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Control of intramolecular singlet fission in a pentacene dimer by hydrostatic pressure†

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Singlet fission (SF), which produces two triplet excitons from a singlet exciton, has been identified as a novel nanointerface for efficient (photo)energy conversion. This study aims to control exciton formation in a pentacene dimer through intramolecular SF using hydrostatic pressure as an external stimulus. We reveal the hydrostatic-pressure-induced formation and dissociation processes of correlated triplet pairs (TT) in SF by means of pressure-dependent UV/vis and fluorescence spectrometry and fluorescence lifetime and nanosecond transient absorption measurements. The photophysical properties obtained under hydrostatic pressure suggested distinct acceleration of the SF dynamics by microenvironmental desolvation, the volumetric compaction of the TT intermediate based on solvent reorientation toward an individual triplet (T_1), and pressure-induced shortening of T_1 lifetimes. This study provides a new perspective on the control of SF by hydrostatic pressure as an attractive alternative to the conventional control strategy for SF-based materials.

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Introduction

Singlet fission (SF) is a reversible photophysical process in which two chromophores in the ground state (S_0) and an excited singlet state (S_1) interact to form correlated triplet pairs (${}^1\text{TT}$ and ${}^5\text{TT}$). These pairs then relax, forming two individual excitons (T_1) with an extremely high T_1 quantum yield (Φ_T) of up to 2 (Fig. 1a).^{1–3} This process is very promising and even more successful than other recently developed photophysical processes, including thermally activated delayed fluorescence (TADF),⁴ aggregation-induced emission (AIE),⁵ and upconversion,⁶ toward the construction of photo-relevant materials and optical chemosensors.^{7–9} For example, the highly efficient generation of multiple excitons allows for a wide array of applications for SF-based materials, including the construction of solar cells,¹⁰ photosensitizers,¹¹ singlet oxygen generators for photodynamic therapy, and relevant biological systems.¹² Nevertheless, the fabrication of SF-based materials requires that the important energy balance at the excited singlet and triplet levels be met ($E(S_1) \geq 2E(T_1)$),^{1–3} which in turn can hamper their wide range of molecular design. One appropriate solution for adjusting such rigidity or non-tunability of SF properties is the flexible/dynamic control achieved by external stimuli such

as temperature,^{13–16} solvent,^{17–19} or reaction (aggregation^{20–23} and supramolecular complexation^{19,24,25}). This control strategy has recently attracted significant attention although it is yet to be demonstrated on a trial-and-error basis.

Hydrostatic pressure is a mechanical isotropic stimulus and one of the most significant state quantities, enabling scientists to precisely control thermodynamic equilibria and kinetic rates, particularly in solutions in both ground and excited states.^{26–28} Thus, studies on hydrostatic pressure have been widely conducted since the early 1960s.²⁹ In the 1980s, advances were made in the understanding of photophysical and photochemical characteristic processes, such as excimer/exciplex formation,^{30,31} photoinduced electron transfer,³² and twisted intramolecular charge transfer³³ upon hydrostatic pressurization. In this way, the effects of hydrostatic pressure on the properties of solutions have been investigated on a long-term basis, in addition to solid-state high-pressure chemistry using a diamond anvil cell (DAC).^{34,35} Although such high-pressure chemistry has been extensively researched, recently, the emergence of mechanochemistry³⁶ and mechanobiology,³⁷ in which “pressure” as a mechanical force plays a key role, has brought solution-state hydrostatic pressure chemistry under the spotlight once again. We recently revealed that the TADF properties in a mechanochromic material³⁸ and the on-off AIE behavior of dynamic AIE-active polymers³⁹ can be controlled by the application of hydrostatic pressure. These results prompted us to examine how hydrostatic pressure affects or “controls” the SF processes in functional molecules.

In this study, to implement the hydrostatic pressure control concept in SF chemistry, we focused on a biphenyl-bridged

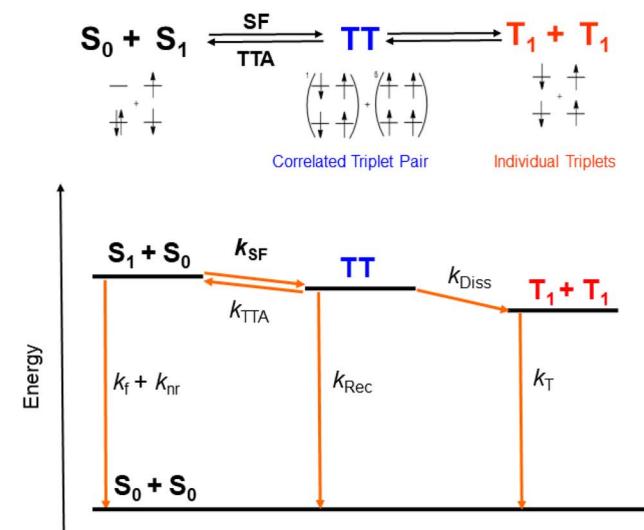
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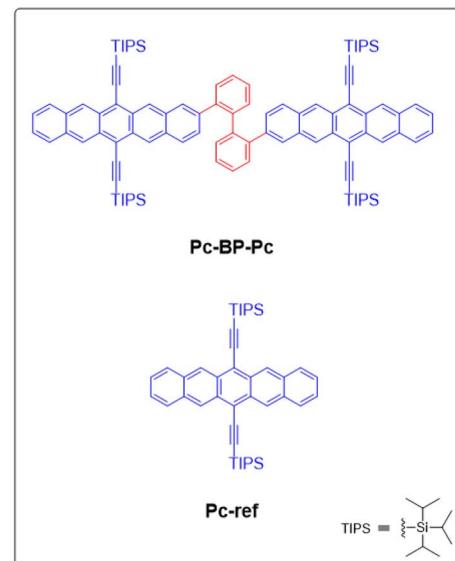
† Electronic supplementary information (ESI) available: Photographs for experimental apparatus; hydrostatic pressure spectroscopic results; ${}^1\text{H}$ NMR spectrum of **Pc-BP-Pc**. See DOI: <https://doi.org/10.1039/d3sc00312d>



(a)



(b)



(c)

