



Cite this: *Soft Matter*, 2025,
21, 9275

Dynamical behavior of compound vesicles in wall-bounded shear flow

Antonio Lamura

We report a numerical study addressing the dynamics of compound vesicles confined in a channel under shear flow. The system comprises a smaller vesicle embedded within a larger one and can be used to mimic, for example, leukocytes or nucleate cells. A two-dimensional model, which combines molecular dynamics and mesoscopic hydrodynamics including thermal fluctuations, is adopted to perform an extended investigation. We are able to vary independently the swelling degree and the relative size of vesicles, the viscosities of fluids internal and external to vesicles, and the Capillary number, so to observe a rich dynamical phenomenology which goes well beyond what observed for single vesicles, matching quantitatively with experimental findings. Tank-treading, tumbling, and trembling motions are enriched by dynamical states where inner and outer vesicles can perform different motions. We show that thermal fluctuations are crucial during trembling and swinging dynamics, as observed in experiments. Undulating motion of the external vesicle, characterized by periodic oscillation of the inclination and buckling of the membrane, is observed at high filling fractions. This latter state exhibits features that are shown to depend on the relative size, the swelling degree of both vesicles as well as on thermal noise lacking in previous analytical and numerical studies.

Received 28th October 2025,
Accepted 17th November 2025

DOI: 10.1039/d5sm01079a

rsc.li/soft-matter-journal

1. Introduction

Compound vesicles, which comprise a smaller vesicle embedded within a larger one, are systems of relevant biological importance since can be seen as biomimetic models for multi-compartmentalized cells such as leukocytes^{1–5} or nucleate cells.^{6–8} Furthermore, vesicles embedding one or even more vesicles, as in the case of vesosomes,^{9,10} are also of interest for drug delivery applications,¹¹ where controlled deformation and compartmentalization can be properly exploited.¹² Compound vesicles under shear flow have attracted attention due to their relevance in physiological conditions and microfluidic systems, where cells or synthetic vesicles experience different dynamical states.¹³ The study of their dynamical behavior requires to take into account inter-membrane interactions, membrane dynamics, and fluid-structure coupling when exposed to external flow.

The dynamical responses of single vesicles in shear flow, including tank-treading (TT), tumbling (TU), and trembling (TR) motions, have been extensively studied theoretically,^{14–18} numerically,^{19–26} and experimentally.^{27–31} However, the introduction of an internal vesicle produces a richer dynamics since extra degrees of freedom are added such as internal vesicle rotation, displacement, and deformation, as well as

hydrodynamic coupling between the internal and external membrane. These features make compound vesicles a system suitable to understand the interplay between membranes dynamics under geometrical constraints and hydrodynamic interactions.

Recent studies have revealed, in addition to motions observed for single vesicles, novel dynamical modes unique to compound vesicles, such as synchronized or asynchronous TU of the internal and external vesicle,^{5,13} swinging (SW) of the internal vesicle, whose main axis oscillates around a positive mean angle, at large shear rates,^{2,13} and undulating (UND) motion where the external vesicle oscillates around a positive mean angle while its shape undergoes large undulations.³ Nevertheless, the parameter space governing these behaviors – including reduced volumes of internal and external vesicles, ratio of the viscosity of the fluid around the internal vesicle to the viscosity of the outer fluid, Capillary number, and occupied fraction of the interior volume – remains only partially mapped, and many questions persist about dynamical behaviors. Indeed, theoretical approaches^{2,5} are limited to consider vesicles in the quasi-spherical approximation, available numerical simulations^{2,3} neglect thermal fluctuations that are known to be relevant in TR,^{13,32–34} and experimentally it is not trivial to control physical and geometrical parameters of compound vesicles.¹³

In the present study, we aim at exploring the behavior of compound vesicles through numerical simulations accessing wide ranges of the parameters controlling the system. Our

Istituto Applicazioni Calcolo, Consiglio Nazionale delle Ricerche (CNR), Via Amendola 122/D, 70126 Bari, Italy. E-mail: antonio.lamura@cnr.it



results are achieved by using mesoscale hydrodynamic simulations of a two-dimensional model system, including thermal fluctuations. These are expected⁵ to be crucial in order to capture features observed only in experiments.¹³ We can vary the swelling degree of both vesicles going from deflated to quasi-circular shapes, to tune independently viscosities for the fluids around and inside each vesicle, to change the Capillary number, and to consider different sizes of the internal vesicle. In this way, we are able to observe a rich phenomenology in the dynamical behavior of compound vesicles which well matches with experimental observations of ref. 13. Moreover, our study allows a full characterization of the UND motion under different conditions to an extent never observed before.

2. Numerical model

We consider a two-dimensional system consisting of a compound vesicle, made of one internal (int) vesicle embedded in an external (ext) one, which is suspended in a shearing fluid (see Fig. 1).

