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Non-invasive depth determination of inclusion in biological tissues using spatially offset Raman spectroscopy with external calibration†

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Spatially offset Raman spectroscopy (SORS) allows chemical characterisation of biological tissues at depths of up to two orders of magnitude greater than conventional Raman spectroscopy. In this study, we demonstrate the use of SORS for the non-invasive prediction of depth of an inclusion within turbid media (e.g. biological tissues) using only external calibration data sets, thus extending our previous approach that required internal calibration. As with the previous methodology, the concept is based on relative changes in Raman band intensities of the inclusion that are directly related to the path length of Raman photons travelling through the medium thereby encoding the information of depth of the inclusion. However, here the calibration model is created using data only from external measurements performed at the tissue surface. This new approach facilitates a fully non-invasive methodology applicable potentially to *in vivo* medical diagnosis without any *a priori* knowledge. Monte Carlo simulations of photon propagation have been used to provide insight into the relationship between the spatial offset and the photon path lengths inside the tissues enabling one to derive a general scaling factor permitting the use of spatial offset measurements for the depth prediction. The approach was validated by predicting the depth of surface-enhanced Raman scattering (SERS) labelled nanoparticles (NPs) acting as inclusions inside a slab of *ex vivo* porcine tissue yielding an average root mean square error of prediction of 7.3% with respect to the overall tissue thickness. Our results pave the way for future non-invasive deep Raman spectroscopy *in vivo* by enabling, for example, the localisation of cancer lesions or cancer biomarkers in early disease diagnosis and targeted treatments.

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Introduction

Recent developments in deep Raman spectroscopy such as spatially offset Raman spectroscopy (SORS) and transmission Raman spectroscopy (TRS)^{1,2} have opened a number of new application areas including non-invasive disease diagnosis with its ability to characterise deep layers inside biological tissues.^{3,4–6} SORS takes advantage of characteristic light propagation in diffusely scattering (turbid) media,⁷ being capable of retrieving information about deep layers by introducing a spatial separation between the illumination zone and the Raman collection area.⁸ Recent medically-oriented research

has also demonstrated the use of SORS and TRS for monitoring physical information such as pH⁹ and temperature¹⁰ and localising and predicting the depth of an inclusion (target) within turbid media.^{11–13} Moreover, surface-enhanced Raman spectroscopy (SERS) nanoparticles (NPs) can be used to target specific biomarkers of interest at depth by combining SERS and SORS techniques (*i.e.* SESORS) permitting, for example, effective *in vivo* imaging of cancer lesions^{14–16} or neurotransmitters through skull.¹⁷ Specifically, in the clinical environment, the *in vivo* determination of the depth of subsurface target (e.g. lesion) within body could potentially improve the effectiveness of diagnosis and subsequent treatments. Recent research has proposed and demonstrated the use of SORS to retrieve the depth information of a single inclusion buried within a diffusively scattering synthetic phantom¹² and *ex vivo* tissues.¹³ The method relies on there being a difference in the optical properties (e.g. absorption coefficient μ_a and scattering coefficient μ'_s) of the diffusive medium as a function of wavelength over the acquired Raman spectral range.¹¹ This condition is satisfied readily by biologically relevant samples (e.g. tissues and organs) where μ_a and μ'_s change significantly over

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the NIR range.¹⁸ When Raman photons of different Raman bands, positioned at different wavelengths, travel through this type of matrix, they undergo a different intensity attenuation by the matrix.¹⁹ Due to this effect, the difference in the relative intensities of two, or more, Raman bands of the inclusion can be used to monitor its depth inside a matrix. In such an application, the prediction of the depth is achieved by evaluating relative Raman intensities of inclusion following a calibration procedure that relates these spectral intensity distortions to the depth based on prior internal calibration measurements.¹² However, such procedures are inherently invasive and as such not readily transferable to *in vivo* diagnosis.

To date the fully non-invasive calibration has only been demonstrated in transmission geometries. In this work, for the first time, we propose a calibration method for depth prediction that relies exclusively on SORS data collected at the sample surface and therefore constituting a fully non-invasive concept applicable in non-invasive fashion, potentially also for use in *in vivo* applications. Monte Carlo modelling (*i.e.* MC) of photon propagation²⁰ was used to derive the relationship between propagation path length, depth and spatial offset.²¹ The proposed method was tested and compared with previous approaches in terms of accuracy of depth prediction of a single inclusion buried up to tens of millimetres deep inside synthetic phantom and *ex vivo* porcine tissue. The concepts constitute a major step towards realizing the potential of SORS for non-invasive *in vivo* localisation of abnormal regions of tissue, such as in disease diagnosis.

