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## Optical study of Te<sub>8</sub> ring clusters: comparison with density functional theory and a step towards materials design using nanoporous zeolite space

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The Te<sub>8</sub> ring molecule (cluster) is poorly investigated due to the lack of experimental data. Here, we report an experimental and theoretical study of a regular array of oriented Te<sub>8</sub> rings formed in the ~1.14 nm diameter cavities of zeolite LTA, which are arranged in a cubic lattice with a spacing of ~1.2 nm. Single crystals of LTA with encapsulated tellurium (LTA-Te) were studied using Raman spectroscopy (RS) and optical absorption spectroscopy (OAS). The experimental LTA-Te spectra were found to be in agreement with those calculated using density functional theory (PBE0 hybrid functional and def2-TZVP basis sets) for the crown-shaped Te<sub>8</sub> ring molecule with *D*<sub>4d</sub> symmetry. Using polarization–orientation RS, we show that the Te<sub>8</sub> rings are oriented by their major axes along the 4-fold axes of cubic LTA. We also show that the site symmetry of Te<sub>8</sub> in LTA-Te is lower than *D*<sub>4d</sub>. Te<sub>8</sub> bond-bending modes are well described in the harmonic approximation, while bond-stretching modes are mixed due to the reduced ring symmetry and, probably, anharmonicity. Importantly, OAS data of LTA-Te display dependence on the Te<sub>8</sub> concentration, implying the interaction of the rings from neighbouring LTA cavities with the generation of the valence and conduction electron bands of such a cluster crystal.

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## Introduction

Tellurium (Te) belongs to the same group of the periodic table as oxygen (O), sulphur (S) and selenium (Se). Since its discovery at the end of the 18<sup>th</sup> century, Te was widely used in metal alloys. Nowadays, Te and its compounds have become a real frontier of science and technology. They have found important applications in electronics, optoelectronics, photonics, phase-change materials, topological insulators, Li-Te batteries, thermoelectrics and solar cells.<sup>1–6</sup> However, elemental Te is mainly known in its chain form analogous to crystalline trigonal tellurium (t-Te) while it is nearly unknown in the ring form.

Zeolites provide a unique opportunity to form and accommodate uniform guest species in their cavities/channels where the species are oriented due to the crystalline nature of zeolites.<sup>7–13</sup> For example, only zeolites allowed obtaining polarized Raman spectra (RS) and optical absorption spectra (OAS) of isolated Te chains and rings.<sup>7–11</sup> Moreover, fabrication of such regular high-density Te species arrays can be considered as an important direction in the development of new functional materials, so-called cluster crystals.<sup>14–16</sup>

Here, we study Te<sub>8</sub> rings regularly arranged in the large cavities of zeolite LTA. Importantly, we are using the polarization–orientation RS examination of zeolite single crystals. On one hand, this method can be considered as an alternative to X-ray diffraction (XRD) for the structural study of zeolite-confined clusters. On the other hand, it provides very useful information about the cluster properties and their interaction with zeolites. We also study the OAS spectra of Te<sub>8</sub> and demonstrate the effect of the interaction between Te<sub>8</sub> rings in neighboring cavities, which is important for the fabrication of real cluster crystals.

The idea of Te<sub>8</sub> ring formation in LTA cavities was proposed in ref. 17 where the only observed RS feature at ~168 cm<sup>-1</sup> was assigned to the ring with an argument that its frequency is different from the dominant trigonal tellurium (t-Te) band frequency at ~120.5 cm<sup>-1</sup>.<sup>18,19</sup> However, this observation was not confirmed later. In ref. 20, the RS of LTA-Te microcrystalline powder was re-examined. A strong band at ~182 cm<sup>-1</sup> and a weaker one at ~45 cm<sup>-1</sup> were observed and attributed rather to Te<sub>12</sub> than to Te<sub>8</sub> since Te<sub>8</sub> requires a relatively strong symmetric bond-bending mode band at around 60 cm<sup>-1</sup>, which was not clearly detected due to laser-induced band broadening, as was found later. The LTA-Te band at ~62 cm<sup>-1</sup> accompanied by ~182 cm<sup>-1</sup> and ~45 cm<sup>-1</sup> bands was observed in ref. 21. This was a rather strong argument in favor of Te<sub>8</sub> ring formation in LTA cavities. The observed bands at ~45,

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$\sim 62$  and  $\sim 182$   $\text{cm}^{-1}$  look similar to the bending  $E_2$  and  $A_1$  mode bands and stretching  $A_1$  mode bands of  $S_8$  ( $\sim 153$ ,  $\sim 222$ ,  $\sim 480$   $\text{cm}^{-1}$ ) and  $Se_8$  ( $\sim 76$ ,  $111$ ,  $267$   $\text{cm}^{-1}$ ) rings.<sup>21</sup> The RS of another zeolite AFI with encapsulated tellurium (AFI-Te) showed similar bands at  $41$ – $47$   $\text{cm}^{-1}$ ,  $65$   $\text{cm}^{-1}$  and  $184$   $\text{cm}^{-1}$ .<sup>7</sup> The bands were attributed to the  $Te_8$  ring. However, no polarized RS study of LTA-Te single crystals has been performed yet, which could provide crucial evidence for  $Te_8$  ring formation. Here, we present this evidence. Importantly, no other species but  $Te_8$  was found in LTA-Te. This is beneficial for realizing the idea of cluster crystal formation.

## Experimental and theoretical methods

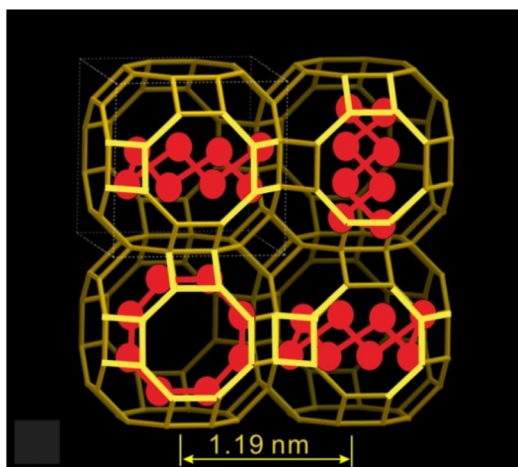
Synthetic zeolite LTA with the composition  $Na_{12}Al_{12}Si_{12}O_{48}$  was used in this work. The crystal growth procedure for LTA is described in ref. 22. The sizes of the cubic LTA crystals were  $\sim 20$   $\mu\text{m}$  along the edge of the cube. LTA has nearly spherical-shaped large cavities with a diameter of  $\sim 1.14$  nm, which are connected through narrow windows of  $\sim 0.42$  nm diameter (Fig. 1). The LTA cavities are arranged in a simple cubic lattice. The zeolite structures are accessible at the International Zeolite Association website [https://www.iza-structure.org/IZA-SC/ftc\\_table.php](https://www.iza-structure.org/IZA-SC/ftc_table.php).

