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## A single cation or anion dendrimer-based liquid electrolyte†

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We propose here a novel liquid dendrimer-based single ion conductor as a potential alternative to conventional molecular liquid solvent–salt solutions in rechargeable batteries, sensors and actuators. A specific change from ester ( $-COOR$ ) to cyano ( $-CN$ ) terminated peripheral groups in generation-one poly(propyl ether imine) ( $G_1$ -PETIM)–lithium salt complexes results in a remarkable switchover from a high cation ( $t_{Li^+} = 0.9$  for  $-COOR$ ) to a high anion ( $t_{PF_6^-} = 0.8$  for  $-CN$ ) transference number. This observed switchover draws an interesting analogy with the concept of heterogeneous doping, applied successfully to account for similar changes in ionic conductivity arising out of dispersion of insulator particle inclusions in weak inorganic solid electrolytes. The change in peripheral group simultaneously affects the effective ionic conductivity, with the room temperature ionic conductivity of PETIM–CN ( $1.9 \times 10^{-5} \Omega^{-1} \text{ cm}^{-1}$ ) being an order of magnitude higher than PETIM–COOR ( $1.9 \times 10^{-6} \Omega^{-1} \text{ cm}^{-1}$ ). Notably, no significant changes are observed in the lithium mobility even following changes in viscosity due to the change in the peripheral group. Changes in the peripheral chemical functionality directly influence the anion mobility, being lower in PETIM–COOR than in PETIM–CN, which ultimately becomes the sole parameter controlling the effective transport and electrochemical properties of the dendrimer electrolytes.

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## Introduction

One of the important strategies towards building stable and safe rechargeable batteries has been to develop newer forms of electrolytes as potential alternatives to conventional liquid molecular solvent–salt solutions.<sup>1,2</sup> In this line of thought, various electrolyte systems ranging from solid crystalline electrolytes to “solid-like” soft organic electrolytes have been explored as alternatives to conventional liquid electrolytes focused mainly for applications in lithium-based battery chemistries.<sup>3–7</sup> Polymer electrolytes, which exhibit interesting compliant mechanical properties in addition to high ionic conductivity, have shown greater potential than solid crystalline electrolytes in diverse electrochemical devices, *viz.* batteries,<sup>8a</sup> fuel cells,<sup>8b</sup> actuators,<sup>8c</sup> sensors.<sup>9</sup> Design of novel polymer architectures (*e.g.* network, branched polymers) has been one of the important strategies for the development of high ion conducting polymer electrolytes.<sup>10</sup> The major drawback with various polymer-based electrolytes is that the cations and anions contribute to the specific conductivity and thus the

specific ion, *i.e.* cation or anion, transference number is not high. In particular the cation ion transference number, which is of practical interest for various rechargeable battery chemistries, is typically low in the range  $t_+ = 0.2–0.5$ .<sup>10,11</sup> There have been several strategies to design polymer electrolytes with high transference number without significantly compromising the effective ionic conductivity. The majority of these approaches have been applied on polymers in the solid form, *viz.* *via* chemical manipulations of the constituting units of the polymer or from single ion conductors where an ion of one type (say anion) is immobilized on the polymer backbone as in block or copolymer units. The other major concern with solid-like electrolytes is related to poor charge transport kinetics at the electrode|electrolyte interface which leads to poor device efficiency. Apart from a few glowing examples, the overwhelming majority of polymer electrolytes and in general solid electrolytes have not been able to successfully transcend the precincts of laboratory-scale demonstrations. This has led to the persistence of conventional liquid electrolytes in the majority of modern-day electrochemical devices, including rechargeable batteries.

Dendrimers are a special class of mono-dispersed branched polymers, containing a large number of branched flexible chain-ends emanating from a core or linker molecules. This unique architecture has attracted considerable attention in biomedicine, catalysis, sensing and energy storage.<sup>12</sup> The mechanical consistency of dendrimers is intermediate between

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low viscosity molecular solvents ( $\eta \approx 10^{-3}$  Pa s) and high viscosity polymer gel or polymer–salt complex (very high  $\eta$ ;  $\eta \rightarrow \infty$ ) electrolytes. Due to the higher viscosities of dendrimers compared to typical molecular solvents, the ionic mobility and hence the effective ionic conductivity of dendrimer electrolytes are expected to be lower compared to molecular solvent–salt liquid solutions. In the context of the vast volume of work accomplished with regard to solid polymers with high ionic conductivity,<sup>13</sup> viscosity cannot be the sole criteria for the determination of ionic mobility and effective ionic conductivity. Similarly, the high viscosity of dendrimers should not be a deterrent for exploring their application in electrochemical devices. The high degree of branching in the dendrimer network leads to multiple advantages, *viz.* larger free volume, higher amorphicity and low glass transition temperatures (implying higher chain flexibility).<sup>14</sup> These, coupled with the flexibility to freely tune the chemical composition and conformation as a function of generation number, also significantly affect the ion-solvating ability and host–guest interactions with various metal salts including alkali-metal salts which are of direct relevance to rechargeable batteries. These advantageous features should make dendrimers an attractive alternative liquid matrix to conventional liquids, ionic liquids or solid polymers for the synthesis of ion conductors tailored to perform specific tasks in various electrochemical devices. Of specific interest is whether a dendrimer can be employed to produce electrolytes with high ion transference number of a single ion type. To the best of our knowledge there have been no efforts undertaken in this direction. Additionally, there have been no detailed and conclusive studies undertaken on the correlation of various chemical functional parameters with the ion transport mechanism in dendrimers, in spite of their anticipated potential in various electrochemical applications. Studies in these directions will be expected to throw more light on the ion transport mechanism in dendrimers and identify key parameters for the development of dendrimer-based ion conductors for various applications such as rechargeable batteries, sensors, and actuators.<sup>14,15</sup> We present here for the first time a detailed study of the influence of the chemical nature of peripheral functional groups on ionic conductivity, diffusion and transference number in generation-1 poly(propyl ether imine) ( $G_1$ -PETIM)–lithium salt mixtures. We demonstrate here that the peripheral chemical functionality is a very important parameter to optimize the effective transport as well as the electrochemical properties of the dendrimer electrolyte.

## Results and discussion

Ion transport in polymer electrolytes largely depends on the chemical characteristics of the polymer, such as branching, network, and functionality of the side chains.<sup>8a,16,17</sup> Different functional groups exhibit different binding abilities due to varying polarity of the groups and this influences the ion solvation *via* dissociation of the salt. On the other hand, their spatial distributions and the sizes of end chains affect both cation and anion mobility, resulting in large differences in the corresponding ion transference numbers.<sup>17</sup> The above stated

issues on ionic conductivity and transference number are systematically probed here in the context of  $G_1$ -PETIM dendrimers. The  $G_1$ -dendrimers are synthesized with different end functional groups ( $-COOH$ ,  $-COOR$ ,  $-OH$ ,  $-CN$ ) maintaining the same linker (ether) and branching points (tertiary amine; *cf.* schematic Fig. 1). The study here focuses only on the first generation dendrimers primarily due to the following reasons. Firstly, the viscosity (0.1–6 Pa s) of the  $G_1$ -PETIM dendrimer, though higher than the viscosity of a typical liquid molecular solvent (0.5–10 mPa s),<sup>18</sup> is much lower than the higher generation  $G_n$ -PETIM dendrimers ( $n = 2–4$ ). So, the influence of viscosity on ionic conductivity will be much less in  $G_1$ -PETIM compared to  $G_n$ -PETIM dendrimer liquid electrolytes. The magnitude of ionic conductivity of some of the  $G_1$  dendrimers will not be significantly lower compared to a typical molecular liquid solvent–salt solution<sup>17</sup> of relevance to rechargeable batteries (typically:  $>10^{-3}$  to  $10^{-2} \Omega^{-1} \text{ cm}^{-1}$ ). Secondly, in  $G_1$  dendrimers the number density of the linker and branching moieties is lower compared to higher generation dendrimers. Hence, the non-trivial influence due to the linker and branching moieties on solvation and ion mobility anticipated in the higher generations will be minimal and assumed constant in the  $G_1$  dendrimers. The ion transport in the  $G_1$  dendrimer electrolytes can then be directly correlated to the nature of the peripheral chemical functionality.

