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# High recorded color rendering index in single Ce,(Pr,Mn):YAG transparent ceramics for high-power white LEDs/LDs†

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Transparent ceramics (TCs) are incredibly promising color converters for high-power white LEDs/LDs. However, the preparation process of multiple structured TCs with a high color rendering index (CRI) is a complicated technical challenge, and the inability of single-structured TCs to achieve a high CRI significantly limits their real applications. In this study, high quality single-structured Ce,(Pr,Mn):YAG TCs with “wide peak” and “narrow peak” red light emissions were designed and fabricated via a solid-state reaction and a vacuum sintering method. Compared with the emission spectra of Ce:YAG TC, the synchronous doping of Pr<sup>3+</sup> and Mn<sup>2+</sup> ions into Ce:YAG TC resulted in inhomogeneous broadening of the full width at half maximum (FWHM) from 91.7 nm to 102.2 nm. Impressively, the CRI of a single Ce,(Pr,Mn):YAG TC-based high-power white LED was as high as 84.8, and the correlated-color temperatures (CCTs) of the white LEDs/LDs were 5450 K and 3550 K, respectively. Furthermore, when the addition amounts of Pr<sup>3+</sup> and Mn<sup>2+</sup> were 0.2 at% and 0.8 at%, respectively, the as-prepared Ce,(Pr,Mn):YAG TC displayed a high quantum efficiency (IQE = 48.14%) and excellent color stability (only 5% fluctuation). Therefore, this study not only demonstrates how to overcome the spectrum deficiencies of single-structured TCs that restrain intrinsic CRI improvement, but also provides a reference for the pursuit of high luminescence properties. This work significantly reinforces the understanding of the CRI problems of TC-based high-power lighting, which is crucial for the real application of white LEDs/LDs.

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## 1 Introduction

In recent years, white lighting emitting diodes and laser diodes (LEDs/LDs) have been considered as a new generation of light sources owing to their long lifetimes, energy savings and environmentally friendly nature.<sup>1–3</sup> Currently, the most commonly used color converting material for white LEDs/LDs is the

Ce<sup>3+</sup>:Y<sub>3</sub>Al<sub>5</sub>O<sub>12</sub> (Ce:YAG) phosphor. Unfortunately, with the rising power of white LEDs and LDs, this phosphor usually suffers from decreasing luminous efficiency and chromaticity shifts due to heat accumulation. Therefore, color converters such as glass ceramics or phosphors in glass (PiG),<sup>4–7</sup> single crystals (SCs),<sup>8–10</sup> and transparent ceramics (TCs)<sup>11–13</sup> have been proposed to replace the applied phosphor due to their facile implementation of the remote encapsulation mode. Compared with PiG and SCs, Ce:YAG TCs are developing rapidly due to their low cost, ease of achieving high doping concentration, short production cycles, high thermal conductivity (about 14 W m<sup>-1</sup> K<sup>-1</sup>) and high mechanical strength.<sup>14–18</sup> However, limited by the proportion mismatch of their emitted light, Ce:YAG TCs can hardly meet the requirement of a high color rendering index (CRI). This restricts the application of this brand-new solid-state lighting technology with TCs as the color converter.

To address this problem, numerous strategies have been proposed to improve the CRI of Ce:YAG TCs. Co-doping red-emitting ions, such as Pr<sup>3+</sup>, Eu<sup>3+</sup>, Sm<sup>3+</sup> or Cr<sup>3+</sup>, into Ce:YAG TC<sup>19–22</sup> has been considered as an effective way to compensate for the loss of the red component. For example, the obtained

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CRI of a reported Ce/Pr-doped YAG TC was only 66.9.<sup>23</sup> A Ce/Pr/Cr tri-doped YAG TC was investigated, and the highest CRI of the encapsulated white LEDs was 78.<sup>20</sup> Although co-doping Pr<sup>3+</sup> ions into Ce:YAG ceramics is a promising strategy to increase their CRI performance, the diverse charge states and the “narrow band” emission of Pr<sup>3+</sup> ion dramatically limit the CRI of the ceramics.<sup>24–27</sup> In addition to Pr<sup>3+</sup> doping, a Ce<sup>3+</sup>/Mn<sup>2+</sup>/Si<sup>4+</sup> tri-doped YAG TC obtained a CRI as high as 82.5.<sup>28</sup> Unfortunately, the luminous efficiency of Ce<sup>3+</sup>/Mn<sup>2+</sup>/Si<sup>4+</sup>:YAG was only 14.0 lm W<sup>-1</sup>. Mn<sup>2+</sup> ion has multiple coordination numbers in the YAG host; however, only 6-coordinated Mn<sup>2+</sup> contributes to the 580 nm emission, which is significant to promote the CRI in white LEDs/LDs. Increasing the Mn<sup>2+</sup> doping concentration would force excess Mn<sup>2+</sup> ions to occupy the 4 coordinated Al<sup>3+</sup> sites (emitted at 740 nm), which exceeds the cutoff wavelength of the CIE 1931 color-matching function  $x$ .<sup>22,29–33</sup> Therefore, the contribution of Mn<sup>2+</sup> doping to the CRI improvement of white LEDs/LDs is still limited.