Dehydration of zeolites in a vacuum was performed in Pyrex ampoules for several hours at  $t \sim 550$   $^\circ\text{C}$ . The design of the ampoules made it possible to add crystalline Te powder to the zeolite after dehydration without exposing the dehydrated samples to air. The subsequent adsorption of the Te vapour was carried out at  $t \sim 550$   $^\circ\text{C}$  for a few days until the brownish colour of the sample was saturated.

RS spectra of the LTA-Te single crystals were studied using a Renishaw micro-Raman spectrometer equipped with Semrock edge-filters and an Olympus micro-objective lens allowing an  $\sim 1$   $\mu\text{m}$  focused laser probe size. A 633 nm wavelength line of the He-Ne laser and a 785 nm line of the light-emitting diode

laser were used for RS excitation. Weak absorption of  $Te_8$  at these wavelengths is beneficial for  $Te_8$  compared to the non-absorbing LTA matrix for obtaining the cluster Raman signal with nearly no contribution of LTA. An additional measurement with 561 nm laser diode line excitation was performed using a Nanofinder-30 Raman/AFM system (Tokyo Instruments Inc.) specially designed for the maximal RS detection sensitivity around this wavelength, which allowed LTA-Te Raman measurement at relatively high  $Te_8$  absorption and very low laser excitation power in order not to strongly heat or destroy the clusters. Moreover, a Nanofinder-30 (Tokyo Instruments Inc.) was equipped with a set of three volume Bragg grating notch filters (OptiGrate Co), allowing the detection of the Raman signal with frequencies down to  $\sim 6$   $\text{cm}^{-1}$ . UV-visible OAS was performed using a Carl Zeiss micro-optical spectrometer with a light probe size of  $\sim 5$   $\mu\text{m}$ . The samples were intentionally broken to minimize their optical density in the ultraviolet (UV) spectral range and placed in glycerol between two cover glasses to minimize surface light scattering.

The  $Te_8$  RS, infrared (IR) spectra and OAS were simulated using density functional theory (DFT).<sup>23–26</sup> The structure and vibrational frequencies of  $Te_8$  were computed using the hybrid PBE0 exchange–correlation functional<sup>26</sup> and the resolution-of-identity approximation for the Coulomb part (RI-J).<sup>27</sup> The def2-TZVP basis sets<sup>28</sup> together with the corresponding auxiliary basis sets<sup>29</sup> and small-core effective core potentials (ECPs)<sup>30</sup> were employed for Te. RS intensities were obtained in the double harmonic approximation using analytical derivatives of static electronic polarizabilities<sup>31</sup> and analytical energy Hessians.<sup>32,33</sup> Calculations of static polarizability derivatives used the PBE0 functional and augmented def2-TZVPD basis sets.<sup>34</sup> IR spectra were computed in the double harmonic approximation from the analytical dipole and energy Hessians.<sup>32,33</sup> The vibrational frequencies were not scaled in this work. OAS was performed with time-dependent DFT (TDDFT)<sup>35</sup> using the PBE0 functional with the RI-J approximation<sup>36</sup> and def2-TZVPD basis sets. All calculations were performed with the Turbomole program package, version 7.5.<sup>23,37</sup>

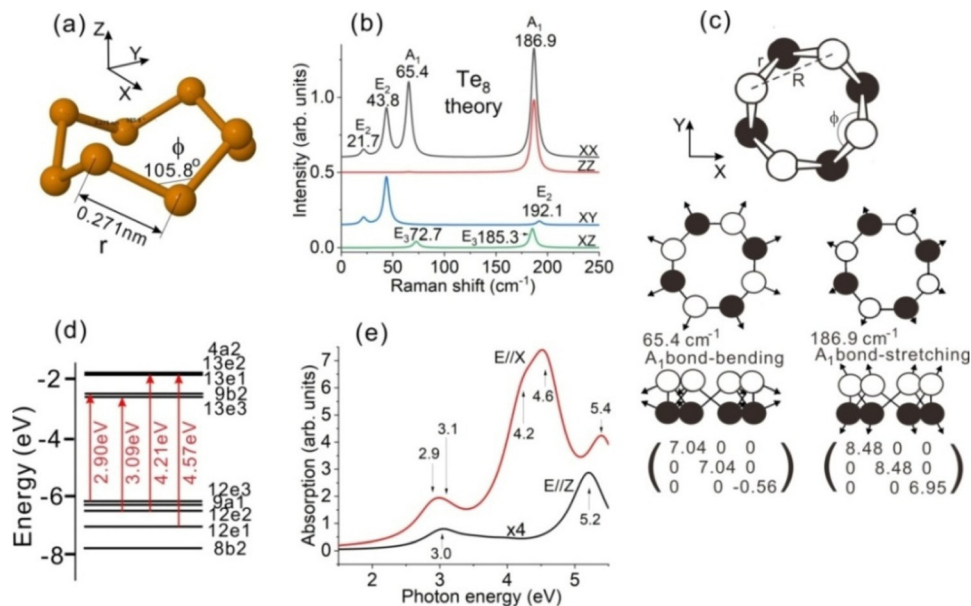


**Fig. 1** Schematic view of the LTA framework with  $Te_8$  rings in large cavities.

## Structure and Raman and absorption spectra of single $Te_8$ rings via DFT calculations

The calculated  $Te_8$  ( $D_{4d}$  point group symmetry) ring structure is shown in Fig. 2(a). The bond length of  $Te_8$  is  $r = 0.271$  nm and the second neighbour distance is  $R = 0.432$  nm, while its bond and dihedral angles are  $105.8^\circ$  and  $101.2^\circ$ , respectively. The obtained structure parameters are in agreement with previous DFT calculations.<sup>38–42</sup> The ring  $Z$  axis is directed along its 4-fold axis, while the  $X$  and  $Y$  axes lie in the molecule plane normal to the  $Z$  axis (Fig. 2(a)).

