



Cite this: *Phys. Chem. Chem. Phys.*,  
2023, 25, 4667

# First principles insights into the relative stability, electronic and catalytic properties of core–shell, Janus and mixed structural patterns for bimetallic Pd–X nano-alloys (X = Co, Ni, Cu, Rh, Ag, Ir, Pt, Au)<sup>†‡</sup>

Soumendu Datta, \* Aishwaryo Ghosh and Tanusri Saha-Dasgupta

The three well-known orderings of the two constituting atomic species in a bimetallic nano-alloy – core–shell, Janus and mixed structural patterns – may be interconvertible depending on the synthesis conditions. Using first principles electronic structure calculations in the present work, we look for the microscopic origin for such structural transformation considering eight Pd-related bimetallic nano-alloys. Our analysis shows that it is the change in atom–atom covalency that is responsible for such structural transformation. Our study also reveals that the three patterns are distinctly identified in terms of total orbital hybridization. Finally, we have analyzed the trend in the relative catalytic activity for the three structures of each bimetallic nano-alloy using the d-band model. Our analysis indicates that the trend in the catalytic activity for the bimetallic Pd–X nano-alloys seems to be intermediate to those of the pristine Pd and Pt nano-clusters possessing similar structure and equal number of total atoms. Among the studied binary nano-alloys, the bimetallic Pd–Ni nano-alloy appears as the most suitable binary pair to develop a non-Pt catalyst.

Received 17th September 2022,  
Accepted 20th January 2023

DOI: 10.1039/d2cp04342d

rsc.li/pccp

## 1 Introduction

Compared to elemental nano-particles, multi-elemental nano-alloys are promising functional materials due to their superior practical uses in diverse areas, such as catalysis,<sup>1</sup> optics,<sup>2</sup> sensing,<sup>3</sup> advanced bio-medicine,<sup>4</sup> magnetic storage,<sup>5</sup> which are basically driven by the synergism of their size, composition and chemical ordering-dependent properties.<sup>6,7</sup> Depending on the ordering of the two species of atoms, the overall structure of a bimetallic nano-particle can be broadly classified into three main structural patterns – core–shell, Janus and mixed patterns.<sup>6</sup> In a binary core–shell nano-alloy, a core of one species of atoms is covered by a shell of another elemental atom. In a Janus structure, two sub-clusters of the atoms from either elemental species are sandwiched side-by-side. Likewise, the atoms of two different species are intermixed throughout the structure of a mixed bimetallic nano-alloy. For a bimetallic nano-alloy, the preference to form either of the three structural

patterns out of the atoms from the two constituting metallic species is primarily driven by several factors characterizing the two elemental species – notably atomic size-mismatch between the two species of atoms, difference of their bulk cohesive energies/surface energies, coordination difference in their bulk crystal structure, bulk miscibility and magnetism.<sup>6,8,9</sup> For a core–shell bimetallic nano-alloy, more control over the ordering of the two species of atoms in the core and shell regions could also be achieved experimentally under kinetically and thermodynamically controlled reaction conditions.<sup>10</sup>

On the other hand, it is interesting to note that the three structural patterns for a bimetallic nano-alloy of a fixed composition may also be interconvertible depending on some factors, especially between the core–shell and Janus structural patterns. Several studies have been carried out in the past to understand such structural inter-conversion. The main driving factors reported so far include – I. competition between bulk-like chemical interactions, surface segregation, and strain-driven segregation;<sup>11</sup> II. the interplay of core size, shell thickness, asymmetric core position, and large difference of surface energies between the two constituting elements;<sup>12,13</sup> III. size and chemical composition of the nano-alloy system;<sup>13,14</sup> IV. synthesis conditions such as temperature and presence of a support;<sup>6</sup> V. competition between the energetic and

Satyendra Nath Bose National Centre for Basic Sciences, JD Block, Sector-III, Salt Lake City, Kolkata 700 106, India. E-mail: soumendu.datta@gmail.com

<sup>†</sup> PACS number: 36.40.Cg, 73.22.–f, 71.15.Mb, 31.10.+z.

<sup>‡</sup> Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d2cp04342d>

configurational entropic contributions.<sup>15</sup> This prompted us to identify some relevant microscopic properties whose change is in synchronization with the simultaneous change in the relative stability among the three structures for a bimetallic nano-alloy of fixed composition.

Pd is a Pt-group metal and Pd-based bimetallic nano-alloys are promising due to increasing demand for developing stable non-Pt electro-catalysts that have comparable catalytic performance due to the low abundance and high cost of Pt. In this regard, Pd-based alloys are particularly preferred due to their similar electronic structure, comparable performance, better methanol tolerance, more resistance towards CO-poisoning and above all the lower cost of Pd than Pt.<sup>16</sup> Alloying of Pd with 3d, 4d and 5d transition metal elements has received particular research attention to enhance the catalytic activity pertaining to different kinds of fuel cells.<sup>17–19</sup> In the present work, we focus on eight Pd-based bimetallic nano-alloys, namely three bimetallic Pd-3d nano-alloys – Pd–Co, Pd–Ni and Pd–Cu; two bimetallic Pd-4d nano-alloys – Pd–Rh and Pd–Ag; and three bimetallic Pd-5d nano-alloys – Pd–Ir, Pd–Pt and Pd–Au. It should be noted that the eight alloying elements encompass magnetic elements – Co and Ni; catalytic elements – Rh, Ir and Pt; and plasmonic elements – Cu, Ag and Au.

To prepare the three structural configurations containing a fixed number of total atoms and the same composition, we started from a large 147-atom sized icosahedral structure of a pristine Pd<sub>147</sub> cluster and constructed the three initial structural patterns considering two compositions – Pd<sub>55</sub>X<sub>92</sub> (*i.e.* Pd-poor) and Pd<sub>92</sub>X<sub>55</sub> (*i.e.* Pd-rich) for each bimetallic Pd–X nano-alloy, where “X” indicates each of the remaining eight elements. Though the physical properties of a nano-system are size and shape dependent,<sup>20</sup> our choice of an atomic shell structured icosahedron as the overall host structure is motivated by its high stability for the noble metal elemental nanoparticles.<sup>21</sup> Moreover, it is rather easy to construct the three structural patterns with the same composition, out of an icosahedron. Then, we optimized the three structures for each bimetallic nano-alloy. Considering the optimized structures, we have studied the relative stability among the three structural patterns for each bimetallic nano-alloy of either composition. Subsequently, we have analyzed the structural and electronic properties of the three patterns using first principles density functional theory (DFT) calculations and their variations along the bimetallic Pd–X nano-alloy series. Our analysis of Crystal Orbital Hamiltonian Populations<sup>22</sup> indicates that the relative stability is mainly governed by the change in atom–atom covalency from one structural pattern to another for each bimetallic Pd–X nano-alloy of fixed composition. We note that the three patterns are also distinctly identified in terms of an orbital hybridization index. Finally, we have analyzed the trend in catalytic activity using the d-band model<sup>23</sup> for the three structures of each bimetallic nano-alloy. Our analysis of the occupied d-band center's position indicates that the catalytic activity for most of the studied bimetallic Pd–X nano-alloys would lie in-between that of a pristine Pt<sub>147</sub> and Pd<sub>147</sub> cluster in either composition limit, except for the Pd-poor

Pd–Ag nano-alloy. The overall analysis of both the relative stability and catalytic properties implies that the bimetallic Pd–Ni nano-alloy would be the most suitable candidate to design a non-Pt catalyst – particularly its Janus and mixed structures for the Pd-poor composition and the core–shell structure for the Pd-rich composition (with the Pd atoms at the surface). To cross-check the size effects within the motif of an atomic shell structured icosahedron, we have carried out additional calculations for a 55 atom sized bimetallic PdNi nano-alloy considering its two compositions, *i.e.* Pd<sub>13</sub>Ni<sub>42</sub> (Pd-poor) and Pd<sub>42</sub>Ni<sub>13</sub> (Pd-rich) nano-alloys, which confirm the trend observed for the 147 atom-sized binary Pd–X nano-alloys. Finally, we have confirmed the predicted trend of catalytic activity for the three structures of the bimetallic Pd–Ni nano-alloy relative to that of the pristine Pd<sub>147</sub> cluster, by rigorous analysis of the absorption reaction with a CO molecule on their surface.

## 2 Construction of three structural patterns and computational details

It should be noted that the icosahedron-like concentric cell structure has been predicted previously as the ground state geometry for larger sized pristine Pd nano-clusters.<sup>24</sup> An icosahedron of 147 atoms forms a third closed shell icosahedron with 55-atoms in the core and 92-atoms at the outer shell. Therefore, an initial core–shell structural pattern for a bimetallic X<sub>55</sub>Pd<sub>92</sub> nano-alloy can be easily constructed from the pristine Pd<sub>147</sub> cluster by replacing its 55 core atoms with the X elemental atoms. Likewise, an initial structure of Janus pattern for a bimetallic X<sub>55</sub>Pd<sub>92</sub> nano-alloy is prepared by replacing a hemispherical fraction of Pd<sub>55</sub> from one side, by 55 atoms of the X element, so that the X<sub>55</sub> and Pd<sub>92</sub> fractions are sandwiched side-by-side. Finally, the initial structure of a mixed pattern for a bimetallic X<sub>55</sub>Pd<sub>92</sub> nano-alloy is formed by random mixing of the 55 X atoms with 92 atoms of Pd throughout the whole structure. Fig. 1 shows the initial structural pattern for a core–shell, Janus and mixed bimetallic nano-alloy with X<sub>55</sub>Pd<sub>92</sub> composition. On the other hand, the initial structural patterns for the core–shell, Janus and mixed structures of a bimetallic Pd<sub>55</sub>X<sub>92</sub> nano-alloy are obtained from those of the X<sub>55</sub>Pd<sub>92</sub> nano-alloy by interchanging between the X atoms and Pd atoms. We have additionally considered 55 atom-sized bimetallic Ni–Pd clusters to cross-check the cluster-size effect.



Fig. 1 Initial structural patterns of core–shell, Janus and mixed configurations of a bimetallic Pd<sub>92</sub>X<sub>55</sub> nano-cluster having an overall icosahedral structure containing 147 atoms. The green balls correspond to Pd-atoms and red balls represent atoms of an X element.