Designing ceramics with multiple structures has been employed to regulate the spectral properties of TCs to achieve high CRI. A Ce:YAG/Ce,Cr:YAG dual-layered phosphor ceramic revealed that the colorimetric parameters could be controlled by the Cr<sup>3+</sup> ion concentration and ceramic thickness; however, its CRI value was not provided.<sup>34</sup> A Ce:YAG/Ce:(Gd,Y)AG dual-layered TC was constructed using a silicone binder, and a CRI of 62.9 was obtained.<sup>35</sup> However, the designed structure was unreliable in white LED applications due to its poor thermal stability and refractive index mismatch. In addition, a Ce:GYAG/Ce:YAG ceramic was fabricated by controlling the powder ratios, and the obtained CRI was only 70.<sup>36</sup> However, the nonuniform linear shrinkage of the composite ceramics inevitably resulted in the generation of micro cracks and defects. Furthermore, the red phosphor Eu<sup>2+</sup>:CaAlSiN<sub>3</sub> (Eu:CASN) was coated on the Ce:YAG TC surface using the spin-coating technique, which effectively increased the red component of the emitted light, and a CRI of up to 85 was obtained.<sup>17</sup> However, the red phosphor is unstable and easily decomposes or oxidizes. Also, the preparation process of the red phosphor is harsh. These reported CRI improvement methods with respect to the composite structure design strategy were achieved merely by spectral superposition and interface combination, which inevitably result in complex fabrication processes, high cost and difficult encapsulation. Efforts have been made by researchers to enhance the CRI of TC; most of the obtained CRIs were lower than 70, except for the red light-emitting ion doping strategy, as shown in Table S1 in the ESI.<sup>†</sup><sup>15,19,20,23,37–46</sup> Therefore, there is no doubt that adjusting the emitted color proportion in single-structured TCs should be considered as an effective strategy to pursue high CRI performance.

Considering the appropriate red-green-blue light proportion along with the simultaneous emission of wide and narrow peaks in the red region, the strategy of synchronous doping of Pr<sup>3+</sup> and Mn<sup>2+</sup> ions into Ce:YAG TC is expected to achieve a high CRI value. In this study, based on the specific design of synchronous doping of Pr<sup>3+</sup> and Mn<sup>2+</sup> ions, the multi-path energy transfer of both “wide peak” and “narrow peak” red light emission in single structured TCs was proposed. The microstructures,

optical properties, quantum efficiencies, color stability and surface temperatures of the Ce,(Pr,Mn):YAG TCs were investigated systematically. In addition, high-power white LED/LD devices were designed and constructed in the remote encapsulation mode, and the luminous performance of the as-prepared TCs was verified with the purpose of providing approaches to fabricate TCs with high CRI. This excellent color performance indicates that this work will open a new perspective to develop TCs for real applications in the lighting, projection and scintillation fields in the future.

## 2 Experimental procedure

### 2.1 Preparation of Ce,(Pr,Mn):YAG ceramics

In this study, commercial Al<sub>2</sub>O<sub>3</sub> (99.99%, Alfa Aesar, Ward Hill, America), Y<sub>2</sub>O<sub>3</sub> (99.99%, Alfa Aesar, Ward Hill, America), CeO<sub>2</sub> (99.995%, Shandong Xiya, Shandong, China), Pr<sub>6</sub>O<sub>11</sub> (99.995%, Shandong Xiya, Ltd, Shandong, China) and MnCO<sub>3</sub> (99.9%, Shandong Xiya, Shandong, China) were selected as the starting materials, and 0.5 wt% tetraethyl orthosilicate (99.99%, Alfa Aesar, Ward Hill, America) was chosen as the sintering additive. These starting powders were mixed precisely using an analytical balance according to the formula (Ce<sub>0.003</sub>Pr<sub>x</sub>Y<sub>0.997–x</sub>)<sub>3</sub>(Al<sub>1–y</sub>Mn<sub>y</sub>)<sub>5</sub>O<sub>12</sub> ( $x = 0, 0.002, 0, 0.0005, 0.001, 0.0015, 0.002, 0.004$  and  $y = 0, 0, 0.004, 0.005, 0.006, 0.007, 0.008, 0.01$ ), denoted as Pr0Mn0, Pr02Mn0, Pr0Mn04, Pr005Mn05, Pr01Mn06, Pr015Mn07, Pr02Mn08, and Pr04Mn10, respectively, and the detailed formula design is listed in Table 1.

The mixed powders were ball milled with high purity alumina balls for 20 h under 200 rpm in ethanol, and the obtained slurry was then dried at 80 °C for 24 h in an oven. The dried powders were ground and sieved through a 150-mesh screen. After pressing the powders into pellets at 20 MPa using a stainless steel mold ( $\Phi = 22$  mm), the obtained green bodies were then cold isostatic pressed (CIPed) at 220 MPa for 10 min. The CIPed green bodies were calcined at 700 °C for 6 h in air to remove the organic residues and then vacuum sintered at 1750 °C for 8 h under 10<sup>-5</sup> Pa in a tungsten mesh heated vacuum furnace (712T, Thermal Technology LLC, USA). Finally, all the sintered ceramics were polished on both surfaces to 1.0 mm for characterization.

### 2.2 Characterization

The phase compositions of all the samples were characterized using a X-ray diffraction (XRD) machine equipped with a copper

Table 1 Ingredients of the Ce,(Pr,Mn):YAG TCs

| Sample no. | Stoichiometry                                                                                                                                        | Doping (at%) |      |     |
|------------|------------------------------------------------------------------------------------------------------------------------------------------------------|--------------|------|-----|
|            |                                                                                                                                                      | Ce           | Pr   | Mn  |
| Pr0Mn0     | (Ce <sub>0.003</sub> Y <sub>0.997</sub> ) <sub>3</sub> Al <sub>5</sub> O <sub>12</sub>                                                               | 0.3          | —    | —   |
| Pr02Mn0    | (Ce <sub>0.003</sub> Pr <sub>0.002</sub> Y <sub>0.995</sub> ) <sub>3</sub> Al <sub>5</sub> O <sub>12</sub>                                           | 0.3          | 0.2  | —   |
| Pr0Mn04    | (Ce <sub>0.003</sub> Y <sub>0.997</sub> ) <sub>3</sub> (Al <sub>0.996</sub> Mn <sub>0.004</sub> ) <sub>5</sub> O <sub>12</sub>                       | 0.3          | —    | 0.4 |
| Pr005Mn05  | (Ce <sub>0.003</sub> Pr <sub>0.0005</sub> Y <sub>0.9965</sub> ) <sub>3</sub> (Al <sub>0.995</sub> Mn <sub>0.005</sub> ) <sub>5</sub> O <sub>12</sub> | 0.3          | 0.05 | 0.5 |
| Pr01Mn06   | (Ce <sub>0.003</sub> Pr <sub>0.001</sub> Y <sub>0.996</sub> ) <sub>3</sub> (Al <sub>0.994</sub> Mn <sub>0.006</sub> ) <sub>5</sub> O <sub>12</sub>   | 0.3          | 0.1  | 0.6 |
| Pr015Mn07  | (Ce <sub>0.003</sub> Pr <sub>0.0015</sub> Y <sub>0.9955</sub> ) <sub>3</sub> (Al <sub>0.993</sub> Mn <sub>0.007</sub> ) <sub>5</sub> O <sub>12</sub> | 0.3          | 0.15 | 0.7 |
| Pr02Mn08   | (Ce <sub>0.003</sub> Pr <sub>0.002</sub> Y <sub>0.995</sub> ) <sub>3</sub> (Al <sub>0.992</sub> Mn <sub>0.008</sub> ) <sub>5</sub> O <sub>12</sub>   | 0.3          | 0.2  | 0.8 |
| Pr04Mn10   | (Ce <sub>0.003</sub> Pr <sub>0.004</sub> Y <sub>0.993</sub> ) <sub>3</sub> (Al <sub>0.99</sub> Mn <sub>0.01</sub> ) <sub>5</sub> O <sub>12</sub>     | 0.3          | 0.4  | 1.0 |