Theoretical point group analysis for the vibration modes of  $Te_8$  is similar to that made for the  $S_8$  molecule with a similar crown shape.<sup>43</sup> The  $Te_8$  vibration irreducible representations



**Fig. 2**  $\text{Te}_8$  ring with the bond length  $r$  and angle  $\phi$  (a); calculated Raman spectra of a single free-standing  $\text{Te}_8$  ring for  $XX$ ,  $ZZ$ ,  $XY$  and  $XZ$  polarization configurations (b);  $\text{Te}_8$  atomic displacements for  $A_1$  bond-bending and  $A_1$  bond-stretching modes with corresponding Raman tensors (c);  $\text{Te}_8$  MO diagram with dominant MO contributions to the strongest optical transitions shown in red, all for  $E//X$  (d); and calculated  $\text{Te}_8$  OAS including transitions with photon energies of  $<5.5$  eV (e).

(symmetries) with calculated frequencies and IR/Raman activities are shown in Table 1. The  $E_1$  symmetry corresponds to the IR activity for the light polarized in the ring plane ( $E//X$  or  $E//Y$ ), while the  $B_2$  symmetry corresponds to the  $E//Z$  IR activity. Calculated  $\text{Te}_8$  Raman intensities for different polarization configurations are equal to their squared Raman tensor components ( $XX$ ,  $ZZ$ ,  $XY$ ,  $XZ$ ) for each mode, which are determined by the corresponding atomic displacements and related molecular polarizability derivatives. The corresponding values in Table 1 are given in atomic units of  $10^{-2}e^2a_0/E_h$ , where  $e$ ,  $a_0$  and  $E_h$  are the electron charge, Bohr radius and Hartree energy, respectively.

The calculated RS spectra of  $\text{Te}_8$  for different polarization configurations, corresponding to the Raman tensor components from Table 1, are shown in Fig. 2(b). Raman bands are represented as Lorentz curves with the amplitudes corresponding to the calculated Raman activities and half-widths of  $5\text{ cm}^{-1}$  for all bands.

The dominant band in the calculated RS spectra of  $\text{Te}_8$ , displaying comparable  $XX$  and  $ZZ$  activities, corresponds to the  $A_1$  symmetric bond-stretching mode at a frequency of  $186.9\text{ cm}^{-1}$ . In contrast, the  $A_1$  bond-bending mode at  $65.4\text{ cm}^{-1}$  displays high  $XX$  activity and nearly negligible  $ZZ$  activity. Fig. 2(c) shows atomic displacements and Raman tensors of these two

**Table 1**  $\text{Te}_8$  normal vibrations: symmetries, computed frequencies, IR intensities and Raman activities

Symmetry	Frequency ( $\text{cm}^{-1}$ )	IR intensity ( $\text{km mol}^{-1}$ )	Raman activity $XX$	Raman activity $ZZ$	Raman activity $XY$	Raman activity $XZ$
$E_2$	21.7	0	-0.49	0	-2.08	0
	21.7	0	-2.08	0	0.49	0
$E_2$	43.8	0	-4.38	0	-3.36	0
	43.8	0	3.36	0	-4.38	0
$E_1$	56.9	0.524	0	0	0	0
	56.9	0.524	0	0	0	0
$A_1$	65.4	0	-7.04	0.56	0	0
$B_2$	68.6	0.128	0	0	0	0
$E_3$	72.7	0	0	0	0	-1.86
	72.7	0	0	0	0	-0.86
$B_1$	179.9	0	0	0	0	0
$E_3$	185.3	0	0	0	0	-3.54
	185.3	0	0	0	0	-0.17
$A_1$	186.9	0	-8.48	-6.95	0	0
$E_1$	188.1	1.29	0	0	0	0
	188.1	1.29	0	0	0	0
$E_2$	192.1	0	1.21	0	-1.05	0
	192.1	0	-1.05	0	-1.21	0

$A_1$  modes, which explain the difference in their  $ZZ$  activities. One can consider the  $A_1$  bond-stretching as a breathing mode since both in-plane and off-plane  $Te_8$  dimensions expand and contract the in-phase, while the  $A_1$  bond bending mode is an important instrument for  $Te_8$  ring orientation determination.

The  $E_2$  bond-bending mode at  $43.8\text{ cm}^{-1}$  is rather active in the  $XX$  and  $XY$  configurations, while inactive in the  $ZZ$  configuration. Similar but weaker activity is displayed by the  $E_2$  torsion and stretching modes at  $21.7\text{ cm}^{-1}$  and  $192.1\text{ cm}^{-1}$ , respectively. The  $E_3$  bond-bending mode at  $72.7\text{ cm}^{-1}$  and the bond-stretching mode at  $185.3\text{ cm}^{-1}$  dominate in the  $XZ$  configuration spectrum.

In addition to the  $Te_8$  Raman-active modes, there are several inactive modes:  $E_1$  bond-bending at  $56.9\text{ cm}^{-1}$ ,  $B_2$  bond-bending at  $68.6\text{ cm}^{-1}$ ,  $B_1$  bond stretching at  $179.9\text{ cm}^{-1}$  and  $E_1$  bond stretching at  $188.1\text{ cm}^{-1}$ , which can become active when the ring  $D_{4d}$  symmetry is reduced. As we mentioned above, the  $E_1$  and  $B_2$  modes are IR-active.

The calculated partial molecular orbital (MO) diagram of  $Te_8$  including the highest occupied MOs and the lowest unoccupied MOs is presented in Fig. 2(d). These MOs consist of combinations of  $Te\ 5p$  orbitals, as determined by Mulliken population analysis. The occupied orbitals correspond to  $Te$  lone pairs, while the unoccupied orbitals have antibonding character with respect to  $Te$ – $Te$  bonds. The largest MO contributions to the allowed electronic transitions with photon energies  $<4\text{ eV}$  and the largest oscillator strengths are shown in red. It should be noted that the electron transition energies take into account excitonic effects. Therefore, the transition energies are less than the orbital energy differences of the dominant MO contributions. All these four transitions have  $E_1$  symmetry, implying the polarization of light  $E//X$ .

There are also  $E//Z$ -allowed transitions of  $B_2$  symmetry. However, their oscillator strengths appeared to be much lower than those of the strongest  $E//X$  transitions (Table 2). The calculated OAS spectra of  $Te_8$  for polarizations of light  $E//X$  and  $E//Z$  are shown in Fig. 2(e). Allowed electronic transitions ( $E_1$

for  $E//X$  and  $B_2$  for  $E//Z$ ) with photon energies of up to  $5.5\text{ eV}$  were computed. Due to the difference in the oscillator strengths, the  $E//X$  absorption is stronger than the  $E//Z$  one. The  $E//X$  absorption is mainly represented by transitions at photon energies of  $\sim 2.9$ ,  $\sim 3.1$ ,  $\sim 4.2$ ,  $\sim 4.6$  and  $\sim 5.4\text{ eV}$ , while the  $E//Z$  one is determined by transitions at photon energies of  $\sim 3$ ,  $\sim 3.6$ ,  $\sim 4$  and  $\sim 5.2\text{ eV}$ . The corresponding absorption bands are approximated by the Lorentz curves with half-widths of  $\sim 0.6\text{ eV}$ .

