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# Prospects of copper–bismuth chalcogenide absorbers for photovoltaics and photoelectrocatalysis

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Wittichenite,  $\text{Cu}_3\text{BiS}_3$  (CBS), is perceived as an appealing material for application as a photoabsorber due to its  $ns^2$  electronic configuration and antibonding character on its upper valence band, which provides defect-tolerance characteristics. Herein, we highlight the high charge carrier mobility and low effective mass due to the high dispersive character of both the valence and conduction bands in CBS. The absorption coefficient of CBS is as high as  $>10^5 \text{ cm}^{-1}$  (energy  $> 1.43 \text{ eV}$ ); however, it is still poorly explored in photovoltaic (PV) (maximum power conversion efficiency of 1.7%) and photoelectrochemical (PEC) devices. The challenges and questions related to the nature of the bandgap, intrinsic defects, and charge transport in CBS need to be addressed to enhance its PV and PEC performance. In this review, we provide an overview of the crystal and energetic structure, optoelectronic properties, recent advances in synthesis methods, and experimental realization of devices based on CBS. Further, we outline the important research directions that are indispensable for the development of CBS-based devices.

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## 1 Introduction

The global temperature rise due to anthropogenic carbon emissions and associated climate change continue to pose a threat to our future, while global energy demands continue to increase. Thus, concerted efforts are required towards decarbonization to

restrain the global temperature rise to less than  $2 \text{ }^\circ\text{C}$  (in accordance with the 2015 Paris Agreement). In this case, an effective solution to the energy requirements and environmental conservation is the use of solar energy, which is a renewable and clean source of energy and nearly inexhaustible.

Solar energy can be harnessed through different approaches based on (i) the direct conversion of sunlight into electricity

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using semiconducting (SC) absorbers in photovoltaic (PV) cells and (ii) solar-driven chemical reactions to store energy in the form of chemical bonds, which are also known as solar fuels. The fundamental component of the aforementioned systems is the semiconducting light absorber, which is capable of efficiently absorbing light and allowing the flow of photogenerated charge carriers. From a material standpoint, an ideal photo-absorber should have superior optoelectronic properties such as high optical absorption in the visible spectrum, suitable energy band gap, long charge carrier diffusion lengths, *etc.* In addition, it should preferably be cheap, environmentally benign, and stable. Over the past few decades, materials systems such as silicon, oxides, nitrides, chalcogenides, and recently hybrid organic–inorganic halide perovskites have been extensively explored for solar cells and photoelectrochemical energy conversion. Among the inorganic absorbers, high-performing thin film solar cells are realized using direct band gap materials with high optical absorption. The typical materials employed for the fabrication of thin film solar cells include  $\text{CuIn}_x\text{Ga}_{1-x}(\text{S,Se})_2$  (CIGS), kesterites, and  $\text{Cu}_2\text{ZnSn}(\text{S,Se})_4$  (CZTS/CZTSe). To date, a power conversion efficiency (PCE) of 23.4% and 22.4% has been achieved for rigid and flexible CIGS solar cells, respectively. However, despite its advantages, CIGS still lags behind Si solar cell technology (26.1% PCE).<sup>1</sup> Moreover, the presence of scarce (In, indium) and toxic elements (Cd, cadmium) poses environmental and cost concerns, which need to be addressed. The latter two features are very important and must be competitive with the primary power sources generated by fossil fuels (*i.e.*, coal, oil, and natural gas)<sup>2</sup> to have global market acceptance and become commercialized. Consequently, most of the proposals offered on this front have focused on searching for alternative materials to address these limitations.

Kesterites are considered interesting candidates to replace CIGS as a photoabsorber layer given that they have nontoxic and earth-abundant constituents, addressing the concerns associated with toxicity and scarcity. However, due to their quaternary nature, kesterites are associated with problems such as secondary phases, intrinsic point defects (*i.e.*, vacancies, interstitials, and anti-sites), and multivalence and structural polymorphism of Sn, making their production a challenge.<sup>3–5</sup> Therefore, more research is

required to for the appropriate material selection that fulfils the above-mentioned criteria.