The fluid in each of the three zones, separated by the vesicle membranes, is described by the multi-particle collision dynamics^{35–39} with viscosities η_{out} , η_{an} , η_{in} for the outer (out), annular (an), and inner (in) areas, respectively. The system of size $L_x \times L_y$ is placed between two horizontal walls sliding along the x -direction (flow direction) with velocities $\pm v_w$. Periodic boundary conditions are imposed along the flow direction and bounce-back boundary conditions are implemented at the walls.^{40,41} This allows us to simulate a linear flow profile $(\dot{\gamma}y, 0)$ with shear rate $\dot{\gamma} = 2v_w/L_y$, y being the vertical coordinate along the shear direction. Each vesicle is modeled by enforcing the conservation of internal area A_0 and perimeter L_0 of the enclosing membrane, and simulated *via* molecular dynamics. Internal and external membranes are characterized by the same bending rigidity κ and experience a mutual repulsion at short distances.⁴² Solvent particles and membranes interact *via* three-body collisions which conserve linear and angular momenta.²² The details for the numerical implementation of fluid, vesicles, and solvent-membrane coupling can be found in ref. 26 and 43. It is useful for the following analysis to introduce

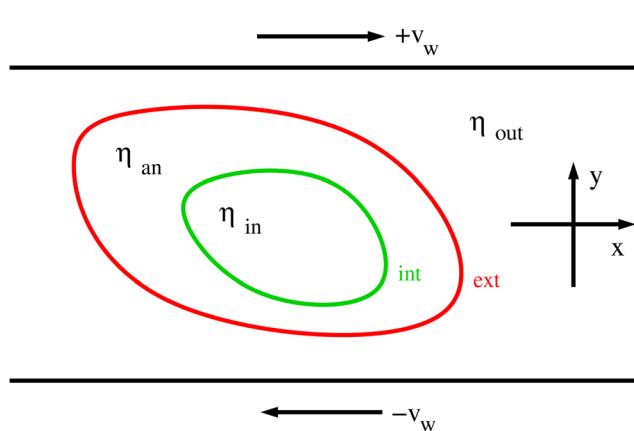


Fig. 1 Schematic layout of the simulated system.

the excess length $\Delta = L_0/\sqrt{A_0/\pi} - 2\pi$. Δ is a measure of the swelling degree, being $\Delta > 0$ in the case of deflated vesicles and $\Delta = 0$ for circular ones. The vesicle dynamical states in shear flow can be characterized by the dimensionless quantities $\Lambda = (32 + 23\lambda)\sqrt{\Delta_{\text{ext}}}/(8\sqrt{30\pi})$ and $S = 14\pi\dot{\gamma}^*/(3\sqrt{3}\Delta_{\text{ext}})^{16}$ where $\lambda = \eta_{\text{an}}/\eta_{\text{out}}$ is the viscosity contrast and $\dot{\gamma}^* = \dot{\gamma}\eta_{\text{out}}R_{0,\text{ext}}^3/\kappa$ is the Capillary number (or reduced shear rate), $R_0 = L_0/(2\pi)$ being the vesicle radius. The use of Λ and S will allow us to make easy comparison with the experimental results.¹³ The compound vesicle dynamics is also affected by the occupied fraction of the interior area of the external vesicle, defined to be $\phi = R_{0,\text{int}}/R_{0,\text{ext}}$,^{3,5,13} and by the excess length Δ_{int} of the internal vesicle.¹³

The novelty of the present study is that we are able to vary all the four parameters Λ , S , ϕ , Δ_{int} in wide ranges, accessing different dynamical states of the compound vesicle. In the following we set $0.16 \leq \Delta_{\text{ext}} \leq 1.23$, $0.16 \leq \Delta_{\text{int}} \leq 0.74$, $0.2 \leq \phi \leq 0.8$, $1 \leq \lambda \leq 15$, $1 \leq \dot{\gamma}^* \leq 20$, $\eta_{\text{in}} = \eta_{\text{out}}$, $L_x/R_{0,\text{ext}} = 12.04$, and $L_y/R_{0,\text{ext}} = 5.76$, while maintaining fixed $R_{0,\text{ext}}$. The bending rigidity is chosen to be $\kappa = 3.27k_B T R_{0,\text{ext}}$, where $k_B T$ is the thermal energy of the system. This value of κ gives rise to a comparable amplitude of undulation modes as for compound vesicles in experiments¹³ (see the following Fig. 3(a)). All simulations are run by keeping the inertial effects negligible: this is assured by making the Reynolds number $\text{Re} = \dot{\gamma}\rho_{\text{out}}R_{0,\text{ext}}^2/\eta_{\text{out}} \lesssim 0.1$, ρ_{out} being the outer fluid density.

3. Results

In this section we show and characterize the different dynamical states of compound vesicles discussing the results and comparing them with previous studies.

3.1. Tank-treading motion

At values $\Lambda \lesssim 1$ and $10 \lesssim S \lesssim 10^3$, compound vesicles are found in a steady TT regime. Both internal and external vesicle keep a steady inclination angle θ , defined as the angle between the vesicle main long axis and the flow direction, while the membranes perform a tank-treading motion with a characteristic frequency ω . This motion is observed for different swelling degrees of internal and external vesicles in a wide range of the filling fraction ($0.4 \leq \phi \leq 0.8$). In order to compute the inclination angle, the two eigenvalues λ_M and λ_m , with $\lambda_M > \lambda_m$, of the vesicle gyration tensor are extracted. The inclination angle is defined as the angle formed by the eigenvector corresponding to λ_M with the flow direction. The average inclination angles $\langle \theta \rangle$ of compound vesicles are shown in Fig. 2(a). Data of the external vesicle (red triangles) show a monotonic decrease with Λ . Moreover, the values of $\langle \theta \rangle_{\text{ext}}$ are similar to the ones observed for single vesicles,²⁶ which are reported in the same figure (black circles) and can be quite well fitted by a straight line. The comparable values of the inclination angles suggest that the transition at $\langle \theta \rangle \simeq 0$ from the TT regime to the TR one occurs at similar value $\Lambda_c \simeq 1.91$, and that the internal vesicle has negligible effect on the external