## Experimental polarized Raman spectra of LTA-Te single crystals

Using  $Te$  vapor adsorption, we obtained LTA-Te samples with  $Te$  loading densities of  $\sim 3$  and  $\sim 8$  atoms per cavity (at. per cav.). Unlike the RS of LTA-Se with different loading densities showing changes in the  $Se_8/Se_{12}$  intensity ratio,<sup>44</sup> LTA-Te RS spectra did not show any qualitative dependence on the  $Te$  loading density except an increase in the Raman signal of the same bands with an increase in the  $Te$  density. This is a strong argument in favour of the formation of only one stable  $Te$  cluster in the LTA cavities.

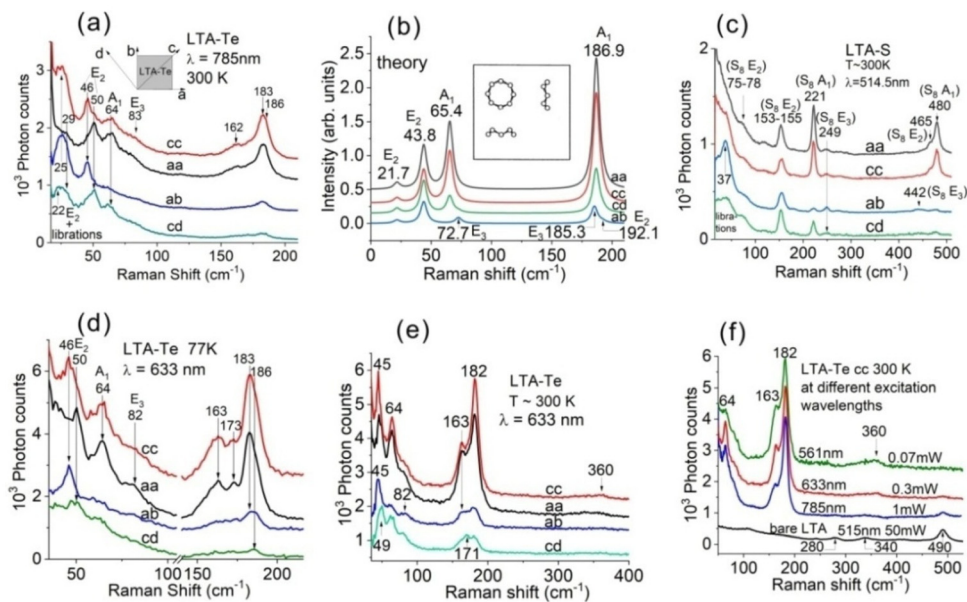
RS spectra of LTA-Te with a loading density of  $\sim 8$  at. per cav. at an excitation wavelength of  $785\text{ nm}$  are shown in Fig. 3(a). We performed measurements in four different polarization configurations:  $aa$ ,  $cc$ ,  $ab$  and  $cd$  (see the inset in Fig. 3(a)) similar to the polarization–orientation Raman study of LTA with sulphur LTA-S<sup>12</sup> and LTA with selenium LTA-Se.<sup>12,44,45</sup> Experimentally, we rotated the LTA-Te crystals in the  $ab$  plane with incident and scattered light polarizations (1) parallel and (2) perpendicular to each other. The procedure was described in detail earlier.<sup>12</sup> Theoretical RS spectra of LTA-Te for the  $aa$ ,  $cc$ ,  $ab$  and  $cd$  polarization configurations were obtained *via* summation of the Raman responses of three  $Te_8$  rings in their three possible orientations in LTA crystals [Fig. 3(b)].

The experimental LTA-Te RS spectra display rather good correspondence with the calculated ones. Indeed, the bands at  $162\text{ cm}^{-1}$  and  $183\text{ cm}^{-1}$  can be assigned to the  $Te_8\ E_3$  and  $A_1$  bond-stretching modes similar to the  $442\text{ cm}^{-1}$  and  $480\text{ cm}^{-1}$  bands of  $S_8$  in LTA-S [Fig. 3(c)]. The correspondence between the experiment and theory looks especially good for the  $A_1$  mode frequency (experimental:  $\sim 183\text{ cm}^{-1}$  vs. theoretical:  $186.9\text{ cm}^{-1}$ ) and intensity (the strongest band in the spectrum). At first glance, the polarization dependence of its intensity also looks reasonable: the band is strong in  $aa$  and  $cc$  polarization configurations, while it is weak in  $ab$  and  $cd$  configurations as expected for the  $A_1$  modes. Actually, as we have shown below, the situation is more complicated for the bond stretching modes due to a possible mode mixing effect but, in a rough approximation, the assignment of the  $183\text{ cm}^{-1}$  band to the  $A_1$  bond-stretching mode looks reasonable.

Polarization dependencies of the LTA-Te bands in the  $Te_8$  bond-bending mode region appear to be qualitatively similar to those of the bond-bending modes of  $S_8$ <sup>12</sup> [Fig. 3(c)] and

**Table 2** Computed dipole-allowed electron transitions of  $Te_8$

Symmetry	Photon energy (eV)	Oscillator strength
$E_1$	2.90	0.104
$B_2$	3.03	0.016
$E_1$	3.09	0.053
$B_2$	3.45	0.003
$E_1$	3.53	0.001
$B_2$	3.61	0.003
$E_1$	3.89	0.012
$B_2$	3.96	0.004
$E_1$	4.11	0.006
$E_1$	4.21	0.357
$E_1$	4.35	0.002
$E_1$	4.57	0.525
$B_2$	5.21	0.071
$E_1$	5.42	0.333
$B_2$	5.77	0.008
$E_1$	5.82	0.229
$B_2$	5.94	0.002



**Fig. 3** LTA-Te Raman spectra of *aa*, *cc*, *ab* and *cd* polarization configurations at room temperature and 785 nm wavelength excitation with the inset schematically showing LTA-Te cubic crystals and the *a*, *b*, *c* and *d* axes (a); summation of the calculated Raman spectra of three  $\text{Te}_8$  orientations in LTA (see the inset) for *aa*, *cc*, *ab* and *cd* polarization configurations (b); LTA-S Raman spectra of *aa*, *cc*, *ab* and *cd* polarization configurations (c); LTA-Te Raman spectra of *aa*, *cc*, *ab* and *cd* polarization configurations at  $T \sim 77$  K and 633 nm wavelength excitation (d); LTA-Te Raman spectra of *aa*, *cc*, *ab* and *cd* polarization configurations at room temperature and 633 nm wavelength excitation (e); and LTA-Te Raman spectra of the *cc* polarization configuration at different excitation wavelengths with the unpolarized spectrum of bare LTA (black curve) (f).