Proposed concept – external SORS calibration

As previously introduced, the prediction of the depth is obtained by modelling the intensity distortion due to diffuse scattering and absorption events inside a turbid matrix of at

least two different Raman bands of the inclusion. This distortion is related to the photon path length inside the medium. The concept of external calibration resides in finding the correspondence between the intensity distortion and the depth. In the TRS concept, for example, fully non-invasive calibration is readily achievable by placing the inclusion on both sides of the sample, *i.e.* into the laser illumination and collection zones in turn and obtaining a TRS spectrum for each. The relative intensity distortion of the inclusion observed can be directly related to sample thickness. Any smaller distortion observed for this inclusion when located inside the matrix at a specific depth can then be assigned to it based on the above calibration measurement using a proportionality relationship. In contrast to TRS,^{10,13,22} placing the inclusion into the laser illumination and Raman collection cannot be assigned straightforwardly any specific depth in SORS as the spatial offset does not equate directly to depth – although it is related to it by some scaling factor. To find this relationship and specifically the scaling factor, we have carried out Monte Carlo (MC) simulations for medically relevant scenarios. Our MC analysis compared two principal SORS scenarios:

I. With an inclusion placed at the surface of the matrix (within the illumination zone only) and with Raman signal collected at a fixed spatial offset $SO = d$ (Fig. 1a – blue dots).

II. With the inclusion buried deep inside a matrix at depth d_0 right below the Raman collection zone (Fig. 1a – cyan dots). (Note that in this case, the only critical parameter is the depth of inclusion below the Raman collection zone – the location of the illumination zone is irrelevant in the first approximation of negligible dimensions of the inclusion as the laser illumination-to-target section of the photon path is traversed by laser photons that themselves do not contribute to the relative Raman intensity distortion.)

From each of the above scenarios, it is possible to reconstruct the statistical distribution of photon path lengths and

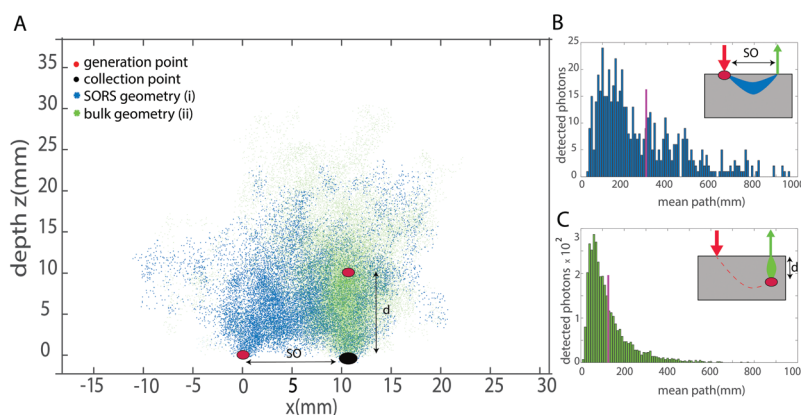


Fig. 1 (a) Monte Carlo modelling of photon trajectory inside a medium in the two different geometries relevant for the calibration plotted as $(x, y = 0, z)$ points: SORS geometry generated with inclusion at the surface (red spot at the origin) and Raman signal collected at 10 mm offset (black spot) giving rise to (x, y, z) locations explored by the photons inside the medium due to propagation (blue dots), and generated with inclusion at 10 mm depth inside the medium and collected at the surface ($SO = d$) giving rise to (x, y, z) locations explored by the photons inside the medium due to propagation (green dots). Mean path length of photons (b) generated at the surface and collected at 10 mm offset (blue bars) and (c) generated in the bulk at a depth of 10 mm (green bars). The magenta bars in (b) and (c) indicate the weighted average for the detected photon.



extrapolate a mean photon path length (Fig. 1b and c). This can then be related to a certain degree of intensity distortion observed. It is anticipated that the imprinting of the same distortion on the spectrum in the two geometries requires the spatial offset to be smaller than the depth of the inclusion by a given scaling factor SF. This is because the spatial offset geometry (with the inclusion at the surface) involves photon trajectories that on average are statistically bent somewhat into the medium, away from the inclusion-Raman collection zone axis, due to the presence of the air-sample interface.

