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Recent advances in layered $Ln_2NiO_{4+\delta}$ nickelates: fundamentals and prospects of their applications in protonic ceramic fuel and electrolysis cells†

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In the past decade, intensive research on proton-conducting oxide materials has provided a basis for the development of intermediate-temperature protonic ceramic electrochemical cells, which constitute a real alternative to conventional cells based on oxygen-conducting electrolytes. To achieve both high efficiency and excellent performance, not only electrolytes but also electrode materials should be carefully selected considering their functional properties. Compared to the traditional ABO₃ perovskite electrode materials, $Ln_2NiO_{4+\delta}$ with a layered structure has unique advantages (high chemical stability, mechanical compatibility, improved oxygen transport, and hydration ability), and thus is now becoming a hot topic in this field, offering both scientific and practical interests. However, a comprehensive and indepth review is still lacking in the literature to date. Accordingly, this work presents a comprehensive overview of the prospects of layered nickelates ($Ln_2NiO_{4+\delta}$, where Ln=La, Nd, and Pr) as one of the most attractive oxygen (steam) electrode materials for protonic ceramic electrochemical cells. In particular, the crystalline features, defect structure, stability, chemical properties, and mechanical compatibility of this class of materials, contributing to their transport functionality, are discussed with the primary emphasis on revealing the relationship between the composition of the materials and their properties. The presented systematic results reveal the main strategies regarding the utilisation of Ln₂NiO_{4+ô}-based electrodes and existing gaps related to fundamental and applied research aspects.

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Introduction

The ever-increasing demand for the development of highly efficient, environmentally friendly, technologically straightforward and economically expedient energy systems has stimulated rapid progress in the hydrogen and electrochemical energy fields.1-4 Accordingly, research and development related to solid oxide electrochemical cells are widespread due to the higher efficiency, flexibility and variety of these cells for use in energy conversion processes compared with other electrochemical cells.5-7 Although conventional solid oxide electrochemical cells have some disadvantages due to their high operating temperatures, the existing problems can be overcome by replacing the oxygen-ionic conductive electrolytes with proton-conducting counterparts. The improved transfer behaviour of the latter is realised due to the greater mobility of the proton charge carriers, which also exhibit a low migration



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barrier.8-11 Consequently, energy can be efficiently converted in solid oxide electrochemical cells based on proton-conducting electrolytes below 700 °C, and at this temperature, degradation and compatibility issues become less problematic.

Among the different types of electrochemical cells, protonic ceramic fuel cells (PCFCs) and protonic ceramic electrolysis cells (PCECs) are of special interest as systems for converting chemical energy into electricity and vice versa. 12-14 These two types of operations can be combined in reversible protonic ceramic cells (rPCCs), which enable power generation or energy conversion depending on the necessary (momentary) requirements. However, although very promising performances have recently been achieved for both PCFCs and PCECs, 15-20 their long-term operation under multiple cycles of temperature or steam and oxygen partial pressure requires further improvement. Problems associated with degradation phenomena may also occur for these systems, where instead of fast microstructural changes, cationic interdiffusion, segregation and poisoning occur at high temperatures,21-24 chemical and thermal incompatibilities in the electrode/electrolyte pair play a key role in maintaining the integrity of PCFCs, PCECs and

Although the transport properties of the corresponding electrochemical devices are regulated by a proton-conducting membrane, other functional materials also affect their target electrochemical characteristics (power density, hydrogen production rate, etc.). In particular, the selection of suitable oxygen electrodes is currently of great importance for fabricating low- and intermediate-temperature electrochemical systems, including PCFCs, PCECs and rPCCs.25-27 There is a significant number of publications on the design of suitable oxygen electrodes, as systematically detailed in the following works.28-31 Among the various classes of electrodes, herein, we focus on the interesting group of oxides based on lanthanide nickelates.



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fuel/electrolysis cells (SOFC/SOEC) and reversible solid electrolyte cells (r-SOC), mixed-conducting ceramic membranes and membrane reactors, energy conversion and storage materials, and energy generation from renewable sources (biogas, biochar). Dr Yaremchenko has co-authored over 160 papers in SCI journals, including several reviews, with an h-index of 41.

The oxide $Ln_2NiO_{4+\delta}$ (Ln = La, Nd, and Pr) materials having a layered structure belong to the Ruddlesden-Popper (RP) family with the general formula $A_{n+1}B_nO_{3n+1}$, where $n \ge 1$, A is a rare-earth or alkaline-earth element and B is a transition metal.³² The properties of $Ln_2NiO_{4+\delta}$ such as excellent oxygen diffusion coupled with high surface exchange coefficients and reduced thermal expansion coefficients (TECs) have resulted in their wide application in solid oxide electrochemical cells, including that based on ZrO₂, 33,34 CeO₂, 35,36 and LaGaO₃ (ref. 37 and 38) oxygen-conducting electrolytes, and proton-conducting analogues. It should be noted that $Ln_2NiO_{4+\delta}$ oxides are considered to show triple-conducting behaviour when protonic transport exists simultaneously with oxygen-ionic and electronic transportation.39 Since this feature can affect the target parameters of electrochemical systems based on protonconducting electrolytes, it should be analysed in detail. Therefore, this review work is devoted to revealing the application peculiarities of nickelates in PCECs, PCFCs and rPCCs and identifying future strategies for the improvement of their efficiency and performance.

2. Structure and defect chemistry of $Ln_2NiO_{4+\delta}$

2.1. Crystal structure and oxygen nonstoichiometry

Ln₂NiO_{4+δ} nickelates belong to a series of perovskite-related layered compounds with the general formula $A_{n+1}B_nO_{3n+1}$, which are referred to as Ruddlesden-Popper (RP) phases. 40,41 The RP structure is built of consecutive perovskite-type (ABO₃)_n blocks alternating with rock-salt-type AO layers along the c crystallographic axis (Fig. 1). Its formula can be represented as $(AO)(ABO_3)_n$, where n is the number of connected layers of vertex-sharing BO₆ octahedra. Ln₂NiO_{4+δ} and its derivatives are n = 1 members of the homologous RP-type nickelate series. Due to their structural similarities, this type of material is also referred to as K2NiF4-type phases. In the stoichiometric Ln₂³⁺Ni²⁺O₄²⁻ compounds, the A-site Ln cations positioned at the boundary of two types of layers are coordinated by 9 oxygen ions, while the B-site Ni cations are located in the center of octahedra formed by 6 oxygen anions. The NiO₆ octahedra share corners in the a-b plane, forming a 2-dimensional network.

The characteristic feature of the $\rm Ln_2NiO_{4+\delta}$ phases is the possibility to accommodate variable oxygen excess of up to $\delta \sim 0.25\text{--}0.30$ in the case of $\rm Ln = \rm La,^{42,43}$ which is required for the stabilization of the RP-type $\rm A_2BO_4$ structure. The structural stability of $\rm A_2BO_4$ phases is governed by the bond length matching between the perovskite and rock-salt-type layers, which can be rationalized by the Goldschmidt tolerance factor: $^{41,42,44-47}$

$$t = \frac{r_{\rm Ln} + r_{\rm O}}{\sqrt{2}(r_{\rm Ni} + r_{\rm O})} \tag{1}$$

where r is the ionic radii of cations and oxygen in the appropriate coordination.

Empirically, the tetragonal RP-type A_2BO_4 (or K_2NiF_4 -type) structure is stable over the approximate range of 0.85 < t <

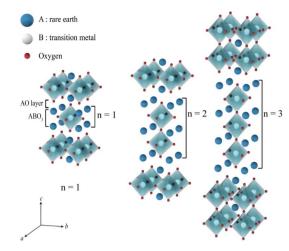


Fig. 1 Idealized representation of the crystal structures of $A_{n+1}B_nO_{3n+1}$ Ruddlesden–Popper phases.

1.00.41,44-46 The perfect matching of two layer types will yield a tolerance factor of t = 1, while any deviation in the t value implies a buildup of internal stress in the a-b plane of the crystal lattice.42,45,47 The ionic radii of rare-earth cations are smaller than the ideal $r_{\rm A}$ of \sim 1.56 Å (estimated employing Shannon's radii). The calculated t = 0.885 in the case of stoichiometric La₂NiO_{4+δ} means that the La-O bonds are under tensile stress and the Ni-O bonds are under compressive stress. Incorporation of interstitial oxygen ions into the rock-salt-type LaO layers relaxes the structure due to several factors as follows: 42,47,49,50 (i) an increase in the average Ln-O bond length as a result of an increase in the Ln cation coordination; (ii) electrostatic repulsion of intercalated oxygen ions, leading to the same effect; and (iii) shortening of the average Ni-O bond length due to the partial oxidation of nickel cations as a result of the following charge compensation:

$$2Ni_{Ni}^{\times}+0.5O_{2} \xleftarrow{T \downarrow pO_{2}\uparrow} 2Ni_{Ni}^{\star}+O_{i}^{''} \tag{2}$$

with electroneutrality condition

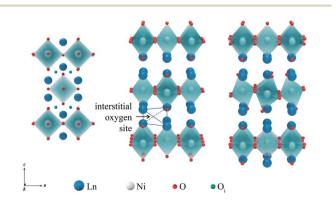


Fig. 2 Crystal structures of A_2BO_4 phases (A = Ln and B = Ni) with tetragonal (center) and orthorhombic (right) lattice symmetry. The a-c projection of the tetragonal lattice (left) illustrates the positions of the interstitial oxygen ions.

$$\left[N_{N_{i}}\right] = 2\left[O_{i}^{"}\right] \quad \text{or} \quad p = 2\delta$$
 (3)

where p is the concentration of electron holes.

Due to geometric constraints, only three rare-earth cations can form RP-type $\text{Ln}_2\text{NiO}_{4+\delta}$ compounds in the undoped form, *i.e.* lanthanum, praseodymium and neodymium. The oxygen excess, δ , in these phases reversibly decreases on heating or reducing the oxygen partial pressure as a result of oxygen release from the lattice and reduction of Ni cations *via* reverse reaction (2), but is frozen at temperatures below $\sim\!350\,^{\circ}\text{C}$ (ref. 51–54) due to kinetic reasons. High-temperature $p\text{O}_2$ –T– δ diagrams and their analysis can be found in ref. 55–65. At elevated temperatures, $\text{Ln}_2\text{NiO}_{4+\delta}$ has a tetragonal structure with the space group $I4/mmm^{55,66,67}$ (Fig. 2) and interstitial oxygen ions in the $\left(\frac{1}{4}, \frac{1}{4}, \frac{1}{4}\right)$ positions in the unit cell. ⁶⁸ A common view is that

on cooling in air, the effect of intercalation of interstitial oxygen ions becomes insufficient for the release of internal strains in the tetragonal lattice due to the different thermal expansion of the Ln-O and Ni-O bonds. Consequently, the mismatch between the bond lengths in the perovskite and rock-salt-type layers is additionally released by titling the NiO6 tetrahedra, leading to a symmetry reduction to orthorhombic symmetry (or even monoclinic).69,70 Aspera et al.71 suggested that increasing the concentration of interstitial oxygen in rock-salt-type LnO layers may be a major factor contributing to the change in structure: the interstitial oxygen ions repulse the neighboring apical oxygen ions in the perovskite layers, thus causing tilting of the octahedra. Different space groups have been proposed to describe the orthorhombic lattice of air-equilibrated nickelates (Fig. 2), including $\mathit{Fmmm}^{53,67,68,70,72-74}$ for $La_2NiO_{4+\delta}$, Fmmm^{73-75} and $Bmab^{72,75}$ for $Pr_2NiO_{4+\delta}$, and Fmmm, $^{72,76-78}$ $Abma^{76}$ and $Bmab^{53,79}$ for $Nd_2NiO_{4+\delta}$. Although praseodymium exists in a mixed 3+/4+ oxidation state in binary PrOx oxides under oxidizing conditions, analysis of the XANES spectra collected at 25-700 °C demonstrated that the Pr cations retain only the 3+ state in the $Pr_2NiO_{4+\delta}$ structure in air.80

The ionic radius of the Ln3+ cations decreases in the sequence La > Pr > Nd, 50 leading to a corresponding increase in oxygen excess content necessary to compensate the structural strain under identical conditions, 53,81-83 and also a shift of the orthorhombic-tetragonal transition to a higher temperature. For air-equilibrated $Ln_2NiO_{4+\delta}$, the reported values of δ at room temperature increase from 0.15–0.18 for $Ln = La^{43,49,66,67,72,73,84,85}$ to 0.20-0.24 for $Ln = Pr^{49,54,72,73,84}$ and to 0.22-0.28 for Ln =Nd. 49,72,76,77,86,87 The orthorhombic-to-tetragonal transition on heating in air was reported to occur at ~50-150 °C in airequilibrated $La_2NiO_{4+\delta}$, 53,66,67,72,84 420-450 °C in Pr_2 - $NiO_{4+\delta}$, 49,54,72,84 and 520–620 °C in $Nd_2NiO_{4+\delta}$. 49,53,72,76,78,86,88,89 In the case of $Pr_2NiO_{4+\delta}$ and $Nd_2NiO_{4+\delta}$, this reversible structural transformation is accompanied by a discrete change in the oxygen nonstoichiometry and occurs when δ in the orthorhombic modification is decreased to a certain level by 0.175-0.20 for $Pr_2NiO_{4+\delta}$ (ref. 54, 72 and 90) and 0.135-0.155 for $Nd_2NiO_{4+\delta}$. ^{72,91} Consequently, the temperature of the transition decreases with reducing pO_2 . 54,72,76,88,90,91 In La₂NiO_{4+ δ} ($\delta \ge 0.15$),

an orthorhombic-to-tetragonal transition occurs at reduced temperatures when oxygen exchange is kinetically frozen, but the temperature of this transition increases with an increase in the frozen-in δ value.^{55,66}

In addition to interstitial oxygen ions, oxygen defects important for transport and electrocatalytic properties include oxygen vacancies in the perovskite layers of the RP-type structure. Oxygen vacancies may form due to intrinsic Frenkel-type disorder

$$O_{O}^{\times} \Leftrightarrow V_{O}^{"} + O_{i}^{"} \tag{4}$$

or as a result of modifications of the cation composition as discussed below.

2.2. A-site deficiency and Ln₂NiO_{4+δ}-based solid solutions

Similar to ABO₃ perovskites, the RP-type A₂BO₄ structure is very flexible to the introduction of A-site cation deficiency and various substitutions in both sublattices, with effects on structural properties, oxygen nonstoichiometry, ionic and electronic transport, electrocatalytic activity, and other relevant properties. The reported solubility of selected cations in the A and B sublattices of $\rm Ln_2NiO_{4+\delta}$ is summarized in Table 1. However, the existing range of $\rm Ln_2NiO_{4+\delta}$ -based solid solutions is not limited to individual doping in the Ln or Ni sites, but also includes a variety of co-substitutions in one or both sublattices.

2.2.1. A-site cation deficiency. Available literature data indicates that the $\text{Ln}_2\text{NiO}_{4+\delta}$ structure can tolerate a certain concentration of cation vacancies in the Ln sublattice. The formation of $\text{Ln}_{2-x}\text{NiO}_{4+\delta}$ solid solutions was reported up to x=0.15 for $\text{Ln}=\text{La},^{51,85,122-124}$ x=0.30 for $\text{Ln}=\text{Pr},^{75,92}$ and x=0.10 for $\text{Ln}=\text{Nd},^{53,87,125,126}$ although a detailed phase and structural analysis was not always provided. Formation of A-site deficient nickelates may also be affected by the employed synthetic procedure. In particular, the presence of phase impurities was reported for $\text{La}_{2-x}\text{NiO}_{4+\delta}$ $(x \ge 0.02-0.03)^{127,128}$ and $\text{Pr}_{2-x}\text{NiO}_{4+\delta}$ $(x \ge 0.10 \text{ and } 0.20).^{75}$

An increase in the concentration of Ln vacancies was reported to gradually decrease the oxygen hyperstoichiometry in the lattice, 51,75,85,87,126 eventually leading to oxygen deficiency even in air. 51,53 However, room-temperature neutron diffraction studies of $\text{Pr}_{2-x}\text{NiO}_{4+\delta}$ demonstrated that an increase in praseodymium deficiency results in an increase in the concentration of oxygen vacancies in the perovskite layer, in the equatorial and, preferentially, apical positions, while the concentration of interstitial oxygen remains essentially unchanged. Thus, the introduction of cation vacancies into the A-sublattice of $\text{Ln}_2\text{NiO}_{4+\delta}$ is charge compensated preferentially by the formation of oxygen vacancies, implying the importance of Frenkel disorder in these phases, and the full electroneutrality condition can be expressed as

$$3\left[\mathbf{V}_{\mathrm{Ln}}^{"'}\right] + 2\left[\mathbf{O}_{\mathrm{i}}^{"}\right] = \left[\mathbf{N}\mathbf{i}_{\mathrm{Ni}}^{\star}\right] + 2\left[\mathbf{V}_{\mathrm{O}}^{\star}\right] \quad \text{or} \quad p = 3x + 2\delta, \quad (5)$$

where $\delta = [O_i''] - [V_O'']$, and $V_{Ln}^{"'}$ denotes Ln vacancy.

2.2.2. Mutual solubility. The full range of solid solutions was reported to exist in $La_2NiO_{4+\delta}-Pr_2NiO_{4+\delta}$, 73,83,93,129 $La_2NiO_{4+\delta}-Nd_2NiO_{4+\delta}$, 89 and $Pr_2NiO_{4+\delta}-Nd_2NiO_{4+\delta}$ (ref. 74 and

Table 1 Reported solubility of substituting cations in the A- and B-sublattices of $Ln_2NiO_{4+\delta}$ under oxidizing conditions

$\frac{\mathrm{Ln}_{2-x}\mathrm{A}_{x}\mathrm{NiO}_{4\pm\delta}}{\mathrm{V}_{\mathrm{Ln}}^{'''}}$ La Pr	0.15 (max)	Ref.	x	Dof		
	0.15 (max)			Ref.	x	Ref.
La		51	0.30	92	0.10	53
Pr	_		2.0	73, 83 and 93	2.0	89
	2.0	73, 83 and 93	_	•	2.0	74 and 94–96
Nd	2.0	89	2.0	74 and 94-96	_	
Sm	1.1	97				
Ca	$0.6 \ (max)$	98	0.7	101	$0.6 \ (max)$	79
	$0.3 < x_{(max)} < 0.4$	99 and 100	0.5 (max)	102	0.5 (max)	78
	,				1.0	103
Sr	1.6 (max)	104 and 105	1.6	106 and 107	1.6	108
	1.5 (max)	98			1.67	104
Ba	1.0	109	0.4	110	0.6 (max)	79
	1.1 (max)	98				
	Ln = La		Ln = Pr		Ln = Nd	
$Ln_2Ni_{1-y}M_yO_{4+\delta}$	у	Ref.	у	Ref.	у	Ref.
Cu	1.0	111-114	0.4 (max)	111	0.3 (max)	115 and 116
			0.5	45	()	
Co	$0.20 \; (max)$	97 and 117	0.1	118	0.1	119
Fe	0.10	120 and 121	0.1	52	0.1	119
Mn					\sim 0.1	119

96) pseudobinary systems, which is expected considering the structural identity between the parent materials. Changing the fractions of two Ln cations in the A sublattice leads to a gradual change in the structural properties, oxygen nonstoichiometry and temperature of the orthorhombic-to-tetragonal transition on heating, in agreement with the properties of individual $\text{Ln}_2\text{NiO}_{4+\delta}$ compounds. In the case of $\text{La}_{2-x}\text{Pr}_x\text{NiO}_{4+\delta}$ solid solutions, the oxygen nonstoichiometry, δ , at temperatures above \sim 400 °C was observed to exhibit a rather weak dependence on the praseodymium content up to $x \sim 1.5$, and then decreased rapidly. Ontrary to other studies, Vibhu *et al.* Precently reported the existence of a biphasic region in the $\text{La}_{2-x}\text{Pr}_x\text{NiO}_{4+\delta}$ system at x = 0.5–1.0 based on roomtemperature high-resolution synchrotron and neutron powder diffraction data.

2.2.3. A-site substitutions by alkaline-earth cations. RP-type nickelates form a wide range of $\text{Ln}_{2-x}\text{A}_x\text{NiO}_{4+\delta}$ solid solutions, where A = $\text{Ca},^{53,78,79,99-102,109,130-132}$ $\text{Sr}^{56-58,79,83,105,106,108,109,133-149}$ or $\text{Ba}.^{79,98,109,110}$ As a common effect, an increase in x results in a gradual decrease in oxygen content from excess $(x \leq 0.3-0.4)$ to oxygen stoichiometry $(\delta \approx 0$ for intermediate x) and eventually to oxygen deficiency $(x \geq 1.0)$. This is accompanied by a decrease in the orthorhombic-totetragonal transition temperature (below room temperature for $x \geq 0.1-0.3$), 78,79,101,102,110,130,132,141,143,144,146,147,149,150 a gradual increase in Ni^{3+} concentration, 57,83,98,105,134,136,151 and even the formation of Ni cations in the formal 4+ oxidation state for Arich compositions. 98,105,108,145,152 Thus, acceptor-type doping by

alkaline-earth cations is charge compensated by a decrease in interstitial oxygen concentration and the formation of electron holes and oxygen vacancies (Fig. 3). The full electroneutrality condition is given by

$$\left[\mathbf{A}_{Ln}^{'}\right] + 2\left[\mathbf{O}_{i}^{"}\right] = \left[\mathbf{N}\mathbf{i}_{Ni}^{"}\right] + 2\left[\mathbf{N}\mathbf{i}_{Ni}^{"}\right] + 2\left[\mathbf{V}_{O}^{"}\right] \quad \text{or} \quad p = x + 2\delta$$
(6)

where $p = [Ni_{Ni}^{"}] + 2[Ni_{Ni}^{"}]$, $\delta = [O_i^{"}] - [V_O^{"}]$, and $O_i^{"}$ and $V_O^{"}$ prevail for small and large A-cation fractions, respectively.

The solubility of alkaline-earth cations decreases in the sequence Ca < Ba < Sr (Table 1), in agreement with considerations of the tolerance factor, t, (eqn (1)) and stability of the A₂BO₄ structure.⁹⁸ Sr²⁺ has a slightly higher ionic radius than that of Ln3+ (coordination number of A-site cations in the idealized Ln₂NiO₄ structure is 9).⁴⁸ Therefore, the incorporation of strontium into the Ln sites effectively releases the tensile stress in the Ln–O bonds, and Sr cations may occupy up to $\sim 5/6$ sites in the A sublattice of $Ln_{2-x}A_xNiO_{4+\delta}$. The ionic radius of Ca²⁺ is between that of La³⁺ and Nd³⁺ and is similar to the ionic radius of Pr³⁺ (ref. 50), thus having the opposite, but comparatively small steric effect on the RP lattice for different Ln. In combination with the shortening of the Ni-O bond lengths due to gradual $Ni^{2+} \rightarrow Ni^{3+}$ oxidation, this results in the comparatively limited solubility of calcium in the A-sublattice. Finally, large Ba²⁺ cations⁴⁸ are even more effective than Sr²⁺ for releasing the tensile stress in the rock-salt-type layers; however, the tolerance factor reaches the upper limit of t = 1 at a lower concentration compared to the Sr-substituted counterparts.

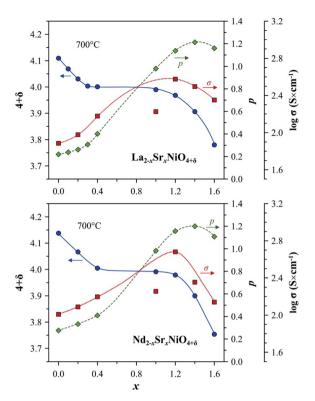


Fig. 3 Trends in the variation of oxygen nonstoichiometry $4 + \delta$, electron-hole concentration, p, and electrical conductivity, σ , with strontium content, x, in $Ln_{2-x}Sr_xNiO_{4+\delta}$ systems at 700 °C. The experimental data on oxygen nonstoichiometry and electrical conductivity are from ref. 56, 59, 108, 142, 145 and 148. The electronhole concentration is calculated using eqn (6). Note that σ of the Srrich compositions is undervalued due to the porosity of samples and microcracking effects. 108,145

The level of A-site doping and, consequently, the level of oxygen content has an impact on the variations of oxygen nonstoichiometry with temperature and pO2 within the phase stability domain of RP-type Ln_{2-x}A_xNiO_{4+δ} phases at temperatures above $\sim 400~^{\circ}\text{C}$ as follows: $^{56-59,81,106,108,131,139,140,142,145,153-155}$

(i) Oxygen-hyperstoichiometric phases ($\delta > 0$ in air). These compositions behave similarly to the parent nickelates. The oxygen content decreases on heating and with a reduction in oxygen partial pressure, but the materials maintain oxygen excess.

(ii) Oxygen-stoichiometric and nearly stoichiometric phases ($\delta \sim$ 0 in air). These materials demonstrate negligible δ variations upon thermal and pO_2 cycling under oxidizing conditions, but tend to exhibit moderate temperature-dependent oxygen deficiency at reduced oxygen pressures. Specifically, increasing x suppresses the formation of interstitial oxygen under oxidizing conditions and promotes the formation of oxygen vacancies under reducing conditions.

(iii) Oxygen-deficient solid solutions ($\delta < 0$ in air). Oxygen deficiency increases on heating and with a decrease in oxygen partial pressure. These phases tend to exhibit a structural transformation to a distorted orthorhombic structure (space group Immm) under moderately reducing conditions (e.g. $pO_2 \sim$

10⁻⁵ atm) accompanied by ordering in the oxygen sublattice and substantial dimensional changes. 108,155 Besides, A-rich phases exhibit strongly anisotropic thermochemical expansion. which complicates high-temperature ceramic processing. 108,138,145

2.2.4. B-site substitutions by transition metal cations. The formation of a full range of solid solutions with orthorhombic or tetragonal structure was reported for the La₂NiO_{4+δ}-La₂- $CuO_{4+\delta}$ pseudobinary system. Due to the stable 2+ oxidation state of the copper cations under oxidizing conditions, increasing y in $La_2Ni_{1-\nu}Cu_{\nu}O_{4+\delta}$ results in a gradual decrease in oxygen excess in the lattice, ^{113,156,157} down to $\delta \leq 0.01$ for y = 0.4–0.8 at room temperature. Since Cu^{2+} has a larger ionic radius than that of Ni2+ (coordination number of B-site cations in the idealized Ln₂NiO₄ structure is 6),⁴⁸ increasing the B-O bond length and a decrease in δ as a result of Cu substitution decrease the tolerance factor111 and destabilize the RP-type A₂BO₄ structure. This leads to the limited solubility of the copper cations in the nickel sublattice of praseodymium and neodymium nickelates, i.e. y = 0.4-0.5 for $Pr_2Ni_{1-\nu}Cu_{\nu}O_{4+\delta}$ (ref. 45, 111, 158 and 159) and y = 0.3 for $Nd_{2}Ni_{1-y}Cu_{\nu}O_{4+\delta}.^{87,115,116}$

Other transition metal cations (such as Co, Fe, Mn, Cr, and Ti) tend to exhibit higher oxidation states in oxide compounds under oxidizing conditions, and therefore have very limited (if any) solubility in the nickel sublattice of $Ln_2NiO_{4+\delta}$. The solid solubility range of cobalt in $La_2Ni_{1-\nu}Co_{\nu}O_{4+\delta}$ was found to be limited to y = 0.2 in air, 97,117,160 although it can be extended under reducing conditions.97,117 The confirmed solubility of cobalt cations in $Ln_2Ni_{1-\nu}Co_{\nu}O_{4+\delta}$ (Ln = Pr and Nd) and iron cations in $Ln_2Ni_{1-\nu}Fe_{\nu}O_{4+\delta}$ (Ln = La, Pr, and Nd) is only 10 at%.52,60,118-121,161 Miyoshi et al.162 reported the results of oxygen permeation studies of $La_2Ni_{0.8}Fe_{0.2}O_{4+\delta}$ and $Pr_2Ni_{0.8}M_{0.2}O_{4+\delta}$ (M = Co, Fe, and Mn), but did not present a detailed analysis of the XRD data, and thus their phase purity is questionable. Phase impurities were detected in the XRD patterns of $Nd_2Ni_{0.9}Mn_{0.1}O_{4+\delta}$, ¹¹⁹ $La_2Ni_{0.9}M_{0.1}O_{4+\delta}$ (M = Ta and V), 163,164 and La₂Ni_{1- ν}Mo $_{\nu}$ O_{4+ δ} (ν = 0.0125-0.05). 165

The impact of partial substitution of nickel by cobalt or iron in Ln₂NiO_{4+δ} in terms of defect chemistry is an increase in the concentration of interstitial oxygen and holes. 60-62,97,120,121,160,161 The corresponding electroneutrality condition is

$$\left[N_{N_{i}}\right] + \left[M_{N_{i}}\right] = 2\left[O_{i}^{"}\right] \quad \text{or} \quad p = 2\delta,$$
 (7)

where $p = [Ni_{Ni}^{\bullet}] + [M_{Ni}^{\bullet}]$.

In the La₂Ni_{0.9}M_{0.1}O_{4+ δ} series, δ was found to increase in the sequence $M = Cu < Ni < Co < Fe.^{61,62,65,120,121}$ A similar tendency was observed for B-site-doped Pr₂NiO_{4+δ}.⁵²

The solid solubility of transition metal cations in the nickel sublattice of Ln₂NiO_{4+δ} under oxidizing conditions can be substantially expanded by co-substitutions with alkaline earthcations into the Ln sublattice. In that case, donor-type doping into the B sublattice is partially charge-compensated by acceptor-type doping in the A sublattice. This leads to a complex interplay between the cation composition,

nonstoichiometry, transport and electrochemical properties. The reported examples of co-substitutions include $La_{2-x}Sr_xNi_{1-y}M_yO_{4+\delta}$ (M = Co, Fe, Mn, Ti, and Mo), $^{45,165-174}La_{2-x}Ca_xNi_{1-y}M_yO_{4+\delta}$ (M = Fe), 131 Nd $_{2-x}Sr_xNi_{1-y}M_yO_{4+\delta}$ (M = Cu and Co), 175,176 and $Pr_{2-x}Sr_xNi_{1-y}M_yO_{4+\delta}$ (M = Co, Mn). 177,178 The solubility of the transition metal typically increases with an increase in the fraction of alkaline-earth cation in the Asublattice, and in the case of Co or Fe, may reach the entire range of B-site concentrations. $^{170,172,174-176}$

The synthesis of Cr-doped $(Ln,A)_2Ni_{1-y}Cr_yO_{4+\delta}$ (A – alkalineearth cation) is only possible at reduced oxygen pressures, under an argon atmosphere. As an exception, Ishihara et al. Proported the preparation of $Pr_2Ni_{0.75}Cu_{0.20}M_{0.05}O_{4+\delta}$ (M = Cr and V) in air, but without accurate structural studies.

2.2.5. B-site substitutions by other metal cations. The data on other B-site substitutions are rather scarce, and in many cases, accurate analysis of the phase composition and structural data are not available.

Ye and Hertz¹⁸⁶ reported that Mg^{2^+} can substitute into the nickel sublattice up to y=0.3 in $La_2Ni_{1-y}Mg_yO_{4+\delta}$ and y=0.2 in $La_{1.85}Ni_{1-y}Mg_yO_{4+\delta}$, and that the solubility can be further extended by strontium co-substitution into the lanthanum sublattice. Zhang $et~al.^{38}$ confirmed that $La_2Ni_{1-y}Mg_yO_{4+\delta}$ is formed at least in the range $0 \le y \le 0.1$. The formation of $La_{1.7}Sr_{0.3}Ni_{1-y}Mg_yO_{4+\delta}$ (y=0–0.5) was reported by Meeporn $et~al.^{187}$ On the other hand, Klande $et~al.^{163}$ failed to prepare single-phase $La_2Ni_{0.9}Mg_{0.1}O_{4+\delta}$.

Ganguli *et al.*¹²³ reported the formation of Zn-substituted $La_{1.9}Ni_{0.9}Zn_{0.1}O_{4+\delta}$ and $La_{1.9}Sr_{0.1}Ni_{0.9}Zn_{0.1}O_{4+\delta}$, while Silva *et al.*¹⁸⁸ claimed the zinc solubility of up to 40 at% in the nickel sublattice.

The confirmed solubility of aluminum in La₂Ni_{1-y}Al_yO_{4+ δ} corresponds to 5 at% in the nickel sublattice, ^{189,190} while La₂Ni_{0.9}Al_{0.1}O_{4+ δ} was reported to comprise of secondary phases. ¹⁶³ Surprisingly, donor-type doping by Al³⁺ was found to reduce the oxygen excess in La₂Ni_{0.95}Al_{0.05}O_{4+ δ} compared to the parent nickelate. ¹⁸⁹

The formation of the full range of LaSrNi_{1-y}Ga_yO_{4± δ} (y=0–1.0) solid solutions with a tetragonal structure was reported by Reinen et~al., ¹⁹¹ whereas the solid solubility of scandium in LaSrNi_{1-y}Sc_yO_{4+ δ} was found to be limited to y=0.2. ¹⁹² Zhang et~al. ¹⁹³ also reported the formation of single-phase La_{1.9}Ni_{0.45}Cu_{0.45}Sc_{0.1}O_{4± δ}.

The synthesis of $La_2Ni_{0.9}Zr_{0.1}O_{4+\delta}$ was attempted in ref. 163 but the product contained $La_2Zr_2O_7$ impurity.

The membrane screening tests by Miyoshi *et al.*¹⁶² included oxygen permeation studies of $Pr_2Ni_{0.8}Zn_{0.2}O_{4+\delta}$ and $Pr_2Ni_{0.9}M_{0.1}O_{4+\delta}$ (M = Mg, Al, and Ga), but the reported information is insufficient to conclude about the phase purity of the samples. A number of works focused on investigating the oxygen permeability or electrochemical properties of co-substituted phases in the $(Ln_{1-x}A_x)_{2-a}Ni_{1-y-z}Cu_zM_yO_{4+\delta}$ (Ln = Pr or Nd; A = La or Sr; M = Al, Ga, In, Zr; x = 0–1.0, a = 0–0.10, y = 0–0.10)^{185,194–200} series. The compositions were claimed to be phase-pure or to contain minor phase impurities, which were ignored. In most cases, again, the quality of the

provided XRD data is not sufficient to accurately conclude about the solid solubility limits of the dopant cations M.

2.2.6. Sm₂NiO_{4+ δ}-based solid solutions. As mentioned above, undoped Ln₂NiO_{4+ δ} exists only for Ln = La, Pr and Nd.¹⁰⁴ Sm³⁺ cations are too small to form the Sm₂NiO_{4+ δ} compound. However, the formation of Sm₂NiO_{4+ δ}-based solid solutions is possible if an appropriate larger cation is incorporated into the Sm sublattice to increase the tolerance factor.

 ${
m Sm_{2-x}La_xNiO_{4+\delta}}$ solid solutions with a tetragonal structure were reported to form in the concentration range of $0.9 \le x \le 2.0.97$ Changing the lanthanum content was observed to have a negligible impact on the low-temperature oxygen non-stoichiometry: for all the compositions, δ varied between 0.11 and 0.13 at room temperature.

