

Cite this: *Chem. Sci.*, 2024, 15, 18364

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Regioisomers containing triarylboron-based motifs as multi-functional photoluminescent materials: from dual-mode delayed emission to pH-switchable room-temperature phosphorescence†

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Triarylboron compounds have been established as promising candidates in optoelectronic applications. However, realizing multi-functional properties in triaryl boron-based materials remains challenging. Herein, we present two regioisomers, **1** and **2**, designed judiciously by connecting a dimethylamino donor and a dimesitylboryl acceptor at 1,4 and 2,6-positions of the naphthalene spacer, respectively. Both compounds **1** and **2** display simultaneous, delayed fluorescence and persistent room-temperature phosphorescence (580 nm, τ_{av} = 168 ms, Φ = 76% for **1**; 550 nm, τ_{av} = 129 ms, Φ = 88% for **2** in 1 wt% PMMA), with the delayed fluorescence bands being sensitive to doping concentration (in PMMA). Notably, compound **1** in 1 wt% PMMA films demonstrates a reversibly switchable single-molecule phosphorescence from orange (580 nm) to green (λ_{Ph} = 550 nm, τ_{av} = 42 ms) in response to pH, which can be utilized for anti-counterfeiting applications. These results were further corroborated by studying the respective cationic salts **1**-OTF and **2**-OTF. Moreover, **1** and **2** exhibited blue-shifted fluorescence in response to mechanical pressure. Compound **2** also showed three-photon (σ 3P) absorption properties which were better compared to those of compound **1**.

Received 23rd August 2024
Accepted 7th October 2024

DOI: 10.1039/d4sc05656f

rsc.li/chemical-science

Introduction

Main group element-containing organic materials have received a lot of attention due to their promising photophysical properties^{1–10} and potential applications in organic light-emitting diodes (OLEDs),^{10–13} organic field effect transistors (OFETs),^{14–18} optical switches,^{19–22} organic solar cells,^{23–26} bio-imaging^{27–29} and sensors.^{30–34} In recent years, research on organic materials that show thermally activated delayed fluorescence (TADF) has been rapidly progressing as it offers 100% internal quantum efficiency which has led to the development of high-performance metal-free OLEDs.^{35–38} TADF occurs from

the excited singlet state (S_1) prior to the intersystem crossing (ISC) and reverse intersystem crossing (rISC) processes between excited singlet S_1 and triplet states (T_n , $n \geq 1$). Another delayed emission that involves triplet states is phosphorescence/afterglow in which radiative decay occurs from excited triplet (T_1) to the singlet ground state (S_0).^{39–41} Numerous ultralong room temperature phosphorescence (RTP) organic materials were reported by employing different strategies such as hydrogen bonding, incorporation of aromatic carbonyls, H-aggregation, doping in polymer matrices, external heavy atom effects, crystal packing, and incorporating heteroatoms with lone pairs of electrons.^{42–47} Such strategies were found to significantly suppress the molecular vibrations and enhance the phosphorescence efficiency from the triple state. Besides, heavy atoms increase the spin-orbit coupling and increase the ISC rate from S_1 to T_n . Among different strategies, doping chromophores in rigid polymer matrices has been extensively used because the polymer matrix can not only suppress/arrest the molecular vibrations significantly and disperse the chromophores but also reduce the triplet oxygen diffusion which is detrimental to the triplet state. RTP materials have found applications in sensors, anti-counterfeiting, information storage, and imaging.^{31,48–51}

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† Electronic supplementary information (ESI) available: General experimental procedure details and photophysical and theoretical data. CCDC 2288523 for **1**, 2288524 for **2** and 2291487 for **1**-OTF. For ESI and crystallographic data in CIF or other electronic format see DOI: <https://doi.org/10.1039/d4sc05656f>

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Among a wide variety of organic RTP materials, triarylboron (TAB) compounds offer intriguing bright fluorescence because of the presence of an empty p-orbital on boron which allows enhanced π -conjugation. In spite of this promise, studies on phosphorescence in triarylboron compounds have been relatively sparse.^{45,52–57} Recently, Marder and co-workers reported TAB compounds that showed efficient RTP resulting from an ISC from (σ , Bp) $\rightarrow$ (π , Bp) and C–H $\cdots$ C interactions in the crystalline state.⁵⁶ Xu *et al.* reported a series of molecules with different Lewis acidity and basicity that showed tunable ultra-long organic phosphorescence.⁵⁸ Zhao and co-workers reported efficient room-temperature phosphorescence (RTP) in the amorphous state from boron-containing biphenyl derivatives owing to intramolecular π – π and B $\cdots$ N electronic interactions and the presence of bromine atoms.⁴⁵ Thilagar *et al.* reported the impact of branching on the RTP and lifetime of TAB-based compounds.⁵⁹ Very recently, Ding *et al.* reported TAB containing donor (D)– π –acceptor (A) materials that showed time-dependent photo-activated room-temperature phosphorescence.⁶⁰

A subset of such interesting photoluminescent molecules exhibited dual delayed emissions such as TADF and RTP. For instance, Wang *et al.* reported a carbazole-type crystalline compound with persistent RTP and TADF emissions or time-dependent afterglow color.⁶¹ Yang *et al.* reported an indolocarbazole-doped polyvinyl alcohol (PVA) film which showed TADF and RTP emissions and then used it as temperature-sensitive security ink for multilevel anti-counterfeiting and information encryption.⁶² Recently, Thilagar and co-workers reported TADF and single molecular RTP in a BN-based molecular system.⁵⁴ Very recently, Wagner and co-workers reported laterally substituted 9,10-dimesityl-9,10-diboranthracenes that showed both TADF and RTP.⁶³ Although dual-mode delayed emission finds fundamental importance and useful applications like anti-counterfeiting and information encryption, however, achieving both TADF and single molecular RTP in a single emitter is a challenging task.

As part of our ongoing exploration of TAB-based RTP materials, we were interested in tuning the excited state properties by molecular engineering, *i.e.*, by decorating the boron acceptor and amino donor on the different positions of the naphthalene spacer. Accordingly, we have designed and synthesized two regioisomers, **1** and **2**, containing dimethylamino donor and dimesitylboryl acceptor at 1,4 and 2,6-positions of the naphthalene spacer, respectively. The boron–nitrogen system is well known for its delayed fluorescence behavior, while naphthalene derivatives are recognized for their inherent phosphorescence. By incorporating both systems into a single unit, dual emission can be achieved, utilizing the unique features of each. Additionally, the two distinct transition dipole moments in naphthalene provide directionality to the substituents, resulting in distinct electronic properties based on the substitution position. As anticipated, these regioisomers exhibited dual-mode delayed fluorescence and room-temperature phosphorescence (RTP) emission, along with multi-stimuli responsive optical properties (mechanical pressure, acid, and base). These properties were utilized to develop a rewritable ink (1 wt% PMMA)

for anti-counterfeiting applications. Furthermore, compound **2** exhibited superior three-photon non-linear optical properties compared to **1**. All these intriguing results are discussed in this manuscript.

