

Cite this: *Mater. Horiz.*, 2022,
9, 500Received 22nd August 2021,
Accepted 19th November 2021

DOI: 10.1039/d1mh01357b

rsc.li/materials-horizons

Backbone-driven host–dopant miscibility modulates molecular doping in NDI conjugated polymers†

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Molecular doping is the key to enabling organic electronic devices, however, the design strategies to maximize doping efficiency demands further clarity and comprehension. Previous reports focus on the effect of the side chains, but the role of the backbone is still not well understood. In this study, we synthesize a series of NDI-based copolymers with bithiophene, vinylene, and acetylenic moieties (P1G, P2G, and P3G, respectively), all containing branched triethylene glycol side chains. Using computational and experimental methods, we explore the impact of the conjugated backbone using three key parameters for doping in organic semiconductors: energy levels, microstructure, and miscibility. Our experimental results show that P1G undergoes the most efficient n-type doping owed primarily to its higher dipole moment, and better host–dopant miscibility with N-DMBI. In contrast, P2G and P3G possess more planar backbones than P1G, but the lack of long-range order, and poor host–dopant miscibility limit their doping efficiency. Our data suggest that backbone planarity alone is not enough to maximize the electrical conductivity (σ) of n-type doped organic semiconductors, and that backbone polarity also plays an important role in enhancing σ via host–dopant miscibility. Finally, the thermoelectric properties of doped P1G exhibit a power factor of $0.077 \mu\text{W m}^{-1} \text{K}^{-2}$, and ultra-low in-plane thermal conductivity of $0.13 \text{ W m}^{-1} \text{K}^{-1}$ at 5 mol% of N-DMBI, which is among the lowest thermal conductivity values reported for n-type doped conjugated polymers.

Introduction

Organic semiconductors (OSCs) hold the potential to enable a whole new generation of electronics such as flexible and

New concepts

In recent studies, modulation of host–dopant miscibility led to enhanced molecular doping of organic semiconductors (OSCs), thus enabling their applicability several organic electronic devices such as organic thermoelectrics. The most common approach includes the addition of oligo-ethylene glycol side chains to the OSC to improve their miscibility with polar dopants such as N-DMBI. Herein we focus on the conjugated polymer (CP) backbone and its impact on the molecular doping efficiency of naphthalene diimide (NDI) based CPs, with various degrees of backbone planarity, rigidity, and dipole moments. Our data show a strong correlation between the backbone dipole moment and the maximum value of the electrical conductivity achieved upon molecular doping with N-DMBI. The most polar backbone in the series (P1G) showed the highest electrical conductivity at the lowest doping concentrations. On the contrary, P3G possessing the least polar backbone showed the lowest electrical conductivity, requiring larger dopant amounts to reach its maximum conductivity. This suggests a backbone-driven miscibility effect that can enable or limit molecular doping efficiency. Consequently, we measured the temperature-dependent thermoelectric properties of P1G, finding an ultra-low in-plane thermal conductivity value of $0.13 \text{ W m}^{-1} \text{K}^{-1}$ enabled by the branched side chains. This opens a new possible strategy to further decrease the thermal conductivity of CPs.

wearable devices.¹ Unlike their inorganic counterparts, OSCs can be solution-processed and possess inherent flexibility.^{2,3} However, they are still limited by their low mobility. Molecular doping has become a major enabler of several technologies such as transport layers in organic photovoltaics and organic light-emitting diodes, organic photodiodes, organic field-effect transistors, and most recently in organic electrochemical