 ${
m Sm}_{2-x}{
m Sr}_x{
m NiO}_{4+\delta}$ solid solutions exist in a wide range of strontium contents from x=0.4–0.5 up to x=1.67 (i.e. 5/6 sites in the Sm sublattice). ${
m ^{104,143,201-204}}$ As in the case of other Ln series, increasing the Sr content decreases the temperature of the orthorhombic-to-tetragonal transition. At room temperature, the structure changes from orthorhombic for x=0.5–0.7 to tetragonal for x=0.8–1.67. ${
m ^{104,143,204}}$ Under these conditions, ${
m Sm}_{2-x}{
m Sr}_x{
m NiO}_{4+\delta}$ is nearly oxygen-stoichiometric for x=0.5–1.0 and tends to exhibit oxygen deficiency at higher Sr concentrations, while the average Ni oxidation state increases monotonically with $x.^{202}$

2.2.7. Modifications of the oxygen sublattice. Doping into the oxygen sublattice of $Ln_2NiO_{4+\delta}$ remains an unexplored area, although additional anionic defects certainly may affect the crystal structure, defect chemistry and transport properties. Only a few studies can be found on this topic. Bhat et al.205 attempted fluorination of oxygen stoichiometric La2NiO4 at 150 °C. They found that the amount of fluorine entering the bulk of the samples is small, but results in a transformation from orthorhombic to tetragonal structure. Elongation along the c-axis upon fluorination was interpreted as the incorporation of interstitial species. Arbuckle et al.206 prepared fluorinated $Nd_{2-x}Sr_xNiO_{4+\delta}F_y$ (x = 0, 0.5, and 1.0) with the fluorine content changing from $y \sim 0.018$ for x = 0 to $y \sim 0.08$ for x = 1.0. Fluorination was observed to have an impact on the crystal lattice parameters. Contrary to $Nd_2NiO_{4+\delta}$, the fluorinated counterpart showed no evidence of phase transition in the temperature range of up to 700 °C (upper limit of thermal stability of the prepared materials). Tarutin et al.207 reported the synthesis and characterization of $Nd_{1.9}Ba_{0.1}NiO_{4+\delta}F_{\gamma}$ ($\gamma=0$ -0.10) solid solutions with an orthorhombic structure. The incorporation of fluorine was found to have only a minor effect on the structural parameters but results in a gradual decrease of oxygen content in the lattice, whereas the overall anion content $(4 + \delta + \gamma)$ remained nearly constant. Wissel et al.²⁰⁸ prepared and characterized the La2NiO3F2 composition derived from the polymer-based fluorination of the $La_2NiO_{4+\delta}$ parent phase. The XRD technique combined with ¹⁹F magic-angle spinning NMR spectroscopy allowed the localization of the fluorine ions to be determined. Particularly, it was found that the fluoride ions occupied the apical anion sites, whereas the oxide ions are located in the interstitial sites. Nowroozi et al.209 studied a fluorine-containing nickelate for utilization in all-solid-state

fluoride ion batteries (FIB). They demonstrated that La₂NiO_{4+δ} $(\delta = 0.13)$ can accumulate a large amount of fluorine ions during FIB charging at 170 °C, with the formation of highlyfluorinated La2NiO4.13F1.59 with nearly full occupation of the interstitial sites in the rock-salt-type layers by anions. The latter phase has an orthorhombic structure (space group Fmmm) with a unit cell significantly elongated along the c-axis ($c \sim 15.2 \text{ Å}$ compared to $c \sim 12.7 \text{ Å for La}_2 \text{NiO}_{4+\delta}$, confirming the incorporation of fluorine ions into the interstitial sites.

Thermodynamic and chemical stability of Ln₂NiO_{4+δ}

The thermal stability of the phases in the pseudobinary La₂O₃-NiO system (Fig. 4) increases in the sequence $LaNiO_{3-\delta} < La_4Ni_3O_{10-\delta} < La_3Ni_2O_{7-\delta} < La_2NiO_{4+\delta},^{210,215-217}$ consistent with the reduction in the average oxidation state of the Ni cations from 3+ for LaNiO₃ to 2+ for La₂NiO₄. $La_2NiO_{4+\delta}$ is the only stable compound in this system at atmospheric oxygen pressure at temperatures above \sim 1250 °C (ref. 210 and 215–217) and exists up to the melting point of 1670-1750 °C.210 Although La₂NiO_{4+δ} easily forms in air, and once formed can maintain a phase-pure state at different temperatures, RP-type $La_{n+1}Ni_nO_{3n+1}$ (n = 2 or 3) and perovskite-like LaNiO_{3- δ} appear to be thermodynamically more favorable at temperatures below 900-1100 °C under ambient pO_2 . The corresponding pO_2 -T stability ranges can be found in available phase diagrams.210,216,217 However, the transformation of $La_2NiO_{4+\delta}$ into other La-Ni-O phases at temperatures below 1100 °C is kinetically suppressed.^{211,217} While it can be promoted by thermal treatments at elevated oxygen pressures (>150 bar) at 825-850 °C, 156,216 no evidence of phase decomposition can be observed after heating La₂NiO_{4+ô} for 100 h in ambient air at different temperatures in the range of 600-1100 °C.21 However, Amow et al.97,218 demonstrated that ${\rm La_2NiO_{4+\delta}}$ does undergo phase decomposition during a longer term: the formation of a secondary phase (presumably $La_3Ni_2O_{7-\delta}$) was detected after heating at 900 °C for two weeks in air. Later, Gauquelin et al.²¹⁹ observed the segregation of nickel-rich La₄Ni₃O_{10-δ} and La₃Ni₂O_{7-δ} phases on the surface of La₂NiO_{4+δ} single crystals after annealing in air at 1000 °C for 13 days. Furthermore, it was

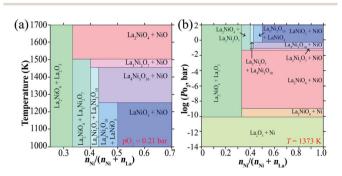
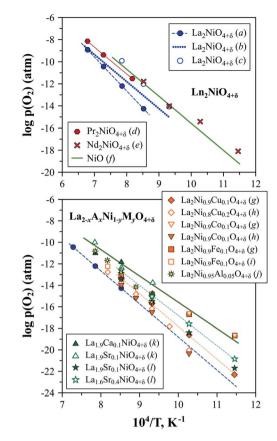


Fig. 4 T- (a) and pO_2 - (b) based phase diagrams in the NiO-LaO_{1.5} oxide systems. Reproduced with permission.²¹⁰ Copyright 2004, Elsevier.



5 Low- pO_2 stability boundaries of Ln₂NiO_{4+ δ} $La_{2-x}A_xNi_{1-v}M_vO_{4+\delta}$ nickelates. Literature sources: (a) ref. 212, (b) ref. 211, (c) ref. 63, (d) ref. 222, (e) ref. 58, (f) ref. 223, (g) ref. 120, (h) ref. 214, (i) ref. 121, (j) ref. 189, (k) ref. 153 and (l) ref. 56.

found that the decomposition of La₂NiO_{4+δ} into higher-order RPtype nickelates is catalyzed and strongly facilitated by contact with Pt at temperatures $\geq 800 \,^{\circ}$ C, ^{220,221} thus raising concerns about the impact of Pt current collectors on the results of electrochemical studies of La2NiO4+6-based electrodes. No such an effect was observed in contact with Au current collectors under similar conditions.220

The stability boundary of $La_2NiO_{4+\delta}$ at reduced oxygen partial pressures at 600–1300 °C was determined in ref. 56, 211 and 212 and corresponds to $pO_2 \sim 10^{-17}$ – 10^{-15} atm at 800 °C (Fig. 5). A decrease in oxygen chemical potential below this boundary results in a reduction to La₂O₃ and metallic Ni. Moderate substitutions (5-20 at%) in the B-sublattice by other transition metal cations (Fe, Co, and Cu)120,121,213,214 or aluminum189 have a rather minor effect of the low-pO2 stability limits at 600-1000 °C (Fig. 5). Nakamura *et al.* ⁵⁶ studied the $La_{2-x}Sr_xNiO_{4\pm\delta}$ (x=0–0.4) system at 600-900 °C and observed only a slight shift in the low $p(O_2)$ stability boundary towards more oxidizing conditions caused by Sr doping for x = 0.2-0.4. Similar results were obtained by Kim et al. 153 for La_{1.9}Sr_{0.1}NiO_{4+ δ} and La_{0.9}Ca_{0.1}NiO_{4+ δ} at 800-1000 °C.

No information on the possible instability of $Nd_2NiO_{4+\delta}$ under oxidizing conditions can be found in the literature. Petrov et al.211 indicated that the low-pO2 stability limits of ${
m Nd_2NiO_{4+\delta}}$ are close to the NiO/Ni boundary. Later, Nakamura $et~al.^{58}$ evaluated the ${
m Nd_{2-x}Sr_xNiO_{4+\delta}}$ (x=0–0.4) system at 600–900 °C and confirmed that the stability boundary of undoped ${
m Nd_2NiO_{4+\delta}}$ is close to that of nickel oxide (Fig. 5), and, similar to the ${
m La_{2-x}Sr_xNiO_{4+\delta}}$ system, the substitution of Nd by Sr has no substantial effect on the stability limits under reduced oxygen pressures.

Thus, electrode materials derived from $\rm Ln_2NiO_{4+\delta}$ (Ln = La or Nd) can be expected to remain tolerant towards reductive decomposition under the typical conditions of oxygen electrode operation, even under very high cathodic polarisation.

The low- $p(O_2)$ stability boundary of $Pr_2NiO_{4+\delta}$ (Fig. 5) is between that of La₂NiO_{4+ δ} and Nd₂NiO_{4+ δ}, ²²² while partial substitutions in the nickel sublattice were reported to shift the stability limits of Pr₂Ni_{0.9}Fe_{0.1}O_{4+\delta} and Pr₂Ni_{0.8}Cu_{0.2}O_{4+\delta} to more oxidising conditions.⁵² Contrary to its La- and Nd-based counterparts, the $Pr_2NiO_{4+\delta}$ phase easily decomposes in the high-p(O2) region. Sullivan et al.222 were the first to show that Pr₂NiO_{4+δ} undergoes a transformation into Ruddlesden-Popper $Pr_4Ni_3O_{10+\delta}$ and PrO_v when exposed to oxidising conditions below a certain temperature (e.g. below ~1030 °C in pure oxygen) (Fig. 6). The instability of the $Pr_2NiO_{4+\delta}$ lattice at higher oxygen pressures was confirmed and assessed in more detail in subsequent studies. $^{52,75,224-229}$ Pr $_2$ NiO $_{4+\delta}$, which is stable in air at $T \ge 925$ °C, can be rapidly quenched down to low temperatures. 52,225 Thermal treatments at $T \le 900$ °C result in a reversible decomposition into a mixture of praseodymium oxide PrO_v, Ruddlesden-Popper Pr₄Ni₃O_{10+δ} and perovskite-like PrNiO₃. Ruddlesden-Popper $Pr_4Ni_3O_{10\pm\delta}$ is a dominant nickelcontaining product at 800-900 °C,75,224-226 while PrNiO3 forms preferentially at lower temperatures ≤ 700 °C.^{228,229} Although the decomposition kinetics is very slow at 580-600 °C (lowest temperature at which the oxidative decomposition was

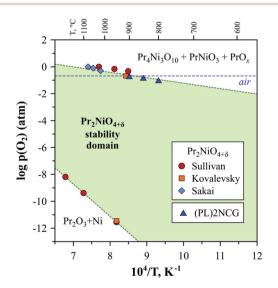


Fig. 6 Approximate stability domain of RP-type $Pr_2NiO_{4+\delta}$. The literature data on low- and high- pO_2 stability limits was taken from Sullivan, ²²² Kovalevsky⁵² and Sakai. ⁹⁵ The data on the high- pO_2 stability boundary of $(Pr_{0.9}La_{0.1})_2Ni_{0.74}Cu_{0.21}Ga_{0.05}O_{4+\delta}$ ((PL)2NCG) was from Xue. ²³⁰

observed),^{75,228,229} it accelerates with temperature, particularly above 700 °C.⁹⁶

The stability field boundary of the Pr₂NiO_{4+δ} lattice is interrelated with the overstoichiometric oxygen content. Partial substitution of nickel by copper was found to decrease δ and expand the stability field in air to lower temperatures, while doping by iron in the Ni sublattice led to the opposite effect.52 However, B-site substitutions have a rather negligible effect on the stability from a practical point of view. 52,111,118,199,230 A more effective approach, involving the suppression of excess oxygen and stabilisation of the $Pr_2NiO_{4+\delta}$ phase under oxidising conditions, is the substitution of praseodymium by alkalineearth cations. In particular, $Pr_{2-x}Ca_xNiO_{4+\delta}$ solid solutions are reported to be stable in air when $x \ge 0.3$ for at least 250 h at 850 °C.101 Moreover, no evidence of phase separation under oxidizing conditions was observed for Pr_{1.35}Sr_{0.65}Ni_{0.75}Co_{0.25}- $O_{4+\delta}$ (ref. 177) and heavily Sr-doped $Pr_{2-x}Sr_xNiO_{4+\delta}$ (x=1.0-1.6).106,107 The incorporation of cation vacancies into the praseodymium sublattice was suggested as a possible alternative to decrease the oxygen excess and expand the stability field.230 Yet another strategy involves the substitution of praseodymium by lanthanum or neodymium with the formation of $(Pr_{1-x}Ln_x)_2NiO_{4+\delta}$ solid solutions. 94-96,228,229 However, it was found that full stabilisation requires the substitution of up to 75% of praseodymium cations. 96,228,229 Although the intermediate compositions with $x \sim 0.5$ were reported to be stable in the short term (≤100 h scale),94,95 they decomposed in the course of long-term thermal treatments (a month scale). 96,228,229

Ln₂NiO_{4+δ}-derived ceramic materials demonstrate good stability in the presence of steam or carbon dioxide in the gas phase at elevated temperatures. In particular, no degradation of the electrical conductivity of $Ln_2NiO_{4+\delta}$ (Ln = La, Pr, and Nd) and LaSrNiO₄ ceramics in moist air, $pH_2O = 0.03$ atm, was observed at 25-630 °C.81 Water partial pressures of up to 0.3 atm were reported to have no detrimental effect on the conductivity of $Pr_2NiO_{4+\delta}$ ceramics at temperatures ≤ 700 °C.²³¹ Li *et al.*²³² reported the phase stability of Pr₂NiO_{4+δ} powder in a 40% steam/60% air mixture at 700 °C (24 h) and the absence of degradation of electrical conductivity of Pr₂NiO_{4+δ} ceramics in a 60% steam/40% air atmosphere at 700 °C (>7 h test). The TGA-MS experiments performed by Egger et al.86 did not reveal any significant influence of humidity or carbon dioxide on the oxygen exchange properties of $Nd_2NiO_{4+\delta}$ in air containing $\sim 1\%$ of H₂O or 5% of CO₂ at 500-900 °C. Upasen et al.²³³ exposed $Ln_2NiO_{4+\delta}$ (Ln = La, Pr, and Nd) ceramics to harsher treatment conditions, consisting of high steam pressure of 40 bar, both with and without CO2 gas dissolved in water, at 550 °C for 5 days. The presence of secondary phases originating from undesirable hydroxylation and carbonation in the course of the treatment was detected only in the near-surface layers, while the $Nd_2NiO_{4+\delta}$ ceramics exhibited the highest structural, mechanical and chemical stability under these conditions. When pure CO₂ was used as the sweep gas, the phase stability upon exposure to CO2 and the absence of oxygen permeation flux degradation through mixed-conducting Ln₂NiO_{4+δ}-based ceramic membranes were demonstrated in numerous studies conducted at 750-950 °C. 163,234-239 At the same time, excessive doping by

alkaline-earth cations into the Ln sublattice may deteriorate the stability of CO2, as was demonstrated in the case of $(La_{1-x}Ca_x)_2(Ni_{0.75}Cu_{0.25})O_{4+\delta}$ for $x \ge 0.2.^{237}$

In addition to chemical stability with respect to the components of the gas phase, electrode materials should preferably exhibit good tolerance towards poisoning by volatile contaminants originating from chromium-based interconnects and silica-based glass-ceramic sealants. The available studies show that Ln₂NiO_{4+δ} electrodes exhibit fairly good tolerance with respect to $Cr^{240-244}$ or $Cr + Si^{245,246}$ poisoning in a dry atmosphere. Lee et al. 240 reported that despite La₂NiO_{4+δ} and Cr₂O₃ readily reacting at 1000 °C with the formation of LaNiO₃ and LaCrO₃, the polarisation resistance of La₂NiO_{4+δ} electrodes at 520-800 °C is insensitive to chromium introduced solution infiltration for Cr concentrations of up to 0.6 wt%. The Nd₂NiO_{4+δ} electrodes did not show an appreciable increase in polarisation resistance in the presence of gaseous chromium species at 850 °C and current density of 500 mA cm⁻² during 800 h of testing.241 Oxygen surface exchange and electrochemical studies demonstrate the better stability of Ln₂NiO_{4+δ} compared to Sr-containing perovskite-type electrode materials^{242,243,245} due to chromium poisoning being correlated with the formation of strontium chromate at the surface. At the same time, deterioration of the electrode performance was observed due to the acceleration of Cr/Si poisoning as a result of humidity;243-246 moreover, the negative effect of humidity becomes stronger when the temperature decreased from 800 $^{\circ}\mathrm{C}$ to 700 °C.245,246

Electrical transport properties of $Ln_2NiO_{4+\delta}$ -based ceramics

4.1. Electronic conductivity

4.1.1. Undoped $Ln_2NiO_{4+\delta}$. The electrical conductivity of Ln₂NiO_{4+δ} and its derivatives is predominantly p-type electronic. Undoped $Ln_2NiO_{4+\delta}$ exhibits thermally activated conductivity in the low-temperature range, but demonstrates a smooth transition from semiconducting to metallic-like behavior above 350-500 °C.51,53,57,81,82,86,93,96,97,116,144,156,247-253 The interpretation of the apparent "semiconductor-metal" transition is controversial. Goodenough proposed the coexistence of the localized d_{z^2} electron state and itinerant $d_{x^2-y^2}$ electron states, forming a narrow in-plane $\sigma_{x^2-y^2}$ band in La₂NiO₄.^{254,255} The transition to semiconducting behavior in the lowtemperature range was interpreted in terms of splitting of the $\sigma_{x^2-y^2}$ band due to antiferromagnetic ordering, 254,255 which was not confirmed, or due to intraatomic exchange interactions with localized d_{z^2} electrons, ⁴⁵ and also was correlated with an orthorhombic-tetragonal transition.²⁵⁶ These considerations did not account for temperature-dependent oxygen nonstoichiometry. Bassat et al.51,54 noted that the apparent "semiconductor-metal" transition coincides with the onset of oxygen loss on heating, and argued that the decrease in conductivity on heating in the high-temperature range is caused by a decrease in oxygen content and the hole concentration. Later, attempts to model the temperature dependence of electrical conductivity

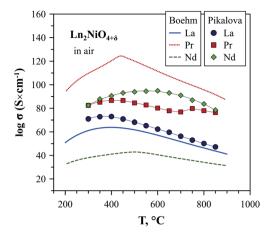


Fig. 7 Electrical conductivity of $Ln_2NiO_{4+\delta}$ ceramics in air. The data is from Boehm⁵³ and Pikalova.⁸²

showed that metallic-like behavior cannot be explained only by the decrease in the concentration of charge carriers. 251,257 In recent works, Ln₂NiO_{4+ô} nickelates are discussed as degenerate p-type semiconductors^{63,64,258} with metal-like band conduction in the high-temperature range^{63,64,142,148,258} and nearly pO₂-independent hole mobility (0.14-0.25 cm² V⁻¹ s⁻¹), which slightly decreases on heating. 60,63,142,148,258

As expected for a layered structure, the electrical properties of Ln₂NiO_{4+δ} nickelates are highly anisotropic.^{71,259-261} Bassat et al.260 demonstrated that the electrical conductivity of $La_2NiO_{4+\delta}$ single crystals along the a-b plane is more than 3 orders of magnitude higher than that along the c-axis at room temperature. Anisotropic electrical properties were obtained for textured Nd2NiO4+6 ceramics:262 the conductivity along and across the a-b plane at 500 °C was 270 and 16 S cm⁻¹, respectively, compared to 117 S cm⁻¹ for ceramics with randomly oriented grains.

The electronic conductivity of $Ln_2NiO_{4+\delta}$ is p-type and decreases with a reduction in the oxygen partial pressure due to oxygen release from the lattice and corresponding decrease in the concentration of electron holes according to eqn (2).52,57,63,139,140,142,148,258,263 Although there is some scattering in the reported conductivity values, which is apparently due to different synthesis techniques and porosity of ceramics, the electrical conductivity of Ln₂NiO_{4+δ} ceramics generally vary in of 40–120 S cm⁻¹ in air at 500– 900 °C. 51,53,57,81,82,86,93,96,97,116,140,144,148,247-250,252,253 Considering the relationship between oxygen nonstoichiometry and concentration of electronic charge carriers, eqn (3), the conductivity of $\text{Ln}_2\text{NiO}_{4+\delta}$ can be expected to increase in the sequence La < Pr < Nd. This was confirmed, in particular, by Pikalova et al.,82 although different relationships were reported in some other works,53,81,96 as shown in Fig. 7.

It is necessary to highlight that due to the phase instability in air at 700-900 °C, the reported data on the transport properties of $Pr_2NiO_{4+\delta}$ in this temperature range are mostly compromised. In fact, phase decomposition to a mixture of $Pr_4Ni_3O_{10+\delta}$ and PrO_x on cooling from higher temperatures was found to result in a sharp increase in electrical conductivity due to the superior conductivity of RP-type $\text{Pr}_4\text{Ni}_3\text{O}_{10\pm\delta}$ nickelate. Thus, one may expect that reliable data on the electrical properties of undoped $\text{Pr}_2\text{NiO}_{4+\delta}$ in air include measurements done using fresh samples within the stability domain (\geq 925 °C in air) or quenched samples in a metastable region (up to 600–700 °C in air). The same considerations relate to solid solutions derived from $\text{Pr}_2\text{NiO}_{4+\delta}$ and their respective stability fields.

4.1.2. Effect of A-site deficiency. Bassat *et al.*^{51,264} reported that the electrical conductivity of $\text{La}_{2-x}\text{NiO}_{4\pm\delta}$ (x=0–0.15) gradually decreases with an increase in the concentration of A-site cation vacancies. In the case of $\text{Nd}_{2-x}\text{NiO}_{4\pm\delta}$, the conductivity in air was found to increase in the sequence x=0 < 0.10 < 0.05, ⁵³ although the conductivity values for the parent material (x=0) were much lower compared to other reports. Therefore, it can be concluded that A-site deficiency is charge-compensated by a decrease in both oxygen content and, to a lesser extent, hole concentration (eqn (5)) and results in a decline in electronic transport.

4.1.3. Mutual solid solutions. In mutual solid solutions, the electrical conductivity generally varies between that of the end members of the series. Increasing praseodymium content in $\text{La}_{2-x}\text{Pr}_x\text{NiO}_{4+\delta}$ was found to enhance the conductivity^{93,265} sharply until x=1.0, and then weakly in the $1.0 \le x \le 2.0$ range. Electronic transport in the $(\text{Pr}_{1-x}\text{Nd}_x)_2\text{NiO}_{4+\delta}$ system was reported to decline gradually with an increase in neodymium fraction.⁹⁶ The electrical conductivity of the $\text{La}_{2-x}\text{Sm}_x\text{NiO}_{4+\delta}$ (x=0-1.1) solid solution tended to increase with an increase in samarium content, although the observed variations of σ vs. x were not very systematic, and the highest conductivity was obtained for x=0.5.⁹⁷

4.1.4. A-site doping by alkaline-earth metal cations. As a general trend, an increase in the content of alkaline-earth metal cation in the A sublattice of Ln_{2-x}A_xNiO_{4+δ} is accompanied by a decrease in the oxygen content, δ , and an increase in the electron-hole concentration according to eqn (6). Consequently, the electronic conductivity in $Ln_{2-x}A_xNiO_{4+\delta}$ (Ln = La, Pr, Nd; A = Ca, Sr, Ba) increases with x within the solid solution formation range (Fig. 3). This trend was observed in the high-temperature range for $\text{La}_{2-x}\text{Sr}_{x}\text{NiO}_{4+\delta_{3}}^{81,133,139,141,142,154,266,267}$ ${\rm La_{2-x}Ca_{x}NiO_{4+\delta},}^{82,99,100,130,131}$ $\text{La}_{2-x}\text{Ba}_x\text{NiO}_{4+\delta_1}^{266} \quad \text{Pr}_{2-x}\text{Ca}_x\text{NiO}_{4+\delta_2}^{82,102,267} \quad \text{Pr}_{2-x}\text{Sr}_x\text{NiO}_{4+\delta_2}^{57,144,268}$ $Nd_{2-x}Ca_xNiO_{4+\delta}$, 53,78,79,82,132 $Nd_{2-x}Sr_xNiO_{4+\delta}$, 79,148 $Nd_{2-x}Ba_xNiO_{4+\delta}$ (ref. 79) systems. Similar to the parent compounds, substituted Ln_{2-x}A_xNiO_{4+δ} exhibit a transition from low-temperature semiconducting to high temperature metallic-like behavior on heating. 53,57,78,79,82,98-100,102,105,131-133,136,144,253,266-268 For Srdoped (including Sm-based)143 systems, this "semiconductormetal transition" was observed to shift to lower temperatures with an increase in the strontium content, and the compositions with show metallic-like behavior above 0.8 - 1.0 $temperature.^{79,81,98,105,133,136,137,143,146,147,152,268}\\$

In the case of Sr-rich compositions ($x \ge 1.0$), the electronhole concentration and electronic conductivity in the high-temperature range reaches the maximum for x = 1.2 (Fig. 3), with conductivity values in the range of 340–470 S cm⁻¹ at 700 °C, and then declines due to an increase in oxygen vacancy concentration. ^{106–108,145} However, the strong anisotropic

expansion of the ${\rm Ln_{2-x}Sr_xNiO_{4+\delta}}$ (x > 1.0) crystal lattice results in microcracking effects in ceramics and, consequently, in an apparent hysteresis in dilatometric and conductivity curves on temperature cycling. ^{106,108,138,145}

Similar to undoped nickelates, $Ln_{2-x}A_xNiO_{4+\delta}$ are p-type electronic conductors. Variations in the electrical conductivity as a function of the oxygen partial pressure is directly correlated with the changes in oxygen nonstoichiometry. A decrease in oxygen content with a reduction in pO2 is accompanied by a decrease in electrical conductivity for compositions with a moderate ($x \le 0.3$) and large ($x \ge 1.0$) substitutional level, while solid solutions with intermediate dopant contents demonstrate nearly pO2-independent oxygen nonstoichiometry ($\delta \sim 0$) and electronic conductivity at 600–900 °C and pO_2 in the range of 10^{-5} to 1.0 atm. 57,106,108,139,142,145,148,154,269 A reduction of Sr-rich $\text{La}_{2-x}\text{Sr}_x\text{NiO}_{4+\delta}\ (x \ge 1.0)$ at $p\text{O}_2 \sim 10^{-5}$ atm and associated structural changes were observed to result in a localization of electronic charge carriers and a transition to semiconducting behavior at $T \le 1000$ °C under mildly reducing conditions.152,155

Shen *et al.*²⁶⁶ reported that the electrical conductivity of $\text{La}_{2-x}\text{A}_x\text{NiO}_{4+\delta}$ (x=0.1,0.3) is higher for A=Sr compared to A=Ba, while the electronic transport in $\text{La}_{1.7}\text{A}_{0.3}\text{NiO}_{4+\delta}$ was found to increase in the sequence $A=\text{Sr}\sim\text{Ba}<\text{Ca.}^{253,267}$ It should also

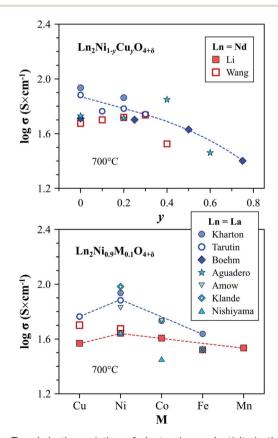


Fig. 8 Trends in the variation of electronic conductivity in the Ln₂-Ni_{1-y}Cu_yO_{4+ δ} and Ln₂Ni_{0.9}M_{0.1}O_{4+ δ} series at 700 °C in air. Literature sources: Kharton,^{60,214} Tarutin,¹⁵⁷ Boehm,¹¹³ Aguadero,¹¹⁴ Amow,⁹⁷ Klande,²⁷⁰ Nishiyama,²⁵⁰ Wang,¹¹⁶ and Li.¹¹⁹

be noted that the conductivity of $Sm_{1.4}Sr_{0.6}NiO_{4+\delta}$ ceramics was reported to be only slightly lower compared to their Prbased analog,73 which seems to imply that the trends in the $\mathrm{Sm}_{2-x}\mathrm{Sr}_x\mathrm{NiO}_{4+\delta}$ system should be similar to other A-site-doped series.

B-site doping by transition metal cations. Substitu-4.1.5. tion of nickel by copper generally results in a gradual moderate decrease in electrical conductivity in La₂Ni_{1-ν}Cu_νO_{4+δ} (ref. 112–114 and 157) and $Pr_2Ni_{1-\nu}Cu_{\nu}O_{4+\delta}$ (ref. 45 and 52) systems within the solid solution formation range (Fig. 8), while a minor increase in the conductivity with an increase in copper content was reported for the $Nd_2Ni_{1-\nu}Cu_{\nu}O_{4+\delta}$ series.¹¹⁶ All the $Ln_2Ni_{1-\nu}Cu_{\nu}O_{4+\delta}$ exhibit a "metal-insulator" transition on heating in air, with metalliclike behavior in the high-temperature range. Analysis of the pO_2 -T- δ diagram of La₂Ni_{0.8}Cu_{0.2}O_{4+ δ} employing a statistical thermodynamic approach showed that the lattice sites occupied by nickel and copper cations are essentially equivalent from an energetic point of view. 61,62,65 Therefore, the decline in electronic transport with copper doping should be attributed to a decrease in oxygen nonstoichiometry and hole concentration.

Although substitution by cobalt or iron has a contrary effect on the δ and concentration of electron holes, as shown in eqn (7), electronic conductivity of $\text{Ln}_2\text{Ni}_{1-\nu}\text{Co}_{\nu}\text{O}_{4+\delta}$ (ref. 60, 97, 118, 119 and 250) and $Ln_2Ni_{1-\nu}Fe_{\nu}O_{4+\delta}$ (ref. 52, 60, 119 and 270) decreases on doping. This is attributed to the hole localization and trapping of p-type charge carriers by the dopant cations, forming stable Co³⁺ and Fe³⁺ states. 60-62,121,271

Upon comparing the effect of different transition metal dopants, the electrical conductivity at 700 °C in air was reported to decrease from 44 to 33 S cm⁻¹ in the sequence M = Ni > Co >Cu > Fe \sim Mn for the Nd₂Ni_{0.9}M_{0.1}O_{4+ δ} series¹¹⁹ and from 87 to 44 S cm⁻¹ in the sequence M = Ni > Co > Fe for La₂Ni_{0.9}M_{0.1}O_{4+ δ} solid solutions⁶⁰ (Fig. 8). As in the case of the parent praseodymium nickelate, the data on the electrical conductivity of the doped Pr₂NiO_{4+δ} can be strongly affected by the instability of the lattice under oxidizing condition: the decomposition into $Pr_4Ni_3O_{10+\delta}$ and PrO_x phases was observed to result in a sharp increase in the conductivity of Pr2Ni0.9Fe0.1O4+8 and Pr₂Ni_{0.9}Cu_{0.1}O_{4+δ} on cooling in air.⁵²

B-site doping by transition metal cations is also known to suppress electronic transport in A-site co-doped compositions. The corresponding high-temperature conductivity data is available for $La_{1.9}Sr_{0.1}Ni_{1-\nu}Fe_{\nu}O_{4+\delta}$ (y = 0.02-0.10), 272,273 $La_{2-x}Sr_xNi_{1-y}Fe_yO_{4+\delta}$ (x = 0.5-0.8, y = 0-0.5), ^{168,172}

Table 2 Ionic conductivity of $Ln_2NiO_{4+\delta}$ -based ceramics in air. Visualisation of these data is presented in Fig. S1, ESI

	$\sigma_{\rm O}$, S cm ⁻¹					
Composition	900 °C	800 °C	700 °C	$t_{ m O} \left(800~^{\circ}{ m C}\right)$	$E_{\rm a}{}^{b}$, eV	Ref.
${ m La_2NiO_{4+\delta}}$	_	0.037	0.015	$8.0 imes 10^{-4}$	0.87	53, 85 and 277
$\text{La}_2 ext{NiO}_{4+\delta}$	0.130	0.053	0.013	7.0×10^{-4}	1.15	263
$\text{La}_2 \text{NiO}_{4+\delta}{}^a$	0.062	0.046	_	$8.4 imes 10^{-4}$		278
$\text{La}_2 ext{NiO}_{4+\delta}$	_	0.010	3.7×10^{-3}	1.6×10^{-4}	1.07	100
$\text{La}_{1.95} \text{NiO}_{4+\delta}$	_	0.019	8.7×10^{-3}	_	0.77	85
$\text{La}_{1.9}\text{Sr}_{0.1}\text{NiO}_{4+\delta}$	_	0.012	_	1.3×10^{-4}		154
$\text{La}_{1.9}\text{Sr}_{0.1}\text{NiO}_{4+\delta}{}^a$	0.017	0.010	_	7.8×10^{-5}	0.63	279
$\text{La}_{1.5}\text{Sr}_{0.5}\text{NiO}_{4+\delta}$	3.2×10^{-4}	$5.8 imes 10^{-5}$	_	_	1.90	280
$La_{1.5}Sr_{0.5}Ni_{0.8}Co_{0.2}O_{4+\delta}$	1.9×10^{-4}	3.6×10^{-5}	$6.0 imes 10^{-6}$	_	1.69	280
$La_{1.5}Sr_{0.5}Ni_{0.6}Co_{0.4}O_{4+\delta}$	1.3×10^{-4}	3.4×10^{-5}	4.8×10^{-6}	_	1.86	280
$La_{1.5}Sr_{0.5}Ni_{0.9}Fe_{0.1}O_{4+\delta}$	3.6×10^{-4}	1.5×10^{-4}	_	1.2×10^{-6}	1.21	172 and 281
$La_{1.5}Sr_{0.5}Ni_{0.6}Fe_{0.4}O_{4+\delta}$	9.7×10^{-4}	2.1×10^{-4}	3.7×10^{-5}	1.3×10^{-5}	1.58	172 and 281
$\text{La}_{1.9}\text{Ca}_{0.1}\text{NiO}_{4+\delta}{}^a$	0.038	0.022	_	1.7×10^{-4}	0.74	279
$\text{La}_{1.9}\text{Ca}_{0.1}\text{NiO}_{4+\delta}$	_	0.022	9.5×10^{-3}	2.8×10^{-4}	0.86	100
$\text{La}_{1.7}\text{Ca}_{0.3}\text{NiO}_{4+\delta}$	_	4.9×10^{-3}	2.7×10^{-3}	4.2×10^{-5}	0.62	100
$\text{La}_{1.7}\text{Ca}_{0.3}\text{NiO}_{4+\delta}$	_	$8.4 imes 10^{-4}$	3.0×10^{-4}	1.0×10^{-4}	1.04	282
$\text{La}_{2}\text{Ni}_{0.5}\text{Cu}_{0.5}\text{O}_{4+\delta}$	_	_	$7 imes 10^{-3}$	_	_	283
$\text{La}_{2} ext{Ni}_{0.9} ext{Fe}_{0.1} ext{O}_{4+\delta}$	0.258	_	_	_	_	60
$\mathrm{La_{2}Ni_{0.95}Al_{0.05}O_{4+\delta}}^{a}$	0.109	0.085	_	2.5×10^{-3}	_	190
$\text{La}_{2}\text{Ni}_{0.95}\text{Al}_{0.05}\text{O}_{4+\delta}{}^{a}$	2.7×10^{-3}		_	_	_	284
$\text{Pr}_2 \text{NiO}_{4+\delta}$	_	0.047	0.024	$5.0 imes 10^{-4}$	0.70	53 and 277
$\text{Pr}_2 \text{NiO}_{4+\delta}$	_	0.057	0.026	7.3×10^{-4}	0.84	282
$Pr_{1.7}Ca_{0.3}NiO_{4+\delta}$	_	7.9×10^{-4}	3.0×10^{-4}	9.4×10^{-6}	0.96	282
$Pr_{1.9}Ni_{0.75}Cu_{0.25}O_{4+\delta}$	_	0.075	0.040	_	0.65	276
$Pr_{1.9}(Ni_{0.75}Cu_{0.25})_{0.95}Ga_{0.05}O_{4+\delta}$	_	0.099	0.061	_	0.59	276
$Nd_2NiO_{4+\delta}$	_	0.045	0.015	1.3×10^{-3}	1.06	53 and 277
$Nd_2NiO_{4+\delta}{}^a$	2.0×10^{-3}	1.1×10^{-3}	4.4×10^{-4}	_	0.76	286
$Nd_2NiO_{4+\delta}$	_	0.076	0.017	9.1×10^{-4}	1.46	282
$Nd_{1.7}Ca_{0.3}NiO_{4+\delta}$	_	6.8×10^{-4}	2.2×10^{-4}	6.0×10^{-6}	1.08	282
$Nd_{1.9}Ba_{0.1}NiO_{4+\delta}$	_	0.040	$9.4 imes 10^{-3}$	4.1×10^{-4}	1.28	207
$Nd_{1.9}Ba_{0.1}NiO_{4+\delta}F_{0.05}$	_	0.081	0.023	9.7×10^{-4}	1.17	207

 $^{^{}a}$ pO₂ = 0.1 atm. b Calculated using Arrhenius model, $\sigma = (A_0/T)\exp(-E_a/(kT))$.