The temperature dependent ionic conductivity and viscosity of first-generation dendrimers (with different functional groups) with 0.1 M LiPF<sub>6</sub> in the temperature range 0–60 °C are shown in Fig. 2a and b respectively. The first-generation nitrile-terminated PETIM dendrimer electrolyte ( $G_1$ -CN–0.1 M LiPF<sub>6</sub>)

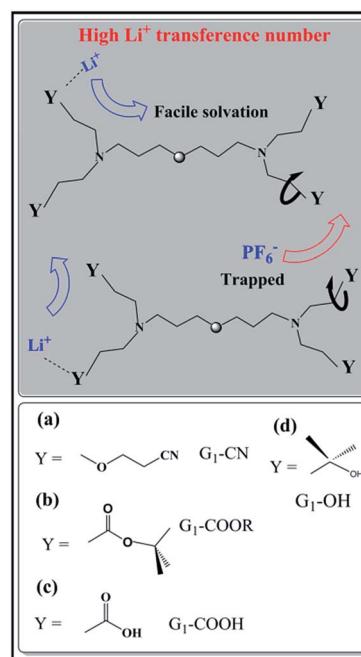


Fig. 1 Schematic representation of ion transport mechanism in  $G_1$ -PETIM dendrimers with different peripheral functional groups: (a)  $G_1$ -CN, (b)  $G_1$ -COOR, (c)  $G_1$ -COOH, (d)  $G_1$ -OH.

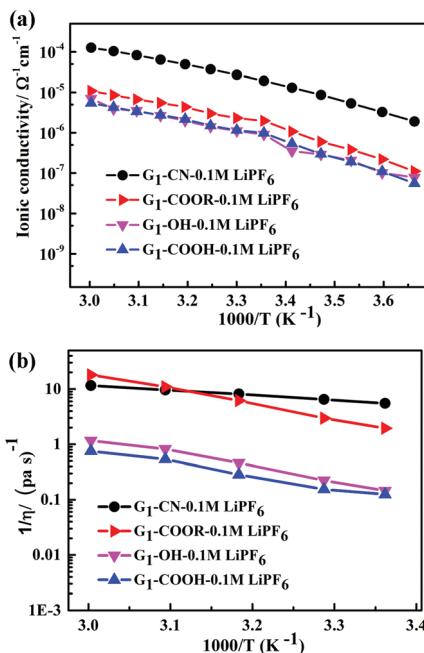


Fig. 2 (a) Temperature dependent ionic conductivity and (b) fluidity ( $\eta^{-1}$ ) of  $\text{G}_1$  dendrimers with different peripheral functional groups.

exhibits the highest room temperature ionic conductivity of  $1.9 \times 10^{-5} \Omega^{-1} \text{cm}^{-1}$ , whereas  $\text{G}_1\text{-OH-0.1 M LiPF}_6$  shows the lowest conductivity of  $9 \times 10^{-7} \Omega^{-1} \text{cm}^{-1}$ . The conductivities of  $\text{G}_1\text{-COOR-0.1 M LiPF}_6$  and  $\text{G}_1\text{-COOH-0.1 M LiPF}_6$  at 25 °C are intermediate to those of  $\text{G}_1\text{-CN-0.1 M LiPF}_6$  and  $\text{G}_1\text{-OH-LiPF}_6$ , being  $1.9 \times 10^{-6} \Omega^{-1} \text{cm}^{-1}$  and  $9.8 \times 10^{-7} \Omega^{-1} \text{cm}^{-1}$  respectively. The measured ionic conductivity values of the various dendrimer electrolytes, which are on a par with many single ion conductors,<sup>10a,19</sup> have the following trend:  $\sigma_{\text{G}_1\text{-CN}} > \sigma_{\text{G}_1\text{-COOR}} > \sigma_{\text{G}_1\text{-COOH}} \sim \sigma_{\text{G}_1\text{-OH}}$ . This variation in conductivity between the various dendrimer electrolytes by more than one order of magnitude indicates a strong correlation between the chemical nature of peripheral functionalization and ion transport. The activation energies of conductivity are obtained by a linear fit of the conductivity data using the Arrhenius equation:  $\sigma = Ae^{-E_a/kT}$ , where  $A$ ,  $E_a$ ,  $k$ ,  $T$  are the pre-exponential factor, activation energy, Boltzmann constant and temperature respectively (cf. fitting parameters in Table ST2†). In spite of notable differences in ionic conductivities between  $\text{G}_1\text{-CN}$ ,  $\text{G}_1\text{-COOR}$  and  $\text{G}_1\text{-COOH}$ , interestingly no significant differences exist in the estimated activation energies between  $\text{G}_1\text{-CN}$  ( $E\sigma_{\text{G}_1\text{-CN}} = 0.54$  eV),  $\text{G}_1\text{-COOR}$  ( $E\sigma_{\text{G}_1\text{-COOR}} = 0.58$  eV) and  $\text{G}_1\text{-COOH}$  ( $E\sigma_{\text{G}_1\text{-COOH}} = 0.58$  eV). The VTF fitting<sup>20</sup> parameters of the conductivity plot (Fig. 2a) are tabulated in Table ST3.† However, it is strongly felt that VTF fitting of the conductivity results is inappropriate for the present study. This is mainly attributed to the simpler chemical structure of the first-generation dendrimers compared to higher generation dendrimers and polymers.

The conductivity behaviour is correlated to the temperature dependent fluidity ( $1/\eta$ , where  $\eta$  is the viscosity) shown in Fig. 2b. The viscosities of PETIM dendrimers, calculated from the static viscosity *versus* shear rate measurements, are shown

in Table ST2.† The viscosity increases by nearly two times from 0.15 Pa s for  $\text{G}_1\text{-CN}$  to 0.30 Pa s for  $\text{G}_1\text{-COOR}$ . The viscosities at 30 °C for  $\text{G}_1\text{-OH-0.10 M LiPF}_6$  and  $\text{G}_1\text{-COOH-0.10 M LiPF}_6$  are even higher, being 4.5 Pa s and 6.5 Pa s respectively. The viscosity values showed the following trend:  $\eta_{\text{G}_1\text{-CN}} < \eta_{\text{G}_1\text{-COOR}} < \eta_{\text{G}_1\text{-OH}} < \eta_{\text{G}_1\text{-COOH}}$ . The higher viscosity of the ester dendrimer ( $\text{G}_1\text{-COOR}$ ) compared to  $\text{G}_1\text{-CN}$  is attributed to the more polar nature of the -COOR group compared to the -CN group and the steric hindrance exerted by the bulkier -COOR ( $R = t\text{-butyl}$ ) group implying a higher dragging force compared to the linear -CN in  $\text{G}_1\text{-CN}$ . Differences in viscosity between  $\text{G}_1$ -carboxyl/hydroxyl groups (*i.e.* COOR, COOH and OH) and  $\text{G}_1\text{-CN}$  can also be accounted for on the basis of the intra- or inter-molecular hydrogen bonding. The strength of hydrogen bonds is significantly higher in  $\text{G}_1\text{-COOH}$  and  $\text{G}_1\text{-OH}$  resulting in significantly higher viscosity compared to  $\text{G}_1\text{-CN}$  and  $\text{G}_1\text{-COOR}$ . An increase in viscosity for carboxyl groups ( $\text{G}_1\text{-COOR}$  and  $\text{G}_1\text{-COOH}$ ) and hydroxyl ( $\text{G}_1\text{-OH}$ ) terminated dendrimers results in a decrease in ionic conductivity compared to the  $\text{G}_1\text{-CN}$  dendrimer. Activation energies of viscosity, obtained by fitting the temperature dependent fluidity ( $1/\eta$ ) (Fig. 2b) using the Arrhenius equation, are tabulated in Table ST2.† Similar trends in activation energies for both temperature dependent conductivity and viscosity indicate that the underlying mechanism for conductivity and viscosity is thermally activated. However, the observed trends in viscosity cannot be correlated one-to-one with the conductivity trends. This is especially true for the -COOH terminated dendrimer. A logical interpretation of the observed trends in conductivity and viscosity can only be achieved by simultaneously studying the diffusion behavior of the various participating entities, *viz.*  $\text{Li}^+$ ,  $\text{PF}_6^-$  and H (dendrimer).  $\text{G}_1\text{-COOH}$  and  $\text{G}_1\text{-OH}$  have been excluded from additional studies as both have lower ionic conductivities and higher viscosities compared to  $\text{G}_1\text{-CN}$  and  $\text{G}_1\text{-COOR}$  in the measured temperature range.  $\text{G}_1\text{-COOR}$  is selected as the representative among the two carboxyl groups as an understanding of the mechanism in -COOR will also aid in accounting for the experimental observations in -COOH.

