comparison.

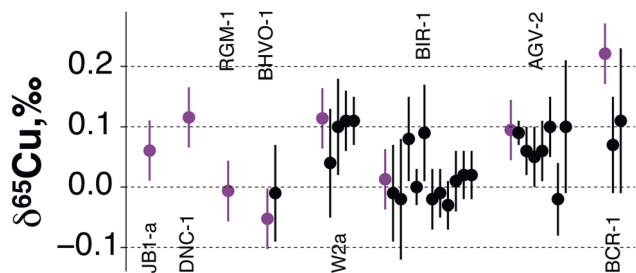


Fig. 9  $\delta^{65}\text{Cu}$  values measured in the present study (purple) and compared with the literature (black). The error bars are  $\pm 2$  SD of the mean (the number of replicates is given in Table S1†).

## Conclusions

Our results indicate that the double-Wien prefiltering system of the Neoma MC-ICPMS/MS produces an asymmetrical mass window transmitted through the slit located between the two filters, relative to the selected axial mass-to-charge value, with removal of low masses being more effective than removal of high masses. This feature is a normal consequence of the ion trajectories in a Wien filter, and taking into account this feature will be necessary during the tuning of the instrument, notably for low masses. Care we be also necessary to avoid trimming the analyte isotope signals, that could possibly produce instrumental instability and isotope fractionation. For well identified reactions and collisions, the possibility of removing part of the mass range will be of significant importance for laser ablation studies, the double-Wien prefiltering will probably allow a neater ion beam to enter the CRC.

## Conflicts of interest

There are no conflicts to declare.

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