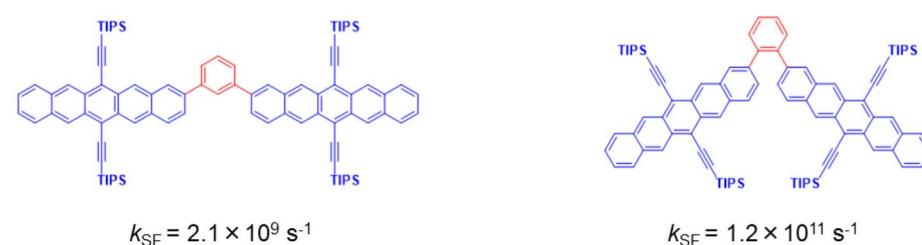


Fig. 1 (a) Schematic of the mechanism of SF in a TIPS-pentacene dimer. k_f and k_{nr} : deactivation processes from S_1 ; k_{SF} : SF from S_1 to TT; k_{Rec} : recombination process from TT to S_0 ; k_{Diss} : dissociation process from TT to $T_1 + T_1$ ($2T_1$); k_T : deactivation process from $2T_1$ to S_0 . Considering the reported energies of S_1 (1.9 eV) and TT (1.6 eV), the reverse triplet-triplet annihilation process ($TT \rightarrow S_0 + S_1$) can be omitted because of the large exergonic trend (~ 0.3 eV).⁴² (b) Target chemical structures of **Pc-BP-Pc** and **Pc-ref**. (c) Structures and their SF rates of TIPS-pentacene dimers that are linked by *m*- and *o*-phenylenes.⁴²

pentacene dimer (**Pc-BP-Pc**, Fig. 1b), which possesses a more flexible biphenyl linker than its corresponding reference monomer (**Pc-ref**, Fig. 1b). Recently, Asbury *et al.* investigated the solid-state pressure effects of **Pc-ref** using the DAC technique, revealing efficient triplet-pair separation.⁴⁰ This sophisticated case provided us with a breakthrough in the design of smart solar cells and relevant materials in the thin film state. However, control of solution-state SF processes upon the application of hydrostatic pressure, for example, photodynamic therapy in physiological (buffer) solutions, remains a major challenge in current multidisciplinary chemistry.

As a guideline for molecular design, the choice of linkers between bichromophores⁴¹ is significant for achieving the stated purpose. This significance is highlighted by the results shown in Fig. 1c: the SF kinetics (k_{SF}) in the *m*- and *o*-phenylene bridges in toluene were $2.1 \times 10^9 \text{ s}^{-1}$ and $1.2 \times 10^{11} \text{ s}^{-1}$,⁴² respectively, indicating that adjusting the distance and angle of each chromophore is critical. These results allowed us to choose a relatively flexible biphenyl linker with a k_{SF} of $1.8 \times 10^9 \text{ s}^{-1}$ in tetrahydrofuran (THF) under atmospheric conditions (0.1 MPa).¹⁷

We now report the unprecedented excited-state dynamics of **Pc-BP-Pc**, in which the SF kinetics are drastically facilitated by

hydrostatic pressurization. This study enabled us to examine the extent to which hydrostatic pressure affects the SF dynamics. The present demonstration of hydrostatic pressure control highlights the potential for a wide variety of attractive applications using SF processes in solution systems.

Results and discussion

Investigation of pressure-induced structural relaxation

First, to investigate the degree of aggregation of the pentacene dimer under hydrostatic pressure (~ 320 MPa), concentration-dependent (14–228 μM) UV/vis spectra were measured at 0.1 (atmospheric pressure), 160, and 320 MPa in toluene. As shown in Fig. S2 in ESI,[†] the absorbance increased at a constant rate with increasing concentration without any spectral changes in each hydrostatic pressure range, with the standard absorbance curves exhibiting strong linear relationships. Therefore, we confirmed that pressure-induced aggregation and crystallization based on intermolecular interactions of **Pc-BP-Pc** do not occur in these concentration ranges, which enables us to treat the pentacene analogs as well-dispersed, “monomeric” states in the following experiments.



Next, through steady-state UV/vis and fluorescence spectrometry under hydrostatic pressure, we investigated the effect of hydrostatic pressure on both ground-state absorption and excited-state fluorescence properties. As a control experiment, we performed similar tests using **Pc-ref**, a monomeric portion of the target pentacene dimer, and compared the pressure-induced spectral changes. As shown in Fig. 2a–f, under hydrostatic pressure, the spectra of both **Pc-BP-Pc** and **Pc-ref** showed gradual increases in absorbance and stepwise bathochromic shifts without significant changes in the spectral shape. This result suggests that intramolecular π -stacking does not occur in the ground state and in the excimer species in the excited state even under high pressure. It is well known that an increase in pressure induces a considerable change in solvent polarizability, causing the absorption and fluorescence peaks to

decrease in energy.²⁹ In addition, the monotonic hyperchromic effect on absorbance is simply due to the increase in the effective concentration upon pressurization. Therefore, such pressure-induced wavelength shifts can provide us with significant information about the pentacene analogs in hydrostatic-pressurized solutions; the absorption and fluorescence slopes in toluene were calculated to be -0.750 and $-0.856\text{ cm}^{-1}\text{ MPa}^{-1}$ for **Pc-BP-Pc**, and -0.848 and $-0.891\text{ cm}^{-1}\text{ MPa}^{-1}$ for **Pc-ref**, respectively, as listed in Table S1.[†] These values appear to be very similar to those observed in other π systems such as anthracene, pyrene, and perylene.²⁸ Hence, these results confirm that the intramolecular stacking behavior (*vide supra*) and substantial conformational relaxation around the biphenyl linker (particularly in the S_1 transition) may not be affected by hydrostatic pressure. Parallel experiments in methylcyclohexane (MCH) and THF showed similar slopes from -0.588 to $-0.755\text{ cm}^{-1}\text{ MPa}^{-1}$, verifying the occurrence of no significant conformational changes even in a wide variety of dipole moments (0.00–1.75 (Table S1 and Fig. S3–S10 in ESI[†])).