The calculations reported in this study were based on DFT within the framework of the pseudo-potential plane wave method, as implemented in the Vienna ab initio Simulation Package (VASP).<sup>25</sup> We used the Projected Augmented Wave (PAW) pseudo-potential<sup>26,27</sup> coupled with the generalized gradient approximation (GGA) to the exchange correlation energy functional as formulated by Perdew, Burke and Ernzerhof (PBE).<sup>28</sup> The 3d as well as 4s electrons for atoms of Co, Ni and Cu; 5s as well as 4d electrons for atoms of Rh, Pd and Ag; and the 6s and 5d electrons for the atoms of Ir, Pt and Au were treated as the valence electrons. The wave functions were expanded in the plane wave basis set with the kinetic energy cut-off of 300 eV. For the cluster calculations, a simple cubic super-cell with each side of 28 Å was used with the periodic boundary conditions, where two neighboring clusters were kept separated by around 12 Å vacuum space. The large cell-size essentially makes the interaction between cluster images negligible. Reciprocal space integrations were carried out at the  $\Gamma$  point. All structures were optimized using the conjugate gradient and the quasi-Newtonian methods until all the force components were less than a threshold value of 0.001 eV Å<sup>-1</sup>. The convergence of the energies with respect to the cut-off as well as the  $k$ -mesh values were checked. Due to the presence of heavy 4d and 5d atoms, non-collinear DFT calculations including spin-orbit interactions were performed for all calculations.

### 3 Results and discussion

The major relevant properties of the constituting elements, which influence the formation of binary nano-alloys, are atomic radius, bulk cohesive energy and surface energy, and electro-negativity. Table 1 contains the values of these quantities for the constituent species. It is seen from Table 1 that the cohesive energy and surface energy of Cu, Ag and Au are lower with respect to the corresponding values of Pd, while the other elements have higher cohesive energy and surface energy. Furthermore, Ir (Cu) possesses the biggest (smallest) atomic radius and Au (Co) is associated with the most (least) electro-negativity. Ir (Ag) also possesses the highest (smallest) cohesive energy and surface energy. In the following, we have analyzed the optimized structures and the trend in relative stability,

**Table 1** Atomic radius ( $r$ ), bulk cohesive energy ( $E_c$ ), surface energy ( $S$ ) and electro-negativity ( $\chi$ ) of the constituent elements<sup>a</sup>

Element	$r$ (Å)	$E_c$ (eV per atom)	$S$ (J m <sup>-2</sup> )	$\chi$ (Pauli scale)
Co	1.52	4.39	2.55	1.88
Ni	1.49	4.44	2.38	1.91
Cu	1.45	3.49	1.79	1.9
Rh	1.73	5.75	2.70	2.28
Pd	1.69	3.89	2.05	2.20
Ag	1.65	2.95	1.25	1.93
Ir	1.80	6.94	3.00	2.20
Pt	1.77	5.84	2.475	2.28
Au	1.74	3.81	1.50	2.54

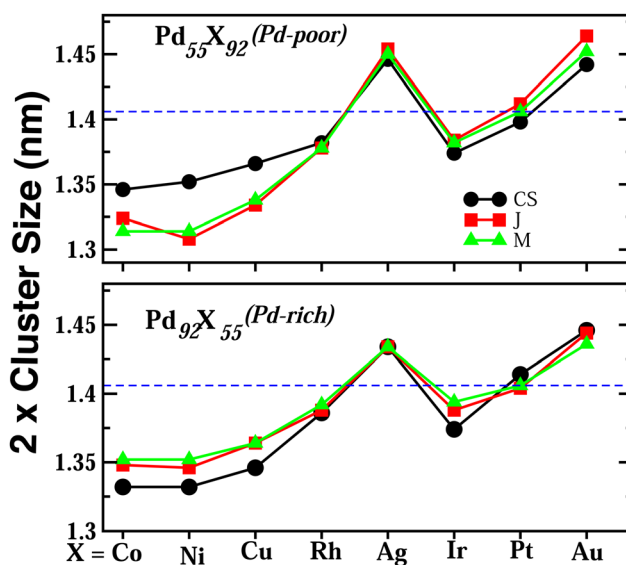
<sup>a</sup> CRC Handbook of Chemistry and Physics, ed. D. R. Lide, CRC Press, Boca Raton, Florida, 84th edn, 2003.

electronic and catalytic properties among the three structural patterns for each bimetallic nano-alloy.

#### 3.1 Optimized structures

Upon structural optimization for the three structures of every bimetallic Pd-X nano-alloy, the overall icosahedral structural pattern remains intact for each structural pattern. However, the effective cluster sizes exhibit some variation along the Pd-X series. We define cluster size as the average distance between the cluster center and the outermost surface atoms. Details of our estimated cluster size for the Pd-X clusters with their respective optimized structure are given in the ESI.† Fig. 2 shows the plot of our estimated cluster size of the three morphologies for each binary nano-cluster in both the composition limits. First, starting from an unoptimized cluster size of 0.703 nm (the same as the cluster size for the optimized pristine Pd<sub>147</sub> cluster), it is seen that the optimized cluster sizes for the binary Pd-X clusters vary within a range of 0.66–0.73 nm. The overall trend of variation remains almost similar for both the compositions. As expected, the size for the Pd-Ni or Pd-Co cluster with their Pd-poor composition remains smaller due to their smaller atomic size. Then, the cluster size shows an increasing trend with the increasing atomic number of the X element. It is to be noted that the cluster size peaks for the binary Pd-Ag clusters. This could result from the smaller atom-atom interaction due to the presence of a filled 4d valence orbital of the Ag atom. Interestingly, we note that the variation of cluster size is qualitatively understandable from the variation of orbital hybridization as discussed in Section 3.3.

In a bimetallic Pd-X nano-alloy, there are three types of nearest neighbor (NN) bonds-Pd-Pd, Pd-X and X-X. The



**Fig. 2** Plot of the estimated cluster size of the optimized structures for the three morphologies of each Pd-X nano-alloy in both compositions. The black colored circles, red colored squares and green colored triangles represent the data for the core-shell (CS), Janus (J) and Mixed (M) morphologies, respectively, for each Pd-X nano-alloy. The horizontal dashed line represents our estimated size of the pristine Pd<sub>147</sub> cluster.

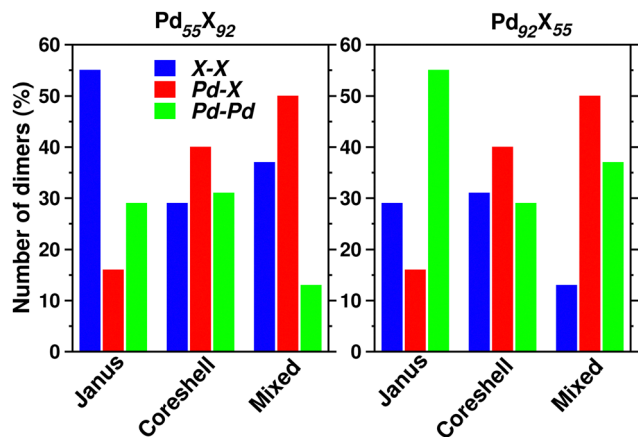


Fig. 3 Bar plot for the overall percentage of the number of X–X, Pd–X and Pd–Pd nearest neighbor dimers in the optimized Janus, core–shell and Mixed structures for each bimetallic  $\text{Pd}_{55}\text{X}_{92}$  and  $\text{Pd}_{92}\text{X}_{55}$  cluster.

number of a similar type of NN bonds for each structural pattern along the Pd–X nanoalloys within a fixed composition also remains approximately the same with varying X element. The plot of the overall percentage for the total number of X–X, Pd–X and Pd–Pd NN bonds with respect to the overall total NN bonds associated with each of the three structures constructed out of a 147 atom icosahedron is shown in Fig. 3 for either composition. Interestingly, we note that the three structural patterns are characterized by the dominance or deficiency of a specific type of NN bond. For the  $\text{Pd}_{55}\text{X}_{92}$ , *i.e.* Pd-poor nanoalloys, it is seen that:

- the optimized core–shell structure involves the most number of Pd–X NN bonds and the least number of X–X NN bonds,
- the optimized Janus structure involves the most number of X–X NN bonds and the least number of Pd–X NN bonds,
- the optimized mixed structure involves the most number of Pd–X NN bonds and the least number of Pd–Pd NN bonds

On the other hand, the sequence of NN bond types for the  $\text{Pd}_{92}\text{X}_{55}$ , *i.e.* Pd-rich nanoalloys, is

- the optimized core–shell structure involves the most number of Pd–X NN bonds and the least number of Pd–Pd NN bonds,
- the optimized Janus structure involves the most number of Pd–Pd NN bonds and the least number of Pd–X NN bonds,
- the optimized mixed structure involves the most number of Pd–X NN bonds and the least number of X–X NN bonds

For a bimetallic  $\text{Pd}_{55}\text{X}_{92}$  or  $\text{Pd}_{92}\text{X}_{55}$  nano-alloy, the NN bond-length for each of the three types of NN bonds is also different in the optimized structures of the three patterns. Fig. 4 shows the variation of only  $\langle\text{Pd–Pd}\rangle$  NN bond-length (“ $\langle-\rangle$ ” indicates averaged value) for the three optimized structures of each Pd–X nano-alloy to compare it with that of the pristine  $\text{Pd}_{147}$  nano-cluster. It is seen that the  $\langle\text{Pd–Pd}\rangle$  NN bond-length for the Janus structures, especially for the Pd-rich composition, is less affected with respect to the  $\langle\text{Pd–Pd}\rangle$  NN bond-length of the pristine  $\text{Pd}_{147}$  nano-cluster. For the other two structural patterns, it shows considerable variation for each Pd–X alloy in

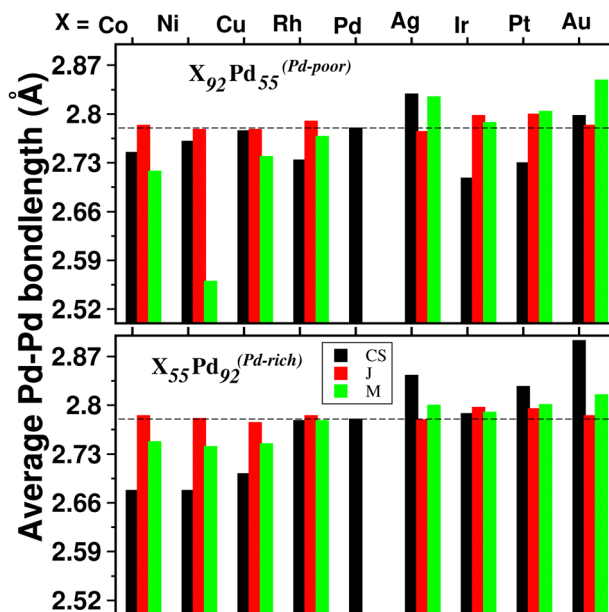


Fig. 4 Bar plot of average Pd–Pd nearest neighbor bond-lengths in the optimized Janus (J), core–shell (CS) and Mixed (M) structures for each bimetallic  $\text{Pd}_{55}\text{X}_{92}$  (top panel) and  $\text{Pd}_{92}\text{X}_{55}$  cluster (bottom panel). The corresponding value for the pristine  $\text{Pd}_{147}$  cluster is indicated by the horizontal dashed line.