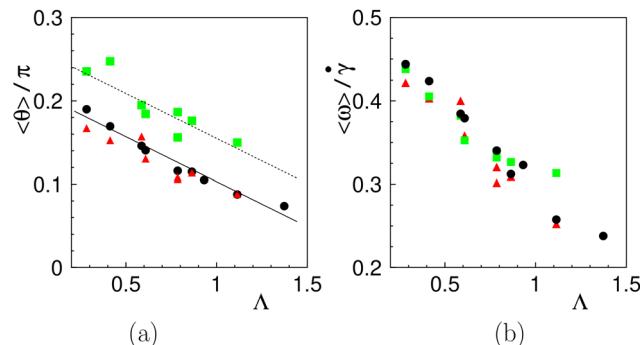


Fig. 2 (a) Dependence of the average inclination angle on Λ for the external (red triangles) and internal (green squares) tank-treading vesicles compared with data of single tank-treading vesicles (black circles).²⁶ The full line is the fit to the data of single vesicles and the dashed line is the same fit line shifted up. The data shown are obtained from different compound vesicles with $0.16 \leq \Lambda_{\text{ext}} \leq 1.23$, $0.16 \leq \Lambda_{\text{int}} \leq 0.74$, and $0.4 \leq \phi \leq 0.8$. In the case of single vesicles it is $0.16 \leq \Lambda \leq 1.23$. (b) Dependence of the average tank-treading frequency on Λ for the external (red triangles) and internal (green squares) vesicles compared with data of single vesicles (black circles).

dynamics. On the other hand, the internal vesicle attains larger inclination angles (green squares) with the same dependence on Λ so that the data can be interpolated by simply shifting up the fit line when using the external flow parameters Λ and S . The larger values of $\langle \theta \rangle_{\text{int}}$ might be due to the confinement exerted by the external vesicle, a phenomenon already observed for single vesicles.⁴⁴ As a consequence, the relative angle $\langle \theta \rangle_{\text{int}} - \langle \theta \rangle_{\text{ext}}$ does not depend on Λ and ϕ . These results are fully consistent with the ones observed in the only available experiments on sheared compound vesicles.¹³ Numerical simulations of a compound vesicle with a solid spherical inclusion² and with an inner vesicle^{3,5} show a dependence of $\langle \theta \rangle_{\text{ext}}$ and $\langle \theta \rangle_{\text{int}}$ on ϕ . Moreover, differences in the values of $\langle \theta \rangle_{\text{ext}}$ and Λ_c , with respect to the case of a single vesicle, are observed. These discrepancies can be attributed to the lack of thermal noise, which is known to affect the TT-TR transition,³⁴ in previous simulations. We can also compute the tank-treading frequencies ω whose average values, rescaled by the shear rate $\dot{\gamma}$, are shown for the external (red triangles) and internal (green squares) membranes of a compound vesicle and for a single vesicle (black circles) in Fig. 2(b). Data points refer to the same systems presented in the panel (a). Remarkably, it appears that the internal and external membranes as well as the membrane of single vesicle rotate with comparable rescaled average frequencies which show the same dependence on Λ , while being not dependent on both the filling fractions and the excess lengths.

In order to quantitatively characterize the shapes of the external vesicle during the TT motion, we parametrize the membrane contour as a function of the polar angle ϕ ($0 \leq \phi \leq 2\pi$) by writing the position $\mathbf{r}(\phi)$ of any point on the membrane as $\mathbf{r}(\phi) = R_{0,\text{ext}} \mathbf{e}_r(\phi) [1 + u(\phi)]$, where $\mathbf{e}_r(\phi)$ is a unit vector along the radial direction and $u(\phi)$ is the dimensionless shape deformation. Since $u(\phi)$ is a real periodic function of ϕ , it can be expanded via Fourier transform as $u(\phi) = \sum_m u_m \exp(im\phi)$.

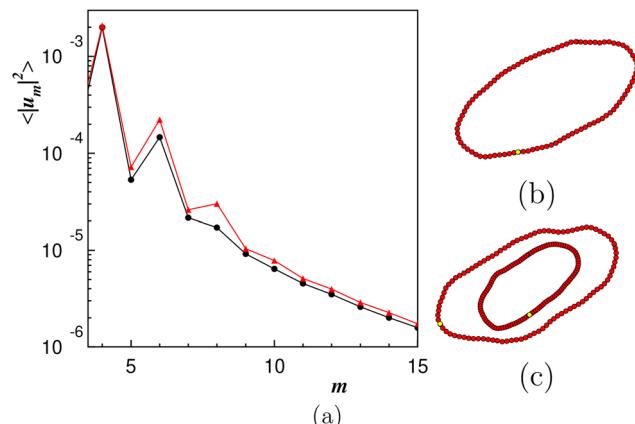


Fig. 3 (a) Comparison of the spectra, averaged in time, of amplitudes of shapes for a single vesicle with $\Lambda = 0.74$ (black circles) and a compound one with $\Lambda_{\text{ext}} = \Lambda_{\text{int}} = 0.74$, $\phi = 0.6$ (red triangles). Both single vesicle and compound one are in tank-treading motion at $(\Lambda, S) = (0.61, 229)$. Corresponding typical configurations are shown in panels (b) and (c), respectively.

The computed stationary correlations of the complex Fourier modes u_m are shown in Fig. 3(a) for a compound vesicle with $\Lambda_{\text{ext}} = \Lambda_{\text{int}} = 0.74$ and $\phi = 0.6$, and for a single vesicle with the same excess length. Both single vesicle and compound one are in tank-treading motion with flow parameters $(\Lambda, S) = (0.61, 229)$. The main feature that appears, is the amplification of the even modes of the external vesicle spectrum at $m \geq 6$ (red triangles) compared with those measured for the single vesicle (black circles). This is due to the fact that the shape of the external vesicle is modified in the compound one, being more parallelogram-like (see Fig. 3(c)), with respect to the ellipsoidal shape of the single vesicle (see Fig. 3(b)).^{18,26} These features match with what observed in experiments¹³ where the measured values of $\langle |u_m|^2 \rangle$ (see Fig. 2 of ref.13) are in a good quantitative agreement with the present ones.