 $\begin{array}{lll} \text{LaSrNi}_{1-y}\text{Mn}_y\text{O}_{4-\delta} \left(y=0\text{-}0.3\right),^{171} \text{La}_{1.8}\text{Sr}_{0.2}\text{Ni}_{1-y}\text{Mo}_y\text{O}_{4+\delta} \left(y=0\text{-}0.05\right),^{165} & \text{La}_{1.6}\text{Ca}_{0.4}\text{Ni}_{1-y}\text{Cu}_y\text{O}_{4+\delta} & \left(y=0\text{-}0.3\right),^{131} \text{La}_{1.6}\text{Ca}_{0.4}\text{Ni}_{1-y}\text{Fe}_y\text{O}_{4+\delta} & \left(y=0\text{-}0.2\right),^{131} \text{Nd}_{0.8}\text{Sr}_{1.2}\text{Ni}_{1-y}\text{M}_y\text{O}_{4-\delta} & \left(\text{M}=\text{Co}, \text{Fe}, y=0\text{-}0.2\right),^{274} & \text{La}_{2-2x}\text{Sr}_{2x}\text{Ni}_{1-x}\text{Ti}_x\text{O}_{4+\delta},^{45} & \text{and} \\ \text{La}_{2-2x}\text{Sr}_{2x}\text{Ni}_{1-x}\text{Mn}_x\text{O}_{4+\delta} & \left(x=0\text{-}0.3\right),^{169} & \text{systems}. & \text{On the contrary}, \\ \text{Shen et $al.$}^{275} & \text{reported an increase in electrical conductivity with an increase in copper content in $\text{La}_{1.7}\text{Ca}_{0.3}\text{Ni}_{1-y}\text{Cu}_y\text{O}_{4+\delta} & \left(y=0.1\text{-}0.3\right)$ series.} \end{array}$

4.1.6. B-site doping by other metal cations. The incorporation of metal cations with a stable oxidation state, such as Mg^{2^+} and Al^{3^+} , into the nickel sublattice of $\mathrm{Ln_2NiO_{4+\delta}}$ reduces the hole concentration, breaks the $\mathrm{Ni}^{n^+}\text{-}\mathrm{O}^{2^-}\text{-}\mathrm{Ni}^{n^+}$ chains, and therefore decreases the electronic conductivity, 38,190 although the hole mobility in $\mathrm{La_2Ni_{0.95}Al_{0.05}O_{4+\delta}}$ was reported to be slightly higher compared to undoped $\mathrm{La_2NiO_{4+\delta}}$. The same trends in the variation of electrical conductivity were observed for co-doped compositions including $\mathrm{LaSrNi_{1-y}Sc_yO_{4+\delta}}$ $(y=0-0.2)^{192}$ and $\mathrm{Pr_{2-x}(Ni,Cu)_{1-y}Ga_yO_{4+\delta}}$ (y=0-0.10). 198,276

4.2. Oxygen-ionic transport

 ${\rm Ln_2NiO_{4+\delta}}$ -based oxide materials are considered mixed ionic-electronic conductors. However, the oxygen-ionic conductivity $(\sigma_{\rm O})$ in these phases is 3 or more orders of magnitude lower compared to electronic conductivity at 500–1000 °C (Table 2). In addition to $\sigma_{\rm O}$, the important parameters used to describe oxygen transport processes include the oxygen diffusion coefficient, D, and oxygen surface exchange coefficient, k (Table 3). The diffusion coefficient is the measure of the oxygen diffusion rate in the oxide bulk and is interrelated with the oxygen-ionic conductivity via the Nernst-Einstein equation (for predominantly electronic conductors) as follows: 190,286,299,307

$$\sigma_{\rm O} = \frac{4F^2c_{\rm O}D_{\rm O}}{RT} \tag{8}$$

where c_0 is the concentration of oxygen ions and D_0 is the oxygen self-diffusion coefficient, while the surface exchange coefficient describes the kinetics of oxygen exchange at the oxide/gas interface.

Three main experimental techniques are employed to study oxygen transport in RP-type nickelates:

(i) Isotope exchange depth-profiling with secondary ion mass-spectrometry (IEDP/SIMS) is based on analysis of the concentration profile of traced $^{18}{\rm O}$ ionic species in an oxide sample (polycrystalline ceramics, $^{160,276,290,292,294,298-301}$ single crystals 277,302,303 or thin films) 304 after exposure to an environment enriched with an $^{18}{\rm O}$ isotopic tracer ($^{18}{\rm O}_2/^{16}{\rm O}_2$ gas mixtures) at different temperatures. This method yields the oxygen tracer diffusion (D^*) and surface exchange (k^*) coefficients. The oxygen tracer and self-diffusion coefficients are related as follows: 85,282,283,289,299

$$D^* = f \times D_{\Omega} \tag{9}$$

where f is the tracer correlation factor (or, more generally, the Haven ratio). The values of f are ≤ 1 and depend on the mechanism of diffusion. In the case of RP-type $\operatorname{Ln_2NiO_{4+\delta}}$ nickelates, it is typically assumed that $f \sim 1.^{85,282,283,289,299}$

(ii) The electrical conductivity relaxation (ECR) method involves the analysis of the transient change in the conductivity of the sample (ceramics $^{86,87,92,118,227,286,289,299,305-307}$ or thin films 308,309) induced by an instant change in pO_2 at a constant temperature. The oxygen chemical diffusion (\tilde{D}_O or $D_{\rm chem}$) and surface exchange ($k_{\rm chem}$) coefficients are obtained by fitting the relaxation data. The chemical diffusion coefficient, $D_{\rm chem}$, is interrelated with the self-diffusion coefficient, D_O , via the thermodynamic enhancement factor, γ_O : 190,286,289,299,306,307

$$\gamma_{\rm O} = \frac{1}{2} \frac{\partial \ln p O_2}{\partial \ln c_{\rm O}} = \frac{D_{\rm chem}}{D_{\rm O}} \tag{10}$$

which can be obtained from the analysis of the δ -pO₂ dependence.

(iii) Oxygen permeability (OP) studies represent the measurements of oxygen permeation flux, j_{O_2} , through a dense ceramic membrane under an applied gradient of oxygen chemical potential. The j_{O_2} values obtained under similar conditions can be used as a measure for the rough comparative analysis of the oxygen transport in mixed-conducting ceramics. 162,163,185,194,195 The oxygen-ionic conductivity, σ_{O} , oxygen self-diffusion coefficient, D_{O} , and surface exchange coefficient, k_0 , can be extracted from the oxygen permeation data employing the Wagner equation for bulk diffusion and considering surface exchange kinetics. 65,154,263,280,281,283,297,301,310,311 For a mixed conductor with predominately p-type electronic conductivity, the exchange coefficients obtained by different methods are related by:^{286,306,307}

$$k_{\rm O} \approx k^* \approx \frac{k_{\rm chem}}{\gamma_{\rm O}}$$
 (11)

The other less commonly used experimental techniques are the electrochemical ion/electron blocking cell methods, $^{100,207,279-281,284,312}$ mass 92 or lattice parameter 92,305,313,314 relaxation after instant $p(O_2)$ change, analysis of electrochemical impedance spectroscopy (EIS) and distribution of relaxation times (DRT) data, 169,296,315,316 and different methods based on the oxygen isotope exchange with the gas phase, including isotope exchange with the gas phase equilibration (IE-GPE), 130,291,317,318 pulse isotopic exchange (PIE), 227,319 and temperature-programmed isotope exchange (TPIE) with $^{18}O_2$ or $C^{18}O_2$, 82,102,276,282,317

4.2.1. Mechanisms of oxygen-ion transport in RP-type $\mathbf{Ln_2NiO_{4+\delta}}$. Similar to the case of electronic conduction, the layered structure of $\mathbf{Ln_2NiO_{4+\delta}}$ oxides determines the strong anisotropy of oxygen-ionic transport in these phases. Studies on $\mathbf{Ln_2NiO_{4+\delta}}$ single crystals and oriented thin films by IEDP/SIMS at 340–900 °C demonstrated that the oxygen diffusivity along the a-b plane is ~ 3 orders of magnitude faster than along c axis. $^{277,302-304}$ The D^*_{a-b} values were found to be comparable (slightly higher) to that of the bulk polycrystalline samples, implying that oxygen transport in $\mathbf{Ln_2NiO_{4+\delta}}$ ceramics is governed by the oxygen diffusion along the basal plane. A large anisotropy was also observed for the surface exchange coefficient with higher k values for the a-b plane. $^{277,302-304}$

Table 3 Oxygen surface exchange and diffusion processes and the corresponding apparent activation energy values for $Ln_2NiO_{4+\delta}$ -based compounds. Visualisation of these data is presented in Fig. S2, ESI

		$D_{\rm chem}$, cm ² s ⁻	1	$k_{\rm chem}$, cm s ⁻¹		$E_{\rm a}$, eV		
Composition	Method ^a	600 °C	700 °C	600 °C	700 °C	$D_{ m chem}$	$k_{ m chem}$	Ref.
La ₂ NiO _{4+δ}	ECR	3.9×10^{-5}	8.3×10^{-5}	2.8×10^{-5}	2.6×10^{-5}	0.53	1.36	287
$\text{La}_2 \text{NiO}_{4+\delta}$	EIS	_	_	_	8.6×10^{-7}	_	1.21	169
$\text{La}_2 \text{NiO}_{4+\delta}$	ECR	_	2.1×10^{-6}	_	1.2×10^{-6}	2.82	1.99	258
$\text{La}_2 ext{NiO}_{4+\delta}$	EIS	_	_	1.6×10^{-7}	8.6×10^{-7}	_	1.21	288
$\text{La}_2 \text{NiO}_{4+\delta}$	ECR	$4.4 imes 10^{-6}$	2.1×10^{-5}	1.2×10^{-4}	$8.2 imes 10^{-4}$	1.06	1.68	289
$\text{La}_{2}\text{Ni}_{0.95}\text{Al}_{0.05}\text{O}_{4.025+\delta}$	ECR	7.9×10^{-6}	2.1×10^{-5}	1.7×10^{-5}	4.4×10^{-5}	0.92	0.84	190
$\text{La}_{1.8}\text{Sr}_{0.2}\text{NiO}_{4+\delta}$	EIS	_	_	4.0×10^{-8}	2.7×10^{-7}	_	1.31	288
$La_{1.6}Sr_{0.4}NiO_{4+\delta}$	EIS	_	_	_	1.7×10^{-7}	_	1.40	288
$La_{1.8}Sr_{0.2}Ni_{0.95}Mo_{0.05}O_{4+\delta}$	ECR	8.0×10^{-6}	12.6×10^{-6}	_	2.0×10^{-4}	0.51	1.31	165
$Nd_2NiO_{4+\delta}$	ECR	_	_	3.2×10^{-8}	1.3×10^{-7}	_	1.10	286
$Nd_2NiO_{4+\delta}$	ECR	_	_	_	1.2×10^{-4}	_	_	295
$Nd_2NiO_{4+\delta}$	EIS	1.8×10^{-8}	$5.6 imes 10^{-8}$	3.0×10^{-6}	1.8×10^{-7}	0.90	1.43	296
$Pr_2NiO_{4+\delta}$	ECR	_	_	2.9×10^{-6}	4.7×10^{-5}	_	1.84	118
$\text{Pr}_{2}\text{Ni}_{0.9}\text{Co}_{0.1}\text{O}_{4+\delta}$	ECR	_	_	3.1×10^{-5}	9.9×10^{-5}	_	0.83	118

		D^* , cm ² s ⁻¹		k^* , cm s ⁻¹		$E_{\rm a}$, eV		
Composition	Method ^a	600 °C	700 °C	600 °C	700 °C	D^*	<i>k</i> *	Ref.
$\mathrm{La_2NiO}_{4+\delta}$	IEDP/SIMS	1.3×10^{-8}	3.0×10^{-8}	2.8×10^{-8}	1.9×10^{-7}	0.59	1.24	160
$\text{La}_2 \text{NiO}_{4+\delta}$	IE-GPE	_	5.5×10^{-10}	_	7.7×10^{-5}	1.04	1.04	82
$\text{La}_2 \text{NiO}_{4+\delta}$	IE-GPE	1.2×10^{-9}	1.1×10^{-8}	$1.5 imes 10^{-8}$	1.4×10^{-7}	1.34	1.39	130
$\text{La}_2 \text{NiO}_{4+\delta}$	IEDP/SIMS	_	3.4×10^{-8}	_	1.8×10^{-7}	0.85	1.61	290
$\text{La}_2 \text{NiO}_{4+\delta}$	IE-GPE	1.1×10^{-9}	9.6×10^{-9}	_	1.4×10^{-8}	1.41	1.43	291
$La_2NiO_{4+\delta}$	IEDP/SIMS	_	4.8×10^{-8}	_	2.1×10^{-7}	0.88	0.42	53
$\mathrm{La_{2}Ni_{0.9}Fe_{0.1}O_{4+\delta}}$	IEDP/SIMS	_	3.5×10^{-8}	_	4.5×10^{-7}	0.99	1.17	53
$\text{La}_{2}\text{Ni}_{0.9}\text{Co}_{0.1}\text{O}_{4+\delta}$	IEDP/SIMS	$\boldsymbol{1.0\times10^{-8}}$	2.6×10^{-8}	$5.9 imes 10^{-7}$	6.4×10^{-7}	0.69	0.57	160
$La_2Ni_{0.8}Co_{0.2}O_{4+\delta}$	IEDP/SIMS	1.3×10^{-8}	4.0×10^{-8}	7.3×10^{-7}	1.4×10^{-6}	0.62	0.50	160
$\text{La}_2 \text{Ni}_{0.5} \text{Co}_{0.5} \text{O}_{4+\delta}$	IEDP/SIMS	2.6×10^{-8}	4.7×10^{-8}	7.1×10^{-7}	$1.4 imes 10^{-6}$	0.62	0.26	160
$La_2Ni_{0.75}Cu_{0.25}O_{4+\delta}$	IEDP/SIMS	_	3.5×10^{-8}	_	1.3×10^{-6}	0.62	1.24	113
$\text{La}_2 \text{Ni}_{0.5} \text{Cu}_{0.5} \text{O}_{4+\delta}$	IEDP/SIMS	_	3.0×10^{-8}	_	6.0×10^{-7}	0.51	1.24	113
$\text{La}_{1.9}\text{Sr}_{0.1}\text{NiO}_{4+\delta}$	IEDP/SIMS	_	1.0×10^{-8}	_	1.7×10^{-7}	0.57	1.29	290
$\text{La}_{1.8}\text{Sr}_{0.2}\text{NiO}_{4+\delta}$	IEDP/SIMS	_	6.1×10^{-10}	_	1.2×10^{-7}	0.60	1.06	293
$\text{La}_{1.8}\text{Sr}_{0.3}\text{NiO}_{4+\delta}$	IE-GPE	_	1.4×10^{-9}	_	$6.1 imes 10^{-6}$	1.04	1.24	267
$\text{La}_{1.9}\text{Ca}_{0.1}\text{NiO}_{4+\delta}$	IE-GPE	_	1.6×10^{-10}	_	1.5×10^{-7}	1.75	0.71	291
$\text{La}_{1.7}\text{Ca}_{0.3}\text{NiO}_{4+\delta}$	IE-GPE	_	1.5×10^{-9}	_	$6.0 imes 10^{-6}$	1.04	1.24	82
$\text{La}_{1.7}\text{Ba}_{0.3}\text{NiO}_{4+\delta}$	IE-GPE	_	3.3×10^{-9}	_	1.5×10^{-5}	1.04	1.24	267
$Nd_2NiO_{4+\delta}$	IE-GPE	_	4.5×10^{-8}	_	_	1.45	_	82
$Nd_2NiO_{4+\delta}$	IEDP/SIMS	_	4.5×10^{-8}	_	3.4×10^{-7}	1.06	0.77	53
$\mathrm{Nd}_{1.95}\mathrm{NiO}_{4+\delta}$	IEDP/SIMS	_	3.5×10^{-8}	_	9.0×10^{-7}	0.65	0.94	53
$Nd_{1.8}Ca_{0.2}NiO_{4+\delta}$	IEDP/SIMS	_	1.0×10^{-8}	_	1.9×10^{-7}	0.72	0.74	53
$Nd_{1.7}Ca_{0.3}NiO_{4+\delta}$	IE-GPE	_	6.3×10^{-10}	_	_	1.09	_	82
$\text{Pr}_2 \text{NiO}_{4+\delta}$	IE-GPE	_	7.0×10^{-8}	_	$9.5 imes 10^{-5}$	0.83	1.19	82
$\text{Pr}_2 \text{NiO}_{4+\delta}$	IEDP/SIMS	_	7.2×10^{-8}	_	1.4×10^{-6}	0.73	1.34	53
$Pr_{1.7}Ca_{0.3}NiO_{4+\delta}$	IE-GPE	_	7.6×10^{-10}	_	$2.4 imes 10^{-4}$	0.93	1.35	82
$Pr_{1.9}Ni_{0.75}Cu_{0.25}O_{4+\delta}$	IEDP/SIMS	4.9×10^{-8}	1.1×10^{-7}	1.8×10^{-7}	$5.0 imes 10^{-7}$	0.64	0.77	276
$Pr_{1.9}Ni_{0.71}Cu_{0.24}Ga_{0.05}O_{4+\delta}$	IEDP/SIMS	4.3×10^{-8}		2.2×10^{-7}	_	_	0.72	297
$Pr_{1.9}(Ni_{0.75}Cu_{0.25})_{0.95}Ga_{0.05}O_{4+\delta}$	IEDP/SIMS	4.2×10^{-8}	1.8×10^{-7}	2.8×10^{-7}	1.9×10^{-6}	0.57	0.62	276
$\mathrm{Pr}_{0.85}\mathrm{La}_{0.85}\mathrm{Ca}_{0.3}\mathrm{NiO}_{4+\delta}$	IE-GPE	_	5.1×10^{-10}	_	7.6×10^{-5}	1.04	1.24	267

^a Abbreviations: ECR – electrical conductivity relaxation; EIS – electrochemical impedance spectroscopy analysis; IEDP/SIMS – isotope exchange depth profiling combined with secondary ion mass-spectrometry; and IE-GPE – isotope exchange with gas phase equilibration.

Oxygen hyperstoichiometry of $\mathrm{Ln_2NiO_{4+\delta}}$ phases under oxidizing conditions implies that oxygen migration in these materials may involve either oxygen interstitials or oxygen vacancies, or both. Furthermore, interstitial oxygen ions are highly mobile. Temperature-programmed oxygen desorption

studies demonstrate that the release of weakly bonded interstitial oxygen on heating in inert gas starts at temperatures as low as \sim 200 °C, while desorption processes associated with the formation of oxygen vacancies in the perovskite layers occur at \geq 700 °C (ref. 129, 140, 317 and 320) (except Sr-rich oxygen-

deficient ${\rm Ln_{2-x}Sr_xNiO_{4-\delta}}$ with the onset of oxygen loss from the perovskite layers at lower temperatures). This agrees with the results of the TPIE investigations, also showing that the exchange of interstitial oxygen initiates at ${\sim}200~{\rm ^{\circ}C.}^{82,102,267}$

Several mechanisms of oxygen migration in $\rm Ln_2NiO_{4+\delta}$ lattice were considered and modeled employing static lattice simulations (SLS), ^{321–324} molecular dynamics simulations (MDS), ^{324–327} and density functional theory (DFT) calculations. ^{47,328–332}

(A) *Direct interstitial diffusion* pathway involving direct jumps of interstitial oxygen ions between vacant adjacent interstitial sites within rock-salt-type LnO layers. SLS and DFT calculations showed that this pathway is characterized by an excessive energy barrier of 1.2–1.4 eV, 47,324 and therefore, is least probable compared to other mechanisms of diffusion along the a–b plane.

(B) *Interstitialcy mechanism*, also referred to as a *push-pull* or *cooperative* mechanism. In this mechanism, the oxygen interstitial displaces an apical oxygen ion from the NiO_6 octahedron, which in turn move to an adjacent oxygen interstitial site. Oxygen diffusion in the a-b plane by the interstitialcy mechanism is considered as the main pathway of oxygen transport in $\text{Ln}_2\text{NiO}_{4+\delta}$ phases. $^{321,324-327}$ The calculated E_a values for oxygen migration by this mechanism were reported to be in the range of 0.50–0.65 eV for $\text{La}_2\text{NiO}_{4+\delta}$ and $\text{Pr}_2\text{NiO}_{4+\delta}$ in the high-temperature range, 324,325,327 which is in good agreement with the oxygen diffusion activation energy values obtained experimentally in most works (Table 3). The interstitialcy diffusion pathway in the a-b plane is also supported by the results of high-temperature neutron diffraction experiments combined with the analysis by the maximum-entropy method. 328,333,334

(C) Vacancy diffusion mechanism involving oxygen vacancies in equatorial and/or apical oxygen positions in the perovskite layers. Anion Frenkel disorder, as shown by eqn (4), was predicted to be a dominant intrinsic defect formation mechanism in La₂NiO₄, with oxygen vacancies at the equatorial sites of the NiO₆ octahedra energetically more favorable than that at the apical sites.321,323,329 Cleave et al.323 performed static lattice simulations of La2NiO4 and reported that all the considered vacancy mechanisms exhibited lower activation energies than the direct interstitial process, and that oxygen anion diffusion between the equatorial positions is energetically most preferable with the calculated $E_a = 0.55$ eV. However, the formation of oxygen vacancies in the perovskite layers of oxygenhyperstoichiometric La₂NiO_{4+δ} via Frenkel disorder is likely to be suppressed due to the saturation of oxygen interstitials. The thermodynamic analysis of the pO_2 -T- δ diagrams of undoped and B-site-substituted $La_2NiO_{4+\delta}$ (ref. 61, 62 and 271) showed that the vacancy formation processes under thermodynamic equilibrium conditions are statistically insignificant when $\delta \ge$ 0. Furthermore, the MD simulations carried out on stoichiometric La₂NiO₄ and Nd₂NiO₄ demonstrated insignificant oxygen diffusivities in the high-temperature range.325,326 Nonetheless, the vacancy mechanism must be involved in oxygen diffusion across the perovskite layers (along the c axis). 277,303 A limited vacancy transport mechanism was identified along the c axis in the MD simulations of $Pr_2NiO_{4+\delta}$. Furthermore, the role of oxygen vacancies is expected to increase as a result of acceptor-type A-site substitutions in $\text{Ln}_{2-x}A_x\text{NiO}_{4+\delta}$ (A = alkaline-earth cation) accompanied by a reduction in the overall oxygen content and a gradual transition to oxygen deficiency with an increase in x (ref. 145 and 329) and with a reduction in pO_2 .³³⁵ This seems to be supported by the results of TPIE studies, which were interpreted as an indication of two distinguishable oxygen diffusion pathways in some $\text{Ln}_{2-x}A_x\text{NiO}_{4+\delta}$ nickelates.^{82,102,276,282,330}

In addition to O^{2-} transport, the involvement of interstitial peroxide O^{-} ions in different diffusion mechanisms is considered.^{277,321,329,331} In particular, the DFT calculations predict that O_{i}^{-} species may be more stable compared to O_{i}^{2-} ions and oxygen vacancies at certain strontium contents in $La_{2-x}Sr_xNiO_{4+\delta}$, ^{329,331} while the diffusion of O_{i}^{-} by the interstitialcy mechanism may be more energetically favorable. ^{321,331}

4.2.2. Undoped $\operatorname{Ln_2NiO_{4+\delta}}$. The IEDP/SIMS studies of single crystals at 450–700 °C demonstrated comparable oxygen diffusivity along the a–b plane for $\operatorname{La_2NiO_{4+\delta}}$ and $\operatorname{Pr_2NiO_{4+\delta}}$ with an activation energy of 0.67 eV for $\operatorname{Ln}=\operatorname{Pr}$, whereas higher activation energy (1.4 eV) in the case of $\operatorname{Nd_2NiO_{4+\delta}}$ results in lower D^* values at temperatures ≤ 600 °C, 302,303 as shown in Fig. 9. All the $\operatorname{Ln_2NiO_{4+\delta}}$ single crystals also have comparable k_{a-b}^* , which is 0.5–1.0 orders of magnitude higher than D_{a-b}^* . 302,303 Similar trends are generally observed for undoped polycrystalline $\operatorname{Ln_2NiO_{4+\delta}}$. The ionic conductivity of $\operatorname{Ln_2NiO_{4+\delta}}$ ceramics is around 0.05 S cm $^{-1}$ at 800 °C in air (Table 2). Furthermore, $\operatorname{La_2NiO_{4+\delta}}$ and $\operatorname{Pr_2NiO_{4+\delta}}$ ceramic membranes show similar oxygen permeability at ≥ 900 °C. 52,225

Compared to the oxygen diffusion coefficients, there is a substantially larger variation in the reported values of k and corresponding activation energies, as noted in a number of works. ^{53,300,302} It was highlighted that the scattering in measured surface exchange coefficients can be caused by the differences in surface chemistry (surface termination, near-surface rearrangement, and presence of extrinsic impurities) as a result of different thermal pre-history. ³⁰² Generally, the oxygen diffusion

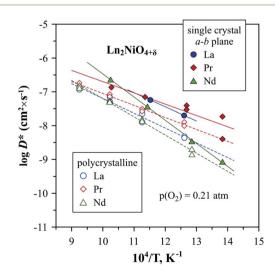


Fig. 9 Temperature dependence of oxygen tracer diffusion coefficient, D^* , reported for $Ln_2NiO_{4+\delta}$ single crystals (a-b plane) and polycrystalline samples.^{53,302,303}

and surface exchange coefficients of $Ln_2NiO_{4+\delta}$ ceramics are higher compared to many perovskite systems, such as La_{0.6}- $Sr_{0.4}Fe_{0.8}Co_{0.2}O_{3-\delta}$ and $La_{0.6}Sr_{0.4}Fe_{0.8}Ni_{0.2}O_{3-\delta}$, but lower compared the best $(La,Sr)CoO_{3-\delta}$ mixed conductors. 53,113,290,307,317

Since interstitial oxygen ions are the dominant ionic charge carriers in $Ln_2NiO_{4+\delta}$ under oxidizing conditions, the ionic conductivity, 65,263,278,279,286,307,311 oxygen diffusion cient^{278,279,301,307,311,312} and surface exchange coefficient^{286,301,307} decrease with a reduction in pO2 due to the decrease in the interstitial oxygen concentration, as shown by eqn (2).

The instability of the RP-type $Pr_2NiO_{4+\delta}$ lattice under oxidizing conditions has a strong impact on the experimentally measured ionic transport parameters in the temperature range of 700–900 °C. In particular, it has been demonstrated that the decomposition of $Pr_2NiO_{4+\delta}$ ceramics into $Pr_4Ni_3O_{10+\delta}$, PrO_x and PrNiO₃ at 750 °C in air results in an increase in the oxygen permeation flux and apparent values of oxygen diffusion and surface exchange coefficients by 1-2 orders of magnitude within 120 h.227 A similar increase in the oxygen permeability of Pr₂NiO_{4+δ} membranes with time caused by phase decomposition at temperatures below 900 °C was reported in ref. 52 and 225. On the contrary, phase separation resulted in a deterioration of the oxygen permeation flux through $(Pr_{0.9}La_{0.1})_{2-x}Ni_{0.74}Cu_{0.21}Ga_{0.05}O_{4+\delta}$ membranes at 800-900 °C.230 Berger et al.118 observed the exsolution of PrOx particles on the surface of $Pr_2NiO_{4+\delta}$ and $Pr_2Ni_{0.9}Co_{0.1}O_{4+\delta}$ ceramic samples in the course of ECR at 600-800 °C and pO₂ $= 10^{-3}$ atm, and correlated this with an increase in the measured oxygen exchange coefficients upon temperature cycling.

A-site cation deficiency. The data on ionic transport in A-site-deficient nickelates is scarce and somewhat contradictory, although one may expect a decline in oxygen diffusivity considering the trends in oxygen nonstoichiometry (see Subsection 2.2.1). Zhao et al. 85 observed that the introduction of A-site cation vacancies in $La_{2-x}NiO_{4+\delta}$ (x = 0.5) results in a slight decrease in the oxygen diffusion coefficient, D*, and ionic conductivity, but also increases the activation energy for k^* , thus suppressing oxygen exchange at temperatures below 800 °C compared to La₂NiO_{4+ δ}. In the Nd_{2-x}NiO_{4+ δ} system,⁵³ a minor neodymium deficiency (x = 0.05) was found to have a negligible effect on the oxygen diffusivity and slightly improved the oxygen exchange, while a further increase in A-site vacancy concentration up to x = 0.10 resulted in a decline in both D* and k* compared to the cation-stoichiometric Nd2- $NiO_{4+\delta}$. Sadykov et al. 317 found no substantial effect of A-site deficiency on the oxygen self-diffusion coefficients and exchange rates of $Pr_{1-x}NiO_{4+\delta}$ at 600–850 °C, but later reported more than an order of magnitude increase in $D_{\rm chem}$ and $k_{\rm chem}$ at 500 °C for $Pr_{1.9}NiO_{4+\delta}$ compared to the parent $Pr_2NiO_{4+\delta}$. The introduction of A-site cation vacancies in Nd_{2-x}Ni_{0.75}Cu_{0.25}O_{4+δ} (x = 0-0.10) was found to improve the oxygen diffusivity and oxygen permeation flux through ceramic membranes at 750-880 °C despite a decrease in oxygen content,87 which was accompanied by some decline in surface exchange rates. Similarly, the oxygen permeability of $(Pr_{0.9}La_{0.1})_{2-x}Ni_{0.74}Cu_{0.21}Ga_{0.05}O_{4+\delta}$ (x = 0

0.10) membranes was reported to increase 2 times with the introduction of A-site cation vacancies, 285,336 while oxygen excess in the lattice was suppressed.²⁸⁵ This effect was attributed to the contribution of highly mobile oxygen vacancies to ionic transport.285

4.2.4. Mutual solid solutions. Vibhu et al. 93 performed IEDP/SIMS studies of $La_{2-x}Pr_xNiO_{4+\delta}$ (x = 0.5, 1.0 and 1.5) ceramics. They found that the oxygen diffusion coefficient, D^* , of solid solutions at 500-700 °C is close to that reported for $La_2NiO_{4+\delta}$ and $Pr_2NiO_{4+\delta}$ and moderately increases with an increase in praseodymium content. The values of the surface exchange coefficient, k^* , did not show a systematic change with the composition and varied in a narrow range.

4.2.5. A-site substitutions. The studies on the oxygen transport parameters using different techniques revealed that the partial substitution of Ln3+ by alkaline-earth cations in $Ln_{2-x}A_xNiO_{4+\delta}$ generally suppresses the oxygen diffusivity, ionic conductivity, surface exchange kinetics, and oxygen permeability (e.g. Tables 2 and 3). The corresponding literature data is available for the $La_{2-x}Sr_xNiO_{4+\delta}$, 140,269,279,290,298,337,338 $\text{La}_{2-x}\text{Ca}_x\text{NiO}_{4+\delta}$, 82,100,130,279,282,291 $\text{Nd}_{2-x}\text{Ca}_x\text{NiO}_{4+\delta}$, 53,82,282 $Pr_{2-x}Ca_xNiO_{4+\delta}^{82,102,276,282,330}$ systems. In particular, the substitution of 5 and 10 at% of lanthanum by strontium in La2-x- $Sr_xNiO_{4+\delta}$ was reported to decrease the oxygen diffusion coefficient at 640-1000 °C by ca. 1 and 2 orders of magnitude, respectively. 279,290,298,337 The negative effect of acceptor-type doping on oxygen ionic transport and exchange is mainly attributed to the decrease in the concentration of ionic charge carriers (interstitial oxygen ions) with an increase in x in $Ln_{2-x}A_xNiO_{4+\delta}$ (see Subsection 2.2.3). Possible steric effects are considered as an additional factor. 53,82,102,279,282 The interstitialcy mechanism involves the movement of interstitial oxygen ions through Ln₃ triangles in the rock-salt-type layers. The substitution of Ln³⁺ by larger A²⁺ cations (e.g. La³⁺ by Sr²⁺) reduces the free volume available for interstitial migration in the LnO layers²⁷⁹ and the radius of diffusion channels through (Ln,A)₃ triangles, 53,82,102,282 thus contributing to the decline in ionic mobility.

Contrary to other works, Shen et al. 100 reported that the ionic conductivity of $La_{2-x}Ca_xNiO_{4+\delta}$ (x = 0-0.3) at 600-800 °C increases with moderate Ca doping (x = 0.1), despite the decrease in oxygen overstoichiometry, and then declines upon further doping. However, it should be noted that the value of $\sigma_{\rm O}$ obtained for undoped La₂NiO_{4+δ} in that work is lower compared to other literature data (Table 2). Zhu et al.339 found that substitution of 5 at% of lanthanum by bismuth enhances the oxygen chemical diffusion and surface exchange coefficient of $La_{1.65}Bi_{0.1}Sr_{0.25}NiO_{4+\delta}$ by a factor of 2-3 compared to $La_{1.75}$ $\mathrm{Sr}_{0.25}\mathrm{NiO}_{4+\delta}$ at 700–800 °C. This was attributed to the minor increase in oxygen excess with bismuth doping and higher oxygen diffusivity in the doped material associated with the high polarizability of bismuth cations.

4.2.6. B-site doping by transition metal cations. IEDP/SIMS and oxygen permeation studies113,283,299,306 demonstrated that the ionic transport parameters in $La_2Ni_{1-y}Cu_yO_{4+\delta}$ (y = 0-0.75) tend to deteriorate with an increase in copper concentration in the nickel sublattice, and consequently, a decrease in oxygen

interstitial concentration (see Subsection 2.2.4). La₂Ni_{0.9}Cu_{0.1}O_{4+ δ} was reported to exhibit a slightly higher oxygen self-diffusion coefficient, $D_{\rm O}$, and a similar surface exchange coefficient, k, compared to undoped La₂NiO_{4+ δ} at 800–950 °C.³⁰⁶ A further increase in copper content ($y \ge 0.25$) results in a gradual decrease in the oxygen tracer diffusion coefficient, D^* , at 490–810 °C, although the changes in oxygen diffusivity do not exceed one order of magnitude^{113,283} and moderate copper doping may still be favorable at lower temperatures of ~500 °C.¹¹³ Simultaneously, the substitution of nickel by copper (y = 0.25–0.75) was found to increase the activation energy for the exchange coefficient, k^* , thus substantially suppressing the oxygen surface exchange at temperatures below 750 °C compared with that for undoped lanthanum nickelate.¹¹³

Moderate copper doping was found to be favorable for oxygen transport in $Pr_2NiO_{4+\delta}$. In particular, Miyoshi *et al.*¹⁶² reported that the oxygen permeability of $Pr_2Ni_{1-y}Cu_yO_{4+\delta}$ (y=0.1–0.5) ceramic membranes at 600–1000 °C reached the maximum at y=0.2 and declines upon further doping. IEDP/SIMS studies revealed that the oxygen tracer diffusion coefficient, D^* , and ionic conductivity, σ_O , of copper-doped $Pr_{1.9}Ni_{0.75}Cu_{0.25}O_{4+\delta}$ at 730–800 °C is \sim 2 times higher compared to undoped praseodymium nickelate.²⁷⁶

Static lattice and molecular dynamics simulations³²⁴ showed that although the incorporation of transition metal cations with the 3+ oxidation state (such as Co3+ or Fe3+) into the nickel sublattice increases the concentration of ionic charge carriers, these dopants simultaneously tend to reduce the ionic mobility. This generally agrees with the reported experimental results. Kilner and Shaw160 performed IEDP/SIMS studies of $La_2Ni_{1-\nu}Co_{\nu}O_{4+\delta}$ ceramics and found that within the solid solution formation range under oxidizing conditions (y = 0-0.2), cobalt doping has a rather minor impact on the oxygen diffusivity with similar values of D^* for all compositions (despite the increase in δ with doping). Simultaneously, the substitution of Ni by Co decreases E_a for the surface exchange coefficient, k^* , and enhances oxygen exchange, especially at lower temperatures. 160 Similarly, the studies of co-substituted La_{1.5}Sr_{0.5}Ni_{1- ν}Co_{ν}O_{4+ δ} (ν = 0-0.4) nickelates by oxygen permeation and Hebb-Wagner polarization methods at 800-1000 °C revealed that although doping by cobalt increases the interstitial oxygen content and improves the surface oxygen exchange, it also deteriorates the oxygen diffusivity, and to a lesser extent, the ionic conductivity.280 Berger et al.118 reported based on the ECR studies at $pO_2 = 10^{-3}$ atm at 600-800 °C that doping with cobalt in $Pr_2Ni_{0.9}Co_{0.1}O_{4+\delta}$ decreases the activation energy for surface oxygen exchange and substantially improves $k_{\rm chem}$ at temperatures below 800 °C, which is ca. one order of magnitude at 600 °C. Later, they confirmed 316 by analysis of the EIS data for $Pr_2Ni_{1-\gamma}Co_{\gamma}O_{4+\delta}$ (y = 0 and 0.1) microelectrodes at 550-850 °C that the substitution by cobalt improves the oxygen exchange at T < 800 °C and $p(O_2) = 10^{-3} - 10^{-2}$ atm, while the effect was opposite at atmospheric oxygen pressure.