Conversely, in the situation where the inclusion is fully buried inside the medium, the photon trajectories are more direct between the inclusion and the Raman collection zone. A scaling factor SF was derived from the Monte Carlo simulations translating a given spatial offset, SO, into a particular depth of inclusion inside the matrix d_{SO} , such that it would induce the same distortion. To create a calibration model, one needs to divide the spatial offset by the scaling factor in order to derive an equivalent depth ($d_{\text{SO}} = \text{SO}/\text{SF}$) that can then be used to predict an unknown depth of the inclusion inside the same matrix.

Experimental

Calibration procedure and measurement protocol

The external calibration procedure is summarized in Fig. 2. It consists of the sequential acquisition of Raman spectra, at a fixed spatial offset SO, of a target placed at the surface in two positions as illustrated in Fig. 2b:

- At the Raman collection zone on the sample surface (minimum intensity distortion $d = 0$ mm).
- At the laser illumination zone on the sample surface (maximum intensity distortion $d = \text{SO}/\text{SF}$ mm).

For both the calibration measurements (Fig. 2a), a Gaussian-shape curve fit was performed on two Raman bands R_1 and R_2 of the target to evaluate the effect of differential absorption. The natural logarithm of the intensity ratio of the two bands (*i.e.* $\text{IRatio} = \ln(R_1/R_2)$) for the measurement (i) corresponds to the minimum intensity distortion for a depth of the target $d = 0$, while the same parameter for measurement (ii) corresponds to the maximum intensity distortion for a depth of the target equal to the spatial offset divided by the scaling factor extrapolated from MC simulation ($d = d_{\text{SO}} = \text{SO}/\text{SF}$). A linear trend was used to create a model between the intensity distortion, evaluated as above, and the depth (Fig. 2c). The prediction of the depth was obtained by reading an unknown intensity distortion (IRatio) and solving the equation for the depth resulting from the calibration procedure ($\text{depth} = (\text{IRatio} - b)/a$). The measurement protocol consisted of the sequential acquisition of Raman spectra from the model phantom with a fixed spatial offset ($\text{SO} = 15$ mm). The first step of the protocol consists in the alignment of the inclusion under the optical axis (z -axis, collection zone). This is carried out by setting a zero spatial offset ($\text{SO} = 0$ mm) and maximising the signal of the inclusion by moving the sample along the (x, y) axes. This preliminary step ensures the presence of the inclusion under the collection path. After that, the experiments start with the inclusion (*e.g.* paracetamol or NPs) at the surface ($z = 0$) and subsequently being moved deeper along the z -axis.

Monte Carlo modelling

The numerical simulation of photon propagation was performed using the Monte Carlo method where photons were individually followed within a three-dimensional space in a random walk-like propagation as in previous work.^{19–21} The

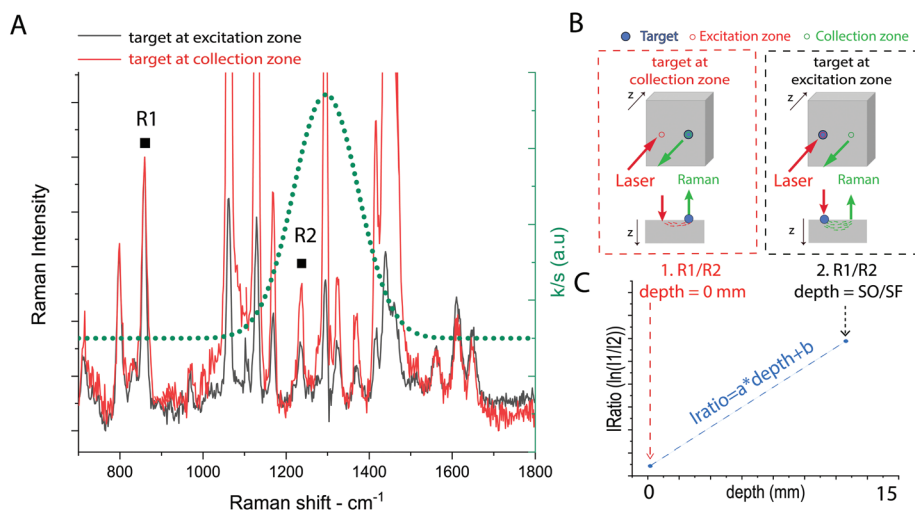


Fig. 2 Calibration procedure: (a) Raman spectra collected with the target (paracetamol) on the sample surface at the laser excitation zone (black curve) and the Raman collection zone (red curve) normalised to the intensity at R_1 . The attenuation profile of the turbid medium (polyethylene) is shown as a green dotted line. Black filled squares denote the two Raman band used for the calibration procedure. (b) Schematic configuration for external calibration procedure based on external measurement in SORS configuration (the inclusion is shown as a blue circle). (c) Calibration model of the intensity distortion (IRatio) vs. depth. I_1 and I_2 are the Raman intensities of R_1 and R_2 , respectively.