3.2. Trembling motion

Upon increasing Λ beyond Λ_c , the external vesicle performs TR motion during which its main axis oscillates around the flow direction while undergoing strong shape modifications, as found for single vesicles.^{45,46} The inner vesicle is observed to perform SW motion, which is characterized by large variations of the inclination angle around a positive mean value, as depicted in Fig. 4(a) for small ϕ at $(\Lambda, S) = (1.94, 1058)$. It can be seen that at $\dot{\gamma} \simeq 48$ an isolated tumbling event occurs. This reminds the intermediate regime between TU and SW observed in simulations of quasi-spherical compound vesicles.⁵ In order to gain information about the shape deformation, the asphericity $A = [(\lambda_M - \lambda_m)/(\lambda_M + \lambda_m)]^2$ is computed, being $A = 0$ for circular shapes. The time evolution of A for the internal and external vesicles is shown in Fig. 4(b) where it can be appreciated that A_{ext} widely oscillates in time, appearing synchronized with θ_{ext} , as in the case of single vesicles.²⁴ The supplementary animation file movie1.mp4 provides a comprehensive view of the TR/SW motion. A similar dynamics has



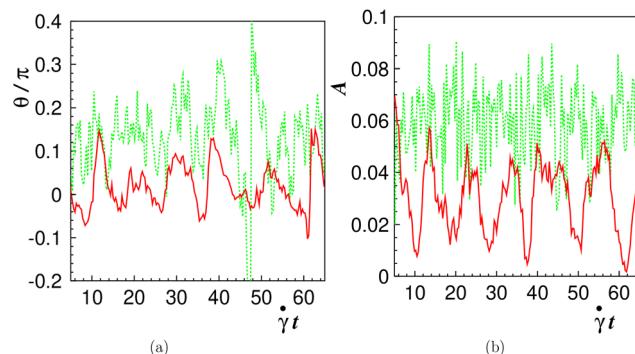


Fig. 4 (a) Time evolution of the inclination angles θ_{ext} (red continuous line) and θ_{int} (green dashed line) of a compound vesicle with $\Delta_{\text{ext}} = \Delta_{\text{int}} = 0.16$, $\phi = 0.2$ at $(\mathcal{A}, S) = (1.94, 1058)$ showing trembling (external vesicle) and swinging (internal vesicle) motion. (b) Time evolution of the asphericities A_{ext} (red continuous line) and A_{int} (green dashed line) corresponding to the case in panel (a).

been observed for comparable values of ϕ , \mathcal{A} , and S in the experiments of ref. 13 (see Fig. 4 therein). The SW motion has been also seen for the external vesicle of a vesicle embedding a solid ellipsoidal inclusion which performs TU motion.² The internal vesicle has been found in SW state also in previous simulations of compound vesicles.^{3,5}

3.3. Tumbling motion

A further increase of \mathcal{A} determines the transition of the compound vesicle to TU motion: both internal and external vesicle rotate as solid-like elongated particles. Their angular velocities are triggered by the fraction ϕ of the occupied interior area. At small ϕ , the internal and external vesicle motions are weakly correlated. Their centers are not stationary, differently from what happens in studies where thermal noise is discarded,^{3,5} and their TU velocities slightly deviate from each other. This desynchronizes $\theta_{\text{int}}(t)$ and $\theta_{\text{ext}}(t)$, as shown in Fig. 5(a) for $\phi = 0.2$ at $(\mathcal{A}, S) = (5.38, 138)$. The relative difference

between the inclination angles is depicted in panel (b) of Fig. 5 in order to quantify the aforementioned desynchronization. Larger values of ϕ allow the internal and external vesicles to tumble in a synchronous way, as illustrated in Fig. 5(c) at $(\mathcal{A}, S) = (4.18, 23)$. These features are in common with the theoretical analysis of ref. 5 for quasi-spherical compound vesicles, and have been as well observed in the experiments of ref. 13 (see Fig. 5 and 6 therein) for comparable values of the control parameters \mathcal{A} , S , ϕ , Δ_{int} .

3.4. Undulating motion

For large values of \mathcal{A} and S and high filling fraction ($\phi \lesssim 0.6$), the UND motion of the external vesicle appears. This is characterized by regular oscillations of the inclination angle $\theta_{\text{ext}}(t)$ around a positive mean angle $\langle \theta_{\text{ext}} \rangle$ with wide shape undulations³ during the rotation of the membrane. By increasing the filling fraction ϕ , the amplitudes of $\theta_{\text{ext}}(t)$ reduce while the average value $\langle \theta_{\text{ext}} \rangle$ slightly increases. ϕ affects the motion of the internal vesicle, as well, that can perform either SW, as shown in Fig. 6(a) for $\phi = 0.6$ (see also the supplementary animation file movie2.mp4), or TT when increasing ϕ , as shown in Fig. 7(a) for $\phi = 0.8$ (see also the supplementary animation file movie3.mp4). In both cases, the flow parameters are $(\mathcal{A}, S) = (5.38, 138)$ and the external vesicle is more deflated than the internal one, being $\Delta_{\text{ext}} = 1.23$ and $\Delta_{\text{int}} = 0.74$. It is interesting to note that during the UND motion, the external membrane buckles assuming lobated shapes which depend on the filling fraction, as shown in Fig. 8(a) and (b). The shape variations are quantified in the panels (b) of Fig. 6 and 7 where the time behaviors of the asphericities A_{ext} and A_{int} are depicted. It is evident that A_{ext} oscillates with the same frequency of θ_{ext} , reaching a local maximum when the inclination angle is at its maximum. It has to be remarked that the amplitudes of $A_{\text{ext}}(t)$ during UND motion are much larger compared to those ones observed in TR motion (see Fig. 4(b)). Moreover, the oscillatory pattern can be used to determine a characteristic frequency $\omega_{\mathcal{A}}$ by performing a spectral analysis of