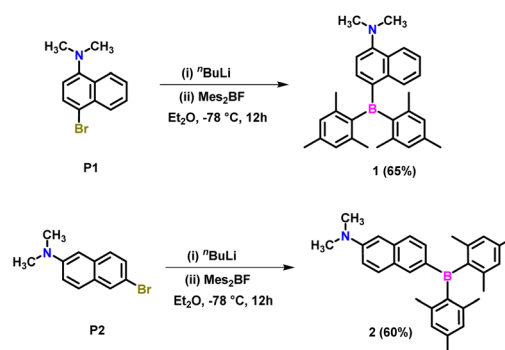
Results and discussion

Synthesis and characterization

The precursors 4-bromo-*N,N*-dimethylnaphthalen-1-amine (**P1**)⁶⁴ and 6-bromo-*N,N*-dimethylnaphthalen-2-amine (**P2**)⁶⁵ were synthesized following a modified literature protocol (ESI Scheme S1†). Lithiation of **P1** and **P2** using *n*-butyllithium in diethyl ether at -78°C , followed by trapping the C-centered carbanion with dimesitylfluoroborane, afforded target compounds **1** and **2** in 65% and 60% yields, respectively (Scheme 1). These compounds were characterized by ^1H , ^{13}C and ^{11}B NMR and high-resolution mass spectrometry (ESI Fig. S1–S8†). The ^1H NMR spectra of both compounds show two singlets corresponding to *ortho* (12H) and *para* (6H) methyl protons. This indicates that the two mesityl groups are in the same chemical environment and are not influenced by the position of the naphthylamine unit. In contrast, the substitution position significantly affects the ^1H resonances of the naphthalene spacer. The ^{11}B resonance in **1** is deshielded (70.2 ppm) compared to **2** (68.5 ppm). The variation of ^{11}B chemical shifts in **1** and **2** suggests that different substitution on the naphthyl amine motif has significant electronic influence and thus, these molecules were expected to show distinct optical features.

Single crystal X-ray diffraction studies

Single crystals of **1** and **2** suitable for X-ray diffraction studies were obtained from a slow evaporation of solutions of these compounds in dichloromethane and *n*-hexane solvent mixture. Compound **1** crystallizes in an orthorhombic lattice with a $P2_12_12_1$ space group, while **2** crystallizes in the triclinic space group $P\bar{1}$ (ESI Table S1†). The molecular structures of **1** and **2** are given in Fig. 1a and b. The boron center in **1** and **2** adopted a trigonal planar arrangement with sums of total bond angles being 360° , which is consistent with the results observed in 4-(dimesitylborenyl)-*N,N*-dimethylaniline (**R1**)⁶⁶ (ESI Table S2 and Fig. S8a†). The nitrogen center in **2** adopted a trigonal



Scheme 1 Synthesis of **1** and **2**.



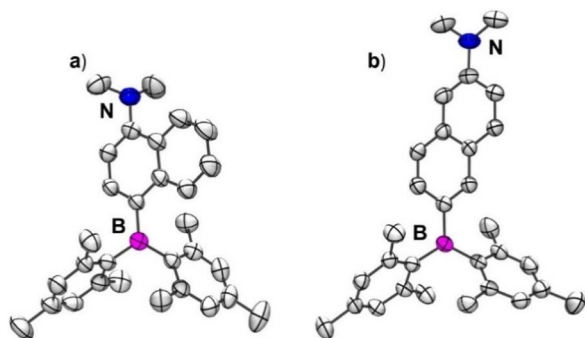


Fig. 1 ORTEP representation of the molecular structure of (a) **1** and (b) **2** with a 50% probability of thermal ellipsoids. Atom color codes: carbon (gray), nitrogen (blue), boron (magenta). All the hydrogen atoms are removed for clarity.

planar arrangement. In contrast, the nitrogen center in **1** adopted a pyramidal geometry with sums of total bond angles being $\sim 343^\circ$ which can be attributed to the steric repulsions between the methyl groups on N and adjacent hydrogen atoms of the naphthalene ring. The B–C(Nap) bond length is nearly the same in **1** and **2**; however, the N–C(Nap) bond length is longer in the case of **1** (1.409(5) Å) compared to **2** (1.371(2) Å). These bond length parameters suggest that the D–A interaction in **1** is expected to be weaker than in **2**. The dihedral angle (θ_{BN}) between the B3C and N3C planes in **2** (32.2°) is smaller than that in **1** (57.4°) (ESI Table S2 and Fig. S9–S11†). Similarly, the dihedral angle between NMe₂ and naphthyl planes in **2** (3.2°) is lower than that in **1** (61.1°). Compound **1** showed weak C–H $\cdots\pi$ interactions between the mesityl proton of one molecule and the mesityl- π cloud of the adjacent molecule (3.140 Å) (ESI Fig. S12†). In contrast, **2** showed dense packing due to its nearly planar arrangement of NMe₂ and the naphthalene ring, enabling two molecules to stack in an antiparallel fashion,

resulting in N $\cdots\pi$ interaction (3.440 Å) in addition to C–H $\cdots\pi$ interactions (3.069 Å and 3.210 Å) (ESI Fig. S13†). These interactions lead to a dense 3D packing in **2** compared to **1**, which can directly affect their photophysical properties in the solid state, as discussed in the following sections.

Photophysical and theoretical studies

Solution state absorption and emission studies

The absorption spectra of **1** and **2** in *n*-hexane revealed two absorption bands (Fig. 2a). Compound **1** showed absorption bands at 370 and 320 nm, while **2** exhibited bands at 380 and 297 nm, along with a shoulder at 323 nm. To understand the electronic nature of these bands, absorption spectra were recorded in solvents of different polarities (ESI Fig. S14†). The results obtained revealed that the high-energy bands were nearly insensitive to solvent polarity, while the broad low-energy bands showed bathochromic shifts in polar solvents. This indicated that the high-energy bands arise from local electronic transitions within the D/A or spacer, while the low-energy bands result from an intramolecular charge transfer (ICT) transition from the NMe₂ donor to the boryl acceptor, which is sensitive to the solvent polarity. The redshift (10 nm) of the low-energy band in **2** compared to **1** can be attributed to the extended conjugation and strong donor–acceptor interaction in the former compared to the latter.

Furthermore, the ground state geometry of **1** and **2** was optimized using density functional theory (DFT), and vertical transitions were simulated using time-dependent DFT (TD-DFT) calculations to shed more light on the nature of absorption bands (Fig. 2b, ESI Fig. S15–S17, Tables S4 and S5†). The ground state-optimized geometries closely resembled the molecular structures obtained by X-ray analysis (ESI Table S3†). The highest occupied molecular orbitals (HOMOs) of both **1** and **2** are localized on the NMe₂ donor and naphthalene spacer, while

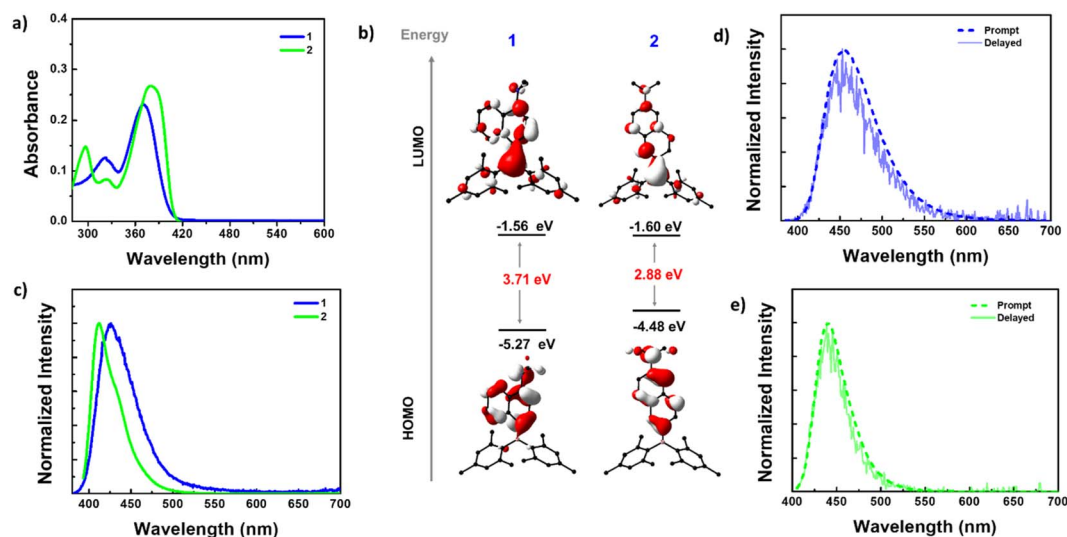


Fig. 2 (a) UV-visible absorption spectra for **1** and **2** in hexane (conc. 10^{-5} M). (b) Frontier molecular orbitals (iso value = 0.04) generated from the ground state [S_0] geometries using the B3LYP/6-31G(d,p) level of theory. (c) Steady-state photoluminescence (PL) spectra for **1** and **2** in hexane (conc. 10^{-5} M) at λ_{ex} = 380 nm. Prompt and delayed (20 μ s delay) PL spectra for (d) **1** and (e) **2** in degassed toluene (conc. 10^{-5} M) at λ_{ex} = 380 nm.

the lowest unoccupied molecular orbitals (LUMOs) are distributed on the entire molecular skeleton with a significant contribution from the boron center. The HOMO energy level of **1** (−5.27 eV) is more stabilized than that of **2** (−4.48 eV). The LUMO energy level of **2** (−1.60 eV) is slightly stabilized compared with that of **1** (−1.56 eV). Consequently, a lower band gap was observed for **2** (2.88 eV) than for **1** (3.71 eV). These bandgap trends were reflected in their experimental absorption spectra.