Substitution of 10 at% of nickel by iron has a comparatively minor effect on the oxygen transport, as revealed by IEDP/SIMS studies at 700–840 $^{\circ}$ C.²⁹⁸ Specifically, D^* and k^* of La₂Ni_{0.9}Fe_{0.1}O_{4+ δ} are similar to that of undoped lanthanum nickelate at higher temperatures (despite the higher δ), but the

oxygen diffusivity tends to decline to some extent upon cooling with respect to $La_2NiO_{4+\delta}$, and the opposite effect was observed for the surface exchange. Analysis of the oxygen permeation data for La₂Ni_{0.9}Fe_{0.1}O_{4+δ} ceramic membranes^{60,161} showed that doping by iron results in a slightly higher ionic mobility compared to $\text{La}_2\text{NiO}_{4+\delta}$ at 900-950 °C, but deteriorates the oxygen transport at lower temperatures. Furthermore, Klande et al. 163 reported that the oxygen permeability of La₂NiO_{4+δ} membranes exceeds that of La₂Ni_{0.9}Fe_{0.1}O_{4+δ} in the entire studied temperature range (750-950 °C). Gilev et al.281 studied the oxygen transport in the $La_{2-x}Sr_xNi_{1-\nu}Fe_{\nu}O_{4+\delta}$ (x = 0.5) and 0.8, y = 0.1-0.4) system employing the oxygen permeation and Hebb-Wagner polarizations techniques and found that co-substitution by iron (up to 40 at% in the nickel sublattice) enhances the oxygen permeability, bulk ionic diffusion and oxygen surface exchange, while the oxygen selfdiffusion coefficient $D_{\rm O}$ is lower compared to that of Fe-free $\text{La}_{1.5}\text{Sr}_{0.5}\text{NiO}_{4+\delta}$. Miyoshi et al. 40 reported that co-doping by iron improves the oxygen permeability of Pr₂Ni_{0.8-ν}Cu_{0.2}Fe_νO_{4+δ} (y = 0-0.20) ceramics at 600-1000 °C with the best results obtained for y = 0.05.

The positive effects of co-substitutions by transition metal cations on oxygen transport were also reported for other (La,Sr)₂(Ni,M)O_{4+δ} systems. Gómez et al. 165 performed ECR measurements of La_{1.8}Sr_{0.2}Ni_{0.95}Mo_{0.05}O_{4+δ} at 600-900 °C and found that D_{chem} is almost an order of magnitude higher than that of La_{1.8}Sr_{0.2}NiO_{4+δ} and comparable to that of undoped $La_2NiO_{4+\delta}$ (although with a two times lower activation energy, leading to a higher oxygen diffusivity at T < 700 °C), while the surface exchange coefficient, k_{chem} , exceeds that of the undoped nickelate by more than an order of magnitude. Li et al. 169 peranalysis of EIS data obtained $\text{La}_{2-2x}\text{Sr}_{2x}\text{Ni}_{1-x}\text{Mn}_x\text{O}_{4+\delta}$ electrodes (x = 0-0.30) at 600-800 °C and demonstrated that co-substitution by manganese promotes the surface exchange kinetics: the surface exchange coefficient, k, reaches the maximum for x = 0.10, but declines upon further doping.

Comparing the available oxygen permeation data for RP-type nickelate membranes with different dopants, one may also find the following trends:

- The oxygen transport in La₂Ni_{0.9}M_{0.1}O_{4+ δ} at temperatures below 850–900 °C decreases in the sequence M = Ni > Co > Fe;^{60,163}
- In the La $_2$ Ni $_{0.8}$ M $_{0.2}$ O $_{4+\delta}$ series, the oxygen transport at 600–1000 $^{\circ}$ C was reported to decrease in the sequence M = Cu > Co > Fe > Mn; 162
- The oxygen permeability of La₂Ni_{0.8}Cu_{0.2}O_{4+ δ} is slightly lower compared to that of La₂Ni_{0.9}Co_{0.1}O_{4+ δ} at 850–950 °C and similar at 700–800 °C;^{214,341}
- Oxygen permeability in the $\text{PrNi}_{0.75}\text{Cu}_{0.20}M_{0.05}\text{O}_{4+\delta}$ series at 600–1000 °C was found to decline in the sequence M = Fe > Cr > V. 185
- **4.2.7. B-site doping by other metal cations.** Jeon *et al.*¹⁹⁰ studied aluminum-doped $La_2Ni_{0.95}Al_{0.05}O_{4+\delta}$ by the ECR method at 800–1000 °C and reported a non-negligible improvement in the oxygen self-diffusion coefficient, D_O , and ionic conductivity, σ_O , compared to the literature data on

undoped La₂NiO_{4+ δ} (Table 2). Later they reinvestigated this material employing the electrochemical blocking cell method at 900–1000 °C.284 Although the updated $D_{\rm chem}$ and $k_{\rm chem}$ values were similar to that in the first report and slightly exceeded the corresponding parameters for undoped lanthanum nickelate, doping with aluminum was shown to suppress the selfdiffusion coefficient, Do, and oxygen-ionic conductivity by \sim 1.5 orders of magnitude (Tables 2 and 3). This was attributed to the steric effects, namely shrinkage of the lattice and reduction of free volume for oxygen diffusion caused by aluminum doping.284 However, it should be noted that the aluminum doping also was found to decrease the concentration of interstitial oxygen.189

Klande et al. 163 reported that partial substitution of nickel by aluminum or magnesium resulted in the lower oxygen permeability of La₂Ni_{0.9}M_{0.1}O_{4+ δ} (M = Al or Mg) ceramic membranes compared to the parent La₂NiO_{4+ δ}, although this may be partly attributed to the presence of phase impurities.

Ishihara et al. 194,195,198 performed extensive oxygen permeability screening tests at 600-1000 °C and found that the introduction of gallium into the nickel sublattice of $Ln_2Ni_{0.75}Cu_{0.25}O_{4+\delta}$ (Ln = Nd or Pr) combined with the introduction of A-site vacancies is a suitable approach to improve the oxygen permeability of nickelate membranes. The optimum composition was claimed to be $\operatorname{Ln}_{2-2x}(\operatorname{Ni}_{0.75}\operatorname{Cu}_{0.25})_{1-x}\operatorname{Ga}_x\operatorname{O}_{4+\delta}$ with x=0.05 and with better transport properties in the case of Ln = Pr. 194,195,198 However, it should be noted that it is hard to distinguish between the effects of cation deficiency and gallium doping without detailed studies. As mentioned above, the oxygen permeability of A-site-deficient $(Pr_{0.9}La_{0.1})_{1.90}Ni_{0.74}Cu_{0.21}Ga_{0.05}O_{4+\delta}$ membranes was found to be more than 2 times that of the cation-stoichiometric analog $(Pr_{0.9}La_{0.1})_2Ni_{0.74}Cu_{0.21}Ga_{0.05}O_{4+\delta}$. Hyodo et al. 276 performed IEDP/SIMS studies of $Pr_{1.9}(Ni_{0.75}Cu_{0.25})_{1-\nu}Ga_{\nu}O_{4+\delta}$ (y = 0 and 0.05) nickelates and found that gallium doping results in a slight improvement in the oxygen tracer diffusion coefficient, D^* , and ionic conductivity (~1.5 times, Table 3), and also results in some enhancement in oxygen exchange. This composition exhibits the highest ionic conductivity among the Ln2NiO4+6-based phases reported thus far.

Based on the oxygen permeation screening tests at 600-1000 °C, the following trends should also be mentioned:

- The oxygen transport in $Pr_2Ni_{0.9}M_{0.1}O_{4+\delta}$ was reported to decrease in the sequence M = Mg > Ga > Al. 162 Besides, the oxygen permeability of Pr₂Ni_{0.9}Mg_{0.1}O_{4+δ} membranes was also higher compared to that of $Pr_2Ni_{0.8}M_{0.2}O_{4+\delta}$ (M = Cu or Zn);¹⁶²
- The oxygen permeability of $PrNi_{0.75}Cu_{0.20}M_{0.05}O_{4+\delta}$ membranes decreases in the sequence Al > In > Zr, in all cases being lower compared to M = Fe;185
- The oxygen transport in $Pr_{1.90}(Ni_{0.75}Cu_{0.25})_{0.95}M_{0.05}O_{4+\delta}$ decreases in the sequence M = Ga > Al > In. 195
- 4.2.8. Substitutions in the oxygen sublattice. Tarutin et al.²⁰⁷ studied $Nd_{1.9}Ba_{0.1}NiO_{4+\delta}F_{\gamma}$ ($\gamma = 0-0.10$) nickelates employing the Hebb-Wagner electron blocking method and found that the moderate introduction of fluorine into the oxygen sublattice may be favorable for oxygen-ionic transport, leading to an increase in ionic conductivity by a factor of 2 for γ = 0.05. This was attributed to the mixed anion lattice effect.342-348

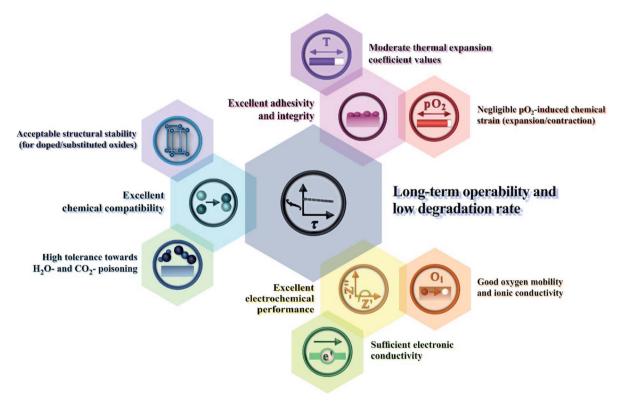


Fig. 10 Main correlations between the performance of $Ln_2NiO_{4+\delta}$ electrode materials and their functional properties.

5. $Ln_2NiO_{4+\delta}$ -based materials as electrodes for proton-conducting electrochemical cells

In this section, we discuss the results of research regarding the utilisation of $\rm Ln_2NiO_{4+\delta}\text{-}based$ materials as electrodes of protonic ceramic electrochemical cells. The success of this application depends on many factors (Fig. 10), where the inherent properties of the RP phases (considered in detail in the previous sections) and their behaviour upon contact with proton-conducting electrolytes, including chemical reactivity, thermal compatibility and electrochemical activity.

5.1. Chemical compatibility with proton-conducting electrolytes

In terms of their thermodynamic stability, $\text{Ln}_2\text{NiO}_{4+\delta}$ -based phases exhibit considerable chemical tolerance towards interaction with state-of-the-art proton-conducting materials. This may be due to two main factors, *i.e.* their wide structural flexibility upon cationic interdiffusion, and the low number of impurity phases capable of being formed at interfaces upon long-term treatment.

Lyagaeva et al. 349 studied the chemical interaction features of 13 cathode materials with $BaCe_{0.9}Y_{0.1}O_{3-\delta}$ (BCY) and $BaZr_{0.8}Y_{0.2}O_{3-\delta}$ (BZY) electrolytes using 50: 50 wt% mixtures calcined at 1100 °C for The Co-containing cathodes and comprising simple cobaltites (Ba_{0.5}Sr_{0.5}CoO_{3-δ}), layered cobaltites $(GdBaCo_2O_{5+\delta}, NdBaCo_2O_{5+\delta}, and Y_{0.8}Ca_{0.2}BaCo_4O_{7+\delta})$, simple cobaltite-ferrites (Ba_{0.5}Sr_{0.5}Co_{0.2}Fe_{0.8}O_{3- δ}, Ba_{0.5}Sr_{0.5}Co_{0.8}Fe_{0.2}O_{3- δ}, $La_{0.6}Sr_{0.4}Co_{0.2}Fe_{0.8}O_{3-\delta}$), layered cobaltite-ferrites $(GdBaCoFeO_{5+\delta} \text{ and } NdBa_{0.5}Sr_{0.5}Co_{1.5}Fe_{0.5}O_{5+\delta})$ and manganites $(La_{0.75}Sr_{0.2}MnO_{3-\delta})$, are chemically incompatible either with one of the BCY and BZY phases or with both of them. However, the studied ferrites (Ba_{0.5}Sr_{0.5}FeO_{3- δ}) and nickelates (LaNi_{0.6}Fe_{0.4}O_{3- δ} and $La_2NiO_{4+\delta}$) showed no interaction with both the BCY and BZY phases.

Tolchard and Grande³⁵⁰ applied more stringent conditions (1100 °C for 72 h) to reveal the interaction abilities of LaMnO₃, LaFeO₃, LaCoO₃, and La₂NiO₄ phases with BaZrO₃. Using XRD and SEM and EDX analyses, they found that LaMnO₃ is the most unsuitable oxide, which results in the formation of 12.2 wt% of an La₂Zr₂O₇ impurity phase according to the simplified reaction:

$$xBaZrO_3 + LaMO_3 + \frac{x}{4}O_2 \rightarrow \frac{x}{2}La_2Zr_2O_7 + La_{1-x}Ba_xMO_3$$
 (12)

where M = Mn, Co or Fe.

For the couples of LaFeO₃/BaZrO₃ and La₂NiO₄/BaZrO₃, this impurity phase was also detected in amounts of 0.7 and 2.2 wt%, respectively. The highest chemical tolerance was reported for LaCoO₃/BaZrO₃, where no La₂Zr₂O₇ was formed after prolonged high-temperature calcination. However, this does not lead to the conclusion that LaCoO3 is the most suitable for BaZrO3 since cationic interdiffusion (without decomposition of the original phases or appearance of new ones) can also occur. For example, cobalt-based solutions in a wide concentration range, *i.e.* BaZr_{1-x}Co_xO_{3- δ} (0 \leq x \leq 0.4), can be formed, 27,351 leading to significant disruption of the chemical compositions for both the basic LaCoO3 and BaZrO3 phases without the formation of any impurities. A similar situation has been described in other works, 352,353 where Ba-enriched $Ba_{0.5}Sr_{0.5}Co_{0.8}Fe_{0.2}O_{3-\delta}/Ba$ -deficient $BaCe_{0.9}Y_{0.1}O_{3-\delta}$ and Coenriched $BaCe_{0.7}Zr_{0.1}Y_{0.2}O_{3-\delta}/Y$ -enriched compositions were observed after prolonged high-temperature calcination.

Considering the model $BaCeO_3/Ln_2NiO_{4+\delta}$ pair as an example, Ni-dissolution in $BaCeO_3$ is limited to ~ 1 mol%;³⁵⁴ Ln-dissolution may be also minimised if $BaCeO_3$ is doped by acceptor dopants; Ba-dissolution in $Ln_2NiO_{4+\delta}$ is possible,³⁵⁵ but this improves the catalytic activity of the modified electrodes³⁵⁶ and is limited in the case of already Ba-doped nickelates (*i.e.* $(Ln,Ba)_2NiO_{4+\delta}$); and finally, Ce-dissolution in the $Ln_2NiO_{4+\delta}$ parent phase is negligible. These features constitute quite

Table 4 Chemical compatibility of Ln₂NiO₄-based electrodes with Ba- and La-based proton-conducting electrolytes

Electrolyte	Electrode	Weight ratio	Calcination condition, T , $^{\circ}$ C/ τ , h	Impurity phases (from XRD data)	Ref.
$BaZrO_3$	${ m La_2NiO_{4+\delta}}$	1:1	1100/72	La ₂ Zr ₂ O ₇ (trace)	350
$La_{28-x}W_{4+x}O_{54+3x/2}$	$\text{Pr}_2 \text{NiO}_{4+\delta}$	1:1	800/2	Pr ₆ O ₁₁ (trace)	357
$La_{5.5}WO_{11.25-\delta}$	$La_2NiO_{4+\delta}$	1:1	1150/5	La ₆ W ₂ O ₁₅ (trace)	265
$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	$\text{Pr}_2 \text{NiO}_{4+\delta}$	1:1	1200/1	(Ce,Pr)O ₂ and NiO	231
$BaCe_{0.7}Zr_{0.1}Y_{0.1}Yb_{0.1}O_{3-\delta}$	$\mathrm{Nd}_{2}\mathrm{NiO}_{4+\delta}$	1:1	1200/3	No interaction	358
$BaCe_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}$	$\text{La}_{1.7}\text{Ba}_{0.3}\text{NiO}_{4+\delta}$	1:1	700/1000	BaO (trace)	359
$BaCe_{0.8}Y_{0.2}O_{3-\delta}$	$La_2NiO_{4+\delta}$	1:1	1100/10	No interaction	349
$BaZr_{0.8}Y_{0.2}O_{3-\delta}$	$\text{La}_2 \text{NiO}_{4+\delta}$	1:1	1100/10	No interaction	349
$BaZr_{0.1}Ce_{0.7}Y_{0.2}O_{3-\delta}$	$Pr_{1.8}La_{0.2}Ni_{0.74}Cu_{0.21}Nb_{0.05}O_{4+\delta}$	1:1	950/3	No interaction	360
$BaCe_{0.5}Zr_{0.3}Dy_{0.2}O_{3-\delta}$	$\mathrm{Nd}_{1.95}\mathrm{Ba}_{0.05}\mathrm{NiO}_{4+\delta}$	1:1	1100/10	No interaction	361
$BaCe_{0.55}Zr_{0.3}Y_{0.15}O_{3-\delta}$	$\text{Pr}_2 \text{NiO}_{4+\delta}$	1:1	800, 900	(Ce,Pr)O $_2$ (at $T \ge 900$ °C)	362
			1000, 1100		
			1200/2		
$BaZr_{0.1}Ce_{0.7}Y_{0.2}O_{3-\delta}$	$La_{1.2}Sr_{0.8}Ni_{0.6}Fe_{0.4}O_{4+\delta}$	1:1	900/5	No interaction	363
$BaCe_{0.5}Zr_{0.3}Dy_{0.2}O_{3-\delta}$	$Pr_{1.9}Ba_{0.1}NiO_{4+\delta}$	1:1	1350/5	NiO	364
$\text{La}_2\text{Ce}_2\text{O}_{7-\delta}$	$\text{Pr}_2 \text{NiO}_{4+\delta}$	1:1	1150	Pr_6O_{11} and $La_2NiO_{4+\delta}$	365

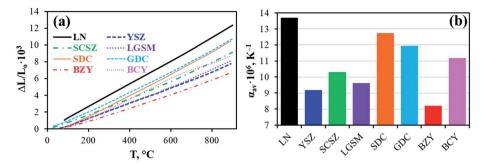


Fig. 11 Thermal expansion behaviour (a) and average TEC values (b) of some oxygen-ionic and proton-conducting electrolytes compared with LN. These data were taken from ref. 157 for LN, ref. 367 for YSZ and LGSM, ref. 368 for SDC, ref. 369 for GDC and SCSZ, and ref. 370 for BZY and $BCY. \ Abbreviations: LN = La_2NiO_{4+\delta}, \ YSZ = Zr_{0.9}Y_{0.1}O_{2-\delta}, \ LGSM = La_{0.9}Sr_{0.1}Ga_{0.8}Mg_{0.2}O_{3-\delta}, \ SDC = Ce_{0.8}Sm_{0.2}O_{2-\delta}, \ GDC = Ce_{0.9}Gd_{0.1}O_{2-\delta}, \ LGSM = La_{0.9}Sr_{0.1}Ga_{0.8}Mg_{0.2}O_{3-\delta}, \ SDC = Ce_{0.8}Sm_{0.2}O_{2-\delta}, \ LGSM = La_{0.9}Sr_{0.1}Ga_{0.8}Mg_{0.2}O_{3-\delta}, \ SDC = Ce_{0.8}Sm_{0.2}O_{2-\delta}, \ LGSM = La_{0.9}Sr_{0.1}O_{2-\delta}, \ LGSM = La_{0.9}Sr_{0.1}O_{3-\delta}, \ LGSM = La_{0.9}Sr_{0$ $SCSZ = (Sc_2O_3)_{0.1}(CeO_2)_{0.01}(ZrO_2)_{0.89}, \ BZY = BaZr_{0.8}Y_{0.2}O_{3-\delta}, \ and \ BCY = BaCe_{0.8}Y_{0.2}O_{3-\delta}.$

encouraging results regarding the chemical compatibility of nickelates with state-of-the-art proton-conducting oxides, as shown in Table 4.

5.2. Mechanical compatibility with proton-conducting electrolytes

Components of solid oxide electrochemical systems are typically subjected to severe conditions under both fabrication of the cells (very high co-sintering temperatures) and their operation (reduced temperatures, but high gradients of electrochemically active components). Thus, to achieve excellent integrity of multi-layered structures with no further unacceptable consequences (electrode delamination, cracking and other failures), materials should be selected based on their thermo-(chemico-) mechanical properties. Generally, this issue has been thoroughly described by Løken et al., 366 who reviewed the thermal and chemical expansion features of proton-conducting electrolytes together with the compatibility aspects of their use with different electrode materials. To briefly summarise their main conclusions, suitable electrodes should exhibit TEC values close to that of the widely studied protonic conductors $((8-12) \times 10^{-6} \text{ K}^{-1})$, featuring a low contribution from chemical expansion.

5.2.1. Thermomechanical expansion. A comparative analysis (Fig. 11) shows that nickelate materials are attractive electrodes, which demonstrate TECs closest to that of various solid oxide electrolytes, including proton-conducting materials. Although the TEC difference comprises about 2-3 respective units for a BCY/LN couple and 5-6 respective units for a BZY/LN couple, further adjustment of their thermal compatibility can be realised through the design of composite materials. These composites (combination of an electrode and an electrolyte with 10-50 wt% of the latter phase) have decreased TECs that follow an additive rule.

$$\alpha_{\text{composite}} = (1 - x)\alpha_{\text{electrode}} + x\alpha_{\text{electrolyte}}$$

where *x* is the weight fraction.

Together with composite materials, the TEC levels of Ln2- $NiO_{4+\delta}$ -based materials can be widely tuned via the doping

approach. Table 5 lists some doped oxides and their average TECs determined within the entire studied temperature range.

5.2.2. Chemical expansion. Fairly low absolute levels of TECs achieved for Ln₂NiO_{4+δ}-based ceramics are combined with the very weak temperature dependence of TECs (as a differential form of dilatometry curve, Fig. 12). This implies that the chemical-induced effects (or an orthorhombic/tetragonal phase transition, if existing) insignificantly affect the overall expansion of nickelate materials compared to some widely studied Co- and Fe-based oxides. The latter can suffer from phase transition(s) or chemical-induced expansion, either leading to the corresponding local TEC burst within a narrow temperature range or increase in TEC in the high-temperature range compared with that at lower temperatures, as shown in Fig. 12b. For example, GdBaCo₂O_{5+δ} exhibits a close-to-linear dilatometry curve with a TEC of $(20 \pm 1) \times 10^{-6} \text{ K}^{-1}$; however, the structure of this double cobaltite is transformed from orthorhombic to tetragonal between 450 °C and 500 °C.386,387 At temperatures above 300 °C, Ba_{0.5}Sr_{0.5}Co_{0.8}Fe_{0.2}O_{3-δ} has a cubic-type perovskite structure,388-390 but chemical expansion occurs upon heating from 500 °C, increasing the TEC value by 2 times from $(11 \pm 1) \times 10^{-6} \,\mathrm{K}^{-1}$ at 100–400 °C to $(22 \pm 1) \times 10^{-6} \,\mathrm{K}^{-1}$ at 650– 850 °C. The complex oxide of Y_{0.8}Ca_{0.2}BaCo₄O_{7+δ} with a swedenborgite structure type shows very low TECs of (8–11) \times 10⁻⁶ K⁻¹; nevertheless, two specific temperatures (around 400 °C and 800 °C) can be distinguished when oxygen adsorption processes with subsequent oxygen desorption occur.391-393 This leads to an increase in TEC of up to 17×10^{-6} and 12×10^{-6} K⁻¹, respectively.

The Ni-based Ruddlesden-Popper phases are attractive from the viewpoint of chemical-induced strain effects (Fig. 13), showing an inappreciable response in dimensional changes towards oxygen deviation (δ) from the stoichiometric values or oxygen partial pressure variations. This unique pO2-tolerance strain originates from the structural features of nickelates. 376 Oxygen desorption of $Ln_2NiO_{4+\delta}$ results in two opposite effects, i.e. contraction of the c-parameters owing to the release of interstitial oxygen from the rock-salt layers and expansion of the perovskite-like layers (a-b plane) due to charge compensation by the partial reduction of existing cations (Ni3+). The combination of these effects results in negligible overall expansion.

Table 5 Average thermal expansion coefficient values (α_{av}) of Ln₂-NiO_{4+ δ}-based materials. In most cases, dilatometry measurements were performed in ambient air atmosphere between RT and 1000 °C. Visualisation of the data is presented in Fig. S3, ESI

$\alpha_{\rm av} imes 10^6, { m K}^{-1}$	Ref.
12.6	114
	53, 72 and 371
	340
	372
	218
	373
	371 and 374
	339
	359
	171
	374
	131
	359
	132
	131
	131
	359
	375
12.8	204 and 205
13.8	376
12.7	377
14.2	376
13.2	378
13.3	372
13.4	72
13.6	53
13.9	379
15.5	101
	373
	380
	268
	201
	268
	177
	381
	382
	132
	101
	118
13.5	379
40 =	
	53
	383
14.5	384
15.0	132
13.4	373
12.9	91
13.0	385
14.4	385
12.3	108
	108 132
12.3	
	12.6 13.0 13.1 13.5 13.8 13.2 12.6 11.2 13.9 15.3 13.8 14.5 13.9 14.2 14.6 14.9 15.2 13.0 12.8 13.8 12.7 14.2 13.2 13.3 13.4 13.6 13.9 15.5 13.9 14.1 12.0 15.0 12.9 13.1 13.8 12.7 14.5 13.7 16.2 13.7 16.2 13.5

Conversely, defect formation and interaction upon oxygen desorption occur within the sole structure of the conventional perovskite- and fluorite-related materials. For Co- and Fe-based oxides, oxygen vacancies are also formed (that should contract the cell), but cation reduction and cation–cation repulsion have a greater influence, leading to overall chemical expansion.

Interesting results were obtained by Flura et al.72 in their study on the structural parameters of basic nickelates $(La_2NiO_{4+\delta}, Nd_2NiO_{4+\delta}, and Pr_2NiO_{4+\delta})$ utilising hightemperature XRD analysis, revealing the effects of temperature, measurement atmosphere and phase transitions on the thermo-chemical response of the corresponding ceramics (Fig. 13c). Based on this analysis, decreasing pO₂ from the conventional value of air to $\sim 1 \times 10^{-4}$ atm (helium atmosphere) resulted in an apparent decrease in TEC across the entire temperature range. Simultaneously, the TEC behaviour of $Pr_2NiO_{4+\delta}$ was distinct from that of other nickelates, where its phase transition (Bmab \rightarrow I4/mmm) decreased the TECs, while the phase transitions for La₂NiO_{4+δ} and Nd₂NiO_{4+δ} (Fmmm \rightarrow I4/mmm occurred around 150 and 500 °C, respectively) were found to increase the TECs with respect to the lowtemperature range. The dilatometry data of Berger et al. 118 confirm that the TECs decreased with a reduction in pO2 for a Pr-based nickelate $(Pr_2Ni_{0.9}Co_{0.1}O_{4+\delta})$ within the lowtemperature range of 12.2 \times 10⁻⁶ K⁻¹ at pO₂ = 1 atm to $10.8 \times 10^{-6} \text{ K}^{-1}$ at $pO_2 = 1 \times 10^{-2}$ atm. However, the TECs increased for the high-temperature range of 15.9 \times 10⁻⁶ to $16.9 \times 10^{-6} \text{ K}^{-1}$ due to the oxygen desorption mechanism. The difference between these two works is the lowtemperature structure of praseodymium nickelate (Bmab72 and Fmmm¹¹⁸), implying that this can also affect the achieved TECs and their relation relative to each other in different temperature regions.

5.3. Proton transportation

The RP phases of the first order, $A_2BO_{4+\delta}$, are typically hyperstoichiometric compounds having an excess of interstitial oxygen in their structure. It is evident that the appearance of proton transport (if existing) in $A_2BO_{4+\delta}$ occurs with different mechanism(s) compared with the conventional acceptor-doped perovskites, where proton defects are formed under dissociative adsorption of water molecules by existing oxygen vacancies:

$$H_2O + V_0^{"} + O_0^{\times} \rightleftarrows 2OH_0^{"}$$
 (13)

In contrast, the interstitial sites in the classical $\text{Ln}_2\text{NiO}_{4+\delta}$ compounds are important for proton transportation since H_2O is more favourable for insertion at the interstitial sites, as shown in Fig. 14.

The protonation of oxides is generally achieved within a hydration process. The hydration equation for Ln_2BO_{4+x} can be written as:¹⁴⁴

$$H_2O + O_0^{\times} \rightleftharpoons OH_0^{\cdot} + OH_i^{\prime}$$
 (14)

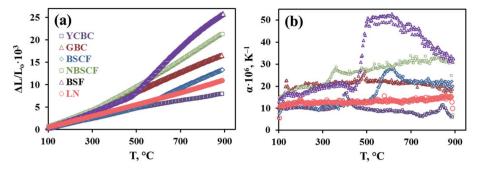


Fig. 12 Thermal expansion (a) and TEC values (b) of some electrode materials as a function of temperature (ambient air atmosphere). Data is presented based on experiments. Abbreviations: LN = $La_2NiO_{4+\delta}$, YCBC = $Y_{0.8}Ca_{0.2}BaCo_4O_{7+\delta}$, GBC = GdBaCo $_2O_{5+\delta}$, BSCF = $Ba_{0.5}Sr_{0.5}Co_{0.8}Fe_{0.2}O_{3-\delta}$, BSF = $Ba_{0.5}Sr_{0.5}FeO_{3-\delta}$, and NBSCF = NdBa $_{0.5}Sr_{0.5}Co_{1.5}Fe_{0.5}O_{5+\delta}$. Reproduced with permission.³⁴⁹ Copyright 2017, Elsevier.

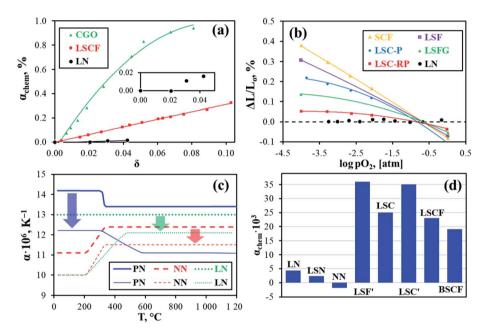


Fig. 13 Chemical strain effects for some oxide materials: (a) chemical expansion coefficient as a function of the oxygen non- (over-) stoichiometry and (b) dimensional change as a function of oxygen partial pressure at 800 °C. These graphs presented in ref. 376 and 394 were prepared based on the data in ref. 395–400. (c) Temperature dependence of TECs for nickelates under oxidising ($pO_2 = 0.2$ atm) and reducing (He, $pO_2 = 1.2$ x 10^{-4} atm) atmospheres. (d) Chemical expansion coefficients for some oxygen electrodes at 800 °C in air. (9.120.401) Abbreviations: LN = La₂-NiO_{4+ δ}, NN = Nd₂NiO_{4+ δ}, PN = Pr₂NiO_{4+ δ}, LSN = La_{1.8}Sr_{0.2}NiO_{4+ δ}, CGO = Ce_{0.9}Gd_{0.1}O_{1.95- δ}, LSCF = La_{0.6}Sr_{0.4}Co_{0.2}Fe_{0.8}O_{3- δ}, SCF = SrCo_{0.8}Fe_{0.2}O_{3- δ}, LSF = La_{0.3}Sr_{0.7}FeO_{3- δ}, LSF' = La_{0.6}Sr_{0.4}FeO_{3- δ}, LSC-P = La_{0.5}Sr_{0.5}CoO_{3- δ}, LSC = La_{0.6}Sr_{0.4}CoO_{3- δ}, LSC' = La_{0.3}Sr_{0.7}FeO_{3- δ}, and BSCF = Ba_{0.5}Sr_{0.5}Co_{0.2}Fe_{0.8}O_{3- δ}.

According to the equation, it seems that H_2O is split into OH and H, while OH occupies the interstitial site and H attaches to the lattice oxygen. It should be noted that this is only an assumption since little research activity has been carried out to determine the exact mechanism for hydration and proton transportation in $Ln_2NiO_{4+\delta}$ materials. However, this assumption is reasonable if one considers the loss of the interstitial oxygen in the classical $Ln_2NiO_{4+\delta}$ (such as La_2NiO_{4+x}) at high temperatures, ⁶⁷ leaving the interstitial sites to be compensated by H_2O . Although the loss of the interstitial oxygen was observed in La_2NiO_{4+x} , no observation of oxygen vacancy was reported in this study. Li *et al.*²³² arrived at a similar conclusion. They studied the hydration ability of Pr_2NiO_{4+x} and found that

the insertion of H_2O at the interstitial sites is the most likely explanation for the protonation of this material. Although there are three possible ways for the formation of proton defects in Pr_2NiO_{4+x} , the very low H_2 partial pressure at the cathode side eliminates the possibility of the direct combination of H_2 with interstitial oxygen. They also revealed that the insertion of H_2O at the oxygen vacancies through eqn (13) is not significant due to the very low oxygen vacancy content in the oxygen-rich environment for Pr_2NiO_{4+x} . Therefore, protonation can only occur with the insertion of H_2O at the interstitial sites. However, differing from the opinion of Grimaud *et al.*¹⁴⁴ that proton defects form at both interstitial sites and with lattice oxygen, Li

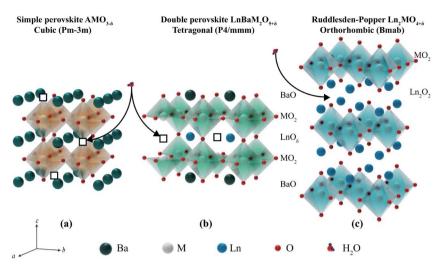


Fig. 14 Possible places for the insertion of water molecules in the different structures

et al. indicated that the H₂O split reaction is only associated with the interstitial sites, with the reaction equation written as:

$$2H_2O + 4O_i' \rightleftharpoons 4OH_i' + O_2 + 4e'$$
 (15)

The difference in the opinion of researchers concerning the hydration process of $\rm Ln_2NiO_{4+\delta}$ materials implies that this can be an interesting question for further study. Nevertheless, both studies agree that $\rm Ln_2NiO_{4+\delta}$ has a different hydration mechanism in comparison with that for traditional ABO_{3-x} perovskite oxides, while the interstitial sites play an important role in the protonation.

For the RP phases with pronounced basic properties (for example, BaLaInO₄), the proton defects may be formed according to the following reactions:^{144,405–408}

$$H_2O + V_O^{"} + O_i^{"} \rightleftarrows OH_O^{'} + OH_i^{'}$$
 (16)

$$H_2O + O_O^{\times} + V_i^{\times} \rightleftarrows OH_O^{\bullet} + OH_i^{\prime}$$
 (17)

$$H_2O + O_i^{"} \rightleftharpoons 2OH_i^{'}$$
 (18)

The first two reactions require vacancy (lattice or interstitial) defects. These defects occur to a considerable degree in nickelates, even those that are not doped by acceptors (eqn (19)). In detail, the heating of $\mathrm{Ln_2NiO_{4+\delta}}$ leads to oxygen desorption, and correspondingly the appearance of the necessary vacancies (eqn (20)), which are theoretically capable of providing hydration. The third reaction does not require crystallographic or interstitial vacancies, allowing the formation of proton defects directly upon water interaction with interstitial oxygen. Tarasova *et al.*⁴⁰⁷ used IR spectroscopy to reveal the two energetically non-equivalent hydroxyl groups in BaLaInO₄ and BaLaIn_{0.9}Nb_{0.1}O_{4+ δ}, confirming the possible realisation of eqn (14) and (15).

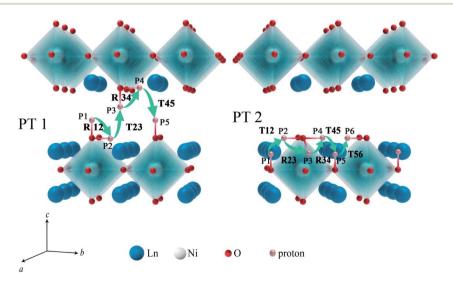


Fig. 15 Schematic representations of two paths for proton transport (PT) between the inter-layer PT1 and inner-layer PT2. Pn (n = 1-6) are the possible proton locations and R and T refer to proton rotation and transfer, respectively.³³²

$$MO \rightarrow M_{Ln}^{'} + V_O^{"} + O_O^{\times}$$
 (19)

$$2Ni_{Ni}^{\times} + 0.5O_{2} \rightleftharpoons 2Ni_{Ni}^{'} + V_{O}^{"} + \frac{1}{2}O_{2}$$
 (20)

Grimaud et al. 144 analysed the weight change of Pr₂NiO_{4+δ} in air with various pH₂O levels (from 0.002 to 0.095 atm). This analysis revealed different weight changes depending on pH2O, but could not determine the actual proton concentration in praseodymium nickelate due to the difficulty associated with two simultaneous processes, hydration/dehydration and oxygen sorption/desorption.