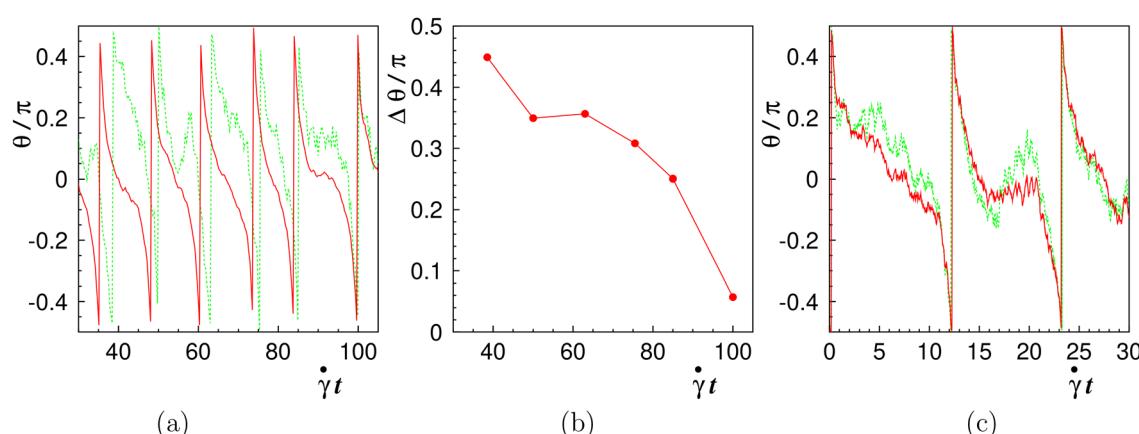


Fig. 5 (a) Time evolution of the inclination angles θ_{ext} (red continuous line) and θ_{int} (green dashed line) of a compound vesicle with $\Delta_{\text{ext}} = 1.23$, $\Delta_{\text{int}} = 0.74$, $\phi = 0.2$ at $(\mathcal{A}, S) = (5.38, 138)$ in desynchronized tumbling motion. (b) Time dependence of $\Delta\theta = \theta_{\text{int}} - \theta_{\text{ext}}$ corresponding to the data in panel (a). (c) Time evolution of the inclination angles θ_{ext} (red continuous line) and θ_{int} (green dashed line) of a compound vesicle with $\Delta_{\text{ext}} = \Delta_{\text{int}} = 0.74$, $\phi = 0.6$ at $(\mathcal{A}, S) = (4.18, 23)$ in synchronized tumbling motion.



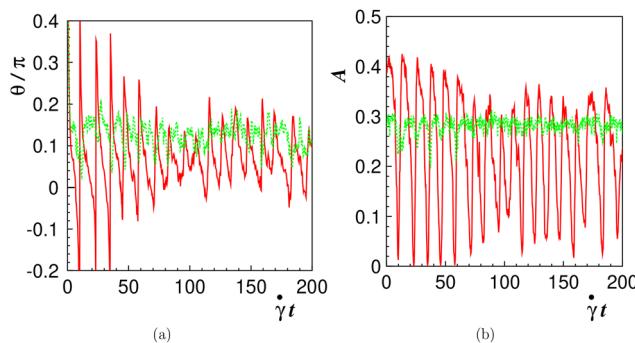


Fig. 6 (a) Time evolution of the inclination angles θ_{ext} (red continuous line) and θ_{int} (green dashed line) of a compound vesicle with $\Delta_{\text{ext}} = 1.23$, $\Delta_{\text{int}} = 0.74$, $\phi = 0.6$ at $(A, S) = (5.38, 138)$ showing undulating (external vesicle) and swinging (internal vesicle) motion. (b) Time evolution of the asphericities A_{ext} (red continuous line) and A_{int} (green dashed line) corresponding to the case in panel (a).

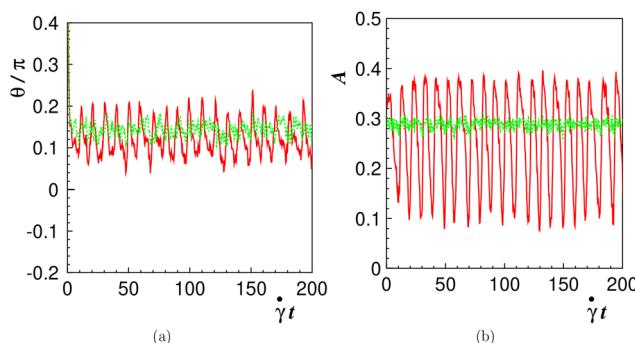


Fig. 7 (a) Time evolution of the inclination angles θ_{ext} (red continuous line) and θ_{int} (green dashed line) of a compound vesicle with $\Delta_{\text{ext}} = 1.23$, $\Delta_{\text{int}} = 0.74$, $\phi = 0.8$ at $(A, S) = (5.38, 138)$ showing undulating (external vesicle) and tank-treading (internal vesicle) motion. (b) Time evolution of the asphericities A_{ext} (red continuous line) and A_{int} (green dashed line) corresponding to the case in panel (a).