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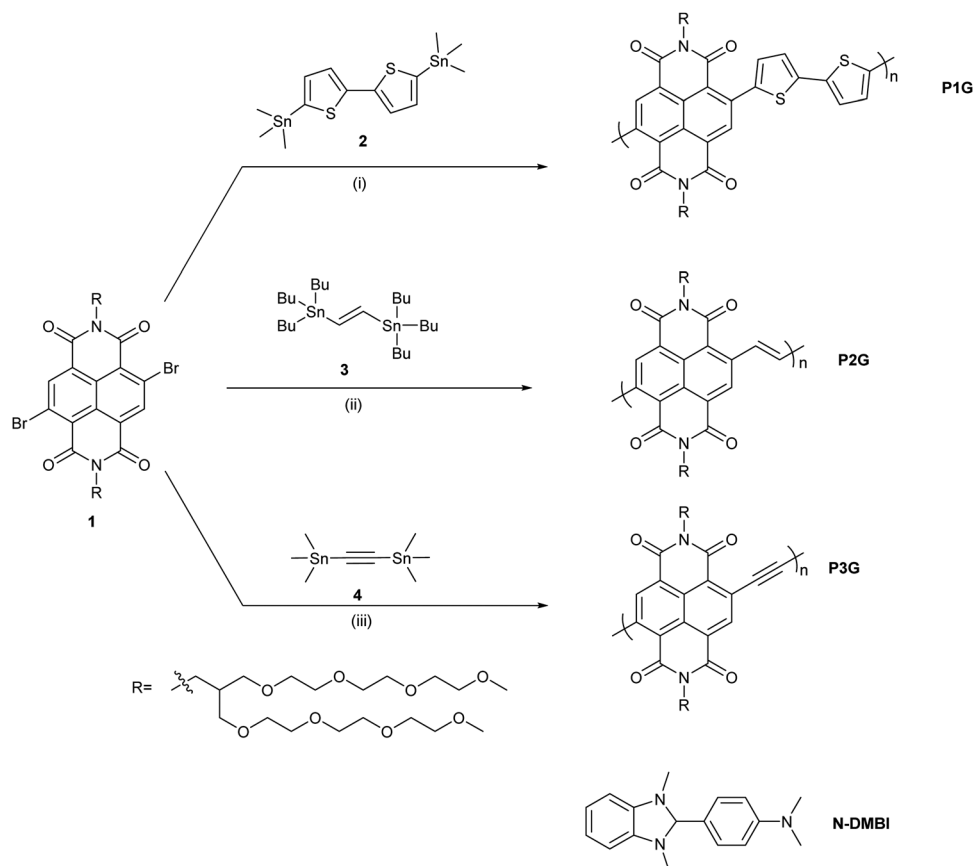
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† Electronic supplementary information (ESI) available. See DOI: 10.1039/d1mh01357b

In this contribution, we explore the backbone modification of NDI-based CPs and its effect on molecular doping.

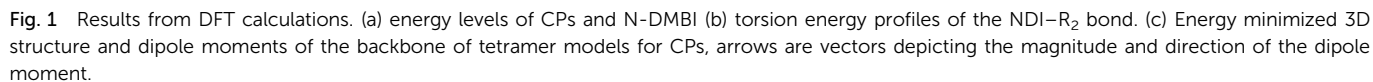
Experimental results showed that all three polymers can be doped by N-DMBI, however, the conductivity of P1G outperformed the other two polymers owed to a higher crystallinity and good miscibility with N-DMBI. While P2G exhibits good miscibility with N-DMBI, the lack of long-range order limits its σ . P3G shows poor host-dopant miscibility and a low degree of molecular order, which make it the polymer with the lowest σ . Finally, we analysed the temperature-dependent thermoelectric properties of P1G as its electrical conductivity is the largest of the three CPs. Our results show that OTEs can benefit from branched TEG side chains, as only 5 mol% of N-DMBI suffices to maximize the power factor (PF). Remarkably, P1G shows an in-plane thermal conductivity ($\kappa_{||}$) of only $0.13 \text{ W m}^{-1} \text{ K}^{-1}$, one of the lowest reported values for CPs. This suggests that branched TEG side chains are a promising strategy to reduce the thermal conductivity of CPs, in addition to improving their miscibility with polar dopants.



Mater. Horiz., 2022, 9, 500–508 | 501

Previous reports suggest that the bithiophene-induced backbone twist in N2200 and perylene diimide-bithiophene CPs limits the polaron delocalization length and the electrical conductivity, showing that more planar backbones possess higher σ .^{24,30} This makes backbone planarization an attractive approach to improve σ of n-type doped CPs. Hence, the designs of P2G and P3G are aimed at increasing the backbone planarity. Previous reports have used non-bonding conformational locks to improve molecular packing and charge transport properties.^{31–33} To better understand the role of the polymer conjugated backbone on molecular doping, we considered three parameters. First, the energetics of the polymers (considering lowest unoccupied molecular orbitals (LUMO)), second, the backbone planarity and rigidity, and third, the dipole moment of the backbone,

