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## Eco-friendly repurposing of by-pass waste for optics and radiation protection: addressing hazardous material challenges

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This study delves into the investigation of the optical and gamma radiation shielding properties of glasses formulated with varying concentrations of by-pass (cement dust) combined with  $\text{Na}_2\text{O}$ ,  $\text{Fe}_2\text{O}_3$ ,  $\text{Bi}_2\text{O}_3$ , and  $\text{P}_2\text{O}_5$ . The specific compositions are represented as  $((x) \text{ by-pass}-(20) \text{ Na}_2\text{O}-(10) \text{ Fe}_2\text{O}_3-(20) \text{ Bi}_2\text{O}_3-(50-x) \text{ P}_2\text{O}_5)$ , with "x" ranging from 0 to 15 wt%. The structural and optical properties of these glasses were analyzed using X-ray diffraction (XRD), Fourier-transform infrared (FTIR) spectroscopy, and UV-visible-NIR spectroscopy techniques. The XRD results reveal the absence of Bragg peaks and highlight the amorphous nature of the prepared samples. Notably, increasing by-pass concentration enhanced UV-visible-NIR transmittance, particularly from 400 nm to 850 nm. Radiation shielding assessments were conducted across photon energies ranging from 80 to 2614 keV. The results highlighted that the linear and mass attenuation coefficients were inversely proportional to the photon energy. The by-pass15 glass, with the highest by-pass concentration, consistently demonstrated superior shielding attributes compared to its counterparts, with the lowest half-value layer ( $G_{\text{HVL}}$ ) values and the most substantial radiation protection efficiency (RPE). Comparative assessments with other glasses and concrete types further accentuated the potential of the by-pass15 glass as a promising candidate for gamma radiation shielding applications. The results obtained in this research pave the way for potentially utilizing by-pass15 glass in environments where radiation protection is imperative. Incorporating by-pass in glass composites enhances their gamma shielding efficiency and promotes the sustainable use of by-pass, underscoring the dual advantages of this innovative approach.

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## 1. Introduction

Environmental preservation is an imperative responsibility for the global community in the current era. The escalating concerns about ecological degradation, mainly driven by improper disposal of industrial wastes and the proliferation of harmful materials, have called for innovative strategies to transform

potential pollutants into beneficial products. This eco-conscious shift from waste generation to sustainable utilization is not just about conservation; it is about ushering in a new epoch of sustainable development.<sup>1,2</sup> One such promising avenue is the transformation of harmful industrial by-products into high-value materials with exceptional properties. Glass, an ancient material with diverse applications, presents a unique opportunity. Traditionally valued for its transparency and malleability, recent advancements have expanded its utility to optics and radiation shielding realms. Modern glass formulations doped with specific elements or compounds can exhibit remarkable optical characteristics, making them indispensable in advanced applications ranging from optoelectronics to photonics.<sup>3-6</sup>

Radioactive isotopes and related technologies significantly influence modern human life. Naturally occurring isotopes serve as a source of radiation and have various applications. Their detectable radioactivity makes them ideal tracers. They can also eliminate harmful agents from food due to the energy they emit. Moreover, isotopes have diverse medical roles, such as using radioactive  $^{131}\text{I}$  to assess thyroid functionality.<sup>7,8</sup>

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Yet, unintended radiation exposure from these isotopes can harm living organisms and the environment. The potential harm depends on the type and amount of energy they release. For example, high doses can lead to conditions like cataracts, DNA damage, infertility, and several diseases.<sup>9</sup> Thus, those handling radioactive materials are advised to use shielding materials to minimize exposure risks.<sup>10–12</sup> Historically, due to their high atomic numbers and density, materials like lead, concrete, and lead composite have been favored for nuclear safety applications. However, concerns like toxicity have limited the use of lead products.<sup>13</sup> This has prompted the search for alternative materials. In this quest, glass systems emerge as potential candidates. They offer advantages over other materials, such as cost-effectiveness, ease of fabrication, recyclability, and light transparency.<sup>14</sup> Specifically, glasses infused with heavy metal oxides, like  $\text{Bi}_2\text{O}_3$ , are gaining traction. Their high density and unique features make them apt for radiation shielding tasks.<sup>15–18</sup>

In recent years, exploring silicate structural, optical, and radiation shielding properties of studies such as the innovative doping of soda-lime-silica glasses with organic compounds like peanut shell powder have opened new avenues for modifying optical properties while promoting sustainability.<sup>19</sup> Further investigations, as documented,<sup>20,21</sup> have delved into the complex interplay between glass composition and its impact on light transmission and radiation protection efficiency. Recent work highlighted<sup>22</sup> has also contributed to our understanding of the structural characteristics crucial for developing non-toxic, eco-friendly radiation shielding materials. These studies underscore the importance of continuous innovation in glass science, particularly in leveraging various dopants and waste materials, to achieve enhanced performance in optical applications and radiation shielding.<sup>23–26</sup>

This research examines the usefulness of prepared glasses in shielding applications at different photon energies. The primary focus of the current study is to experimentally investigate the radiation attenuation characteristics of by-pass (cement dust) glasses, which comprise  $\text{SiO}_2$ ,  $\text{Al}_2\text{O}_3$ ,  $\text{Fe}_2\text{O}_3$ ,  $\text{TiO}_2$ ,  $\text{CaO}$ ,  $\text{MgO}$ ,  $\text{Na}_2\text{O}$ , and  $\text{K}_2\text{O}$ . For gamma, the effect of systematically replacing by-pass with  $\text{P}_2\text{O}_5$  on the attenuation characteristics was dissected in great depth. This study also provides a wide-scope comparison of the shielding performance of the current recommended system with that of the typically employed systems. The selection of by-pass waste as a primary material in the present study was driven by multiple factors. First and foremost, by-pass waste, a by-product of the cement industry, poses significant environmental challenges due to its large volume generation and disposal issues.<sup>22,27,28</sup> Its composition, rich in silica and alumina, alongside other oxides such as calcium oxide and iron oxide, makes it a potential candidate for glass synthesis, leveraging its waste-to-wealth conversion potential.<sup>29,30</sup> Furthermore, incorporating by-pass waste into glass composites aligns with global efforts towards sustainable development by reducing landfill use and promoting the recycling of industrial by-products.<sup>31</sup> Additionally, previous studies have indicated that certain waste materials can enhance the optical and radiation-shielding properties of glasses.

Therefore, by-pass waste was selected to investigate its efficacy in improving these properties while addressing environmental conservation and waste management challenges. Beyond optics, there is an emergent demand for materials that can shield against harmful radiation, especially in medical, industrial, and nuclear domains. With the growing reliance on nuclear energy and medical radiation, ensuring safety has become paramount. Herein lies the novelty of re-engineering harmful by-products into specialized glass composites that can attenuate harmful radiations. Such advancements address the pressing issue of waste management and provide solutions to critical challenges in radiation protection.

In essence, the convergence of environmental protection with the creation of high-performance materials represents a pioneering approach. It signifies the dawning of an era where harmful residues are no longer discarded but are transformed, through the alchemy of science and innovation, into materials that protect, serve, and inspire. This research underscores the importance and potential of such transformations, highlighting the dual benefits: safeguarding our environment while meeting the demands of cutting-edge technological applications.

## 2. Materials and methods

### 2.1. Sample preparation

Cement dust, characterized as by-pass, was sourced from cement factories located in Beni Suef (Egypt). Recognized as industrial waste, this waste (dust) was first analyzed using X-ray fluorescence (XRF) to determine its chemical composition. The predominant elements in the by-pass included 47.67%  $\text{SiO}_2$ , 4.24%  $\text{Al}_2\text{O}_3$ , 2.14%  $\text{Fe}_2\text{O}_3$ , 0.20%  $\text{TiO}_2$ , 37.95%  $\text{CaO}$ , 3%  $\text{MgO}$ , 2.95%  $\text{Na}_2\text{O}$ , and 1.85%  $\text{K}_2\text{O}$ , cumulatively accounting for 100% of its total oxide content.