In *n*-hexane solution, **1** exhibited a single broad emission with a maximum at 423 nm (FWHM = 52 nm), while **2** displayed a sharper emission at 412 nm (FWHM = 37 nm) with a shoulder at 430 nm (Fig. 2c). Unlike the absorption spectra, the PL maxima of **2** are blue-shifted compared to **1**, which can be attributed to the larger structural reorganization of **1** in the excited state. The PL spectra of both the compounds showed a bathochromic shift in polar solvents indicating that the emission originates from a polar charge transfer state (ESI Fig. S18–S20†). From time-resolved PL decay, compound **2** showed a mono-exponential decay with a lifetime (τ) of 2.94 ns @ 412 nm while **1** displayed bi-exponential decay with an average lifetime (τ_{av}) of 0.73 ns @ 423 nm (ESI Table S6†). The PLQY (Φ) of these compounds in different solvents is in line with the intensity variation [Φ_{PL} = 6.90% for **1** and 71.60% for **2** in hexane, 33.25% for **1** and 86.74% for **2** in DMSO] (ESI Table S6†). The excitation spectra recorded for both compounds reproduced the absorption spectra in respective solvents, indicating that the same species are present in the ground and excited states (ESI Fig. S21†).

We have compared the absorption and emission features of the precursors (**P1** and **P2**) with those of compounds **1** and **2** to understand the changes in their absorption and emission features upon the covalent linking of boryl units to the precursors (ESI Fig. S22†). Interestingly, the emission and absorption characteristics of **P1** differ significantly from those of **1**, while **P2** has almost identical features to **2**, in dilute hexane solutions. For **P1**, the absorption spectrum shows a single broadband around 320 nm, similar to the higher energy band in compound **1** (ESI Fig. S22a†). These findings indicate that the higher energy absorption band corresponds to a local transition within the naphthylamine unit, while the lower energy band arises due to CT from the donor naphthylamine to the boryl acceptor. These differences are also evident in the emission spectra, where compound **1** shows a redshift (λ_{em} = 423 nm) compared to **P1** (λ_{em} = 400 nm). In contrast, the changes in **2** compared to **P2** are minimal. Both **P2** and compound **2** exhibit two sets of absorption bands, with the higher energy bands being identical. However, the lower energy band and emission band in **P2** are blue-shifted by 10 nm compared to **2** (ESI Fig. S22b†). These results are further supported by the computational results; the band gap found for **P1** (4.29 eV) and **1** (3.71 eV) differs largely, whereas **P2** (2.72 eV) and **2** (2.88 eV) have very minimal changes (ESI Fig. S22c†).

To validate the geometrical reorganization in the excited state, the PL spectra of these compounds were recorded in various viscous solvents, as viscosity is a key factor influencing molecular motion. PL studies were performed in *n*-hexane (non-

polar and non-viscous) and paraffin (non-polar and viscous), methanol (polar and non-viscous), and ethylene glycol (polar and viscous) (ESI Fig. S23†). It is evident from the spectra that there is a significant increase in PL intensity in paraffin over *n*-hexane in the case of **1**. However, this increment is very minimal in the case of compound **2**. Conversely, upon changing solvents from methanol to ethylene glycol, the emission intensity of compound **1** is reduced, while compound **2** shows a slight rise in intensity. These results indicate that considerable structural reorganization happens for **1** in the excited state in comparison to **2**.

Furthermore, the excited singlet state (S_1) geometry of **1** and **2** was optimized and the structural parameters were compared with the ground state structures (ESI Table S3†). The results suggest that the molecules undergo significant geometrical changes, which can be understood by monitoring the change in the dihedral angle between D–A/A–S/D–S (*S* = spacer) and bond parameters as tabulated in ESI Table S3.† Compound **1** exhibits a highly twisted geometry in the ground state compared to **2**. Upon excitation, both molecules undergo twisting around the B–C bond, creating a greater twist angle between the acceptor and the naphthalene–NMe₂ (S–D) group. These significant changes in compound **2** result in a decrease in the co-planarity that the molecule has in the ground state, leading to destabilization of the state and resulting in a slightly blue-shifted emission. These molecular motions are arrested in the viscous solvent, resulting in an enhancement of intensity in the case of paraffin compared to hexane. In hexane, the molecule can dissipate its excited energy non-radiatively due to intramolecular motion.

The D– π –A systems are well known for their delayed luminescence (DF) behavior, so we conducted delayed luminescence studies of **1** and **2** in the solution state. Time-gated spectra (delayed time = 20 μ s) and lifetimes were collected for these compounds in degassed toluene (Fig. 2d and e). As anticipated, the prompt fluorescence (PF) spectra of both compounds perfectly matched the delayed spectra, indicating the delayed fluorescence nature of these emissions. Furthermore, the lifetime of the DF is in the microsecond range. These longer lifetimes of the DF can be attributed to the involvement of the triplet state, as observed in many DF systems (ESI Fig. S24b, d, S25b and d†).

To confirm the involvement of the triplet excited state, the PL spectra were collected under different atmospheres (ESI Fig. S24 and S25†). The results indicate that both the prompt and delayed spectra exhibit a drastic decrease in intensity in the presence of oxygen, which is attributed to the quenching of the triplet state by oxygen. Furthermore, prompt lifetime and total photoluminescence quantum yield (PLQY) also reduced drastically under oxygen-rich conditions [$^N_2\Phi_{PL}$ = 47.99% for **1** and 93.48% for **2**, $^{O_2}\Phi_{PL}$ 24.04% for **1**, and 53.34% for **2**] (ESI Table S7†). These findings support the delayed fluorescence characteristics of these compounds.

To comprehensively understand the impact of the positions of D and A on the naphthalene spacer, we thoroughly investigated emission parameters such as lifetime, quantum yield, and various radiative and non-radiative decay rates for **1** and **2**



Table 1 Lifetime, ϕ and corresponding radiative and nonradiative decay rate for **1** and **2** in degassed toluene (conc. 10^{-5} M)^a

Compound	λ_{ex} [nm]	λ_{em} [nm]	τ_{PF} [ns]	τ_{DF} [μ s]	ϕ_{total} [%]	ϕ_{PF} [%]	ϕ_{DF} [%]	k_{PF} [10^7 s^{-1}]	k_{DF} [10^4 s^{-1}]	k_{ICS} [10^7 s^{-1}]	k_{RCS} [10^4 s^{-1}]	$k_{\text{nt}}(S_1)$ [10^6 s^{-1}]
1	375	455	4.57	5.23	47.99	24.04	23.95	5.26	4.58	10.94	38.09	56.82
2	395	440	3.52	5.49	93.48	53.34	40.14	15.05	7.28	12.27	18.21	10.89

^a τ_{PF} and τ_{DF} are the average fluorescence and delayed fluorescence lifetime. ϕ_{total} is total ϕ , ϕ_{PF} , and ϕ_{DF} for prompt and delayed fluorescence quantum yields, k_{PF} , k_{DF} , k_{ICS} , k_{RCS} , $k_{\text{nt}}(S_1)$ are the fluorescence decay, delayed fluorescence decay, intersystem crossing, reverse intersystem crossing, non-radiative decay rate constants respectively. λ_{ex} and λ_{em} are the excitation and emission wavelengths respectively.