target X-ray tube (D8 Avance, Bruker, Karlsruhe, Germany) in the scanning range of 10–80°. The crystal structure refinement was performed by the Rietveld method using the General Structure Analysis System (GSAS) software.<sup>47–49</sup> The morphologies of the sintered ceramics were characterized by scanning electron microscopy (SEM; JSM-6510, JEOL, Tokyo, Japan). The transmittance spectra of the polished samples were tested using an UV-Vis-NIR spectrophotometer (Lambda 950, PerkinElmer, USA). Photoluminescence (PL), photoluminescence excitation (PLE) and fluorescence decay spectra were measured using a fluorescence spectrophotometer (FLS920P, Edinburgh, UK) with a scintillating xenon lamp. The chromaticity parameters of the samples were measured using an integrating sphere (R98, Everfine, Hangzhou, China) excited by a 460 nm blue chip as well as a 455 nm laser source. All these measurements were carried out at room temperature.

### 3 Results and discussion

#### 3.1 Phase structures and microstructures

The XRD patterns of the sintered Ce,(Pr,Mn):YAG TCs are shown in Fig. 1. The diffraction peaks of all the Ce,(Pr,Mn):YAG ceramics matched well with the crystal structure of YAG (JCPDS 033-0040), and no impurity phases (e.g. YAP, YAM, CeO<sub>2</sub>, Pr<sub>6</sub>O<sub>11</sub> or MnCO<sub>3</sub>) were detected; this indicates that the Ce<sup>3+</sup>, Mn<sup>2+</sup> and Pr<sup>3+</sup> ions were completely solid-soluted into the YAG lattice. With Pr<sup>3+</sup> and Mn<sup>2+</sup> doping, the diffraction peaks of

the ceramics shifted to lower angles (in the right part of Fig. 1(a), around 33.5°). This demonstrated that Pr<sup>3+</sup> and Mn<sup>2+</sup> doping resulted in a larger unit cell. Generally, due to their similar ion radii, Pr<sup>3+</sup> (0.096 nm, CN = 8) occupies the dodecahedral Y<sup>3+</sup> (0.1019 nm, CN = 8) position, and Mn<sup>2+</sup> (0.066 nm, CN = 6) occupies the octahedral Al<sup>3+</sup> (0.0535 nm, CN = 6) site in the YAG lattice.<sup>29,30,50–52</sup> A schematic of the crystal structure of the Ce,(Pr,Mn):YAG TCs is displayed in Fig. 1(b). Therefore, the concentration variation of the quantity of Mn<sup>2+</sup> ion substituting Al<sup>3+</sup> ion in this study was greater than that of Pr<sup>3+</sup> ion substituting Y<sup>3+</sup> ion, resulting in the lower angle shift of the main diffraction peak (420) in the XRD pattern.

The SEM micrographs of the polished surfaces and fracture surfaces of the prepared ceramics are exhibited in Fig. 2. The average grain size was determined by the linear intercept method. All the samples showed regular grains and clean grain boundaries. Notably, the samples Pr0Mn0 to Pr005Mn05 exhibited completely dense microstructures, leading to their high transmittance. However, with increasing Mn and Pr concentrations up to 0.4 at% and 0.1 at%, respectively, micro-pores could be easily observed (Fig. 2(e)–(g)), resulting in decreased transparency. From the fracture surfaces of the ceramics (Fig. 2(e')–(g')), it could be found that their fracture modes changed from intergranular to transgranular with increasing Mn and Pr doping concentrations. This was remarkably similar to the results obtained by Jiang *et al.*<sup>53</sup> The average grain sizes of Pr0Mn0–Pr04Mn10 were 5.34 μm, 5.98 μm, 6.26 μm, 6.28 μm, 9.91 μm, 6.08 μm, 6.34 μm and 6.12 μm, respectively, indicating that addition of a small amount of active ions did not greatly affect the grain size of the ceramics. The particle size distributions of all the samples can be found in Fig. S1 (ESI†). Moreover, it has been pointed out that the second phases and micro-pores can act as scattering centers in phosphor ceramics,<sup>6,11,45,54–58</sup> resulting in increased extraction efficiency of the incident blue light, which has a favorable effect on the luminous performance of TCs. Actually, in the application