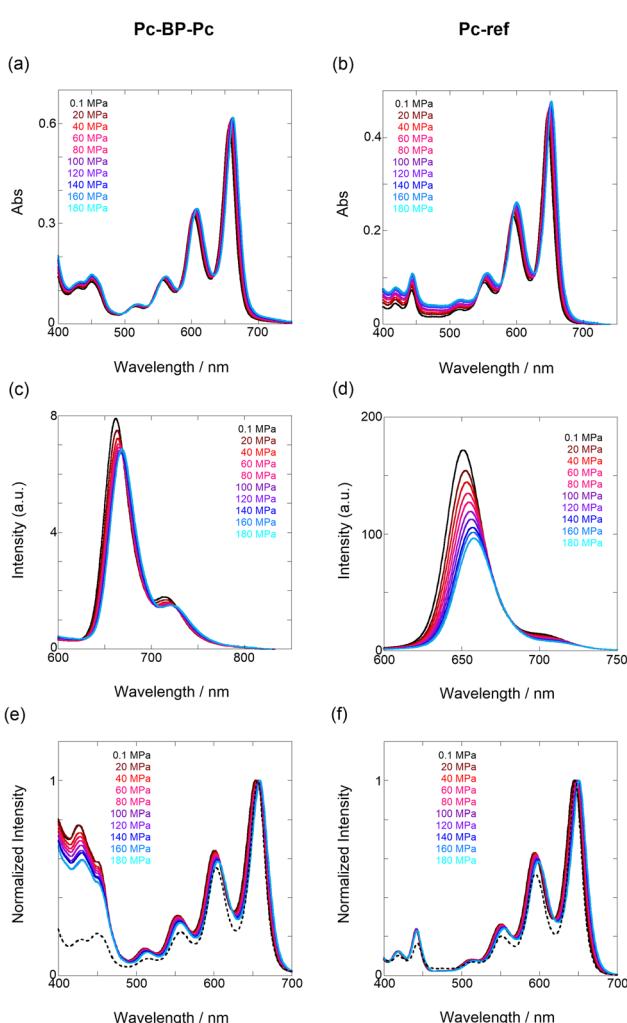


Fig. 2 (a–f) Pressure-dependent UV/vis spectra of (a) **Pc-BP-Pc** (82 μM) and (b) **Pc-ref** (91 μM); fluorescence spectra of (c) **Pc-BP-Pc** (67 μM) ($\lambda_{\text{ex}} = 580\text{ nm}$) and (d) **Pc-ref** (91 μM) ($\lambda_{\text{ex}} = 543\text{ nm}$); and excitation spectra of (e) **Pc-BP-Pc** (82 μM) ($\lambda_{\text{em}} = 660\text{ nm}$) and (f) **Pc-ref** (91 μM) ($\lambda_{\text{em}} = 650\text{ nm}$) in toluene at room temperature at 0.1, 20, 40, 60, 80, 100, 120, 140, 160, and 180 MPa (from black to sky blue), measured in a high-pressure cell. The black dotted lines represent the normalized UV/vis spectra at 0.1 MPa.

Intramolecular SF dynamics

We estimated the effects of hydrostatic pressure on the intramolecular SF kinetics of the pentacene analogs. The fluorescence lifetimes of **Pc-BP-Pc** and **Pc-ref** under hydrostatic pressure were measured in toluene (Fig. 3a and b), with the resulting decay profile obtained for **Pc-BP-Pc** containing multiple components. The profile reasonably fitted to the sum of two exponential functions, in contrast to that of the monoexponential function observed for **Pc-ref** (see Fig. S11–S16 and Tables S2–S4 in ESI[†]). Very short-lived decays (τ_2) in **Pc-BP-Pc** are observed in the enlarged figure. The flexible **Pc-BP-Pc** dimer adopts some conformers,^{17,43} in which the τ_1 species emits the fluorescence as a monomer (without the SF process) and the τ_2 species mainly decays as a deactivation path (involved the SF process). Indeed, the strong fluorescence quenching of **Pc-BP-Pc** at 0.1 MPa, rather than **Pc-ref**, was clearly observed.¹⁷ Certainly, other **Pc** dimers also showed the same decay behavior; the long-lived species (monomer conformer) and the short-lived species (TT process).⁴⁴ Thus, the rate constant of SF, $k_{\text{SF,app}}$, for the generation of correlated TTs can be written as eqn (1):

$$k_{\text{SF,app}} = (k_{\text{SF,app}} + k_0) - k_0 = \frac{1}{\tau_{2,\text{Pc-BP-Pc}}} - \frac{1}{\tau_{1,\text{Pc-ref}}} \quad (1)$$

where k_0 represents the decay rate constant from the excited singlet state **Pc-ref**, which can further be divided into the sum of k_f (fluorescence rate constant) and k_{nr} (radiationless deactivation rate constant). According to our previous study,¹⁷ the rate constant in the SF process should be expressed $k_{\text{SF,app}}$ as an apparent rate because the reverse TTA process has not been proved in this dimer. As listed in Tables 1 and S5,[†] the $k_{\text{SF,app}}$ values increased from $1.65 \times 10^9\text{ s}^{-1}$ to $1.96 \times 10^9\text{ s}^{-1}$, consistent with the elevated pressures (0.1–180 MPa) in toluene. Interestingly, a clear increasing trend of $k_{\text{SF,app}}$ from $1.23 \times 10^9\text{ s}^{-1}$ to $1.67 \times 10^9\text{ s}^{-1}$ (40–160 MPa) in the more polar THF upon hydrostatic pressurization was also observed, whereas the $k_{\text{SF,app}}$ values in nonpolar MCH were almost constant in the range of pressures