Recently, an emerging class of bismuth-based chalcogenide materials have been gaining considerable attention due to their interesting optoelectronic properties and performance together with enhanced understanding assisted by theory. They include  $\text{AgBiS}_2$ ,  $\text{NaBiS}_2$ ,  $\text{BiS}_3$ , and  $\text{Cu}_3\text{BiS}_3$ . The attractiveness of bismuth (Bi)-based compounds is due to the fact that Bi is an incredibly versatile metal found in ores as bismuthinite (bismuth sulfide) and bismite (bismuth oxide).<sup>6</sup> It is considered a “green metal” given that it is a harmless element compared to the toxic heavy metals in the periodic table,<sup>7</sup> and compared to In, Bi has greater world ore reserves (estimated 320 000 metric tons)<sup>8</sup> and lower cost (5.51 \$ per kg in 2019, ~68% cheaper than in 2013),<sup>9,10</sup> which indicates its distinct superiorities in abundance and price. In this context, the ternary compound of I<sub>3</sub>-V-VI<sub>3</sub>,  $\text{Cu}_3\text{BiS}_3$  (referred to as CBS in this work), also known as bismuth sulfosalt of copper or wittichenite, is an earth-abundant, nontoxic candidate, which remains largely unexplored. According to statistics from Scopus, 42 papers were published in the period 2019 to 2023 (assessed 02 June 2023), = using  $\text{Cu}_3\text{BiS}_3$  as the keyword, as shown in Fig. 1a.

CBS is a stable phase with a low melting point (~527 °C), demonstrating the feasibility for grain growth at a low temperature.<sup>11</sup> CBS exhibits remarkable optoelectronic properties such as a strong absorption coefficient ( $\alpha > 10^5 \text{ cm}^{-1}$ ) in the visible region (>1.6 eV), p-type conductivity, and an energy band gap value of ~1.10–1.86 eV<sup>12</sup> in the optimum range for different applications. Accordingly, bismuth-based chalcogenides have emerged as promising light absorbers with high optical absorption coefficient; namely  $\text{AgBiS}_2$  and  $\text{NaBiS}_2$ . Fig. 1b shows the optical absorption coefficient of established semiconducting light absorbing materials. The optical absorption of CBS shows a steep increase close to the band edge, similar to halide perovskites and  $\text{NaBiS}_2$ , and further increases  $>10^5 \text{ cm}^{-1}$  above 1.7 eV, which is presumed to be the direct band edge of CBS.<sup>13–15</sup> Moreover, it also indicates the peculiar high optical absorption behavior of Bi-based chalcogenide absorbers, which requires further attention.

The absorption coefficient together with the refractive index can be used to estimate the absorber layer thickness required for a given absorbance using the relation shown in eqn (1).

$$a = (1 + [4n^2\alpha d]^{-1})^{-1}, \text{ (in the presence of light scattering)} \quad (1)$$

where  $a$  is the absorbance,  $n$  is the refractive index and  $d$  is the absorber layer thickness. Thus, for 90% light absorption (and  $n = 2.6$  (at 800 nm)),<sup>16</sup> an absorber thickness of 40 nm is sufficient. However, in the absence of scattering, the layer thickness should be around 294 nm. The ultrathin layer relaxes the constraints posed by the charge transport, given that the carriers in the bulk have to cover a shorter distance to reach the surface before their recombination. Besides, the cost related to material consumption is also minimized. A detailed discussion on the optical absorption of CBS is presented in Section 2.2.

Although the above-mentioned properties make CBS a promising candidate for PV and PEC application, they are not necessarily sufficient to achieve a superior performance. Indeed, CBS-



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Fig. 1 (a) Number of publications on  $\text{Cu}_3\text{BiS}_3$ -related papers in the period 2019 to 2023 (source: Scopus; keyword:  $\text{Cu}_3\text{BiS}_3$ , assessed on 02 June 2023). (b) Optical absorption coefficient of Bi-based chalcogenide materials. References inserted from Wang *et al.*,<sup>13</sup> Huang *et al.*,<sup>14</sup> and Whittles *et al.*<sup>15</sup>

based solar cells have demonstrated a power conversion efficiency of only 1.7%, which is far from its Shockley–Queisser limit ( $\sim 32.9\%$  efficiency, 1.4 eV).<sup>17,18</sup> Alternatively, the impressive solar-to-hydrogen (STH) conversion efficiency of 2.04% has been reported for a CBS photocathode-based PEC device, demonstrating excellent prospects for photoelectrocatalysis.<sup>19</sup> Recently, Moon *et al.*<sup>20</sup> achieved a solar-to-hydrogen efficiency of 2.33% under unassisted conditions using a  $\text{Bi}_2\text{S}_3$ – $\text{Cu}_3\text{BiS}_3$  mixed phase interlayer. The results displayed a high onset potential of 0.9 V (*vs.* RHE) and three-times higher photocurrent density ( $7.8 \text{ mA cm}^{-2}$ , at 0 V *vs.* RHE) compared to the single-phase  $\text{Cu}_3\text{BiS}_3$ -photocathode. Besides, the photocathode showed great stability for over 20 h. Notably, there is no report on photoelectrochemical  $\text{CO}_2$  reduction to date. This highlights the need to understand the property-performance gap in CBS compounds. Accordingly, more research is needed to enhance the functionality and versatility of the CBS photoabsorber, which can be achieved by understanding its optoelectronic properties by combining theoretical and experimental results. In this review, we employ a concerted and integrated theoretical and experimental approach to present an overall assessment of CBS compounds and their underlying challenges.