either composition. For example, in the case of alloying of Pd with the 3d element, for the Pd-poor composition, the  $\langle\text{Pd–Pd}\rangle$  NN bond-length is reduced (compared to that of the pristine  $\text{Pd}_{147}$  nano-cluster) for the optimized mixed structure and the reduction is the most for the optimized mixed structure of the  $\text{Pd}_{55}\text{Ni}_{92}$  nano-alloy. In the case of the Pd-rich composition of the Pd–3d nano-alloys, both the optimized core–shell and mixed structures have lower  $\langle\text{Pd–Pd}\rangle$  NN bond-length and the reduction is comparatively more for the core–shell structure. In the case of alloying of Pd with the plasmonic elements like Ag and Au, it is seen that the  $\langle\text{Pd–Pd}\rangle$  NN bond-length for both compositions has increased for the optimized core–shell and mixed structures and the enhancement is the most for the core–shell structure of the  $\text{Au}_{55}\text{Pd}_{92}$  nano-alloy. However, the core–shell structure for the bimetallic  $\text{Au}_{55}\text{Pd}_{92}$  nano-alloy with the Au-atoms at the core is not a stable structure as will be seen in the subsequent section.

### 3.2 Trend in relative stability

To analyze the relative stability among the three structures of a binary Pd–X nano-alloy with a fixed composition, we have calculated the total energy difference ( $\Delta E$ ) for the optimized Janus/mixed structure with respect to the optimized core–shell structure, *i.e.*  $\Delta E = E_{\text{Janus/mixed}} - E_{\text{core-shell}}$ ;  $E$  being the DFT calculated total energy for the optimized geometry. A +ve  $\Delta E$  means that the optimized core–shell structure with that chosen ordering for the two species of atoms is more stable, and –ve  $\Delta E$  means that the optimized core–shell structure with such ordering, has lower stability than that of the optimized Janus/mixed structure. Therefore, if  $\Delta E$  is +ve for both the optimized Janus and mixed structures, then the core–shell structure of



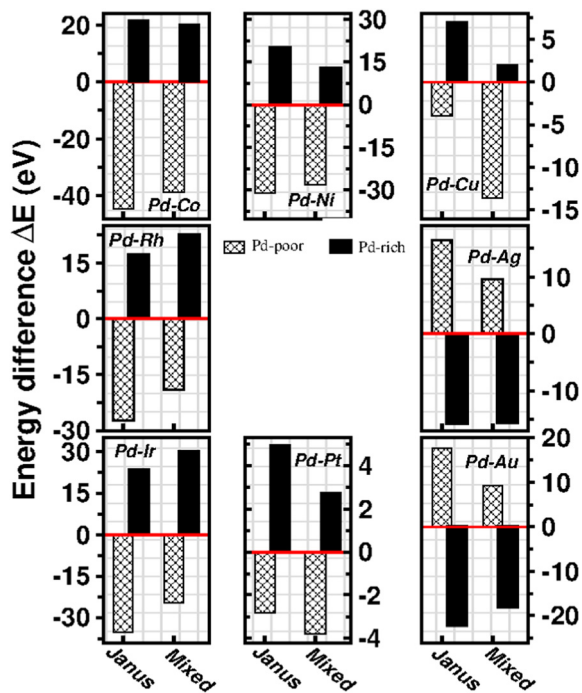


Fig. 5 Bar plot of total energy difference for the optimized Janus and Mixed structures with respect to the optimized core-shell structural pattern for each of the  $\text{Pd}_{55}\text{X}_{92}$  clusters (represented by patterned bars) and  $\text{Pd}_{92}\text{X}_{55}$  clusters (represented by filled bars). The energy of the optimized core-shell structure is set at zero energy as represented by the horizontal red line passing through zero. The panels in the top, middle and bottom rows correspond to the binary Pd-3d, Pd-4d and Pd-5d clusters, respectively.

that composition with that chosen ordering of the two atomic species, is the most stable structure for that binary nano-alloy. On the other hand, if  $\Delta E$  is  $-ve$  for both the optimized Janus and mixed structures, then the core-shell structure of that composition and ordering, will be the least stable. Fig. 5 shows the plot of our calculated  $\Delta E$  for both the compositions of each bimetallic Pd-X nano-alloy. It is seen that Pd-rich binary Pd-Co, Pd-Ni, Pd-Cu and Pd-Pt nano-alloys have a similar trend in the relative stability with respect to the order of preference among the three structural patterns, which is core-shell (Pd-shell) > mixed > Janus. Likewise, the order of the relative stability among the three structural patterns in case of the Pd-rich Pd-Rh and Pd-Ir nano-alloys is core-shell (Pd-shell) > Janus > mixed and that for the Pd-rich binary Pd-Ag and Pd-Au nano-alloys is Janus > mixed > core-shell (Pd-shell). Inversion of the composition affects the order of the relative stability among the three optimized structural patterns as seen for most of the Pd-poor binary nano-alloys. It is seen from Fig. 5 that the order of the relative stability for the Pd-poor bimetallic Pd-Co, Pd-Ni, Pd-Rh and Pd-Ir nanoalloys is similar to Janus > mixed > core-shell (Pd-core); for the Pd-poor bimetallic Pd-Cu and Pd-Pt nano-alloys, it is similar to mixed > Janus > core-shell (Pd-core); for the Pd-poor Pd-Ag and Pd-Au nano-alloys, the order of relative stability is core-shell (Pd-core) > mixed > Janus. For further analysis of the relative stability among the three

structural patterns for each binary Pd-X nano-cluster of fixed composition, we have also calculated their formation energies,<sup>29,30</sup> as given in Section S2 of the ESI.† According to our formulation, a negative formation energy implies a mixing trend between the two species of constituting atoms, while its positive value implies a segregated trend. This analysis shows that the formation energy for each bimetallic Pd-X cluster is most negative for the most stable of the three structures. Our next analysis aims to understand this observed trend in the relative stability in terms of elemental, atomic and electronic properties of the binary nano-alloys.

**3.2.1 Understanding from elemental and atomic properties.** In order to understand the trend in the relative stability, we have first analyzed the variation of the surface energies among the three structures for each bimetallic nano-alloy studied in the present work. Surface energy is the excess energy of the nano-particle surface with respect to its bulk. We have formulated the surface energy per atom<sup>31</sup> for a bimetallic  $A_mB_n$  nano-alloy as  $E_{\text{sur}}(A_mB_n) = \frac{E(A_mB_n) - m\varepsilon_A - n\varepsilon_B}{N_s}$ , where  $E(A_mB_n)$  is the total energy of the binary  $A_mB_n$  nano-alloy and  $\varepsilon_A$  and  $\varepsilon_B$  are the bulk cohesive energies per atom of the A and B elements, respectively;  $N_s$  is the total number of surface atoms in the optimized structure. As we are dealing with icosahedral symmetry-based structures containing an equal number of surface atoms for each of the three structures in every binary nano-cluster studied in this work, the division by  $N_s$  is, therefore, immaterial in the present work. Regarding the effects of surface energy contribution on the relative stability, a higher surface energy usually corresponds to a lower stability and *vice versa*. Our calculated surface energy for the three structures of each bimetallic Pd-X nano-alloy is shown in Fig. 6. The order of energetic stability among the three structures is also mentioned

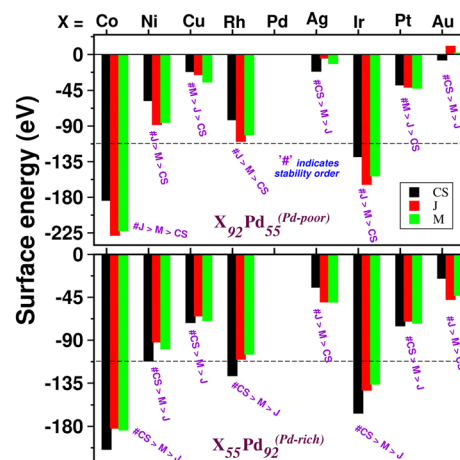


Fig. 6 Bar plot of surface energy for the three structures of each bimetallic Pd-X nano-alloy for both the Pd-poor (top panel) and Pd-rich (bottom panel) compositions. As the number of surface atoms is the same for each system, we have considered here the total surface energy. The surface energy for the pristine  $\text{Pd}_{147}$  cluster is marked with the horizontal dashed line. For each bimetallic nano-alloy, the order of energetic stability is also mentioned by a # symbol.

in Fig. 6 for each binary Pd-X nano-alloy for comparison purposes. It is seen that the structure of the highest stability is associated with the least surface energy for most of the bimetallic nano-alloys in either composition. It should also be noted from Fig. 6 that the magnitude of the surface energy for the bimetallic Pd-Ag and Pd-Au nano-alloys is much smaller, especially for their Pd-poor composition, compared to that of the other bimetallic nano-alloys. This is in accordance with a known fact that Ag and Au atoms tend to segregate mostly at the surface in the case of the Ag or Au-related bimetallic nano-alloys.<sup>32</sup>

In order to clarify the importance of two factors for a bimetallic Pd-X nanoalloy, namely the difference between the two constituting elemental species of their atomic radii (*i.e.*  $\Delta r = r(X) - r(\text{Pd})$ ) and bulk cohesive energies (*i.e.*  $\Delta E_c = E_c(X) - E_c(\text{Pd})$ ), Fig. 7(a) displays the color-coded plot of the most stable structure for each bimetallic nano-alloy of either composition and Fig. 7(b) represents the color-coded plot of the values for the  $\Delta r$  and  $\Delta E_c$  of the corresponding (X,Pd) pair. In the case of either plot, each box except for the central box represents a binary Pd-X nano-alloy. Furthermore, each box is divided into two triangular parts which represent two different entities. In the case of Fig. 7(a), the lower-left triangle and the upper right-triangle of each box represent the stable structure for the Pd-poor and Pd-rich compositions, respectively, of a bimetallic