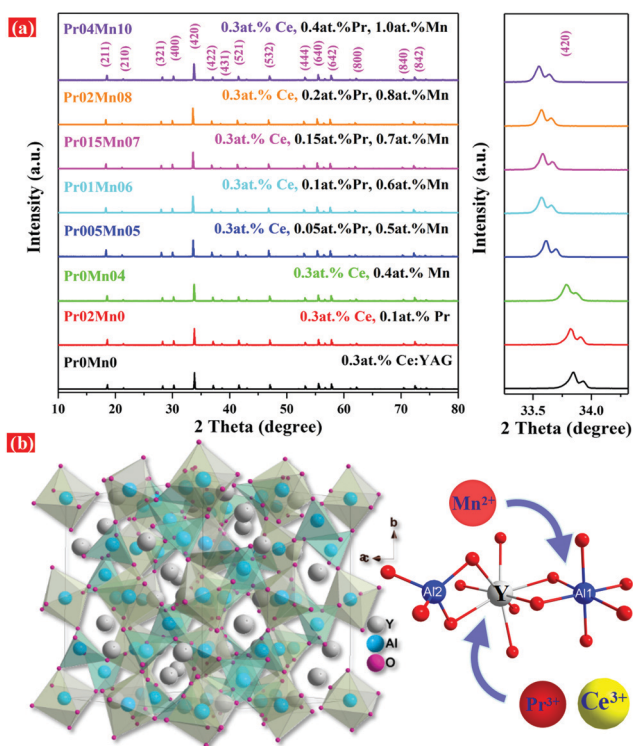


Fig. 1 XRD patterns of (a) TCs with different Pr<sup>3+</sup> and Mn<sup>2+</sup> doping concentrations and magnified XRD patterns around 33°, (b) schematic of the crystal structure of the Ce,(Pr,Mn):YAG TCs.

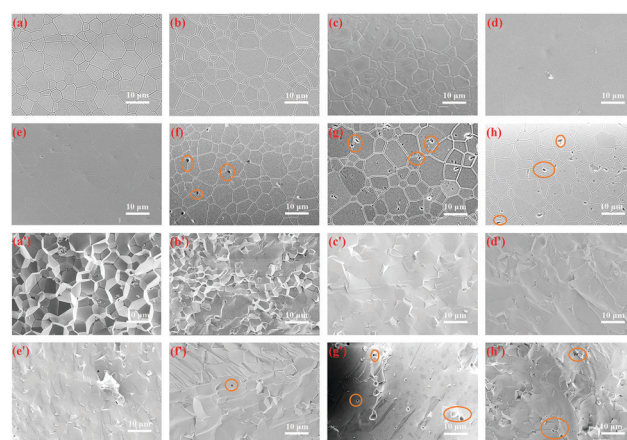


Fig. 2 SEM images of the thermal etched surfaces and the fracture surfaces of (a) and (a') Pr0Mn0, (b) and (b') Pr02Mn0, (c) and (c') Pr0Mn04, (d)–(h) and (d')–(g') Pr005Mn05, Pr01Mn06, Pr015Mn07, Pr02Mn08, and Pr04Mn10.

of white LED lighting, the second phase indeed enhances the utilization of incident blue light and weakens the TIR (total interface refraction). However, further research should be carried out on white LD lighting.

### 3.2 Optical properties

The in-line transmission spectra and appearances of the Ce,(Pr,Mn):YAG TCs are presented in Fig. 3. In the inset, the words behind the TCs can be clearly recognized by the naked eye, indicating good transparency of the ceramics. Obviously, with increasing Pr<sup>3+</sup> and Mn<sup>2+</sup> concentrations, the color of the samples varied from yellow to orange-red. Therefore, it can be concluded that the emission of the TCs was tuned effectively by doping with Pr<sup>3+</sup> and Mn<sup>2+</sup> ions, which is important to enhance the red light emission of white LEDs/LDs. The Ce:YAG TC Pr0Mn0 possessed the best transmittance, which reached 80.7% at 800 nm. Absorption bands around at 330 to 360 nm and 440 to 470 nm could be observed for all the ceramics, originating from the 4f-5d<sup>1</sup> and 4f-5d<sup>2</sup> transitions of Ce<sup>3+</sup> ions. The characteristic absorption bands of Mn<sup>2+</sup> (440 nm/534 nm) and Pr<sup>3+</sup> (450 nm/605 nm) ions were also observed from all the Pr/Mn-doped samples. The absorption bands around 450 nm and 605 nm were attributed to the <sup>3</sup>H<sub>4</sub> → <sup>3</sup>P<sub>0</sub> and <sup>3</sup>H<sub>4</sub> → <sup>1</sup>D<sub>2</sub> transitions of Pr<sup>3+</sup> ion,<sup>59</sup> whereas the 440 nm and 534 nm bands were due to the <sup>6</sup>A<sub>1</sub> → <sup>4</sup>T<sub>2</sub> and <sup>4</sup>T<sub>1</sub> → <sup>4</sup>T<sub>2</sub> transitions of Mn<sup>2+</sup> ions.<sup>22,32,57,60,61</sup> With increasing Mn<sup>2+</sup> and Pr<sup>3+</sup> concentrations, the transmittance at 800 nm of the TCs declined sharply (*T* = 80.7–44.0%). The detailed transmittances of the ceramics located at 800 nm and 400 nm are shown in Fig. S2 (ESI<sup>†</sup>). The decrease in transmittance can be attributed to the sintering promotion effects of the Pr<sub>6</sub>O<sub>11</sub> and MnCO<sub>3</sub> raw powders as sintering additives, which accelerated the grain boundary migration rate of the ceramics during sintering. As a result, a portion of the intergranular pores could be enclosed into the grains to form intragranular pores, which can act as scattering centers to decrease the optical transmittance of the sintered ceramics.