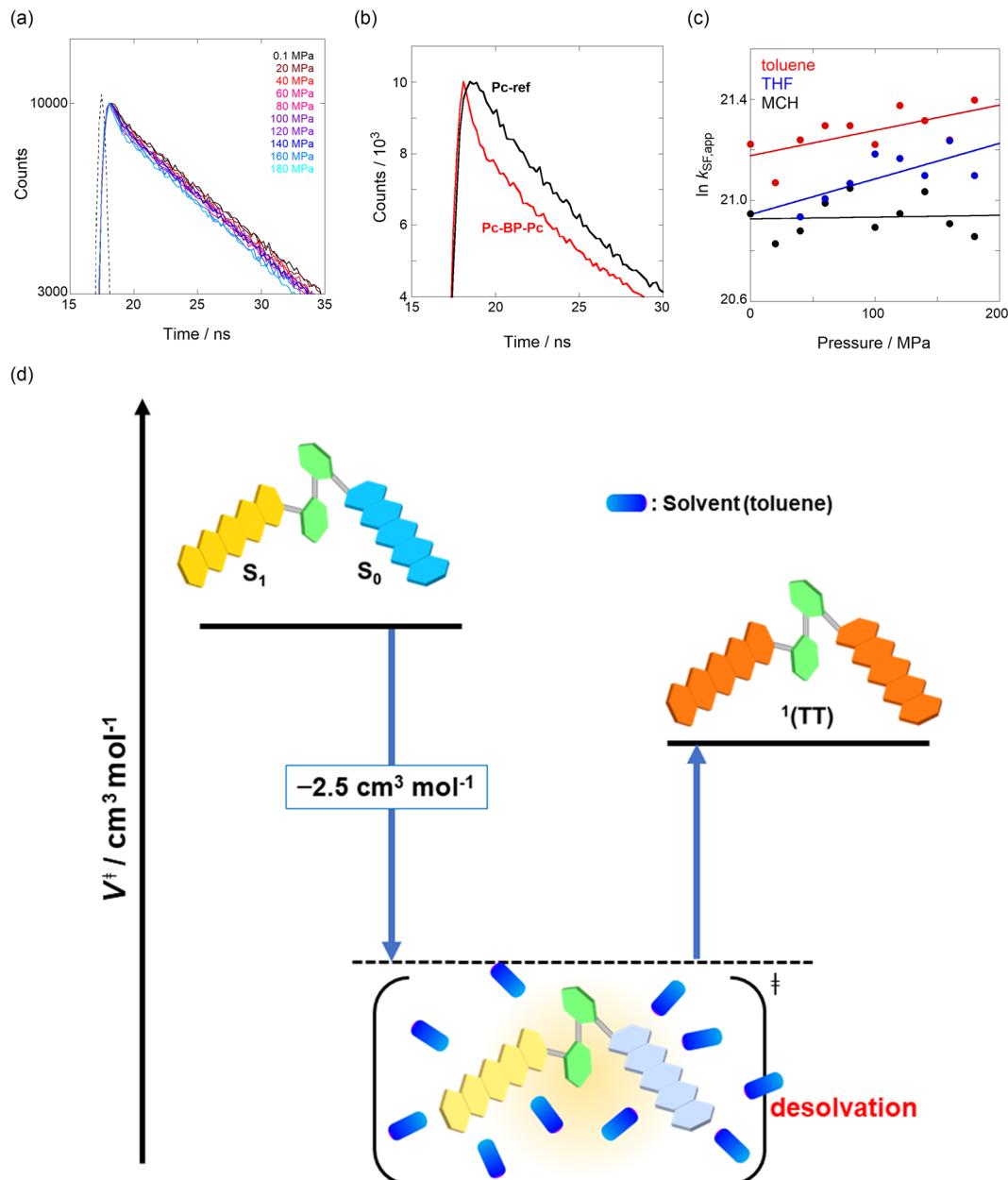


Fig. 3 (a and b) Enlarged time-correlated fluorescence decays ($\lambda_{\text{ex}} = 405 \text{ nm}$) of (a) **Pc-BP-Pc** ($74 \mu\text{M}$) ($\lambda_{\text{em}} = 660 \text{ nm}$) in toluene at room temperature at 0.1, 20, 40, 60, 80, 100, 120, 140, 160, and 180 MPa (from black to sky blue) and (b) **Pc-BP-Pc** ($74 \mu\text{M}$) ($\lambda_{\text{em}} = 660 \text{ nm}$, red) and **Pc-ref** ($89 \mu\text{M}$) ($\lambda_{\text{em}} = 650 \text{ nm}$, black) in toluene at room temperature at 180 MPa, measured in a high-pressure cell. The black dotted line represents the instrument response function. (c) Pressure dependence of SF rate constants ($k_{\text{SF,app}}$) of **Pc-BP-Pc** in toluene (red, correlation coefficient $r = 0.655$), THF (blue, $r = 0.706$), and MCH (black, $r = 0.059$) at room temperature; the separate plots with error bars are given in Fig. S26.[†] The large deviations for r are simply based on the roughness of the number of the digits in the lifetime apparatus. (d) Schematic diagram of **Pc-BP-Pc** for volumetric changes in the transition state via the intramolecular SF process.

examined. These findings indicate that hydrostatic pressure critically affects the SF process with solvent polarity.

To further elucidate the hydrostatic-pressure-induced SF dynamics more quantitatively, we calculated the activation volume change (ΔV^\ddagger) in the transition state according to eqn (2):

$$\Delta V^\ddagger = -RT \left(\frac{\partial \ln k_{\text{SF,app}}}{\partial P} \right) \quad (2)$$

The natural logarithm of $k_{\text{SF,app}}$ was plotted against pressure, with a linear relationship obtained in each solvent (Fig. 3c). This indicates that the excited-state reaction toward the transition state proceeds by only a single mechanism, that is, the SF process, in the range of the hydrostatic pressures examined. A distinct difference in the negative ΔV^\ddagger was observed for the different solvents: $-2.5 \text{ cm}^3 \text{ mol}^{-1}$ in toluene and $-3.5 \text{ cm}^3 \text{ mol}^{-1}$ in THF, compared to the almost zero ΔV^\ddagger ($-0.2 \text{ cm}^3 \text{ mol}^{-1}$) in MCH (see Table 1). Considering that ΔV^\ddagger is the degree



Table 1 Apparent rate constants of singlet fission ($k_{SF,app}$) and activation volume changes (ΔV^\ddagger) of **Pc-BP-Pc** under hydrostatic pressure^a

Solvent	Dipole moment/D	$k_{SF,app}/10^9 \text{ s}^{-1}$					
		0.1 MPa	60 MPa	120 MPa	180 MPa	$\Delta V^\ddagger/\text{cm}^3 \text{ mol}^{-1}$	
MCH	0.00	1.25	1.30	1.25	1.14	-0.2 ± 1.1	
Toluene	0.38	1.65	1.77	1.92	1.96	-2.5 ± 1.0	
THF	1.75	^b	1.33 (1.23) ^c	1.56	1.45 (1.67) ^d	-3.5 ± 1.4	

^a Measured at 298 K. ^b Not determined. ^c At 40 MPa. ^d At 160 MPa.

of compactness in a transition state depending on not only structural changes, such as differences in bond lengths and angles,²⁶ but also solvation,^{28,45} in this case relating to the intramolecular SF system, the latter solvation contribution seems to be preferable. Hence, these results can be explained in terms of the intramolecular charge density change in the correlated TT. As the desolvation process proceeds, the transition-state structure becomes more compact in polarized toluene (0.38 D) or THF (1.75 D), as shown in Fig. 3d. This scenario may be supported by a very small contribution of ΔV^\ddagger in nonpolarized MCH (0.00 D). Here, as the volume of one toluene or THF molecule is $106 \text{ cm}^3 \text{ mol}^{-1}$ or $81 \text{ cm}^3 \text{ mol}^{-1}$, the desolvation contribution to each ΔV^\ddagger is approximately 2% in one toluene molecule or 4% in one THF molecule. This indicates that intramolecular SF processes can be accelerated by critical desolvation, which allows the transition-state assembly involving the solvent cluster to be much more compact. Therefore, the solvation/desolvation contribution plays a decisive role in the precise control of intramolecular SF dynamics.

Hydrostatic pressure effects of triplets

We elucidated the hydrostatic-pressure-induced intramolecular SF processes, at which the active exciton, TT, eventually becomes T₁. Nanosecond transient absorption (nsTA) spectrometry is an effective analytical tool for elucidating the photophysical properties of T₁ as a final step of SF.¹⁷ The construction of a newly designed optical system for hydrostatic pressure nsTA spectrometry is detailed in the Materials and methods section. This system was used to measure the solvent-dependent spectra of toluene, MCH, and THF solutions in the triplet absorption band (Fig. 4a). As shown in Fig. 4b and S17a–c,[†] spectral measurements at 0.1 (atmospheric pressure), 160, and 320 MPa also showed pressure-dependent decays and slight pressure-induced bathochromic shifts, similar to those observed in the steady-state UV-vis absorption spectra upon hydrostatic pressurization. In addition, the nsTA decay profiles of the generated individual T₁ at each hydrostatic pressure in Fig. 4c and S17d[†] can be reasonably fitted to a monoexponential function with triplet lifetimes (τ_T) of 0.7–1.6 μs (Fig. S18–S19 and Table S6[†]). So far, we have revealed that τ_T of **Pc-BP-Pc** can be prolonged from 0.36 μs in THF to 1.0 μs by mixing paraffin into THF, based on the suppression of collisional deactivation by the polar solvent.¹⁷ We achieved precise hydrostatic pressure control of τ_T without changing the solvent, in which the shortened lifetimes in response to hydrostatic pressurization are responsible for T₁ deactivation under higher

viscosity solvent conditions, that is, acceleration of solvent interactions with T₁ excitons (*vide infra*: result of the bulkier T₁ structure). As shown in Fig. S27,[†] the pressure-induced viscosity changes play significant roles in the excited-state processes rather than the polarizability, density, and polarity; the latter ones almost affect to the absorption behavior.²⁹ Indeed, Lacour and Vauthey *et al.* demonstrated the SF control by solvent viscosity.¹⁹