Zhang et al.332 studied both oxygen migration and proton diffusivity in undoped and doped La2NiO4+ô using firstprinciple calculations, including analyses of relative energy, effective charge and electron density. They proposed a possible mechanism (Fig. 15) and found that proton transportation is unfavourable in a perfect La₂NiO_{4+δ} crystal; however, fast proton transport may occur for Co- and Cu-doped derivatives. It should be noted that unlike oxygen transport, proton migration is proposed to occur between the apical and equatorial oxygen ions with a resulting energy barrier of \sim 1.1 eV, having five potential insertion sites on oxygen ions. This fact implies that the migration of protons between two rock-salt layers is unfavourable in nickelates in terms of a very large energy barrier, which is in accordance with the hydration of other RP phases. 409,410 However, it is also noted that the calculated proton transport route, whether rotating or jumping, is associated with the lattice oxygen only, and no interstitial oxygen is considered. As discussed above, the interstitial oxygen can play an important role in the protonation for $Ln_2NiO_{4+\delta}$, while the participation of interstitial oxygen may change the distance of proton migration, and thus alter the optimised migration routes. Therefore, it will be desirable to have a more comprehensive view concerning the first-principle calculations of proton transportation by also considering the interstitial oxygens. Although this aspect has yet to be considered in detail, it can be beneficial for an understanding of proton transportation in the $Ln_2NiO_{4+\delta}$ class of materials.

As can be seen above, the issue of the proton transportation in $Ln_2NiO_{4+\delta}$ phases remains open and should be evaluated in detail in future research.

5.4. Electrochemical activity towards oxygen- and protoninvolving electrode processes

The electrochemical activity of electrodes depends on different parameters affecting the active zone where the electrochemical reactions occur. Among them, a number of parameters have a technology-related nature (particle size, tortuosity, and porosity), which regulates the triple phase boundary (TPB) length.411-415 Another set is related to the inherent properties of the phases, including oxygen and electron transport, which determine the activity of TPB. Although the complete separation of composition from technological parameters is a difficult task, some regularities can be revealed.

Oxygen electrodes of SOFCs or PCFCs allow the realisation of the oxygen reduction reaction (ORR):

For SOFCs:

$$O_{2(gas)} + 4e_{(electrode)}^{-} \rightleftharpoons 2O_{(electrolyte)}^{2-}$$
 (21)

For PCFCs:

$$O_{2(gas)} + 4e_{(electrode)}^- + 4H_{(electrolyte)}^+ \rightleftharpoons H_2O_{(gas)}$$
 (22)

This reaction occurs in the case of predominantly electronic conductors, such as Pt and LSM.416 When materials can also conduct oxygen-ions, they become mixed ionic-electronic conductors (MIECs).417-422 MIECs have a distinct advantage over electronic conductors since the ORR can occur at the surface of the ceramic electrode together with TPB:

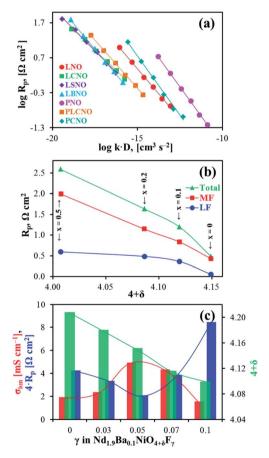


Fig. 16 Correlation of polarisation resistances of $Ln_2NiO_{4+\delta}$ -based electrodes (R_p) measured on the symmetrical cells with proton-conducting electrolytes with kinetic parameters (kD), oxygen content (4 + δ) and ionic conductivity (σ_{ion}) of nickelates. Panels (a), (b) and (c) were reproduced from works by Sadykov et al., 267 Grimaud et al. 144 and Tarutin et al., 207 respectively. In panel (a) the following abbreviations are used: LNO = $La_2NiO_{4+\delta}$, LCNO = $La_{1.7}Ca_{0.3}NiO_{4+\delta}$, LSNO = $\mathsf{La}_{1.7}\mathsf{Sr}_{0.3}\mathsf{NiO}_{4+\delta},\,\mathsf{LBNO}=\mathsf{La}_{1.7}\mathsf{Ba}_{0.3}\mathsf{NiO}_{4+\delta},\,\mathsf{PNO}=\mathsf{Pr}_2\mathsf{NiO}_{4+\delta},\,\mathsf{PCNO}=\mathsf{Pr}_2\mathsf{NiO}_{4+\delta}$ $Pr_{1.7}Ca_{0.3}NiO_{4+\delta}$, and $PLCNO = Pr_{0.85}La_{0.85}Ba_{0.3}NiO_{4+\delta}$. Panel (b) highlights the results obtained for the PSN|BCY|PSN symmetrical cells at 600 °C and $pH_2O = 0.03$ atm. Here, PSN = $Pr_{2-x}Sr_xNiO_{4+\delta}$, and BCY = BaCe_{0.9} $Y_{0.1}O_{3-\delta}$. In panel (c), all the data are provided for 600 °C and the R_p values were obtained for symmetrical cells based on protonconducting BaCe_{0.5}Zr_{0.3}Y_{0.1}Y_{0.1}O_{3- δ} electrolyte.

$$For \ SOFCs: \begin{cases} O_{2(gas)} \rightleftarrows O_{2(ads)} \rightleftarrows 2O_{(ads)} \\ O_{(ads)} + 2e_{(electrode)}^{-} \rightleftarrows O_{(ads)}^{2-} \\ O_{(ads)}^{2-} \rightleftarrows O_{(electrode)}^{2-} \rightleftarrows O_{(electrolyte)}^{2-} \end{cases} \qquad (23) \qquad For \ PCFCs: \begin{cases} O_{2(gas)} \rightleftarrows O_{2(ads)} \rightleftarrows 2O_{(ads)} \\ O_{(ads)} + 2e_{(electrode)}^{-} \rightleftarrows O_{(ads)}^{2-} \\ O_{(ads)}^{2-} \rightleftarrows O_{(electrode)}^{2-} \rightleftarrows O_{(ads)}^{2-} \\ O_{(ads)}^{2-} \rightleftarrows O_{(electrode)}^{2-} \rightleftarrows O_{(electrode)}^{2-} \end{cases} \qquad (24)$$

Table 6 Overall polarisation resistances (R_p) and corresponding activation energy (E_a) of nickelate-based electrodes for symmetrical cells, electrode|electrolyte|electrode. Visualisation of the data is presented in Fig. S4, ESI

		$R_{\rm p}$, Ω cm ²	2			
Electrode	Electrolyte	500 °C	600 °C	700 °C	$E_{\rm a}$, eV	Ref.
$La_2NiO_{4+\delta}$	$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	91.8	14.1	_	1.01	429
$La_2NiO_{4+\delta}$	$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	41.6	3.88	_	1.01	81
$\text{La}_2 \text{NiO}_{4+\delta}$	$\text{La}_{5.5}\text{WO}_{11.25-\delta}$	_	35	11.5	1.36	265
$La_2NiO_{4+\delta}$	$BaCe_{0.8}Y_{0.2}O_{3-\delta}$	_	20.8	4.94	1.01	349
$\text{La}_2 \text{NiO}_{4+\delta}$	$BaCe_{0.5}Zr_{0.3}Dy_{0.2}O_{3-\delta}$	4.5	0.64	0.10	1.24	157
$La_{2}Ni_{0.9}Cu_{0.1}O_{4+\delta}$	$BaCe_{0.5}Zr_{0.3}Dy_{0.2}O_{3-\delta}$	4.7	0.92	0.21	0.96	157
$\text{La}_2\text{Ni}_{0.8}\text{Cu}_{0.2}\text{O}_{4+\delta}$	$BaCe_{0.5}Zr_{0.3}Dy_{0.2}O_{3-\delta}$	13.1	2.06	0.47	1.01	157
$\text{La}_{2}\text{Ni}_{0.7}\text{Cu}_{0.3}\text{O}_{4+\delta}$	$BaCe_{0.5}Zr_{0.3}Dy_{0.2}O_{3-\delta}$	42.3	7.90	1.20	1.08	157
$La_{1.5}Pr_{0.5}NiO_{4+\delta}$	$\text{La}_{5.5} \text{WO}_{11.25-\delta}$	_	18.4	2.52	1.39	265
$70\%\text{La}_2\text{NiO}_{4+\delta}$ $-30\%\text{LaNi}_{0.6}\text{Fe}_{0.4}\text{O}_{3-\delta}$	$BaZr_{0.1}Ce_{0.7}Y_{0.2}O_{3-\delta}$	_	13.4	2.09	1.24	430
$50\%\text{La}_2\text{NiO}_{4+\delta}$ -50%LaNi _{0.6} Fe _{0.4} O _{3-\delta}	$BaZr_{0.1}Ce_{0.7}Y_{0.2}O_{3-\delta}$		15.5	2.34	1.26	430
$30\%\text{La}_2\text{NiO}_{4+\delta}$ -70%LaNi _{0.6} Fe _{0.4} O _{3-\delta}	$BaZr_{0.1}Ce_{0.7}Y_{0.2}O_{3-\delta}$		19.9	2.92	1.27	430
$50\%\text{La}_2\text{NiO}_{4+\delta}$ $-50\%\text{LaNi}_{0.6}\text{Fe}_{0.4}\text{O}_{3-\delta}$	$BaZr_{0.1}Ce_{0.7}Y_{0.2}O_{3-\delta}$		5.42	0.99	1.08	431
$\text{La}_{1.95} \text{Ba}_{0.05} \text{NiO}_{4+\delta}$	$BaCe_{0.5}Zr_{0.3}Dy_{0.2}O_{3-\delta}$	37.9	5.56	1.11	1.08	432
$La_{1.7}Ca_{0.3}NiO_{4+\delta}$	$BaCe_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}$	_	10.8	1.88	1.24	433
$\text{La}_{1.7}\text{Sr}_{0.3} ext{NiO}_{4+\delta}$	$BaCe_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}$		5.99	0.74	1.47	433
$La_{1.7}Ba_{0.3}NiO_{4+\delta}$	$BaCe_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}$		4.38	0.64	1.44	433
$La_{1.7}Ba_{0.3}NiO_{4+\delta}$ -Ba $Ce_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}$	$BaCe_{0.7}Zr_{0.1}Y_{0.2}O_{3-\delta}$		1.52	0.21	1.35	433
$La_{1.7}Ba_{0.3}NiO_{4+\delta}$ -BaCe _{0.89} Gd _{0.1} Cu _{0.01} O _{3-\delta}	$BaCe_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}$		4.64	0.62	1.35	359
$La_{1.7}Ca_{0.3}NiO_{4+\delta}$ -Ba $Ce_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}$	$BaCe_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}$		6.14	1.14	1.21	359
$La_{1.7}Sr_{0.3}NiO_{4+\delta}$ -Ba $Ce_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}$	$BaCe_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}$		9.59	0.82	1.34	359
$La_{1.7}Ba_{0.3}NiO_{4+\delta}$ -BaCe _{0.89} Gd _{0.1} Cu _{0.01} O _{3-\delta}	$\text{CaZr}_{0.95}\text{Sc}_{0.05}\text{O}_{3-\delta}$		278	45	1.14	434
$La_{1.7}Ca_{0.3}NiO_{4+\delta}$ -Ba $Ce_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}$	$\text{CaZr}_{0.95}\text{Sc}_{0.05}\text{O}_{3-\delta}$	_	666	128	1.20	434
$La_{1.7}Sr_{0.3}NiO_{4+\delta}$ -Ba $Ce_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}$	$\text{CaZr}_{0.95}\text{Sc}_{0.05}\text{O}_{3-\delta}$		213	30.0	1.24	434
$La_{1.2}Sr_{0.8}NiO_{4+\delta}$	$BaCe_{0.7}Zr_{0.1}Y_{0.2}O_{3-\delta}$	1.58	0.29	0.07	0.93	435
$BaCe_{0.68}Zr_{0.1}Y_{0.1}Yb_{0.1}Cu_{0.02}O_{3-\delta} \ infiltrated$	$BaCe_{0.68}Zr_{0.1}Y_{0.1}Yb_{0.1}Cu_{0.02}O_{3-\delta}$	1.34	0.20	0.04	1.15	436
with 42.2 wt% $La_{1.2}Sr_{0.8}NiO_{4-\delta}$						
LaSrNiO _{4+δ}	$BaCe_{0.9}Y_{0.1}O_{3-\delta}$		2.04	0.32	1.73	81
$Nd_2NiO_{4+\delta}$	$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	278	29.9	4.9	1.06	429
$Nd_2NiO_{4+\delta}$	$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	81.1	6.75	3.5	1.06	81
$Nd_{1.95}Ba_{0.05}NiO_{4+\delta}$	$BaCe_{0.5}Zr_{0.3}Dy_{0.2}O_{3-\delta}$	24.6	3.42	0.45	1.25	432
$60\% Nd_{1.95}NiO_{4+\delta}$ -40% $BaZr_{0.1}Ce_{0.7}Y_{0.1}Yb_{0.1}O_{3-\delta}$	$BaZr_{0.1}Ce_{0.7}Y_{0.1}Yb_{0.1}O_{3-\delta}$	_	15.4	1.87	1.33	358
$50\% Nd_{1.95}NiO_{4+\delta} - 50\% \ BaZr_{0.1}Ce_{0.7}Y_{0.1}Yb_{0.1}O_{3-\delta}$	$BaZr_{0.1}Ce_{0.7}Y_{0.1}Yb_{0.1}O_{3-\delta}$	_	9.55	1.47	1.24	358
$50\% {\rm Nd_{1.95}NiO_{4+\delta}} - 50\% \ {\rm BaZr_{0.1}Ce_{0.7}Y_{0.1}Yb_{0.1}O_{3-\delta}}$	$BaZr_{0.1}Ce_{0.7}Y_{0.1}Yb_{0.1}O_{3-\delta}$	_	7.06	1.21	1.23	358
$Nd_{1.9}Ba_{0.1}NiO_{4+\delta}$	$BaZr_{0.3}Ce_{0.5}Y_{0.1}Yb_{0.1}O_{3-\delta}$	136	17.9	2.67	1.30	207
$\mathrm{Nd}_{1.9}\mathrm{Ba}_{0.1}\mathrm{NiO}_{4+\delta}$	$BaCe_{0.5}Zr_{0.3}Dy_{0.2}O_{3-\delta}$	115	15.9	1.8	1.40	437
$Nd_{1.9}Ba_{0.2}NiO_{4+\delta}$	$\mathrm{BaCe}_{0.5}\mathrm{Zr}_{0.3}\mathrm{Dy}_{0.2}\mathrm{O}_{3-\delta}$	212	21.4	3.5	1.30	437
$Nd_{1.9}Ba_{0.1}NiO_{4+\delta}F_{0.05}$	$BaZr_{0.3}Ce_{0.5}Y_{0.1}Yb_{0.1}O_{3-\delta}$	63.1	7.82	0.41	1.40	207
$Nd_{1.9}Ba_{0.1}NiO_{4+\delta}F_{0.1}$	$BaZr_{0.3}Ce_{0.5}Y_{0.1}Yb_{0.1}O_{3-\delta}$	259	33.1	5.09	1.34	207
$\mathrm{NdSrNiO}_{4+\delta}$	$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	93.2	1.62	0.27	1.54	81
$Pr_2NiO_{4+\hat{\delta}}$	$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	6.71	1.7	0.38	0.87	429
$\text{Pr}_2 \text{NiO}_{4+\delta}$	$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	3.7	0.54	0.18	1.02	144
$\text{Pr}_{2}\text{NiO}_{4+\delta}$	$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	6.75	0.49	0.14	0.86	81
$\text{Pr}_2 \text{NiO}_{4+\delta}$	$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	1.83	0.32	0.13	0.97	231
$Pr_2NiO_{4+\delta}$	$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	9.17	0.89	0.19	0.96	252
$(Pr_{0.9}La_{0.1})_2(Ni_{0.74}Cu_{0.21}Nb_{0.05})O_{4+\delta}$	$BaCe_{0.7}Zr_{0.1}Y_{0.2}O_{3-\delta}$	5.07	0.77	0.17	1.03	360
$Pr_{1.9}Ca_{0.1}NiO_{4+\delta}$	$BaCe_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}$	_	3.67	0.51	1.47	433
$Pr_{1.95}Ba_{0.05}NiO_{4+\delta}$	$BaCe_{0.5}Zr_{0.3}Dy_{0.2}O_{3-\delta}$	10.0	0.62	0.19	1.19	434
$Pr_{1.9}Sr_{0.1}NiO_{4+\delta}$	$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	9.92	1.16	0.21	1.22	144
$Pr_{1.8}Sr_{0.2}NiO_{4+\delta}$	$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	14.5	1.92	0.39	1.18	144
$Pr_{1.5}Sr_{0.5}NiO_{4+\delta}$	$\mathrm{BaCe}_{0.9}\mathrm{Y}_{0.1}\mathrm{O}_{3-\delta}$	30.1	2.7	0.45	1.26	144
$\text{Pr}_{1.2} \text{Sr}_{0.8} \text{NiO}_{4+\delta}$	$BaCe_{0.7}Zr_{0.1}Y_{0.2}O_{3-\delta}$	3.47	0.57	0.13	0.98	435

Therefore, the TPB length is not limited by the contact of three phases, and this length can be extended due to additional pathways for oxygen molecules.

Despite having an ionic electrolyte type, the ORR comprises a combination of two consequent macroscopic processes, namely oxygen exchange at the surface of the electrodes and oxygen diffusion within the ceramic phase. On the one hand, these processes are characterised by kinetic parameters (oxygen exchange constant, k, and oxygen diffusion coefficient, D); on the other hand, they determine the polarisation resistance of the electrodes (R_p) , as shown in particular by the ALS (Adler, Lane, Steele) model.423 Therefore, it is evident that there is a direct relation between $R_{\rm p}$, k, and D. As can be seen from Fig. 16a, this correlation indeed exists, indicating that the lowest R_p values can be achieved for materials having the highest product of kD.

Grimaud et al. evaluated the electrochemical activity of Srdoped $Pr_2NiO_{4+\delta}$ electrodes in the system with a $BaCe_{0.9}Y_{0.1}O_{3-\delta}$ protonic conductor, and found that the optimal electrode performance was observed for the undoped Pr₂NiO_{4+δ} having a higher 4 + δ level. In detail, both partial (middle- and lowfrequency) resistances, related with the charge transfer through the interface and the electrode reactions, respectively, decreased

with 4 + δ growth (Fig. 16b). Most probably, the " R_p -(4 + δ)" relation is a special case that overlays with other tendencies. For example, no direct " R_p -(4 + δ)" correlation was revealed in the work of Tarutin et al. 207 Instead, the electrochemical activity of the Nd₂NiO_{4+ô}-based electrodes was regulated by the oxygen-ionic conductivity (Fig. 16c) or oxygen diffusion coefficient, according to the Nernst-Einstein relationship.

Table 6 summarises the Ln₂NiO_{4+δ}-based electrodes obtained for symmetrical cells with proton-conducting electrolytes. Although the polarisation resistances vary across a wide range (\sim 3 orders of magnitude), the target values, after excluding the extreme examples, can be estimated to be \sim 10, 3 and 0.5 Ω cm² at 500, 600 and 700 °C, respectively, with a corresponding activation energy of \sim 1 eV. Analysis of the activation energies shows that they drop in the range of ~ 0.9 –1.5 eV, implying that kinetic parameters (see Table 3) regulate the overall electrochemical activity of the electrodes.

The high noise of the listed R_p results is due to the technological features of the various materials. As can be seen, the compositions of the electrodes include basic nickelates and their doped or co-doped analogues and composites, which can be prepared by conventional means or utilising attractive techniques for the improvement of electroactivity (for example,

Table 7 Literature survey on the analysis of impedance spectra performed for the symmetrical ED|ET|ED cells, where ED is the electrode and ET is the electrolyte

ET composition	ED composition	Conditions	Equivalent circuit	Description	Ref.
$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	$La_2NiO_{4+\hat{\delta}}$	Wet air, $pH_2O = 0.03$ atm	$(RQ)_{\mathrm{MF}}$ - $(RQ)_{\mathrm{LI}}$	F MF: ionic transfer reaction at the ED/ET interface LF: oxygen dissociation and oxygen reduction reaction	e 81
$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	$\mathrm{Pr}_{2}\mathrm{NiO}_{4+\delta}$	Wet air, $pH_2O = 0.03$ atm	$(RQ)_{\mathrm{MF}}$ - $(RQ)_{\mathrm{LI}}$	F MF: H ⁺ charge transfer at the ED/ET interface LF: oxygen adsorption, oxygen dissociation and molecular diffusion	439
$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	$\mathrm{Pr}_{2}\mathrm{NiO}_{4+\delta}$	Wet air, $pH_2O = 0.03-0.3$	$(RQ)_{\mathrm{MF}}$ - $(RQ)_{\mathrm{LI}}$	F MF: H ⁺ charge transfer at the ED/ET interface LF: water formation or water gas diffusion	144
$BaCe_{0.9}Y_{0.1}O_{3-\delta}$	$\text{Pr}_{2}\text{NiO}_{4+\delta}$		$(RQ)_{\mathrm{MF}}$ – $(RQ)_{\mathrm{LI}}$	F MF: H ⁺ charge transfer at the ED/ET interface LF: oxygen dissociative adsorption or oxygen reduction reaction	231
$BaZr_{0.7}Ce_{0.2}Y_{0.1}O_{3-\delta}$	${ m La_2NiO_{4+\delta}}$	Wet air, $pH_2O = 0.03$	$(RQ)_3$	 charge transfer at the ED/ET interface oxygen adsorption molecular oxygen diffusion 	440
$\begin{array}{l} BaCe_{0.7}Zr_{0.1}Y_{0.2}O_{3-\delta}\\ \left(BCZY\right) \end{array}$	$\begin{array}{l} La_2NiO_{4+\delta} \ and \ La_2NiO_{4+\delta} \\ -LaNi_{0.6}Fe_{0.4}O_{3-\delta} \ composites \end{array}$	_	$(RQ)_{\mathrm{HF}}$ - $(RQ)_{\mathrm{LF}}$	HF: O ²⁻ charge transfer at the ED/ET interface LF: oxygen adsorption, dissociation, charge transfer and molecular diffusion	430
BCZY	BCZY–Ndr $_{1.95}$ NiO $_{4+\delta}$	Wet air, $pH_2O = 0.03$		HF: H ⁺ charge transfer at the ED/ET interface F MF: oxygen reduction reaction LF: diffusion process of O ⁻ to TPB	358
BCZY	${\rm La_2NiO_{4+\delta}\text{-}LaNi_{0.6}Fe_{0.4}O_{3-\delta}}$	_	$(RQ)_{\mathrm{HF}}$ - $(RQ)_{\mathrm{LH}}$	HF: O ²⁻ charge transfer at the ED/ET interface LF: oxygen adsorption, dissociation, charge transfer and molecular diffusion	431
BCZY	$\big(Pr_{0.9}La_{0.1}\big)_2Ni_{0.74}Cu_{0.21}Nb_{0.05}O_{4+\delta}$	_	$(RQ)_{\mathrm{HF}}$ – $(RQ)_{\mathrm{LF}}$	F HF: H ⁺ charge transfer LF: oxygen surface exchange	360
$BaCe_{0.5}Zr_{0.3}Dy_{0.2}O_{3-4}$	$_{\delta}~\mathrm{Nd}_{1.95}\mathrm{Ba}_{0.05}\mathrm{NiO}_{4+\delta}$	Dry air	$(RQ)_{\mathrm{HF}}^{-}$ $(RQ)_{\mathrm{MF}}^{-}$ $(RQ)_{\mathrm{LI}}^{-}$	HF: charge transfer across the electrode/electrolyte	e 361
BCZY	$Ln_{1.2}Sr_{0.8}NiO_{4+\delta}$ (Ln = La and Pr)	Ambient air	$(RQ)_{\mathrm{HF}}$ – $(RQ)_{\mathrm{LF}}$	F: bulk charge transfer process LF: —	435
$\begin{array}{l} BaCe_{0.6}Zr_{0.2}Y_{0.2}O_{3-\delta} \\ \left(BCZY'\right) \end{array}$	$Pr_2NiO_{4+\delta}$ -BCZY'	Wet air, $pH_2O = 0.6$	$(RQ)_{\mathrm{HF}}$ - $(RQ)_{\mathrm{LF}}$	 -	232

infiltration). Moreover, depending on the compositions of the cells and external conditions, the $\rm Ln_2NiO_{4+\delta}$ -based phases are reported to show proton transportation, thus constituting the class of triple-conducting materials. $^{424-428,438}$ Thus, considering the above-mentioned factors, further elaboration will specify the important details affecting the electrochemical performance of the considered electrode materials.

Analysing the nature of ongoing electrode processes, it is rational to consider an algorithm for decoding the impedance spectra collected for symmetrical cells. According to Table 7, the corresponding spectra are described by two or three RQ-combinations (R is the partial resistance of R_p , while Q is the constant phase element). If the first of these combinations is mostly due to an ionic charge transfer between the electrolyte and electrode phases, the second (or third) combination(s) may be attributed to various processes, starting from the reduction of oxygen atoms and ending with molecular diffusion of gas components within a porous electrode. Again, this uncertainty in physical models occurs due to completely different electrochemical systems that have varying functional parameters.

Among the numerous publications, there are a few works in which the authors tried to distinguish the factors associated with the inherent properties of nickelate phases and their microstructural parameters from each other.

For example, Grimaud *et al.*¹⁴⁴ studied the simultaneous effects of pO_2 and pH_2O on the R_p of the $Pr_2NiO_{4+\delta}$ electrode at 600 °C. According to their results, R_p comprises the sum of two

(medium- (MF) and low- (LF) frequency) processes. $R_{\rm MF}$ is not affected by $p{\rm H}_2{\rm O}$, but depends on the variation of $p{\rm O}_2$ according to the following relation: $R_{\rm MF} \sim (p{\rm O}_2)^{-1/4}$. On the contrary, the power function, n, for the $R_{\rm LF} \sim (p{\rm O}_2)^{-n}$ relation changes from 0.44 for dry air ($p{\rm H}_2{\rm O} = 0.006$ atm) to 0.61 for wet air ($p{\rm H}_2{\rm O} = 0.20$ atm) due to the involved $p{\rm H}_2{\rm O}$ -associated processes.

In the next work of Grimaud $et\ al.$, ⁴⁴¹ special attention was paid to the electrode performance of $\Pr_2 \text{NiO}_{4+\delta}$ depending on the variation in $p\text{H}_2\text{O}$ and microstructural parameters. For a porous state of $\Pr_2 \text{NiO}_{4+\delta}$, the following tendencies were observed: $R_{\text{MF}} \sim (p\text{H}_2\text{O})^{-1/2}$ and $R_{\text{LF}} \sim (p\text{H}_2\text{O})^{-1}$. Their study indicated that the H_2O -assisted mechanism occurred for the considered system, namely proton charge transfer in the former case and water formation in the latter case. When the $\Pr_2 \text{NiO}_{4+\delta}$ electrode layer was prepared in a dense form, the first relation was again observed, confirming the ability of the nickelate to demonstrate proton transportation. However, no dependence (n=0) was observed between R_{LF} and $p\text{H}_2\text{O}$ that could be explained in terms of the limited diffusion of water molecules inside the electrodes.

Quarez *et al.*³⁵⁷ considered the effects of several parameters on the electrochemical activity of the same electrode, *i.e.* calcination temperature to fabricate the PN|LW|PN symmetrical cells, the thickness of the PN electrodes, and weight amount of LW phases in the PN–LW electrode composites (here PN = $Pr_2NiO_{4+\delta}$ and LW = $La_{27.15}W_{4.85}O_{55+\delta}$). Temperature- and thickness-related effects were predictable, where an

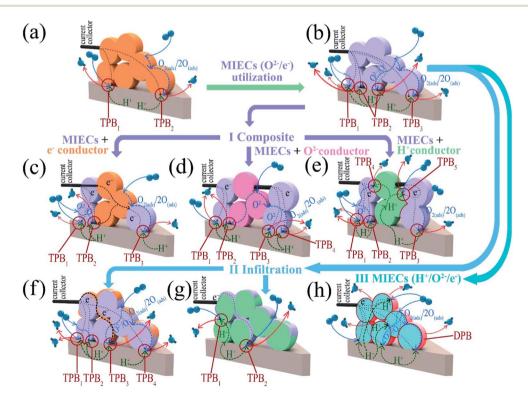


Fig. 17 Possible ways of electrochemical reactions in systems with proton-conducting electrolytes depending on the nature of the electrodes: (a) single-phase electronic conductor (EC), (b) single-phase MIEC, (c) MIEC + EC composite, (d) MIEC + O^{2-} -electrolyte, (e) MIEC + O^{2-} -electrolyte, (f) MIEC infiltrated with EC, (g) O^{2-} -electrolyte infiltrated with MIEC and (h) ideal single-phase triple-conducting O^{2-} -electrolyte infiltrated with MIEC and (h) ideal single-phase triple-conducting O^{2-} -electrolyte infiltrated with MIEC and (h) ideal single-phase triple-conducting O^{2-} -electrolyte infiltrated with MIEC and (h) ideal single-phase triple-conducting O^{2-} -electrolyte infiltrated with MIEC and (h) ideal single-phase triple-conducting O^{2-} -electrolyte, (e) MIEC + O^{2-} -electrolyte, (f) MIEC + O^{2-} -electrolyte, (g) MIEC + O^{2-} -electrolyte, (e) MIEC + O^{2-} -electrolyte, (f) MIEC + O^{2-} -electrolyte, (g) MIEC + O^{2-} -electrolyte, (h) MIEC + O^{2-} -electroly

improvement in electrochemical activity occurred for the highly porous electrodes (prepared at the lowest sintering temperature of 800 °C) and with the lowest thickness (12 μm against 24 and 36 µm for the other cases), implying that the determining role is played by molecular diffusion. However, no positive composition effect was observed upon the addition of the LW phase; therefore, composite materials do not always provide an incremental improvement in electrode functionality.

Technological effects were also evaluated by Solís et al.265 Here it was found that among three sintering temperatures (1050 °C, 1100 °C and 1150 °C), the best electrochemical activity for La_{1.5}Pr_{0.5}Ni_{0.8}Co_{0.2}O_{4+δ} electrodes was achieved for the lowest temperature, confirming the crucial effect of electrode porosity.

From the provided examples, it is clear that the electrochemical activity of nickelates can be quite easily tuned by technological factors, especially sintering temperature, ensuring a certain electrode porosity. However, the performance of these electrodes can be also regulated by means of engineering techniques aimed at extending the surface of the electrochemically active zone.

Fig. 17 shows some of the main strategies associated with increasing the number of active parts in electrochemical reactions. These strategies aim at introducing an additional charge carrier together with the electron carrier, i.e. either oxygen-ions, protons or both.

A gradual complication of the provided scheme consists of substituting a pure electron conductor (Fig. 17a) with a mixed ionic-electronic conductor (MIEC) with O²⁻-transportation (Fig. 17b). Many complex oxides belong to these MIECs, including cobaltites, 442-444 ferrites, 434-448 nickelates 52,234,379 having simple and layered structures, and solid solutions based on these classes. 449-452 Utilisation of MIECs deposited on protonconducting electrolytes allows the TPB to be extended by an MIEC phase perimeter limited by the electrolyte surface.

Another strategy for increasing the TPB involves the development of composite systems. Considering (O²⁻/h*)-MIEC as an example, three different types of composites can be formed (Fig. 17c-e), consisting of an MIEC combined with an electronic, oxygen-ionic or protonic conductor. The last combination seems to represent the most promising variant since it allows transportation of three different types of charge carriers, some of which exist within one phase, while others exist within the second phase.

An increase in the TPB can be realised by means of nanoscaled sediments formed purposefully on the surface of the main porous phases (Fig. 17f and g). This method, which is known as infiltration, is widely used to achieve high performance electrodes. 25,453-455 However, it should be noted that the number of active parts determining the TPB length may be lower than that for single-phase MIECs (see the comparison of Fig. 17b and g). Nevertheless, in most cases, this accelerates the

Table 8 Methods for improving the electrochemical performances of Ln₂NiO₄-based electrodes, where R_p is the polarisation resistance determined for the symmetrical ED|ET|ED cells at 600 °C

		Performance		
Case (see Fig. 17)	Example, $ED^a ET^b$	$R_{\rm p}$, $\Omega~{\rm cm}^2$	Ref.	Remarks
Single-phase materials				
(b)	LN BCY20	20.8	349	
	NBN BCZD	13.9	361	
Composite materials				
(c)	LN-LNF BCZY	13.4	430	LN : LNF = 7 : 3 (wt ratio)
	LN-LNF BCZY	19.9	430	LN : LNF = 3 : 7 (wt ratio)
	LN-LNF BCZY	5.4	431	LN : LNF = 1 : 1 (wt ratio)
(d)	LNF-SDC BCZY	7.2	431	LNF : SDC = 1 : 1 (wt ratio)
	LBN-SDC BCGC	7.8	433	LBN : SDC = 1 : 1 (wt ratio) + LNF current collecto
(e)	NN-BCZYYb BCZYYb	7.24	358	NN : BCZYYb = 1 : 1 (wt ratio)
	LBN-BCGC BCGC	4.6	359	LBN : BCGC = 1 : 1 (wt ratio)
	LBN-BCGC BCGC	4.3	433	LBN : BCGC = 1 : 1 (wt ratio)
	·			+ LNF current collector
Infiltration				
(f)	LNF _{inf} -LN _{bb} BCZY	0.31	459	31 wt% of infiltration
	III 1881 -			component
(g)	LN _{inf} -BZCY _{bb} BZCY	29.3	440	1
(8)	PLNCN _{inf} -BCZY _{bb} BCZY	0.76	360	46.1 wt% of infiltration component
Single-phase triple-cond	ucting material			
(h)	PN BCY	0.89	144	$pH_2O = 0.03$ atm (0.31 Ω cm ² at $pH_2O = 0.30$ atm)

 $^{^{}a} \ \ Electrode \ (ED) \ abbreviations: PSN = Pr_{1.4}Sr_{0.6}NiO_{4+\delta}, LN = La_{2}NiO_{4+\delta}, NBN = Nd_{1.95}Ba_{0.05}NiO_{4+\delta}, LNF = LaNi_{0.6}Fe_{0.4}O_{3-\delta}, NN = Nd_{1.95}NiO_{4+\delta}, LBN = La_{1.7}Ba_{0.3}NiO_{4+\delta}, PLNCN = (Pr_{0.9}La_{0.1})_{2}Ni_{0.74}Cu_{0.21}Nb_{0.05}O_{4+\delta}, PN = Pr_{2}NiO_{4+\delta}, \\ b \ Electrolyte \ (ET) \ abbreviations: BCY = BaCe_{0.9}Y_{0.1}O_{3-\delta}, BCY20 = BaCe_{0.8}Y_{0.2}O_{3-\delta}, BCZP = BaCe_{0.7}Zr_{0.1}Y_{0.2}O_{3-\delta}, SDC = Ce_{0.8}Sm_{0.2}O_{2-\delta}, BCGC = BaCe_{0.89}Gd_{0.1}Cu_{0.01}O_{3-\delta}, BCZYYb = BaCe_{0.8}Y_{0.1}O_{3-\delta}, BCZY = BaCe_{0.8}Y_{0$ $BaCe_{0.7}Zr_{0.1}Y_{0.2}Yb_{0.1}O_{3-\delta},\ BZCY=BaZr_{0.7}Ce_{0.2}Y_{0.1}O_{3-\delta},\ inf\ is\ the\ infiltration,\ and\ bb\ is\ the\ backbone.$

electrochemical reactions, thus promoting the higher performance of electrodes and cells based thereupon.

The last example represents (Fig. 17h) a case when the electrochemically active zone for triple-conducting materials is dramatically increased, where a double phase boundary (DPB) replaces the TPB, which is limited by the contact of three phases. Specifically, the entire surface of the single-phase triple-conducting electrode becomes active towards the oxygen reduction and water evolution reactions. In this case, there are three types of charge carriers (as in the case of the composites presented in Fig. 17e), which simultaneously exist within the same phase.