Ion solvation in  $\text{G}_1\text{-CN-LiPF}_6$  and  $\text{G}_1\text{-COOR-LiPF}_6$  electrolytes was studied by FTIR spectroscopy at various salt concentrations (ranging from 0 to 0.1 M) and temperatures (from room temperature to 70 °C). The FTIR spectra for both the electrolytes in the wavenumber region 800–920 cm $^{-1}$  (normalized with respect to the highest intensity peak at 1118 cm $^{-1}$  for  $\text{G}_1\text{-CN}$  and 1154 cm $^{-1}$  for  $\text{G}_1\text{-COOR}$ ) are shown in Fig. 3. The area under the peak is calculated by fitting the spectra with Gaussian function. Fig. 3a and b depicts the FTIR spectra of  $\text{G}_1\text{-CN-}x\text{ M LiPF}_6$  and  $\text{G}_1\text{-COOR-}x\text{ M LiPF}_6$  respectively with varying salt concentration ( $x$ ) ranging from 0 to 0.10 M. Both pristine  $\text{G}_1\text{-COOR}$  and  $\text{G}_1\text{-CN}$  exhibit the characteristic symmetric stretching vibration of C–O–C groups of the aliphatic ether present in the dendrimer core at 848 cm $^{-1}$  (splitting in the C–O–C symmetric stretch in  $\text{G}_1\text{-CN}$  leads to an additional band at 830 cm $^{-1}$ , possibly due to the presence of two types of ether groups in the core and at the periphery).<sup>21</sup>

This band is observed to merge with the strong band also appearing at 848 cm $^{-1}$  corresponding to the P–F vibration of the



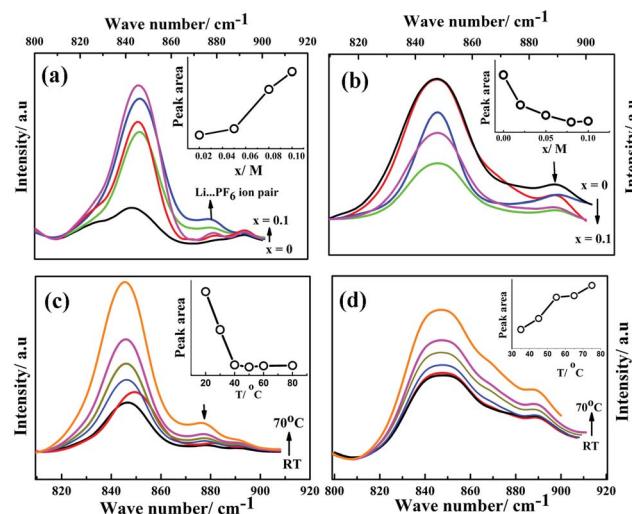


Fig. 3 FTIR spectra of  $\text{G}_1\text{-CN-LiPF}_6$  (a) and  $\text{G}_1\text{-COOR-LiPF}_6$  (b) at different salt concentrations [ $x = 0$  (black), 0.01 (red), 0.02 (magenta), 0.05 (green), 0.1 (blue)]. FTIR spectra of  $\text{G}_1\text{-CN-LiPF}_6$  (c) and  $\text{G}_1\text{-COOR-LiPF}_6$  (d) at different temperatures ( $T = \text{RT}$  to  $70\text{ }^\circ\text{C}$ , with  $x = 0.1$ ) [ $T = \text{RT}$  (black),  $30\text{ }^\circ\text{C}$  (red),  $40\text{ }^\circ\text{C}$  (blue),  $50\text{ }^\circ\text{C}$  (green),  $60\text{ }^\circ\text{C}$  (magenta),  $70\text{ }^\circ\text{C}$  (orange)].

free  $\text{PF}_6^-$  ion of the salt<sup>22</sup> in  $\text{G}_1\text{-CN-}x\text{ M LiPF}_6$ . A weak band appearing at  $875\text{ cm}^{-1}$  for  $\text{G}_1\text{-CN-}x\text{ M LiPF}_6$  is assigned to the associated tri-dentate ion pair of the  $\text{LiPF}_6$  salt and this is in good agreement with previous reports.<sup>22</sup> With increasing salt concentration, the band area of the weaker band at  $875\text{ cm}^{-1}$  is observed to intensify (inset of Fig. 3a) with respect to the stronger band (at  $848\text{ cm}^{-1}$ ) for  $\text{G}_1\text{-CN-}x\text{ M LiPF}_6$ . This suggests the presence of ion pairs in  $\text{G}_1\text{-CN-}x\text{ M LiPF}_6$ . The intensity of this shoulder band at  $875\text{ cm}^{-1}$  decreases with increasing temperature from RT to  $70\text{ }^\circ\text{C}$  (Fig. 3c), signifying the dissociation of ion pairs with increasing temperature. In the case of pure  $\text{G}_1\text{-COOR}$ , the observed IR band at  $898\text{ cm}^{-1}$  (Fig. 3b) is attributed to the C-C stretching frequencies of ester ( $\text{O}-\text{C}-\text{C}$ ) or ether groups,<sup>21</sup> which is affected by addition of the salt as well as temperature. No additional bands corresponding to ion pairs are observed in  $\text{G}_1\text{-COOR-}x\text{ M LiPF}_6$  (Fig. 3b) at various salt concentrations, signifying facile salt dissociation in ester functionalized dendrimers. Stronger interaction of the oxygen atom of the ester group with the  $\text{Li}^+$  ion, compared to the  $-\text{CN}$  group, leads to higher dissociation of ion pairs in  $\text{G}_1\text{-COOR}$ . To support the FTIR observations, the Stokes equations (cf. ESI†) for both dendrimers have been investigated (Fig. S3†). Fig. S3† shows the product of dc conductivity (related to the number of free charges from the Stokes equation) and viscosity at various temperatures for both dendrimers. The constancy of this product (i.e.  $Nq^2/6\pi r_s$  in eqn (SE1) in ESI†) at various temperatures for  $\text{G}_1\text{-COOR}$  signifies no change in free charges with an increase of temperature, which is true for electrolytes with a fully dissociated salt. In the case of  $\text{G}_1\text{-CN}$ , an increase in free charges is observed with an increase in temperature (in Fig. S3 in ESI†), which is attributed to an increase in salt dissociation with increasing temperature.

The self-diffusion coefficients, characterizing long-range macroscopic transport of  ${}^7\text{Li}$ ,  ${}^{19}\text{F}$  and  ${}^1\text{H}$  nuclei, are obtained from multinuclear PFG-NMR experiments using  $\ln I/I_0 = -D^{\text{NMR}}\gamma^2(\Delta - \delta/3)\delta^2g^2$ ,<sup>23</sup> where  $I$  and  $I_0$  are the signals in the presence and absence of the gradient respectively,  $\gamma$  is the gyromagnetic ratio of the nucleus studied,  $\Delta$  is the interval between the gradient pulses,  $\delta$  is the length of the gradient pulse, and  $g$  is the magnitude of the gradient pulse. Self-ionic diffusion coefficients in the temperature range  $0\text{--}60\text{ }^\circ\text{C}$  are shown in Fig. 4a.