Pd-X nano-alloy. In the case of Fig. 7(b), each box represents the same Pd-X pair as that of Fig. 7(a). However, it is important to note that the lower-left triangle and upper right triangle now represent the values of  $\Delta E_c$  and  $\Delta r$ , respectively, for the corresponding pair. It is seen from Fig. 7(b) that the binary 3d-Pd pairs have large  $-ve \Delta r$  and this is the main driving factor to decide the stable structure – in the Pd-rich composition, the core-shell structure with larger sized Pd-atoms in the shell acquired the most stability, while in the Pd-poor composition, a core-shell structure with the Pd-atoms in the shell is not allowed to form with our chosen composition in the icosahedral structure. Rather, the relative stability is mostly decided by  $\Delta E_c$  for the Pd-poor composition. Large  $\Delta E_c$  favors the Janus structure for the Pd-poor Pd-Co and Pd-Ni nano-alloys, while small  $\Delta E_c$  favors a mixed structure for the Pd-poor Pd-Cu nano-alloy. For the Rh-Pd, Ir-Pd and Pt-Pd pairs, there is competition between  $\Delta E_c$  and  $\Delta r$  as both are  $+ve$ . For Pd-Ir and Pd-Rh pairs,  $\Delta E_c$  dominates, which favors the core-shell structure with Pd-atoms at the surface for the Pd-rich composition and favors the Janus structure for the Pd-poor composition. For the Pd-rich Pd-Pt nano-alloy,  $\Delta E_c$  wins and stabilizes the core-shell structure with Pd-atoms at the surface. For the Pd-poor Pd-Pt combination, the situation becomes frustrated and the mixed structure gets stabilized. On the other hand, in the case of the Pd-poor Pd-Ag, Pd-Au nano-alloys, the relative stability is dominated by surface segregation of Ag/Au due to their lower surface energy as seen above. Here, the term “surface segregation” implies that the element between the two constituting species, possessing lower surface energy, tends to segregate at the surface of the nano-alloy.

**3.2.2 Understanding from electronic properties.** Further qualitative understanding about the relative stability could be perceived by comparing the relative bond strength and the number of three types of NN dimers present in the three structures of a bimetallic Pd-X nano-alloy for a fixed composition. Regarding the relative strength of the homo-atomic dimers for different elements, the experimental results of dimer binding energies indicate their order of strength as  $\text{Ir}_2$  (3.50 eV)<sup>33</sup> >  $\text{Co}_2$  (3.44 eV)<sup>34</sup> >  $\text{Pt}_2$  (3.14 eV)<sup>35</sup> >  $\text{Rh}_2$  (2.77 eV)<sup>36</sup> >  $\text{Au}_2$  (2.29 eV)<sup>37</sup> >  $\text{Cu}_2$  (2.08 eV)<sup>38</sup> >  $\text{Ni}_2$  (2.06 eV)<sup>39</sup> >  $\text{Ag}_2$  (1.66 eV)<sup>40</sup> >  $\text{Pd}_2$  (1.03 eV).<sup>41</sup> This implies that the Ir-Ir dimer has the strongest binding and the Pd-Pd dimer has the least binding strength. As the Pd-Pd dimer is the weakest homo-atomic dimer, the order of relative bond strength for the three types of dimers in a bimetallic Pd-X nano-alloy is expected to be mostly X-X > Pd-X > Pd-Pd. On the other hand, the number of the three types of NN dimers depends on both the structural pattern and composition. It is seen from Fig. 3 that the order of the number of X-X NN dimers ( $N_{X-X}$ ), Pd-X NN dimers ( $N_{\text{Pd-X}}$ ) and Pd-Pd NN dimers ( $N_{\text{Pd-Pd}}$ ) in the three structures of a bimetallic Pd-X nano-alloy, varies as given below.

For the bimetallic  $\text{Pd}_{55}\text{X}_{92}$  (*i.e.* Pd-poor) nano-alloys, it is seen that:

- $N_{X-X}^{\text{I}} > N_{X-X}^{\text{M}} > N_{X-X}^{\text{C-S}}$
- $N_{\text{Pd-X}}^{\text{M}} > N_{\text{Pd-X}}^{\text{C-S}} > N_{\text{Pd-X}}^{\text{I}}$
- $N_{\text{Pd-Pd}}^{\text{C-S}} > N_{\text{Pd-Pd}}^{\text{I}} > N_{\text{Pd-Pd}}^{\text{M}}$

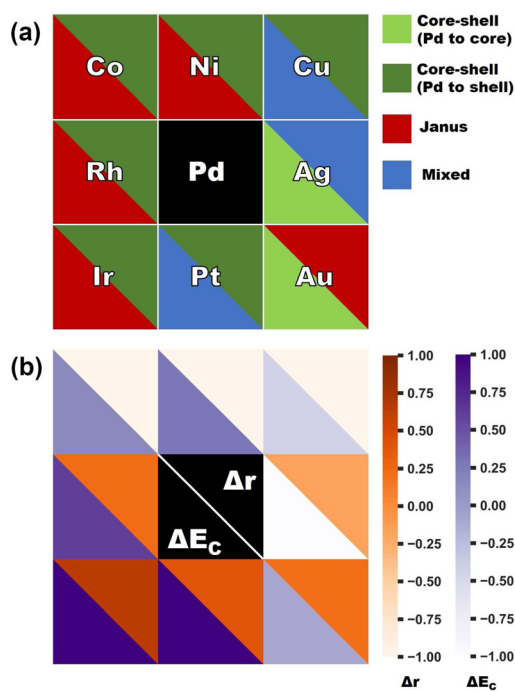


Fig. 7 Color coded plots of (a) the most stable structure for each bimetallic Pd-X nano-alloy of either composition – the lower triangle of each box corresponds to the Pd-poor composition and the upper triangle corresponds to its Pd-rich counterpart, and (b) the difference of atomic radii ( $\Delta r = r(X) - r(\text{Pd})$ ) and bulk cohesive energies ( $\Delta E_c = E_c(X) - E_c(\text{Pd})$ ) between the two constituent elemental species of each Pd-X binary alloy. Each box in either plot represents a binary Pd-X system. Only the “X” element is mentioned in each box of the top figure.

For the bimetallic Pd<sub>92</sub>X<sub>55</sub> (*i.e.* Pd-rich) nano-alloys, it is seen that:

- $N_{X-X}^{C-S} > N_{X-X}^J > N_{X-X}^M$
- $N_{X-Pd}^M > N_{Pd-X}^{C-S} > N_{Pd-X}^J$
- $N_{Pd-Pd}^J > N_{Pd-Pd}^M > N_{Pd-Pd}^{C-S}$

We note that the relative bond strength and the relative number of various types of dimers in a binary nano-alloy can give a qualitative understanding about the trend in the calculated relative stability among the three structural patterns for that nano-alloy system. For example, consider the case of the bimetallic Pd–Co nano-alloy for which the order of bond strength for the three types of NN dimers is Co–Co > Pd–Co > Pd–Pd. In the case of the Pd-poor bimetallic Pd–Co nano-alloy, the order of dominance for the number of strongest Co–Co NN bonds in the three structures is  $N_{Co-Co}^J > N_{Co-Co}^M > N_{Co-Co}^{CS}$ . Interestingly, the order of relative stability follows the same trend, *i.e.* J > M > C–S for the Pd-poor Pd–Co nano-alloy. On the other hand, for the Pd-rich binary Pd–Co nano-alloy, we find that the total number of Pd–Co and Co–Co NN bonds in the three structures varies as  $N_{(Co-Co)+(Pd-Co)}^{C-S} > N_{(Co-Co)+(Pd-Co)}^M > N_{(Co-Co)+(Pd-Co)}^J$ . Again, the trend in the relative stability for the Pd-rich Pd–Co nano-alloy is C–S > M > J. We figure out such similarity for all the bimetallic nano-alloys of either composition, in the trend of variation between the order of relative stability and the order of dominance for a combination of some specific types of NN dimers. Does this similarity in the variation of the two quantities among the three structural patterns indicate some sort of relation between the relative stability and NN bond strength?

In order to investigate the change of atom–atom co-valency among the three structural patterns upon alloying of Pd nano-cluster with the atoms of different elements, we have analyzed the Integrated Crystal Orbital Hamiltonian Population (ICOHP)<sup>22</sup> for the three structures of each bimetallic nano-alloy. COHP is the hopping-weighted density of states between two adjacent atoms and its negative value indicates bonding interactions. ICOHP is the energy integration of all COHP up to

the energy of the highest occupied molecular orbital (HOMO) and it measures the strength of covalency for the bonds. We have calculated ICOHP averaging over all the NN interactions for each bimetallic Pd–X nano-alloy including the pristine Pd<sub>147</sub> cluster. Our calculated ICOHP for the optimized structure of three patterns is shown in Fig. 8 as a function of electron energy for every bimetallic nano-alloy of both compositions. The plots are obtained by averaging the ICOHP values over both the spin channels for every nano-alloy. It is seen from the plot of ICOHP that the structure corresponding to the highest stability is mostly associated with the enhanced covalency (*i.e.* more –ve ICOHP) throughout the bonding region. To compare the values of ICOHP at the Fermi energy among the three structures for every bimetallic nano-alloy, we have calculated the difference of ICOHP for the optimized Janus/mixed structure with respect to the ICOHP for the optimized core–shell structure of the same system, *i.e.*  $\Delta\text{ICOHP} = \text{ICOHP}(\text{Janus/mixed}) - \text{ICOHP}(\text{core-shell})$ . A +ve  $\Delta\text{ICOHP}$  means that the optimized core–shell structure with our chosen ordering for the two species of atoms possesses enhanced covalency, and –ve  $\Delta\text{ICOHP}$  means that the optimized core–shell structure with such ordering, has lower covalency than that of the optimized Janus/mixed structure. It is to be noted that the formulation of  $\Delta\text{ICOHP}$  and the interpretation of its sign have similarities with those of  $\Delta E$  as defined in Section 3.2. Fig. 8 also shows the plot of our calculated  $\Delta\text{ICOHP}$  for all the Pd-poor (patterned bars) and Pd-rich (filled bars) nano-alloys. From the plot of  $\Delta\text{ICOHP}$ , it is interesting to note that the  $\Delta\text{ICOHP}$  for each bimetallic nano-alloyed cluster exhibits a similar trend of variation as seen in the variation of  $\Delta E$  in Fig. 5. It, therefore, clearly indicates a one-to-one correspondence between the relative stability and a change of atom–atom covalency among the three structures. That is, a change in the relative stability among the three structures is intrinsically driven by the change of the atom–atom covalency. The more stable structure is associated with a higher atom–atom covalency and it is the smallest for the least stable structure.