The ability of these phases to show triple-conducting behaviour is presently under discussion. For example, there are reports confirming the hydration of the nickelate phases: 144,441 however, no direct confirmation of the data has been

presented. Moreover, the term "triple-conducting" should mean that all three partial (electronic, oxygen-ionic and protonic) conductivities are comparable. This has not been achieved for any state-of-the-art electrode system at 600–800 °C, where only temperatures as low as 200–400 °C allow the promotion of sufficient mobility/concentration of protons. 456–458

To move from theoretical to experimental aspects, the research data corresponding to the above-mentioned examples of electrode systems is presented in Table 8. According to these observations, strategies related to infiltration and the design of triple-conducting materials are identified as the most reliable, resulting in the achievement of a polarisation resistance of less than 1 Ω cm² at 600 °C. It is important to note that Section 5.4 reports data obtained for symmetrical cells, when the opposite electrode sides operate under the same conditions. In the real

Table 9 Design and performance of PCFCs and PCECs with nickelate-based oxygen (steam) electrodes under OCV conditions at 600 °C, where $U_{\rm OC}$ is the open-circuit voltage value, $R_{\rm o}$, $R_{\rm p}$, and $R_{\rm t}$ are the ohmic, polarisation and total resistance, respectively, and σ is the conductivity of the thin-film electrolyte. Visualisation of the data is presented in Fig. S5, ESI

			Electrochemical characteristics						
Fuel electrode	Electrolyte ^a (thickness)	Oxygen electrode ^b	U_{OC} , V	$R_{\rm O}$, $\Omega~{\rm cm}^2$	$R_{\rm p}$, $\Omega~{\rm cm}^2$	$R_{\rm t}$, $\Omega~{\rm cm}^2$	σ , mS cm ⁻¹	Year	Ref.
Ni-BCY	BCY (40 μm)	PN	1.15	1.84	2.76	4.60	2.2	2010	252
Ni-BCY	BCY (27 μm)	NN	1.06	_	_	4.8	_	2013	460
Ni-BCY	BCY (50 μm)	PN	1.09	0.81	0.8	1.61	6.2	2014	461
Ni-BCZY	BCZY (20 μm)	LN	1.04	0.64	0.96	1.60	3.1	2014	430
Ni-BCZY	BCZY (20 μm)	70%LN-30% LNF	1.04	0.44	0.49	0.93	4.5	2014	430
Ni-BCZY	BCZY (20 μm)	50% LN-50% LNF	1.05	0.50	0.51	1.01	4.0	2014	430
Ni-BCZY	BCZY (20 μm)	30% LN-70% LNF	1.04	0.53	0.57	1.10	3.8	2014	430
Ni-BCZY	BCZY (24 μm)	50% LN-50% LNF	1.04	0.54	0.59	1.13	4.4	2015	431
Ni-BZY	BCZY44 (5 μm)	PN	1.03	0.77	1.00	1.77	0.6	2015	462
Ni-BCZY	BCZY (12 μm)	PLNCN	1.01	0.33	0.32	0.65	3.6	2017	360
Ni-BCZD	BCZD (15 µm)	NBN	0.97	0.89	0.44	1.33	1.7	2018	361
Ni-BCZY	BCZY (15 μm)	PS0.8N	1.03	0.40	2.17	2.57	3.7	2018	435
Ni-BCZY	BCZY (15 μm)	LS0.8N	1.03	0.34	1.48	1.82	4.4	2018	435
Ni-BCZY15	BCZY15 (5 μm)	PN	1.02	0.08	0.28	0.36	6.2	2018	362
Ni-BCZY	BCZY (20 µm)	LN	1.03	0.36	0.88	1.24	5.5	2018	459
Ni-BCZY	BCZY (20 μm)	LN-13 wt% LNF	1.03	0.37	0.36	0.73	5.4	2018	459
Ni-BCZY	BCZY (20 μm)	LN-23 wt% LNF	1.01	0.37	0.34	0.71	5.4	2018	459
Ni-BCZY	BCZY (20 μm)	LN-31 wt% LNF	0.99	0.29	0.31	0.60	6.8	2018	459
Ni-BCZY	BCZY (20 μm)	LN-37 wt% LNF	0.99	0.41	0.54	0.95	4.9	2018	459
Ni-BCZY62	BCZY62 (20 μm)	PN	0.98	0.22	0.71	0.93	9.0	2018	232
Ni-BCGC	BCGC (25 µm)	PCN-BCGC LNF c	1.08	_	_	1.86	_	2018	463
Ni-BCZY	BCZY (15 μm)	LSNF	1.06	0.29	0.41	0.70	5.1	2019	363
Ni-BCZD	BCZD (25 µm)	PB10N-BCZD	1.08	0.52	0.39	0.91	4.8	2019	364
$ PB10N-BCZD^d $									
Ni-BCZYYb	BCZYYb (15 μm)	PN	0.99	0.28	0.58	0.86	5.3	2019	365
Ni-BCZYYb	LC BCZYYb (20 μm)	PN	0.95	0.49	1.11	1.60	4.0	2019	365
Ni-BCZYYb	BCZYYb (35 µm)	LNC	1.05	0.72	0.70	1.42	4.8	2019	157
Ni-BCZD	BCZD (25 µm)	PB5N LNF	1.02	0.37	0.25	0.62	6.7	2020	464
Ni-BCZYYC	BCZYYC (13 µm)	LS0.8N-BCZYYC	1.01	0.19	0.36	0.55	6.8	2020	436
Ni-BCZY53	BCZY53 (—)	LCN	1.02	0.18	0.28	0.46	_	220	465
Ni-BCZD	BCZD (30 μm)	PB10N-BCZD	1.03	0.43	0.18	0.61	8.3	2020	466

regimes of PCFCs or PCECs, only one electrode side operates under oxidising conditions, while the other is fed by reducing gases. Therefore, the performance of Ln₂NiO_{4+δ}-based electrodes also needs to be considered under real conditions.

Applied peculiarities of Ln₂NiO_{4+δ}based electrodes in proton-conducting electrochemical cells

This section highlights the features of applying Ln₂NiO₄-based electrodes in PCFCs or PCECs. Firstly, open-circuit voltage (OCV) conditions will be considered, and then the performance of the electrodes will be evaluated under electrochemical cell operating conditions.

6.1. OCV conditions

The OCV conditions of any electrochemical cell represent important information since they allow conclusions to be formulated regarding cell integrity and electrolyte electron conductivity. Although OCV measurements ($U_{\rm OC} = \sim 1 \, {
m V}$) are far from the conditions corresponding to the maximum power density (~ 0.5 V) or thermoneutral voltage (~ 1.3 V), the

electrochemical characterisation of PCFCs/PCECs is almost always performed at OCV. In detail, the impedance spectra are measured at OCV applying a low acceleration voltage (20-50 mV). The spectra obtained in this way provide the corresponding information on ohmic resistance (R_0) and polarisation resistance of the electrodes (R_p) .

Considering R_0 , the quality of the electrolytes used (in terms of composition and fabrication technique) can be estimated as follows:

$$\sigma = \frac{h}{R_{\rm O}} \tag{25}$$

where h is the thickness of the electrolyte. It is clear that lower ohmic loss can be achieved by developing thin electrolytes. However, the conductivity of the Ba(Ce,Zr)O₃-based electrolytes is not higher than 10 mS cm⁻¹ at 600 °C (Table 9). Therefore, performance improvement is also limited in terms of developing new compositions fabricated at an optimal thickness (10-30 μm).

Considering R_p , the main strategies for improving the electrode performance were described in Section 5.4. As shown in Table 9, the polarisation resistances of Ln₂NiO₄-based electrodes are quite low, mostly reaching values below 1 Ω cm² at

Table 10 Design and performance of PCFCs with nickelate-based cathodes. Visualisation of the data is presented in Fig. S5, ESI

				P_{max} , n	nW cm	2		
Fabrication details	Anode	Electrolyte ^a (thickness)	Cathode ^b	500 °C	600 °C	700 °C	Year	Ref.
Co-pressing/screen-printing	Ni-BCY	BCY (40 μm)	PN	_	96	_	2010	252
Tape-casting/screen-printing	Ni-BCY	BCY (27 μm)	NN	_	60	_	2013	460
Co-pressing/screen-printing	Ni-BCY	BCY (50 μm)	PN	_	184	_	2014	461
Co-pressing/painting	Ni-BCZY	BCZY (20 μm)	LN	_	196	398	2014	430
Co-pressing/painting	Ni-BCZY	BCZY (20 μm)	70% LN-30% LNF	_	298	590	2014	430
Co-pressing/painting	Ni-BCZY	BCZY (20 μm)	50% LN-50% LNF	_	273	532	2014	430
Co-pressing/painting	Ni-BCZY	BCZY (20 μm)	30% LN-70% LNF	_	264	486	2014	430
Co-pressing/painting	Ni-BCZY	BCZY (24 µm)	50% LN-50% LNF	_	266	490	2015	431
Co-pressing/spin coating	Ni-BZY	BCZY44 (5 μm)	PN	_	102	234	2015	462
Tape-calendering/painting	Ni-BCZD	BCZD (30 µm)	LN	_	130	215	2016	467
Co-pressing/screen-printing	Ni-BCZY	BCZY (12 µm)	PLNCN	_	420	770	2017	360
Co-pressing/screen-printing	Ni-BCZY	BCZY (15 µm)	PS0.8N	_	127	352	2018	435
Co-pressing/screen-printing	Ni-BCZY	BCZY (15 µm)	LS0.8N	_	223	461	2018	435
Co-pressing/screen-printing	Ni-BCZY15	BCZY15 (5 μm)	PN	240	560	_	2018	362
Co-pressing/painting	Ni-BCZY	BCZY (20 µm)	LN	_	210	461	2018	459
Co-pressing/painting	Ni-BCZY	BCZY (20 μm)	LN-13 wt% LNF	_	361	623	2018	459
Co-pressing/painting	Ni-BCZY	BCZY (20 µm)	LN-23 wt% LNF	_	369	877	2018	459
Co-pressing/painting	Ni-BCZY	BCZY (20 μm)	LN-31 wt% LNF	_	369	969	2018	459
Co-pressing/painting	Ni-BCZY	BCZY (20 μm)	LN-37 wt% LNF	_	302	802	2018	459
Dry-pressing/dip-coating/screen-printing	Ni-BCZY	BCZY (16 μm)	LS0.8N	_	220	460	2018	468
Dry-pressing/dip-coating/screen-printing	Ni-BCZYYC	BCZYYC (13 µm)	LS0.8N	_	250	680	2018	469
Tape-calendering/painting	Ni-BCGC	BCGC (25 µm)	PCN-BCGC LNF ^c	_	61	132	2018	463
Co-pressing/painting	Ni-BCZY	BCZY (15 µm)	LSNF	139	421	782	2019	363
Tape-calendering	Ni-BCZD PB10N -BCZD	BCZD (25 μm)	PB10N-BCZD	_	305	395	2019	364
Tape-calendaring/spraying	Ni-BCZYYb	BCZYYb (35 µm)	LNC	_	200	340	2019	157
Co-pressing/screen-printing	Ni-BCZYYC	BCZYYC (13 μm)	LS0.8N-BCZYYC	_	540	1220	2020	436

^a Electrolyte abbreviations: BCY = BaCe_{0.9}Y_{0.1}O_{3−δ}, BCZY = BaCe_{0.7}Zr_{0.1}Y_{0.2}O_{3−δ}, BZY = BaZr_{0.85}Y_{0.15}O_{3−δ}, BCZY44 = BaCe_{0.4}Zr_{0.4}Y_{0.2}O_{3−δ}, BCZD = BaCe_{0.5}Zr_{0.3}Dy_{0.2}O_{3−δ}, BCZY15 = BaCe_{0.55}Zr_{0.3}Y_{0.15}O_{3−δ}, BCZY62 = BaCe_{0.6}Zr_{0.2}Y_{0.2}O_{3−δ}, BCC = BaCe_{0.89}Gd_{0.1}Cu_{0.01}O_{3−δ}, BCZYYb = BaCe_{0.7}Zr_{0.1}Y_{0.2}Yb_{0.1}O_{3−δ}, LC = La₂Ce₂O₃, and BCZYYC = BaCe_{0.68}Zr_{0.1}Y_{0.1}Yb_{0.1}Cu_{0.02}O_{3−δ}, believed abbreviations: PN = Pr₂NiO_{4+δ}, NN = BaCe_{0.68}Zr_{0.1}Y_{0.1}Yb_{0.1}Cu_{0.02}O_{3−δ}, believed abbreviations: PN = Pr₂NiO_{4+δ}, NN = BaCe_{0.68}Zr_{0.1}Y_{0.1}Yb_{0.1}Cu_{0.02}O_{3−δ}, believed abbreviations: PN = Pr₂NiO_{4+δ}, NN = Pr₂NiO_{4+δ}, anode (reduced PB10N-BCZD) is used together with the supported anode layer (Ni-BCZD).

Table 11 Design and performance of PCECs with nickelate-based anodes

			$U_{\rm OC},{ m V}$		i _{@1.3 V} , m	nA cm ⁻²		
Anode ^a	Electrolyte ^b (thickness)	Cathode	600 °C	700 °C	600 °C	700 °C	Year	Ref.
LN	BCZD (30 μm)	Ni-BCZD	1.06	1.00	180	300	2016	467
NBN	BCZD (15 μm)	Ni-BCZD	0.97	0.94	155	400	2018	361
PS0.8N	BCZY (15 μm)	Ni-BCZY	1.03	0.97	350	1120	2018	435
LS0.8N	BCZY (15 µm)	Ni-BCZY	1.01	0.98	420	1410	2018	435 and 468
NBN-BCZD35	BCZD35 (50 μm)	Ni-BCZD35	_	0.9	_	542	2018	470
PN	BCZY62 (20 μm)	Ni-BCZY62	0.98	0.93	345	950	2018	232
LS0.8N	BCZYYC (13 µm)	Ni-BCZYYC	_	1.01	590	1960	2018	469
PB10N-BCZD	BCZD (25 μm)	Ni-BCZD PB10N-BCZD ^c	1.08	1.01	295	535	2019	364
PN	BCZYYb (15 µm)	Ni-BCZYYb	0.99	0.95	620	1643	2019	365
PN	LC BCZYYb (20 μm)	Ni-BCZYYb	0.95	0.90	330	975	2019	365
$PB5N LNF^d$	BCZD (25 μm)	Ni-BCZD	1.02	0.94	610	855	2020	464
NBNF	BCZYYb53 (25 µm)	BCZYYb53	0.99	0.93	360	1370	2020	207
LS0.8N-BCZYYC	BCZYYC (13 µm)	Ni-BCZYYC	1.01	0.99	1040	3020	2020	436

 $[^]a \text{ Anode abbreviations: LN} = \text{La}_2\text{NiO}_{4+\delta}, \text{NBN} = \text{Nd}_{1.95}\text{Ba}_{0.05}\text{NiO}_{4+\delta}, \text{PS0.8N} = \text{Pr}_{1.2}\text{Sr}_{0.8}\text{NiO}_{4+\delta}, \text{LS0.8N} = \text{La}_{1.2}\text{Sr}_{0.8}\text{NiO}_{4+\delta}, \text{PN} = \text{Pr}_2\text{NiO}_{4+\delta}, \text{PB10N} = \text{Pr}_{1.9}\text{Ba}_{0.1}\text{NiO}_{4+\delta}, \text{PB5N} = \text{Pr}_{1.85}\text{Ba}_{0.05}\text{NiO}_{4+\delta}, \text{LNF} = \text{LaNi}_{0.6}\text{Fe}_{0.4}\text{O}_{3-\delta}, \text{and NBNF} = \text{Nd}_{1.9}\text{Ba}_{0.1}\text{NiO}_{4+\delta}\text{Fo}_{0.05}. \\ & \text{BcZy} = \text{BaCe}_{0.5}\text{Zr}_{0.3}\text{Dy}_{0.2}\text{O}_{3-\delta}, \text{BCZY} = \text{BaCe}_{0.7}\text{Zr}_{0.1}\text{Y}_{0.2}\text{O}_{3-\delta}, \text{BCZD35} = \text{BaCe}_{0.3}\text{Zr}_{0.5}\text{Dy}_{0.2}\text{O}_{3-\delta}, \text{BCZY62} = \text{BaCe}_{0.6}\text{Zr}_{0.2}\text{Y}_{0.2}\text{O}_{3-\delta}, \\ & \text{BcZyyb} = \text{BaCe}_{0.7}\text{Zr}_{0.1}\text{Y}_{0.1}\text{O}_{1.02}\text{O}_{3-\delta}, \text{BCZYYb} = \text{BaCe}_{0.7}\text{Zr}_{0.1}\text{Y}_{0.2}\text{V}_{0.1}\text{O}_{3-\delta}, \\ & \text{CF}_{0.5}\text{LC} = \text{La}_2\text{Ce}_2\text{O}_7, \text{and BCZYYb53} = \text{BaZr}_{0.3}\text{Ce}_{0.5}\text{Y}_{0.1}\text{Yb}_{0.1}\text{O}_{3-\delta}, \\ & \text{CF}_{0.5}\text{LC} = \text{La}_2\text{Ce}_2\text{O}_7, \\ & \text{CF}_{0.5}\text{LC} = \text{La}_2\text{Ce}_2\text{Ce}_7, \\ & \text{CF}_{0.5}\text{LC} = \text{LA}_2$

600 °C. The best results were achieved by An et al. 362 utilising an undoped Pr₂NiO_{4+δ} electrode and Sun et al. 436 utilising an infiltrated $La_{1.2}Sr_{0.8}NiO_{4-\delta} - BaCe_{0.68}Zr_{0.1}Y_{0.1}Yb_{0.1}Cu_{0.02}O_{3-\delta}$ electrode. Looking ahead, it can be noted that these cells also exhibit the highest performance in fuel cell (Table 10) and electrolysis cell (Table 11) operation modes. However, these correlations are not always observed since the electrode response may be considerably (and differently) altered with a variation in the bias. A vivid example of this conclusion was shown in the work by Danilov et al.,471 where the impedance data for a protonic ceramic electrochemical cell with a Pr_{1.95}- $Ba_{0.05}NiO_{4+\delta}$ (PBN) electrode was collected not only for OCV, but for the fuel cell (FC) and electrolysis cell (EC) regimes. As can be seen, the overall polarisation resistance of the electrodes monotonously decreased with an increase in bias at low measured temperatures, having a maximum OCV at higher temperatures. This peculiarity indicates that the impedance spectroscopy measurements should be performed under real voltage regimes instead of OCV.

6.2. Fuel cell and electrolysis cell operation modes

Tables 10 and 11 list the evolution of existing results on the application of $\rm Ln_2NiO_4\text{-}based$ electrodes with proton-conducting electrolytes. According to the data, it can be seen that the intensive activity initiated in 2017 formed the basis for the work published in the following year. Although some research groups have demonstrated superior results for both PCFCs and PCECs, a significant number of works showed that the maximum power densities range between 200 and 400 mW cm $^{-2}$ (fuel cell mode) at 600 °C, while the current densities range between 200 and 600 mA cm $^{-2}$ (electrolysis mode) at 1.3 V and the same temperature.

Principally, the literature data reports more promising performances, which can be associated with the utilisation of highly conductive electrode materials, such as simple and double cobaltites and their derivatives. Tonsidering the current distribution, this indeed promotes the efficient supply or removal of electrons from the TPB regions to an external electrical circuit. However, as was mentioned in Section 5.2, even if excellent results are achieved for the first time, Co-based oxide-based electrodes suffer from thermo-mechanical incompatibility. Consequently, thermo-cycling and the long-term operation of cells featuring these electrodes may be unsuccessful.

The problem of the low conductivity of nickelates (below 100 S cm⁻¹ for a gas-tight body) can be solved using additional current collector layers (for example, (La,Sr)MnO₃, La(Ni,Fe)O₃, Pt), which improve the current distribution characteristics by covering the surface of the origin electrode.

Generally, the achieved average levels are quite promising for the development of intermediate-temperature PCFCs and PCECs in terms of their up-scaling and durability. For example, Dailly and Marrony showed that a 3.5 \times 3.5 cm² PCFC with an Nd2NiO4+ δ cathode exhibits no visible degradation at 600 °C for over 950 h.

7. Conclusion and future perspectives

The materials of the $\rm Ln_2NiO_{4+\delta}$ -related class represent a promising alternative to the conventional ABO₃ electrodes in terms of their application in solid oxide electrochemical cells, including protonic ceramic fuel cells (PCFCs) and protonic ceramic electrolysis cells (PCECs). Exhibiting high structural flexibility towards different types of doping, these layered nickelates enable purposeful tailoring of functional properties important

for prospective application, including high phase stability, chemical and mechanical compatibility, and good mixed ionicelectronic conductivity and electrochemical performance.

Numerous experimental results published over the last five years have shown that lanthanide nickelate-based electrodes demonstrate quite good performances, with polarisation resistance values as low as 0.3–1 Ω cm² at 600 °C (open circuit voltage regime of PCFCs and PCFCs), which can remain stable over long term operation. These results constitute a possible basis for the scalable fabrication of durable PCFCs/PCECs.

However, although this review describes many advantages of layered nickelates in comparison with other possible oxygen electrodes, a number of pending issues still need to be solved/ clarified:

- (1) Degree of hydration/protonation of Ln₂NiO_{4+δ}-based materials. Revealing the effect of type and concentration of dopant(s) on the concentration of protons. Identification of mechanisms of proton transport in the layered phases. This information is needed for the design of triple-conducting materials having a high active electrode zone towards occurring (oxygen- and hydrogen-involved) electrochemical reactions. In addition, the mechanism studies for proton migration in layered nickelates are interesting and potentially important for proton transportation in oxides. Since oxygen vacancies are usually regarded as the prerequisite for hydration/protonation, the interstitial oxygens in layered nickelates may play a different role in the proton migration procedure in comparison with that in oxygen vacancy-containing oxides. This difference not only leads to a different strategy for designing materials, but also may inspire scientists to have a better understanding of the mechanism for proton migration since the whole landscape has not been revealed to date.
- (2) Rational design of microstructural parameters. The provided analysis was carried out considering macroscopic parameters such as polarisation resistance. On the one hand, these parameters depend on internal solid phase properties; on the other hand, the performance is also affected by the electrode structure. Currently, no general information regarding the morphological properties of Ln₂NiO_{4+δ}-based electrodes exists in the literature. Although nano-structuring of electrodes is a widely employed approach for improving the electrode performance, the high calcination temperatures for achieving the desired $Ln_2NiO_{4+\delta}$ can be detrimental for the formation and/ or maintenance of the nanostructure of the electrode. Thus, the use of advanced preparation techniques to lower the phase formation temperatures and shorten the dwell time can be beneficial for the formation of nanostructured electrodes with improved TPBs. Moreover, advanced electrode development techniques (three-dimensional ordered pore, finger-like, and nanowires-based structures),477-480 which deserve further studies, have not been used for Ln₂NiO_{4+δ}.
- (3) New fabrication approaches to improve the electrochemical activity of Ln₂NiO_{4+δ}-based electrodes, in particular, adaptation of the plasma spraying technique, pulsed laser deposition, and exsolution, or optimisation of the impregnation method.481-484

- (4) Further tailoring the functionality of Ln₂NiO_{4+δ}-based materials by introducing new dopants or their combination in appropriate concentrations at the Ln-, Ni- or O-sublattices (see, for example, the latest works485-489). The (co)doping strategy has great potential due to the previously-mentioned flexibility of the RP structures. The design of suitable composite materials also belongs to this direction. 315,490,491
- (5) Long-term features of $Ln_2NiO_{4+\delta}$ -based electrodes. To fabricate low-cost and performance-competitive PCFCs and PCECs, special attention should be devoted to the kineticrelated processes taking place over longer than 1000 h, including particle coarsening, porosity changes and chemical reactivity with CO₂, high concentrations of H₂O (in the case of electrolysis cells), and electrolyte components. All these factors affect the TPB length and its corresponding catalytic activity.

In summary, layered $Ln_2NiO_{4+\delta}$ oxides show potential suitability for application in protonic ceramic electrochemical cells as electrode materials due to their unique features. However, the application of layered Ln₂NiO_{4+δ} oxides has only emerged in recent years and many issues are still not clear, leaving great room for further exploration. The investigation of layered Ln2- $NiO_{4+\delta}$ oxides not only provides a solution for achieving high cell performance, which is the primary goal for practical applications, but also leads to an in-depth understanding of the scientific issues for the community, possibly opening a new door for the development of protonic ceramic electrochemical

Conflicts of interest

There are no conflicts to declare.

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References

- 1 J. Speirs, P. Balcombe, P. Blomerus, M. Stettler, P. Achurra-Gonzalez, M. Woo, D. Ainalis, J. Cooper, A. Sharafian, W. Merida, D. Crow, S. Giarola, N. Shah, N. Brandon and A. Hawkes, Prog. Energy, 2020, 2, 012002.
- 2 F. Dawood, M. Anda and G. M. Shafiullah, Int. J. Hydrogen Energy, 2020, 45, 3847-3869.

- 3 Z. Abdin, A. Zafaranloo, A. Rafiee, W. Mérida, W. Lipiński and K. R. Khalilpour, *Renewable Sustainable Energy Rev.*, 2020, **120**, 109620.
- 4 A. Martin, M.-F. Agnoletti and E. Brangier, *Int. J. Hydrogen Energy*, 2020, **45**, 11889–11900.
- 5 G. Yang, C. Su, H. Shi, Y. Zhu, Y. Song, W. Zhou and Z. Shao, Energy Fuels, 2020, DOI: 10.1021/acs.energyfuels.0c01887.
- V. Venkataraman, M. Pérez-Fortes, L. Wang,
 Y. S. Hajimolana, C. Boigues-Muñoz, A. Agostini,
 S. J. McPhail, F. Maréchal, J. V. Herle and P. V. Aravinda,
 J. Energy Storage, 2019, 24, 100782.
- 7 A. M. Abdalla, S. Hossain, A. T. Azad, P. M. I. Petra, F. Begum, S. G. Eriksson and A. K. Azad, *Renewable Sustainable Energy Rev.*, 2018, 82, 353–368.
- 8 Y. Meng, J. Gao, Z. Zhao, J. Amoroso, J. Tong and K. S. Brinkman, *J. Mater. Sci.*, 2019, **54**, 9291–9312.
- 9 J. Kim, S. Sengodan, S. Kim, O. Kwon, Y. Bu and G. Kim, Renewable Sustainable Energy Rev., 2019, 109, 606-618.
- 10 N. L. R. M. Rashid, A. A. Samat, A. A. Jais, M. R. Somalu, A. Muchtar, N. A. Baharuddin, W. N. Roslam and W. Isahak, Ceram. Int., 2019, 45, 6605-6615.
- 11 W. Wang, D. Medvedev and Z. Shao, *Adv. Funct. Mater.*, 2019, 28, 1802592.
- 12 D. Medvedev, Int. J. Hydrogen Energy, 2020, 44, 26711-26740.
- 13 S. Hossain, A. M. Abdalla, S. N. BintiJamain, J. H. Zaini and A. K. Azad, *Renewable Sustainable Energy Rev.*, 2017, **79**, 750–764.
- 14 D. A. Medvedev, J. G. Lyagaeva, E. V. Gorbova, A. K. Demin and P. Tsiakaras, *Prog. Mater. Sci.*, 2016, 75, 38–79.
- 15 C. Duan, R. J. Kee, H. Zhu, C. Karakaya, Y. Chen, S. Ricote, A. Jarry, E. J. Crumlin, D. Hook, R. Braun, N. P. Sullivan and R. O'Hayre, *Nature*, 2018, 557, 217–222.
- 16 C. Duan, J. Tong, M. Shang, S. Nikodemski, M. Sanders, S. Ricote, A. Almansoori and R. O'Hayre, *Science*, 2015, 349, 1321–1326.
- 17 S. Choi, C. J. Kucharczyk, Y. Liang, X. Zhang, I. Takeuchi, H.-I. Ji and S. M. Haile, *Nat. Energy*, 2018, 3, 202–210.
- 18 S. Choi, T. C. Davenport and S. M. Haile, *Energy Environ. Sci.*, 2019, **12**, 206–215.
- 19 K. Bae, D. Y. Jang, H. J. Choi, D. Kim, J. Hong, B.-K. Kim, J.-H. Lee, J.-W. Son and J. H. Shim, *Nat. Commun.*, 2017, 8, 14553.
- 20 H. An, H.-W. Lee, B.-K. Kim, J.-W. Son, K. J. Yoon, H. Kim, D. Shin, H.-I. Ji and J.-H. Lee, *Nat. Energy*, 2018, 3, 870–875.
- 21 Y. Wang, W. Li, L. Ma, W. Li and X. Liu, *J. Mater. Sci. Technol.*, 2020, **55**, 35–55.
- 22 M. Reisert, A. Aphale and P. Singh, *Materials*, 2018, 11, 2169.
- 23 B. Koo, K. Kim, J. K. Kim, H. Kwon, J. W. Han and W. Jung, *Joule*, 2018, **2**, 1476–1499.
- 24 Z. Yang, M. Guo, N. Wang, C. Ma, J. Wang and M. Han, *Int. J. Hydrogen Energy*, 2017, 42, 24948–24959.
- 25 L. Fan, B. Zhu, P.-C. Su and C. He, *Nano Energy*, 2018, 45, 148–176
- 26 W. Zhou, X. Wang, Y. Zhu, J. Dai, Y. Zhu and Z. Shao, *Mater. Rev.*, 2018, 32, 337–356.

- 27 A. V. Kasyanova, L. R. Tarutina, A. O. Rudenko, J. G. Lyagaeva and D. A. Medvedev, *Russ. Chem. Rev.*, 2020, **89**, 667–692.
- 28 P. Kaur and K. Singh, Ceram. Int., 2020, 46, 5521-5535.
- 29 S. Afroze, A. Karim, Q. Cheok, S. Eriksson and A. K. Azad, Front. Energy, 2019, 13, 770–797.
- 30 S. S. Hashim, F. Liang, W. Zhou and J. Sunarso, *ChemElectroChem*, 2019, **6**, 3549–3569.
- 31 M. Rafique, H. Nawaz, M. S. Rafique, M. B. Tahir, G. Nabi and N. R. Khalid, *Energy Res.*, 2019, **43**, 2423–2446.
- 32 G. Nirala, D. Yadav and S. Upadhyay, *J. Adv. Ceram.*, 2020, 9, 129–148.
- 33 D. Lee, Y.-L. Lee, A. Grimaud, W. T. Hong, M. D. Biegalski, D. Morgane and Y. Shao-Horn, *J. Mater. Chem. A*, 2014, 2, 6480–6487.
- 34 S. Choi, S. Yoo, J.-Y. Shin and G. Kim, *J. Electrochem. Soc.*, 2011, **158**, B995–B999.
- 35 A. R. Gilev, E. A. Kiselev and V. A. Cherepanov, *Solid State Ionics*, 2019, **339**, 115001.
- 36 M. Garali, M. Kahlaoui, B. Mohammed, A. Mater, C. B. Azouz and C. Chefi, *Int. J. Hydrogen Energy*, 2019, 44, 11020–11032.
- 37 B. Philippeau, F. Mauvy, C. Mazataud, S. Fourcade and J.-C. Grenier, *Solid State Ionics*, 2013, 249–250, 17–25.
- 38 X. Zhang, L. Zhang, J. Meng, W. Zhang, F. Meng, X. Liu and J. Meng, *Int. J. Hydrogen Energy*, 2017, 42, 29498–29510.
- 39 V. A. Sadykov, E. M. Sadovskaya, N. F. Eremeev, P. I. Skriabin, A. V. Krasnov, Y. N. Bespalko, S. N. Pavlova, Y. E. Fedorova, E. Y. Pikalova and A. V. Shlyakhtina, *Russ. J. Electrochem.*, 2019, 55, 701–718.
- 40 S. N. Ruddlesden and P. Popper, *Acta Crystallogr.*, 1958, **11**, 54–55.
- 41 M. Greenblatt, Curr. Opin. Solid State Mater. Sci., 1997, 2, 174–183.
- 42 A. Demourgues, A. Wattiaux, J. C. Grenier, M. Pouchard, J. L. Soubeyroux, J. M. Dance and P. Hagenmuller, *J. Solid State Chem.*, 1993, **105**, 458–470.
- 43 A. Aguadero, J. A. Alonso, M. J. Martínez-Lope, M. T. Fernández-Díaz, M. J. Escudero and L. Daza, J. Mater. Chem., 2006, 16, 3402–3408.
- 44 J. F. Ackerman, Mater. Res. Bull., 1979, 14, 487-491.
- 45 K. K. Singh, P. Ganguly and J. B. Goodenough, *J. Solid State Chem.*, 1984, **52**, 254–273.
- 46 P. Ganguly and C. N. R. Rao, *J. Solid State Chem.*, 1984, 53, 193–216.
- 47 C. Frayret, A. Villesuzanne and M. Pouchard, *Chem. Mater.*, 2005, **17**, 6538–6544.
- 48 R. D. Shannon, Acta Crystallogr., Sect. A: Found. Crystallogr., 1996, A32, 751–767.
- 49 M. T. Fernández-Díaz, J. L. Martínez and J. Rodríguez-Carvajal, *Solid State Ionics*, 1993, **63–65**, 902–906.
- 50 C. Allançon, J. Rodríguez-Carvajal, M. T. Fernández-Díaz, P. Odier, J. M. Bassat, J. P. Loup and J. L. Martínez, Z. Phys. B: Condens. Matter, 1996, 100, 85–90.
- 51 J. M. Bassat, P. Odier and J. P. Loup, *J. Solid State Chem.*, 1994, **110**, 124–135.

- 52 A. V. Kovalevsky, V. V. Kharton, A. A. Yaremchenko, Y. V. Pivak, E. V. Tsipis, S. O. Yakovlev, A. A. Markov, E. N. Naumovich and J. R. Frade, J. Electroceram., 2007, 18, 205-218.
- 53 E. Boehm, J.-M. Bassat, P. Dordor, F. Mauvy, J.-C. Grenier and P. Stevens, Solid State Ionics, 2005, 176, 2717-2725.
- 54 C. Allançon, A. Gonthier-Vassal, J. M. Bascat, J. P. Loup and P. Odier, Solid State Ionics, 1994, 74, 239-248.
- 55 H. Tamura, A. Hayashi and Y. Ueda, Solid State Ionics, 1996, 258, 61-71.
- 56 T. Nakamura, K. Yashiro, K. Sato and J. Mizusaki, Solid State Ionics, 2009, 180, 368-376.
- 57 V. Vashook, E. Girdauskaite, J. Zosel, T.-L. Wen, H. Ullmann and U. Guth, Solid State Ionics, 2006, 177, 1163-1171.
- 58 T. Nakamura, K. Yashiro, K. Sato and J. Mizusaki, J. Solid State Chem., 2009, 182, 1533-1537.
- 59 T. Nakamura, K. Yashiro, K. Sato and J. Mizusaki, Solid State Ionics, 2009, 180, 1406-1413.
- 60 V. V. Kharton, E. V. Tsipis, E. N. Naumovich, A. Thursfield, M. V. Patrakeev, V. A. Kolotygin, J. C. Waerenborgh and I. S. Metcalfe, J. Solid State Chem., 2008, 181, 1425-1433.
- 61 E. N. Naumovich, M. V. Patrakeev, V. V. Kharton, A. Yaremchenko, D. I. Logvinovich F. M. B. Marques, Solid State Sci., 2005, 7, 1353-1362.
- 62 E. N. Naumovich, M. V. Patrakeev, V. V. Kharton, A. Yaremchenko, E. V. Tsipis, J. R. D. I. Logvinovich and F. M. B. Marques, Mater. Sci. Forum, 2006, 514-516, 397-401.
- 63 H.-S. Kim and H. Yoo, Phys. Chem. Chem. Phys., 2010, 12, 4704-4713.
- 64 S.-Y. Jeon, M.-B. Choi, J.-H. Hwang, E. D. Wachsman and S.-J. Song, J. Solid State Electrochem., 2012, 16, 785-793.
- 65 Y. Naumovich, Quantitative description of oxygen nonstoichiometry in mixed ionic and electronic conductors based on a non-ideal solution approach, Institute of Power Engineering, Warsaw, 2019.
- 66 H. Tamura, A. Hayashi and Y. Ueda, Phys. C, 1993, 216, 83-
- 67 S. J. Skinner, Solid State Sci., 2003, 5, 419-426.
- 68 J. D. Jorgensen, B. Dabrowski, S. Pei, D. R. Richards and D. G. Hinks, Phys. Rev. B: Condens. Matter Mater. Phys., 1989, 40, 2187.
- 69 T. Broux, C. Prestipino, M. Bahout, S. Paofai, E. Elkaïm, V. Vibhu, J.-C. Grenier, A. Rougier, J.-M. Bassat and O. Hernandez, Dalton Trans., 2016, 45, 3024-3033.
- 70 V. Vibhu, M. R. Suchomel, N. Penin, F. Weill, J.-C. Grenier, J.-M. Bassat and A. Rougier, Dalton Trans., 2019, 48, 266-
- 71 S. M. Aspera, M. Sakaue, T. D. K. Wungu, M. Alaydrus, T. P. T. Linh, H. Kasai, M. Nakanishi and T. Ishihara, J. Phys.: Condens. Matter, 2012, 24, 405504.
- 72 A. Flura, S. Dru, C. Nicollet, V. Vibhu, S. Fourcade, E. Lebraud, A. Rougier, J.-M. Bassat and J.-C. Grenier, J. Solid State Chem., 2015, 228, 189-198.
- 73 R. K. Sharma, S.-K. Cheah, M. Burriel, L. Dessemond, J.-M. Bassat and E. Djurado, J. Mater. Chem. A, 2017, 5, 1120-1132.