model considers the samples to be a semi-infinite turbid medium (with specific absorption and scattering properties, μ_a and μ'_s , respectively) with an air-medium interface at the surface ($z = 0$, Cartesian coordinate); with photons propagating in a straight line over a distance l_s and thereafter its direction is fully randomized at the next scattering event. The propagation distance over which the photon direction is randomized can be approximated as the photon transport length of the scattering medium ($l_s = 1/\mu'_s$).²³ After each interaction, the weight of the photon is reduced by a factor equal to the albedo ($W' = W\mu'_s/(\mu'_s + \mu_a)$) according to the Albedo-weight method²⁴). After a boundary-interaction check, a roulette algorithm for photon termination is implemented, leading to photons eventually terminating in an unbiased manner by conserving energy.²⁰ The simulation parameters were chosen to mimic the experimental condition to enable the comparison between experimental and theoretical data. Point-like excitation and collection geometry were used with excitation and collection radius of $r_{\text{ex}} = 0.5$ mm and $r_{\text{coll}} = 1$ mm, respectively. The following two configurations, mimicking the calibration measurements, were studied:

I. Photon generated at the surface (0,0,0) and collected at the surface with a spatial offset $SO = 10$ mm.

II. Photon generated at a depth $d = 10$ mm inside the turbid media and collected at the surface.

Different pairs of values of $\mu_a = (0, 0.005, 0.01, 0.05, 0.1)$ mm⁻¹ and $\mu'_s = (0.5, 1, 1.5, 2)$ mm⁻¹ were simulated to represent the typical optical properties of tissues and organs in the NIR region of interest.¹⁸ The numerical code was written in Matlab (2018b, the MathWorks). For each geometry and the pair of optical properties, one million (10^6) photons were propagated for a maximum number of steps of 5×10^3 .

Model samples

Firstly, a heterogeneous synthetic phantom made of polyethylene stacks (PE) with dimensions of 50 mm $\times$ 50 mm $\times$ 3 mm per layer and a paracetamol tablet (2 mm radius, 3 mm thickness) embedded at the centre of one of the PE layers were used to mimic the presence of an inclusion (paracetamol) at depth in a turbid medium (PE). The multi-layer phantom¹² was an assembly with an overall thickness of 24 mm and the mass of the inclusion was ~ 125 mg. The inclusion layer was moved through the turbid medium along the z -axis by varying its depth from the surface ($d = 3, 6, 9, 12$ mm).

The second type of multi-layer phantom was made of *ex vivo* animal tissue (turbid matrix) and SERS nanoparticles (NPs) (inclusion).¹³ The *ex vivo* tissue consisted of a stack of sliced back bacon from pork loin (45 $\times$ 45 $\times$ 2 mm or 45 $\times$ 45 $\times$ 3 mm per slice) purchased fresh from a local store and frozen until used. The SERS NPs used as an inclusion consisted of gold NPs (approximately 100 nm diameter) labelled with Raman reporters (d8-4,40-dipyridyl-d8DIPY) and coated with a silica shell (acquired from Oxonica²⁵) already characterised and used previously.^{13,16} The SERS NPs were transferred into a quartz cuvette (2 mm optical path length; $H \times W \times D = 2 \times 10 \times 2$ mm³) and then multiple layers of porcine tissue were added

either side of the cuvette. The sample used for the study consisted of approximately 2.8×10^{14} NPs ml⁻¹. The total thickness of the bilayer sample was kept constant, around 36 mm, while the depth of the cuvette containing the NPs was varied ($d = 0, 4, 7, 10$ mm) by adding tissue layers or moving the quartz cuvette along the optical axis (z -axis) within the stratified tissue slab.