By-pass and transparent glasses were synthesized using a specific composition:  $((x) \text{ by-pass}-(20) \text{ Na}_2\text{O}-(10) \text{ Fe}_2\text{O}_3-(20) \text{ Bi}_2\text{O}_3-(50-x) \text{ P}_2\text{O}_5)$ , where “ $x$ ” varied as 0 (by-pass0), 3 (by-pass3), 6 (by-pass6), 10 (by-pass10), and 15 (by-pass15) wt%. The glass samples were created through the conventional melt-quenching technique. Five distinct concentrations of the by-pass mixture were initially prepared according to the given composition formula, with “ $x$ ” ranging from 0 to 15 wt%. The mixtures were continuously stirred at room temperature to achieve homogeneity. Subsequently, the samples were heated to 1450 °C and melted for two hours. The molten samples were poured onto a stainless-steel bar to form the desired shapes. This was followed by a transfer to a secondary furnace set at 300 °C for annealing. The samples were gradually cooled to room temperature, ensuring their structural integrity and amorphous nature.

### 2.2. Characterizations and measurements

The structural and optical characteristics of the prepared samples were meticulously examined through X-ray diffraction (XRD) using an AXS D8 Advance apparatus (Bruker, Germany), utilizing  $\text{Cu-K}_\alpha$  radiation at a wavelength of 1.54 Å ( $\lambda$ ), within



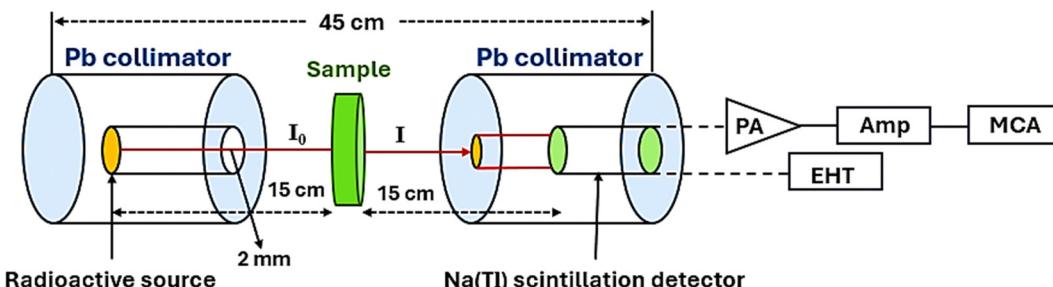


Fig. 1 Experimental radiation measurement setup.

the  $2\theta$  range of  $10^\circ$  to  $89.6^\circ$ . In parallel, Fourier-transform infrared (FTIR) spectroscopy was performed using a Bruker Tensor 27 system to map the vibrational spectra of the prepared samples across the spectral range of  $1700\text{ cm}^{-1}$  to  $450\text{ cm}^{-1}$ . Furthermore, the optical transmission properties were rigorously quantified using a Shimadzu UV-2450 (UV-visible-NIR) spectrophotometer, covering a spectral domain from 200 to 2500 nm. The narrow beam approach, combined with a lead-collimator, was utilized to acquire experimentally established linear attenuation coefficients ( $G_{\text{LAC}}$ ). The values were acquired using a NaI(Tl) scintillation detector with a measurement area of 3 inches by 3 inches coupled to a multichannel analyzer. Throughout the experiment, the following radioactive sources were utilized:  $^{133}\text{Ba}$  (5  $\mu\text{Ci}$ , 80 keV),  $^{137}\text{Cs}$  (5  $\mu\text{Ci}$ , 662 keV),  $^{60}\text{Co}$  (10  $\mu\text{Ci}$ , 1173 and 1333 keV), and  $^{232}\text{Th}$  (20  $\mu\text{Ci}$ , 238, 911, and 2614 keV). Fig. 1 illustrates the experimental setup, which includes the source, the sample, and the detector in their appropriate locations. Calculations were performed to estimate the photon intensity without and with the absorber for each gamma line using the region underneath the photopeak. These calculations were essential to determine the photon intensity. The process was repeated five times, and 10 minutes was set aside for each iteration of the method. The level of uncertainty was far lower than one percent.

### 3. Results and discussion

#### 3.1. XRD and FTIR analyses

The XRD patterns of the prepared by-pass glass samples exhibited a remarkable similarity, characterized by the absence of Bragg peaks, as illustrated in Fig. 2(a). This absence of peaks confirms the amorphous nature of the glassy phase of the samples, where the amorphous regions would contribute a broad, low-intensity background. Additionally, no discernible peak corresponding to by-pass was observed, which reveals that the added quantity of by-pass has been fully incorporated into the glass matrix.<sup>32,33</sup>

The FTIR spectra of by-pass cement dust incorporated into phosphate glass matrices reveal significant structural modifications with varying by-pass contents, as shown in Fig. 2(b). The spectra, ranging from  $600\text{ cm}^{-1}$  to  $1600\text{ cm}^{-1}$ , display distinct absorbance bands associated with Si-O-Si, P-O-P, and P-O vibrational modes. As the by-pass concentration increases from 0% to 15%, notable shifts and intensity changes in these bands

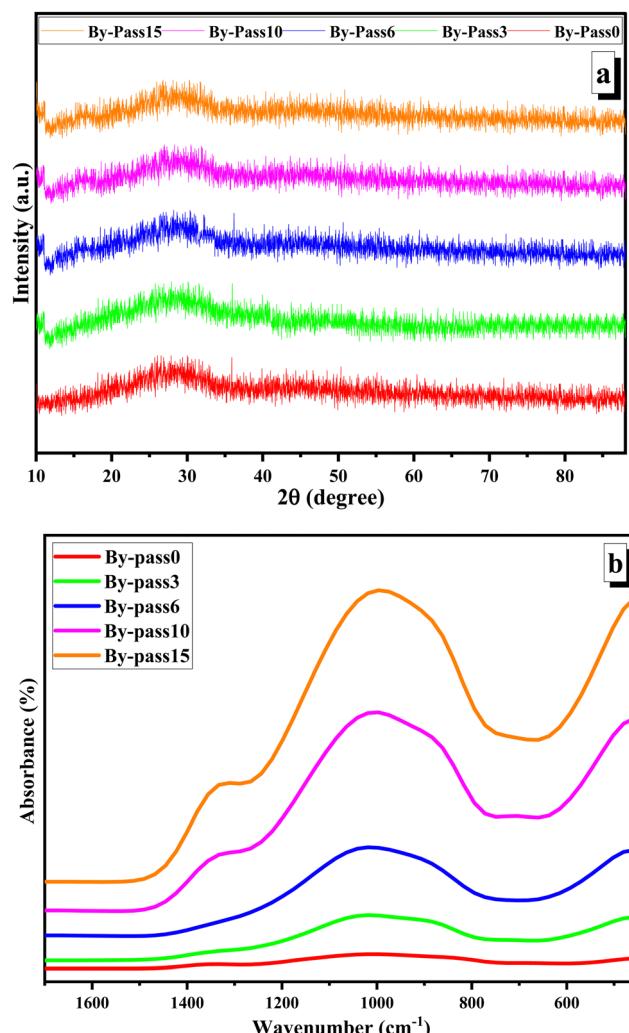


Fig. 2 (a) XRD patterns and (b) FTIR spectra for different by-pass concentrations.

occur, which indicate progressive integration of cement dust components into the glass network. These modifications suggest the formation of new structural units or phases, particularly due to incorporating elements such as CaO, MgO, and  $\text{Al}_2\text{O}_3$ .<sup>22</sup> The observed peaks between  $800$  and  $1300\text{ cm}^{-1}$  may be due to P-O-P and Si-O-Si asymmetric stretching.<sup>30</sup> A peak around  $1333\text{ cm}^{-1}$  belongs to the asymmetric stretching of



carbonate groups,<sup>34</sup> which only appears in by-pass10 and by-pass15, which is well-matched with a previous study.<sup>31</sup> The systematic spectral variations highlight the effectiveness of the melt-quenching technique in achieving homogeneity and structural integrity in the synthesized glasses. Understanding these changes provides insights into developing advanced glass materials with tailored properties, potentially enhancing their waste management and construction applications.