the time series $A_{\text{ext}}(t)$. Interestingly, it is found that $\omega_A \simeq 2\langle\omega\rangle_{\text{ext}}$ where $\langle\omega\rangle_{\text{ext}}$ is the average rotation frequency of the external vesicle. This result corresponds to the fact that the external vesicle experiences stretching along the $\pi/4$ direction and compression along the $-\pi/4$ direction, due to the elongational component of the shear flow,⁴⁷ alternating strong shape fluctuations and relaxation into more circular shapes. This mechanism occurs twice during every rotation cycle of the membrane, as can also be appreciated in the supplementary movie files movie2.mp4 and movie3.mp4. In order to characterize the UND motion, it may be useful to plot the asphericity as a function of the inclination angle, as reported in Fig. 9. When $\phi = 0.6$, the trajectories for the external vesicle are symmetric with respect to a vertical axis at a value $\langle\theta\rangle_{\text{ext}}^{\text{a}} > 0$. For clarity, we recall that in the TR case, the pattern of trajectories is similar but centered at $\langle\theta\rangle_{\text{ext}}^{\text{a}} \simeq 0$. The internal vesicle undergoes SW with large variations of the inclination angle while the asphericity weakly varies. By increasing the filling fraction, the trajectories of the external vesicle assume the characteristic shape of a closed ring, elongated along the vertical direction, while the pattern of the internal vesicle is typical of TT

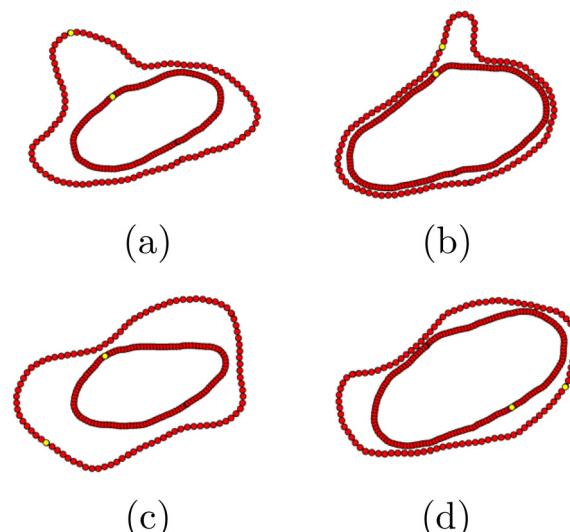


Fig. 8 Configurations of compound vesicles with $\Delta_{\text{ext}} = 1.23$, $\Delta_{\text{int}} = 0.74$, $\phi = 0.6$ (a), 0.8 (b) at $(A, S) = (5.38, 138)$ and with $\Delta_{\text{ext}} = \Delta_{\text{int}} = 0.74$, $\phi = 0.6$ (c), 0.8 (d) at $(A, S) = (4.18, 229)$.

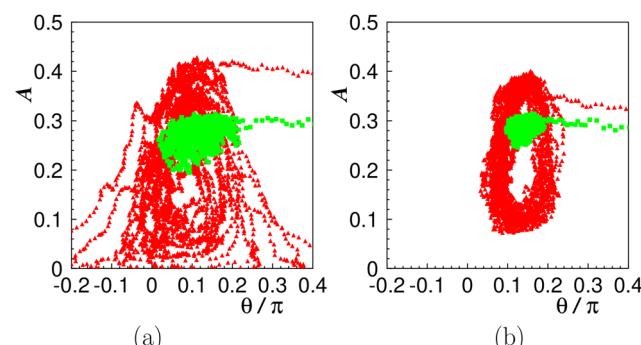


Fig. 9 Asphericity A as a function of the inclination angle θ for the external (red triangles) and internal (green squares) vesicles corresponding in panels (a) and (b) to the cases shown in Fig. 6 and 7, respectively.

motion. The relative positions of the centers of mass form closed fluctuating trajectories whose amplitudes decrease with ϕ , as illustrated in Fig. 10 for the cases with $\Delta_{\text{ext}} = 1.23$, $\Delta_{\text{int}} = 0.74$, $\phi = 0.6$, 0.8 at $(A, S) = (5.38, 138)$ of Fig. 6 and 7.

Lastly, when the external and internal vesicle have the same swelling degree, we still observe UND motion for the external vesicle with the formation of more lobes. The motion of the internal vesicle is again controlled by the filling fraction resulting in SW at $\phi = 0.6$ and TT at $\phi = 0.8$. However, the amplitudes in the oscillations of $\theta_{\text{ext}}(t)$ and $A_{\text{ext}}(t)$ are reduced with respect to the cases having $\Delta_{\text{ext}} > \Delta_{\text{int}}$. Typical configurations of vesicles with $\Delta_{\text{ext}} = \Delta_{\text{int}} = 0.74$, $\phi = 0.6$, 0.8 at $(A, S) = (4.18, 229)$ are shown in Fig. 8(c) and (d), respectively. The presence of UND motion with four rotating lobes in the external vesicle has been also observed in previous simulations with $\Delta_{\text{ext}} = \Delta_{\text{int}}$ at $\phi = 0.85$ and large shear rates.³ In that case, however, due to the lack of thermal fluctuations, external and internal vesicles show radial symmetry, their centers are

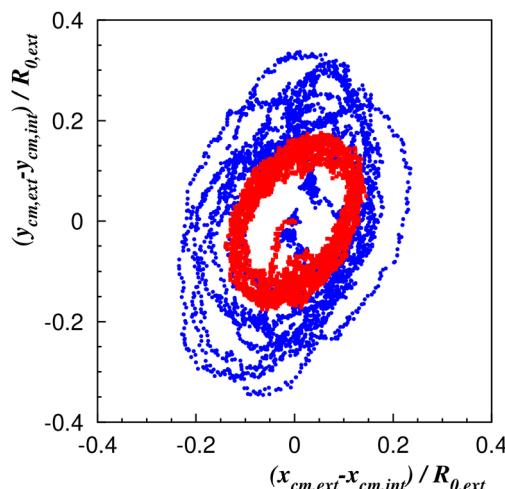


Fig. 10 Probability distribution of the center-of-mass relative positions for a compound vesicle corresponding to the cases shown in Fig. 6 (blue circles) and in Fig. 7 (red squares).