Fig. 3 (a) Prompt and delayed (50 μ s delay) PL spectra for **1** and **2** in solid state (as prepared) [λ_{ex} = 380 nm], (b) time-gated PL spectra (50 μ s delay) and (c) delayed fluorescence lifetime decay [λ_{em} = 520 nm] for **1** (as prepared) [λ_{ex} = 380 nm] at 298 K under vacuum (absence of oxygen) and ambient conditions (presence of oxygen). (d) Fluorescence spectra and (e) time-gated spectra [50 μ s delay] for **1** under a vacuum atmosphere at 298 K and 77 K upon λ_{ex} = 380 nm. (f) Normalized prompt and delayed spectra [50 μ s delay] for **2** under a vacuum atmosphere at 298 K and 77 K upon λ_{ex} = 380 nm.

(Table 1). The average fluorescence lifetime for compound **2** [τ_{av} = 3.52 ns] is slightly lower compared to that for compound **1** [τ_{av} = 4.57 ns], while the delayed fluorescence lifetime is almost the same (Table 1). Interestingly, the photoluminescence quantum yield for compound **2** (93.48%) is significantly higher compared to that for compound **1** (47.99%). The higher ϕ for **2** is due to the smaller non-radiative decay rate constant of $10.89 \times 10^6 \text{ s}^{-1}$ than that of **1** ($56.82 \times 10^6 \text{ s}^{-1}$) (Table 1).

Solid-state emission studies

In contrast to the solution state, the solid-state photoluminescence spectra of **1** exhibit dual emission bands: a narrow emission band at 450 nm and a broad emission band at 520 nm, while compound **2** shows a broad emission band centered at 470 nm with a minor shoulder at 520 nm in the solid state (Fig. 3a, ESI Fig. S26 and S27[†]). The FWHM of the emission bands is found to be 38 nm for **1** and 50 nm for **2**. Compound **1** has the average lifetime τ_{av} = 3.36 ns at 450 nm and 16.23 ns at 520 nm while **2** has τ_{av} = 5.13 ns at 470 nm and 11.60 ns at 520 nm indicating that the emission originates from the singlet state (ESI Fig. S31, Tables S8 and S9[†]).

To resolve any ambiguity regarding the origin of the emission, we conducted delayed emission studies in the solid state. The delayed emission spectra of both compounds overlapped smoothly with the prompt (fluorescence) spectra and showed a lifetime in the microsecond range (Fig. 3a and ESI Fig. S30 and S32[†]). Additionally, to confirm the involvement of the triplet state, we performed both steady-state and time-gated measurements under oxygen-free (vacuum) conditions (Fig. 3b, c and S30–S33[†]). For compound **1**, the band at $\sim$ 520 nm experiences significant emission quenching in the presence of oxygen, whereas the emission band at 450 nm is minimally affected. A similar trend is observed in the delayed emission lifetime as well. However, we were unable to record the emission lifetime of the band centered at 450 nm under ambient conditions. The delayed lifetime found at the 450 nm

band [$\tau_{av} = 9.43 \mu\text{s}$] is considerably shorter compared to the lifetime found for the band at 520 nm [$\tau_{av} = 206.11 \mu\text{s}$] under vacuum. These results indicate that the two emission bands observed for **1** in the solid have different origins, with the lower energy band having a significant contribution from the triplet state. For compound **2**, the emission band centered at 470 nm shows considerable enhancement in intensity, and an additional red-shifted band arises at 610 nm under vacuum (in the absence of oxygen). Lifetime measurements conclude that the newly arisen emission bands have very high lifetimes in milliseconds [$\tau_{av} = 363.17 \text{ ms}$] and are attributed to phosphorescence originating from the triplet state. On the other hand, the 470 nm band showed a lifetime in microseconds [$5.93 \mu\text{s}$] under ambient conditions and increased more than thrice under vacuum (absence of oxygen) [$\tau_{av} = 17.60 \mu\text{s}$]. Thus, the band at 470 nm is ascribed to DF.

To gain further insight into the emission process, we investigated the emission spectra and lifetimes of these compounds in the solid state at different temperatures. Upon lowering the temperature to 77 K, the PL intensity of the high-energy band ($\sim 450 \text{ nm}$) in **1** increased with a strong vibrational progression, while the lower-energy broadband at 520 nm drastically reduced its emission intensity in the steady-state measurement, which was reflected in the lifetime as well [$\tau_{av} = 14.70 \text{ ns}$ at 77 K and 16.58 ns at 298 K] (Fig. 3d and ESI Fig. S34†). However, in the time-gated spectra, the lower-energy band at 520 nm completely vanished, indicating that the lower-energy band observed for compound **1** originated from a thermally assisted process (TADF) (Fig. 3e). Additionally, there is a weak redshifted band at 620 nm, and the lifetime corresponding to this band is in milliseconds [$\tau_{av} = 242.82 \text{ ms}$], proving that the redshifted bands are phosphorescent in nature (ESI Fig. S35, S36, Tables S8 and S9†). On the other hand, for compound **2**, both the prompt and delayed bands (delayed fluorescence band at 470 nm and phosphorescence band centered at 645 nm) increased their emission intensity with structured patterns, along with a sharp rise in their lifetime values (Fig. 3f, ESI Fig. S37–S39, Tables S8 and S9†).

In short, compound **1** exhibits dual emission behavior in the solid state, with emission bands observed at 450 nm and 520 nm. Among these two bands, the emission at 450 nm is fluorescence in nature, originating from the locally excited singlet state (S), which is less affected by the presence of oxygen and increases in intensity with a decrease in temperature. The second band at 520 nm is a thermally activated delayed fluorescence (TADF) band, originating from a stable singlet state (S^*) following efficient thermally assisted reverse intersystem crossing (rISC) from the triplet state. This band is highly sensitive to oxygen and is quenched at low temperatures, allowing the triplet excitons to decay radiatively, resulting in a phosphorescence band at 620 nm at 77 K. On the other hand, compound **2** exhibits a weak room temperature phosphorescence band at 610 nm from a low-energy triplet state with a lifetime in milliseconds under vacuum conditions, along with a delayed fluorescence band at 470 nm. Both these bands are found to increase their intensity upon decreasing in temperature. Hence, the delayed fluorescence emission is likely to occur

by utilizing isoenergetic high-lying triplet states (T_n) for rISC, which do not require temperature assistance. These interpretations are supported by comprehensive theoretical analyses, which will be discussed later in the manuscript.

Single-molecular persistent room temperature phosphorescence (pRTP)

The pristine samples of **1** and **2** exhibited no or very weak phosphorescence, which was attributed to crystallization-induced triplet quenching or overall concentration quenching. However, recent studies indicated that naphthalene derivatives could display efficient single molecular persistent room-temperature phosphorescence (RTP).⁵⁴ In this context, we conducted PL studies of these compounds doped in the PMMA matrix with varying concentrations (1, 10, and 50 wt%) of **1** and **2** (Fig. 4 and ESI S40–S43†). It was observed that both compounds displayed efficient persistent phosphorescence (pRTP) at room temperature in the absence of oxygen (vacuum), with lifetimes in the millisecond range. This phosphorescence was quenched as the doping concentration increased from 1 to 50 wt% and was absent in the neat film, indicating the single molecular origin of the phosphorescence emission. The prompt spectra for compound **1** (450 nm) in the 1 wt% thin film showed a single emission band, similar to the high-energy band observed in the solid state (450 nm) (see Fig. 4a). With increasing concentration, a slight redshift occurred, accompanied by broadening of the spectra, particularly at the tail end (Fig. 4a). These findings provided insight into the dual-band origin for compound **1** in the solid state, indicating that the low-energy band stemmed from a single molecule, exhibiting fluorescence, while the redshifted broadband is of intermolecular origin, demonstrating TADF nature as discussed *vide supra*.