SORS setup

Raman measurements were carried out using a custom-built Raman system for conventional point-like spatial offset Raman spectroscopy (SORS) measurements.^{12,13} A schematic of the set-up is shown in ESI (see Fig. SI-1†). Briefly, the laser light (from a 830 nm diode laser, IPS, Monmouth Junction, NJ) is delivered to the sample through an optical imaging system to a spot size of ~ 500 μ m at the sample surface. The scattered light is collected from a zone of approximately 1.5 mm diameter on the sample surface at a 30° angle relative to the laser beam and coupled through an optical fibre to a Raman spectrometer (Kaiser HoloSpec f/1.8i) and a deep-depletion CCD camera (Andor iDus-420A-BR-DD). The spatial offset is achieved by moving the entire excitation path assembly along the plane parallel to the sample surface. Raman spectra were recorded at a spatial offset of 15 mm with 200 mW laser power at the sample for the synthetic phantom and 300 mW for the *ex vivo* porcine tissue, with 30 s $\times$ 10 accumulations in the spectral range ~ 100 to 1900 cm⁻¹ at a spectral resolution of ~ 8 cm⁻¹.

Results and discussion

Firstly, we studied the photon propagation in terms of mean path length (ESI, Fig. SI-2†) for the two geometries introduced in the MC section, for a range of μ_a and μ'_s relevant to biological tissues in the NIR spectral range.¹⁸ Fig. 3 shows the scaling factor SF calculated as the ratio between the mean paths of the two geometries (*i.e.* (i) and (ii) in the section Proposed Concept – External calibration in SORS) resulting from the MC modelling. The results of MC modelling indicate that the mean path length of photons generated when at the surface and collected with a spatial offset SO (mp_{SORS}) is 2–3 times larger than the mean path of photons generated from an inclusion deep inside the medium at depth equal to the spatial offset SO (mp_{DEPTH}). The same mean photon path length is expected to lead to similar distortion of Raman signals. This implies that, to imprint the same distortion in the two geometries (*i.e.* (i) and (ii) in Fig. 2), the spatial offset needs to be scaled down (by using a proper scaling factor) to create a model that links the depth of the inclusion with the intensity distortion.

To validate the proposed concept, different sets of SORS measurements at 15 mm offset were performed on a multi-layer system based on a turbid medium (PE or *ex vivo* tissue) with an inclusion (paracetamol or NPs) buried at different depths.



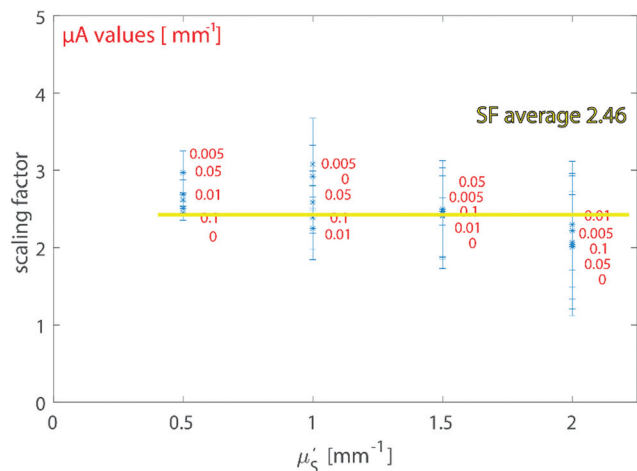


Fig. 3 The scaling factor (SF) resulting from MC modelling for different scattering (μ'_s on x-axis) and absorption (μ'_a ; red label) coefficients. The horizontal yellow line marks the average SF from the MC modelling of all different scattering and absorption coefficients. Vertical error-bars depict the standard error resulting from the ratio of the weighted average.

The results on the synthetic stratified phantom are shown in Fig. 4. The Raman spectra collected at various depths ($d = 3, 6, 9, 12$ mm) for paracetamol buried inside the PE matrix (Fig. 4a) show the differential absorption behaviour of the two paracetamol Raman bands at 857 and 1235 cm^{-1} (the band at 1235 cm^{-1} falls under the PE absorption peak at 930 nm). A scaling factor $SF = 2.45$, derived from MC modelling for $\mu'_s = 0.5$ mm^{-1} and $\mu'_a = 0$ (ESI-2†), was used to create the calibration model for depth prediction based on the intensity distortion defined as the natural logarithm of the ratio of the intensities (857 to 1235 cm^{-1}) for a depth equal to the $SO/SF = 6.12$ mm as in our previous work¹² (red circles in Fig. 4b). Fig. 4c shows the predicted depths of the inclusion resulting from the measured intensity ratio. Similarly, the results on the *ex vivo* tissue-based bilayer phantom are shown in Fig. 5. The Raman signal of SERS