Finally, the S₁, TT, and T₁ quantum yields (Φ_S , Φ_{TT} , and Φ_{T_1}) were calculated (see the Materials and methods section), enabling us to estimate the SF thermodynamics, that is, the reaction volume changes as $\Delta V_{S \rightarrow TT}^\ddagger$ in the equilibrium between S₁ and TT (eqn (3)) and $\Delta V_{TT \rightarrow T_1}^\ddagger$ in the equilibrium between TT and T₁ (eqn (4)).

$$\left(\frac{\partial \ln(\Phi_{TT}/\Phi_S)}{\partial P} \right) = -\frac{\Delta V_{S \rightarrow TT}^\ddagger}{RT} \quad (3)$$

$$\left(\frac{\partial \ln(\Phi_{T_1}/\Phi_{TT})}{\partial P} \right) = -\frac{\Delta V_{TT \rightarrow T_1}^\ddagger}{RT} \quad (4)$$

As shown in Table 2 and Fig. S22,[†] the values of $\Delta V_{S \rightarrow TT}^\ddagger$ in toluene and MCH were negative and almost zero, respectively, indicating that the TT structure comprising the solvent core was thermodynamically more compact than the S₁ structure upon desolvation. This is consistent with the fact that the kinetically formed TT transition-state complex is also more compact owing to the desolvation-driven behavior (*vide supra*). In contrast, the values in Table 2 and Fig. S24[†] for $\Delta V_{TT \rightarrow T_1}^\ddagger$ were positive. Strangely, these values in both toluene and MCH are approximately $+5.8 \text{ cm}^3 \text{ mol}^{-1}$; such a solvent-independent trend suggests that a different mechanism is at play here, rather than the aforementioned solvation/desolvation. These results are likely attributable to the change in conformers in **Pc-BP-Pc** during the TT dissociation process (TT \rightarrow 2T₁), as shown in Fig. 4d. In a previous study on the intramolecular SF of pentacene dimers using time-resolved electron paramagnetic resonance measurements, TT dissociation motion was observed, and the pentacene dimer exhibited a dynamic change in the dihedral angle between chromophores.⁴³ Similarly, **Pc-BP-Pc** is highly likely to undergo a conformational change associated with the dihedral angle change around the biphenyl linker during the TT dissociation process even under hydrostatic pressure. This was the case with the D-A-D triad that showed the excited-state conformational change under hydrostatic pressure.³⁸ Hence, the nature of $\Delta V_{TT \rightarrow T_1}^\ddagger$ may be attributed to the solvent reorientation/interaction occurring after the



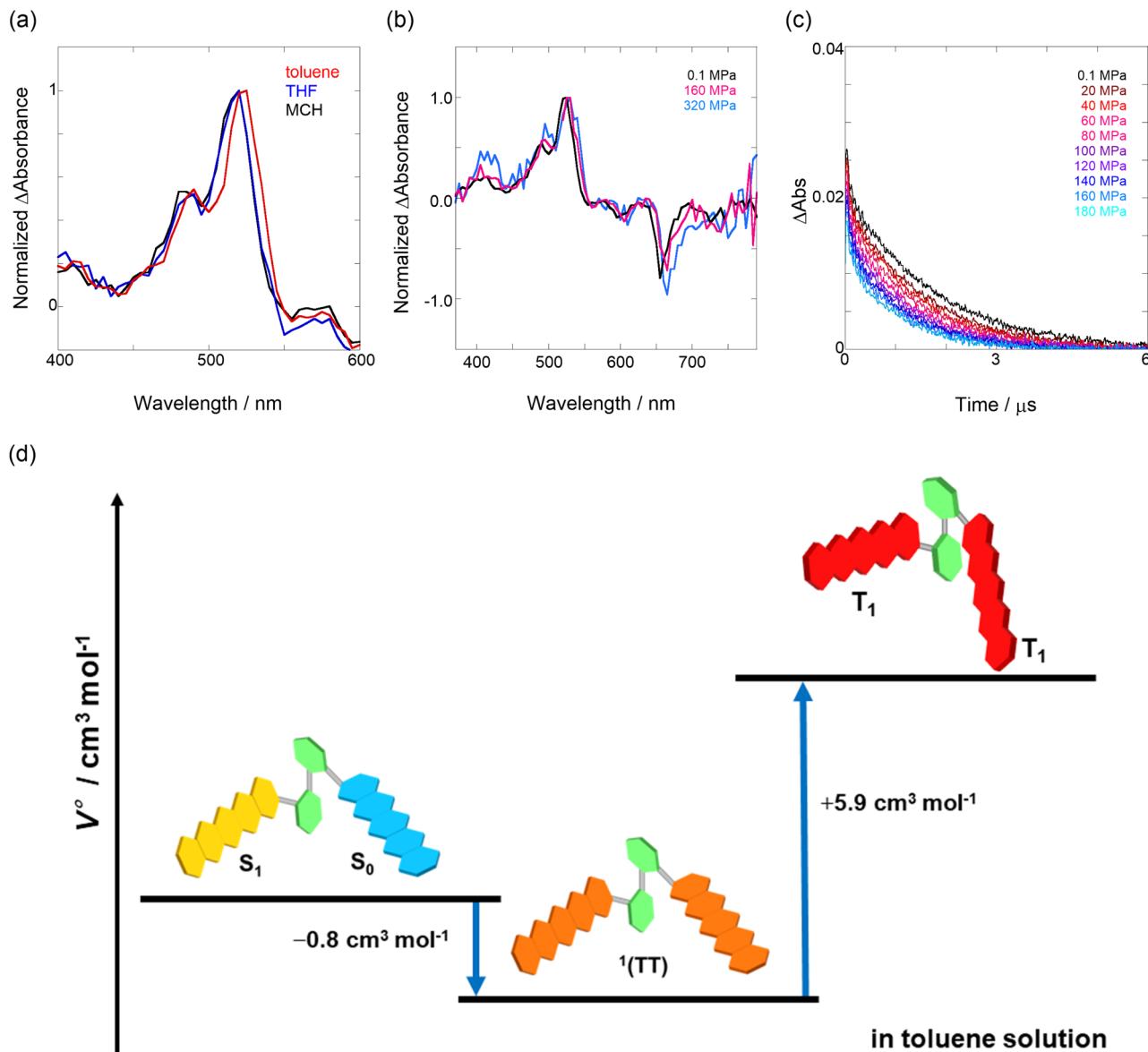


Fig. 4 (a and b) Normalized nsTA spectra of **Pc-BP-Pc** (81 μM) in (a) toluene (red), THF (blue), and MCH (black) at 0.1 MPa and (b) toluene at 0.1 (black), 160 (pink, the line at 510–520 nm was omitted due to an artifact), and 320 MPa (blue) at room temperature. (c) nsTA decay profiles of **Pc-BP-Pc** in toluene (81 μM , $\lambda_{\text{ex}} = 532 \text{ nm}$, $\lambda_{\text{obs}} = 521 \text{ nm}$) at 0.1, 20, 40, 60, 80, 100, 120, 140, 160, and 180 MPa (from black to sky blue). (d) Schematic of the volumetric changes in the equilibria *via* the intramolecular SF and dissociation processes for **Pc-BP-Pc**.