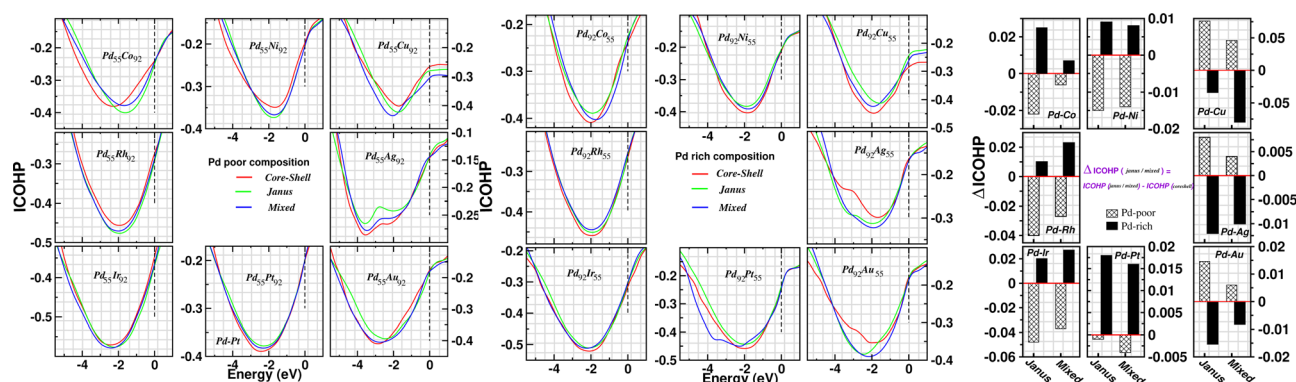


Fig. 8 Plot of ICOHP for the optimized core–shell, Janus and Mixed structures for each bimetallic Pd<sub>55</sub>X<sub>92</sub> cluster (left box) and Pd<sub>92</sub>X<sub>55</sub> cluster (middle box). In the ICOHP plots, the Fermi energy of the corresponding binary nano-cluster is fixed at zero energy as indicated by a vertical dashed line. The right box shows the plot of the difference of ICOHP (*i.e.*  $\Delta\text{ICOHP}$ ) values at the Fermi energy for the optimized Janus and Mixed structures with respect to that of their core–shell structure. The ICOHP of the optimized core–shell structure is set at zero value as represented by the horizontal red line passing through zero. The left, middle and right panels in this box correspond to the binary Pd–3d, Pd–4d and Pd–5d clusters, respectively.

### 3.3 Microscopic characterization of the three morphologies

To characterize the three structural patterns for a bimetallic nano-alloy, one well-known quantity is the radial distribution function (RDF). The RDF for a bimetallic nano-structure is defined as the number of atoms  $N_{A,B}(r)$  of species A or B in a shell at a distance between  $r$  and  $r + \Delta r$  from the center of the cluster<sup>42</sup> *i.e.*  $\rho_{A,B}(r) = \frac{N_{A,B}(r)}{4\pi r^2 \Delta r}$ . In the case of core-shell/Janus binary A-B nano-structures, the first peaks in the plots of  $\rho_A$  and  $\rho_B$  are well separated by a radial distance. In other words, the probability distribution for bimetallic core-shell/Janus nanoparticles exhibits bimodal distributions, while bimetallic alloyed nano-particles display narrow size distributions in the plot of RDF.<sup>43</sup> Furthermore, the ordering of the two elemental species in the core-shell structure could be perceived by analyzing the segregation energy.<sup>44</sup> Likewise, a mixed structure is usually characterized by enhanced mixing energy and hetero-atomic bonds.<sup>6,45</sup> It is, therefore, seen that the core-shell and Janus structural patterns of a bimetallic nano-alloy are not distinctly identified in terms of RDF or mixing energy. In this context, we find that the ‘‘Orbital Hybridization Index’’ could be a useful microscopic quantity that captures a clear distinction among the three structural patterns. As the constituting atomic species have nearly filled valence orbitals, then it is expected that the systems are associated with a large orbital hybridization. To quantify the total orbital hybridization, we have calculated the orbital hybridization index,<sup>45–47</sup> defined as  $h_{kl} = \sum_{l=1}^{147} \sum_{i=1}^{\text{occ}} w_{i,k}^{(l)} w_{i,l}^{(l)}$ ; where  $k$  and  $l$  correspond to the orbital indices – s, p and d, and  $w_{i,k}^{(l)}$  ( $w_{i,l}^{(l)}$ ) is the projection of  $i$ -th Kohn-Sham orbital onto the  $k$  ( $l$ ) spherical harmonic centered at atom  $I$ , integrated over a sphere of specified radius. The spin index is implicit in the summation.

Fig. 9 shows the plot of total hybridization index ( $h_{sd} + h_{pd} + h_{sp}$ ) for each of the  $\text{Pd}_{55}\text{X}_{92}$  (left) and  $\text{Pd}_{92}\text{X}_{55}$  (right) clusters. The value of the total orbital hybridization index for the pristine

$\text{Pd}_{147}$  cluster is also marked with the dashed line for comparison with those of the bimetallic nano-alloys. It is seen that the total hybridization for each of the three structures for every bimetallic nano-alloy has enhanced in either composition compared to that of the pristine  $\text{Pd}_{147}$  cluster, due to the increase of itinerant electrons in the bimetallic nano-alloys. The enhancement of the total hybridization is the most for the bimetallic Pd–Ir nano-alloy due to the presence of more itinerant valence d-electrons of Ir, which results in more p–d and s–d hybridizations, and thereby the total hybridization. On the other hand, the enhancement of orbital hybridization for the Pd–Ag nano-alloy is the least due to reduced itinerancy of the valence d-electrons of Ag. It is seen that the variation of total hybridization index in Fig. 9 is well correlated with the variation of cluster size in Fig. 2 in the sense that a lower hybridization results a higher cluster size and *vice versa*. The interesting point to note from the plot of the total hybridization index within a fixed composition is that the optimized Janus structure for most of the studied bimetallic nano-alloys possesses relatively higher value compared to that of the optimized core-shell and mixed structures for the respective system. In addition, the optimized core-shell structure has the least total orbital hybridization index, while the optimized mixed structure has intermediate total orbital hybridization for each bimetallic cluster. The three structural patterns for a bimetallic nano-alloy, therefore, possess a clear trend in the variation of total orbital hybridization index – most for Janus, intermediate for mixed and least for core-shell patterns.

### 3.4 Trend in catalytic properties

The catalytic activity of the catalysts for a given reaction relies on the activation barrier of the rate-determining step or absorption ability of an intermediate included in the rate-determining step.<sup>48,49</sup> For a catalytic process over a transition metal, the adsorption energies of different molecules scale with each other as do the transition state energies and reaction energies.<sup>50–54</sup> This underlying concept is the basis of the d-band model. Instead of determining the energy of the intermediates and transition states for each catalyst, the celebrated d-band center model proposed by Hammer and Norskov<sup>23</sup> gives a straight forward way to relate the rate of catalytic activity to the absorption energy and thereby the electronic structure of the transition metal catalyst’s surface. Therefore, the position of the d-band center with respect to the Fermi energy is regarded as a first order descriptor of the propensity of a transition metal catalyst to bind to adsorbates. According to this model, the trend in binding of the adsorbate on the surface of the catalyst is an effective indicator about the rate of the catalytic activity for the catalytic surface. For an efficient catalytic process, this bond-strength of the adsorbate with the catalytic surface will not be too strong or too weak. This is due to the fact that if this bonding is very strong, then the barrier height for the desorption process, *i.e.* the removal of the adsorbate from the catalytic surface after the reaction is very high. Likewise, if the bonding is very weak, then the barrier height for the absorption process is very high which limits the

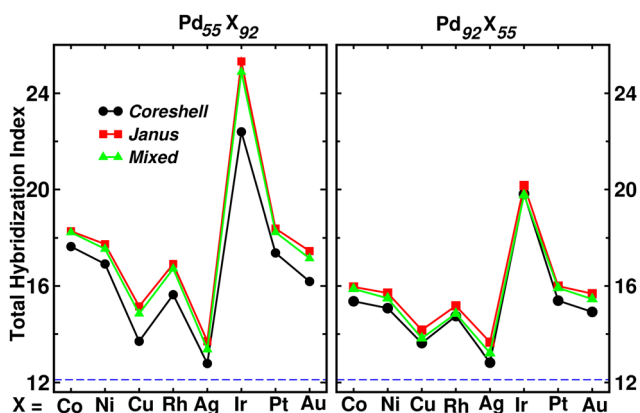


Fig. 9 Plot of total orbital hybridization index corresponding to the optimized core-shell (black curve), Janus (red curve) and Mixed (green curve) structures of each  $\text{Pd}_{55}\text{X}_{92}$  (left) and  $\text{Pd}_{92}\text{X}_{55}$  (right) cluster. The value for the pristine  $\text{Pd}_{147}$  cluster is indicated by the horizontal dashed line.



overall reaction kinetics too. In fact, this is readily understandable as the plot of the catalyst's activity against the binding of the adsorbate results mostly in a so-called *Volcano curve* for efficient catalytic reactions.<sup>55</sup> It should be mentioned that the *Volcano plot* gives an effective measure for catalytic activity, which has been extensively used for various heterogeneous catalytic reactions.<sup>56</sup> The d-band center model is a more fundamental approach to predict catalytic activity as it correlates the adsorbate–surface bond strength, *i.e.* the macroscopic catalytic performance, with the microscopic electronic structure of the surface atoms for the transition metal catalysts. According to this model, an up-shifting of the occupied d-band center towards the Fermi energy corresponds to a stronger absorption strength and *vice versa*. For our binary nano-alloys, we have calculated the position of the occupied d-band center as  $d_{\text{center}} = \frac{\int_0^{E_{\text{HOMO}}} \rho E dE}{\int_0^{E_{\text{HOMO}}} \rho dE}$ , where  $\rho$  = d-density of states,  $E$  = the d-level energy and  $\rho dE$  = the number of d-states in the

energy interval  $E$  and  $E + dE$ . Our calculated position for the occupied d-band center of the three structures for each bimetallic nano-alloy, with respect to the respective  $E_{\text{HOMO}}$  (fixed at zero), is plotted in Fig. 10. It should be noted here that the calculated d-band center position is averaged over both the spin channels. The corresponding values for the pristine Pd<sub>147</sub> and Pt<sub>147</sub> clusters are indicated with horizontal dashed lines to compare them with the d-center positions of the studied bimetallic nano-alloys. For further clarification, the d-orbital projected density of states (d-PDOS) for the three structures of each bimetallic nano-alloy is also shown in the right panel of Fig. 10 and the occupied d-states for the structure having the maximum up-shifting of the occupied d-states is marked with the shaded region.