### 3.2. Optical properties

**3.2.1. UV-visible-NIR spectroscopy analysis.** The UV-visible-NIR spectroscopy analysis, conducted on samples with a thickness of 0.43 cm, is illustrated in Fig. 3(a). Analysis of the UV-visible-NIR transmittance ( $T$ ) spectra revealed a significant increase in transmittance values, specifically between 400 nm and 850 nm, as the concentration of by-pass in the composite progressively increases. This observation strongly supports the notion that including a by-pass markedly enhances the transmittance of the composite. This enhancement is attributed to the incident photon beam possessing sufficient energy to facilitate

electron transfer from the valence band to the conduction band, with the remainder of the photon energy being absorbed by the samples. Notably, this effect is most pronounced in the by-pass15 sample, which boasts the highest by-pass concentration and consistently exhibits superior transmittance compared to the other samples. However, as we shift to higher wavelengths, the peak positions remain remarkably consistent across all concentrations, signifying the unaltered nature of electronic transitions.

Nevertheless, there is a slight variation in the intensity of these peaks associated with changing concentration. As we approach the wavelength range of 2100 nm to 2500 nm, the incident photon beam carries lower energy, reducing interactions with atomic structures within the composite. This suggests that by-pass plays a crucial role in enhancing the light–material interaction in the composite. Consequently, the composite with the highest by-pass concentration exhibits the strongest interactions with UV-visible-NIR light among all the samples examined, aligning well with the results reported in a previous study.<sup>22</sup>

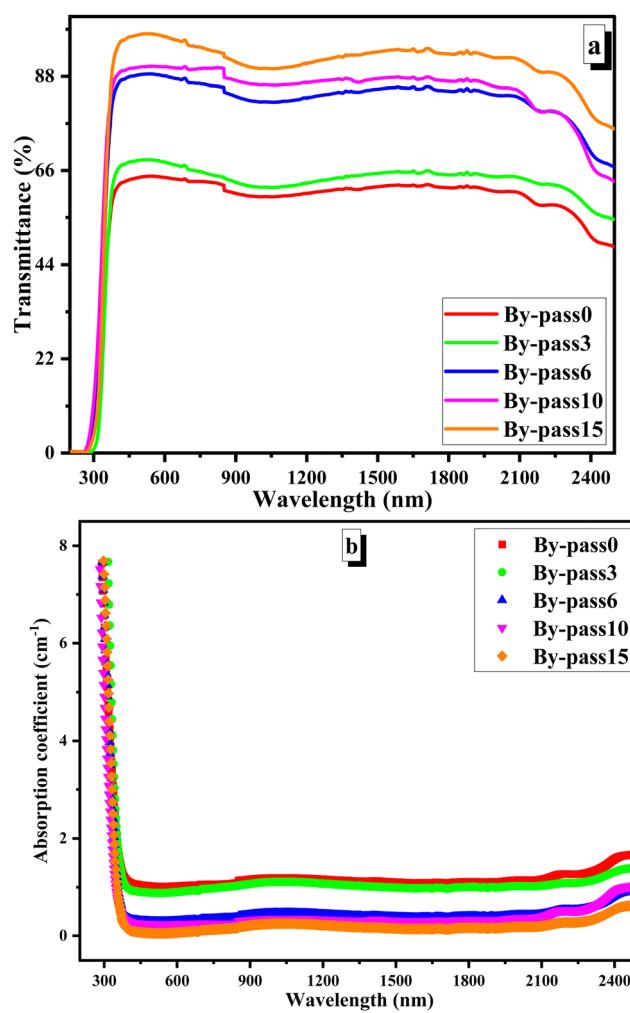


Fig. 3 (a) UV-visible-NIR transmittance profiles and (b) the absorption coefficient of the prepared samples.

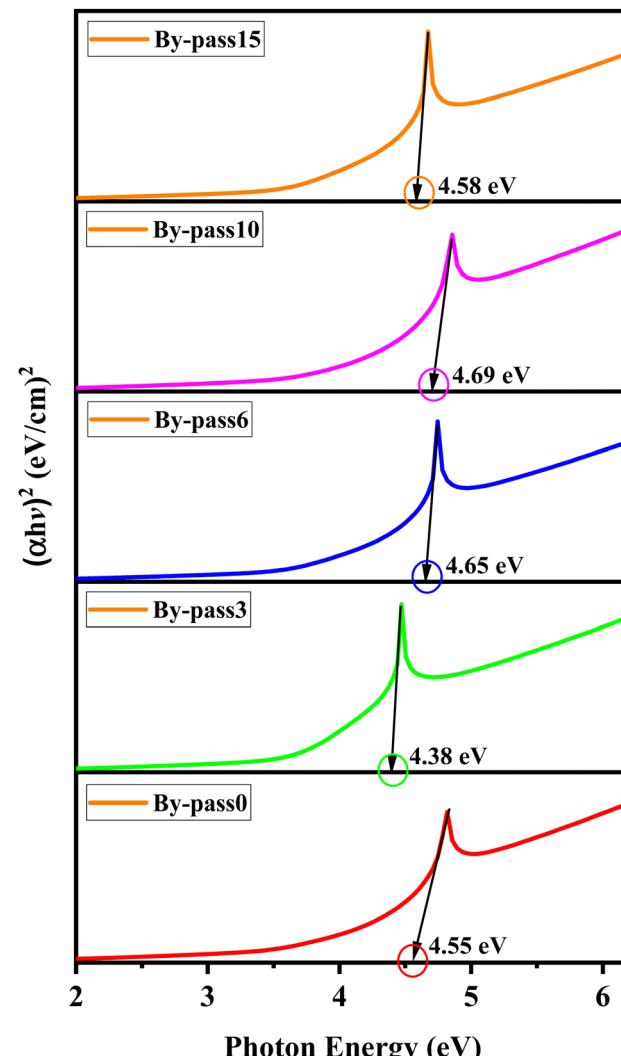


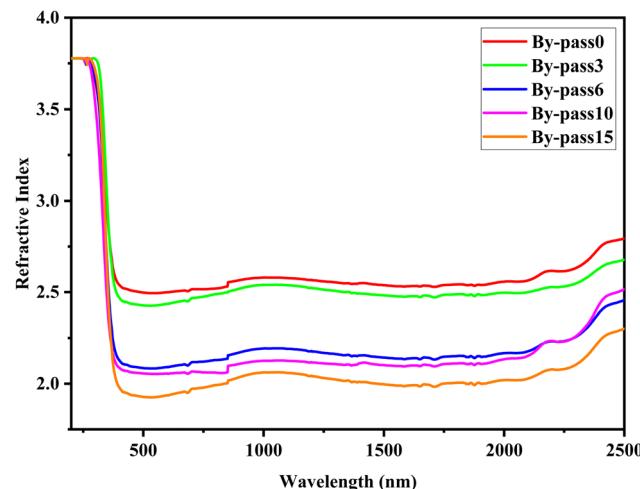
Fig. 4 Estimated direct band gap for different by-pass concentrations.