stationary, and the internal vesicle performs SW motion, differently from the present work. Finally, we add that the UND regime has never been observed in experiments, probably due to the difficulty of accessing high values of ϕ .¹³

4. Conclusions

We have studied the dynamics of a compound vesicle in a wall-bounded shear flow. Encapsulating a vesicle into an external one adds extra degrees of freedom with respect to the case of a single vesicle, due to internal vesicle rotation, displacement, deformation, and to the hydrodynamic coupling between the internal and external membranes. Our outcomes have been obtained by using mesoscale hydrodynamic simulations of a two-dimensional model system, including thermal fluctuations. This latter feature, lacking in all previous analytical and numerical approaches, has revealed to be fundamental in order to observe characteristic trembling motion, as pointed out in experiments.¹³ By varying the flow parameters Λ and S , the filling fraction ϕ , and the internal excess length Λ_{int} , the system is shown to manifest a rich phenomenology in the dynamical behavior. At values $\Lambda \lesssim 1$, vesicles exhibit tank-treading motion that is fully characterized. The external vesicle shows the same inclination angle of a single vesicle at fixed Λ and assumes a parallelogram-like shape, as quantified by the spectrum of shape amplitudes. The internal vesicle shows a larger inclination angle whose difference with respect to the external one does not depend on Λ and ϕ . By increasing Λ slightly beyond the value Λ_c of the TT-TR transition, the external vesicle manifests trembling motion while the internal one swings. Higher values of Λ trigger the transition to the tumbling motion of both vesicles. At small filling fraction, the two vesicles appear desynchronized while larger values of ϕ cause the mutual synchronization of vesicles. For large values of Λ , S , and ϕ , the external vesicle shows undulating motion,³

characterized by periodic oscillations of inclination angle and asphericity while the membrane buckles. This dynamic state exhibits unique patterns in the $\Lambda-\theta$ space that are shown to depend on the relative size and the swelling degree of both vesicles. The present numerical model, though limited to consider a two-dimensional system, is able to provide quantitative agreement with the experimental results of ref. 13. This is a remarkable feature that was previously shown to hold in describing correctly the flow patterns of single vesicles in shear flow.³¹ In the present investigation, it was set $\eta_{\text{in}} = \eta_{\text{out}}$ and $\eta_{\text{an}} \geq \eta_{\text{out}}$ which implies that $\eta_{\text{in}} \leq \eta_{\text{an}}$. In the case of a single vesicle, this would correspond to a tank-treading condition. However, we find that the internal vesicle can be found in different dynamical states due to the complex overall dynamics. It would be interesting to consider which effects might be triggered by setting $\eta_{\text{in}} > \eta_{\text{an}}$. Finally, we comment about the presence of boundary walls. In our case, the degree of confinement, defined to be $\chi = 2R_{0,\text{ext}}/L_y$,²³ is 0.35. For this value of χ , the study of single vesicles with the same numerical model^{26,42,43} did not show relevant differences with respect to the analytical solutions of Keller and Skalak,¹⁴ corresponding to unbounded shear flow ($\chi = 0$). We do not expect that for the considered value of χ , walls can quantitatively affect the present results with respect to unbounded shear flows. Indeed, our results agree quantitatively with the ones of Ref. 13 whose experimental setup is less confined than ours, being $\chi \simeq 0.25$. It would be interesting to consider effects of strong confinement ($\chi > 0.4$) on the present phenomenology, as done for single vesicles.²³

In conclusion, our model has allowed us to explore a large part of the space of controlling parameters so to both obtain results that are fully consistent with experiments, and characterize the undulating motion³ providing a comprehensive study under different conditions. We hope that this latter result can stimulate further experimental investigations capable to consider systems with high filling fractions at large values of Λ and S . Moreover, we believe that the present approach can reveal very useful in order to address the dynamic behavior of compound vesicles in oscillatory shear flow⁴⁸ as well as of vesicles enclosing two or more vesicles, as in the case of vesosomes.¹⁰

Conflicts of interest

There are no conflicts to declare.

Data availability

All data that support the findings of this study are included within the article (and any supplementary files).

Acknowledgements

The work of AL was performed under the auspices of GNFM-INDAM.