NPs at various depths ($0, 4, 7, 10$ mm) was retrieved with a spatial offset of 15 mm through 36 mm thick porcine tissue (Fig. 5a). Likewise, as observed in previous studies,¹³ the intensity of the Raman component at 1580 cm^{-1} (which falls under an absorption band of water, a significant component of the *ex vivo* tissue) is dramatically reduced relative to the 930 cm^{-1} band as a result of increasing depth. A scaling factor $SF = 2.61$, obtained from MC modelling for $\mu'_s = 0.5$ mm^{-1} and $\mu'_a = 0.05$ mm^{-1} (similar value for porcine muscle at 950 nm was reported in ref. 18), was used to relate the intensity distortion (defined as the natural logarithm of the ratio of the intensities at 930 and 1580 cm^{-1}) measured with the target placed on the surface above the excitation area and the depth ($SO/SF = 5.75$ mm) for creating the calibration model for depth prediction (Fig. 5b).

The predicted depths of the NPs resulting from the measured intensity ratio of R_1 and R_2 (*i.e.* I_1/I_2) are shown in Fig. 5c. The effectiveness of the new calibration approach based on SORS measurements at the surface (*i.e.* external calibration, red symbols in Fig. 4b–c and 5b–c), is compared with the previous approach (Table 1) that used a calibration model for the distortion based on internal measurements¹² (*i.e.* internal calibration, black symbols in Fig. 4b–c and 5b–c). Taking into account all the depths for the inclusion, with only the measurement at 15 mm spatial offset, it was possible to predict the depth of paracetamol and NPs in their respective matrices with a root-mean-square error of prediction (RMSEP) of 10.7% and 7.3% respectively (*i.e.* ± 1.28 mm and ± 0.73 mm, respectively). As expected, the calibration model based on internal measurement resulted in a more accurate prediction with an RMSEP of 9.5% for the synthetic phantom and 7.2% for the *ex vivo* tissue sample. Nevertheless, the proposed methodology based on SORS at a fixed offset with a calibration procedure that relies only on surface measurements proved to be still very effective with the added benefit of being entirely applied in a non-invasive manner (as opposed to the internal calibration procedure).

It should be noted that specific values of SF were used to analyse the samples according to the *a priori* knowledge of the

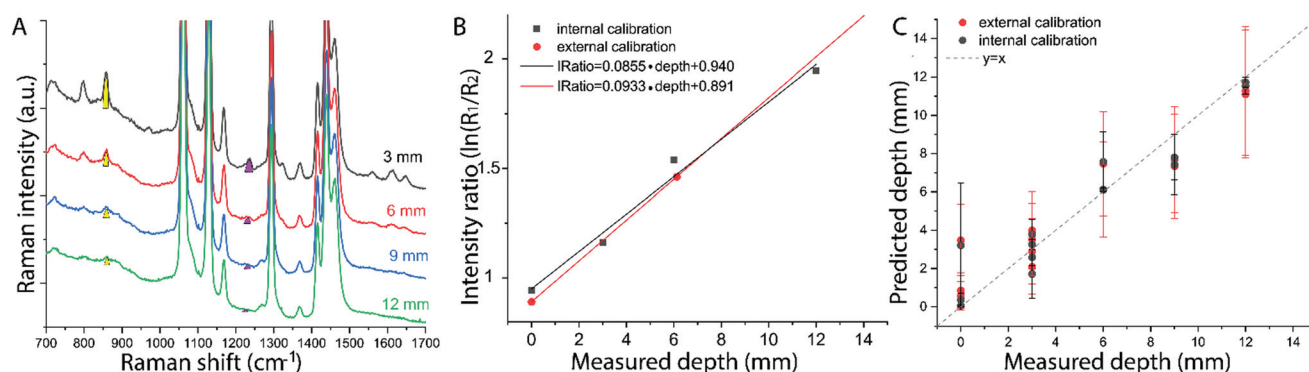


Fig. 4 (a) Raman spectra of the synthetic bilayer system (PE and paracetamol) collected at a fixed spatial offset $SO = 15$ mm with the inclusion (paracetamol) placed at different depths. Yellow and magenta filled areas highlight the two paracetamol Raman bands that undergo differential absorption and are used for depth prediction. (b) Intensity distortion vs. depth used for creating the calibration model based on external measurements (red circles) in comparison with the internal measurements (black squares). (c) Prediction dataset: measured vs. predicted with the two calibration model: external (red circles) and internal (black circles).