This review consists of three sections, as follows: (i) discussion on the crystallographic phase and electronic properties, (ii) synthetic methods (thin films and nanostructure), and (iii) assessment of photovoltaic and photo(electro)catalysis devices. Finally, a summary and outlook are presented.

## 2 Crystallographic structure and electronic properties of CBS with particular interest in solar energy conversion

### 2.1 Crystal structure and phase boundaries

The ternary copper chalcogenide system of  $\text{Cu}_x\text{Bi}_y\text{S}_z$  (CBS) is comprised of various phases, where among them,  $\text{Cu}_3\text{BiS}_3$  is the

most interesting, but still lacks extensive investigation. Historically, wittichenite was discovered in the early 19th century, but only artificially synthesized in 1947.<sup>21</sup> The first fundamental work regarding the characterization of the crystal structure of wittichenite was published by Kocman and Nuffield.<sup>22</sup> CBS crystallizes in an orthorhombic structure with the crystal symmetry belonging to the  $P2_12_12_1$  space group, which differs from the traditional cubic and hexagonal structures of the most important semiconductors.<sup>23</sup> The details regarding its crystal structure parameters are shown in Table 1. All the lattice constants presented are experimental values and are in good agreement with other experimental and theoretical studies.<sup>24–27</sup>

Fig. 2a depicts the crystal structure of wittichenite. Although Bi is trivalent in the CBS structure, similar to In and Ga in CIGS, it presents a distorted lattice. CBS shows a distorted tetrahedral structure with  $\text{BiS}_3$  units. In this structure, the Bi atom is located at one vertex of the tetrahedron and is bonded with three sulfur atoms. The Cu atom also has three-fold coordination to sulfur, which forms a trigonal distorted structure near the planar  $\text{CuS}_3$  unit. This distortion is due to the presence of a post-transition-metal cation ( $\text{Bi}^{3+}$ ) with an  $ns^2$  electronic configuration.<sup>28</sup> Considering the Gillespie–Nyholm symmetry theory, also known as the valence-shell electron-pair repulsion (VSEPR) rule,  $s^2$  induces the distortion due to the presence of Bi

Table 1 Crystal structure parameters of CBS compound. Kocman and Nuffield<sup>22</sup>

Material		CBS
Parameter		
Crystal structure		Orthorhombic
Space group		$P2_12_12_1$
Lattice constant	$a$ (Å)	7.723
	$b$ (Å)	10.395
	$c$ (Å)	6.716
	$\alpha = \beta = \gamma = 90^\circ$	





Fig. 2 (a) Orthorhombic crystal structure of wittichenite (CBS), trigonal planar  $\text{CuS}_3$  unit showing distances between Cu (red) and S (yellow) atoms and distorted square pyramidal  $\text{BiS}_5$  unit showing distances between Bi (grey) and S (yellow) atoms. Adapted with permission from ref. 31 Copyright 2012 Elsevier B.V. (b) Triangle phase diagram of the Cu–Bi–S system at 300 °C. Adapted from the thesis of Chamorro *et al.*<sup>32</sup> (c) Binary phase diagram of the  $\text{Cu}_2\text{S}$ – $\text{Bi}_2\text{S}_3$  system. wit = CBS, h-wit = high form CBS, emp = emplectite, cpb =  $\text{Cu}_4\text{Bi}_7\text{S}_{12}$ , X =  $\text{Cu}_9\text{BiS}_6$ , D =  $\text{Cu}_3\text{Bi}_5\text{S}_9$ , E =  $\text{CuBi}_3\text{S}_5$  cc =  $\text{Cu}_2\text{S}$ , and bm =  $\text{Bi}_2\text{S}_3$ .<sup>33</sup>

$6s^2$  lone pairs, which causes the distortion of the crystal structure based on the interaction of empty states with the antibonding orbitals, resulting in an asymmetric electron density.