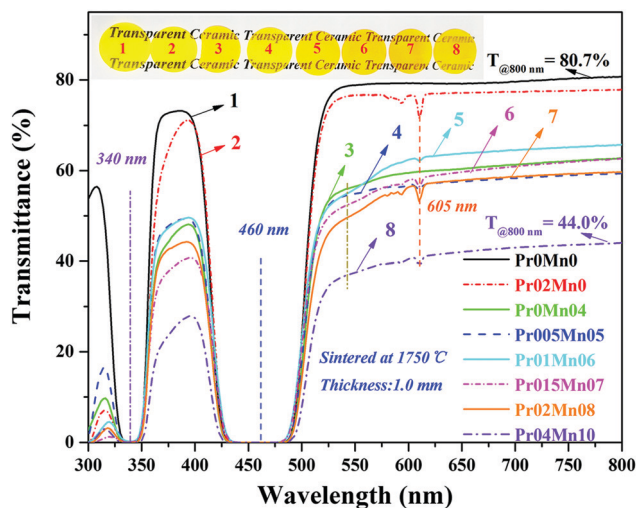


Fig. 3 Optical transmission spectra of the Ce,(Pr,Mn):YAG TCs with 1.0 mm thickness and their appearances (inset).

A similar phenomenon was observed for reported (Y,Gd)AG:Ce ceramics, in which intragranular pores were observed in the heavily Gd-doped samples.<sup>16,18,39,62</sup>

The PL and PLE spectra of the as-prepared ceramics are depicted in Fig. 4. From the PLE spectra ( $\lambda_{em} = 545$  nm), it can be clearly seen that all the samples exhibited two broad excitation bands located at 330–360 nm and 440–470 nm, respectively, originating from the 4f-5d<sup>2</sup> and 4f-5d<sup>1</sup> transitions of Ce<sup>3+</sup> ion (Fig. 4a). The characteristic emission peaks of Ce<sup>3+</sup> (545 nm), Pr<sup>3+</sup> (609 nm) and Mn<sup>2+</sup> (580 nm) ions can be observed, respectively, in Fig. 4b. According to the optimization of the applied sintering additive and sintering method, divalent Mn was obtained in the Ce,(Pr,Mn):YAG TCs. This can be attributed to three main reasons. Firstly, MnCO<sub>3</sub> raw powders (Mn<sup>2+</sup> ions) were applied as the raw material, which prevented the adverse charge state conversion of Mn<sup>2+</sup> ions from the source because the Mn<sup>2+</sup> state is the most stable charge state. Secondly, the sintering process was carried out under a high vacuum atmosphere without a subsequent annealing process; thus, it was impossible to oxidize the Mn<sup>2+</sup> ions into higher charge states. Finally, the applied sintering additive was SiO<sub>2</sub>, and this also prevented the conversion of Mn<sup>2+</sup> ions to adverse charge states. Therefore, the charge state conversion of Mn ion could be directly observed from the spectral properties of all the Ce,(Pr,Mn):YAG TCs. Additionally, the PL spectrum of Pr0Mn04 was distinctly expanded compared to that of Pr0Mn0, which was caused by the emission of Mn<sup>2+</sup> ions (580 nm). Meanwhile, the characteristic emission peaks of Ce<sup>3+</sup>, Pr<sup>3+</sup> and Mn<sup>2+</sup> ions could be observed for all the Ce,(Pr,Mn):YAG TCs (Fig. 4c). Notably, inhomogeneous spectral broadening could be achieved through the synergistic effect of synchronous doping of Pr<sup>3+</sup> and Mn<sup>2+</sup> ions, and the corresponding FWHM values increased from 91.7 to 102.2 nm with increasing Pr<sup>3+</sup> and Mn<sup>2+</sup> concentrations (Fig. 4d). The structural parameters are listed in Table S2 and Fig. S3 (ESI<sup>†</sup>). The polyhedron distortion ratio (*D*) value increased with increasing addition of Pr<sup>3+</sup> and Mn<sup>2+</sup> ions.

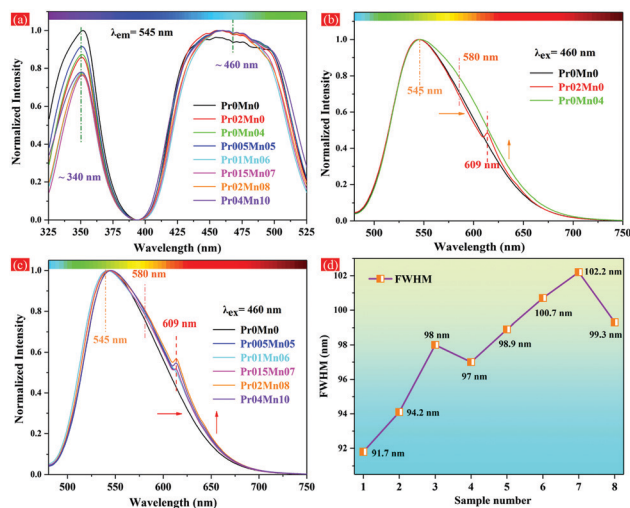


Fig. 4 Normalized (a) PLE and (b and c) PL spectra and (d) the FWHM evolution of the PL spectra of the prepared ceramics.



Therefore, incorporating  $\text{Pr}^{3+}$  and  $\text{Mn}^{2+}$  can reduce the local symmetry of the dodecahedron  $\text{CeO}_8$  in the  $\text{Ce},(\text{Pr},\text{Mn})\text{:YAG}$  TC. Generally, the FWHM is dependent on the local coordination environment of  $\text{Ce}^{3+}$ , which is affected by the host composition.<sup>63–66</sup> The increase of the FWHM can be ascribed to the increased polyhedron distortion ratio and the reduced structural symmetry in the  $\text{PrO2Mn08}$  TC. A PL band with a wide FWHM is advantageous in promoting the color-conversion in  $\text{Ce:YAG}$  TC, which is beneficial to achieve warm light emission with a high CRI.