From our plot of the occupied d-center position in Fig. 10, it is seen that the positions of the occupied d-band center for either composition of the bimetallic Pd–Ni nano-alloys, *i.e.* Pd<sub>55</sub>Ni<sub>92</sub> and Pd<sub>92</sub>Ni<sub>55</sub> nano-clusters, have up-shifted for each structure with respect to that of the pristine Pd<sub>147</sub> cluster. This

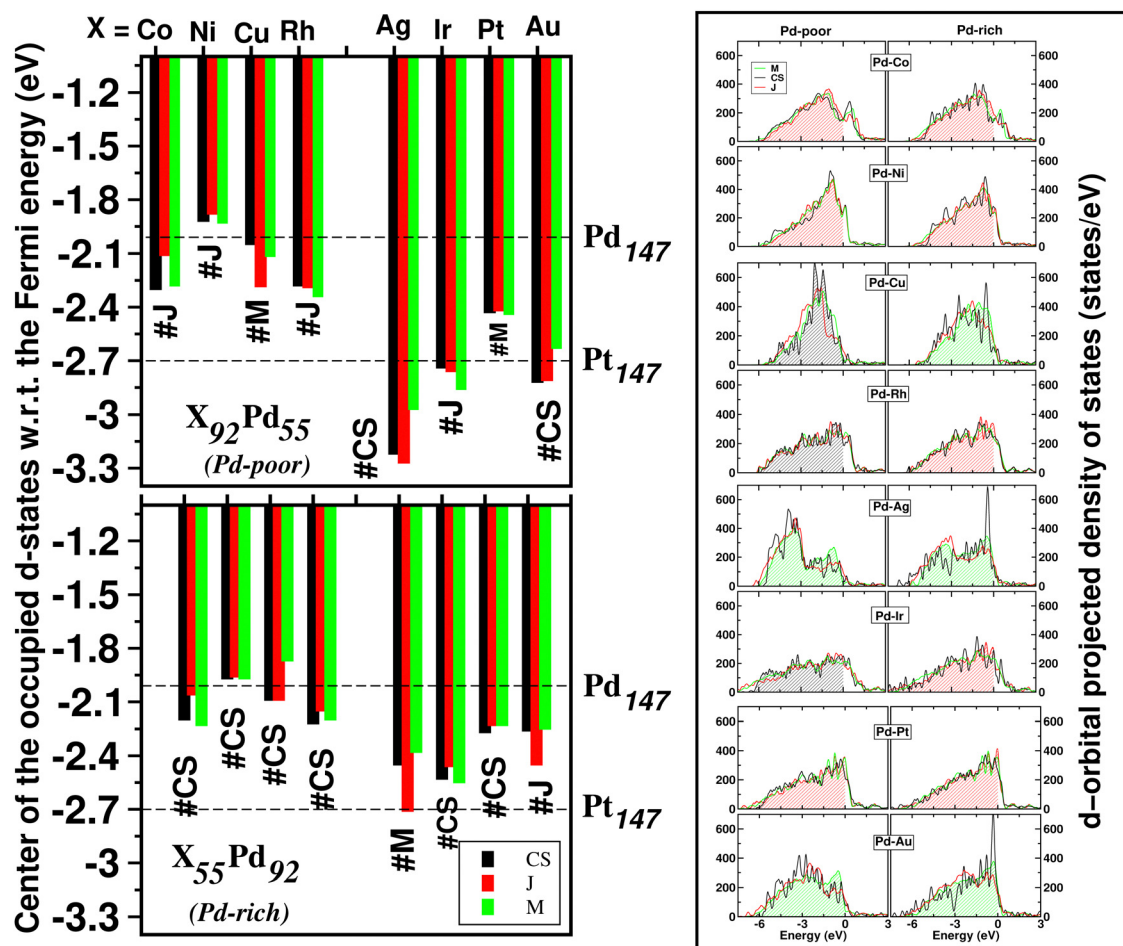


Fig. 10 (left panel) Bar plot for the position of the center of gravity of the occupied d-band with respect to the energy of the highest occupied level. The corresponding values for the pristine icosahedral Pt<sub>147</sub> and Pd<sub>147</sub> clusters are marked by dashed horizontal lines for comparison with the other systems. The calculated value for each structure is averaged over both the spin channels. The most stable structure among the three structural patterns is also mentioned for each bimetallic nano-alloy of both compositions. The right panel shows the d-PDOS for the three structures for each bimetallic nano-alloy. The shaded region indicates the occupied d-states for the structure having the most up-shifting of the occupied d-center. In this d-DOS plot, the Fermi energy is fixed at the energy and the smearing width is fixed at 0.1 eV.

implies that the bimetallic Pd–Ni nano-alloy will have stronger interaction with the adsorbate. Similar up-shifting is also found for the mixed structure of the Pd-rich bimetallic Pd–Cu nano-alloy, *i.e.* for the Pd<sub>92</sub>Cu<sub>55</sub> cluster. For the rest of the studied bimetallic nano-alloys, the occupied d-band center positions fall in-between those of the pristine Pd<sub>147</sub> and Pt<sub>147</sub> clusters. Consequently, their catalytic activity is also expected to be intermediate of those of the pristine Pt<sub>147</sub> and Pd<sub>147</sub> clusters. Furthermore, it is clearly seen from Fig. 10 that the difference in the positions of the occupied d-band center among the three structures is very small for the bimetallic Pd–Ni and Pd–Pt nano-alloys in either composition limits. Therefore, it indicates that the catalytic activity for these systems is not affected much with the change of their structural patterns. In reality, the catalytic activity, however, depends mostly on the availability of the catalytically active sites on the surface. With respect to the importance of the structural pattern on the catalytic performance, our analysis of d-band center position suggests that:

**A. For the Pd-poor composition.** I. The core–shell structure is predicted to have stronger reaction with the adsorbate molecule in the cases of the bimetallic Pd–Cu, Pd–Rh and Pd–Ir nanoalloys,

II. the Janus structure will have strong adsorbate–cluster interaction in the cases of the bimetallic Pd–Co, Pd–Ni and Pd–Pt nanoalloys,

III. the mixed structure in the cases of the bimetallic Pd–Ag, Pd–Au nanoalloys.

Considering the structure which has the most up-shifting of the occupied d-band center position, among the Pd-poor nano-alloys, *i.e.* the X<sub>92</sub>Pd<sub>55</sub> nanoalloys, the predicted up-shifting

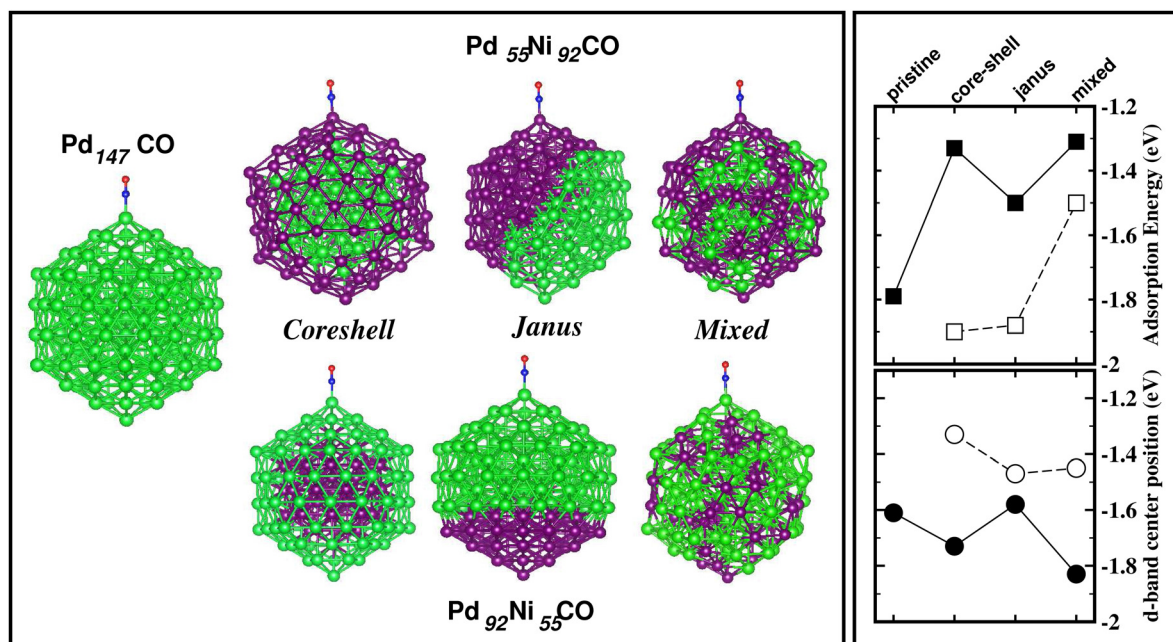
order would be X = Ni(J) > Cu(C–S) > Co(J) > Rh(C–S) > Pt(J) > Au(M) > Ir(C–S) > Ag(M). Here “C–S” stands for core–shell structure, “J” for Janus structure and “M” for Mixed structure.

**B. For the Pd-rich composition.** I. The Janus structure is predicted to have stronger reaction with the adsorbate molecule in the cases of the bimetallic Pd–Co, Pd–Ni, Pd–Rh, Pd–Ir and Pd–Pt nanoalloys,

II. the mixed structure in the cases of the bimetallic Pd–Cu, Pd–Ag and Pd–Au nanoalloys.

Considering the structure which has the most up-shifting of the occupied d-band center position among the Pd-rich nano-alloys, *i.e.* the X<sub>55</sub>Pd<sub>92</sub> nanoalloys, the predicted up-shifting order would be X = Cu(M) > Ni(J) > Co(J) > Rh(J) > Pt(J) > Au(M) > Ag(M) > Ir(J).