The energy levels of the frontier molecular orbitals for the tetramers of the CPs were calculated by DFT (Fig. S4, ESI[†]), the energy levels are represented in Fig. 1a. P1G shows the shallowest LUMO (−3.46 eV) followed by P2G (−3.72 eV) and P3G (−3.98 eV), suggesting that all three polymers should be efficiently doped by the singly occupied molecular orbital (SOMO) of N-DMBI at −2.36 eV (Fig. S5, ESI[†]). For P1G, the HOMO is mainly delocalized on the bithiophene donor subunit, and the LUMO is delocalized on the NDI acceptor subunit, similarly to previously reported D–A CPs.³⁴ In contrast, P2G and P3G show both HOMO and LUMO delocalized along the whole tetramer backbone, such delocalization might be beneficial for charge transport by extending the probability of both intra- and intermolecular hopping, consequently improve σ .^{35,36} Additionally, we calculated the torsion potential for all three polymers (Fig. 1b). P1G showed a twisted backbone with an angle of $\sim 48^\circ$ for the bond between the NDI and 2T subunit, along with a moderate torsion energy barrier of 0.13 eV. P2G showed a more planar backbone with an angle of $\sim 25^\circ$ between NDI and vinylene subunits, it also possesses the largest torsion energy barrier (0.27 eV) making it the most rigid backbone of the three CPs, owing to the interaction of the carbonyl group in the NDI and the hydrogen of the vinylene which work as a non-bonding conformational lock.^{32,37} P3G shows the most planar backbone,



The electrical conductivity of pristine and doped polymers was compared by spin coating films on top of glass substrates with pre-patterned gold electrodes (Fig. 2a-c). All samples showed improved σ upon addition of N-DMBI, however, this change is more significant in the case of P1G, where σ of the doped films is about three orders of magnitude higher compared to neat polymer. In the case of P2G and P3G, the increment was about two orders of magnitude for each polymer. Additionally, P1G reached its maximum σ at rather low dopant concentrations, *i.e.* 10 to 15 mol% N-DMBI. The addition of higher dopant concentrations led to a reduction of σ , possibly due to a decrease in mobility (μ) originated by carrier-carrier scattering.³⁹ In contrast P2G needed 25 mol%, and P3G more than 30 mol% dopant concentration to achieve their maximum conductivity. Further increasing the dopant ratio

Figure 2 consists of four panels (a, b, c, d) showing the effect of N-DMBI concentration on the conductivity and spin density of three polymers: P1G, P2G, and P3G.

(a) Conductivity of P1G vs N-DMBI mol %: A log-linear plot showing conductivity (S/cm) on the y-axis (log scale from 10^{-5} to 1) versus N-DMBI mol % in P1G on the x-axis (Pristine, 5%, 10%, 15%, 20%, 25%, 30%). Conductivity increases from ~ 2×10^{-5} S/cm at 5% to a peak of ~ 5×10^{-2} S/cm at 15%, then decreases to ~ 5×10^{-3} S/cm at 30%.

(b) Conductivity of P2G vs N-DMBI mol %: A log-linear plot showing conductivity (S/cm) on the y-axis (log scale from 10^{-5} to 0.01) versus N-DMBI mol % in P2G on the x-axis (Pristine, 5%, 10%, 15%, 20%, 25%, 30%). Conductivity increases from ~ 5×10^{-5} S/cm at 5% to a peak of ~ 1.5×10^{-3} S/cm at 25%, then decreases to ~ 1×10^{-3} S/cm at 30%.

(c) Conductivity of P3G vs N-DMBI mol %: A log-linear plot showing conductivity (S/cm) on the y-axis (log scale from 10^{-6} to 0.001) versus N-DMBI mol % in P3G on the x-axis (Pristine, 5%, 10%, 15%, 20%, 25%, 30%). Conductivity increases from ~ 3×10^{-6} S/cm at 5% to a peak of ~ 3×10^{-4} S/cm at 20%, then decreases to ~ 4×10^{-5} S/cm at 30%.