The pseudo binary phase diagram of the  $\text{Cu}_2\text{S}$ – $\text{Bi}_2\text{S}_3$  system and ternary phase diagram of the Cu–Bi–S system are depicted in Fig. 2b and c, respectively. To understand the phase transformation of the Cu–Bi–S system, they are discussed in terms of composition and temperature. In the CBS system, stoichiometric changes result in synthetic compounds with different phases. Among them, the  $\text{Cu}_2\text{S}$  (chalcocite),  $\text{Cu}_3\text{BiS}_3$  (wittichenite),  $\text{Cu}_4\text{Bi}_7\text{S}_{12}$  (cuprobismutite),  $\text{CuBiS}_2$  (emplectite) and  $\text{Bi}_2\text{S}_3$  (bismuthinite) compounds are found in nature, while the  $\text{Cu}_9\text{BiS}_6$ ,  $\text{Cu}_3\text{Bi}_5\text{S}_9$  and  $\text{CuBi}_3\text{S}_5$  compounds have not been found to date. The former two present thermal stability in the range of 390–650 °C and 442–620 °C, respectively (Sugaki & Shima, 1972). Regarding the latter,  $\text{CuBi}_3\text{S}_5$ , to date there is no theoretical report and only one experimental report related to the synthesis of this phase<sup>29</sup> showing n-type behavior with a carrier concentration and band gap of  $3.75 \times 10^{-17} \text{ cm}^{-3}$  and 0.77 eV, respectively.  $\text{CuBi}_3\text{S}_5$  exists as a stable and pure phase in the temperature range of 280–649 °C and it is expected to gain more attention in the future. CBS is thermally stable in the range of 350–527 °C, while  $\text{CuBiS}_2$  presents thermal stability below 317 °C, while at higher temperatures, it decomposes to wittichenite, covellite ( $\text{CuS}$ ) and cuprobismutite, as shown in the reaction in

eqn (2), demonstrating the challenge to obtain this phase.<sup>30</sup> Consequently, while  $\text{CuBiS}_2$  is only stable at  $T < 319$  °C, cuprobismutite is synthesized at higher temperatures.



Although the compounds present in the Cu–Bi–S system are stable in a very narrow range, they exhibit interesting optoelectronic properties and show great potential to be applied in various energy systems. Therefore, further studies are needed to controllably synthesize and characterize these distinct observed phases. Regarding this, more information will be described in the following sections.

## 2.2 Electronic structure and optoelectronic properties: character of band edges, optical absorption, effective masses, and bandgaps

The electronic structure of CBS is governed by the fundamental differences in the chemical bonding between Cu, Bi and S in  $\text{Cu}_x\text{Bi}_y\text{S}_z$  compounds. The bonding and antibonding hybrid states form the conduction band minimum (CBM) and the valence band maximum (VBM), respectively. The antibonding character is formed by the interaction of the active  $ns^2$  lone pairs. This electronic structure is unusual compared to other semiconductor materials (bonding VBM and antibonding



CBM). For this type of semiconductor, the electronic energy levels of an element are in resonance with the band edges. Accordingly, the removal or insertion of an element in the crystal structure leads to the formation of shallow states close to the band edges instead of deep defect states in the gap (*i.e.*, trap centers, which restrict the charge transport). For this reason, the semiconductors containing antibonding states at the top of the valence band are also classified as defect tolerant.<sup>34</sup> This implies that defect-tolerant semiconductors are insensitive to the presence of crystallographic defects (*i.e.*, vacancies, interstitials) and have the capability to maintain their optoelectronic properties. Fig. 3a shows the orbital contributions from the individual atoms forming the electronic density of states for CuS, Bi<sub>2</sub>S<sub>3</sub> and Cu<sub>3</sub>BiS<sub>3</sub> compounds. In the case of CuS, the valence band maximum (VBM) and the conduction band minimum (CBM) consist of Cu d and S p states and Cu s with S p states, respectively.<sup>35</sup> While for Bi<sub>2</sub>S<sub>3</sub>, the VBM and CBM are derived from Bi s with S p states and Bi p with S p states, respectively.<sup>36,37</sup> Interestingly, in CBS, which contains contributions from both Cu and Bi, its VBM results from the coupling of the Cu d and S p states and its CBM is characterized by Bi p and S p states.<sup>15,25</sup> According to the density of states (DOS) of CBS, as shown in Fig. 3b, it is possible to observe that the predominant contribution to the valence band comes from the Cu d orbital with some S p contribution. In CBS, the Bi s orbitals are fully occupied; hence, the empty Bi p orbitals are free to hybridize with the S p orbitals, forming the conduction band.