The delayed spectra for doped matrices (1 and 10 wt% in PMMA) exhibited a redshifted structured band at 580 nm, along with a high energy band (450 nm) overlapping with the prompt fluorescence (Fig. 4b). Moreover, the 580 nm band is quenched

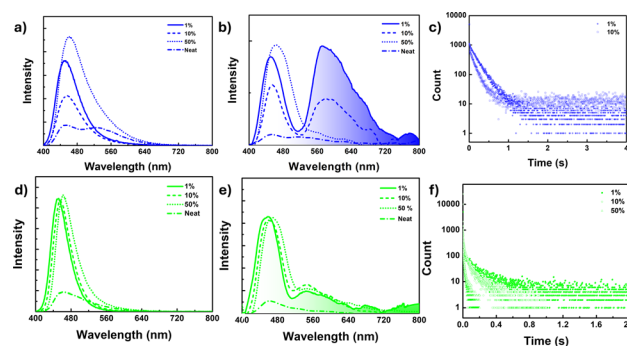


Fig. 4 (a and d) Fluorescence spectra, (b and e) delayed luminescence spectra [50 μs delay] of the neat film and doped PMMA film ($\lambda_{ex} = 380 \text{ nm}$) under vacuum at 298 K for **1** (a and b) and **2** (d and e) [spectra smoothed up to 50 pt for clarity, unsmoothed spectra in ESI Fig. S40c and 42c†], (c) phosphorescence lifetime decay for the **1** doped PMMA film (1 wt%, 10 wt%) at $\lambda_{ex} = 380 \text{ nm}$, $\lambda_{em} = 580 \text{ nm}$. (f) **2** doped PMMA film (1 wt%, 10 wt% and 50 wt%) at $\lambda_{ex} = 380 \text{ nm}$, $\lambda_{em} = 540 \text{ nm}$ under vacuum at 298 K.



at the 50 wt% doped matrix level and disappears in the neat film. A similar trend was observed in their decay profiles, with slightly higher lifetimes observed for **1** at the 1 wt% doping level [$\tau_{av} = 168.1$ ms for 1 wt% and $\tau_{av} = 133.1$ ms for 10 wt%] (Fig. 4c). Additionally, the excitation spectra corresponding to the emission band at 1 and 10 wt% doping levels reproduced their absorption spectra, further confirming their single molecular origin (ESI Fig. S40b†).

A similar phosphorescence pattern was observed for compound **2** as well; however, the observed PL band was weak and slightly blue-shifted ($\lambda_{em} = 550$ nm) compared to compound **1** (580 nm) (Fig. 4d–f). Furthermore, for compound **2**, phosphorescence bands were present even at 50 wt%, exhibiting marginal changes in the intensity with variations in doping concentration and disappearing in the neat film (Fig. 4e). A slight redshift in delayed fluorescence spectra was noted for **2** with increasing doping concentration. However, compound **2** did not fully restore its pristine spectra even at 50 wt%; consequently, the delayed fluorescence spectra in the neat film was slightly blue-shifted compared to 50 wt%. While there was only a marginal intensity difference observed for the phosphorescence band, a notable decrease in the lifetime decay was observed upon increasing the doping concentration ($\tau_{av} = 129.26$ ms at 1 wt%, 67.40 ms at 10 wt%, 8.72 ms at 50 wt%) (Fig. 4f and ESI Fig. S43†).

To gain an unambiguous understanding of the involvement of the triplet state, we conducted PL measurements for these films (both doped and neat) under air and vacuum at 298 K and low temperature (77 K) under vacuum for the 1 wt% doped film for both compounds (ESI Fig. S44–S56†). The phosphorescence band of **1** and **2** (1 wt%, 10 wt%, and 50 wt%) completely disappeared when exposed to air, with only the delayed fluorescence band persisting, albeit with a slightly reduced lifetime.

The oxygen sensitivity of these bands (delayed fluorescence and phosphorescence bands) indicated the involvement of the

triplet state in the emission process (ESI Tables S10–S13†). Furthermore, at low temperatures (77 K), the 1 wt% doped films of **1** and **2** showed stronger luminescence and slower decay than under ambient conditions. The increment in the phosphorescence intensity was almost double in the case of **1**, whereas for **2**, these changes were marginal (ESI Fig. S51c, S55c and Tables S10–S13†). The lowering of the temperature arrested molecular motion, which reduced the nonradiative decay channels from the excited state (both from singlet and triplet), resulting in the enhancement of both fluorescence and phosphorescence emission intensity as well as lifetime.

In summary, compound **1** exhibited a dual emission band in the solid state, with the higher energy band (~ 450 nm) being fluorescence in nature ($\tau_{av} = 3.36$ ns) and originating from a single molecular species, as evident in the 1 wt% doped PMMA. These bands remained consistent with the emission bands observed in toluene, allowing us to assign their origin to the singlet locally excited (1LE) state (see Fig. 5a–c). The lower energy band (~ 520 nm) was only present in the solid or neat film, with its appearance observed as a broad tailing of spectra upon increasing the compound's doping concentration in the PMMA matrix. This indicated an intermolecular origin due to close packing in the solid/neat film. Additionally, this band experienced significant intensity loss with decreased temperature or in the presence of oxygen, possessing a lifetime in microseconds ($\tau_{av} = 206.11$ μ s, and 15.99 μ s in the absence and presence of oxygen, respectively) at room temperature (ESI Table S9†). These findings suggested the emission band at ~ 520 nm is TADF, involving efficient triplet harvesting *via* rISC. Furthermore, these bands closely resembled the emission bands in polar solvents such as methanol and exhibited a significantly lower fluorescence lifetime than the higher energy band, collectively pointing to the CT nature of the state [1CT] (Fig. 5a–c). Moreover, compound **1**, upon lowering the temperature, resulted in the appearance of redshifted



Fig. 5 (a and d) Fluorescence spectra, (b and e) time gated spectra [50 μ s delay] [smoothened up to 50 pt for clarity, unsmoothed spectra in ESI Fig. S34, S37, S40, and S42†] for (a and c) **1** and (b and d) **2** in the solid state (as prepared solid) at 298 K and 77 K, as a neat film and a doped film in the PMMA matrix (1 wt%) under vacuum [$\lambda_{ex} = 380$ nm]. Plausible emission mechanism involved in (c) **1** and (f) **2** respectively.

phosphorescence, which disappeared with the concomitant rise of the TADF band at 520 nm, suggesting probable triplet-up conversion. This up-conversion was arrested at low temperatures, leading to phosphorescence emission from the triplet state ($\Delta E_{ST} = 0.39$ eV). Additionally, the weak and redshifted nature of this phosphorescence band (630 nm) indicated its origin from a CT triplet state of intermolecular origin (3CT). On the other hand, the stronger and more structured phosphorescence band observed at 540 nm in the doped film (1 and 10 wt%) originated mainly from a triplet state of single molecular origin (3LE) (Fig. 5c).

Compound 2 exhibited distinct emission characteristics, manifesting a broad single emission band centered at 470 nm in the solid state. Additionally, it displayed a weak RTP band at 580 nm, which became more prominent with vibronic progression at low temperatures (Fig. 5d–f). Furthermore, the time-gated PL spectra confirmed the presence of a DF band at 470 nm, showing greater intensity and lifetime at low temperatures, thereby ruling out TADF.^{63,67} This DF process likely occurred *via* vibronic coupling between the emissive singlet state and the higher energy triplet manifold with a small energy gap.^{63,67} Similar to doped thin films of compound 1, compound 2 also exhibited a phosphorescence band in the region of 560–700 nm, corresponding to the single molecular RTP peaks observed for reported naphthalene-based phosphors. However, the neat film of these compounds could not replicate their solid-state PL features or delayed fluorescence nature, attributed to the distinct tight molecular packing in the solid state compared to the thin film (Fig. 5f).