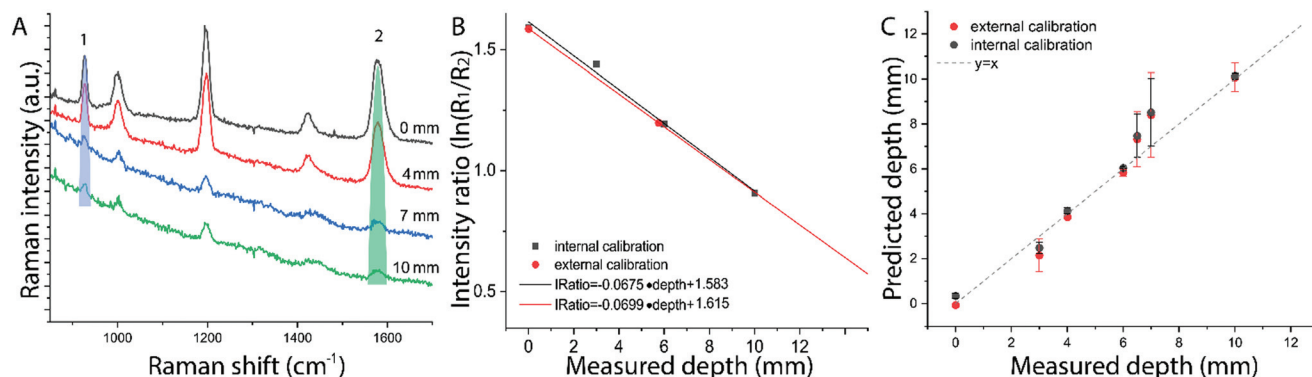


Fig. 5 (a) Raman spectra of the *ex vivo* animal tissue bilayer system collected at a fixed offset $\text{SO} = 15$ mm, with the SERS NPs at different depths. Blue and green filled areas highlight the inclusion (NPs) Raman bands undergoing differential absorption. (b) Calibration model of the intensity distortion vs. depth based on external measurements (red dot) and internal measurements (black square). (c) Depth prediction of the target with the two models.

Table 1 Root mean square error of prediction (RMSEP, mm (%)) of depth for the two different bilayer phantoms using the external and internal calibration model estimated with n number of measurements

RMSEP	External mm (%)	Internal mm (%)
Synthetic phantom	1.28 (10.7%)	1.15 (9.5%)
PE – paracetamol ($n = 14$)		
<i>Ex vivo</i> tissue phantom	0.73 (7.3%)	0.72 (7.2%)
Porcine tissue – NP _S ($n = 7$)		

optical properties. When the values of μ'_s and μ_a are unknown but expected to be within the above-studied range, a rough approximation can be made by considering the average SF resulting from the above MC carried out for typical optical properties of tissue in the NIR region (average value $\text{SF}_{\text{av}} = 2.46$ in Fig. 2). To test and exemplify this, the average value SF_{av} was used to create a calibration model for the *ex vivo* tissue dataset (see ESI-3†). Since the variation of the scaling factor is moderate-low within the calculated scenario, the alteration on the calibration model resulted in a modified RMSEP of only around 2% (new RMSEP = 9.24% with $\text{SF}_{\text{av}} = 2.46$ – see ESI-3†).

The above results are applicable to semi-homogeneous samples and for turbid media, with large scale spatial heterogeneities (e.g. multi-layer systems) but cannot be applied directly if the photon path length varies through more complex dependencies and the relation between intensity distortion and depth of inclusion is not readily scalable with a single scaling factor. Overall, the presented results pave the way towards the use of SORS for non-invasive *in vivo* identification and depth prediction of lesions within biological tissues.

Conclusions

We have proposed and demonstrated the use of SORS for non-invasive prediction of the depth of inclusions within a homogeneous turbid medium without recourse to internal cali-

bration models as required by previous methodologies. The depth of paracetamol tablet and SERS-active NPs buried as single entities inside either a polyethylene matrix or porcine tissue, respectively, have been determined with the non-invasive measurement using only externally acquired calibration measurements. The performance of the method was compared with that of our previous approach based on internal calibration measurements. Although the internal calibration model yielded somewhat more accurate depths, with a lower RMSEP, the external measurement concept proved to be still highly effective (RMSEP 7.3% and 10.7%, instead of 7.2% and 9.5% yielded by the internal calibration method) with a critical advantage of being completely non-invasive and as such highly suitable for *in vivo* uses. Similar results were reported with the use of TRS measurements.¹³ Here, by applying Monte Carlo modelling of photon propagation, it was possible to extend the concept to the SORS geometry creating a model that can be used in situations where TRS cannot be applied. This leads to a wider range of applications where depth determination of a single inclusion is required. In particular, the external calibration concept was validated with *ex vivo* porcine tissue to model typical situations that can be present in *in vivo* scenarios in terms of tissue components (e.g. water, lipid and blood) and geometrical factors. Within this context, the capability of the detection and depth prediction of SERS labelled NPs in biological tissues opens the way for future applications for non-invasive *in vivo* disease diagnostics in clinical settings.