High optical absorption is a characteristic trait of CBS. The origin of its high optical absorption is related to the peculiar nature of its electronic bands. The optical absorption is

governed by the transition matrix  $\langle \text{VB} | \hat{H} | \text{CB} \rangle$  and joint density of states (JDOS) at a given photon energy, which are related by Fermi's Golden rule, as shown in eqn (3):<sup>38</sup>

$$\alpha \propto \frac{4\pi^2}{h} \left| \langle \text{VB} | \hat{H} | \text{CB} \rangle \right|^2 \cdot \text{JDOS} \quad (3)$$

where  $h$  is Planck's constant and VB and CB represent the eigen states of the valence and conduction band, respectively. The transition matrix depends on the coupling of the initial and final transition state and is higher for a direct-allowed transition, *i.e.*, for direct bandgap semiconductors. The JDOS is dictated by the density of states (DOS), and hence the atomic orbitals responsible for the formation of the CB and VB. To understand the nature of the optical absorption of CBS, we compared the DOS of CBS to that of conventional highly absorbing PV materials, *i.e.*, MAPbI<sub>3</sub>, CuInSe<sub>2</sub> and CdTe.

Perovskites show excellent absorption characteristics owing to their direct band gap and  $p \rightarrow p$  transition. The less dispersive  $p$  character of the perovskite CB, unlike the  $s$ -like more dispersive CB of CIS and CdTe, favors a high absorption coefficient.<sup>39</sup> Thus, the  $p$ - $p$  transition is highly desirable to achieve efficient optical absorption, which is possible for compounds having cations with a lone pair of  $s$  electrons,<sup>40</sup> *e.g.*, the Pb<sup>2+</sup> lone pair in perovskite. Interestingly, in the case of CBS, its optical transition is dominated by the  $p$ - $p$  transition from the S  $p$  to Bi  $p$  states, which may be responsible for its high absorption. The presence of the Bi 6s<sup>2</sup> lone pair promotes the  $p$  character of the VB similar to perovskite. Fig. 4 shows the density of states and corresponding schematic optical absorption from several materials including MAPbI<sub>3</sub>, CuInSe<sub>2</sub>, CdTe,



Fig. 3 (a) Illustration of the main orbitals and their occupancy in the valence and conduction band for CuS, Bi<sub>2</sub>S<sub>3</sub> and CBS. Blue and red colors correspond to occupied orbitals in VB and CB, respectively. (b) Density of states (DOS) curves for CBS material. Adapted with permission from ref. 15 Copyright 2019, the American Chemical Society. The top part shows the total and partial electronic DoS curves, where the black curve with gray shading is the total summed DoS. The middle part presents the zoomed region showing the underlying orbitals. The bottom part is the configuration energy (CE) for the valence orbitals, containing the bonding hybridizations.





Fig. 4 Density of states (DOS) and corresponding schematic of the optical absorption for (a) perovskite (MAPbI<sub>3</sub>). Reproduced with permission from ref. 61 Copyright 2019, The Royal Society of Chemistry. (b) Chalcopyrite CuInSe<sub>2</sub>. Reproduced with permission from ref. 62 Copyright 2010, John Wiley & Sons, Ltd. (c) CdTe. Reproduced under the terms of the CC-BY-4.0 license from ref. 63 Copyright 2021, the American Chemical Society. (d) Cu<sub>3</sub>BiS<sub>3</sub>.<sup>15</sup> The dashed line indicates the valence band (zero energy). The grey arrow indicates the strength of the optical absorption. (e) Calculated electronic band structure for CBS using HSE06 + SOC DFT, VB and CB, as depicted in blue and orange, respectively. VBM is set as a reference in both cases. Figures (b), (c), and (d) reproduced with permission from ref. 15 Copyright 2019, the American Chemical Society. (f) Temperature-dependent optical absorption spectra of CBS (at 80 and 280 K, dark red and green lines, respectively). Inserted figure corresponds to DFT DoS at the top of the VB.<sup>15</sup> (g) Bandgap energy values reported for CBS in the literature.<sup>15,26,45,53,57,58,60,64–67</sup>

and Cu<sub>3</sub>BiS<sub>3</sub>. As shown in Fig. 4d, it is clear that the VB of CBS resembles that of CIS, while its CB shows similar character to the perovskite CB. This makes it ideal for efficient optical transition, explaining the excellent absorption coefficient of CBS. The optical absorption of CBS originates from the dipole transition between the Cu-d/S-p orbitals (with a small contribution from the Bi s state due to the presence of the lone pair mechanism)<sup>15</sup> of the VBM states and Bi-p/S-p of the CBM states. It is worth noting that a strong spin-orbit coupling (SOC) is expected in CBS due to the presence of the heavy Bi. However, the presence of SOC and its impact on the optical absorption of CBS are unknown.