Excited state theoretical studies

To understand the nature of the excited states involved in DF/RTP processes, we conducted a detailed theoretical investigation including geometry optimization in an excited state (S_1 and T_1), spin-orbital coupling constants (ξ_{SOC}) between singlet and triplet manifolds, vertical transitions, and natural transition orbitals (NTOs) [S_1 , T_n ($n = 1-5$)] (ESI Fig. S57–S63 and Tables S14–S17†). For the S_1 state of compound 1, the hole is located on the naphthalene and NMe_2 moieties while the particle is located largely on boron with a minor contribution from the naphthalene spacer, indicating that the CT character with the very minor LE character dominates the S_1 state. In the case of T_1 , the hole and particle are concentrated on the naphthalene and NMe_2 moieties with additional contribution from B–C(Nap) indicative of the LE character (ESI Fig. S62 and S63†). Higher triplet states possess an LE character. A similar kind of NTO distribution is observed for 2 also, where the S_1 state illustrates its CT character, the T_1 state as LE, and the higher triplet state as LE.

There was considerable overlap between the highest occupied molecular orbital (HOMO) and lowest unoccupied molecular orbital (LUMO), resulting in a relatively high $\Delta E_{S_1-T_1}$ for 1 (0.87 eV) and 2 (0.83 eV), making efficient rISC realization impossible. Nonetheless, there were high-lying triplet states, such as T_2 or T_3 , which were energetically closer to S_1 , enabling efficient intersystem crossing. For compound 1, the results

revealed that S_1 to both T_2 and T_3 transitions were within the energy constraints to realize efficient spin crossover, and both were of LE nature. However, the spin-orbital coupling (SOC) constant values for the S_1 to T_2 transition were high (1.228 cm^{-1}), suggesting a higher priority for T_2 over T_3 . There is a large energy gap (0.74 eV) between the T_1 and T_2 states. Therefore, there is a competing pathway between DF and phosphorescence. Presumably, these energy level alignments could be the reason for the simultaneous realization of DF and RTP from a mono-molecular species (ESI Fig. S60, S62 and Table S15†). Compound 2 exhibited a similar scenario, with an inaccessible T_1 but high-lying T_2 and T_3 in energetically favorable positions, as seen in the ESI Fig. S59, S61 and Table S15.†

Mechanochromism

Mechanochromic materials have potential applications in optical storage, security papers, memory devices, anti-counterfeiting and sensors.^{31,48–51} Generally, color change or emission shift is due to the disruption of ordered molecular arrangement (molecular packing) upon mechanical force. The distinct crystal packing of compounds 1 and 2 prompted us to investigate their optical response toward mechanical force as it disrupts molecular arrangements. Interestingly, crystals of 1 and 2 showed blue-shifted emission, which is quite uncommon in mechanochromic materials. The greenish yellow color crystals of 1 showed dual emission with λ_{PL} of 460 ($\tau_{av} = 3.11$ ns) and 530 ($\tau_{av} = 19.60$ ns) nm, upon grinding, the low energy band disappeared while the high energy band was slightly blue-shifted (453 nm, $\tau_{av} = 3.09$ ns) which led to blue color of the ground form under UV light irradiation as seen in Fig. 6b. The lower energy band (530 nm) stems from the intermolecular origin (*vide infra*) in the solid state, and mechanical force disrupts such sort of molecular arrangement, leading to a single molecular emission (460 nm) as it is close to the 1 wt% doped PMMA film (450 nm). Attempts to check the reversibility from ground form to crystals either by heating or solvent fuming were unsuccessful. The ground form was converted back to crystals through recrystallization from slow evaporation of the DCM/hexane solvent mixture.

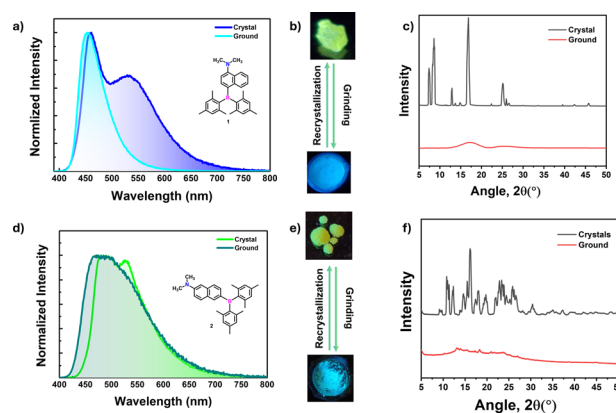


Fig. 6 Steady-state PL spectra for (a) 1 and (d) 2 in crystal and grounded form along with their (b and e) photographs before and after grinding. PXRD pattern for (c) 1 and (f) 2 before and after grinding.



Powder X-ray measurements were recorded to understand the mechanochromic behavior. From PXRD spectra, the crystals showed fine reflection peaks, indicating the well-ordered molecular arrangement, while the ground form showed diffuse spectra with no reflections, indicating the amorphous nature (Fig. 6c). The result further confirms that the emission color change is due to a phase transition from crystalline to amorphous. Akin to **1**, compound **2** also showed similar behavior. The yellowish green color crystals showed emission maxima at 490 nm ($\tau_{av} = 4.51$ ns) with λ_{PL} of 526 nm ($\tau_{av} = 8.83$ ns) shoulder band switches to a slightly blue-shifted broad band with an emission maximum at 480 nm ($\tau_{av} = 5.16$ ns) (ESI Fig. S64, S65 and Table S18†). Similar to compound **1**, the mechanochromic behavior of compound **2** is also irreversible. The PXRD pattern for crystals and ground forms of **2** confirms the phase transition from crystalline to amorphous like **1** (Fig. 6f).

Demonstration of afterglow and its response towards acid/base and anti-counterfeiting applications

After successfully demonstrating pRTP from the doped film, we have exploited the pRTP nature of **1** as a phosphorescence ink for data storage and other applications. The 1 wt% doped PMMA ink of **1** was used to write the letters “TIFR IPC” on a glass slide. These letters exhibited blue color fluorescence under UV light. Once the UV light was turned off, the letters showed persistent orange color room temperature phosphorescence ($\lambda_{ex} = 365$ nm) under vacuum (Fig. 7a). Furthermore, we have used 1 wt% and 50 wt% doped PMMA ink for anti-counterfeiting applications because the former exhibited bright afterglow which can be seen by the naked eye while the latter cannot be seen. For the encryption process, in digital number code format “8888” the alphanumeric secret code “293P” is encrypted using 1 wt% doped PMMA ink while the rest of the area in “8888” is written using 50 wt% doped PMMA ink. Under UV light, all digits “8888” showed bright blue color fluorescence (Fig. 7b). For the decryption, just turning off the UV light resulted in the appearance of an orange color alphanumeric secret code “293P”, showcasing its potential in anti-counterfeiting applications.

Numerous chromophores exhibited fluorescence on and off behavior or fluorescence emission change upon exposure to acid and base. However, the phosphorescence response towards

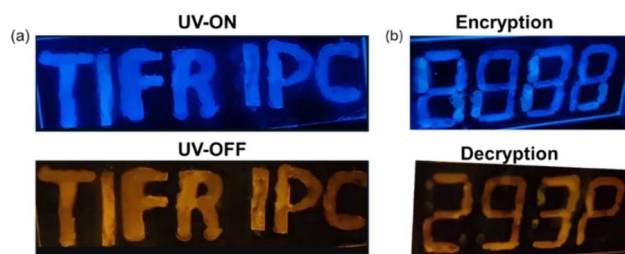


Fig. 7 Digital photographs of (a) writing done with ink made from a mixture of PMMA and **1** (1 wt%) under UV on and off conditions. (b) Anti-counterfeiting demonstration using ink made from a mixture of PMMA and **1** (1 wt% and 50 wt%) compound **1** utilizing single molecular pRTP emission features of **1**.

acid/base was underexplored. The acid-sensitive amino donors in our compounds inspired us to explore their fluorescence and phosphorescence response towards acid and base. The PMMA film (1 wt%) of **1** showed blue fluorescence ($\lambda_{PL} = 450$ nm, $\tau_{av} = 5.3$ ns) and orange color phosphorescence under vacuum ($\lambda_{Ph} = 580$ nm, $\tau = 168$ ms) (Fig. 8 and ESI Fig. S66 and Table S19†). When the films were exposed to trifluoroacetic acid vapors (TFA), it was found that the fluorescence emission ($\lambda_{PL} = 400$ nm, $\tau_{av} = 3.1$ ns) is blue-shifted with the reduced intensity, as seen in Fig. 8. This is due to the protonation of amino donors, which results in a decrease in intramolecular charge transfer. Interestingly, the protonated form showed persistent greenish-yellow RTP emission ($\lambda_{Ph} = 550$ nm, $\tau = 42$ ms) under the vacuum. Further exposure to ammonia vapor immediately brings it back to its original form. This pH-sensitive behavior is illustrated by writing the letters “11B” on the glass slide, as shown in Fig. 8e.