(d) Spin density of P1G, P2G, and P3G vs N-DMBI concentration: A log-linear plot showing spin density (cm^{-3}) on the y-axis (log scale from 10^{18} to 10^{20}) versus N-DMBI concentration (mol %) on the x-axis (0, 5, 10, 15, 20, 25, 30). Three data series are shown: P1G (blue dashed line with circles), P2G (red dashed line with circles), and P3G (purple dashed line with circles). All three polymers show an increase in spin density with increasing N-DMBI concentration, with P2G reaching the highest spin density of ~ $1.2 \times 10^{20} \text{ cm}^{-3}$ at 30%.

Mater. Horiz., 2022, 9, 500–508 | 503

for P3G proved too challenging owed to the poor solubility of P3G, which limited the formation of continuous films and the accurate measurement of σ . These observations correlate with the calculated dipole moments. The polymer with the largest dipole moment (P1G) requires the least amount of N-DMBI to reach its maximum conductivity, and the one with the smallest dipole moment (P3G) requires more than twice the amount to achieve its maximum conductivity. Previous studies in fullerene systems have shown that increasing host-dopant miscibility can decrease the amount of dopant needed to achieve maximum conductivity.⁴⁰ We further investigated these phenomena using AFM topography and phase imaging. In the case of P1G and P2G (Fig. S8 and S9, ESI[†]), no dopant clustering appeared, even at high doping concentrations. In the case of P3G, the topography images (Fig. S10, ESI[†]) reveal the formation of several aggregates only present in the doped polymer films. These aggregates show a strong contrast in the AFM phase image indicating such clusters are formed by phase-segregated dopant and implying poor host-dopant miscibility. Such phase segregation explains in part the much lower conductivity observed in P3G. Additionally, we measured the contact angle between water and the undoped CPs, N-DMBI, as well as the blended host-dopant polymer films, to determine their hydrophobic properties. Results are shown in Fig. S11 (ESI[†]). P1G and P2G show values of 54.1° and 31.6° respectively. These values are closer to the contact angle of N-DMBI (41.6°) than that of P3G (64.9°) suggesting better miscibility in the systems P1G:N-DMBI and P2G:N-DMBI, than in P3G:N-DMBI. This further confirms the observations carried out in AFM. Finally, the addition of N-DMBI to either P1G or P2G resulted in a decrease of the contact angle, whereas P3G remained unchanged.

To further elaborate on the differences in electronic properties of the three CPs, we proceeded to perform advanced spectroscopic analysis. Electron paramagnetic resonance (EPR) provides a quantitative analysis of unpaired electrons in the form of radical anions within a material, which can be correlated to the charge carrier density and offer an understanding of the differences observed in

σ . We performed EPR in pristine and doped films of the NDI copolymers, the resulting spectra are shown in Fig. S12a–c (ESI[†]). The undoped polymer films show very weak EPR signals. The addition of N-DMBI leads to the generation of radical anions and thus large increments of the signal intensity in all polymers. Quantitative analysis of the signals provided the spin density of the samples (Fig. 2d). We observe that the undoped P1G film has a larger spin density ($5 \times 10^{18} \text{ cm}^{-3}$) than the undoped films of P2G and P3G ($1 \times 10^{18} \text{ cm}^{-3}$). As N-DMBI is added to the samples, the spin density increases, and both P1G and P2G show spin densities of about $1 \times 10^{20} \text{ cm}^{-3}$ while P3G shows only $3 \times 10^{19} \text{ cm}^{-3}$. Although this partially explains the lower σ of P3G, it does not explain the difference in σ of P1G and P2G at their respective maximums.

The ionization energy (IE) and electron affinity (EA) of all polymers P1G, P2G, and P3G were directly measured on thin films using ultraviolet photoelectron spectroscopy (UPS) and low-energy inverse photoelectron spectroscopy (LE-IPES) (Fig. 3a and b respectively). The IE and EA of P1G and P2G were found to be 5.58 eV/3.62 eV and 5.67/3.38 eV, respectively. Thin films of P3G were found to have an IE of 5.41 eV and the deepest EA of 3.73 eV among all three polymers. It must be noted that the DFT calculations estimated the orbital energies of tetramers, whereas the spectroscopic studies were carried out in thin solid films of polymers. The orbital energies of the isolated molecules and those in thin solid films are different,⁴¹ and would explain the differences in the calculated HOMO/LUMO and the measured IE/EA values. Among all the three polymers, P1G was found to have the most prominent peak at the onset of both UPS and IPES spectra, corresponding to the density of states (Fig. 3a and b) in the IE and EA, respectively. We attribute these features to an additional density of states (DOS) introduced by the ICT between the NDI and 2T subunits of P1G. These features can also be correlated to the peak at 725 nm in the absorption spectrum of P1G.³⁸ Since the electronic states in the IE and EA regions contribute to the electrical conductivity of the materials, P1G is expected to have the