- 74 A. Montenegro-Hernández, J. Vega-Castillo, A. Caneirom and L. Mogni, I. Solid State Chem., 2019, 276, 210-216.
- 75 D. Ning, A. Baki, T. Scherb, J. Song, A. Fantin, X. Liu, G. Schumacher, J. Banhart and H. J. M. Bouwmeester, Solid State Ionics, 2019, 342, 115056.
- 76 Y. Toyosumi, H. Ishikawa and K. Ishikawa, J. Alloys Compd., 2006, 408-412, 1200-1204.
- 77 M. Zaghrioui, F. Giovannelli, N. P. D. Brouri and I. Laffez, J. Solid State Chem., 2004, 177, 3351-3358.
- 78 S. M. Pikalov, L. B. Vedmid, E. A. Filonov, E. Y. Pikalova, J. G. Lyagaeva, N. A. Danilov and A. A. Murashkina, J. Alloys Compd., 2019, 801, 558-567.
- 79 Y. Takeda, M. Nishijima, N. Imanishi, R. Kanno, O. Yamamoto and M. Takano, J. Solid State Chem., 1992,
- 80 T. Ogier, C. Prestipino, S. Figueroa, F. Mauvy, J. Mougin, J. C. Grenier, A. Demourgues and J. M. Bassat, Chem. Phys. Lett., 2019, 727, 116-120.
- 81 J. Dailly, S. Fourcade, A. Largeteau, F. Mauvy, J. C. Grenier and M. Marrony, Electrochim. Acta, 2010, 55, 5847-5853.
- 82 E. Y. Pikalova, A. A. Kolchugin, V. A. Sadykov, E. M. Sadovskaya, E. A. Filonova, N. F. Eremeev and N. M. Bogdanovich, Int. J. Hydrogen Energy, 2018, 43, 17373-17386.
- 83 C. Allançon, P. Odier, J. M. Bassat and J. P. Loup, J. Solid State Chem., 1997, 131, 167-172.
- 84 A. D. Rougier, A. Flura, C. Nicollet, V. Vibhu, S. Fourcade, E. Lebraud, J.-M. Bassat and J.-C. Grenier, ECS Trans., 2015, 68, 817-823.
- 85 H. Zhao, F. Mauvy, C. Lalanne, J.-M. Bassat, S. Fourcade and J.-C. Grenier, Solid State Ionics, 2008, 179, 2000–2005.
- 86 A. Egger, E. Bucher, W. Sitte, C. Lalanne and J.-M. Bassat, ECS Trans., 2009, 25, 2547-2556.
- 87 B. Gedziorowski, K. Cichy, A. Niemczyk, A. Olszewska, Z. Zhang, S. Kopec, K. Zheng, M. Marzec, M. Gajewska, Z. Du, H. Zhao and K. Swierczek, J. Eur. Ceram. Soc., 2020, 40, 4056-4066.
- 88 E. Niwa, T. Nakamura, J. Mizusaki and T. Hashimoto, Thermochim. Acta, 2011, 523, 46-50.
- 89 H. Ishikawa, Y. Toyosumi and K. Ishikawa, J. Alloys Compd., 2006, 406-412, 1196-1199.
- 90 E. Niwa, K. Wakai, T. Hori, K. Yashiro, J. Mizusaki and T. Hashimoto, Thermochim. Acta, 2014, 575, 129-134.
- 91 T. Nakamura, K. Yashiro, K. Sato and J. Mizusaki, Solid State Ionics, 2010, 181, 402-411.
- 92 V. Sadykov, Y. Okhlupin, N. Yeremeev, Z. Vinokurov, A. Shmakov, V. Belyaev, N. Uvarov and J. Mertens, Solid State Ionics, 2014, 262, 918-922.
- 93 V. Vibhu, A. Rougier, C. Nicollet, A. Flura, J.-C. Grenier and J.-M. Bassat, Solid State Ionics, 2015, 278, 32-37.
- 94 A. Murata, C. Hai and M. Matsuda, Mater. Lett., 2014, 136, 292-294.
- 95 M. Sakai, C. Wang, T. Okiba, H. Soga and T. Hashimoto, J. Therm. Anal. Calorim., 2020, 142, 139-147.
- 96 E. Dogdibegovic, Q. Cai, N. S. Alabri, W. Guan and X.-D. Zhou, J. Electrochem. Soc., 2016, 164, F99-F106.

- 97 G. Amow and S. J. Skinner, *J. Solid State Electrochem.*, 2006, 10, 538–546.
- 98 J. P. Tang, R. I. Dass and A. Manthiram, *Mater. Res. Bull.*, 2000, 35, 411–424.
- 99 K. Ruck, G. Krabbes and I. Vogel, *Mater. Res. Bull.*, 1999, 34, 1689–1697.
- 100 Y. Shen, H. Zhao, X. Liua and N. Xu, *Phys. Chem. Chem. Phys.*, 2010, **12**, 15124–15131.
- 101 E. Y. Pikalova, D. A. Medvedev and A. F. Khasanov, *Phys. Solid State*, 2010, 59, 694–702.
- 102 V. A. Sadykov, E. Y. Pikalova, A. A. Kolchugin, E. A. Filonova, E. M. Sadovskaya, N. F. Eremeev, A. V. Ishchenko, A. V. Fetisov and S. M. Pikalov, *Solid State Ionics*, 2018, 317, 234–243.
- 103 S. W. Li and Y. F. Ren, Mater. Res. Bull., 1995, 30, 1505-1511.
- 104 M. James and J. P. Attfield, J. Mater. Chem., 1996, 6, 57-62.
- 105 Y. Takeda, R. Kanno, M. Sakano and O. Yamamoto, *Mater. Res. Bull.*, 1990, 25, 293–306.
- 106 E. S. Kravchenko, K. V. Zakharchuk, V. V. Pankov and A. A. Yaremchenko, *J. Belarus. State Univ. Chem.*, 2017, 1, 43–49.
- 107 A. Yaremchenko, E. Kravchenko, K. Zakharchuk, J. Grins, G. Svensson and V. Pankov, Oxygen-deficient Ruddlesden–Popper $\Pr_{1-x} Sr_{1+x} NiO_{4-\delta}$ as prospective oxygen electrode materials for SOFC/SOEC, 21^{st} International Conference on Solid State Ionics (SSI-21), Padua, Italy, 18–23 June 2017, Program Guide and Abstracts, abstract I-12_26/O, p. 365.
- 108 E. Kravchenko, D. Khalyavin, K. Zakharchuk, J. Grins, G. Svensson, V. Pankov and A. Yaremchenko, *J. Mater. Chem. A*, 2015, 3, 23852–23863.
- 109 A. B. Austin, L. G. Carreiro and J. V. Marzik, *Mater. Res. Bull.*, 1989, 24, 639–646.
- 110 Y. K. Chung, Y.-U. Kwon and S. H. Byeon, *Bull. Korean Chem. Soc.*, 1995, **16**, 120–125.
- 111 H. Soga, C. Wang, T. Hayashi, T. Morise, E. Niwa and T. Hashimoto, ECS Trans., 2017, 78, 613–622.
- 112 K. K. Singh, P. Ganguly and C. N. R. Rao, *Mater. Res. Bull.*, 1982, 17, 493–500.
- 113 E. Boehm, J.-M. Bassat, M. C. Steil, P. Dordor, F. Mauvy and J.-C. Grenier, *Solid State Sci.*, 2003, 5, 973–981.
- 114 A. Aguadero, J. A. Alonso, M. J. Escudero and L. Daza, *Solid State Ionics*, 2008, **179**, 393–400.
- 115 M. Sakai, C. Wang, T. Okiba, H. Soga, E. Niwa and T. Hashimoto, *J. Therm. Anal. Calorim.*, 2019, **135**, 2765–2774.
- 116 C. Wang, H. Soga, T. Okiba, E. Niwa and T. Hashimoto, *Mater. Res. Bull.*, 2019, **111**, 61–69.
- 117 G. Amow, P. S. Whitfield, I. J. Davidson, R. P. Hammond, C. N. Munnings and S. J. Skinner, *Ceram. Int.*, 2004, **30**, 1635–1639.
- 118 C. Berger, E. Bucher, A. Egger, A. T. Strasser, N. Schrödl, C. Gspan, J. Hofer and W. Sitte, *Solid State Ionics*, 2018, 316, 93–101.
- 119 S. Li, H. Tu, F. Li, M. T. Anwar and L. Yu, *J. Alloys Compd.*, 2017, **694**, 17–23.

- 120 T. Nakamura, Y. Ling and K. Amezawa, *J. Mater. Chem. A*, 2015, 3, 10471–10479.
- 121 E. V. Tsipis, E. N. Naumovich, M. V. Patrakeev, J. C. Waerenborgh, Y. V. Pivak, P. Gaczyński and V. V. Kharton, *J. Phys. Chem. Solids*, 2007, **68**, 1443–1455.
- 122 M.-L. Fontaine, C. Laberty-Robert, A. Barnabé, F. Ansart and P. Tailhades, *Ceram. Int.*, 2004, **30**, 2087–2098.
- 123 A. K. Ganguli, R. Nagarajan, G. R. Rao, N. Y. Vasanthacharya and C. N. R. Rao, Solid State Commun., 1989, 72, 195–197.
- 124 R. Benloucif, N. Nguyen, J. M. Greneche and B. Raveau, *J. Phys. Chem. Solids*, 1989, **50**, 435–440.
- 125 C. Lalanne, G. Prosperi, J.-M. Bassat, F. Mauvy, S. Fourcade, P. Stevens, M. Zahid, S. Diethelm, J. Van herle and J.-C. Grenier, *J. Power Sources*, 2008, 185, 1218–1224.
- 126 F. Mauvy, C. Lalanne, J.-M. Bassat, J.-C. Grenier, A. Brisse, A.-L. Sauvet, C. Barthet and J. Fouletier, *Solid State Ionics*, 2009, 180, 1183–1189.
- 127 Y. Chen, B. Qian, G. Yang, D. Chen and Z. Shao, *J. Mater. Chem. A*, 2015, 3, 6501–6508.
- 128 M. Ferkhi and H. A. Yahia, *Mater. Res. Bull.*, 2016, **83**, 268–274.
- 129 A. Usenka, V. Pankov, V. Vibhu, A. Flura, J.-C. Grenier and J.-M. Bassat, *ECS Trans.*, 2019, **91**, 1341–1353.
- 130 A. A. Kolchugin, E. Y. Pikalova, N. M. Bogdanovich, D. I. Bronin, S. M. Pikalov, S. V. Plaksin, M. V. Ananyev and V. A. Eremin, *Solid State Ionics*, 2016, 288, 48–53.
- 131 A. R. Gilev, E. A. Kiselev, D. M. Zakharov and V. A. Cherepanov, *J. Alloys Compd.*, 2018, 753, 491–501.
- 132 E. Pikalova, A. Kolchugin, E. Filonova, N. Bogdanovich, S. Pikalov, M. Ananyev, N. Molchanova and A. Farlenkov, *Solid State Ionics*, 2018, **319**, 130–140.
- 133 M. Khairy, P. Odier and J. Choisnet, J. Phys., Colloq., 1986, 47, C1-831–C1-835.
- 134 I. Gopalakrishnan, G. Colsmann and B. Reuter, *J. Solid State Chem.*, 1977, 22, 145–149.
- 135 T. Nitadori, M. Muramatsu and M. Misono, *Bull. Chem. Soc. Jpn.*, 1988, **61**, 3831–3837.
- 136 K. Sreedhar and C. N. R. Rao, *Mater. Res. Bull.*, 1990, 25, 1235–1242.
- 137 X. Granados, J. Fontcuberta, M. Vallet-Regi, M. J. Sayague and J. M. Gonza-Calbet, *J. Solid State Chem.*, 1993, 102, 455–464.
- 138 L. V. Makhnach, V. V. Pankov and P. Strobel, *Mater. Chem. Phys.*, 2008, **111**, 125–130.
- 139 V. V. Vashook, S. P. Tolochko, I. I. Yushkevich, L. V. Makhnach, I. F. Kononyuk, H. Altenburg, J. Hauck and H. Ullmann, *Solid State Ionics*, 1998, **110**, 245–253.
- 140 V. V. Vashook, I. I. Yushkevich, L. V. Kokhanovsky, L. V. Makhnach, S. P. Tolochko, I. F. Kononyuk, H. Ullmann and H. Altenburg, *Solid State Ionics*, 1999, 119, 23–30.
- 141 A. Aguadero, M. J. Escudero, M. Perez, J. A. Alonso, V. Pomjakushin and L. Daza, *Dalton Trans.*, 2006, 4377– 4383
- 142 T. Nakamura, K. Yashiro, K. Sato and J. Mizusaki, *Phys. Chem. Chem. Phys.*, 2009, **11**, 3055–3062.

- 143 S. C. Chen, K. V. Ramanujachary and M. Greenblatt, J. Solid State Chem., 1993, 105, 444-457.
- 144 A. Grimaud, F. Mauvy, J. M. Bassat, S. Fourcade, M. Marrony and J. C. Grenier, J. Mater. Chem., 2012, 22, 16017-16025.
- 145 E. Kravchenko, K. Zakharchuk, A. Viskup, J. Grins, G. Svensson, V. Pankov and A. Yaremchenko, ChemSusChem, 2017, 10, 600-611.
- 146 S. M. Doyle, M. P. S. Kumar and D. M. Paul, J. Phys.: Condens. Matter, 1992, 4, 3559-3568.
- 147 B. W. Arbuckle, K. V. Ramanujachary, Z. Zhang and M. Greenblatt, J. Solid State Chem., 1990, 88, 278-290.
- 148 T. Nakamura, K. Yashiro, K. Sato and J. Mizusaki, Mater. Chem. Phys., 2010, 122, 250-258.
- 149 J. Alonso, M. Vallet-Regi and J. M. Gonzalez-Calbet, Solid State Ionics, 1993, 66, 219-223.
- 150 M. Hücker, K. Chung, M. Chand, T. Vogt, J. M. Tranquada and D. J. Buttrey, Phys. Rev. B: Condens. Matter Mater. Phys., 2004, 70, 064105.
- 151 M. Jimenez-Ruiz, C. Prieto, J. L. Martinez and J. M. Alonso, J. Solid State Chem., 1998, 140, 278-284.
- 152 C. J. Liu, M. D. Mays, D. O. Cowan and M. G. Sanchez, Chem. Mater., 1991, 3, 495-500.
- 153 H.-S. Kim and H.-I. Yoo, Solid State Ionics, 2013, 232, 129-
- 154 Z. Li, R. Haugsrud, J. B. Smith and T. Norby, Solid State Ionics, 2009, 180, 1433-1441.
- 155 E. Kravchenko, A. Neagu, K. Zakharchuk, J. Grins, G. Svensson, V. Pankov and A. A. Yaremchenko, Eur. J. Inorg. Chem., 2018, 3320-3329.
- 156 A. Aguadero, M. Pereza, J. A. Alonso and L. Daza, J. Power Sources, 2005, 151, 52-56.
- 157 A. P. Tarutin, J. G. Lyagaeva, A. S. Farlenkov, A. I. Vylkov and D. M. Medvedev, Ceram. Int., 2019, 45, 16105-16112.
- 158 K. Zheng and K. Swierczek, Mater. Res. Bull., 2016, 84, 259-
- 159 Y. Wang, J. Cheng, Q. Jiang, J. Yang and J. Gao, J. Power Sources, 2011, 196, 3104-3108.
- 160 J. A. Kilner and C. K. M. Shaw, Solid State Ionics, 2002, 154-**155**, 523-527.
- 161 E. V. Tsipis, E. N. Naumovich, A. L. Shaula, M. V. Patrakeev, J. C. Waerenborgh and V. V. Kharton, Solid State Ionics, 2008, 179, 57-60.
- 162 S. Miyoshi, T. Furuno, O. Sangoanruang, H. Matsumoto and T. Ishihara, J. Electrochem. Soc., 2007, 154, B57-B62.
- 163 T. Klande, K. Efimov, S. Cusenza, K.-D. Becker and A. Feldhoff, J. Solid State Chem., 2011, 184, 3310-3318.
- 164 S. Crapanzano, I. V. Babich and L. Lefferts, Appl. Catal., A, 2010, 378, 144-150.
- 165 S. Y. Gómez, J. Gurauskis, V. Øygarden, D. Hotza, T. Grande and K. Wiik, Solid State Ionics, 2016, 292, 38-44.
- 166 J. Zhu, D. Xiao, J. Li and X. Yang, Catal. Lett., 2009, 129, 240-246.
- 167 H. El Shinawi and C. Greaves, J. Mater. Chem., 2010, 20,
- 168 S. Yi, Y. Shen, H. Zhao, Z. Du, N. Chen and B. Huang, Electrochim. Acta, 2016, 219, 394-400.

- 169 W. Li, B. Guan, J. Yan, N. Zhang, X. Zhang and X. Liu, J. Power Sources, 2016, 318, 178-183.
- 170 L. Y. Gavrilova, T. V. Aksenova, L. A. Bannykh, Y. V. Teslenko and V. A. Cherepanov, J. Struct. Chem., 2003, 44, 248-251.
- 171 A. R. Gilev, E. A. Kiselev and V. A. Cherepanov, Solid State Ionics, 2015, 279, 53-59.
- 172 A. R. Gilev, E. A. Kiselev and V. A. Cherepanov, RSC Adv., 2016, 6, 72905-72917.
- 173 R. Benloucif, N. Nguyen, J. M. Greneche and B. Raveau, J. Phys. Chem. Solids, 1991, 52, 381-387.
- 174 R. P. Forslund, W. G. Hardin, X. Rong, A. M. Abakumov, D. Filimonov, C. T. Alexander, J. T. Mefford, H. Iyer, A. M. Kolpak, K. P. Johnston and K. J. Stevenson, Nat. Commun., 2018, 9, 3150.
- 175 H. Chaker, T. Roisnel, M. Potel and R. B. Hassen, J. Solid State Chem., 2004, 177, 4067-4072.
- 176 H. Chaker, I. Raies, A. Chouket, T. Roisnel and R. B. Hassen, Ionics, 2017, 23, 2229-2240.
- 177 S. Y. Istomin, O. M. Karakulina, M. G. Rozova, S. M. Kazakov, A. A. Gippius, E. V. Antipov, I. A. Bobrikov, A. M. Balagurov, A. A. Tsirlin, A. Michau, J. J. Biendicho and G. Svensson, RSC Adv., 2016, 6, 33951-33958.
- 178 A. Chouket, O. Bidault, L. Combemale, O. Heintz, M. Khitouni and V. Optasanu, J. Alloys Compd., 2018, 732, 149-159.
- 179 J. E. Millburn and M. J. Rosseinsky, Chem. Mater., 1997, 9, 511-522.
- 180 F. Tonus, M. Bahout, P. D. Battle, T. Hansen, P. F. Henry and T. Roisnel, J. Mater. Chem., 2010, 20, 4103-4115.
- 181 M. Jammali, R. B. Hassen and J. Rohlicek, *Powder Diffr.*, 2012, 27, 184-188.
- 182 H. Chaker, T. Roisnel, M. Ceretti and R. B. Hassen, *Powder* Diffr., 2010, 25, 241-246.
- 183 M. Jammali, H. Chaker, K. Cherif and R. B. Hassen, J. Mater. Chem., 2010, 183, 1194-1199.
- 184 S. E. Dutton, M. Bahout, P. D. Battle, F. Tonus and V. Demange, J. Solid State Chem., 2008, 181, 2217-2226.
- 185 T. Ishihara, S. Miyoshi, T. Furuno, O. Sanguanruang and H. Matsumoto, Solid State Ionics, 2006, 177, 3087-3091.
- 186 N. Ye and J. L. Hertz, Acta Mater., 2014, 63, 123-129.
- 187 K. Meeporn, N. Chanlek and P. Thongbai, RSC Adv., 2016, 6, 91377-91385.
- 188 C. A. Silva, J. B. Silva, M. C. Silva-Santana, P. B. Silva and N. O. Moreno, Adv. Mater. Res., 2014, 975, 75-80.
- 189 S.-Y. Jeon, B. Singh, Y.-S. Yoo, J.-H. Hwang and S.-J. Song, J. Am. Ceram. Soc., 2014, 97, 1489-1496.
- 190 S.-Y. Jeon, B. Singh, H.-N. Im, K.-P. Seong and S.-J. Song, J. Alloys Compd., 2014, 589, 572-578.
- 191 D. Reinen, U. Kesper and D. Belder, J. Solid State Chem., 1995, 116, 355-363.
- 192 V. F. Savchenko, L. V. Makhnach, I. I. Emelyanova and V. V. Pankov, Inorg. Mater., 2006, 42, 788-793.
- 193 Z. Zhang, Z. Du, A. Niemczyk, K. Li, H. Zhao and K. Swierczek, Solid State Ionics, 2018, 317, 26-31.
- 194 A. Kawahara and T. Ishihara, Electrochem. Solid-State Lett., 2010, 13, B76-B78.

- 195 T. Ishihara, N. Sirikanda, K. Nakashima, S. Miyoshi and H. Matsumoto, *J. Electrochem. Soc.*, 2010, 157, B141–B146.
- 196 Q. Zhou, L. Qu, T. Zhang, Y. He, C. Zhao, M. Wang, T. Wei and Y. Zhang, *J. Alloys Compd.*, 2020, **824**, 153967.
- 197 M. Bansod, A. P. Khandale and S. S. Bhoga, *Ionics*, 2017, 23, 2561–2570.
- 198 T. Ishihara, K. Nakashima, S. Okada, M. Enoki and H. Matsumoto, *Solid State Ionics*, 2008, **179**, 1367–1371.
- 199 J. Mao, S. Peng, C. Zhang, S. Qi, J. Cui, Y. Gong, S. Wang, C. Wu and Q. Zhou, *J. Alloys Compd.*, 2019, **793**, 519–525.
- 200 Q. Zheng, J. Xue, Q. Liao, Y. Wei, Z. Li and H. Wang, *Chem. Eng. Sci.*, 2013, **101**, 240–247.
- 201 H. W. Nie, T.-L. Wen, S. R. Wang, Y. S. Wang, U. Guth and V. Vashook, *Solid State Ionics*, 2006, **177**, 1929–1932.
- 202 H. Lou, Y. Ge, P. Chen, M. Mei, F. Ma and G. Lü, *J. Mater. Chem.*, 1997, 7, 2097–2101.
- 203 Q. Li, Y. Fan, H. Zhao, L. P. Sun and L. H. Huo, *J. Power Sources*, 2007, **167**, 64–68.
- 204 V. N. Chaudhari, A. P. Khandale and S. S. Bhoga, *J. Power Sources*, 2014, **248**, 647–654.
- 205 V. Bhat, C. N. R. Rao and J. M. Honig, *Solid State Commun.*, 1992, **81**, 751–756.
- 206 B. W. Arbuckle, K. V. Ramanujachary, A. M. Buckley and M. Greenblatt, *J. Mater. Chem.*, 1992, **97**, 274–282.
- 207 A. P. Tarutin, G. K. Vdovin, D. A. Medvedev and A. A. Yaremchenko, *Electrochim. Acta*, 2020, 337, 135808.
- 208 K. Wissel, J. Heldt, P. B. Groszewicz, S. Dasgupta, H. Breitzke, M. Donzelli, A. I. Waidha, A. D. Fortes, J. Rohrer, P. R. Slater, G. Buntkowsky and O. Clemens, *Inorg. Chem.*, 2018, 57, 6549–6560.
- 209 M. Ali Nowroozi, K. Wissel, M. Donzelli, N. Hosseinpourkahvaz, S. Plana-Ruiz, U. Kolb, R. Schoch, M. Bauer, A. M. Malik, J. Rohrer, S. Ivlev, F. Kraus and O. Clemens, *Commun. Mater.*, 2020, 1, 27.
- 210 M. Zinkevich and F. Aldinger, *J. Alloys Compd.*, 2004, 375, 147–161.
- 211 A. N. Petrov, V. A. Cherepanov, A. Y. Zuyev and V. M. Zhukovsky, *J. Solid State Chem.*, 1988, 77, 1–14.
- 212 D. E. Rice and D. J. Buttrey, *J. Solid State Chem.*, 1993, **105**, 197–210.
- 213 A. A. Yaremchenko, V. V. Kharton, M. V. Patrakeev and J. R. Frade, *J. Mater. Chem.*, 2003, **13**, 1136–1144.
- 214 V. V. Kharton, A. A. Yaremchenko, A. L. Shaula, M. V. Patrakeev, E. N. Naumovich, D. I. Logvinovich, J. R. Frade and F. M. BMarques, *J. Solid State Chem.*, 2004, 177, 26–37.
- 215 P. Odier, Y. Nigara, J. Coutures and M. Sayer, *J. Solid State Chem.*, 1985, 56, 32–40.
- 216 M. Zinkevich, N. Solak, H. Nitsche, M. Ahrens and F. Aldinger, J. Alloys Compd., 2007, 438, 92–99.
- 217 Y. Adachi, N. Hatada, K. Hirota, M. Kato and T. Uda, *J. Am. Ceram. Soc.*, 2019, **102**, 7077–7088.
- 218 G. Amow, I. J. Davidson and S. J. Skinner, *Solid State Ionics*, 2006, **177**, 1205–1210.
- 219 N. Gauquelin, T. E. Weirich, M. Ceretti, W. Paulus and M. Schroeder, *Monatsh. Chem.*, 2009, **140**, 1095–1102.

- 220 A. Montenegro Hernández, L. Mogni and A. Caneiro, *Int. J. Hydrogen Energy*, 2010, 35, 6031–6036.
- 221 R. Sayers and S. J. Skinner, *J. Mater. Chem.*, 2011, **21**, 414–419.
- 222 J. D. Sullivan, D. J. Buttrey, D. E. Cox and J. Hriljac, J. Solid State Chem., 1991, 94, 337–351.
- 223 H. S. C. O'Neill and M. I. Pownceby, *Contrib. Mineral. Petrol.*, 1993, **114**, 296–314.
- 224 P. Odier, C. Allançon and J. M. Bassat, *J. Solid State Chem.*, 2000, **153**, 381–385.
- 225 A. V. Kovalevsky, V. V. Kharton, A. A. Yaremchenko, Y. V. Pivak, E. N. Naumovich and J. R. Frade, *J. Eur. Ceram. Soc.*, 2007, 27, 4269–4272.
- 226 A. Montenegro-Hernandez, J. Vega-Castillo, L. Mogni and A. Caneiro, *Int. J. Hydrogen Energy*, 2011, 36, 15704–15714.
- 227 S. Saher, J. Song, V. Vibhu, C. Nicollet, A. Flura, J.-M. Bassat and H. J. M. Bouwmeester, *J. Mater. Chem. A*, 2018, **6**, 8331–8339.
- 228 V. Vibhu, J.-M. Bassat, A. Flura, C. Nicollet, J.-C. Grenier and A. Rougier, *ECS Trans.*, 2015, **68**, 825–835.
- 229 V. Vibhu, A. Flura, A. Rougier, C. Nicollet, S. Fourcade, T. Hungria, J.-C. Grenier and J.-M. Bassat, *J. Energy Chem.*, 2020, 46, 62–70.
- 230 J. Xue, A. Schulz, H. Wang and A. Feldhoff, *J. Membr. Sci.*, 2016, **497**, 357–364.
- 231 P. Batocchi, F. Mauvy, S. Fourcade and M. Parco, *Electrochim. Acta*, 2014, 145, 1–10.
- 232 W. Li, B. Guan, L. Ma, S. Hu, N. Zhang and X. Liu, *J. Mater. Chem. A*, 2018, **6**, 18057–18066.
- 233 S. Upasen, P. Batocchi, F. Mauvy, A. Slodczyk and P. Colomban, J. Alloys Compd., 2015, 622, 1074–1085.
- 234 S. Engels, T. Markus, M. Modigell and L. Singheiser, J. Membr. Sci., 2011, 370, 58-69.
- 235 Q. Wei, S. Zhang, B. Meng, N. Han, Z. Zhu and S. Liu, *Mater. Lett.*, 2018, **230**, 161–165.
- 236 N. Han, Q. Wei, H. Tian, S. Zhang, Z. Zhu, J. Liu and S. Liu, Energy Technol., 2018, 7, 1800701.
- 237 G. Chen, M. Widenmeyer, B. Tang, L. Kaeswurm, L. Wang, A. Feldhoff and A. Weidenkaff, Front. Chem. Sci. Eng., 2020, 14, 405–414.
- 238 K. S. Yun, J. H. Park, Y. Kwon, D. Y. Kim, C.-Y. Yoo, J. H. Yu and J. H. Joo, *J. Mater. Chem. A*, 2016, 4, 13549–13554.
- 239 Y. Kwon, J. H. Park, S. M. Kang, G. D. Nam, J. W. Lee, J. H. Kim, D. Kim, S. M. Jeong, J. H. Yu and J. H. Joo, *Energy Environ. Sci.*, 2019, 12, 1358–1368.
- 240 S.-N. Lee, A. Atkinson and J. Kilner, *ECS Trans.*, 2013, 57, 605–613.
- 241 K. J. Lee, J. H. Chung, M. J. Lee and H. J. Hwang, *J. Korean Ceram. Soc.*, 2019, **56**, 160–166.
- 242 M. Yang, E. Bucher and W. Sitte, *J. Power Sources*, 2011, **196**, 7313–7317.
- 243 Y. Gong, R. Wang, J. Banner, S. N. Basu, U. B. Pal and S. Gopalan, *JOM*, 2019, 71, 3848–3858.
- 244 A. Egger, N. Schrodl, C. Gspan and W. Sitte, *Solid State Ionics*, 2017, **299**, 18–25.

- 245 N. Schrodl, E. Bucher, A. Egger, P. Kreiml, C. Teichert, T. Hoschen and W. Sitte, *Solid State Ionics*, 2015, **276**, 62–71.
- 246 N. Schrodl, A. Egger, C. Gspan, T. Hoschen, F. Horer and W. Sitte, *Solid State Ionics*, 2018, 322, 44–53.
- 247 D.-P. Huang, Q. Xu, W. Chen, F. Zhang and H.-X. Liu, *Ceram. Int.*, 2008, 34, 651–655.
- 248 M. Chen, B. H. Moon, S. H. Kim, B. H. Kim, Q. Xu and B. -G. Ahn, *Fuel Cells*, 2012, **12**, 86–96.
- 249 D. Huang, Q. Xu, F. Zhang, W. Chen, H. Liu and J. Zhou, Mater. Lett., 2006, 60, 1892–1895.
- 250 S. Nishiyama, D. Sakaguchi and T. Hattori, *Solid State Commun.*, 1995, **94**, 279–282.
- 251 K. Ishikawa, W. Shibata, K. Watanabe, T. Isonaga, M. Hashimoto and Y. Suzuki, *J. Solid State Chem.*, 1997, 131, 275–281.
- 252 G. Taillades, J. Dailly, M. Taillades-Jacquin, F. Mauvy, A. Essouhmi, M. Marrony, C. Lalanne, S. Fourcade, D. J. Jones, J. -C. Grenier and J. Roziere, *Fuel Cells*, 2010, 10, 166–173.
- 253 A. A. Kolchugin, E. Y. Pikalova, N. M. Bogdanovich, D. I. Bronin and E. A. Filonova, *Russ. J. Electrochem.*, 2017, 53, 826–833.
- 254 J. B. Goodenough, Mater. Res. Bull., 1973, 8, 423-431.
- 255 J. B. Goodenough and S. Ramasesha, *Mater. Res. Bull.*, 1982, 17, 383–390.
- 256 J. B. Goodenough, J. Less-Common Met., 1986, 116, 83-93.
- 257 N. Poirot, P. Odier, P. Simon and F. Gervais, *Solid State Sci.*, 2003, **5**, 735–739.
- 258 S.-Y. Jeon, M.-B. Choi, J.-H. Hwang, E. D. Wachsman and S.-J. Song, *J. Electrochem. Soc.*, 2011, **158**, B476–B480.
- 259 C. N. R. Rao, D. J. Buttrey, N. Otsuka, P. Ganguly, H. R. Harrison, C. J. Sandberg and J. M. Honig, *J. Solid State Chem.*, 1984, **51**, 266–269.
- 260 J. M. Bassat, F. Gervais, P. Odier and J. P. Loup, *Mater. Sci. Eng.*, *B*, 1989, 3, 507–514.
- 261 K. Sugiyama, H. Nozaki, T. Takeuchi and H. Ikuta, *J. Phys. Chem. Solids*, 2002, **63**, 979–982.
- 262 A. Murata, T. Uchikoshi and M. Matsuda, *J. Power Sources*, 2015, **293**, 95–100.
- 263 A. L. Shaula, E. N. Naumovich, A. P. Viskup, V. V. Pankov, A. V. Kovalevsky and V. V. Kharton, *Solid State Ionics*, 2009, **180**, 812–816.
- 264 J. M. Bassat, J. P. Loup and P. Odier, *J. Phys.: Condens. Matter*, 1994, **6**, 8285–8293.
- 265 C. Solís, L. Navarrete and J. M. Serra, *J. Power Sources*, 2013, **240**, 691–697.
- 266 Y. Shen, H. Zhao, J. Xu, X. Zhang, K. Zheng and K. Swierczek, *Int. J. Hydrogen Energy*, 2014, 39, 1023–1029.
- 267 V. A. Sadykov, E. M. Sadovskaya, E. Y. Pikalova, A. A. Kolchugin, E. A. Filonova, S. M. Pikalov, N. F. Eremeev, A. V. Ishchenko, A. I. Lukashevich and J. M. Bassat, *Ionics*, 2018, 24, 1181–1193.
- 268 J. Yang, J. Cheng, Q. Jiang, Y. Wang, R. Wang and J. Gao, *Int. J. Hydrogen Energy*, 2012, **37**, 1746–1751.
- 269 V. V. Vashook, N. E. Trofimenko, H. Ullmann and L. V. Makhnach, *Solid State Ionics*, 2000, **131**, 329–336.

- 270 T. Klande, S. Cusenza, P. Gaczynski, K.-D. Becker, L. Dorrer, G. Borchardt and A. Feldhoff, *Solid State Ionics*, 2012, 222–223, 8–15.
- 271 M. V. Patrakeev, E. N. Naumovich, V. V. Kharton, A. A. Yaremchenko, E. V. Tsipis, P. Nunez and J. RFrade, *Solid State Ionics*, 2005, **176**, 179–188.
- 272 V. V. Kharton, A. P. Viskup, E. N. Naumovich and F. M. B. Marques, *J. Mater. Chem.*, 1999, **9**, 2623–2629.
- 273 V. V. Kharton, A. P. Viskup, A. V. Kovalevsky, E. N. Naumovich and F. M. B. Marques, *Solid State Ionics*, 2001, **143**, 337–353.
- 274 B. I. Arias-Serrano, E. Kravchenko, K. Zakharchuk, J. Grins, G. Svensson, V. Pankov and A. Yaremchenko, ECS Trans., 2019, 91, 2387–2397.
- 275 Y. Shen, H. Zhao, K. Świerczek, Z. Du and Z. Xie, *J. Power Sources*, 2013, **240**, 759–765.
- 276 J. Hyodo, K. Tominaga, Y.-W. Ju, S. Ida and T. Ishihara, *Solid State Ionics*, 2014, **256**, 5–10.
- 277 J. M. Bassat, P. Odier, A. Villesuzanne, C. Marin and M. Pouchard, *Solid State Ionics*, 2004, **167**, 341–347.
- 278 H.-S. Kim and H.-I. Yoo, *Phys. Chem. Chem. Phys.*, 2011, **13**, 4651–4658.
- 279 H.-S. Kim and H.-I. Yoo, Phys. Chem. Chem. Phys., 2014, 16, 16595–16605.
- 280 A. R. Gilev, E. A. Kiselev, D. S. Chezganov and V. A. Cherepanov, *Ceram. Int.*, 2020, **46**, 17553–17560.
- 281 A. R. Gilev, E. A. Kiselev and V. A. Cherepanov, *J. Mater. Chem. A*, 2018, **6**, 5304–5312.
- 282 V. A. Sadykov, E. Y. Pikalova, A. A. Kolchugin, A. V. Fetisov, E. M. Sadovskaya, E. A. Filonova, N. F. Eremeev, V. B. Goncharov, A. V. Krasnov, P. I. Skriabin, A. N. Shmakov, Z. S. Vinokurov, A. V. Ishchenko and S. M. Pikalov, *Int. J. Hydrogen Energy*, 2020, 45, 13625–13642.
- 283 F. Mauvy, E. Boehm, J. M. Bassat, J. C. Grenier and J. Fouletier, *Solid State Ionics*, 2007, **178**, 1200–1204.
- 284 S.-Y. Jeon, Y.-S. Yoo, B. Singh, H.-N. Im and S.-J. Song, *J. Electrochem. Soc.*, 2016, **163**, F1302–F1307.
- 285 J. Xue, Q. Liao, W. Chen, H. J. M. Bouwmeester, H. Wang and A. Feldhoff, *J. Mater. Chem. A*, 2015, **3**, 19107–19114.
- 286 T. Ina, Y. Orikasa, T. Masese, T. Nakao, A. Mineshige, K. Amezawa, H. Tanida, T. Uruga and Y. Uchimoto, *Electrochemistry*, 2014, **82**, 875–879.
- 287 A. Egger and W. Sitte, Solid State Ionics, 2014, 258, 30-37.
- 288 W. Li, B. Guan, X. Zhang, J. Yan, Y. Zhoua and X. Liu, *Phys. Chem. Chem. Phys.*, 2016, **18**, 8502–8511.
- 289 Z. Li and R. Haugsrud, Solid State Ionics, 2012, 206, 67-71.
- 290 S. J. Skinner and J. A. Kilner, *Solid State Ionics*, 2000, 135, 709–712.
- 291 E. S. Tropin, M. V. Ananyev, A. S. Farlenkov, A. V. Khodimchuk, A. V. Berenov, A. V. Fetisov, V. A. Eremin and A. A. Kolchugin, *J. Solid State Chem.*, 2018, 262, 199–213.
- 292 J. A. Kilner, S. J. Skinner and H. H. Brongersma, *J. Solid State Electrochem.*, 2011, 15, 861–876.
- 293 T. Inprasit, S. Wongkasemjit, S. J. Skinner, M. Burriel and P. Limthongkul, *RSC Adv.*, 2015, 5, 2486–2492.