$\text{G}_1\text{-CN-}0.1\text{ M LiPF}_6$  exhibits much higher  ${}^1\text{H}$  diffusion coefficients (varying from  $1.9 \times 10^{-12}$  to  $6.3 \times 10^{-11}\text{ m}^2\text{ s}^{-1}$  between 0 and  $60\text{ }^\circ\text{C}$ ), nearly one order in magnitude higher compared to  $\text{G}_1\text{-COOR-}0.1\text{ M LiPF}_6$  ( $2.5 \times 10^{-13}$  to  $1.8 \times 10^{-11}\text{ m}^2\text{ s}^{-1}$  between 10 and  $60\text{ }^\circ\text{C}$ ). Following the Stokes-Einstein equation ( $D = kT(6\pi\eta r_s)^{-1}$ , where  $\eta$ ,  $D$  and  $r_s$  are viscosity, self-diffusion coefficient and effective hydrodynamic (Stokes) radius respectively), the higher viscosity (0.3 Pa s) of  $\text{G}_1\text{-COOR-}0.1\text{ M LiPF}_6$  compared to  $\text{G}_1\text{-CN-}0.1\text{ M LiPF}_6$  (0.15 Pa s) results in lower  ${}^1\text{H}$  diffusion coefficients for the ester dendrimer.  ${}^{19}\text{F}$  self-diffusion coefficients for both  $\text{G}_1\text{-CN-}0.1\text{ M LiPF}_6$  ( $1.36 \times 10^{-12}$  to  $7.8 \times 10^{-11}\text{ m}^2\text{ s}^{-1}$  between 0 and  $60\text{ }^\circ\text{C}$ ) and  $\text{G}_1\text{-COOR-}0.1\text{ M LiPF}_6$  ( $3.5 \times 10^{-13}$  to  $1.1 \times 10^{-11}\text{ m}^2\text{ s}^{-1}$  between 20 and  $60\text{ }^\circ\text{C}$ ) are found to be in close proximity to their respective  ${}^1\text{H}$  diffusion coefficient values. The similarities in diffusion coefficient values between  ${}^{19}\text{F}$  and  ${}^1\text{H}$  nuclei for both  $\text{G}_1\text{-COOR}$  and  $\text{G}_1\text{-CN}$

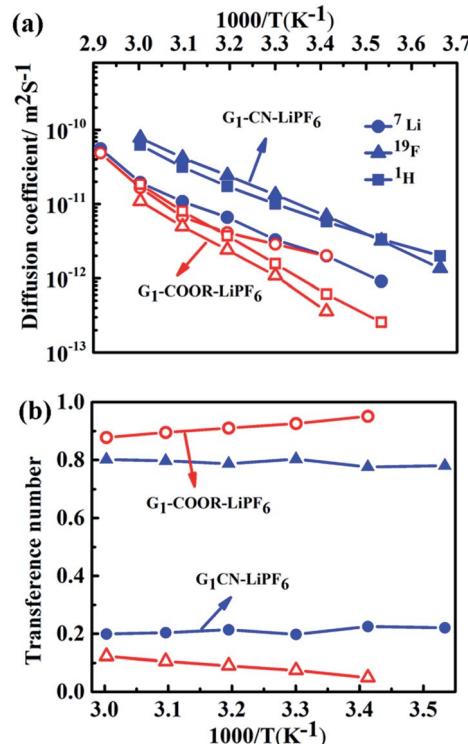


Fig. 4 (a) Self-diffusion coefficients of  ${}^1\text{H}$  (squares),  ${}^{19}\text{F}$  (triangles),  ${}^7\text{Li}$  (circles) at various temperatures for  $\text{G}_1\text{-CN-}0.1\text{ M LiPF}_6$  (blue, with closed symbols) and  $\text{G}_1\text{-COOR-}0.1\text{ M LiPF}_6$  (red, with open symbols). (b)  ${}^7\text{Li}$  and  ${}^{19}\text{F}$  transference numbers of  $\text{G}_1\text{-CN-}0.1\text{ M LiPF}_6$  (blue) and  $\text{G}_1\text{-COOR-}0.1\text{ M LiPF}_6$  (red) at various temperatures.

signify correlated  $\text{PF}_6^-$  anion motion with the dendrimer molecules. The lower viscosity of  $\text{G}_1\text{-CN}$ -0.1 M  $\text{LiPF}_6$  results in nearly one order of magnitude higher  $^{19}\text{F}$  diffusion coefficient for  $\text{G}_1\text{-CN}$ -0.1 M  $\text{LiPF}_6$  ( $1.3 \times 10^{-11} \text{ m}^2 \text{ s}^{-1}$  at 30 °C) compared to that of  $\text{G}_1\text{-COOR}$ -0.1 M  $\text{LiPF}_6$  ( $1.1 \times 10^{-12} \text{ m}^2 \text{ s}^{-1}$  at 30 °C). The estimated  $R_{\text{PF}_6^-}$  ( $= D_{\text{H}}/D_{\text{F}}$ ) for  $^{19}\text{F}$  in  $\text{G}_1\text{-CN}$  (0.6) and  $\text{G}_1\text{-COOR}$  (0.8) supports a stronger solvent-correlated  $\text{PF}_6^-$  motion in  $\text{G}_1\text{-COOR}$  compared to  $\text{G}_1\text{-CN}$ . The reason behind this correlation is possibly due to stronger steric hindrance posed by the bulkier *t*-butyl groups to  $\text{PF}_6^-$  mobility<sup>24</sup> in  $\text{G}_1\text{-COOR}$  compared to linear -CN in  $\text{G}_1\text{-CN}$ . In contrast, lithium-ion diffusion is not at all influenced by the viscosity. Similar values of  $^7\text{Li}$  self-diffusion coefficients are observed in both cases ( $2.0 \times 10^{-12}$  to  $5 \times 10^{-11} \text{ m}^2 \text{ s}^{-1}$  between 30 and 60 °C). The contributions of  $\text{Li}^+$  and  $\text{PF}_6^-$  towards the total effective conductivity for  $\text{G}_1\text{-CN}$ , calculated following the Nernst-Einstein equation<sup>23,25a</sup> ( $\sigma_i = Nq^2(kT)^{-1}D_i$ , where  $\sigma_i$  and  $D_i$  are the dc ionic conductivity and diffusion coefficient of the *i*<sup>th</sup> ion type), are  $1.2 \times 10^{-5} \Omega^{-1} \text{ cm}^{-1}$  and  $4.9 \times 10^{-5} \Omega^{-1} \text{ cm}^{-1}$  respectively. The higher value of anion conductivity strongly suggests that  $\text{G}_1\text{-CN}$  is predominantly an anion conductor. On the other hand, the lower contribution of  $^{19}\text{F}$  ( $\sigma = 3.9 \times 10^{-6} \Omega^{-1} \text{ cm}^{-1}$  from NMR data) to the total conductivity in  $\text{G}_1\text{-COOR}$  compared to its  $^7\text{Li}$  diffusion ( $1.1 \times 10^{-5} \Omega^{-1} \text{ cm}^{-1}$ ) strongly suggests cation transport in the ester dendrimer. Thus, trapping of anions by the peripheral bulkier ester group in  $\text{G}_1\text{-COOR}$  makes it a single cationic conductor. Thus, the difference in effective ionic conductivity values between  $\text{G}_1\text{-CN}$  and  $\text{G}_1\text{-COOR}$  is mainly due to the difference in the anionic conductivity between them. Manipulation of the chemical constitution of dendrimers *via* variations in the peripheral group, which exert varying degrees of steric hindrance to the mobility of anions, results in a transformation from an anionic to a cationic conductor. This approach is interesting as it becomes the organic analog to the concept of heterogeneous doping introduced by Maier.<sup>26</sup> Heterogeneous doping has been in the past successfully implemented to account for changes in effective ionic conductivity of solid-solid composites comprised of dispersions of nanometer- to micrometer-sized oxide additives (e.g.  $\text{Al}_2\text{O}_3$ ,  $\text{SiO}_2$ ) in a weak solid electrolyte, e.g.  $\text{LiI}$ ,  $\text{TiCl}_2$ .<sup>27</sup> In this concept, the changes in conductivity have been attributed to the space-charge layer formed at the interface of the weak electrolyte-oxide insulator which directly influences the transition from an anion to a cation conductor as demonstrated in  $\text{TiCl}_2\text{-Al}_2\text{O}_3$ . The concept with limited success was later extended to liquids where dispersions of fine oxide particles in liquid electrolytes lead to modest enhancements in the effective conductivity of the liquid.<sup>28,29</sup> To the best of our knowledge the heterogeneous doping concept has not yet been utilized to transform the nature of ion transport in liquids. It is envisaged that the approach presented here is the first of its kind to be adopted in the realm of liquids. This adopted approach is simpler and is expected to be highly efficient and reproducible compared to the addition of oxides, which display considerable non-uniformities in size and chemical functionality of the oxide additive.