### 3.3 Luminous performance of white LEDs/LDs

To further evaluate the luminous properties of the ceramics and validate their availability as white light converters, white LED devices were assembled by combining the ceramics with a 20 W blue LED chip (460 nm) using the remote encapsulation mode. Also, a white LD using a 455 nm laser emitting source was constructed by the transmission mode.<sup>67</sup> The lighting photographs and infrared thermal images of the TC-based white LEDs/LDs are shown in Fig. 5. With increasing  $\text{Pr}^{3+}$  and  $\text{Mn}^{2+}$  concentrations, the emitting color of the assembled white LEDs distinctly varied from yellow-white to pure white, and the highest color rendering property was reached (CRI = 84.8) when the  $\text{Pr}^{3+}$  and  $\text{Mn}^{2+}$  concentrations were 0.2 at% and 0.8 at%, respectively. However, the emitting color of the  $\text{PrO4Mn10}$  sample transformed into yellowish-white owing to the mismatched  $\text{Pr}^{3+}/\text{Mn}^{2+}$  concentration (Fig. 5(h)). Therefore, the hue of the TCs could be regulated by the synchronous doping effect.

Additionally, the temperature distributions of all the ceramics were uniform, indicating their good thermal conductivity

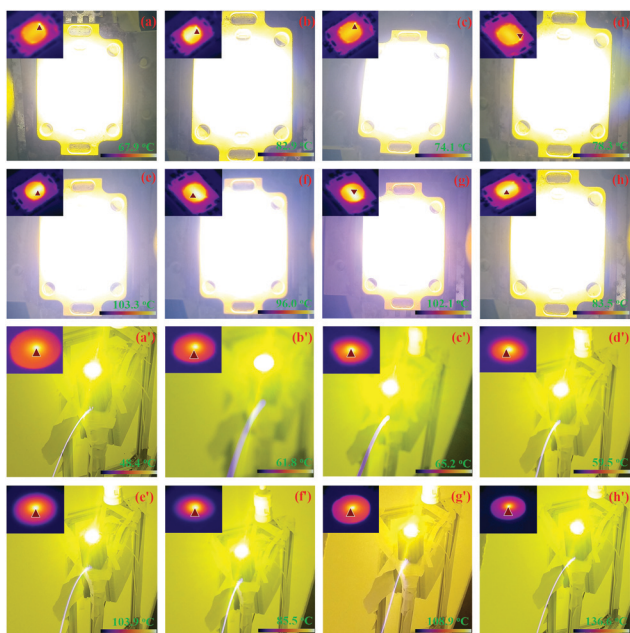


Fig. 5 Photographs of the lighting of white LEDs/LDs devices and corresponding infrared thermal images of (a/a')  $\text{PrO Mn0}$ , (b/b')  $\text{PrO}_2\text{Mn0}$ , (c/c')  $\text{PrO Mn04}$ , (d/d')  $\text{PrO05Mn05}$ , (e/e')  $\text{PrO1Mn06}$ , (f/f')  $\text{PrO15Mn07}$ , (g/g')  $\text{PrO2Mn08}$ , and (h/h')  $\text{PrO4Mn10}$  excited by blue LEDs/LDs.

(see the inset of Fig. 5). The operating temperature increased from 67.9 °C ( $\text{PrO Mn0}$ ) to 103.3 °C ( $\text{PrO1Mn06}$ ) and then decreased to 83.5 °C ( $\text{PrO4Mn10}$ ). It was found that the  $\text{PrO Mn0}$  sample exhibited the lowest temperature; this illustrates that energy loss, which will be transferred into residual heat, is inevitable during the energy transfer process among  $\text{Ce}^{3+}\text{-Pr}^{3+}\text{-Mn}^{2+}$  ions.

Warm yellow light emission could be observed from all the LD-excited TC devices, and no obvious change was observed with respect to the hue (Fig. 5(a'–h')). Notably, the  $\text{PrO2Mn08}$  sample exhibited red-yellow emission (CRI = 70, CCT = 3550 K), indicating that its red emission component was higher than those of the other samples. As a result, the TC-based white LD in this study could be used as a warm light source with high CRI.

Interestingly, for the temperature distribution of the white LDs, the measured temperatures matched well with the optical transmittance of the samples (Fig. 3), confirming that the residual pores reduced the thermal conductivity of the ceramics. Compared with the LED light source, the LD source is a point light source with a high energy density ( $12 \text{ W mm}^{-2}$ ). Accordingly, the effect of the residual pores on the heat dissipation of ceramics under LD excitation is more obvious than that under LED excitation. It has been reported that heat accumulation decreases the luminance performance of white LDs.

The electroluminescence spectra and chromaticity parameters are shown in Fig. 6. The CIE color coordinates ranged from (0.3837, 0.4358) to (0.3327, 0.2951), corresponding to the yellow and white areas, respectively. Amazingly, the CRI of the  $\text{PrO2Mn08}$  sample was as high as 84.8. To the best of our knowledge, this is the highest recorded CRI value obtained from a single-structured TC material. Meanwhile, the CCT of the white LED devices varied from 4272 to 5456 K, which are relatively ideal values on the Planck trajectory. As a result, the obtained ceramics in this study exhibit desirable luminance performance, which addresses the widespread challenge of CRI improvement in single-structured TCs.

Fig. 7(a and b) shows the electroluminescence spectra of the white LEDs and LDs. The characteristic yellow-red bands

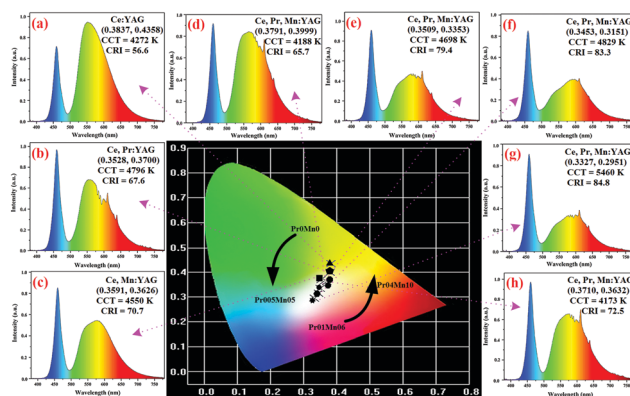


Fig. 6 Electroluminescence spectra and chromaticity parameters of the corresponding white LEDs driven by a 350 mA current: (a)  $\text{PrO Mn0}$ , (b)  $\text{PrO}_2\text{Mn0}$ , (c)  $\text{PrO Mn04}$ , (d–h)  $\text{PrO05Mn05}$ ,  $\text{PrO1Mn06}$ ,  $\text{PrO15Mn07}$ ,  $\text{PrO2Mn08}$ ,  $\text{PrO4Mn10}$ .