**3.2.2. Bandgap energy.** Determining the band gap of a material from its UV-visible-NIR transmission data involves plotting the transformed data to reveal the band gap as an intercept on the  $x$ -axis.

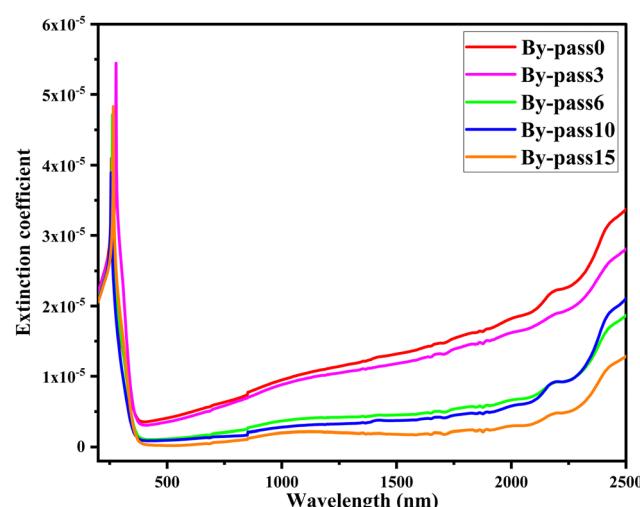
The Tauc plot is commonly used for this purpose. First, we need to approximate the absorption coefficient ( $\alpha$ ) from the

**Table 1** The direct band gaps for the composite at different concentrations of by-pass ( $x$ )

$x$ (wt%)	Direct band gap (eV)
0	4.55
3	4.38
6	4.65
10	4.69
15	4.58



**Fig. 5** The refractive index against the wavelength for different concentrations of by-pass in the studied samples.

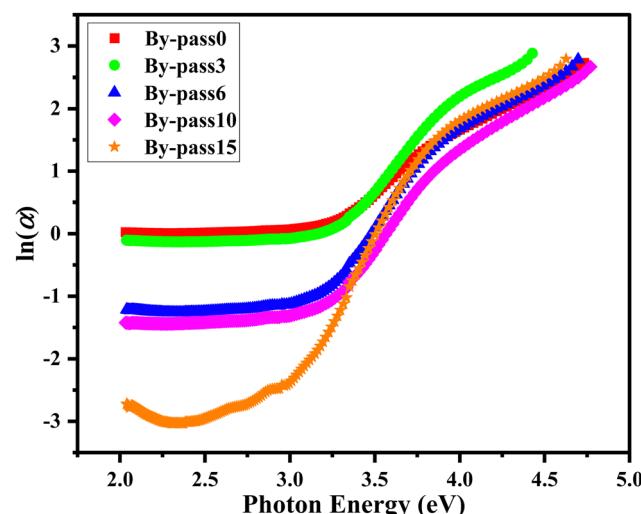


**Fig. 6** The extinction coefficient against wavelength for different concentrations of by-pass in the studied samples.

transmittance ( $T$ ) data using eqn (1):<sup>35</sup>

$$\alpha = -\left(\frac{\ln T}{t}\right) \quad (1)$$

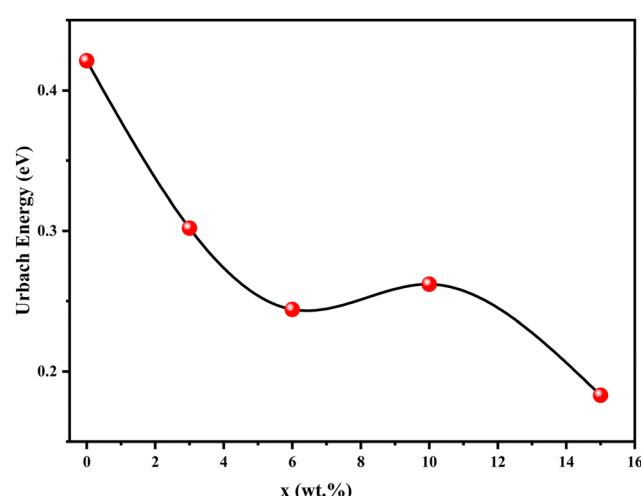
where  $t$  is the thickness of the glass sample (cm), which in our work is  $t = 0.43$  cm. As shown in Fig. 3(b), the absorption



**Fig. 7** The relationship between  $\ln(\alpha)$  and photon energy for different concentrations of by-pass in the studied samples.

**Table 2** The Urbach energies ( $E_U$ ) for the composite with different concentrations of by-pass ( $x$ )

$x$ (wt%)	$E_U$ (eV)
0	~ 0.421
3	~ 0.302
6	~ 0.244
10	~ 0.262
15	~ 0.183



**Fig. 8** The Urbach energy for different concentrations of by-pass ( $x$ ) in the studied samples.



coefficient decreases sharply in the UV region but stabilizes in the visible-NIR areas, making this material ideal for optoelectronic applications.

The Tauc relationship is given by eqn (2):<sup>36-38</sup>

$$\alpha h\nu = A(h\nu - E_g)^n \quad (2)$$

where  $h\nu$  is the photon energy (which can be calculated as  $(1240/\lambda)$ ), where 1240 is a constant in eV and  $\lambda$  is the wavelength in nm,  $E_g$  is the optical band gap,  $A$  is a constant, and  $n = 2$  for indirect transitions and  $n = 1/2$  for direct transitions.

To determine the direct transition band gap, we can plot  $(\alpha h\nu)^2$  vs.  $(h\nu)$ . The direct optical band gap values are