Table 2 Singlet, correlated triplet pair, triplet quantum yields (Φ_S , Φ_{TT} , and Φ_T) and reaction volume changes (ΔV°) of **Pc-BP-Pc** in toluene and MCH under hydrostatic pressure^a

Pressure/MPa	Toluene			MCH		
	Φ_S	Φ_{TT}^b	Φ_T^c	Φ_S	Φ_{TT}^b	Φ_T^c
0.1	0.014	0.955	1.762	0.051	0.950	1.120
60	0.013	0.958	1.671	0.050	0.951	1.055
120	0.013	0.961	1.398	0.050	0.951	0.912
180	0.013	0.962	1.153	0.051	0.947	0.717
$\Delta V^\circ_{S \rightarrow TT} / \text{cm}^3 \text{ mol}^{-1}$	-0.8 ± 0.3			-0.1 ± 0.2		
$\Delta V^\circ_{TT \rightarrow T} / \text{cm}^3 \text{ mol}^{-1}$	$+5.9 \pm 0.4$			$+5.8 \pm 0.5$		

^a Measured at 298 K. ^b Maximum $\Phi_{TT} = 1$. ^c Maximum $\Phi_T = 2$.

conformational change. Because the TT dissociation process is spatially divided by spin interactions, in contrast to the TT formation (*via* the exciton coupling of S_0 and S_1), a possible interpretation of this result is that the dependence of the solvation term (toluene *vs.* MCH) may be extremely small. The contribution of the conformational change during dissociation is further related to the stepwise reduction of Φ_T with increasing hydrostatic pressure (Table 2), which is responsible for the gradually enlarged solvent interactions in **Pc-BP-Pc** due to the increase in hydrostatic-pressure-driven solvent viscosity. This fact strongly supports the relationship among thermodynamically expanded or “bulkier” structures, positive $\Delta V^\circ_{TT \rightarrow T}$, and the gradual τ_T shortening trend. In detail, the thermodynamically bulkier T_1 exciton by clustering a large number of solvent

molecules is strongly deactivated by the solvent attack in the solvent core to shorten τ_T .

Conclusions

In this study, for the first time, we realized the significance of hydrostatic-pressure-induced intramolecular SF behavior. This was demonstrated using the biphenyl-linked pentacene dimer as a model SF material, in which hydrostatic-pressure-based solvent-induced property changes are key factors. Although remarkable conformational changes in the dimer were not observed, through hydrostatic pressure steady-state spectrometry, the fluorescence and triplet absorption lifetime measurements enabled us to recognize that the rates of the correlated triplet pairs from the singlet were greatly accelerated by changing the hydrostatic pressure in toluene only. This revealed that the desolvation process in polar solvents plays an important role in the SF dynamics. More importantly, we have shown that the entire process of SF involving fission ($S_0 + S_1 \rightarrow TT$) and dissociation ($TT \rightarrow 2T_1$) under hydrostatic pressure can be precisely controlled by not only kinetics in the transition state but also thermodynamics in the equilibria on the basis of microenvironmental desolvation and solvent reorientation. Finally, it should be emphasized that by using hydrostatic pressure as an external stimulus, the dynamic control concept of intramolecular SF kinetics observed in this study can be further expanded to other SF scaffolds and relevant systems that are difficult to control in both ground and excited states.

Materials and methods

Materials

All commercial reagents and solvents were used without further purification. Sample solutions dissolved in spectroscopic grade toluene, methylcyclohexane (MCH), and tetrahydrofuran (THF) were deaerated by five freeze–pump–thaw cycles saturated with N_2 for fluorescence lifetime measurements or by Ar bubbling for nanosecond transient absorption (nsTA) measurements. The SF-based material (**Pc-BP-Pc**) was synthesized according to the literature.¹⁷ **Pc-ref** was commercially available.

Instruments

UV/vis and fluorescence spectra were recorded in a high-pressure cell (path length: 2 mm) by using a JASCO V-650 or a JASCO FP-8500. Fluorescence lifetimes were measured in a high-pressure cell by a Hamamatsu Quantaurus-Tau single photon counting apparatus fitted with an LED light source ($\lambda_{ex} = 405$ nm). Nanosecond transient absorption (nsTA) measurements were performed by using a Unisoku TSP-2000 flash spectrometer-pump pulse source: Surelite-I Nd:YAG (Continuum, 4–6 ns fwhm) laser with the second harmonic at 532 nm, monitor light source: xenon lamp (150 W), and detector: photomultiplier tube. 1H NMR spectrum of **Pc-BP-Pc** was recorded on an ECS-400 spectrometer.

Hydrostatic pressure spectroscopy

Steady-state UV/vis absorption/excitation fluorescence spectroscopy and fluorescence lifetime decay measurement were conducted using a custom-built high-pressure apparatus.^{28,45} As the process has previously reported in detail in a previous study, here, we briefly describe this. A quartz inner cell was filled with the sample solution, and then the cell was set into the outer cell, where sapphire or quartz windows were fitted. A tightly closed outer cell, which was hydrostatically pressurized by water, was placed in the spectroscopic apparatus (Fig. S1a–e†).

Hydrostatic pressure nsTA

The outer cell was placed on the nsTA optical system, which allowed us to apply hydrostatic pressure in an appropriate manner (Fig. S1f†). The T_1 quantum yield (Φ_T) of a toluene solution of **Pc-BP-Pc** measured in this optical system under an atmospheric pressure (0.1 MPa) was 176%, which is identical to that ($\Phi_T = 176\%$)¹⁷ of a THF solution observed in a regular cuvette at 0.1 MPa. This proves the validity of the nsTA measurements under hydrostatic pressure using the new optical system prepared in this study.