between 500 nm and 700 nm correspond to the emission bands of  $\text{Ce}^{3+}$ ,  $\text{Mn}^{2+}$  and  $\text{Pr}^{3+}$  ions. The decreased blue light intensity for the  $\text{Ce},(\text{Pr},\text{Mn})\text{:YAG}$  TCs was attributed to their increased scattering loss compared with that of  $\text{Ce:YAG}$ . Obviously, the emission intensity of the samples decreased markedly with increasing concentration of  $\text{Pr}^{3+}$  and  $\text{Mn}^{2+}$  ions, resulting in decreased intensity of the blue-green component and the increased yellow-red proportion.

A schematic of the designed white LD device is shown in the insert of Fig. 7(b). Our previous research indicated that the transmission mode of white LDs has better color stability than their reflection mode.<sup>67</sup> A metal heat sink with the embedded ceramic was employed to improve the heat dissipation of the TC-based LD device, and all the measurements were carried out using the same structured heat sink. The measured colorimetric parameters are listed in Table S3 (ESI<sup>†</sup>).

The luminous efficiency of radiation (LER) was calculated by formula S1 (ESI<sup>†</sup>). As shown in Fig. 7(c), the LER of the ceramics gradually declined from  $\text{Pr0Mn0}$  to  $\text{Pr04Mn10}$ . There are two reasons for this phenomenon. One is that with the increasing  $\text{Pr}^{3+}$  and  $\text{Mn}^{2+}$  concentrations, the energy transfer-induced energy loss increases synchronously. The other factor is the decreased transmission of TCs with increasing  $\text{Pr}^{3+}$  and  $\text{Mn}^{2+}$  concentration.

Compared with  $\text{Ce:YAG}$  TC, the energy loss in the energy transfer process of the  $\text{Ce},(\text{Pr},\text{Mn})\text{:YAG}$  TCs will inevitably decrease the LER. In other words, a trade-off exists between LER and CRI. Yang *et al.* mixed green  $\text{Eu}^{2+}\text{:}(\text{Ba},\text{Sr})_2\text{SiO}_4$  and red  $\text{Eu}^{2+}\text{:CaAlSi}_3\text{N}_3$  phosphors as a color convertor, and the achieved CRI was as high as 89, with only  $56.5 \text{ lm W}^{-1}$  LER obtained.<sup>68</sup>

In addition to strict control of the light proportion, the heat dissipation capability is a key factor that affects the luminance property of TCs, especially in high power white LED/LD lighting applications. Xie *et al.* demonstrated that the better the heat

dissipation, the smaller the thermal accumulation, which is beneficial to mitigate the “thermal run-away” phenomenon of TCs to improve the resulting CRI.<sup>12</sup> Considering that the metal heat sink could effectively improve the heat dissipation, the plate-like metal heat sink was designed to solve the heat accumulation problem under LD excitation with high energy density. As can be seen from Fig. 7(d), the heat accumulation was controlled effectively by the plate-like metal heat sink, and a similar operation temperature was obtained compared with that of the LED excitation source.

To further explore the energy transfer of the  $\text{Ce},(\text{Pr},\text{Mn})\text{:YAG}$  TCs, the fluorescence decay behaviors of  $\text{Ce}^{3+}$  at 545 nm were investigated under 460 nm excitation, as shown in Fig. 8(a) and (b). Single exponential decay was observed in all samples. With increasing  $\text{Mn}^{2+}$  and  $\text{Pr}^{3+}$  concentrations, the lifetimes of the samples displayed a decreasing trend from 69.47 ns ( $\text{Pr0Mn0}$ ) to 60.66 ns ( $\text{Pr0Mn04}$ ), revealing energy transfer between the  $\text{Ce}^{3+}\text{-Pr}^{3+}$  ions and  $\text{Ce}^{3+}\text{-Mn}^{2+}$  ions. Meanwhile, from Fig. 8(b), the lifetimes of the doped TCs decreased with increasing  $\text{Ce}^{3+}$ ,  $\text{Pr}^{3+}$  and  $\text{Mn}^{2+}$  concentrations, manifesting the increased energy transfer efficiency. Moreover, the increased lifetime of the  $\text{Pr04Mn10}$  sample was attributed to the concentration quenching effect.

A schematic of the energy transfer process of the  $\text{Ce}^{3+}\text{-Pr}^{3+}\text{-Mn}^{2+}$  system in the YAG host is provided in Fig. 8(c). The 5d state of the  $\text{Ce}^{3+}$  ions splits into the 5d<sup>1</sup> and 5d<sup>2</sup> states, respectively. Electrons are excited to the upper 5d<sup>1</sup> state of the  $\text{Ce}^{3+}$  ions by the 460 nm excitation; next, they are transferred to the <sup>4</sup>T<sub>1</sub> state of  $\text{Mn}^{2+}$  ions after a rapid relaxation process and are then transferred to the <sup>6</sup>A<sub>1</sub> ground state of the  $\text{Mn}^{2+}$  ions to produce a 580 nm light emission. Simultaneously, the energy transfer process between  $\text{Ce}^{3+}$  and  $\text{Pr}^{3+}$  ions (<sup>5</sup>d<sub>1</sub> → <sup>1</sup>G<sub>4</sub>) also takes place, resulting in the 609 nm red emission. Fig. 8(d) presents the IQEs of the as-prepared ceramics. With increasing Pr and Mn concentrations, the IQE of the corresponding TCs declined gradually from 86.73%