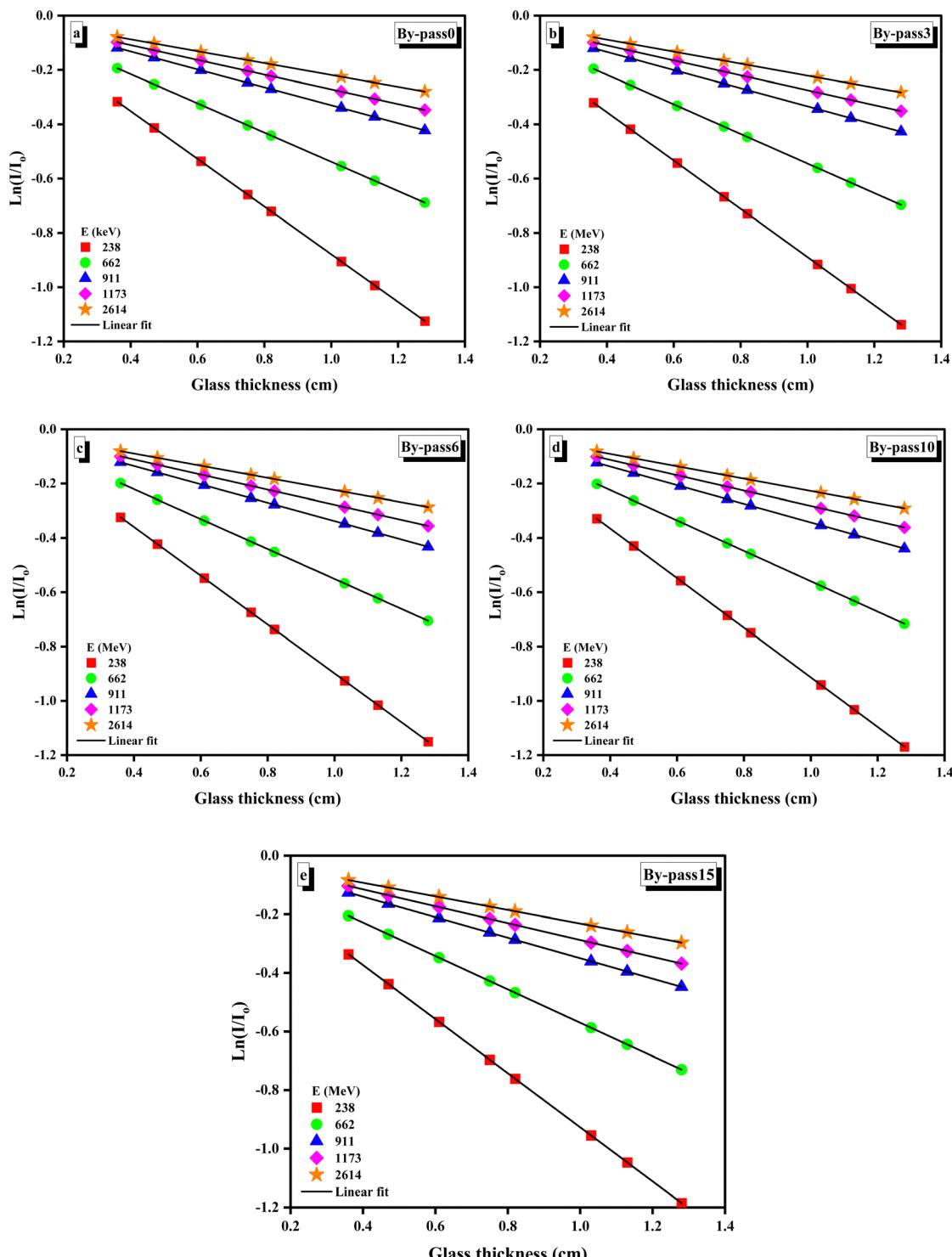


Fig. 9 (a)–(e) Transmission factor ( $\ln(I/I_0)$ ) against glass thickness at a selected photon energy.

determined by identifying the intercepts where the linear portion of the curve extrapolates to meet the linear axis representing photon energy. Fig. 4 shows that as the by-pass concentration increases from 0 to 15 wt%, the optical band gap energy changes, and the values are tabulated in Table 1. This increase implies that adding by-pass atoms to glass fabricates new energy levels in the band gap. These new energy levels act as traps for valence electrons, making the transition to high energy states in the conduction band more difficult.<sup>39</sup>

**3.2.3. Refractive index.** The refractive index ( $n$ ) is a fundamental property of materials, indicating how much light is refracted as it passes through the material, which can be determined from the transmittance ( $T$ ) data. The relationship of the refractive index is given by eqn (3):<sup>40</sup>

$$n = \left( \frac{1}{T-1} \right)^{\frac{1}{2}} + \frac{1}{T} \quad (3)$$

Fig. 5 illustrates the variation of the refractive index ( $n$ ) as a function of wavelength. As the wavelength increases, “ $n$ ” exhibits a general decreasing trend, which can be attributed to increased transmittance and reduced absorption due to traps and intraband transitions across varying concentrations of by-pass.

**3.2.4. Extinction coefficient.** The extinction coefficient, often represented by  $k$ , measures the amount of absorption loss when light passes through a material. It is closely related to the absorption coefficient ( $\alpha$ ) and provides insights into how efficiently a material can absorb light at a specific wavelength or photon energy. The relationship between “ $\alpha$ ” and “ $k$ ” is given by eqn (4):<sup>41</sup>

$$k = \frac{\alpha\lambda}{4\pi} \quad (4)$$

where  $\alpha$  is the absorption coefficient, which has been previously determined, and  $\lambda$  is the wavelength of light. Using eqn (4), the “ $k$ ” value can be computed for each wavelength and different concentrations of by-pass in the composite. Fig. 6 shows “ $k$ ” as a function of wavelength and by-pass concentration. The “ $k$ ” value increases with wavelength, indicating that the glass absorbs and scatters more light at longer wavelengths. These results are well-matched with those obtained in previous studies.<sup>22,42</sup>

**3.2.5. Urbach energy.** Urbach energy ( $E_U$ ) measures the width of the band tail states in disordered materials, which arise due to disorder and local potential fluctuations, leading to a non-zero density of states in the energy gap of a semiconductor. This parameter provides insights into the degree of disorder in the material. The “ $\alpha$ ” near the band edge of disordered materials can be described by the Urbach rule (eqn (5)):<sup>43</sup>

$$\alpha(E) = \alpha_0 \exp\left(\frac{E}{E_U}\right) \quad (5)$$

where  $\alpha(E)$  is the absorption coefficient at photon energy ( $E$ ),  $\alpha_0$  is a constant, and  $E_U$  can be determined by plotting  $\ln(\alpha)$  versus  $E$  and extracting the reciprocal of the slope, as shown in Fig. 7. A steeper slope indicates greater disorder in the material. The  $E_U$  values for the different concentrations of by-pass in the

composite, extracted from the slope of the linear fit to  $\ln(\alpha)$  vs. photon energy, are tabulated in Table 2.

As shown in Fig. 8 and Table 2, the  $E_U$  has a notable trend observed where the  $E_U$  decreases from 0.421 eV at 0 wt% to 0.183 eV at 15 wt%. This indicates a significant reduction in disorder as the concentration of the by-pass component increases up to 15 wt%. The initial high value of  $E_U$  at 0 wt% suggests a high level of disorder in the pure composite material, likely due to inherent structural imperfections. The slight increase at 10 wt% may indicate a point of optimal concentration where additional by-pass does not further reduce disorder but instead introduces new forms of structural irregularities. As the by-pass concentration increases, these imperfections are gradually mitigated, leading to a more ordered structure and thus a lower  $E_U$ .<sup>44,45</sup>

The reduction in  $E_U$  with increasing by-pass concentration suggests an improvement in the structural order of the

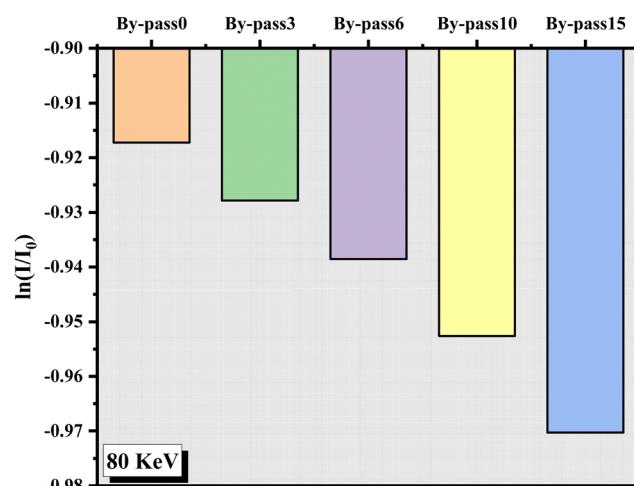


Fig. 10  $\ln(I/I_0)$  values against glass concentration at 80 keV.

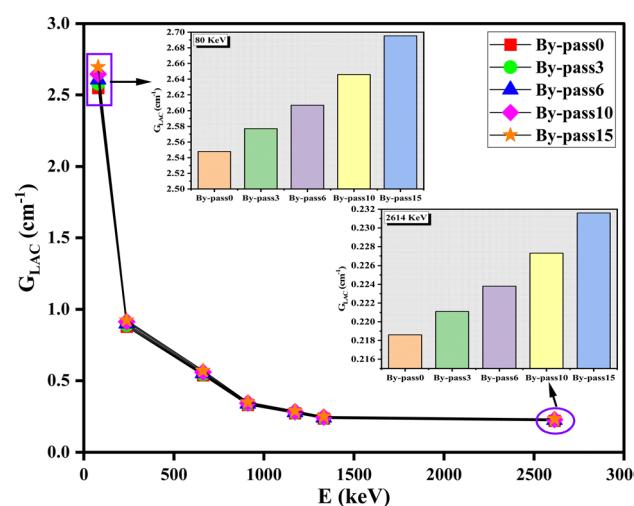


Fig. 11 Linear attenuation coefficient ( $G_{LAC}$ ) against photon energy for all studied samples at a selected photon energy.