In the case of  $\text{G}_1\text{-CN}$ -0.1 M  $\text{LiPF}_6$ , the observed  $^7\text{Li}$  diffusion coefficient is almost an order of magnitude lower compared to its corresponding  $^{19}\text{F}$  or  $^1\text{H}$  diffusion coefficients, signifying uncorrelated motion of lithium ions with the dendrimer molecules, as expected in the case of weak interactions between -CN and  $\text{Li}^+$ . On the other hand, stronger binding between -COOR and  $\text{Li}^+$  in  $\text{G}_1\text{-COOR}$ - $\text{LiPF}_6$  leads to close proximity of  $^7\text{Li}$  to the  $^1\text{H}$  diffusion coefficients at all temperatures, signifying a higher correlated motion of lithium ions with the ester molecules. The correlated motion of  $\text{Li}^+$  in  $\text{G}_1\text{-COOR}$  is further supported by the estimate of the Stokes radius for lithium ( $R_{\text{Li}} = D_{\text{H}}/D_{\text{Li}}$ ),<sup>25b</sup> which is equal to 1.4, signifying an almost 1 : 1 co-ordination between  $\text{G}_1\text{-COOR}$  and  $\text{Li}^+$ .

The temperature dependent ionic diffusivity (Fig. 4a) is fitted using the Arrhenius equation and the activation energies for  $^7\text{Li}$ ,  $^1\text{H}$  and  $^{19}\text{F}$  diffusion are tabulated in Table ST4 (ESI†). The observed trend in activation energies for  $\text{Li}^+$  and  $\text{F}^-$  diffusion is as follows:  $E_{\text{D}(\text{Li})}(\text{G}_1\text{-CN})$  ( $= 0.53 \text{ eV}$ )  $\sim E_{\text{D}(\text{Li})}(\text{G}_1\text{-COOR})$  ( $= 0.53 \text{ eV}$ ) and  $E_{\text{D}(\text{F})}(\text{G}_1\text{-CN})$  ( $= 0.50 \text{ eV}$ )  $< E_{\text{D}(\text{F})}(\text{G}_1\text{-COOR})$  ( $= 0.70 \text{ eV}$ ). This trend further suggests that the viscosity mainly influences the activation energy of  $^{19}\text{F}$  diffusion and not  $^7\text{Li}$  diffusion. Similarly, the  $\text{Li}^+$  diffusion activation energy ( $E_{\text{D}(\text{Li})}$ ) between  $\text{G}_1\text{-CN}$  and  $\text{G}_1\text{-COOR}$  leads to similar activation energy of conductivity, as obtained from ac impedance spectroscopy (Fig. 2a). Hence, the underlying mechanism of lithium conduction is similar to the lithium diffusion mechanism, and the difference in the effective conductivity is mainly determined by the differences in anion mobility. The temperature dependent viscosity (Fig. 2b) and diffusion (Fig. 4a) do not show a clear VTF-like behaviour (*i.e.* curvature-like profile). So, employing the VTF analysis will not be appropriate for analyzing and correlating the diffusion, viscosity and conductivity data together. Additionally, first-generation dendrimers are considered as viscous liquids with considerably simpler molecular architectures than the higher generation dendrimers ( $\text{G}_n$ ,  $n > 2$ ) or polymers. So, based on these aspects we considered thermally activated diffusion, viscosity and conductivity (*i.e.* Arrhenius), rather than segmental motion-driven ion transport where VTF fitting would be more appropriate.

The transference numbers of  $^7\text{Li}(t_+)$  and  $^{19}\text{F}(t_-)$  are calculated from ionic diffusion coefficients following the equation:  $t_+ = (1 - t_-) = D_+(D_+ + D_-)^{-1}$  where  $D_+$  and  $D_-$  are the cationic and anionic diffusion coefficients respectively. Fig. 4b shows the temperature dependent cationic ( $\text{Li}^+$ ) and anionic ( $\text{PF}_6^-$ ) transference numbers for both  $\text{G}_1\text{-COOR}$ -0.1 M  $\text{LiPF}_6$  and  $\text{G}_1\text{-CN}$ -0.1 M  $\text{LiPF}_6$ .  $\text{G}_1\text{-COOR}$ -0.1 M  $\text{LiPF}_6$  exhibits an extremely high lithium transference number ( $t_+$ ) of 0.9 at all experimental temperatures, whereas the  $\text{Li}^+$  transference number of  $\text{G}_1\text{-CN}$ -0.1 M  $\text{LiPF}_6$  is observed to be as low as 0.2, almost comparable to conventional PEO-based polymer electrolytes. Lower values of  $^{19}\text{F}$  diffusion coefficient compared to  $\text{Li}^+$  diffusion coefficient in the  $\text{G}_1\text{-COOR}$  electrolyte lead to an extremely high  $\text{Li}^+$  transference number, suggesting predominantly a cationic conductor. In comparison,  $\text{G}_1\text{-CN}$  presents the opposite scenario, where a very high anionic transference number ( $t_- = 0.8$ ) is observed compared to lithium, implying favorable anion transport. As discussed earlier, the viscosity and steric

hindrance of bulky peripheral  $-COOR$  groups affect the mobility of larger anions, and this trapping effect results in the extremely high  $Li^+$  transference number in  $G_1$ -COOR, in spite of the lower conductivity compared to the  $G_1$ -CN electrolyte. A significant difference in cationic transference number between  $G_1$ -CN ( $t_+ = 0.2$ ) and  $G_1$ -COOR ( $t_+ = 0.9$ ) further suggests that the present approach is highly effective in manipulating the nature of ion transport in dendrimer electrolytes. We attempted the estimation of the  $Li^+$  ion transference number using the electrochemical method proposed by Evans, Vincent and Bruce.<sup>20a,30</sup> This method has been predominantly employed to estimate the cation transference numbers of liquid and polymer electrolytes. Molecular solvent-based liquid and polymer electrolytes exhibit both cation and anion conductivity in one system. The transference number of one ion type is usually greater than the other; however, both are appreciably high and contributions to conductivity from the minority carrier cannot be neglected. So, the  $G_1$  dendrimer electrolytes do not exactly match the criteria for applicability of this method. The  $t_+$  for  $G_1$ -CN-0.1 M LiPF<sub>6</sub> is estimated to be  $\approx 0.4$ , which was higher than our estimates from NMR ( $t_+ = 0.2$ ). Thus, both electrochemical and NMR measurements conclude that the  $G_1$ -CN is an anion conductor. On the other hand, the estimated  $t_+$  of  $G_1$ -COOR-LiPF<sub>6</sub> was 0.3, instead of 0.9 as predicted from the diffusion NMR measurements. The reason for the large discrepancy between the values in  $G_1$ -COOR is due to a combination of various factors. Higher viscosity leads to slower anion kinetics during polarization. An extremely slow anion diffusion coefficient also leads to uncertainties in maintaining the necessary condition of zero anion flux in the steady state. At this juncture, the possibility of an imminent application of the novel  $G_1$  dendrimer electrolyte in an electrochemical device is remote and non-trivial. However, it is strongly envisaged that the present dendrimer and similar systems will have strong implications in various applications such as rechargeable batteries, sensors and actuators. The present dendrimers exhibit very high anion and cation transference numbers, with the conductivities being on a par with those of many polymer-based single ion conductors.<sup>10a,19</sup> The remarkably high cation transference number ( $t_+ = 0.9$ ) of  $G_1$ -COOR-LiPF<sub>6</sub> prompted us to perform electrochemical characterizations for potential application as an electrolyte or as a co-solvent in rechargeable batteries. We discuss here some of the studies which may trigger electrolyte designs based on dendrimers, in general polymeric systems for rechargeable batteries based on lithium. The cell configurations and electrode assemblies employed for the studies are exactly similar to those used for molecular based solvent electrolytes.