Notes and references

1 G. Schmid-Schonbein, Y. Shih and S. Chien, *Blood*, 1980, **56**, 866–875.

2 S. K. Veerapaneni, Y.-N. Young, P. M. Vlahovska and J. Bławdziewicz, *Phys. Rev. Lett.*, 2011, **106**, 158103.

3 B. Kaoui, T. Krüger and J. Harting, *Soft Matter*, 2013, **9**, 8057–8061.

4 T.-W. Pan, S. Zhao, X. Niu and R. Glowinski, *J. Comput. Phys.*, 2015, **300**, 241–253.

5 K. P. Sinha and R. M. Thaokar, *Soft Matter*, 2019, **15**, 6994–7017.

6 R. A. Harouaka, M. Nisic and S.-Y. Zheng, *J. Lab. Autom.*, 2013, **18**, 455–468.

7 C. Alix-Panabières and K. Pantel, *Nat. Rev. Cancer*, 2014, **14**, 623.

8 R. Kumar, R. Chakrabarti and R. M. Thaokar, *Soft Matter*, 2024, **20**, 6995–7011.

9 S. A. Walker, M. T. Kennedy and J. A. Zasadzinski, *Nature*, 1997, **387**, 61.

10 E. T. Kisak, B. Coldren, C. A. Evans, C. Boyer and J. A. Zasadzinski, *Curr. Med. Chem.*, 2004, **11**, 199–219.

11 C. B. Giuliano, N. Cvjetan, J. Ayache and P. Walde, *Chem. SystemsChem*, 2021, **3**, e2000049.

12 I. M. Saadeldin, S. Ehab and J. Cho, *Biomarker Res.*, 2023, **11**, 77.

13 M. Levant and V. Steinberg, *Phys. Rev. Lett.*, 2014, **112**, 138106.

14 S. R. Keller and R. Skalak, *J. Fluid Mech.*, 1982, **120**, 27–47.

15 C. Misbah, *Phys. Rev. Lett.*, 2006, **96**, 028104.

16 V. V. Lebedev, K. S. Turitsyn and S. S. Vergeles, *Phys. Rev. Lett.*, 2007, **99**, 218101.

17 G. Danker, T. Biben, T. Podgorski, C. Verdier and C. Misbah, *Phys. Rev. E: Stat., Nonlinear, Soft Matter Phys.*, 2007, **76**, 041905.

18 R. Finken, A. Lamura, U. Seifert and G. Gompper, *Eur. Phys. J. E: Soft Matter Biol. Phys.*, 2008, **25**, 399.

19 H. Noguchi and G. Gompper, *Phys. Rev. Lett.*, 2004, **93**, 258102.

20 H. Noguchi and G. Gompper, *Phys. Rev. E*, 2005, **72**, 011901.

21 H. Noguchi and G. Gompper, *Phys. Rev. Lett.*, 2007, **98**, 128103.

22 S. Meßlinger, B. Schmidt, H. Noguchi and G. Gompper, *Phys. Rev. E*, 2009, **80**, 011901.

23 B. Kaoui, J. Harting and C. Misbah, *Phys. Rev. E*, 2011, **83**, 066319.

24 H. Zhao and E. S. G. Shaqfeh, *J. Fluid Mech.*, 2011, **674**, 578–604.

25 A. Nait-Ouhra, A. Guckenberger, A. Farutin, H. Ez-Zahraouy, A. Benyoussef, S. Gekle and C. Misbah, *Phys. Rev. Fluids*, 2018, **3**, 123601.

26 A. Lamura, *Mathematics*, 2022, **10**, 3570.

27 V. Kantsler and V. Steinberg, *Phys. Rev. Lett.*, 2005, **95**, 258101.

28 V. Kantsler and V. Steinberg, *Phys. Rev. Lett.*, 2006, **96**, 036001.

29 V. Vitkova, M.-A. Mader, B. Polack, C. Misbah and T. Podgorski, *Biophys. J.*, 2008, **95**, L33.

30 N. J. Zabusky, E. Segre, J. Deschamps, V. Kantsler and V. Steinberg, *Phys. Fluids*, 2011, **23**, 041905.

31 E. Afik, A. Lamura and V. Steinberg, *Europhys. Lett.*, 2016, **113**, 38003.

32 M. Levant and V. Steinberg, *Phys. Rev. Lett.*, 2012, **109**, 268103.

33 D. Abreu and U. Seifert, *Phys. Rev. Lett.*, 2013, **110**, 238103.

34 D. Abreu, M. Levant, V. Steinberg and U. Seifert, *Adv. Colloid Interface Sci.*, 2014, **208**, 129–141.

35 A. Malevanets and R. Kapral, *J. Chem. Phys.*, 1999, **110**, 8605–8613.

36 A. Malevanets and R. Kapral, *J. Chem. Phys.*, 2000, **112**, 7260–7269.

37 R. Kapral, *Multiparticle Collision Dynamics: Simulation of Complex Systems on Mesoscales*, John Wiley & Sons, Ltd, 2008, pp. 89–146.

38 G. Gompper, T. Ihle, D. M. Kroll and R. G. Winkler, in *Multiparticle Collision Dynamics: A Particle-Based Mesoscale Simulation Approach to the Hydrodynamics of Complex Fluids*, ed. C. Holm and K. Kremer, Springer Berlin Heidelberg, Berlin, Heidelberg, 2009, pp. 1–87.

39 M. P. Howard, A. Nikoubashman and J. C. Palmer, *Curr. Opin. Chem. Eng.*, 2019, **23**, 34–43.

40 A. Lamura, G. Gompper, T. Ihle and D. M. Kroll, *Europhys. Lett.*, 2001, **56**, 319.

41 A. Lamura and G. Gompper, *Eur. Phys. J. E: Soft Matter Biol. Phys.*, 2002, **9**, 477.

42 A. Lamura and G. Gompper, *Europhys. Lett.*, 2013, **102**, 28004.

43 A. Lamura and G. Gompper, *Proc. IUTAM*, 2015, **16**, 3–11.

44 B. Kaoui, T. Krüger and J. Harting, *Soft Matter*, 2012, **8**, 9246–9252.

45 V. Kantsler, E. Segre and V. Steinberg, *Phys. Rev. Lett.*, 2007, **99**, 178102.

46 M. Levant, D. Abreu, U. Seifert and V. Steinberg, *Europhys. Lett.*, 2014, **107**, 28001.

47 M. Abkarian and A. Viallat, *Soft Matter*, 2008, **4**, 653–657.

48 H. Noguchi, *J. Phys. Soc. Jpn.*, 2010, **79**, 024801.