Compound **2** also exhibits a response to changes in pH similar to **1**. While the fluorescence response of compound **2** resembles that of compound **1**, the phosphorescence bands remain unchanged upon exposure to TFA. However, the phosphorescence lifetime values of compound **2** decreased significantly upon exposure to acid. In addition, neither the fluorescence emission nor the phosphorescence lifetime returns completely to their original states after treatment with a base (ESI Fig. S67 and Table S20†). Due to the weak phosphorescence emission of compound **2**, images under UV-off conditions could not be captured, and therefore, only images depicting fluorescence switching are included in Fig. S67.†

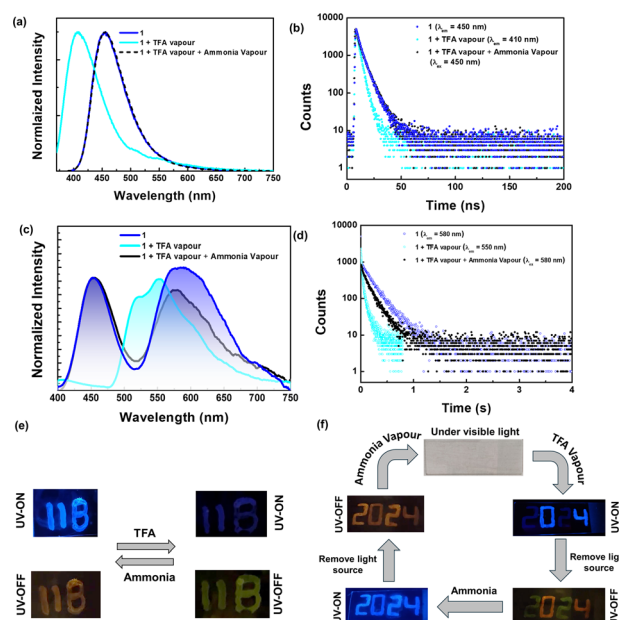


Fig. 8 (a) Fluorescence spectra and (c) phosphorescence spectra [smoothened up to 50 pt for clarity, unsmoothed spectra in ESI Fig. S66d†] (b and d) time-resolved decay for **1** doped in the PMMA film (1 wt%). Photographs illustrating (e) pH-sensitive emission features of the doped film (1 wt% in PMMA) and (f) the anti-counterfeiting application of the **1** film (1 wt% in PMMA) by exploring pH-sensitive emission features.



Additionally, we attempted emission switching using inorganic acid vapor, such as HCl.

We observed similar changes in emission, but the process took significantly longer (>2 h) compared to TFA vapor (<5 min) (ESI Fig. S68†). The pH-sensitive phosphorescence of **1** is used for anti-counterfeiting applications. The 1 wt% doped PMMA ink is used to write the number code “2024”. Under daylight, this code is not visible; however, it showed blue fluorescence under UV light ($\lambda_{\text{ex}} = 365$ nm). Then, the numbers 0 and 4 are masked with aluminum foil and the numbers 2 and 2 are exposed to acid vapour. Under UV light, only the numbers 0 and 4 showed intense fluorescence, while 2 and 2 appeared with faded color. Once the UV light was turned off, 0 and 4 showed orange color phosphorescence, and 2 and 2 showed green color room temperature phosphorescence under vacuum. Upon base treatment, all numbers showed orange color room temperature phosphorescence and blue fluorescence.

The different fluorescence and phosphorescence responses of this compound can be used for real-time anti-counterfeiting applications.

To strongly advocate for the switchable room temperature phosphorescence (RTP) triggered by acid vapor, we endeavored to produce protonated derivatives of compounds **1** and **2**. Unfortunately, our attempts were unsuccessful. Instead, we achieved the successful synthesis of methylated derivatives, referred to as compounds **1-OTF** and **2-OTF** henceforth. Compounds **1-OTF** and **2-OTF** were synthesized by methylating the NMe₂ donor for **1** and **2** using MeOTF in diethyl ether. Detailed synthesis schemes and other characterization data are included in the ESI (Scheme S2, Fig. S69–S75 and Table S21).† Prior to the comparison of photophysical studies, we compare the structural/geometrical parameters of **1-OTF** and **2-OTF** with those of **1** and **2**, respectively, using the data obtained either from crystal (in the case of **1-OTF** and **1**) or from ground state optimized geometry (in the case of **2-OTF** and **2**). It is clear from the crystal/ground optimized structure that the boron center is affected very minimally upon methylation on the N center, as we can see from the marginal geometrical changes around the B center [$\theta_{\text{S-A}} = 43.65^\circ$, 27.13° and B–C bond length 1.590 Å, 1.578 Å for **1-OTF** and **2-OTF** respectively]. However, considerable geometrical changes occur around the nitrogen center as reflected in the higher N–C bond length (1.526 Å, 1.514 Å for **1-OTF** and **2-OTF**) and dihedral angle related to donor unit [$\theta_{\text{S-D}} = 85.59^\circ$, 87.97° and $\theta_{\text{D-A}} = 77.34^\circ$, 89.11° for **1-OTF** and **2-OTF** respectively] (ESI Table S22†). These geometrical changes make the molecules (**1-OTF**, **2-OTF**) more twisted than those of their neutral derivatives (**1** and **2**) in *vide supra*.

Furthermore, these geometrical changes perturbed their electronic structure. As seen from the FMO diagram, the salts (**1-OTF** and **2-OTF**) have distinct FMOs compared to neutral derivatives (**1** and **2**), where the HOMO is solely concentrated on the OTF moiety and LUMO on the boryl aniline unit. The new orbital distribution changes D–A architecture, where the OTF unit acts as a donor and whole boryl aniline as an acceptor, resulting in a stronger charge transfer character in the salt compared to the neutral analogs. The same can be confirmed by the drastic increase in the ground state dipole moment in **1-OTF**

(14.9 D) and **2-OTF** (15.2 D) compared to **1** (1.7160 D) and **2** (4.1916 D) respectively. Upon comparing FMO's energy, both the HOMO and LUMO are stabilized; however, the extent of stabilization in the LUMO is more compared to the HOMO, though the overall stabilization occurs in **1-OTF** and **2-OTF** compared to **1** and **2**. In addition, the band gap is slightly reduced in cases **1-OTF** (3.58 eV) compared to **1** (3.71 eV), and the reverse trend is observed for **2** (2.88 eV) and its **2-OTF** (3.66 eV).