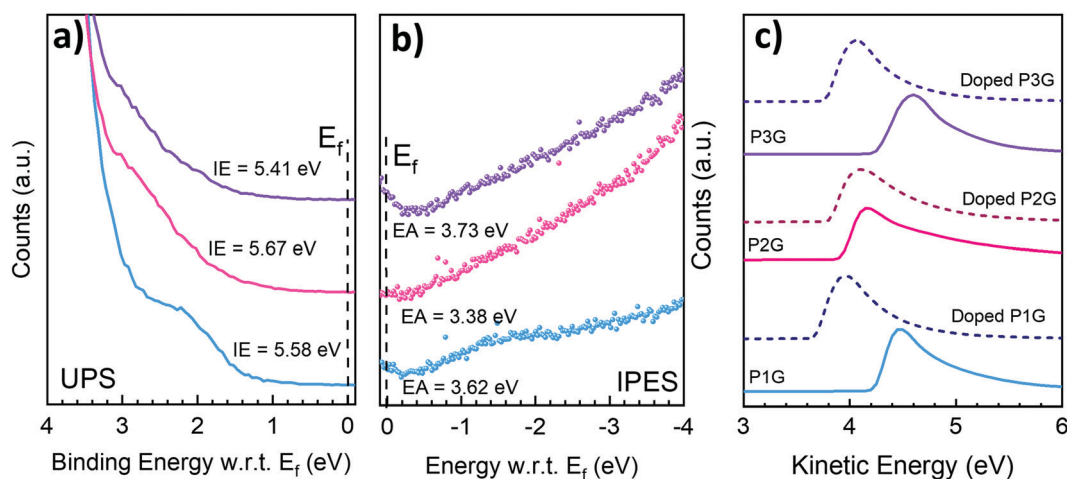
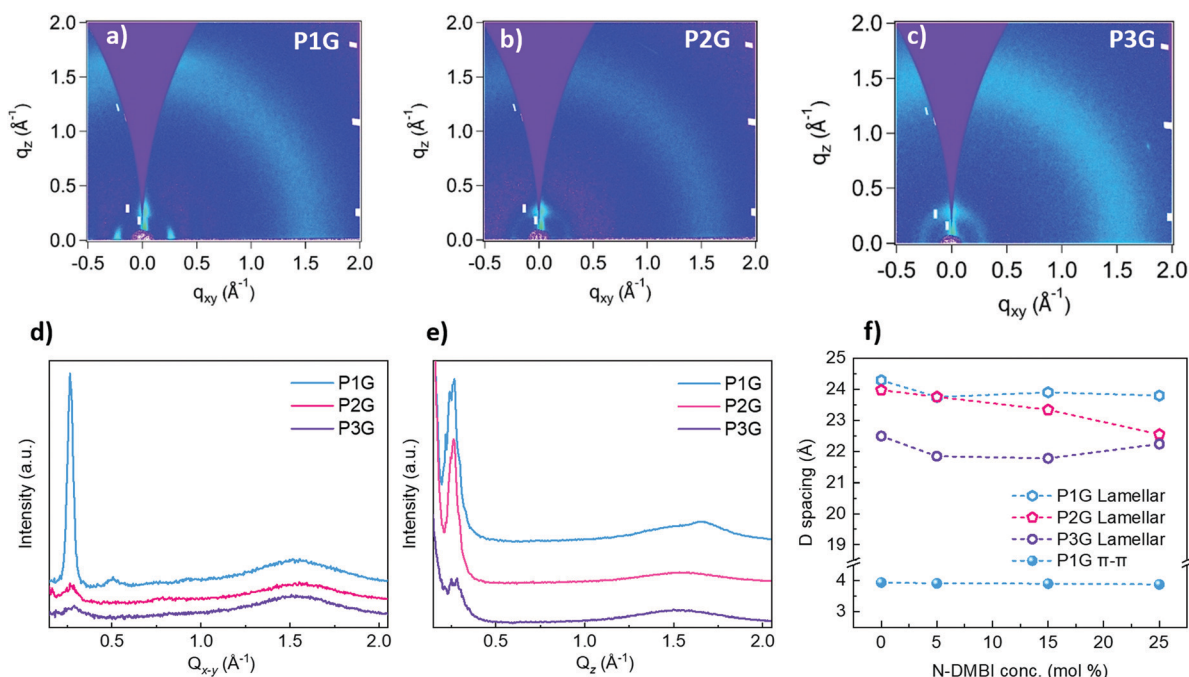