The energetically most stable structure for each bimetallic nano-alloy of either composition is also mentioned in Fig. 10 with the # symbol. In most of the cases, the catalytically active structure does not correspond to the energetically most stable structure. This is in accordance with the fact that a non-equilibrium system is prone to enhanced catalytic activity.<sup>57</sup> Moreover, it is seen that the Janus structure appears more frequently as the catalytically most active structure, which may result from the presence of a large strained interface between two species of atoms in the Janus structure.<sup>58</sup> According to our analysis of the occupied d-band center position, the bimetallic Pd–Ni nano-alloy has the most up-shifting of the occupied d-center towards its Fermi level for both compositions. For direct comparison of the trend towards catalytic activity relative to that of the pristine Pd<sub>147</sub> cluster, we performed rigorous analysis of the absorption of a CO molecule on each structure



**Fig. 11** Plot of (left panel) the optimized configurations for CO-absorption on a surface Pd-atom at the atop position of the pristine Pd<sub>147</sub> cluster and the three structures of the Pd<sub>92</sub>Ni<sub>55</sub> cluster; (right panels) our calculated absorption energy for the four systems (top) and the position of the d-band center with respect to the Fermi energy for the respective absorption site (bottom). In the right panels, the filled symbols correspond to the pristine Pd<sub>147</sub> cluster and the Pd-rich Pd–Ni nano-alloy, the empty symbols correspond to the Pd-poor Pd–Ni nano-alloy.

of the bimetallic Pd–Ni nano-alloy of either composition. To simplify our analysis, we consider absorption of a carbon-monoxide molecule at the atop position on a least coordinated vertex site: on a surface Pd-site in the case of the pristine Pd<sub>147</sub> cluster and the three structures of Pd-rich bimetallic Pd–Ni cluster, *i.e.* Pd<sub>92</sub>Ni<sub>55</sub> cluster. In the case of the three structures of Pd<sub>55</sub>Ni<sub>92</sub> cluster, we consider absorption on a surface Ni-site at the vertex position. The optimized CO-absorbed structures are shown in Fig. 11. It should be noted here that previous works also found that absorption of a CO molecule at the atop position on a Pd-surface, is energetically more favorable.<sup>59</sup>

Then, we calculated the absorption energy for the CO-absorbed systems using the formulation  $E_{\text{abs}} = E(\text{Ni}_m\text{Pd}_n\text{CO}) - E(\text{Ni}_m\text{Pd}_n) - E(\text{CO})$ . To understand the trend in the calculated absorption energy, we have also analyzed the occupied d-band center position for the absorption site. Our calculated absorption energy and the position of the occupied d-band center for the absorption site, are also shown in Fig. 11. First, it should be mentioned that our calculated absorption energy for the pristine Pd<sub>147</sub> cluster [ $E_{\text{abs}}(\text{Pd}_{147}) = 1.79$  eV] is comparable with the previously calculated absorption energy of a CO molecule on a Pt(111) surface [ $E_{\text{abs}}(\text{Pt}(111)/\text{CO}) = 1.64$  eV]<sup>60</sup> or Pt<sub>55</sub> nano-cluster [ $E_{\text{abs}}(\text{Pt}_{55}\text{CO}) = 2.04\text{--}2.37$  eV] by DFT-PBE calculations.<sup>59</sup> From our plot of the calculated absorption energy in Fig. 11, it is seen that the overall magnitude of the absorption energy for the Pd-poor Pd–Ni nano-alloy is more than that of the Pd-rich Pd–Ni nano-alloy. This is due to the fact that the bond-strength of the CO molecule with a surface Ni atom is more than that with a surface Pd atom. The right panel of Fig. 11 also shows the position of the occupied d-band center relative to the Fermi energy for the atom at the absorption site. Interestingly, the occupied d-band center is up-shifted more for each structure of Pd-poor bimetallic Pd–Ni nano-alloy compared to that of the Pd-rich Pd–Ni nano-alloy, which is consistent with the enhanced absorption trend for the Pd-poor bimetallic Pd–Ni nano-alloy.

### 3.5 Exploring size effects

So far, the calculated results as discussed above are based on the cluster size of 147 atoms. In order to explore the possible cluster size effects on the studied properties, we have repeated the analysis of the relative stability in terms of energy difference, formation energy and surface energy; the trends in the orbital hybridization and catalytic activity considering a 55-atom sized icosahedral geometry only for the bimetallic Pd–Ni nano-alloy. It should be mentioned here that a 55-atom sized icosahedron is the second closed-shell structure containing 13 atoms inside and 42 atoms on the outermost surface. Therefore, we focus on the three structural patterns for two compositions – Pd<sub>13</sub>Ni<sub>42</sub> (Pd-poor) and Pd<sub>42</sub>Ni<sub>13</sub> (Pd-rich) nano-alloys. The calculated results are plotted in Fig. 12. It is seen that the predicted order of relative stability from the analysis of  $\Delta E$  and formation energy, is  $J > M > \text{CS}$  for the Pd<sub>13</sub>Ni<sub>42</sub> cluster and  $\text{CS} > M > J$  for the Pd<sub>42</sub>Ni<sub>13</sub> nano-cluster. In addition, the observed trend in the relative stability is in accordance with the variation of their surface energy. Furthermore, the variation of orbital hybridization among the three structures of the 55-atom

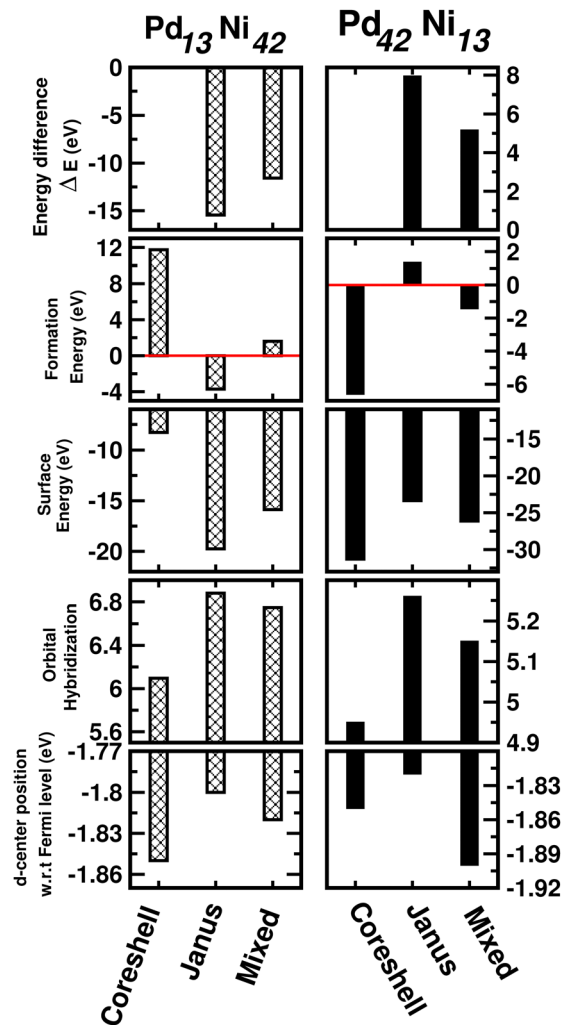


Fig. 12 Plot of total energy difference as defined in Section 3.2, formation energy as defined in the ESI,† surface energy as defined in Section 3.2.1, orbital hybridization as defined in Section 3.3 and position of the d-band center as defined in Section 3.4 for the optimized core–shell, Janus and Mixed structures of the 55-atom sized bimetallic Pd–Ni cluster in both compositions, *i.e.* Pd<sub>13</sub>Ni<sub>42</sub> (left panel) and Pd<sub>42</sub>Ni<sub>13</sub> clusters (right panel).

sized binary Pd–Ni nano-cluster also indicates the most total hybridization for the Janus structure and least for the core–shell structure. Moreover, the position of the occupied d-center with respect to the Fermi level indicates the same trend of catalytic activity among the three structures of the 55-atom PdNi nano-alloy as seen for the 147-atom binary PdNi cluster with either composition limit. Therefore, this analysis indicates that the cluster size has no significant effect on the overall trend of variation for the studied properties among the three structural patterns for the binary Pd–X nano-alloys.

## 4 Summary and conclusions

Using DFT calculations, we have analyzed the relative structural stability, electronic and catalytic properties among the core–shell, Janus and mixed structural patterns for eight bimetallic Pd–X nano-alloys (X = Co, Ni, Cu, Rh, Ag, Ir, Pt, and Au)

considering two compositions – a Pd-poor and a Pd-rich compositions. It is seen that the relative stability depends on the composition and the type of “X” element. The driving factors to determine the relative stability are found to be: the atomic-radius difference of the two constituting elemental species in the cases of Pd-3d pairs; competition of the differences in their atomic radii and also in their bulk cohesive energies for the bimetallic Pd–Rh, Pd–Ir and Pd–Pt pairs; and difference of surface energies in the case of Pd–Ag and Pd–Au pairs. Our analysis of Integrated Crystal Orbital Hamiltonian Population indicates that the relative structural stability among the three patterns is intrinsically governed by the change of atom–atom covalency – a more stable structure is associated with higher atom–atom covalency and it is smaller for the least stable structure. Furthermore, the three structures of each bimetallic Pd–X nano-alloy can be distinctly identified in terms of an orbital hybridization index – our analysis of total orbital hybridization index, shows that the Janus, mixed and core-shell structures are associated with the highest, intermediate and least total orbital hybridization, respectively. The trend in the catalytic activity for the bimetallic Pd–X nano-alloys is predicted by analyzing the shifting of the occupied d-band center position. We find that most of the Pd–X nano-alloys are expected to have catalytic activity in-between those of pristine Pt and Pd clusters of the same size. Among the three structures, the catalytically most active structure for the Pd–X nano-alloys is mostly different from the energetically most stable structure. The bimetallic Pd–Ni nano-alloy exhibits the most up-shifting of the occupied d-band center position, thereby implying stronger interaction with an adsorbate molecule. Our analysis of absorption reaction for a CO molecule on a Pd–Ni cluster exhibits a similar trend of variation for the absorption energy as predicted from the analysis of the occupied d-band center shifting. To verify the cluster size effect on the studied properties, we have further repeated our analysis for two compositions of a 55-atom sized binary Pd–Ni cluster. Our study in the present work indicates that the bimetallic Pd–Ni nano-alloys are expected to be the best combination to develop stable and catalytically active non-Pt nano-alloys.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

T. S.-D. and S. D. acknowledge funding from J. C. Bose National Fellowship (JCB/2020/000004). AG acknowledges a DST-INSPIRE Fellowship for funding.