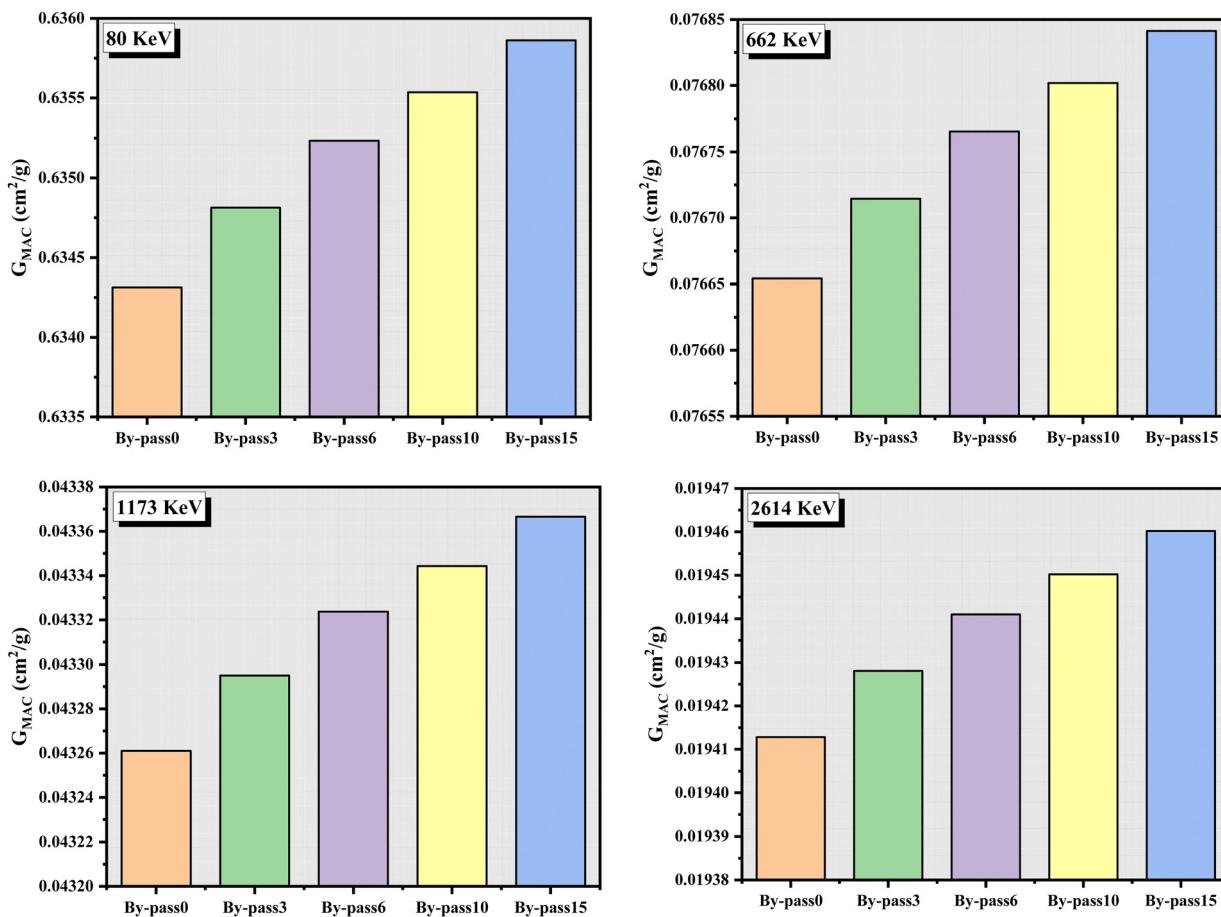


Fig. 12 Mass attenuation coefficient ( $G_{MAC}$ ) against photon energy for all studied glasses at a selected photon energy.

Table 3 Mass attenuation coefficient for all glasses at selected photon energies

Sample	80 keV	238 keV	662 keV	911 keV	1173 keV	1333 keV	2614 keV
By-pass0	0.6343	0.2132	0.0767	0.0557	0.0433	0.0381	0.0194
By-pass3	0.6348	0.2134	0.0767	0.0557	0.0433	0.0381	0.0194
By-pass6	0.6352	0.2135	0.0768	0.0558	0.0433	0.0381	0.0194
By-pass10	0.6355	0.2136	0.0768	0.0558	0.0433	0.0381	0.0195
By-pass15	0.6359	0.2137	0.0768	0.0558	0.0434	0.0382	0.0195

composite material. This can have significant implications for applications that require high optical clarity and minimal electronic defects, such as photovoltaic devices, optical sensors, and transparent conductors. Future work should aim to understand the precise mechanisms by which the by-pass component mitigates disorder and explore the potential to optimize the concentration further to achieve the lowest possible  $E_U$ .

### 3.3. Radiation shielding properties

Fig. 9(a)–(e) shows a plot of  $\ln(I/I_0)$  as a function of glass thickness for all the glasses studied at various energies. According to what is shown in this figure, the  $\ln(I/I_0)$  values for by-pass0, by-pass3, by-pass6, by-pass10, and by-pass15 increase when the

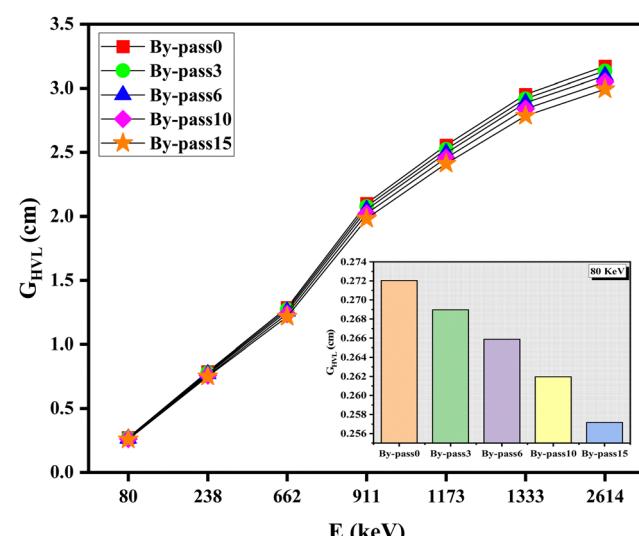


Fig. 13 Half value layer ( $G_{HVL}$ ) against photon energy for all studied samples.

photon energy increases. In addition, at the studied energies,  $\ln(I/I_0)$  values decrease with increasing by-pass.

For example, at 662 keV  $[\ln(I/I_0)]_{\text{by-pass}0} > [\ln(I/I_0)]_{\text{by-pass}3} > [\ln(I/I_0)]_{\text{by-pass}6} > [\ln(I/I_0)]_{\text{by-pass}10} > [\ln(I/I_0)]_{\text{by-pass}15}$ . At a certain energy, for example, 80 keV, the  $\ln(I/I_0)$  values for all glasses decrease as a function of glass compression. The  $\ln(I/I_0)$  values are  $-0.92$ ,  $-0.93$ ,  $-0.94$ ,  $-0.95$ , and  $-0.97$  for by-pass0, by-pass3, by-pass6, by-pass10 and by-pass15 glasses, respectively (Fig. 10). This is because of the shift in density ( $D$ ) that occurs while going from by-pass0 to by-pass3 to by-pass6 to by-pass10 to by-pass15. Moreover, the crystalline phases generated in by-pass15 are denser than those created in by-pass0, making the network structure more compact. The behavior of  $\ln(I/I_0)$  with energy and glass composition also indicates that the photon attenuation lowers with increasing photon energy and increases with altering the amount of by-pass content. The slope between  $\ln(I/I_0)$  and glass thickness has been used to measure the  $G_{\text{LAC}}$  values for by-pass0, by-pass3, by-pass6, by-pass10, and by-pass15 glass samples at different energies.