- 294 K. Develos-Bagarinao, J. D. Vero, H. Kishimoto, T. Ishiyama, K. Yamaji, T. Horita and H. Yokokawa, *Phys. Chem. Chem. Phys.*, 2019, **21**, 7183–7195.
- 295 J. Railsback, G. Hughes, L. Mogni, A. Montenegro-Hernandez and S. Barnett, *J. Electrochem. Soc.*, 2016, **163**, F1433–F1439.
- 296 K. Yakal-Kremski, L. V. Mogni, A. Montenegro-Hernandez, A. Caneiro and S. A. Barnett, *J. Electrochem. Soc.*, 2014, **161**, F1366–F1374.
- 297 J. Hyodo, K. Tominaga, Y.-W. Ju, S. Ida and T. Ishihara, *ECS Trans.*, 2014, **61**, 123–129.
- 298 S. J. Skinner and J. A. Kilner, *Ionics*, 1999, 5, 171–174.
- 299 F. Mauvy, J. M. Bassat, E. Boehm, P. Dordor and J. P. Loup, *Solid State Ionics*, 2003, **158**, 395–407.
- 300 R. Sayers, R. A. D. Souza, J. A. Kilner and S. J. Skinner, *Solid State Ionics*, 2010, **181**, 386–391.
- 301 P.-M. Geffroy, L. Guironnet, H. J. M. Bouwmeester, T. Chartier, J.-C. Grenier and J.-M. Bassat, *J. Eur. Ceram. Soc.*, 2019, **39**, 59–65.
- 302 M. Burriel, H. Tellez, R. J. Chater, R. Castaing, P. Veber, M. Zaghrioui, T. Ishihara, J. A. Kilner and J.-M. Bassat, J. Phys. Chem. C, 2016, 120, 17927–17938.
- 303 J.-M. Bassat, M. Burriel, O. Wahyudi, R. Castaing, M. Ceretti, P. Veber, I. Weill, A. Villesuzanne, J.-C. Grenier, W. Paulus and J. A. Kilner, J. Phys. Chem. C, 2013, 117, 26466–26472.
- 304 M. Burriel, G. Garcia, J. Santiso, J. A. Kilner, R. J. Chater and S. J. Skinner, *J. Mater. Chem.*, 2008, **18**, 416–422.
- 305 J. A. Lane and J. A. Kilner, *Solid State Ionics*, 2000, **136–137**, 997–1001.
- 306 Z. Li, T. Norby and R. Haugsrud, J. Am. Ceram. Soc., 2012, 95, 2065–2073.
- 307 A. Egger, E. Bucher and W. Sitte, *J. Electrochem. Soc.*, 2011, **158**, B573–B579.
- 308 G. Kim, S. Wang, A. J. Jacobson and C. L. Chen, *Solid State Ionics*, 2006, 177, 1461–1467.
- 309 G. Garcia, M. Burriel, N. Bonanos and J. Santiso, J. Electrochem. Soc., 2008, 155, P28–P32.
- 310 J. B. Smith and T. Norby, *J. Electrochem. Soc.*, 2006, **153**, A233–A238.
- 311 S.-Y. Jeon, M.-B. Choi, H.-N. Im, J.-H. Hwang and S.-J. Song, *J. Phys. Chem. Solids*, 2012, 73, 656–660.
- 312 S.-Y. Jeon, H.-N. Im, B. Singh, S.-K. Hong and S.-J. Song, *Ceram. Int.*, 2014, **40**, 16785–16790.
- 313 R. Moreno, J. Zapata, J. Roqueta, N. Bagues and J. Santiso, *J. Electrochem. Soc.*, 2014, **161**, F3046–F3051.
- 314 R. Moreno, P. Garcia, J. Zapata, J. Roqueta, J. Chaigneau and J. Santiso, *Chem. Mater.*, 2013, 25, 3640–3647.
- 315 M. Ghamarinia, A. Babaei and C. Zamani, *Electrochim. Acta*, 2020, 353, 136520.
- 316 C. Berger, A. Egger, R. Merkle, E. Bucher, B. Stuhlhofer, N. Schrodl, J. Lammer, C. Gspan, G. Logvenov and J. Maier, J. Electrochem. Soc., 2019, 166, F1088–F1095.
- 317 V. A. Sadykov, N. F. Eremeev, V. V. Usoltsev, A. S. Bobin, G. M. Alikina, V. V. Pelipenko, E. M. Sadovskaya, V. S. Muzykantov, N. N. Bulgakov and N. F. Uvarov, *Russ. J. Electrochem.*, 2013, **49**, 645–651.

- 318 N. M. Porotnikova, A. V. Khodimchuk, M. V. Ananyev, V. A. Eremin, E. S. Tropin, A. S. Farlenkov, E. Y. Pikalova and A. V. Fetisov, *J. Solid State Electrochem.*, 2018, 22, 2115–2126.
- 319 H. J. M. Bouwmeester, C. Song, J. Zhu, J. Yi, M. S. Annalanda and B. A. Boukamp, *Phys. Chem. Chem. Phys.*, 2009, **11**, 9640–9643.
- 320 A. Boreave, H. Tan, V. Roche, P. Vernoux and J.-P. Deloume, *Solid State Ionics*, 2008, **179**, 1071–1075.
- 321 L. Minervini, R. W. Grimes, J. A. Kilner and K. E. Sickafus, *J. Mater. Chem.*, 2000, **10**, 2349–2354.
- 322 M. S. D. Read, M. S. Islam, F. King and F. E. Hancock, *J. Phys. Chem. B*, 1999, **103**, 1558–1562.
- 323 A. R. Cleave, J. A. Kilner, S. J. Skinner, S. T. Murphy and R. W. Grimes, *Solid State Ionics*, 2008, **179**, 823–826.
- 324 E. N. Naumovich and V. V. Kharton, *J. Mol. Struct.: THEOCHEM*, 2010, **946**, 57–64.
- 325 A. Chroneos, D. Parfitt, J. A. Kilner and R. W. Grimes, *J. Mater. Chem.*, 2010, **20**, 266–270.
- 326 A. Piovano, A. Perrichon, M. Boehm, M. R. Johnson and W. Paulus, *Phys. Chem. Chem. Phys.*, 2016, **18**, 17398–17403.
- 327 D. Parfitt, A. Chroneos, J. A. Kilner and R. W. Grimes, *Phys. Chem. Chem. Phys.*, 2010, **12**, 6834–6836.
- 328 M. Yashima, H. Yamada, S. Nuansaeng and T. Ishihara, *Chem. Mater.*, 2012, **24**, 4100–4113.
- 329 W. Xie, Y.-L. Lee, Y. Shao-Horn and D. Morgan, *J. Phys. Chem. Lett.*, 2016, 7, 1939–1944.
- 330 V. Sadykov, E. Pikalova, N. Eremeev, A. Shubin, I. Zilberberg, I. Prosvirin, E. Sadovskaya and A. Bukhtiyarov, Solid State Ionics, 2020, 344, 115155.
- 331 S. Xu, R. Jacobs and D. Morgan, *Chem. Mater.*, 2018, **30**, 7166-7177.
- 332 L. Zhang, F. Yao, J. Meng, W. Zhang, H. Wang, X. Liu, J. Meng and H. Zhang, J. Mater. Chem. A, 2019, 7, 18558– 18567.
- 333 M. Yashima, M. Enoki, T. Wakita, R. Ali, Y. Matsushita, F. Izumi and T. Ishihara, *J. Am. Chem. Soc.*, 2008, **130**, 2762–2763.
- 334 M. Yashima, N. Sirikanda and T. Ishihara, *J. Am. Chem. Soc.*, 2010, **132**, 2385–2392.
- 335 M. Schroeder and M.-A. Dragan, *J. Mater. Sci.*, 2007, **42**, 1972–1983.
- 336 Q. Liao, L. Zhuang, Y. Wei, J. Xue and H. Wang, *Ceram. Int.*, 2018, 44, 10852–10857.
- 337 Z. Li, R. Haugsrud and T. Norby, *Solid State Ionics*, 2011, **184**, 42–46.
- 338 Z. Li, R. Haugsrud, J. B. Smith and T. Norby, *J. Electrochem. Soc.*, 2009, **156**, B1039–B1044.
- 339 Z. Zhu, M. Li, C. Xia and H. J. M. Bouwmeester, *J. Mater. Chem. A*, 2017, 5, 14012–14019.
- 340 S. Miyoshi, T. Furuno, H. Matsumoto and T. Ishihara, *Solid State Ionics*, 2006, **177**, 2269–2273.
- 341 V. V. Kharton, A. A. Yaremchenko, E. V. Tsipis and J. R. Frade, in *Proceedings of Eighth International Symposium on Solid Oxide Fuel Cells (SOFC-VIII)*, ed. S. C. Singhal and M. Dokiya, The Electrochemical Society Inc., Pennington, NJ, 2003, PV 2003-07, pp. 561–570.

- 342 N. A. Tarasova, A. O. Galisheva and I. E. Animitsa, Russ. J. Electrochem., 2019, 55, 756-761.
- 343 N. Tarasova and I. Animitsa, C. R. Chim., 2019, 22, 363-368.
- 344 N. Tarasova and I. Animitsa, Solid State Sci., 2019, 87, 87-
- 345 M. I. Vlasov, N. A. Tarasova, A. O. Galisheva, I. E. Animitsa and M. V. Ananyev, Phys. Chem. Chem. Phys., 2019, 21, 23459-23465.
- 346 N. Tarasova and I. Animitsa, J. Fluorine Chem., 2018, 216, 107-111.
- 347 N. Tarasova and I. Animitsa, Solid State Ionics, 2018, 317,
- 348 N. Tarasova and I. Animitsa, J. Alloys Compd., 2018, 739, 353-359.
- 349 J. Lyagaeva, D. Medvedev, E. Pikalova, S. Plaksin, A. Brouzgou, A. Demin and P. Tsiakaras, Int. J. Hydrogen Energy, 2017, 42, 1715-1723.
- 350 J. R. Tolchard and T. Grande, Solid State Ionics, 2007, 178, 593-599.
- 351 Y. Rao, S. Zhong, F. He, Z. Wang, R. Peng and Y. Lu, Int. J. Hydrogen Energy, 2012, 37, 12522-12527.
- 352 Y. Lin, R. Ran, Y. Zheng, Z. Shao, W. Jin, N. Xu and J. Ahn, J. Power Sources, 2008, 180, 15-22.
- 353 Y. Lin, R. Ran, C. Zhang, R. Cai and Z. Shao, J. Phys. Chem. A, 2010, **114**, 3764–3772.
- 354 D. Medvedev, A. Murashkina, E. Pikalova, A. Demin, A. Podias and P. Tsiakaras, Prog. Mater. Sci., 2014, 60, 72-129.
- 355 Y. K. Chung, Y.-U. Kwon and S. H. Byeon, Bull. Korean Chem. Soc., 1995, 16, 120-125.
- 356 Y. Zhu, D. Wang, F. Yuan, G. Zhang and H. Fu, Appl. Catal., B, 2008, 82, 255-263.
- 357 E. Quarez, Y. Oumellal and O. Joubert, Fuel Cells, 2013, 13,
- 358 C. Yang, X. Zhang, H. Zhao, Y. Shen, Z. Du and C. Zhanga, Int. J. Hydrogen Energy, 2015, 40, 2800-2807.
- 359 E. Y. Pikalova and A. A. Kolchugin, Eurasian Chem.-Technol. J., 2016, 18, 3-11.
- 360 G. Li, H. Jin, Y. Cui, L. Gui, B. He and L. Zhao, J. Power Sources, 2017, 341, 192-198.
- 361 N. Danilov, J. Lyagaeva, G. Vdovin, E. Pikalova and D. Medvedev, Energy Convers. Manage., 2018, 172, 129-137.
- 362 H. An, D. Shin and H.-I. Ji, J. Korean Ceram. Soc., 2018, 55, 358-363.
- 363 L. Miao, J. Hou, Z. Gong, Z. Jin and W. Liu, Int. J. Hydrogen Energy, 2019, 44, 7531-7537.
- 364 A. Tarutin, J. Lyagaeva, A. Farlenkov, S. Plaksin, G. Vdovin, A. Demin and D. Medvedev, *Mater*, 2019, 12, 118.
- 365 W. Li, B. Guan, L. Ma, H. Tian and X. Liu, ACS Appl. Mater. Interfaces, 2019, 11, 18323-18330.
- 366 A. Løken, S. Ricote and S. Wachowski, Crystals, 2018, 8, 365.
- 367 R. Raghvendra, R. K. Singh and P. Singh, J. Mater. Sci., 2014, 49, 5571-5578.
- 368 D. Tsvetkov, N. Tsvetkova, I. Ivanov, D. Malyshkin, V. Sereda and A. Zuev, *Energies*, 2019, **12**, 417.
- 369 P. Gao, A. Bolon, M. Taneja, Z. Xie, N. Orlovskaya and M. Radovic, Solid State Ionics, 2017, 300, 1-9.

- 370 Y. G. Lyagaeva, D. A. Medvedev, A. K. Demin, P. Tsiakaras and O. G. Reznitskikh, Phys. Solid State, 2015, 57, 285-289.
- 371 Y.-P. Wang, Q. Xu, D.-P. Huang, K. Zhao, M. Chen and B.-H. Kim, Int. J. Hydrogen Energy, 2016, 41, 6476-6485.
- 372 V. Vibhu, A. Rougier, J.-C. Grenier and J.-M. Bassat, ECS Trans., 2013, 57, 2093-2100.
- 373 Y. G. Lyagaeva, N. A. Danilov, M. Y. Gorshkov, G. K. Vdovin, B. D. Antonov, A. K. Demin and D. A. Medvedev, Russ. J. Appl. Chem., 2018, 91, 583-590.
- 374 Y.-P. Wang, Q. Xu, D.-P. Huang, K. Zhao, M. Chen and B.-H. Kim, Int. J. Hydrogen Energy, 2017, 42, 6290-6302.
- 375 A. Aguadero, J. A. Alonso, M. T. Fernandez-Diaz, M. J. Escudero and L. Daza, J. Power Sources, 2007, 169, 17-24.
- 376 V. V. Kharton, A. V. Kovalevsky, M. Avdeev, E. V. Tsipis, M. V. Patrakeev, A. A. Yaremchenko, E. N. Naumovich and J. R. Frade, Chem. Mater., 2007, 19, 2027-2033.
- 377 J. Yeyongchaiwat, K. Nonthawissarut, S. Charojrochkul and N. Sukpirom, Adv. Appl. Ceram., 2014, 114, 1-8.
- 378 X.-D. Zhou, J. W. Templeton, Z. Nie, H. Chen, J. W. Stevenson and L. R. Pederson, Electrochim. Acta, 2012, 71, 44-49.
- 379 J. Cheng, S. Zhang, B. Meng, J. Ding and X. Tan, J. Alloys Compd., 2018, 742, 966-976.
- 380 T.-W. Chiu, M.-X. Lin, H.-Y. Shih, B. Hwang, H.-Y. Chang and Y.-M. Wang, Ceram. Int., 2017, 43, S700-S704.
- 381 M. Li, J. Cheng, Y. Gan, S. Li, B. He and W. Sun, J. Power Sources, 2015, 275, 151-158.
- 382 X. Meng, S. Lü, S. Liu, X. Liu, Y. Sui, X. Li, M. Pang, B. Wang, Y. Ji and M. Z. Hu, Ceram. Int., 2015, 41, 12107-12114.
- 383 T. Chen, Y. Zhou, C. Yuan, M. Liu, X. Meng, Z. Zhan, C. Xia and S. Wang, J. Power Sources, 2014, 269, 812-817.
- 384 M. Letilly, A. L. G. L. Salle, M. Caldes, M. Marrony and O. Joubert, Fuel Cells, 2009, 9, 622-629.
- 385 L.-P. Sun, Q. Li, H. Zhao, L.-H. Huo and J.-C. Grenier, J. Power Sources, 2008, 183, 43-48.
- 386 R. Pelosato, G. Cordaro, D. Stucchi, C. Cristiani and G. Dotelli, J. Power Sources, 2015, 298, 46-67.
- 387 L. Mogni, F. Prado, C. Jimenez and A. Caneiro, Solid State Ionics, 2013, 240, 19-28.
- 388 B. Sikder, A. Chanda, S. Goswami, D. Bhattacharya and S. Velaga, Mater. Chem. Phys., 2019, 236, 121770.
- 389 L.-S. Unger, C. Niedrig, S. F. Wagner, W. Menesklou, S. Baumann, W. A. Meulenberg and E. Ivers-Tiffee, J. Eur. Ceram. Soc., 2018, 38, 2378-2387.
- 390 F. Liang, Z. Yang, H. Deng, J. Sunarso, L. Yang and J. Mao, J. Mater. Sci. Technol., 2019, 35, 1184-1191.
- 391 H. Hao, L. Zhao, J. Hu, X. Hu and H. Hou, J. Rare Earths, 2009, 27, 815-818.
- 392 N. A. Danilov, A. P. Tarutin, J. G. Lyagaeva, E. Y. Pikalova, A. A. Murashkina, D. A. Medvedev, M. V. Patrakeev and A. K. Demin, Ceram. Int., 2017, 43, 15418-15423.
- 393 W. H. Kan, K.-Y. Lai, A. Huq and A. Manthiram, J. Power Sources, 2016, 307, 454-461.
- 394 S. R. Bishop, D. Marrocchelli, C. Chatzichristodoulou, N. H. Perry, M. B. Mogensen, H. L. Tuller and E. D. Wachsman, Annu. Rev. Mater. Res., 2014, 44, 205-239.

- 395 S. R. Bishop, K. L. Duncan and E. D. Wachsman, *J. Am. Ceram. Soc.*, 2010, **93**, 4115–4121.
- 396 S. R. Bishop, K. L. Duncan and E. D. Wachsman, *Acta Mater.*, 2009, 57, 3596–3605.
- 397 T. Nakamura, K. Yashiro, K. Sato and J. Mizusaki, *Solid State Ionics*, 2010, **181**, 292–299.
- 398 S. R. Bishop, D. Marrocchelli, W. Fang, K. Amezawa, K. Yashiro and G. W. Watson, *Energy Environ. Sci.*, 2013, **6**, 1142–1146.
- 399 V. V. Kharton, A. A. Yaremchenko, M. V. Patrakeev, E. N. Naumovich and F. M. B. Marques, *J. Eur. Ceram. Soc.*, 2003, 23, 1417–1426.
- 400 V. V. Vashook, DSc. thesis, Institute of General and Inorganic Chemistry, National Belarus Academy of Sciences, Minsk, Belarus, 2000.
- 401 H. Ullmann, N. Trofimenko, F. Tietz, D. Stöver and A. Ahmad-Khanlou, *Solid State Ionics*, 2000, **138**, 79–90.
- 402 L. P. Putilov, N. A. Shevyrev, A. M. Mineev, A. S. Farlenkov, D. A. Medvedev and V. I. Tsidilkovski, *Acta Mater.*, 2020, 190, 70–80.
- 403 K. D. Kreuer, Annu. Rev. Mater. Res., 2003, 33, 333-359.
- 404 E. Fabbri, Daniele Pergolesi and Enrico Traversa, *Chem. Soc. Rev.*, 2010, **39**, 4355–4369.
- 405 N. Tarasova, A. Galisheva and I. Animitsa, *Ionics*, 2020, **26**, 5075–5088.
- 406 N. Tarasova, I. Animitsa and A. Galisheva, *J. Solid State Electrochem.*, 2020, **24**, 1497–1508.
- 407 N. A. Tarasova, A. O. Galisheva, I. E. Animitsa and D. V. Korona, *Russ. J. Phys. Chem. A*, 2020, **94**, 818–821.
- 408 N. Tarasova, I. Animitsa, A. Galisheva and V. Pryakhina, *Solid State Sci.*, 2020, **101**, 106121.
- 409 M. Chen, Y. Xuan, F. Zhang, L. He, X. Wang, H. Pan, J. Ren and Z. Lin, *Int. J. Hydrogen Energy*, 2020, **45**, 14964–14971.
- 410 W. Tan, D. Huan, W. Yang, N. Shi, W. Wang, R. Peng, X. Wu and Y. Lu, *RSC Adv.*, 2018, **8**, 26448–26460.
- 411 G. Rinaldi, A. Nakajo, P. Burdet, M. Cantoni, W. K. S. Chiu and J. Van herle, *Acta Mater.*, 2019, 178, 194–206.
- 412 H. Moussaoui, R. K. Sharma, J. Debayle, Y. Gavet, G. Delette and J. Laurencin, *J. Power Sources*, 2019, **412**, 736–748.
- 413 W. Kong, M. Zhang, Z. Han and Q. Zhang, *Appl. Sci.*, 2019, **9**, 493.
- 414 P. Vijay, M. O. Tadéa, Z. Shao and M. Ni, *Int. J. Hydrogen Energy*, 2017, **42**, 28836–28851.
- 415 X. Lu, T. M. M. Heenan, J. J. Bailey, T. Li, K. Li, D. J. L. Brett and P. R. Shearing, *J. Power Sources*, 2017, 365, 210–219.
- 416 E. V. Tsipis and V. V. Kharton, *J. Solid State Electrochem.*, 2008, **12**, 1039–1060.
- 417 X.-Y. Wu and A. F. Ghoniem, *Prog. Energy Combust. Sci.*, 2019, 74, 1–30.
- 418 A. A. Plazaola, A. C. Labella, Y. Liu, N. B. Porras, D. A. P. Tanaka, M. V. S. Annaland and F. Gallucci, *Processes*, 2019, 7, 128.
- 419 M. Acosta, F. Baiutti, A. Tarancón and J. L. MacManus-Driscoll, *Adv. Mater. Interfaces*, 2019, **6**, 1900462.
- 420 C. Li, J. J. Chew, A. Mahmoud, S. Liu and J. Sunarso, *J. Membr. Sci.*, 2018, **567**, 228–260.

- 421 C. Zhang, J. Sunarso and S. Liu, *Chem. Soc. Rev.*, 2017, **46**, 2941–3005.
- 422 V. L. Kozhevnikov, I. A. Leonidov and M. V. Patrakeev, *Russ. Chem. Rev.*, 2013, **82**, 772–782.
- 423 S. B. Adler, J. A. Lane and B. C. H. Steele, *J. Electrochem. Soc.*, 1996, 143, 3554.
- 424 Z. Zhao, J. Cui, M. Zou, S. Mu, H. Huang, Y. Meng, K. He, K. S. Brinkman and J. Tong, *J. Power Sources*, 2020, 450, 227609.
- 425 C. Xia, Y. Mi, B. Wang, B. Lin, G. Chen and B. Zhu, *Nat. Commun.*, 2019, **10**, 1707.
- 426 Y. Song, Y. Chen, W. Wang, C. Zhou, Y. Zhong, G. Yang, W. Zhou, M. Liu and Z. Shao, *Joule*, 2019, 3, 2842–2853.
- 427 K. Wei, N. Li, Y. Wu, W. Song, X. Wang, L. Guo, M. Khan, S. Wang, F. Zho and Y. Ling, *Ceram. Int.*, 2019, **45**, 18583–18591.
- 428 R. Ren, Z. Wang, C. Xu, W. Sun, J. Qiao, D. W. Rooneyc and K. Sun, *J. Mater. Chem. A*, 2019, 7, 18365–18372.
- 429 J. C. Grenier, F. Mauvy, C. Lalanne, J.-M. Bassat, F. Chauveau, J. Mougin, J. Dailly and M. Marrony, *ECS Trans.*, 2009, 25, 2537–2546.
- 430 J. Hou, Z. Zhu, J. Qian and W. Liu, *J. Power Sources*, 2014, **264**, 67–75.
- 431 J. Hou, J. Qian, L. Bi, Z. Gong, R. Penga and W. Liu, *J. Mater. Chem. A*, 2015, 3, 2207–2215.
- 432 Y. G. Lyagaeva, N. A. Danilov, M. Y. Gorshkov, G. K. Vdovin, B. D. Antonov, A. K. Demin and D. A. Medvedev, *Russ. J. Appl. Chem.*, 2018, 91, 583–590.
- 433 E. P. Antonova, A. A. Kolchugin, E. Y. Pikalova, D. A. Medvedev and N. M. Bogdanovich, *Solid State Ionics*, 2017, **306**, 55–61.
- 434 E. Y. Pikalova, N. M. Bogdanovich and A. V. Kuzmin, *Russ. J. Electrochem.*, 2017, 53, 752–760.
- 435 S. Yang, Y. Wen, J. Zhang, Y. Lu, X. Ye and Z. Wen, *Electrochim. Acta*, 2018, **267**, 269–277.
- 436 C. Sun, S. Yang, Y. Lu, J. Wen, X. Ye and Z. Wen, *J. Power Sources*, 2020, **449**, 227498.
- 437 A. P. Tarutin, M. Y. Gorshkov, I. N. Bainov, G. K. Vdovin, A. I. Vylkov, J. G. Lyagaeva and D. A. Medvedev, *Ceram. Int.*, 2020, 46, 24355–24364.
- 438 T. Park, J. Korean Soc. Precis. Eng., 2018, 35, 1141-1146.
- 439 J. Dailly, F. Mauvy, M. Marrony, M. Pouchard and J.-C. Grenier, *J. Solid State Electrochem.*, 2011, **15**, 245–251.
- 440 S. M. Babiniec, S. Ricote and N. P. Sullivan, *J. Electrochem. Soc.*, 2014, **161**, F717–F723.
- 441 A. Grimaud, F. Mauvy, J. M. Bassat, S. Fourcade, L. Rocheron, M. Marrony and J. C. Grenier, *J. Electrochem. Soc.*, 2012, **159**, B683–B694.
- 442 A. Y. Zuev, V. V. Sereda and D. S. Tsvetkov, *J. Electrochem. Soc.*, 2012, **159**, F594–F599.
- 443 N. E. Volkova, V. A. Kolotygin, L. Y. Gavrilova, V. V. Kharton and V. A. Cherepanov, *Solid State Ionics*, 2014, **260**, 15–20.
- 444 E. V. Tsipis, E. N. Naumovich, M. V. Patrakeev, A. A. Yaremchenko, I. P. Marozau, A. V. Kovalevsky, J. C. Waerenborgh and V. V. Kharton, *Solid State Ionics*, 2011, 192, 42–48.

Review

- 445 E. V. Tsipis, M. V. Patrakeev, P. V. Anikina, E. N. Naumovich and V. V. Kharton, Mater. Lett., 2020, 265, 127425.
- 446 A. A. Markov, O. V. Merkulov, M. V. Patrakeev and I. A. Leonidov, Int. J. Hydrogen Energy, 2019, 44, 26807-
- 447 C. Li, W. Li, J. J. Chew, S. Liu, X. Zhu and J. Sunarso, Sep. Purif. Technol., 2020, 235, 116224.
- 448 L. R. Tarutina, G. K. Vdovin, J. G. Lyagaeva and D. A. Medvedev, J. Alloys Compd., 2020, 831, 154895.
- 449 W. Deibert, M. E. Ivanova, S. Baumann, O. Guillon and W. A. Meulenberg, J. Membr. Sci., 2017, 543, 79-97.
- 450 M. Li, H. Niu, J. Druce, H. Téllez, T. Ishihara, J. A. Kilner, H. Gasparyan, M. J. Pitcher, W. Xu, J. F. Shin, L. M. Daniels, L. A. H. Jones, V. R. Dhanak, D. Hu, M. Zanella, J. B. Claridge and M. J. Rosseinsky, Adv. Mater., 2020, 32, 1905200.
- 451 K. Ishii, C. Matsunaga, K. Kobayashi, A. J. Stevenson, C. Tardivat and T. Uchikoshi, J. Eur. Ceram. Soc., 2019, 39, 5292-5297.
- 452 E. V. Shubnikova, M. P. Popov, S. F. Bychkov, S. A. Chizhik and A. P. Nemudry, Chem. Eng. J., 2019, 372, 251-259.
- 453 D. Ding, X. Li, S. Y. Lai, K. Gerdes and M. Liu, Energy Environ. Sci., 2014, 7, 552-575.
- 454 S. P. Jiang, Int. J. Hydrogen Energy, 2012, 37, 449-470.
- 455 C. Zhao, Y. Li, W. Zhang, Y. Zheng, X. Lou, B. Yu, J. Chen, Y. Chen, M. Liu and J. Wang, Energy Environ. Sci., 2020, 13, 53-85.
- 456 R. Zohourian, R. Merkle, G. Raimondi and J. Maier, Adv. Funct. Mater., 2018, 28, 1801241.
- 457 R. Zohourian, R. Merkle and J. Maier, Solid State Ionics, 2017, 299, 64-69.
- 458 H. T. Lozano, J. Druce, S. J. Cooper and J. A. Kilner, Sci. Technol. Adv. Mater., 2017, 18, 977-986.
- 459 H. Tang, Z. Gong, Y. Wu, Z. Jin and W. Liu, Int. J. Hydrogen Energy, 2018, 43, 19749-19756.
- 460 J. Dailly and M. Marrony, J. Power Sources, 2013, 240, 323-327.
- 461 J. Dailly, M. Marrony, G. Taillades, M. Taillades-Jacquin, A. Grimaud, F. Mauvy, E. Louradour and J. Salmi, J. Power Sources, 2014, 255, 302-307.
- 462 N. Nasani, D. Ramasamy, S. Mikhalev, A. V. Kovalevsky and D. P. Fagg, J. Power Sources, 2015, 278, 582-589.
- 463 E. Pikalova, A. Kolchugin, N. Bogdanovich, D. Medvedev, J. Lyagaeva, L. Vedmid, M. Ananyev, S. Plaksin and A. Farlenkov, Int. J. Hydrogen Energy, 2020, 45, 13612-13624.
- 464 A. Tarutin, A. Kasyanova, J. Lyagaeva, G. Vdovin and D. Medvedev, J. Energy Chem., 2020, 40, 65-74.
- 465 X. Li, D. Huan, N. Shi, Y. Yang, Y. Wan, C. Xia, R. Peng and Y. Lu, Int. J. Hydrogen Energy, 2020, 45, 17736-17744.
- 466 A. Tarutin, N. Danilov, J. Lyagaeva and D. Medvedev, Appl. Sci., 2020, 10, 2481.
- 467 J. Lyagaeva, N. Danilov, G. Vdovin, J. Bu, D. Medvedev, A. Demin and P. Tsiakaras, J. Mater. Chem. A, 2016, 4, 15390-15399.
- 468 S. Yang, Y. Lu, Q. Wang, C. Sun, X. Ye and Z. Wen, Int. J. Hydrogen Energy, 2018, 43, 20050-20058.
- 469 S. Yang, S. Zhang, C. Sun, X. Ye and Z. Wen, ACS Appl. Mater. Interfaces, 2018, 10, 42387-42396.

- 470 N. Danilov, A. Tarutin, J. Lyagaeva, G. Vdovin and D. Medvedev, J. Mater. Chem. A, 2018, 6, 16341-16346.
- 471 N. Danilov, J. Lyagaeva, G. Vdovin and D. Medvedev, Appl. Energy, 2019, 237, 924-934.
- 472 J. Li, J. Hou, Y. Lu, Q. Wang, X. Xi, Y. Fan, X.-Z. Fu and J.-L. Luo, J. Power Sources, 2020, 453, 227909.
- 473 Y. Bu, S. Joo, Y. Zhang, Y. Wang, D. Meng, X. Ge and G. Kim, J. Power Sources, 2020, 451, 227812.
- 474 S. Sun and Z. Cheng, J. Electrochem. Soc., 2020, 167, 024514.
- 475 D. Xie, K. Li, J. Yang, D. Yan, L. Jia, B. Chi, J. Pu and J. Li, Int. Hydrogen 2020, DOI: 10.1016/ Energy, j.ijhydene.2020.01.014.
- 476 K. Miyazaki, Y. Ding, H. Muroyama, T. Matsui and K. Eguchi, Electrochemistry, 2020, 88, 28-33.
- 477 C. Ren, Y. Zhang, O. Xu, T. Tian and F. Chen, Int. J. Hydrogen Energy, 2020, 45, 6926-6933.
- 478 M. Ahn, S. Hwang, S. Han, M. Choi, D. Byun and W. Lee, Korean J. Chem. Eng., 2020, 37, 1371-1378.
- 479 H. Tang, Z. Jin, Y. Wu, W. Liu and L. Bi, Electrochem. Commun., 2019, 100, 108-112.
- 480 A. He, J. Onishi, J. Gong and N. Shikazono, J. Power Sources, 2020, 478, 228771.
- 481 N. A. Baharuddin, N. F. A. Rahman, H. A. Rahman, M. R. Somalu, M. A. Azmi and J. Raharjo, Int. J. Energy Res., 2020, 44, 8269-8313.
- 482 J. H. Shim, G. D. Han, H. J. Choi, Y. Kim, S. Xu, J. An, Y. B. Kim, T. Graf, T. D. Schladt, T. M. Gür and F. B. Prinz, Int. J. Precis. Eng. Manuf.- Green Technol., 2019, 6,629-646.
- 483 E. G. Kalinina and E. Y. Pikalova, Russ. Chem. Rev., 2019, 88, 1179-1219.
- 484 P. A. Connor, X. Yue, C. D. Savaniu, R. Price, G. Triantafyllou, M. Cassidy, G. Kerherve, D. J. Payne, R. C. Maher, L. F. Cohen, R. I. Tomov, B. A. Glowacki, R. V. Kumar and J. T. S. Irvine, Adv. Energy Mater., 2018, 8, 1800120.
- 485 V. A. Tsvinkinberg, A. S. Tolkacheva, E. A. Filonova, O. I. Gyrdasova, S. M. Pikalov, V. A. Vorotnikov, A. I. Vylkov, N. I. Moskalenko and E. Y. Pikalova, J. Alloys Compd., 2020, 853, 156728.
- 486 V. A. Sadykov, E. M. Sadovskaya, E. A. Filonova, N. F. Eremeev, V. D. Belyaev, V. A. Tsvinkinberg and E. Y. Pikalova, Solid State Ionics, 2020, 357, 115432.
- 487 Q. Zhou, T. Zhang, C. Zhao, L. Qu, Y. He, T. Wei and X. Tong, Mater. Res. Bull., 2020, 131, 110986.
- 488 Y. Sadia, Y. Gelbstein and S. J. Skinner, J. Solid State Chem., 2020, 290, 121556.
- 489 E. A. Filonova, E. Yu Pikalova, T. Yu Maksimchuk, A. I. Vylkov, S. M. Pikalov and A. Maignan, Int. J. 2020, DOI: Hydrogen Energy, 10.1016/ j.ijhydene.2020.10.243.
- 490 M. Khoshkalam, D. Tripković, X. Tong, M. A. Faghihi-Sani, M. Chen and P. V. Hendriksen, J. Power Sources, 2020, 457, 228035.
- 491 T. Wan, A. Zhu, Y. Guo, C. Wang, S. Huang, H. Chen, G. Yang, W. Wang and Z. Shao, J. Power Sources, 2017, 348, 9-15.