Fig. 3 Photoelectron spectroscopy of the CPs. (a) shows the IE, (b) corresponds to the EA. (c) SECO showing the work function of the undoped films (continuous lines) and their shift upon doping (dashed lines).



GIWAXS was carried out to study the morphology of the CP films. The 2D data and linecuts are displayed in Fig. 4, (see

P1G shows higher σ at very low dopant loading concentrations suggesting its suitability for thermoelectric applications. We measured its temperature-dependent thermoelectric (TE) properties with a chip-based technique. σ and Seebeck coefficient (α) were measured using Van der Pauw configuration to measure temperature-dependent σ and α ; temperature-dependent $\kappa_{||}$ was measured in the same chip using 3ω method. In the case of P2G and P3G, σ is lower than the measurement



Mater. Horiz., 2022, 9, 500–508 | 505

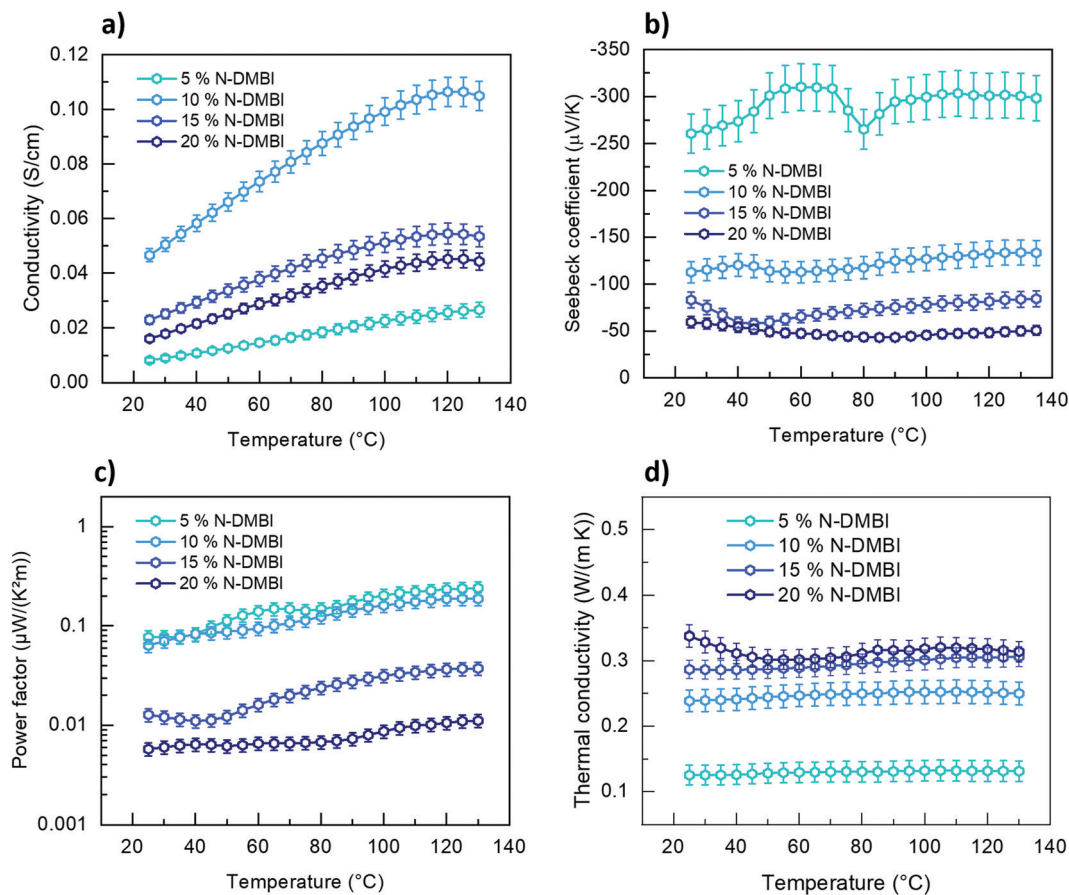


Fig. 5 Temperature-dependent thermoelectric parameter of P1G (a) σ , (b) α , (c) PF, and (d) $\kappa_{\parallel}$.

limit of the instrument. The temperature-dependent σ (Fig. 5a) shows that P1G follows typical semiconductor behaviour *i.e.* increasing proportionally with temperature. The highest α was measured at 5 mol% N-DMBI, at higher dopant concentrations this value decreases rapidly, as increased dopant concentrations reduce the average energy of the charge carriers, also known as transport energy (E_t).⁴² α shows negligible dependency on the temperature. The highest power factor for P1G was also achieved at this doping concentration ($0.077 \mu\text{W K}^{-2} \text{m}^{-1}$ at room temperature and $0.25 \mu\text{W K}^{-2} \text{m}^{-1}$ at 130 $^{\circ}\text{C}$). Finally, P1G shows an ultra-low $\kappa_{\parallel}$ of $0.13 \text{ W m}^{-1} \text{K}^{-1}$, which is one of the lowest values known for CPs so far. N2200, the alkyl-substituted analogue of P1G, possesses an in-plane thermal conductivity of $0.27 \text{ W m}^{-1} \text{K}^{-1}$.⁴³ $\kappa_{\parallel}$ of P1G shows a low dependency on temperature within the studied temperature window. However, increasing the dopant concentration led to increased $\kappa_{\parallel}$, hence, the highest figure of merit (ZT) of 2×10^{-4} was achieved at 5 mol% of N-DMBI (Fig. S16, ESI†). This suggests that TEG side chains are a promising approach to reduce the thermal conductivity of CPs, a key factor to improve their figure of merit for thermoelectric applications.⁴⁴