Following successful synthesis and geometrical analysis, we compared the emission properties of doped films of **1-OTF** and **2-OTF** (1 wt% and neat) with those of **1** and **2** (ESI Fig. S76–S85, Tables S23 and S24†). As expected, the 1 wt% doped film of compound **1-OTF** exhibits almost similar phosphorescence features ($\lambda_{\text{max}} = 560$ nm) to that of the TFA-exposed 1 wt% doped film of compound **1**. However, a significant increase in phosphorescence lifetime is observed for **1-OTF** ($\tau_{\text{av}} = 154.58$ ms) compared to **1** ($\tau_{\text{av}} = 41.98$ ms) (Fig. 9 and ESI Fig. S76b, S78 and S80†). Conversely, a slight redshift is observed for the fluorescence band of **2-OTF** ($\lambda_{\text{max}} = 440$ nm), attributed to the change in donor strength from amine (NMe₂) to ammonium salt (NMe₃⁺) that results from the reduction in band gap as mentioned above (Fig. 9 and ESI Fig. S81b, S83, S85, Tables S23 and S24†). In contrast, the photoluminescence (PL) features of neat films of **1-OTF** differ distinctly from those of **1**. The fluorescence spectra of **1-OTF** ($\lambda_{\text{max}} = 430$ nm) in the neat film are blue-shifted compared to those of the neat film of **1** ($\lambda_{\text{max}} = 450$ nm) (ESI Fig. S76a, d, S77 and S79†). Notably, phosphorescence bands in **1-OTF** are observed even in the neat film ($\lambda_{\text{max}} = 610$ nm) compared to **1**, with a slight redshift and lower lifetime ($\tau_{\text{av}} = 4.04$ ms) compared to the 1 wt% doped film of **1-OTF**. The appearance of phosphorescence bands in the neat film and the longer phosphorescence lifetime in compound **1-OTF** can be attributed to the efficient spin-flip process occurring within the cationic salt. This is due to its elevated spin–orbit coupling (SOC), a trait commonly found in emissive organic ionic species (ESI Table S25†). Similar changes were observed for **2-OTF** compared to **2** (ESI Fig. S81–S85 and Table S25†). Furthermore, to support the claim, we have compared the energy levels of the singlet and triplet manifolds involved in the emission process (ESI Fig. S86†). Both the singlet and triplet manifolds are greatly stabilized in the case of **1-OTF**; however, changes are minimal for **2-OTF** compared to **2**.

Non-linear optical properties

TAB-containing compounds with donor substituents showed prominent non-linear optical properties because of their intra-molecular CT characteristics.^{68–72} For example, Marder and co-workers have made a seminal contribution to the two-photon absorption properties of several TAB-based donor–acceptor compounds.^{66,73–77} The present two compounds (**1** and **2**) have different degrees of CT characteristics which inspired us to investigate their non-linear optical properties, in particular, three-photon absorption properties. Fig. 10a–d illustrate the open-aperture and closed-aperture Z-scan data of compounds **1** and **2**. Fig. 10a shows both two-photon absorption (2PA; green-





Fig. 9 (a) Molecular structure for compound 1-OTF, (b) ground state optimized geometry for compound 2-OTF and (c) FMO's for compounds 1-OTF and 2-OTF generated from geometry optimized structure (corresponding coordinates are taken from the crystal structure in the case of 1-OTF) using DFT at B3LYP and 631g(d,p) level of theory. (d) Prompt and delayed spectra under vacuum [absence of oxygen] for 1-OTF, 1, and 1(TFA) [1 after being exposed to TFA] in the PMMA matrix (1 wt%) [$\lambda_{\text{ex}} = 380$ nm], (e) 2-OTF, 2, and 2(TFA) [1 after being exposed to TFA] in the PMMA matrix (1 wt%) [$\lambda_{\text{ex}} = 380$ nm].

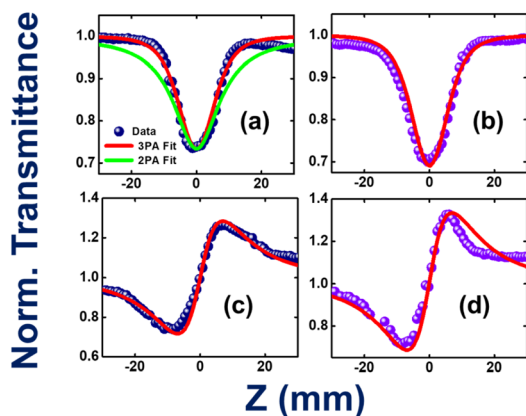


Fig. 10 (a) Open aperture Z-scan data of compound 1. Both two-photon absorption (green-colored line) and three-photon absorption (red-colored line) fits are included in the graph. (b) Open aperture Z-scan data and fits of compound 2. (c) Closed aperture Z-scan data of compound 1 (d) closed aperture Z-scan data of compound 1. Solid symbols represent the experimental data, while the solid lines represent the theoretical fits to the data. 800 nm, 1 kHz, ~ 70 femtosecond pulses were used to obtain this data.

colored line) and three-photon absorption (3PA; red-colored line) fits. The 3PA fits are a perfect match to the experimental data. The magnitudes of the extracted 3PA coefficients were $3.0 \times 10^{-23} \text{ cm}^3 \text{ W}^{-2}$ and $3.9 \times 10^{-23} \text{ cm}^3 \text{ W}^{-2}$, respectively (ESI Table S26[†]). Complete details of the calculations of various NLO coefficients have been presented in our earlier works.^{78,79} The corresponding 3PA cross-sections were calculated to be $\sim 10^{-78} \text{ cm}^6 \text{ s}^2$. Compound 2 has depicted a decisive enhancement of

$\sim 33\%$ when compared to the magnitude of compound 1. This can be attributed to both molecules' ground state dipole moments of 1.72 D for compound 1 and 4.19 D for compound 2 obtained from theoretical calculations (ESI Table S14[†]).

Zhu *et al.* have investigated the structure–property relationships of stilbene-based dipolar/quadrupolar chromophores using essential-state models for 3PA.⁸⁰ They claim that for dipolar systems, they observed an important role played by the change in dipole moment (D_m). This dipole moment change occurred between the ground and 3PA active states. They predicted 3PA cross-sections of 10^{-78} to $10^{-80} \text{ cm}^6 \text{ s}^2$ in their case. Fig. 10c and d depict the closed aperture data of compounds 1 and 2, respectively.

The magnitude of NLO coefficients depends inherently on the input pulse duration, excitation wavelength, and the availability of the high-lying excited states. In the present case, both the compounds have high-lying states conducive for 3PA [excitation at 800 nm (2.55 eV) to 3PA near 266 nm (4.66 eV)]. The magnitudes of the nonlinear refraction (n_2) coefficients were extracted to be $4.9 \times 10^{-16} \text{ cm}^2 \text{ W}^{-1}$ and $5.6 \times 10^{-16} \text{ cm}^2 \text{ W}^{-1}$, respectively. The 3PA coefficients of the present compounds are superior to those of some recently reported novel organic compounds, such as porphyrin–phthalimide-based donor–acceptor systems,⁸¹ and comparable to structurally unconstrained green fluorescence protein (GFP) chromophores.⁸² The solvent contribution (dichloromethane, DCM) was also investigated in similar experimental conditions, and the data is presented in ESI Fig. S87.[†] The magnitude of solvent NLO coefficients was much smaller (more than one order of magnitude lower than the solute) and was, therefore, ignored. The errors in the above measurements were estimated to be $<10\%$

India for the research fellowship (09/079(2806)/2019-EMR-1). Prof. PT is thankful to SERB (EEQ/2023/000903) and IISc for the financial support.

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The data that support the findings of this study are available in the ESI† of this article. Details on the synthesis and structural characterization, photo-physical, results from DFT and TD-DFT and RTP crystallographic data (2288523 for **1**, 2288524 for **2**, and 2291487 for **1-OTf**) are provided.

The project was designed by VC and PT. RA and MMAT contributed equally. Synthesis and characterization by RA and RK. Photophysical characterization and mechanochromism studies by RA, MMAT, and PS, Nonlinear Optical properties by DB and VR Soma. All authors contributed to the writing of the manuscript.

There are no conflicts to declare.

Prof. V. C. is thankful to the Department of Science and Technology, New Delhi, India, for a National J. C. Bose Fellowship. R. A., R. K., and P. S. are thankful to the TIFR Hyderabad from the Department of Atomic Energy (DAE), India. MMAT thanks CSIR

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