The  $G_{\text{LAC}}$  values for each studied glass are shown in Fig. 11 at the chosen photon energy. As can be seen from this figure, the  $G_{\text{LAC}}$  values for all glass ceramics decrease as the photon energy increases. This might be because there are three different ways in which photons can impart energy to matter. Glasses interact with gamma rays through the photoelectric effect, Compton scattering, and pair creation at low, medium, and high photon energies.<sup>46–49</sup> When measured at a specific energy (inset figures for 80 and 2614 keV, Fig. 11), the  $G_{\text{LAC}}$  values increase as the percentage of by-pass increases. The glass with the smallest  $G_{\text{LAC}}$  value is by-pass0, while the glass with the largest is by-pass15, which is well-matched with those obtained in a previous study.<sup>50</sup> The mass attenuation coefficient ( $G_{\text{MAC}}$ ) results for all glasses, measured using glass density and  $G_{\text{LAC}}$  values, have been plotted at 80, 662, 1173, and 2614 keV in Fig. 12. All results are listed in Table 3. This table indicates that the  $G_{\text{MAC}}$  behaviors with photon energy and sample concentration are identical to  $G_{\text{LAC}}$  behaviors for all glasses.<sup>51–54</sup>

The half-value layers, denoted by  $G_{\text{HVL}}$ , and the mean free path, denoted by  $G_{\text{MFP}}$ , have been computed using the  $G_{\text{LAC}}$  values.<sup>55</sup> The values of  $G_{\text{HVL}}$  and  $G_{\text{MFP}}$  for all the tested glass samples are displayed in Fig. 13 and 14 at the chosen photon energy. As can be observed from these figures, the  $G_{\text{HVL}}$  and  $G_{\text{MFP}}$  values go up as the photon energy increases, but they go down when by-pass effects are considered. Compared to other glass samples, the by-pass15 sample, shown to have the maximum density among the glass ceramics researched, has the lowest  $G_{\text{HVL}}$  values at the tested photon energy, which is well-matched with those obtained in a previous study.<sup>50</sup>

Table 4 shows that the by-pass15 sample has the smallest  $G_{\text{HVL}}$  values, at 662, 1173, and 1333 keV than S1,<sup>56</sup> S2,<sup>57</sup> S3,<sup>58</sup> PCNKBi7.5,<sup>59</sup> Pb20,<sup>60</sup> PbG,<sup>61</sup> and S5<sup>62</sup> and different concrete samples (ordinary concrete (OC), hematite-serpentine (HSC), ilmenite-limonite (ILC), basalt-magnetite (BMC), and steel-scrap (SSC)).<sup>63</sup> This indicates that by-pass15 glass provides superior radiation protection, making it a highly efficient material for shielding applications. Fig. 15 and 16 present a comprehensive analysis of the radiation protection efficiency (RPE) for glasses with varying by-pass contents. Fig. 15 illustrates

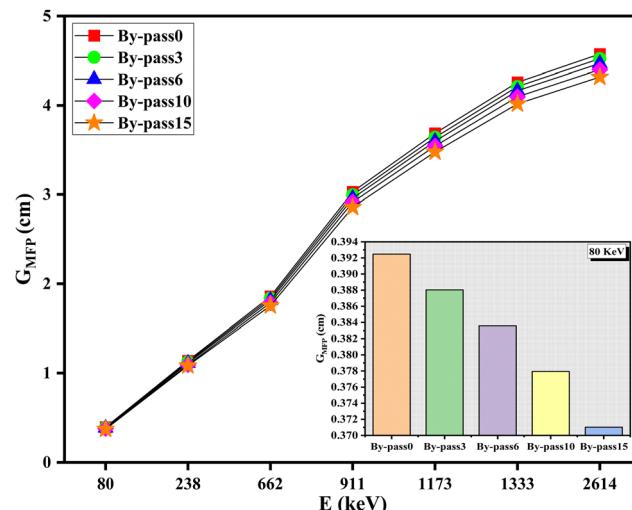


Fig. 14 Mean free path ( $G_{\text{MFP}}$ ) against photon energy for all studied samples.

Table 4 Half value layer values of the by-pass15 glass sample compared to other glasses and concrete

Sample	Half value layer ( $G_{\text{HVL}}$ ) (cm)			Ref.
	662 keV	1173 keV	1333 keV	
By-pass15	1.22	2.41	2.78	This work
S1	4.48	5.87	6.27	56
S2	3.80	4.99	5.32	57
S3	3.06	4.04	4.31	58
PCNKBi7.5	3.92	5.14	5.49	59
Pb20	3.85	5.13	5.48	60
PbG	3.77	4.95	5.28	61
S5	3.36	4.51	4.82	62
OC	3.80	4.98	5.31	63
HSC	3.62	4.75	5.07	
ILC	3.19	4.19	4.47	
BMC	2.98	3.91	4.17	
SSC	2.32	3.05	3.25	

the RPE as a function of glass thickness across a spectrum of photon energies for each by-pass concentration. It is immediately apparent that the RPE increases with increasing glass thickness, which is consistent with the expected behavior of gamma radiation attenuation. Notably, this increase in RPE is more pronounced in glasses with higher by-pass content, indicating that by-pass plays a significant role in enhancing the shielding properties of the glass. Fig. 16 compares the RPE at 80 and 2614 keV at two specific photon energies for the different by-pass concentrations. At a lower energy of 80 keV, the RPE is substantially higher across all samples, which indicates the more effective interaction of low-energy photons with the glass matrix. The by-pass15 sample exhibits the highest RPE at this energy, reinforcing that the by-pass content correlates with improved attenuation.

Conversely, at a higher energy of 2614 keV, there is an overall reduction in RPE for all samples, which aligns with the decreased efficiency of gamma-ray attenuation at higher energies due to the predominance of less interactive processes such as pair production. However, even at this higher energy, the