Determination of S_1 quantum yield (Φ_S) of **Pc-BP-Pc**

Absolute S_1 quantum yields Φ_{P_0} (Φ_S at 0.1 MPa; 0.0140 in toluene, 0.0514 in MCH) were evaluated using a spectrofluorometer (FP-8500) fitted with an integrating sphere at 0.1 MPa, and then relative S_1 quantum yields can be determined using the following eqn (5):⁴⁶

$$\Phi_S = \Phi_{P_0} \frac{A_{P_0}}{A_P} \frac{I_{P_0}}{I_P} \frac{n_P^2}{n_{P_0}^2} \frac{D_P}{D_{P_0}} \quad (5)$$

where A , I , n , and D are the absorbance at the excitation wavelength, intensity of the excitation light, refractive index, and area ratio in the fluorescence spectrum, respectively (Tables S7 and S8†). The data for A were extracted from Fig. 2a and S5a.† The data for I_P/I_{P_0} were 1 because the measurements were performed at the same excitation wavelength (Fig. 2c and S6a†). The refractive indices used under atmospheric pressure (n_{P_0}) were identical to the values reported in a previous study,⁴⁷ and the refractive indices under pressure (n_P) were calculated using the Eykman equation (eqn (6)) from the change in the density of each solvent with respect to pressure:^{48,49}

$$\frac{n^2 - 1}{n + 0.4} \frac{1}{d} = C \quad (6)$$

Constant C was calculated using a known value.⁴⁷ The data for D were obtained by fitting the area of the fluorescence spectra in toluene (Fig. 2c) and MCH (Fig. S6a†) using the MATLAB software with a three-component Gaussian function (eqn (7)) (see Fig. S20 and S21 and Tables S9 and S10†):

$$A_P(x) = \int \sum_{i=1}^n a_i \exp\left(-((x - b_i)/c_i)^2\right) \quad (7)$$

Determination of the TT quantum yield (Φ_{TT}) of **Pc-BP-Pc**

The Φ_{TT} values were calculated using eqn (8) (see Table S11[†]):

$$\frac{k_{\text{SF,app}}}{k_0 + k_{\text{SF,app}}} = \Phi_{\text{TT}} \quad (8)$$

where the data for k_0 and $k_{\text{SF,app}}$ were extracted from Tables S2, S3 and S5,[†] respectively.

Determination of T_1 quantum yield (Φ_{T}) of **Pc-BP-Pc**

The $\Phi_{\text{T,Pc-BP-Pc}}$ value was calculated as the relative Φ_{T} quantum yield (see Tables S12 and S13[†]) in relation to that of zinc tetraphenylporphyrin (ZnTPP) or zinc tetratolylporphyrin (ZnTPP), using the following eqn (9):¹⁷

$$\Phi_{\text{T,Pc-BP-Pc}} = \frac{\Delta A_{\text{Pc-BP-Pc}}}{\Delta A_{\text{ZnTPP}}} \frac{\varepsilon_{\text{T,ZnTPP}}}{\varepsilon_{\text{T,Pc-BP-Pc}}} \frac{\text{Abs}_{(532 \text{ nm}, \text{ZnTPP})}}{\text{Abs}_{(532 \text{ nm}, \text{Pc-BP-Pc})}} \Phi_{\text{T,ZnTPP}} \quad (9)$$

where Φ_{T} , ε_{T} , ΔA and Abs are the T_1 quantum yield, excitation coefficient, delta absorbance in nsTA, and absorbance at 532 nm in the steady-state absorption measurements, respectively. $\Phi_{\text{T,ZnTPP}}$ and $\varepsilon_{\text{T,ZnTPP}}$ were acquired from a previous study⁵⁰ ($\lambda_{\text{ex}} = 532 \text{ nm}$, $\lambda_{\text{obs}} = 470 \text{ nm}$ in toluene; $\Phi_{\text{T,ZnTPP}} = 0.88$, $\varepsilon_{\text{T,ZnTPP}} = 69000 \text{ M}^{-1} \text{ cm}^{-1}$), and $\varepsilon_{\text{T,Pc-BP-Pc}}$ from another paper¹⁷ ($\lambda_{\text{ex}} = 532 \text{ nm}$, $\lambda_{\text{obs}} = 516 \text{ nm}$ in THF; $\varepsilon_{\text{T,ZnTPP}} = 54900 \text{ M}^{-1} \text{ cm}^{-1}$). ΔA_{ZnTPP} was observed at 470 nm and $\Delta A_{\text{Pc-BP-Pc}}$ was observed at 521 nm for toluene and 516 nm for MCH, taking into account the solvent-induced shift of the excited-state absorption T_1-T_n band (Fig. 4a). Data of $\text{Abs}_{(532 \text{ nm}, \text{ZnTPP})}$ and $\text{Abs}_{(532 \text{ nm}, \text{Pc-BP-Pc})}$ are shown in Fig. S23b, 2a, and S5a,[†] respectively.

Data availability

The data supporting this article have been uploaded as part of the ESI.[†]

Author contributions

G. F. initiated and supervised the whole study. G. F. and T. H. designed the project and the experiments. T. K. performed the spectroscopic experiments. S. N. prepared the materials. M. H. designed and prepared the foundation for the transient absorption spectrometry under hydrostatic pressure. All authors contributed to writing the manuscript.

Conflicts of interest

There are no conflicts to declare.

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References

- 1 M. B. Smith and J. Michl, *Chem. Rev.*, 2010, **110**, 6891–6936.
- 2 K. Miyata, F. S. Conrad-Burton, F. L. Geyer and X.-Y. Zhu, *Chem. Rev.*, 2019, **119**, 4261–4292.
- 3 T. Ullrich, D. Munz and D. M. Guldi, *Chem. Soc. Rev.*, 2021, **50**, 3485–3518.
- 4 Z. Yang, Z. Mao, Z. Xie, Y. Zhang, S. Liu, J. Zhao, J. Xu, Z. Chi and M. P. Aldred, *Chem. Soc. Rev.*, 2017, **46**, 915–1016.
- 5 J. Mei, N. L. C. Leung, R. T. K. Kwok, J. W. Y. Lam and B. Z. Tang, *Chem. Rev.*, 2015, **115**, 11718–11940.
- 6 N. Yanai and N. Kimizuka, *Acc. Chem. Res.*, 2017, **50**, 2487–2495.
- 7 T. W. Bell and N. M. Hext, *Chem. Soc. Rev.*, 2004, **33**, 589–598.
- 8 L. You, D. Zha and E. V. Anslyn, *Chem. Rev.*, 2015, **115**, 7840–7892.
- 9 G. Fukuhara, *J. Photochem. Photobiol. C*, 2020, **42**, 100340.
- 10 M. Einzinger, T. Wu, J. F. Kompalla, H. L. Smith, C. F. Perkinson, L. Nienhaus, S. Wieghold, D. N. Congreve, A. Kahn, M. G. Bawendi and M. A. Baldo, *Nature*, 2019, **571**, 90–94.
- 11 P. D. Reusswig, D. N. Congreve, N. J. Thompson and M. A. Baldo, *Appl. Phys. Lett.*, 2012, **101**, 113304.
- 12 T. Saegusa, H. Sakai, H. Nagashima, Y. Kobori, N. V. Tkachenko and T. Hasobe, *J. Am. Chem. Soc.*, 2019, **141**, 14720–14727.
- 13 J. J. Burdett, D. Gosztola and C. J. Bardeen, *J. Chem. Phys.*, 2011, **135**, 214508.
- 14 M. W. B. Wilson, A. Rao, K. Johnson, S. Gélinas, R. di Pietro, J. Clark and R. H. Friend, *J. Am. Chem. Soc.*, 2013, **135**, 16680–16688.
- 15 S. Nakamura, H. Sakai, M. Fuki, Y. Kobori, N. V. Tkachenko and T. Hasobe, *J. Phys. Chem. Lett.*, 2021, **12**, 6457–6463.
- 16 X. Zhao, J. P. O'Connor, J. D. Schultz, Y. J. Bae, C. Lin, R. M. Young and M. R. Wasielewski, *J. Phys. Chem. B*, 2021, **125**, 6945–6954.
- 17 T. Sakuma, H. Sakai, Y. Araki, T. Mori, T. Wada, N. V. Tkachenko and T. Hasobe, *J. Phys. Chem. A*, 2016, **120**, 1867–1875.
- 18 A. M. Alvertis, S. Lukman, T. J. H. Hele, E. G. Fuemmeler, J. Feng, J. Wu, N. C. Greenham, A. W. Chin and A. J. Musser, *J. Am. Chem. Soc.*, 2019, **141**, 17558–17570.
- 19 A. Aster, F. Zinna, C. Rumble, J. Lacour and E. Vauthey, *J. Am. Chem. Soc.*, 2021, **143**, 2361–2371.
- 20 C. Grieco, G. S. Doucette, R. D. Pensack, M. M. Payne, A. Rimshaw, G. D. Scholes, J. E. Anthony and J. B. Asbury, *J. Am. Chem. Soc.*, 2016, **138**, 16069–16080.
- 21 M. Dvořák, S. K. K. Prasad, C. B. Dover, C. R. Forest, A. Kaleem, R. W. MacQueen, A. J. Petty II, R. Forecast, J. E. Beves, J. E. Anthony, M. J. Y. Tayebjee, A. Widmer-Cooper, P. Thordarson and T. W. Schmidt, *J. Am. Chem. Soc.*, 2021, **143**, 13749–13758.
- 22 J. Zhang, H. Sakai, K. Suzuki, T. Hasobe, N. V. Tkachenko, I.-Y. Chang, K. Hyeon-Deuk, H. Kaji, T. Teranishi and M. Sakamoto, *J. Am. Chem. Soc.*, 2021, **143**, 17388–17394.