Fig. 4e shows the band structure of CBS, which was calculated using HSE06 + spin-orbit coupling (SOC). Both the VBM and CBM are located at the  $\Gamma$  point of the Brillouin zone for

CBS. For most copper-based chalcogenide materials, such as CBS, s-p or p-d hybridization is present. The effective mass ( $m_e^*$ ) carriers are related to the curvature of the individual bands. Thus, strongly curved wide bands (VBM and CBM more dispersive) suggest a lower effective mass, and hence higher carrier mobility.<sup>41</sup> Also, the overlapping orbitals on neighboring atoms promotes larger delocalization, which leads to extended states/band, favoring the transport of free charge carriers and resulting in higher carrier mobility (specially desirable for photoabsorbers with an indirect bandgap).

The carrier mobility ( $\mu$ ) of a given material is related to the width and intrinsic dispersion of its conduction and valence bands. Eqn (4) describes the relationship between  $\mu$  and effective mass ( $m^*$ ). In the case of n-type (or p-type) conductivity,  $m^*$  is the electron effective mass  $m_e^*$  (or hole effective mass ( $m_h^*$ )).



Both the carrier concentrations and mobilities can be determined from electrical Hall and thermoelectric measurements. Regarding ( $m_e^*$ ) there are only a few reports on CBS in the literature, and it was not discussed in detail, where  $m_h^*$  was estimated to be 0.35 by Hall-voltage and thermoelectric Seebeck-coefficient measurements.<sup>42,43</sup> Besides, the effective mass can play a crucial role in the net photo-current, and hence the cell efficiency for different applications (PV or PEC).<sup>44</sup>

$$m^* = \frac{\hbar^2}{\partial^2 E / \partial k^2}, \quad \mu = \frac{q\tau}{m^*} \quad (4)$$

Table 2 shows the carrier concentration and carrier mobility values for CBS materials prepared by different synthesis methods. Most of the reports indicate p-type conductivity, similar carrier concentration ( $10^6$ – $10^7$  cm<sup>-3</sup>) and low carrier mobility (0.4–4.0 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>), except for the study by Li *et al.*,<sup>45</sup> Fazal *et al.*,<sup>46</sup> and Lee *et al.*<sup>47</sup> A low carrier mobility indicates the active role of electronic defects in the material. Defects are known to act as recombination centers and affect the carrier transport in different ways through the scattering process within the grain interior and/or at the grain boundaries (GB).<sup>48</sup> It is known that the grain boundaries can be regulated by changing the synthesis method and minimizing the transport barriers, limiting the conductivity of the film.<sup>42</sup> The fact that low mobility values are obtained regardless of the synthesis route (solution or vacuum processed) indicates the intrinsic nature of the defects in CBS. Although most of the studies report that the presence of defects has a negative impact on the electronic properties of materials, defect engineering can be an optimal approach for some specific applications, specifically photoelectrocatalysis.<sup>49–51</sup>

The fundamental bandgap energy ( $E_g$ ) is established as the energetic difference between the VBM and CBM. Upon illumination, the photogenerated charge carriers at the CBM relax to the VBM through radiative transitions in direct bandgap

semiconductors. In contrast, for indirect bandgap semiconductors, the absorbed light is emitted as a nonradiative transition *via* phonons, thus losing its energy. The bandgap of CBS is under debate in the literature due to the ill-defined nature of its band edges and possibly the presence of defect states, emphasizing the necessity for further investigations. In the literature, there is a spread in the experimental value of the reported data (1.0–1.8 eV), showing a direct bandgap for CBS. In contrast, theoretical studies predict an indirect band gap for CBS. Fig. 4g shows both the experimental and theoretical band gap energy values reported for CBS. Notably, a significant difference is observed in the value of the bandgap, which may be due to the synthesis method (hence defect) adopted. Most of the studies rely on optical absorption spectroscopy to determine the bandgap. The variation in the bandgap value may originate from the measurement itself given that it is difficult to unambiguously ascertain the bandgap for indirect semiconductors from their absorption onset alone. In the case of a small energy difference between the direct and indirect band edge and/or the presence of defect states near the band edge, the bandgap appears to be pseudo-direct. Additionally, this variation can be due to the crystallinity of the material or the crystal orientation. For instance, a single crystal exhibits a lower band gap compared to thin films.<sup>52</sup>