## References

- M. B. Gawande, A. Goswami, T. Asefa, H. Guo, A. V. Biradar, D.-L. Peng, R. Zboril and R. S. Varma, *Chem. Soc. Rev.*, 2015, **44**, 7540–7590.
- J. Lim, B. G. Jeong, M. Park, J. K. Kim, J. M. Pietryga, Y.-S. Park, V. I. Klimov, C. Lee, D. C. Lee and W. K. Bae, *Adv. Mater.*, 2014, **26**, 8034–8040.
- J. Rick, M.-C. Tsai and B. J. Hwang, *Nanomaterials*, 2015, **6**, 5, DOI: [10.3390/nano6010005](https://doi.org/10.3390/nano6010005).
- K. McNamara, S. A. M. Tofail, N. D. Thorat, J. Bauer and J. J. E. Mulvihill, In *Biomedical Applications of Nanoalloys. Nanoalloys: From Fundamentals to Emergent Applications*, 2nd edn, ed. F. Calvo, Elsevier, 2020, pp. 381–432.
- S. T. Selvan, P. K. Patra, C. Y. Ang and J. Y. Ying, *Angew. Chem., Int. Ed.*, 2007, **46**, 2448–2452; S. Duan and R. Wang, *Prog. Nat. Sci.: Mater. Int.*, 2013, **23**(2), 113–126.
- R. Ferrando, J. Jellinek and R. L. Johnston, *Chem. Rev.*, 2008, **108**, 845–910.
- I. Chakraborty and T. Pradeep, *Chem. Rev.*, 2017, **117**, 8208–8271.
- L.-L. Wang and D. D. Johnson, *J. Am. Chem. Soc.*, 2009, **131**, 14023–14029.
- A. Ghosh, S. Datta and T. Saha-Dasgupta, *J. Phys. Chem. C*, 2022, **126**, 6847–6853.
- S. F. Tan, G. Bisht, U. Anand, M. Bosman, X. E. Yong and U. Mirsaidov, *J. Am. Chem. Soc.*, 2018, **140**, 11680–11685.
- J. M. Rahm and P. Erhart, *J. Phys. Chem. C*, 2018, **122**, 28439–28445.
- C. Langlois, Z. L. Li, J. Yuan, D. Alloyeau, J. Nelayah, D. Bochicchio, R. Ferrando and C. Ricolleau, *Nanoscale*, 2012, **4**, 3381–3388.
- H. Peng, W. Qi, S. Li and W. Ji, *J. Phys. Chem. C*, 2015, **119**, 2186–2195.
- D. Bochicchio and R. Ferrando, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2013, **87**, 165435.
- E. Maras, F. Berthier and B. Legrand, *J. Phys. Chem. C*, 2016, **120**, 22670–22680.
- T. Wang, A. Chutia, D. J. L. Brett, P. R. Shearing, G. He, G. Chai and I. P. Parkin, *Energy Environ. Sci.*, 2021, **14**, 2639–2669.
- D. N. Oleksyszyn, B. L. Albuquerque, D. de O. Silva, G. L. Tripodi, D. C. de Oliveira and J. B. Domingos, *Nanoscale*, 2020, **12**, 1171; S. V. Myers, A. I. Frenkel and R. M. Crooks, *Chem. Mater.*, 2009, **21**, 4824–4829.
- C. Zhu, J. Zeng, J. Tao, M. C. Johnson, I. Schmidt-Krey, L. Blubaugh, Y. Zhu, Z. Gu and Y. Xia, *J. Am. Chem. Soc.*, 2012, **134**, 15822–15831; N. A. Khan, A. Uhl, S. Shaikhutdinov and H.-J. Freund, *Surf. Sci.*, 2006, **600**, 1849–1853.
- A. F. Carlsson, M. Baumer, T. Risse and H.-J. Freund, *J. Chem. Phys.*, 2003, **119**, 10885; R. Raja, V. B. Golovko, J. M. Thomas, A. Berenguer-Murcia, W. Zhou, S. Xie and B. F. G. Johnson, *Chem. Commun.*, 2005, 2026–2028, DOI: [10.1039/B418273A](https://doi.org/10.1039/B418273A).
- S. Datta, M. Kabir and T. Saha-Dasgupta, *J. Phys. Chem. C*, 2008, **112**, 8206–8214; S. Datta, M. Kabir and T. Saha-Dasgupta, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2012, **86**, 115307.
- T. Teranishi and M. Miyake, *Chem. Mater.*, 1998, **10**(2), 594–600; A. Front and C. Mottet, *Theor. Chem. Acc.*, 2022, **141**, 2.



- 22 S. Maintz, V. L. Deringer, A. L. Tchougreff and R. Dronskowski, *J. Comput. Chem.*, 2016, **37**, 1030–1035.
- 23 J. K. Nørskov, F. Abild-Pedersen and T. Bligaard, *Proc. Natl. Acad. Sci. U. S. A.*, 2011, **108**, 937.
- 24 S. Zhou, C. Wang and Y. Xia, *Chem. Mater.*, 2022, **34**(11), 5065–5073.
- 25 *Vienna ab initio simulation package (VASP)*, Technische Universität Wien, 1999; G. Kresse and J. Hafner, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1993, **47**, 558; G. Kresse and J. Furthmuller, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1996, **54**, 11169.
- 26 P. E. Blöchl, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1994, **50**, 17953.
- 27 G. Kresse and D. Joubert, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1999, **59**, 1758.
- 28 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865.
- 29 L. G. Wang and A. Zunger, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2003, **67**, 092103.
- 30 M. H. F. Sluiter, C. Colinet and A. Pasturel, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2006, **73**, 174204.
- 31 S. Datta, M. Kabir, T. Saha-Dasgupta and D. D. Sarma, *J. Phys. Chem. C*, 2008, **112**(22), 8206; Z. Valizadeh and M. Abbaspour, *J. Phys. Chem. Solids*, 2020, **144**, 109480.
- 32 J. M. Rahm and P. Erhart, *J. Phys. Chem. C*, 2018, **122**(49), 28439–28445.
- 33 Z. J. Wu, B. Han, Z. W. Dai and P. C. Jin, *Chem. Phys. Lett.*, 2005, **403**, 367–371.
- 34 A. Kant and B. Strauss, *J. Chem. Phys.*, 1964, **41**, 3806.
- 35 A. Crushow and K. M. Ervin, *J. Chem. Phys.*, 1997, **106**, 9580.
- 36 C.-H. Chien, E. Blaisten-Barojas and M. R. Pederson, *J. Chem. Phys.*, 2000, **112**, 2301.
- 37 G. A. Bishea and M. D. Morse, *J. Chem. Phys.*, 1991, **95**, 5646.
- 38 E. A. Rohlfing and J. J. Valentini, *J. Chem. Phys.*, 1986, **84**, 6560.
- 39 E. Curotto, A. Matro, D. L. Freeman and J. D. Doll, *J. Chem. Phys.*, 1998, **108**, 729.
- 40 R. Fournier, *J. Chem. Phys.*, 2001, **115**, 2165.
- 41 J. Ho, K. M. Ervin, M. L. Polak, M. K. Gilles and W. C. Lineberger, *J. Chem. Phys.*, 1991, **95**, 4845.
- 42 F. Baletto, C. Mottet and R. Ferrando, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2002, **66**, 155420.
- 43 G. D. Förster, M. Benoit and J. Lam, *Phys. Chem. Chem. Phys.*, 2019, **21**, 22774.
- 44 L.-L. Wang and D. D. Jonson, *J. Am. Chem. Soc.*, 2009, **131**, 14023–14029.
- 45 S. Datta and T. Saha-Dasgupta, *J. Phys.: Condens. Matter*, 2013, **25**, 225302.
- 46 H. Häkkinen, M. Moseler and U. Landman, *Phys. Rev. Lett.*, 2002, **89**, 033401; S. Datta, *AIP Adv.*, 2019, **9**, 115316.
- 47 C. M. Chang and M. Y. Chou, *Phys. Rev. Lett.*, 2004, **93**, 133401; S. Datta, M. Kabir and T. Saha-Dasgupta, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2011, **84**, 075429.
- 48 C. Stegelmann, A. Andreasen and C. T. Campbell, *J. Am. Chem. Soc.*, 2009, **131**, 8077.
- 49 R. Pestman, W. Chen and E. Hensen, *ACS Catal.*, 2019, **9**, 4189.
- 50 J. K. Nørskov, T. Bligaard, A. Logadottir, S. Bahn, L. B. Hansen, M. Bollinger, H. Bengaard, B. Hammer, Z. Sljivancanin, M. Mavrikakis, Y. Xu, S. Dahl and C. J. H. Jacobsen, *J. Catal.*, 2002, **209**, 275.
- 51 A. Michaelides, Z.-P. Liu, C. J. Zhang, A. Alavi, D. A. King and P. Hu, *J. Am. Chem. Soc.*, 2003, **125**, 3704.
- 52 T. Bligaard, J. K. Nørskov, S. Dahl, J. Matthiesen, C. H. Christensen and J. Sehested, *J. Catal.*, 2004, **224**, 206.
- 53 F. Abild-Pedersen, J. Greeley, F. Studt, J. Rossmeisl, T. R. Munter, P. G. Moses, E. Skúlason, T. Bligaard and J. K. Nørskov, *Phys. Rev. Lett.*, 2007, **99**, 016104.
- 54 E. M. Fernandez, P. G. Moses, A. Toftelund, H. A. Hansen, J. I. Martinez, F. Abild-Pedersen, J. Kleis, B. Hinnemann, J. Rossmeisl, T. Bligaard and J. K. Nørskov, *Angew. Chem., Int. Ed.*, 2008, **47**, 4683.
- 55 Z. W. Seh, J. Kibsgaard, C. F. Dickens, I. Chorkendorff, J. K. Nørskov and T. F. Jaramillo, *Science*, 2017, **355**(6321), 379.
- 56 A. J. Medford, P. G. Moses, K. W. Jacobsen and A. A. Peterson, *ACS Catal.*, 2022, **12**(15), 9679.
- 57 R. Chen, S. Neri and L. J. Prins, *Nat. Nanotechnol.*, 2020, **15**, 868–875.
- 58 T. Kim, J. Park, Y. Hong, A. Oh, H. Baik and K. Lee, *ACS Nano*, 2019, **13**, 11834; H.-Y. He, Z. He and Q. Shen, *Colloid Interface Sci. Commun.*, 2021, **45**, 100509.
- 59 D. Guedes-Sobrinho, R. L. H. Freire, A. S. Chaves and J. L. F. Da Silva, *J. Phys. Chem. C*, 2017, **121**(49), 27721.
- 60 L. Schimka, J. Harl, A. Stroppa, A. Gruneis, M. Marsman, F. Mittendorfer and G. Kresse, *Nat. Mater.*, 2010, **9**, 741; Y. Wang, S. de Gironcoli, N. S. Hush and J. R. Reimers, *J. Am. Chem. Soc.*, 2007, **129**, 10402.