The electrochemical potential windows of  $G_1$ -CN-LiPF<sub>6</sub> and  $G_1$ -COOR-LiPF<sub>6</sub> electrolytes were studied *via* cyclic voltammetry with  $Li|G_1$ -dendrimer-LiPF<sub>6</sub>||SS (stainless steel) cell configuration at a scan rate of 1 mV s<sup>-1</sup>, and the results are represented in Fig. 5a and b respectively. The cyclic voltammograms clearly shows that the  $G_1$ -CN system does not support stable reversible cycling of Li, whereas  $G_1$ -COOR shows a lithium stripping peak (at  $\sim 2$  V) followed by reductive deposition too. The higher cathodic (deposition) currents observed for  $G_1$ -CN as compared to  $G_1$ -COOR are consistent with the

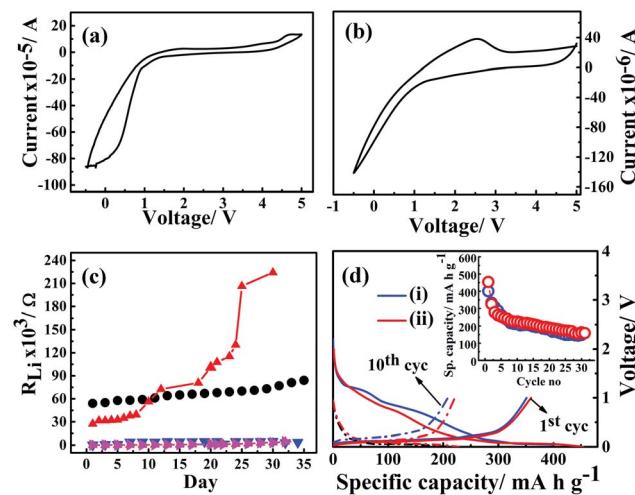


Fig. 5 Cyclic voltammograms of  $G_1$ -CN-0.1 M LiPF<sub>6</sub> (a) and  $G_1$ -COOR-0.1 M LiPF<sub>6</sub> (b) with stainless steel as working and lithium as reference and counter electrodes. (c) Lithium interface stability of  $G_1$ -CN (black circles),  $G_1$ -COOR (red triangles), 50%  $G_1$ -CN-(EC-DMC) (pink triangles), 50%  $G_1$ -COOR-(EC-DMC) (blue triangles) dendrimer electrolytes. (d) Battery cycling performance of ternary 50%  $G_1$ -CN-(EC-DMC-LiPF<sub>6</sub>) (i) (blue) and 50%  $G_1$ -COOR-(EC-DMC-LiPF<sub>6</sub>) (ii) (red) electrolytes with Li/electrolyte/graphite cell configuration at  $C/10$  constant current rate. Inset shows specific capacity ( $mA\ h\ g^{-1}$ ) vs. cycle number for both ternary electrolytes.

higher conductivity of  $G_1$ -CN compared to  $G_1$ -COOR. Improved reversibility in lithium cycling for  $G_1$ -COOR compared to  $G_1$ -CN is a consequence of the higher lithium transference number and faster ionic diffusion at the lithium electrode. Both of these dendrimers show a stable electrochemical window of 4 V as observed from cyclic voltammetry.

The electrochemical stability of  $G_1$ -COOR-LiPF<sub>6</sub> and  $G_1$ -CN-LiPF<sub>6</sub> electrolytes at lithium metal interfaces were investigated over a period of 35 days by ac impedance spectroscopy in a symmetrical  $Li|G_1$ -dendrimer|Li cell configuration, as shown in Fig. 5c. The lithium interface resistance ( $R_{Li}$ ) was evaluated from the Nyquist plots as shown in Fig. S4.<sup>†</sup> The lithium interfacial resistance ( $R_{Li}$ ) of  $G_1$ -COOR-0.1 M LiPF<sub>6</sub> on the first day ( $5.4 \times 10^4 \Omega$ ) is higher than that of  $G_1$ -CN-0.1 M LiPF<sub>6</sub> ( $2.7 \times 10^4 \Omega$ ). However, a sudden increase in interfacial resistance is observed in the case of  $G_1$ -CN-LiPF<sub>6</sub> after the 15<sup>th</sup> day. At the 30<sup>th</sup> day,  $R_{Li}$  of  $G_1$ -CN-0.1 M LiPF<sub>6</sub> increased to  $2.2 \times 10^5 \Omega$  (10 times increase in magnitude compared to day 1). On the other hand, the  $G_1$ -COOR electrolyte displayed a marginal increase over the same period reaching the value of  $7.7 \times 10^4 \Omega$  on the 30<sup>th</sup> day (1.3 times increase in magnitude). This result indicates a slower rate of growth of the passivation layer at the lithium interface for  $G_1$ -COOR-0.1 M LiPF<sub>6</sub> compared to  $G_1$ -CN-0.1 M LiPF<sub>6</sub>. The improved stability can be directly attributed to the high  $t_{Li^+}$  in  $G_1$ -COOR-0.1 M LiPF<sub>6</sub> which improves the charge transfer kinetics at the electrode|electrolyte interface. Following this, galvanostatic charge/discharge cycling measurements were performed (rate =  $C/10$ ). The pristine dendrimer-salt system *i.e.*  $G_1$ -COOR-LiPF<sub>6</sub> (and  $G_1$ -CN-LiPF<sub>6</sub>) exhibited poor galvanostatic cycling. The capacity faded to very

low values within a few cycles. We attribute the failure to the high viscosity of the pristine dendrimer–salt system which resulted in poor charge kinetics at the electrode|electrolyte interface.

Following the unsatisfactory battery cycling of the pristine dendrimer electrolytes, the cells were assembled with a mixture (by volume) containing 50% ethylene carbonate (EC)–dimethyl carbonate (DMC) (EC : DMC = 1 : 1 by v/v) and 50% of  $G_1$ –COOR ( $G_1$ –CN) and LiPF<sub>6</sub>.  $G_1$ –COOR–EC–DMC–LiPF<sub>6</sub> (as well as  $G_1$ –COOR–EC–DMC–LiPF<sub>6</sub>) exhibited a voltage stability of 3 V with stainless steel (as working electrode) and lithium foil (Aldrich) as the counter and reference electrodes (Fig. S5†). The lithium metal is passivated even better in the case of  $G_1$ –COOR–EC–DMC (or  $G_1$ –CN–EC–DMC) compared to the pristine dendrimers. No significant change in the interface resistance is observed in the case of the ternary mixture  $G_1$ –COOR–EC–DMC (and  $G_1$ –CN–EC–DMC) (Fig. 5c). This strongly suggests that the  $G_1$  dendrimers with further chemical design modifications (leading to lower viscosity) can be employed as both alternative electrolytes and electrolyte additives in conventional liquid electrolytes. With regard to the latter issue, there have been a few interesting reports on boron-based additives<sup>31</sup> aimed at stabilizing both the cathode/anode|electrolyte interfaces. While the boron-based additives aid in electron transport, the present dendrimers aid in ion conductivity. Hence, to the best of our knowledge the PETIM dendrimers are the first of their kind where the additive stabilizes the electrode|electrolyte interface *via* promotion of ion transport.