$\text{Se}_8$ <sup>12,44,45</sup> rings in LTA. Indeed, the 46–50  $\text{cm}^{-1}$  band polarization–orientation dependence is similar to that of the  $E_2$  band at 74–76  $\text{cm}^{-1}$  or 77–79  $\text{cm}^{-1}$  of  $\text{Se}_8$ <sup>45</sup> and the  $E_2$  band at 153–155  $\text{cm}^{-1}$  of  $S_8$  [Fig. 3(c)]. The calculated  $\text{Te}_8$   $E_2$  bond-bending mode frequency of 43.8  $\text{cm}^{-1}$  is in good agreement with the observed 46–50  $\text{cm}^{-1}$  band frequencies of LTA-Te. The splitting of the degenerate  $E_2$  mode indicates structural distortion of  $\text{Te}_8$  in LTA and therefore a reduced symmetry compared to the initial  $D_{4d}$  symmetry of the ring. For example, a slight reduction of the  $\text{Te}_8$  symmetry from  $D_{4d}$  to  $C_{4v}$  would split the  $E_2$  mode into  $B_1$  and  $B_2$  modes, active in *XX* and *XY* polarization configurations, respectively. This is exactly what we experimentally observed. Interestingly, the splitting of the  $\text{Te}_8$   $E_2$  mode in LTA is very strong,  $\sim 8.3\%$ , in contrast to the weaker splitting of the corresponding  $E_2$  modes of  $S_8$ ,  $\sim 1.3\%$ , and  $\text{Se}_8$ ,  $\sim 2.6\%$ .

The LTA-Te 64  $\text{cm}^{-1}$  band displays polarization dependence similar to the  $A_1$  symmetric bond-bending mode band of  $S_8$  at  $\sim 221$   $\text{cm}^{-1}$  [Fig. 3(c)] and the analogous band of  $\text{Se}_8$  at  $\sim 112$   $\text{cm}^{-1}$ .<sup>12,44,45</sup> The frequency of 64  $\text{cm}^{-1}$  is very close to the calculated frequency of 65.4  $\text{cm}^{-1}$  of the  $\text{Te}_8$   $A_1$  bond-bending mode [Fig. 2(b); 3(b)]. The experimental 64  $\text{cm}^{-1}$  band is rather strong in the *aa* and *cc* polarization configurations and it is still not weak in the *cd* configuration, while the band completely disappears in the *ab* configuration. This is a consequence of the 2-dimensional character of the Raman tensor of this mode with only two strong diagonal tensor components in contrast to the rather 3-dimensional tensor of the  $A_1$  bond-stretching mode [Fig. 2(c)]. Correspondingly, the single free-

standing  $\text{Te}_8$   $A_1$  bond-bending mode band is strong in the *XX* polarization configuration and negligibly weak in the *ZZ* configuration with the zero *XY* and *XZ* activities [Fig. 2(b)]. Theoretical calculation of this mode activity in the RS spectra of LTA-Te [Fig. 3(b)] is also in very good agreement with the experiment, namely high *aa* activity and no *ab* activity. This is a clear indication of the orientation of the  $\text{Te}_8$  4-fold axis along the 4-fold axis of LTA.

The experimental LTA-Te  $\sim 83$   $\text{cm}^{-1}$  band looks similar to the 248  $\text{cm}^{-1}$   $E_3$  bond-bending mode band of  $S_8$  [Fig. 3(c)]. It can be attributed to the  $E_3$  bond-bending mode of  $\text{Te}_8$  with a calculated frequency of 72.7  $\text{cm}^{-1}$ . The  $\sim 29$   $\text{cm}^{-1}$  band of LTA-Te can be attributed to the  $E_2$  torsional mode of  $\text{Te}_8$  with a calculated frequency of 21.7  $\text{cm}^{-1}$ . A broad  $\sim 25$   $\text{cm}^{-1}$  band, probably, originates from the superposition of the second component of the  $E_2$  torsional mode and the ring librations in the LTA cavity.

As we mentioned above, the bands of LTA-Te in the bond-stretching mode region require more detailed consideration. Fig. 3(d) shows LTA-Te RS in this region taken at temperature  $T \sim 77$  K and an excitation wavelength of 633 nm with better resolved bands due to their temperature-induced narrowing. Fig. 3(e) shows room temperature LTA-Te RS at the same excitation wavelength in a wider spectral range. Four bands at  $\sim 163$   $\text{cm}^{-1}$ ,  $\sim 173$   $\text{cm}^{-1}$ , and  $\sim 183$   $\text{cm}^{-1}$  and its shoulder at  $\sim 186$   $\text{cm}^{-1}$  can be observed in Fig. 3(d). The shoulder is, probably, associated with the  $\text{Te}_8$   $E_2$  bond-stretching mode (theoretical frequency of  $\sim 192$   $\text{cm}^{-1}$ ). Taking into account our preliminary assignment of the 182–183  $\text{cm}^{-1}$  and 162–163  $\text{cm}^{-1}$

bands to the  $\text{Te}_8$   $A_1$  and  $E_3$  modes, respectively, we have to attribute the 171–173  $\text{cm}^{-1}$  band to the forbidden  $B_1$  (theoretical frequency of  $\sim 180 \text{ cm}^{-1}$ ) mode which can be activated due to the reduced symmetry of  $\text{Te}_8$  in the LTA cavity, and we found that the bands in the bond-stretching mode region can be mixed.