Conclusions

A series of NDI-based copolymers with branched TEG side chains were synthesized, proposing conjugated vinylene and

acetylenic linkers to reduce steric hindrance and improve backbone planarity and rigidity. While these spacer groups can effectively produce planar backbones, only P1G results in oriented molecular packing, P2G and P3G on the other hand are mostly amorphous suggesting backbone planarization alone is not enough to produce long-range ordering. While n-type molecular doping was demonstrated for all three polymers, P1G outperformed its analogues owing to its more ordered and long-range molecular packing and better host-dopant miscibility. P2G showed lower electrical conductivity values, with good host-dopant miscibility, but no long-range molecular packing. Finally, N-type doped P3G films showed the lowest conductivity of the three polymers due to its amorphous nature and poor host-dopant miscibility. This study showcases the impact of the backbone orientation on the efficiency of molecular doping and in the miscibility of molecular dopants where low polarity and miscibility of P3G overcame the presence of branched TEG side chains, limiting molecular doping with N-DMBI. Finally, the thermoelectric properties of P1G showed a maximum power factor at 5 mol% of N-DMBI, with a particularly low in-plane thermal conductivity. We believe this low value is enabled by the disordered nature of the branched TEG side chains, which opens a new path to discovering new conjugated polymers with ultra-low thermal conductivity values.



Conflicts of interest

The authors declare no conflicts of interest.

Acknowledgements

This publication is based upon work supported by the King Abdullah University of Science and Technology (KAUST) Office of Sponsored Research (OSR) under Award No. OSR-CRG2018-3737 and British Council Newton Fund Institutional Links (ref: 337067). B.C.S. acknowledges the UK Research and Innovation for Future Leaders Fellowship no. MR/S031952/1 and EPSRC grant (EP/P007767/1). L. G. and X. G. thanks U.S. Department Energy provides partial support for the scattering experiments in this work under award No. DE-SC0022050. The work by J. L. and L. J. A. K. was supported by a grant from STW/NWO (VIDI 13476).

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