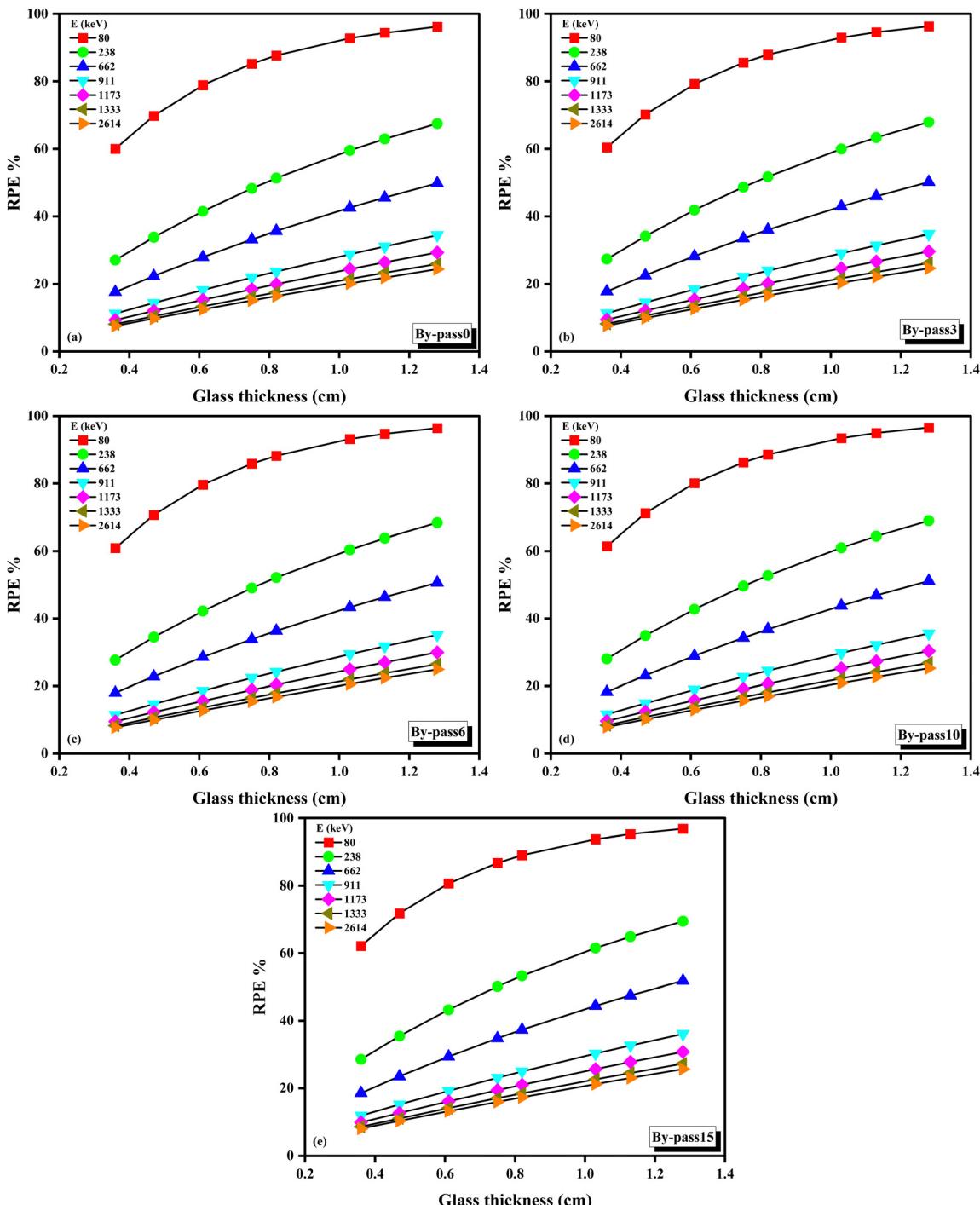


Fig. 15 (a)–(e) Radiation protection efficiency (RPE) for all studied samples at a selected photon energy.

by-pass15 sample maintains a higher RPE than the by-pass0 sample, though the difference is less marked than at 80 keV. This trend confirms the beneficial addition of by-pass to the glass, which not only contributes to increased density but may also introduce structural variations that favor photon attenuation. The interplay between by-pass concentration and RPE is complex and reflects the intricate balance of multiple factors. The compositional variation affects the density and atomic

number, critical parameters for gamma-ray shielding.<sup>64–67</sup> By-pass components such as silica ( $\text{SiO}_2$ ) and alumina ( $\text{Al}_2\text{O}_3$ ) provide additional electron density, which is beneficial for the photoelectric effect at lower energies, while the increase in the overall atomic number enhances the probability of Compton scattering. As the photon energy escalates, these effects become less pronounced, yet the compositional benefits of by-pass are still evident. The data in Fig. 15 and 16 showcase the



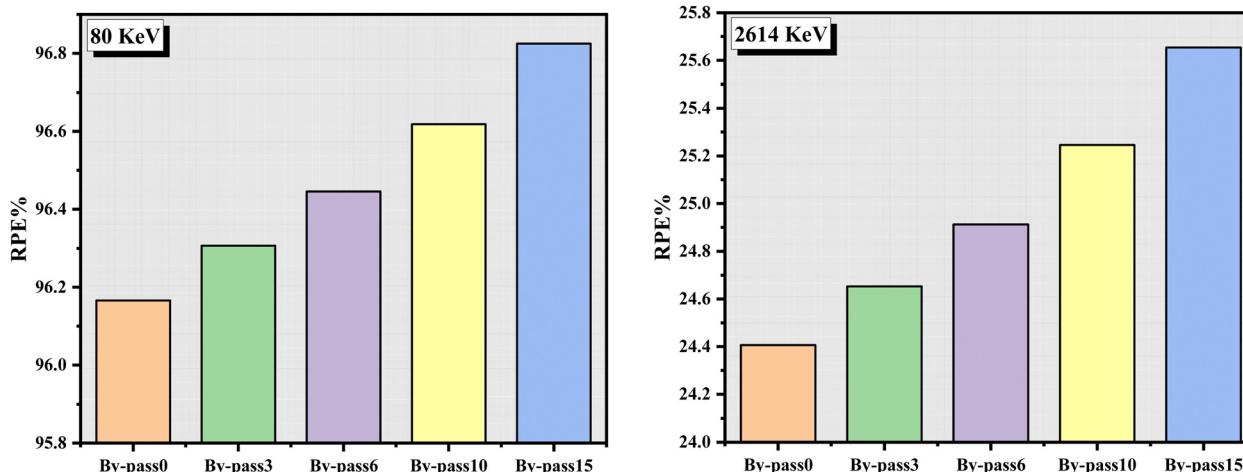


Fig. 16 Radiation protection efficiency (RPE) for all studied glasses at 80 and 2614 keV.

significance of glass composition, particularly the by-pass content, in dictating the RPE. The consistent pattern reinforces our hypothesis that by-pass waste, when integrated into glass matrices, can significantly improve the ability of the material to shield against gamma radiation, offering a dual benefit of waste utilization and enhanced protection against ionizing radiation.

## 4. Conclusions

The results present insightful data that could revolutionize the application of such glasses in gamma shielding endeavors. The XRD patterns of the glass samples confirmed their amorphous nature, denoting a broad scope of potential applications. In the optical domain, UV-visible analysis indicated an absorption surge between 400 and 850 nm as the by-pass concentration increased. Specifically, the by-pass15 sample displayed superior absorption, suggesting the pronounced role of by-pass in enhancing the absorption capacity of the composite. From the radiation shielding analysis, the by-pass15 sample outperformed its counterparts. The by-pass15 glass composition notably achieved the lowest half-value layer ( $G_{HVL}$ ) measurements across the photon energies tested, with values at 662 keV, 1173 keV, and 1333 keV being lower than those of conventional materials like lead-based glass (Pb20 and PbG), indicating a superior radiation shielding capability. Furthermore, the radiation protection efficiency (RPE) of the by-pass15 glass showed a significant increase, outperforming the baseline by-pass0 composition by displaying RPE values that were consistently higher across all thicknesses, with the greatest differential observed at a lower photon energy of 80 keV. In light of these findings, the glasses, especially by-pass15, emerge as potential game-changers in the domain of radiation shielding. Their amorphous nature, enhanced absorption capabilities, and superior shielding properties position them as strong contenders in gamma radiation protection. The practical implications of our findings highlight the dual advantages of utilizing by-pass waste in glass composites: addressing environmental

concerns and creating high-performance materials for optics and radiation protection. The versatility and effectiveness of these composites make them suitable for a wide range of applications, from medical and industrial uses to advanced technological and architectural implementations. Future research can further optimize these materials, potentially paving the way for safer and more efficient environments in industries relying heavily on radioactive isotopes.

## Data availability

The data supporting the findings of this study, including all data presented in figures and tables, are available from the corresponding author, Dr Moustafa A. Darwish, upon reasonable request.

## Conflicts of interest

The authors declare no competing interests.

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