For materials targeted for application in solar conversion, *i.e.*, as semiconductors or photovoltaics, strong optical absorption is crucial. Cu<sub>x</sub>Bi<sub>y</sub>S<sub>z</sub> shows an ideal photo absorption in the visible wavelength region with an optimal absorption coefficient ( $>10^5$  cm<sup>-1</sup>) compared to other Cu-S-based materials, such as CIGS and CZTS. However, according to the band structure (Fig. 4e), from the  $\Gamma$  and X points of the Brillouin zone, it is observed that the low density of states (DoS) reduces the absorption extent of the CBS material. The fundamental bandgap energy for CBS is 1.372 eV and indirect;<sup>15</sup> however, its absorption is dominated by photon absorption with higher energy direct transitions of  $>1.4$  eV (V2  $\rightarrow$  C, inset in Fig. 4e)

Table 2 Reported carrier mobility, concentration, and conductivity type of CBS materials. \*DMSO = dimethyl sulfoxide RT = room temperature

Ref.	Chemistry	Composition/synthesis method	Carrier type	Carrier concentration (cm <sup>-3</sup> )	Carrier mobility (cm <sup>2</sup> V <sup>-1</sup> s <sup>-1</sup> )
27	Solution processing	CBS/CuO and Bi <sub>2</sub> S <sub>3</sub> dissolved in thiol-amine and dried at 400 °C	p-type (Hall measurements)	$4.2 \times 10^{17}$	0.4
42	Vapor phase synthesis	CBS/Co-evaporation of Cu, Bi, and S onto a substrate at 300 °C	p-type (Hall measurements)	$2.00 \times 10^{16}$	4.00
45	Solution processing	CBS/DMSO*-based solution processed, short-time drying of spin-coated precursor film	p-type (Hall measurements)	$5.65 \times 10^{16}$	52.83
46	Chemical bath deposition	CBS/deposited for 12 h	p-type (four-probe method)	$6.07 \times 10^{16}$	34.90
57	Solution processing	CBS/Spin-coating Cu and Bi oxides followed by sulfurization	p-type (Hall measurements)	$5.10 \times 10^{16}$	3.73
58	Vapor phase synthesis	CBS/co-evaporation precursors followed by annealing process	p-type (Hall measurements)	$2.80 \times 10^{16}$	1.50
59	Vapor phase synthesis	CBS/one-stage co-evaporation process	p-type (Hall measurements)	$7.48 \times 10^{16}$	3.95
60	Solution processing	CBS/spin coating precursor solution, followed by annealing on a hotplate	p-type (Hall measurements)	$6.39 \times 10^{17}$	3.09
47	Solution processing	CBS/chemical bath deposition technique	p-type (Hall measurements)	$2.57 \times 10^{16}$	71.20



due to the large contribution given by the presence of density of states at the valence band maximum.

Besides, to measure the bandgap, very thin films of CBS are currently being used in several experimental studies. In the case of indirect bandgap materials, a large thickness is required to realize the complete absorption of the incident light, which can also be another reason for the huge disparity among the reported results.

Unfortunately, the methods used to measure the band gap energy are not suitable to determine its nature given that it gets non-trivial due to the presence of many peaks close to the VB. Thus, other techniques must be explored to reliably determine the bandgap such as photoluminescence (PL), photothermal deflection spectroscopy (PDS), and quantum efficiency measurements. There are very few studies reporting a measurable PL for CBS,<sup>53–56</sup> given that the PL quantum yield of the material is expected to be quite low due to its indirect bandgap or presence of deep defects, which is highly likely in this compound. Considering the above-mentioned arguments, the defect tolerance of CBS has to be scrutinized and more studies are required to understand the nature of its intrinsic defects.

### 2.3 Energy band alignment

The other essential criteria for the optoelectronic properties of materials towards device applications are the energetic band edges offsets and Fermi level ( $E_F$ ) (*i.e.*, ionization potential, work function, and electron affinity). In a device, its interface/junction determines the transport of charge carriers (electrons and holes). Therefore, it is important to have both materials energetically aligned for efficient charge transfer. For most CBS studies reported in the literature, CdS (n-type) is commonly used as an n-type buffer layer to form a heterojunction, given that it shows great results for CIGS and CZTS. However, as shown in Fig. 5a, the energy band positions of CdS are not properly aligned with respect to CBS. Thus, it can be inferred

that the low solar cell efficiencies observed in CBS devices are due to the lack of understanding of the above-mentioned issue and the improper choice of n-type layer. Similarly, hole-selective layers should also be investigated. Fig. 5b illustrates an ideal band alignment scheme for the CBS material.