Fig. 5d represents the galvanostatic cycling performance of  $G_1$ –CN–EC–DMC–LiPF<sub>6</sub> and 50%  $G_1$ –COOR–EC–DMC–LiPF<sub>6</sub> with graphite as working electrode and lithium metal as reference and counter electrodes, respectively. The charge and discharge cycling were done at a constant current rate of  $C/10$  over a 0–2.5 V voltage range for the dendrimers. In Fig. 5b, the first discharge curve shows two distinct reductive plateaux in the ranges 0.5–0.8 V and 0.9–1.5 V corresponding to reductive degradation of EC solvent (SEI formation) *via* single and double reduction processes,<sup>32</sup> which vanish on further cycling in the lithium insertion process. The broad plateau at 0.5–0.8 V signifies decomposition of the  $G_1$ –CN molecule at the graphite surface. The charge plateau appears at 0.18 V corresponding to lithium de-insertion processes. The appearance of reductive and oxidative peaks agrees with the cyclic voltammetry results (ESI Fig. S6†).  $G_1$ –CN–EC–DMC–LiPF<sub>6</sub> shows a 1<sup>st</sup> discharge capacity of 400 mA h g<sup>−1</sup> which decreases to 330 mA h g<sup>−1</sup> in the 2<sup>nd</sup> cycle. In the 30<sup>th</sup> cycle, the capacity stabilized at 150 mA h g<sup>−1</sup>. The 1<sup>st</sup> charge capacity is equal to 354 mA h g<sup>−1</sup> and this stabilized to 148 mA h g<sup>−1</sup> in the 30<sup>th</sup> cycle. Coulombic efficiency increases from 88% (1<sup>st</sup> cycle) to 89% (5<sup>th</sup> cycle) and stabilized at 99% in the 30<sup>th</sup> cycle. The low coulombic efficiency in the 1<sup>st</sup> cycle is a consequence of irreversible capacity loss during the formation of the SEI film, which stabilizes over successive cycling leading to higher coulombic efficiency over successive cycles. Similar cycling behavior for ternary  $G_1$ –COOR–EC–DMC–LiPF<sub>6</sub> is observed except that the additional plateau at 0.9–1.5 V is absent in this case. The disappearance of this reductive plateau at 0.9–1.5 V indicates lesser

decomposition of EC (suppression of the two-electron transfer process of EC)<sup>32</sup> and better stability of the  $G_1$ –COOR dendrimer at the graphite electrode surface (clear from cyclic voltammetry in Fig. S5†). The  $G_1$ –COOR–EC–DMC–LiPF<sub>6</sub> specific capacity in the 1<sup>st</sup> discharge cycle is 453 mA h g<sup>−1</sup> which decreases to 325 mA h g<sup>−1</sup> in the 2<sup>nd</sup> cycle. In the 30<sup>th</sup> cycle, the capacity stabilized at 160 mA h g<sup>−1</sup>. The 1<sup>st</sup> charge capacity is 359 mA h g<sup>−1</sup> and stabilized at 153 mA h g<sup>−1</sup> in the 30<sup>th</sup> cycle. Coulombic efficiency increases from 79% (1<sup>st</sup> cycle) to 85% (5<sup>th</sup> cycle) and eventually stabilized at 96% in the 30<sup>th</sup> cycle. Thus, the ternary dendrimer electrolyte can be successfully cycled with a graphite electrode and may hold promise in lithium battery applications.

## Conclusion

In conclusion, we have demonstrated here a novel dendrimer–salt based ion conductor with high ion transference for prospective applications as an electrolyte in diverse devices such as rechargeable batteries, sensors and actuators. The transference number achieved here is the highest reported so far in dendrimers and polymer electrolytes. We have comprehensively demonstrated for the first time that ion conductivity and transference number can be manipulated by varying the chemical nature of the dendrimer peripheral group. The chemical nature of the peripheral group completely determines the solvation, *i.e.* quantum of free charge carriers, the mobility of the free charge carriers and the electrochemical properties. It is interesting to note that changes in peripheral chemical functional groups which affect the viscosity of the solution do not at all influence the cation diffusivity. We anticipate similar observations for other alkali ions provided the anion remains the same or bulkier than the PF<sub>6</sub><sup>−</sup> anion. In the case of other monovalent cations (K<sup>+</sup>, Rb<sup>+</sup>, Cs<sup>+</sup>), factors specific to the metal ions need to be considered to achieve similar trends in ion transport in dendrimer electrolytes. Given the advancements in polymer processing, it is strongly proposed that similar strategies as discussed here can also be adopted in designing novel solid polymers with a variety of metal salt single-ion conducting polymer electrolytes.

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## References

- 1 J. B. Goodenough and K. S. Park, *J. Am. Chem. Soc.*, 2013, 135, 1167–1176.



2 J. M. Tarascon and M. Armand, *Nature*, 2001, **414**, 359–367.

3 J. Kalhoff, G. G. Eshetu, D. Bresser and S. Passerini, *ChemSusChem*, 2015, **8**, 2154–2175.

4 C. Berthier, W. Gorecki, M. Minier, M. Armand, J. M. Chabagno and P. Rigaud, *Solid State Ionics*, 1983, **11**, 91–95.

5 D. R. Macfarlane, J. Huang and M. Forsyth, *Adv. Mater.*, 2001, **139**, 57–96.

6 P. G. Bruce, B. Scrosati and J. M. Tarascon, *Angew. Chem., Int. Ed.*, 2008, **47**, 2930–2946.

7 M. A. Ratner and D. F. Shriver, *Chem. Rev.*, 1988, **88**, 109–124.

8 (a) R. Bouchet, S. Maria, R. Meziane, A. Aboulaich, L. Lienafa, J. P. Bonnet, T. N. T. Phan, D. Bertin, D. Gigmes, D. Devaux, R. Denoyel and M. Armand, *Nat. Mater.*, 2013, **12**, 452–457; (b) P. P. Mukherjee, C. Y. Wanga and Q. Kang, *Electrochim. Acta*, 2009, **54**, 6861–6875; (c) S. Imaizumi, Y. Kato, H. Kokubo and M. Watanabe, *J. Phys. Chem. B*, 2012, **116**, 5080–5089.

9 M. J. Madou and S. R. Morrison, *Chemical Sensing with Solid State Devices*, United Kingdom, 1989.

10 (a) J. L. Schaefer, D. A. Yanga and L. A. Archer, *Chem. Mater.*, 2013, **25**, 834–839; (b) N. Boaretto, A. Bittner, C. Brinkmann, B. E. Olsowski, J. Schulz, M. Seyfried, K. Vezzu, M. Popall and V. D. Noto, *Chem. Mater.*, 2014, **26**, 6339–6350; (c) A. Ghosh, C. Wang and P. Kofinasb, *J. Electrochem. Soc.*, 2010, **157**, A846–A849.

11 (a) K. Hayamizu, Y. Aihara and W. S. Price, *J. Chem. Phys.*, 2000, **113**, 4785–4793; (b) M. Doyle, T. F. Fuller and J. Newman, *Electrochim. Acta*, 1994, **39**, 2073–2081; (c) K. E. Thomas, S. E. Sloop, J. B. Kerr and J. Newman, *J. Power Sources*, 2000, **89**, 132–138.

12 (a) K. Inoue, *Prog. Polym. Sci.*, 2000, **25**, 453–571; (b) U. P. Thankappan, S. N. Madhusudana, A. Desai, G. Jayamurugan, Y. B. R. D. Rajesh and N. Jayaraman, *Bioconjugate Chem.*, 2011, **22**, 115–119; (c) Z. Wang, M. Ikeda, N. Hirata, M. Kubo, T. Itoh and O. Yamamoto, *J. Electrochem. Soc.*, 1999, **146**, 2209–2215; (d) C. S. Harris, M. A. Ratner and D. F. Shriver, *Macromolecules*, 1987, **20**, 1778–1781; (e) A. Gong, C. Liu, Y. Chen, C. Chen and F. Xi, *Polymer*, 2000, **41**, 6103–6111.