The band at 182–183  $\text{cm}^{-1}$  shows slightly stronger activity in the  $cc$  polarization configuration than that in the  $aa$  configuration, which is clearly seen in the room temperature RS of LTA-Te obtained by excitation with 633 nm light, as shown in Fig. 3(e). This can be associated with the mode mixing due to the reduced symmetry of the  $\text{Te}_8$  rings. Interestingly, the 2<sup>nd</sup> order Raman band of  $\text{Te}_8$  can be recognized in the  $cc$  configuration spectrum at  $\sim 360 \text{ cm}^{-1}$  [Fig. 3(e)]. An even stronger 2<sup>nd</sup> order Raman band can be clearly seen at an excitation wavelength of 561 nm [Fig. 3(f), green curve] due to the resonant Raman enhancement. The band is very broad covering the spectral range from  $\sim 300 \text{ cm}^{-1}$  to  $\sim 375 \text{ cm}^{-1}$ . Importantly, LTA bands do not contribute to the 2<sup>nd</sup> order Raman band of  $\text{Te}_8$  since even the strongest LTA band at  $\sim 490 \text{ cm}^{-1}$  [Fig. 3(f), black curve] does not contribute to LTA-Te RS obtained by excitation with 561 nm light.

Summarizing this section, the polarization dependence of the  $A_1$  bond-bending mode Raman band of  $\text{Te}_8$  indicates the orientation of the ring by its 4-fold axis along the LTA 4-fold axis, while that of the  $E_2$  bond-bending mode suggests a noticeable symmetry reduction of the ring from the ideal  $D_{4d}$  point group symmetry.  $\text{Te}_8$  bond-stretching modes are mixed due to the ring symmetry reduction. The Raman bands at 162–163  $\text{cm}^{-1}$  and 171–173  $\text{cm}^{-1}$ , probably, originate from the  $E_3$  mode and formerly inactive  $B_1$  mode while the Raman bands at 182–183  $\text{cm}^{-1}$  and 186  $\text{cm}^{-1}$  originate from the  $A_1$  and  $E_2$  modes of  $\text{Te}_8$ . The originally inactive  $E_1$  bond-stretching mode may also contribute to the observed Raman bands.

## Optical absorption spectra of LTA-Te single crystals

Fig. 4(a and b) show the OAS of LTA-Te with the loading densities of  $\sim 3$  atoms per cavity [LTA-Te(3)] and  $\sim 8$  atoms per cavity [LTA-Te(8)]. Absorption bands at  $\sim 2.8$  eV, 3.2 eV, 4.3 eV and 4.6 eV are observed in the OAS of LTA-Te(3), while absorption bands at  $\sim 2.6$  eV, 3.2 eV and 4.4 eV are observed in the OAS of LTA-Te(8). The spectra indicate that the energy band gap decreases when the Te loading density increases. Moreover, LTA-Te(8) bands look slightly broader than those of LTA-Te(3). In particular, the absorption band doublet at  $\sim 4.3$  eV and 4.6 eV, resolved in LTA-Te(3), looks like a single band at  $\sim 4.4$  eV.

First, the observed experimental OAS spectra of LTA-Te [Fig. 4(a and b)] appeared to be in nearly perfect agreement with the theoretical OAS spectra of  $\text{Te}_8$  [Fig. 2(e)]. The agreement with theoretical OAS for a single  $\text{Te}_8$  ring is especially good for LTA-Te(3) with approximately two rings per five cavities. We have to note that surface light scattering, which

increases with the photon energy, contributes to the experimental spectra. In contrast, theoretical spectra do not include this effect and do not include high-energy transitions ( $>5.5$  eV in Fig. 2(e)). This may produce an impression that the experimental and theoretical spectra are different. However, actually, they are very similar.

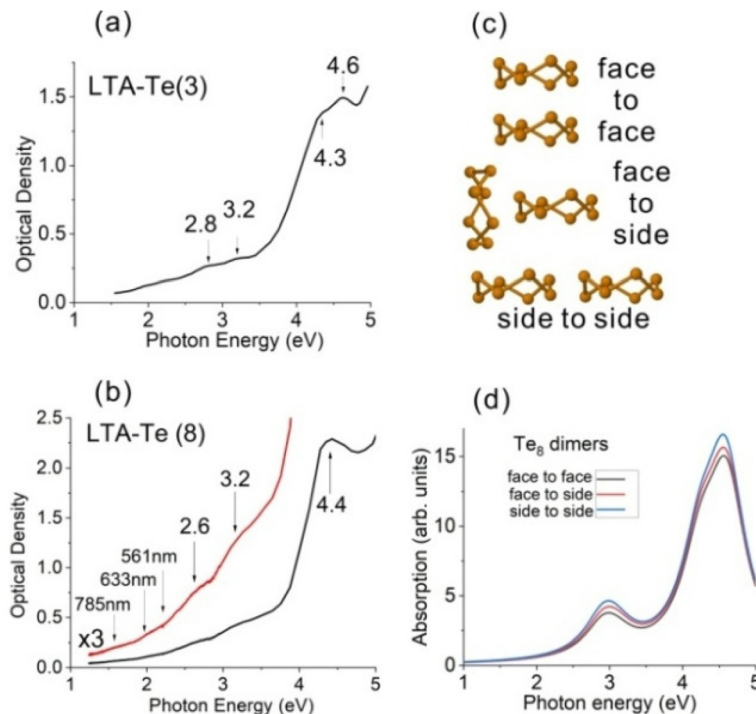
Secondly, differences between the OAS of LTA-Te(3) and LTA-Te(8) indicate a noticeable interaction between the  $\text{Te}_8$  rings in neighbouring LTA cavities, in particular, for LTA-Te(8) with one ring per cavity. In Fig. 4(b), we show the spectral positions of the Raman excitation laser lines, clearly indicating rather strong absorption of LTA-Te(8) at a wavelength of 561 nm corresponding to the resonant Raman effect [Fig. 3(f)]. To avoid laser-induced heating/destruction of  $\text{Te}_8$ , we used a very low 561 nm laser power of  $\sim 0.07$  mW. At weaker LTA-Te absorption corresponding to wavelengths of 785 nm and 633 nm [Fig. 4(b)], we were able to use higher laser powers of  $\sim 1$  mW and  $\sim 0.3$  mW, respectively [Fig. 3(f)].

In order to examine a possible effect of the ring interaction on the OAS spectrum of LTA-Te, we theoretically considered three possible  $\text{Te}_8$  ring dimers with a separation of  $\sim 1.2$  nm between the ring centres [Fig. 4(c)]. The corresponding calculated dimer OAS spectra are shown in Fig. 4(d). The effect of the ring interaction is indicated by the absorption enhancement with an increase in the ring interaction, which is the weakest in the “face-to-face” configuration and the strongest in the “side-to-side” configuration. Due to the inter-ring interaction, the unoccupied MOs of the  $\text{Te}_8$  rings organize conduction bands of the LTA-Te cluster crystals, while the highest occupied MOs organize its valence band. The conduction and valence bands should be rather narrow. Therefore, a variety of non-linear electric conductance effects predicted for narrow semiconductor superlattice mini-bands<sup>46</sup> can be expected in the LTA-Te electric properties. A negative differential electric resistance of LTA-Te observed in ref. 47 confirms this.