23 H. Sakai, K. Yoshino, Y. Shoji, T. Kajitani, J. Pu, T. Fukushima, T. Takenobu, N. V. Tkachenko and T. Hasobe, *J. Phys. Chem. C*, 2022, **126**, 9396–9406.

24 A. J. Tilley, R. D. Pensack, E. L. Kynaston, G. D. Scholes and D. S. Seferos, *Chem. Mater.*, 2018, **30**, 4409–4421.

25 L. Catti, H. Narita, Y. Tanaka, H. Sakai, T. Hasobe, N. V. Tkachenko and M. Yoshizawa, *J. Am. Chem. Soc.*, 2021, **143**, 9361–9367.

26 A. Drlića, C. D. Hubbard, R. van Eldik, T. Asano, M. V. Basilevsky and W. J. le Noble, *Chem. Rev.*, 1998, **98**, 2167–2289.

27 J. L. Silva, A. C. Oliveira, T. C. R. G. Vieira, G. A. P. de Oliveira, M. C. Suarez and D. Foguel, *Chem. Rev.*, 2014, **114**, 7239–7267.

28 H. Mizuno and G. Fukuhara, *Acc. Chem. Res.*, 2022, **55**, 1748–1762.

29 F. A. Bovey and S. S. Yanari, *Nature*, 1960, **186**, 1042–1044.

30 P. C. Johnson and H. W. Offen, *J. Chem. Phys.*, 1972, **56**, 1638–1642.

31 K. Hara and K. Obara, *Chem. Phys. Lett.*, 1985, **117**, 96–98.

32 K. Hara, T. Arese and J. Osugi, *J. Am. Chem. Soc.*, 1984, **106**, 1968–1972.

33 K. Hara, H. Suzuki and W. Rettig, *Chem. Phys. Lett.*, 1988, **145**, 269–272.

34 R. Lee, J. A. K. Howard, M. R. Probert and J. W. Steed, *Chem. Soc. Rev.*, 2014, **43**, 4300–4311.

35 E. F. O'Bannon III, Z. Jenei, H. Cynn, M. J. Lipp and J. R. Jeffries, *Rev. Sci. Instrum.*, 2018, **89**, 111501.

36 Y. Sagara and T. Kato, *Nat. Chem.*, 2009, **1**, 605–610.

37 M. Krieg, G. Fläschner, D. Alsteens, B. M. Gaub, W. H. Roos, G. J. L. Wuite, H. E. Gaub, C. Gerber, Y. F. Dufrêne and D. J. Müller, *Nat. Rev. Phys.*, 2019, **1**, 41–57.

38 Y. Takeda, H. Mizuno, Y. Okada, M. Okazaki, S. Minakata, T. Penfold and G. Fukuhara, *ChemPhotoChem*, 2019, **3**, 1203–1211.

39 K. Nakasha and G. Fukuhara, *ACS Appl. Polym. Mater.*, 2020, **2**, 2303–2310.

40 G. S. Doucette, H.-T. Huang, J. M. Munro, K. T. Munson, C. Park, J. E. Anthony, T. Strobel, I. Dabo, J. V. Badding and J. B. Asbury, *Cell Rep. Phys. Sci.*, 2020, **1**, 100005.

41 S. Nakamura, H. Sakai, H. Nagashima, M. Fukui, K. Onishi, R. Khan, Y. Kobori, N. V. Tkachenko and T. Hasobe, *J. Phys. Chem. C*, 2021, **125**, 18287–18296.

42 H. Sakai, R. Inaya, H. Nagashima, S. Nakamura, Y. Kobori, N. V. Tkachenko and T. Hasobe, *J. Phys. Chem. Lett.*, 2018, **9**, 3354–3360.

43 Y. Kobori, M. Fukui, S. Nakamura and T. Hasobe, *J. Phys. Chem. B*, 2020, **124**, 9411–9419.

44 S. Lukman, K. Chen, J. M. Hodgkiss, D. H. P. Turban, N. D. M. Hine, S. Dong, J. Wu, N. C. Greenham and A. J. Musser, *Nat. Commun.*, 2016, **7**, 13622.

45 A. J.-L. Ayitou, G. Fukuhara, E. Kumarasamy, Y. Inoue and J. Sivaguru, *Chem.-Eur. J.*, 2013, **19**, 4327–4334.

46 J. N. Demas and G. A. Crosby, *J. Phys. Chem.*, 1971, **75**, 991–1024.

47 B. González, I. Domínguez, E. J. González and Á. Domínguez, *J. Chem. Eng. Data*, 2010, **55**, 1003–1011.

48 M. J. Assael, H. M. T. Avelino, N. K. Dalaouti, J. M. N. A. Fareleira and K. R. Harris, *Int. J. Thermophys.*, 2001, **22**, 789–799.

49 M. Dakkach, G. Rubio-Pérez, F. E. M. Alaoui, N. Muñoz-Rujas, F. Aguilar and E. A. Montero, *J. Chem. Eng. Data*, 2020, **65**, 4892–4904.

50 T. H. Tran-Thi, C. Desforge, C. Thiec and S. Gaspard, *J. Phys. Chem.*, 1989, **93**, 1226–1233.