13 (a) V. D. Noto, M. Piga, G. A. Giffin, K. Vezzu and T. A. Zawodzinski, *J. Am. Chem. Soc.*, 2012, **134**, 19099–19107; (b) V. D. Noto, N. Boaretto, E. Negro, P. E. Stallworth, S. Lavina, G. A. Giffin and S. G. Greenbaum, *Int. J. Hydrogen Energy*, 2012, **37**, 6215–6227.

14 (a) U. Boas and P. M. H. Heegard, *Chem. Soc. Rev.*, 2004, **33**, 43–63; (b) K. E. Broaders, S. Grandhe and J. M. J. Fréchet, *J. Am. Chem. Soc.*, 2011, **133**, 756–759; (c) J. S. Moore, *Acc. Chem. Res.*, 1997, **30**, 402–413; (d) M. Garzoni, N. Cheval, A. Fahmi, A. Danani and G. M. Pavan, *J. Am. Chem. Soc.*, 2012, **134**, 3349–3357; (e) D. A. Tomalia, *Soft Matter*, 2010, **6**, 456–474; (f) J. Hu, T. Xu and Y. Cheng, *Chem. Rev.*, 2012, **112**, 3856–3891.

15 (a) E. A. Dillon and D. F. Shriver, *Chem. Mater.*, 2001, **13**, 1369–1373; (b) H. Aydn, M. Senel and A. Bozkurt, *Solid State Ionics*, 2012, **226**, 1–6; (c) M. Li, X. Yang, X. Liu and X. Wang, *J. Appl. Polym. Sci.*, 2006, **101**, 317–322; (d) S. Das, B. R. Suresh, N. Jayaraman and A. J. Bhattacharyya, *J. Polym. Res.*, 2012, **19**, 9924.

16 (a) E. Staunton, Y. G. Andreev and P. G. Bruce, *Faraday Discuss.*, 2007, **134**, 143–156; (b) J. Maier, *Nat. Mater.*, 2005, **4**, 805–815; (c) S. C. Mullin, G. M. Stone, A. Panday and N. P. Balsara, *J. Electrochem. Soc.*, 2011, **158**, A619–A627.

17 (a) G. Jo, H. Ahn and M. J. Park, *ACS Macro Lett.*, 2013, **2**, 990–995; (b) A. Chiappone, S. Jeremias, R. Bongiovanni and M. Schönhoff, *Polym. Phys.*, 2013, **51**, 1571–1580; (c) B. H. Susan, A. Kaneko, A. Noda and M. Watanabe, *J. Am. Chem. Soc.*, 2005, **127**, 4976–4983; (d) D. F. Miranda, C. Versek, M. T. Tuominen, T. P. Russell and J. J. Watkins, *Macromolecules*, 2013, **46**, 9313–9323; (e) M. Kunze, Y. Karatas, H. D. Wiemhofer, H. Eckerta and M. Schonhoff, *Phys. Chem. Chem. Phys.*, 2010, **12**, 6844–6851; (f) Y. Saito, M. Okano, T. Sakai and T. Kamada, *J. Phys. Chem. C*, 2014, **118**, 6064–6068.

18 (a) J. Read, *J. Electrochem. Soc.*, 2006, **153**, A96–100; (b) M. Galiński, L. Andrzej and I. Stepienak, *Electrochim. Acta*, 2006, **51**, 5567–5580; (c) C. O. Laoire, S. Mukerjee, E. J. Plichta, M. A. Hendrickson and K. M. Abraham, *J. Electrochem. Soc.*, 2011, **158**, A302–308; (d) K. Xua and C. A. Angell, *J. Electrochem. Soc.*, 2002, **149**, A920–926.

19 (a) E. Tsuchida, N. Kobayashi and H. Ohno, *Macromolecules*, 1988, **21**, 96–100; (b) D. R. Sadoway, B. Huang, P. E. Trapa, P. P. Soo, P. Banerjee and A. M. Mayes, *J. Power Sources*, 2001, **97–98**, 621–623.

20 (a) F. Bertasi, E. Negro, K. Vezzù, G. Nawna, G. Pagot and V. D. Noto, *Electrochim. Acta*, 2015, **175**, 113–123; (b) F. Bertasi, E. Negro, K. Vezzù and V. D. Noto, *Int. J. Hydrogen Energy*, 2014, **39**, 2896–2903.

21 B. C. Smith, *Infrared Spectral Interpretation: A Systematic Approach*, United States of America, 1998.

22 (a) C. M. Burba and R. Frech, *J. Phys. Chem. B*, 2005, **109**, 15161–15164; (b) R. A. Nazri, G. A. Nazri, A. J. Camargo and M. Trisc, *J. Solution Chem.*, 2000, **29**, 1047–1060; (c) I. S. Pereygin, M. A. Klimchuk, L. V. Rabchuk and L. V. Chaukina, *Russ. J. Phys. Chem.*, 1994, **68**, 44.

23 H. Zhu, U. Rana, V. Ranganathan, J. Jin, L. A. O'Dell, D. R. MacFarlane and M. Forsyth, *J. Mater. Chem. A*, 2014, **2**, 681–691.

24 (a) K. Hayamizu, Y. Aihara and W. S. Price, *Electrochim. Acta*, 2001, **46**, 1475–1485; (b) C. Capiglia, Y. Saito, H. Yamamoto, H. Kageyama and P. Mustarelli, *Electrochim. Acta*, 2000, **45**, 1341–1345.

25 (a) M. B. Armand, P. G. Bruce, M. Forsyth, B. Scrosati and W. Wieczorek, *Polymer Electrolytes*, in *Energy Materials*, ed. D. W. Bruce, D. O'Hare and R. I. Walton, John Wiley & Sons, Ltd, Chichester, UK, 2011; (b) K. Hayamizu, Y. Aihara, S. Arai and C. G. Martinez, *J. Phys. Chem. B*, 1999, **103**, 519–524; (c) B. Fritzinger and U. S. Macromol, *Chem. Phys.*, 2005, **206**, 1288–1291.

26 (a) J. Maier, *J. Phys. Chem. Solids*, 1985, **46**, 309–320; (b) J. Maier, *Solid State Ionics*, 1987, **23**, 59–67; (c) J. Maier, *Mater. Res. Bull.*, 1985, **20**, 383–392.



27 (a) C. C. Liang, *J. Electrochem. Soc.*, 1973, **120**, 1289–1292; (b) J. Maier, *Prog. Solid State Chem.*, 1995, **23**, 171–263; (c) H. Yamada, A. J. Bhattacharyya and J. Maier, *Adv. Funct. Mater.*, 2006, **16**, 525–530; (d) J. Maier and B. Reichert, *Ber. Bunsen-Ges. Phys. Chem.*, 1986, **90**, 666–670.

28 (a) A. J. Bhattacharyya and J. Maier, *Adv. Mater.*, 2004, **16**, 811–814; (b) A. J. Bhattacharyya, J. Fleig, G. Y. Guo and J. Maier, *Adv. Mater.*, 2005, **17**, 2630–2634.

29 (a) S. K. Das and A. J. Bhattacharyya, *J. Phys. Chem. C*, 2009, **113**, 6699–6705; (b) N. Kaskhedikar, Y. Karatas, G. Cui, J. Maiera and H. D. Wiemhofer, *J. Mater. Chem.*, 2011, **21**, 11838–11843.

30 J. Evans, C. Vincent and P. G. Bruce, *Polymer*, 1987, **28**, 2324–2328.

31 Y. Qin, Z. Chen, H. S. Lee, X. Q. Yang and K. Amine, *J. Phys. Chem. C*, 2010, **114**, 15202–15206.

32 (a) S. D. Xu, O. C. Zhuang, J. Wang, Y. Q. Xu and Y. B. Zhu, *Int. J. Electrochem. Sci.*, 2013, **8**, 8058–8076; (b) C. Liu, X. Maa, F. Xu, L. Zheng, H. Zhang, W. Feng, X. Huang, M. Armand, I. Nie, H. Chen and Z. Zhou, *Electrochim. Acta*, 2014, **149**, 370–385.