## $\text{Te}_8$ interaction with LTA and its librations vs. temperature

According to the calculated  $\text{Te}_8$  structure [Fig. 2(a)], the diameter of the ring, including the van der Waals radius of  $\sim 0.2$  nm of the Te atom, is  $\sim 1$  nm, which is slightly smaller than the LTA large cavity diameter. Therefore, one can expect that the  $\text{Te}_8$  ring is slightly shifted from the cavity center along the 4-fold LTA axis similar to the  $S_8$  and  $Se_8$  rings<sup>12,44,45</sup> but not as much as in the cases of  $S_8$  and  $Se_8$ . This should cause  $\text{Te}_8$  symmetry reduction from  $D_{4d}$  at least to  $C_{4v}$  [Fig. 5(a)]. We imposed the bond angle changes in order to simulate such structural distortion. The angles of type  $\phi_1$  were enlarged while those of type  $\phi_2$  were reduced. The second nearest neighbor distances  $R = 0.418$  nm and 0.446 nm in Fig. 5(a) correspond to  $\Delta\phi = \phi_1 - \phi_2 = 5^\circ$  ( $R = 0.432$  nm for undisturbed  $\text{Te}_8$ ).

The experimentally observed splitting of the  $E_2$  bond-bending mode [Fig. 3(a, d and e)] is, definitely, a sign of  $D_{4d}$



**Fig. 4** Experimental absorption spectra of LTA-Te(3) (a) and LTA-Te(8) (b); three  $\text{Te}_8$  dimers considered for the calculation of the effect of ring interaction on the absorption spectrum (c); and calculated absorption spectra of the considered  $\text{Te}_8$  dimers including transitions with photon energies of  $<5$  eV (d).

symmetry reduction. It is important to verify whether the observed  $E_2$  mode splitting fits the ring structure distortion shown in Fig. 5(a) or not. Due to the  $\text{Te}_8$  symmetry reduction from  $D_{4d}$  to  $C_{4v}$ , the  $E_2$  bond-bending mode splits into  $B_1$  and  $B_2$  modes, displaying  $XX$  and  $XY$  Raman activities, respectively. We performed calculations of the  $B_1$  and  $B_2$  mode frequencies and intensities as a function of the distortion  $\Delta\phi$ . The results are shown in Fig. 5(b). The  $B_1$  mode frequency increases while the  $B_2$  mode frequency decreases with  $\Delta\phi$ . This corresponds well to our experimental observation for the  $E_2$  bond-bending mode with the LTA-Te  $aa$ -active component frequency higher than that of the  $ab$ -active component. Thus, we conclude that the  $\text{Te}_8$  structural distortion shown in Fig. 5(a) reasonably describes the experimentally observed splitting of the  $E_2$  bond-bending mode of  $\text{Te}_8$ . However, even stronger symmetry reduction of  $\text{Te}_8$  is possible.

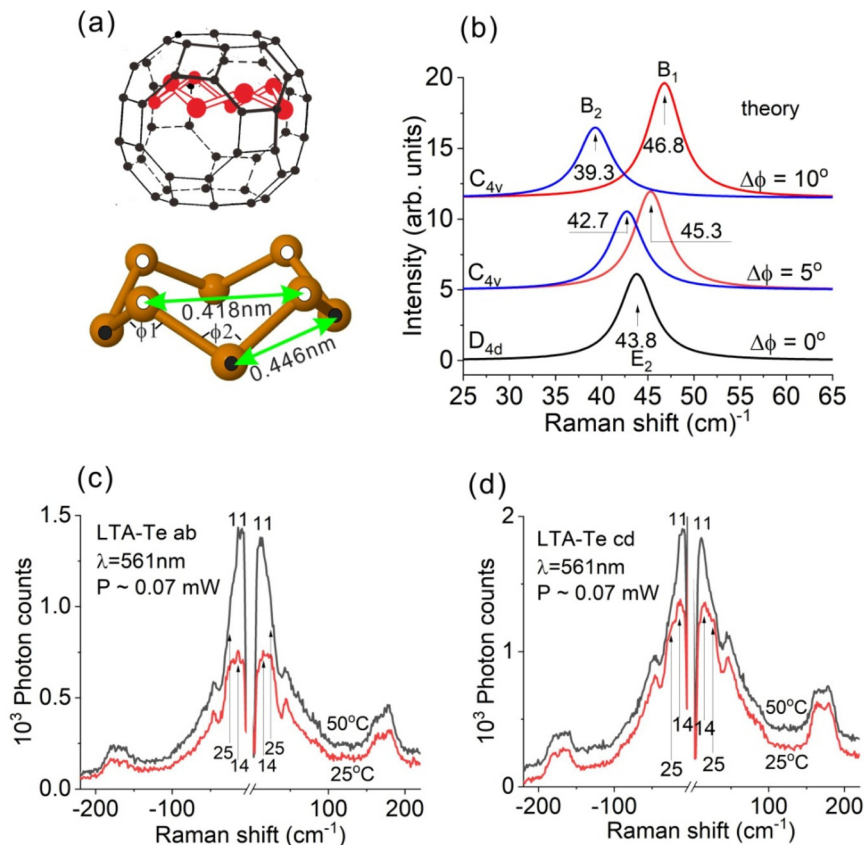
As we mentioned in the Introduction, Raman polarization-orientation spectroscopy in application to the zeolite-confined cluster structure study can be considered in some sense as an alternative to XRD. Indeed, using Raman polarization-orientation spectroscopy, we can determine a type of cluster (molecule) and its orientation in a zeolite. Then we can suggest its possible location in the zeolite when we know its size. However, we can only guess what sort of interaction occurs between the Te atoms and the zeolite. For understanding this point, the XRD data would be very useful. However, a structural analysis of LTA-Te by XRD is rather difficult because Te is distributed statistically inside the cavities. The resulting scatter-

ing power of Te sites with partial occupancy is comparable to those of  $\text{Na}^+$  cations and water molecules filling spaces not occupied by Te.