Conflicts of interest

There are no conflicts to declare.

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References

- 1 P. Matousek, I. P. Clark, E. R. C. Draper, M. D. Morris, A. E. Goodship, N. Everall, M. Towrie, W. F. Finney and A. W. Parker, *Appl. Spectrosc.*, 2005, **59**, 393–400.
- 2 P. Matousek, *Appl. Spectrosc.*, 2006, **60**, 185–187.
- 3 M. D. Keller, S. K. Majumder and A. Mahadevan-Jansen, *Opt. Lett.*, 2009, **34**, 926–928.
- 4 P. Matousek and N. Stone, *Analyst*, 2009, **134**, 1058–1066.
- 5 P. Matousek and N. Stone, *J. Biophotonics*, 2013, **6**, 7–19.
- 6 P. Matousek and N. Stone, *Chem. Soc. Rev.*, 2016, **45**, 1794–1802.
- 7 F. Martelli, T. Binzoni, A. Pifferi, L. Spinelli, A. Farina and A. Torricelli, *Sci. Rep.*, 2016, **6**, 1–14.
- 8 P. Matousek, M. D. Morris, N. Everall, I. P. Clark, M. Towrie, E. Draper, A. Goodship and A. W. Parker, *Appl. Spectrosc.*, 2005, **59**, 1485–1492.
- 9 B. Gardner, N. Stone and P. Matousek, *J. Raman Spectrosc.*, 2020, 1–5.
- 10 B. Gardner, N. Stone and P. Matousek, *Faraday Discuss.*, 2016, **187**, 329–339.
- 11 B. Gardner, N. Stone and P. Matousek, *Anal. Chem.*, 2017, **89**, 9730–9733.
- 12 S. Mosca, P. Dey, T. A. Tabish, F. Palombo, N. Stone and P. Matousek, *Anal. Chem.*, 2019, **91**, 8994–9000.
- 13 S. Mosca, P. Dey, T. A. Tabish, F. Palombo, N. Stone and P. Matousek, *J. Biophotonics*, 2020, **13**, 1–7.
- 14 F. Nicolson, B. Andreiuk, C. Andreou, H.-T. Hsu, S. Rudder and M. F. Kircher, *Theranostics*, 2019, **9**, 5899–5913.
- 15 A. I. Henry, B. Sharma, M. F. Cardinal, D. Kurouski and R. P. Van Duyne, *Anal. Chem.*, 2016, **88**, 6638–6647.
- 16 N. Stone, M. Kerssens, G. R. Lloyd, K. Faulds, D. Graham and P. Matousek, *Chem. Sci.*, 2011, **2**, 776.
- 17 A. S. Moody, P. C. Baghernejad, K. R. Webb and B. Sharma, *Anal. Chem.*, 2017, **89**, 5688–5692.
- 18 S. Mosca, P. Lanka, N. Stone, G. Valentini, S. Konugolu Venkata Sekar, P. Matousek and A. Pifferi, *Biomed. Opt. Express*, 2020, **11**, 1697–1706.
- 19 K. L. Bechtel, W. Shih and M. S. Feld, *Opt. Express*, 2008, **16**, 12737–12745.
- 20 S. L. Jacques and L. Wang, in *Optical-Thermal Response of Laser-Irradiated Tissue. Lasers, Photonics, and Electro-Optics*, Springer, Boston, MA, 1995, pp. 73–100.
- 21 P. Matousek, M. D. Morris, N. Everall, I. P. Clark, M. Towrie, E. R. C. Draper, A. E. Goodship and A. W. Parker, *Appl. Spectrosc.*, 2005, **59**, 1485.
- 22 M. Z. Vardaki, H. Sheridan, N. Stone and P. Matousek, *Appl. Spectrosc.*, 2017, **71**, 1849–1855.
- 23 B. B. Das, F. Liu and R. R. Alfano, *Rep. Prog. Phys.*, 1997, **60**, 227–292.
- 24 A. Sassaroli and F. Martelli, *J. Opt. Soc. Am. A*, 2012, **29**, 2110.
- 25 W. E. Doering, M. E. Piotti, M. J. Natan and R. G. Freeman, *Adv. Mater.*, 2007, **19**, 3100–3108.