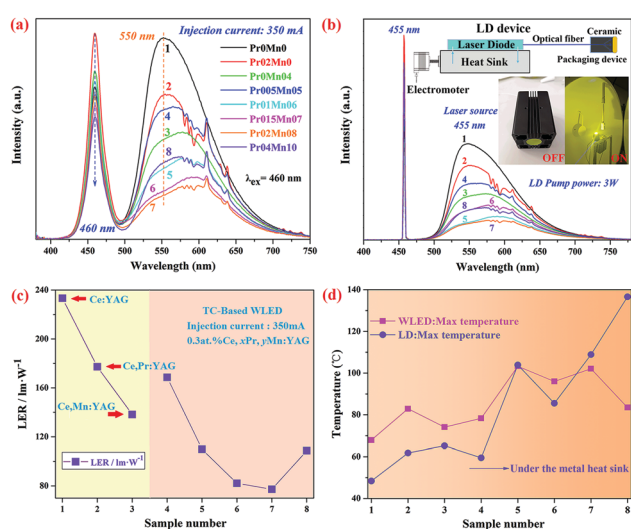


Fig. 7 The electroluminescence spectra of (a) the white LEDs and (b) the LDs and schematic of the white LD device (inset); (c) LER variation of the white LEDs, (d) temperature contrast between the white LEDs and LDs under stable operation.

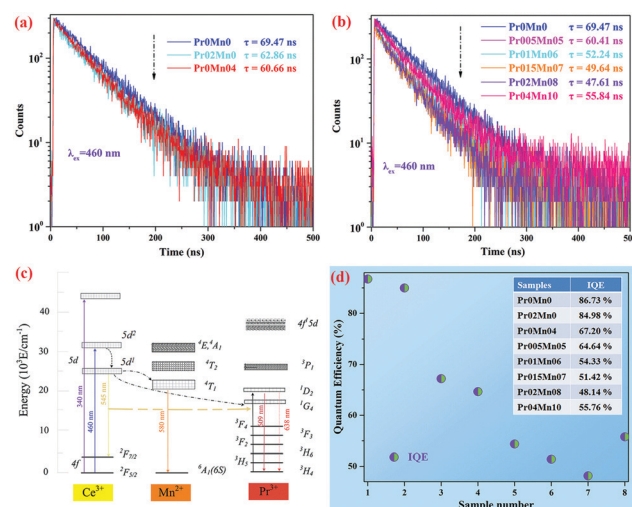


Fig. 8 Fluorescence decay curves: (a)  $\text{Pr0Mn0}$ ,  $\text{Pr0Mn04}$ . (b)  $\text{Pr0Mn0}$ ,  $\text{Pr005Mn05}$ ,  $\text{Pr01Mn06}$ ,  $\text{Pr015Mn07}$ ,  $\text{Pr02Mn08}$ ,  $\text{Pr04Mn10}$ . (c) Schematic of the energy transfer process among  $\text{Ce}^{3+}\text{-Pr}^{3+}\text{-Mn}^{2+}$  ions. (d) Internal quantum efficiencies (IQEs) of the TCs under 460 nm excitation.

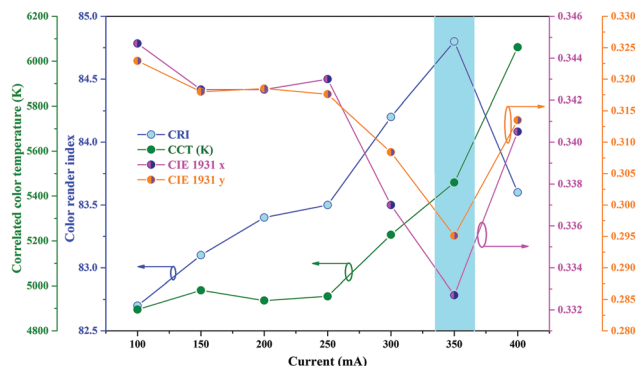


Fig. 9 The colorimetric parameters of Pr02Mn08 under blue LED excitation as a function of the driving current.

to 48.14%, and then increased to 55.76%. The decreased quantum efficiency can be attributed to the energy transfer among the  $\text{Ce}^{3+}$ - $\text{Pr}^{3+}$ - $\text{Mn}^{2+}$  ions in the  $\text{Ce}(\text{Pr}, \text{Mn})\text{:YAG}$  TCs.

The chromaticity parameters of the Pr02Mn08-based LED as a function of the working current are represented in Fig. 9. Obviously, with increasing injection current from 100 mA to 400 mA, the CRI increased from 82.7 to 84.8 and then decreased to 83.5. Additionally, the CCT increased from 4850 K to 6100 K with increasing current due to the increased blue light proportion. However, the variations of  $D_x$  and  $D_y$  were merely 0.031 and 0.046, respectively, which are almost negligible. Notably, the fluctuations of both CRI and CCT were only 5%, demonstrating the excellent color stability of the TCs. The detailed CCT, CRI and CIE color coordinate parameters are listed in Tables S4–S7 (ESI<sup>†</sup>). In general, the enhanced heat accumulation and blue light emission at a high driving current result in lower photon-to-photon transfer efficiency.<sup>7,60,69–73</sup> As a result, the chromaticity parameters changed due to the color component mismatch, which affected the luminous performance of the white LEDs.

## 4 Conclusions

In this paper, by synchronous doping with  $\text{Pr}^{3+}$  and  $\text{Mn}^{2+}$  ions, a recorded CRI of 84.8 and natural white light emission with a CCT of 5456 K were obtained in single-structured  $\text{Ce}(\text{Pr}, \text{Mn})\text{:YAG}$  TCs, and an inhomogeneous broadening of the FWHM from 91.7 nm to 102.2 nm was realized. In addition, the fluctuations of the obtained CRI and CCT under different driving currents of the blue LED were only 5%, demonstrating its excellent color stability. Also, a 3550 K CCT of the warm white light emission under self-constructed LD excitation was realized, and the obtained CRI was 70. Therefore, this study paves a way for the commercialization of TC-based white LEDs/LDs in the fields of high CRI lighting and displays.

## Conflicts of interest

The authors declare that we do not have any commercial or associative interest that represents a conflict of interest in connection with the work submitted.

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