The  $\sim 20$  micron size of LTA-Te is insufficient for single-crystal XRD, while powder XRD of zeolites with clusters is, definitely, not effective for the examination of the cluster structure. It shows that the zeolite filling with clusters or molecules causes just changes in the relative intensities of the XRD peaks like, for example, in ref. 48 and its ESI, where a number of different types of zeolites, including LTA, with/without sulphur clusters were examined using powder XRD. In contrast, XRD analysis of  $\sim 65$  micron size LTA single crystals containing sulphur clearly showed double  $S_8$  ring cluster formation in each large cavity,<sup>49</sup> which is in agreement with the published RS results<sup>12,50,51</sup> and this work [Fig. 3(c)].

An XRD study of  $\sim 60$  micron size LTA-Te single crystals was performed in ref. 52. The XRD analysis suggested the interaction of Te atoms with the zeolite oxygen atoms and cations, while it showed low occupancy and high temperature factors for Te atomic sites. Therefore, the authors could not make any univocal conclusion about the Te cluster structure and they proposed several options including  $\text{Te}_8$  rings.

Using the RS of LTA-Te, the interaction between the  $\text{Te}_8$  ring and the zeolite can be probed *via* observation of the Te ring librations *vs.* temperature. Taking into account the  $S_8$  libration frequency of  $\sim 37$   $\text{cm}^{-1}$  [Fig. 3(c)] and assuming approximately the same van der Waals interaction force with LTA for both  $S_8$  and  $\text{Te}_8$ , we can roughly estimate the expected



**Fig. 5**  $\text{Te}_8$  ring shifted up from the centre of the LTA large cavity (top) with the  $C_{4v}$  structural distortion corresponding to the bond angle difference  $\Delta\phi = \phi_1 - \phi_2 = 5^\circ$  (bottom), the white and black spots show Te atoms in the top and bottom layers, respectively (a); calculated RS spectra of the  $\text{Te}_8$  ( $D_{4d}$ )  $E_2$  bond-bending mode (black) and  $\text{Te}_8$  ( $C_{4v}$ )  $B_1$  (red) and  $B_2$  (blue) bond-bending modes (b); and experimental low-frequency RS of LTA-Te at stage temperatures of  $\sim 25^\circ\text{C}$  and  $\sim 50^\circ\text{C}$  for *ab* (c) and *cd* (d) polarization configurations.

$\text{Te}_8$  room temperature libration frequency as  $\sim 18.5\text{ cm}^{-1}$  using the square root of the Te and S mass ratio, which is  $\sim 2$ .

In the LTA-Te RS spectra obtained by excitation with 785 nm light,  $\text{Te}_8$  librations, probably, contribute to the broad  $\sim 22\text{--}25\text{ cm}^{-1}$  band [Fig. 3(a)]. Instrumental limitations did not allow RS detection at frequencies  $< 20\text{ cm}^{-1}$  at this excitation wavelength. In contrast, such measurements were available at 561 nm light excitation. A weak point of this wavelength is that the corresponding  $\text{Te}_8$  absorption and light-induced heating are not negligible.

Fig. 5(c and d) show the Stokes and anti-Stokes low-frequency RS of LTA-Te in *ab* and *cd* polarization configurations at  $\sim 25^\circ\text{C}$  and  $\sim 50^\circ\text{C}$  temperatures of the heat-controlling table. We roughly estimated the laser-induced heating from the Stokes/anti-Stokes intensity ratio as  $\sim 20\text{--}30^\circ\text{C}$ , the value that should be added to the heating table temperature. In the RS taken at  $\sim 25^\circ\text{C}$  table temperature, one can see a  $\sim 25\text{ cm}^{-1}$  band in combination with a  $\sim 14\text{ cm}^{-1}$  band for both configurations. The last band, definitely, originates from the  $\text{Te}_8$  librations. With the increase of the table temperature to  $\sim 50^\circ\text{C}$ , it shifts to a lower frequency of  $\sim 11\text{ cm}^{-1}$  and displays significant enhancement. This corresponds to the intensification of librations and the reduction of the ring interaction with LTA.

The position of the  $\sim 25\text{ cm}^{-1}$  feature remains unchanged, confirming its origin from the internal torsional  $E_2$  mode of  $\text{Te}_8$ .

Thus, the  $\text{Te}_8$  libration amplitudes strongly increase with a moderate increase in the LTA temperature. This corresponds to a rather weak van der Waals interaction of  $\text{Te}_8$  with LTA. At room temperature, the libration amplitudes are, probably, quite strong as well. This explains high temperature factors observed for Te sites,<sup>52</sup> which complicates the XRD analysis of LTA-Te. In contrast, this effect does not cause any trouble for polarization-orientation Raman spectroscopy, which appears to be a very fruitful method of zeolite-confined cluster study providing information on the cluster structure with distortions, orientation and interaction with zeolites.

## Conclusions

In summary, the electronic and vibrational properties of  $\text{Te}_8$  rings were studied experimentally using Raman and absorption optical spectroscopy and theoretically using density functional theory. A good agreement is obtained between (1) the theoretical optical absorption and Raman spectra of isolated  $\text{Te}_8$  and (2) those of  $\text{Te}_8$  formed in the large cavities of the



zeolite LTA. Owing to zeolite utilization, this is a big step in the experimental study of  $\text{Te}_8$ , which is limited by the fabrication of exotic compounds like  $\text{Cs}_3\text{Te}_{22}$ <sup>53</sup> or cluster registration in molecular beams.<sup>54</sup> Polarization–orientation Raman spectra of LTA-Te show that the rings are oriented by their major axes along the 4-fold axes of LTA. The observed splitting of the  $\text{Te}_8$  bond-bending modes suggests the reduction of the ring  $D_{4d}$  symmetry in the LTA cavity. The bond-stretching modes demonstrate mixing due to the reduced symmetry of the ring and, probably, the essential anharmonicity of the vibrations. Even moderate heating of LTA-Te causes significant intensification of  $\text{Te}_8$  librations in the LTA cavities, which corresponds to a rather weak van der Waals interaction between the ring and the zeolite. Experimental optical absorption spectra of LTA-Te display dependence on the Te loading density due to the electronic interaction of  $\text{Te}_8$  from neighbouring cavities, which increases with an increase in the loading density. The effect is confirmed using DFT calculations of the absorption spectra of  $\text{Te}_8$  dimers. The observed  $\text{Te}_8$  ring–ring interaction is beneficial for the formation of cluster crystals with narrow electron bands and predicted non-linear electric properties. On the other hand, our results can be interesting in the sense of relationships with other inorganic 8-membered ring clusters such as  $\text{Pd}_8$ .<sup>55</sup>

## Conflicts of interest

There are no conflicts to declare.

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