## 3 Synthesis methods

### 3.1 Nanocrystals

The synthesis of Cu–Bi–S nanocrystals is challenging due to the high level of difficulty in controlling their growth and their features, such as size, shape, and configuration. In colloidal synthesis, specific parameters (*i.e.*, nucleation and growth kinetics) lead to the formation of uniform nanocrystals. This is done by optimizing the material preparation conditions and the type of capping ligands used.

The first report on the synthesis of CBS nanocrystals by Yan *et al.*<sup>68</sup> employed a hot-injection method based on colloidal synthesis using copper(II) acetylacetonate, bismuth(III) nitrate pentahydrate [atomic ratio Cu : Bi = 3] and oleylamine (OLA) ligand. High-quality CBS nanocrystals showed good crystallinity and an average particle size in the range of 8–11 nm. Bi<sub>2</sub>S<sub>3</sub> and CuS phases were identified as key intermediate steps to form CBS NCs and the shape of the nanocrystal could be tuned by varying the precursor concentration and reaction temperature. Besides, the optical bandgap of the CBS nanocrystal (1.56 eV) was found to be larger than that of the CBS bulk (1.40 eV) due to the quantum confinement effect. Also, the CBS nanocrystals showed a good photo response, signifying their potential to be used as a thin film for solar energy conversion.

In another study, Chakraborty *et al.*<sup>65</sup> obtained pure CBS semiconducting nanoflowers (450–500 nm) and nanospheres (6–9 nm) using a facile hot injection method. The precursors, reaction temperature and ligand type play a key role in dictating the morphology and the proposed reaction mechanism is shown in Fig. 6a. Specifically, 1-DDT (dodecanethiol) is the

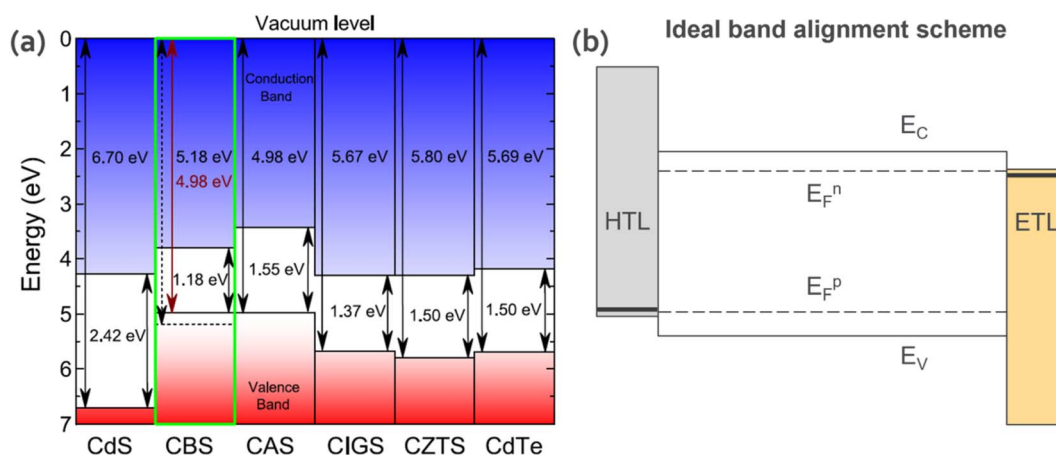


Fig. 5 (a) Vacuum-aligned band diagram between CBS and other common absorber layers. Reproduced with permission from ref. 15 Copyright 2019, the American Chemical Society. (b) Ideal band alignment scheme for CBS. HTL = hole transport layer; ETL = electron transport layer;  $E_F^p$  = Fermi energy level for p-type material;  $E_F^n$  = Fermi energy level for n-type material;  $E_C$  = conduction band energy level; and  $E_V$  = valence band energy level.





reported. Mesa and Gordillo<sup>46</sup> prepared CBS thin films *via* a two-step co-evaporation process. The first step involved the formation of a  $\text{Bi}_x\text{S}_y$  layer by the co-evaporation of Bi and S elements (substrate temperature of 300 °C). Subsequently, Cu was evaporated in the second step to convert  $\text{Bi}_x\text{S}_y$  to CBS (at around 300 °C). The results showed a high optical absorption ( $>10^4 \text{ cm}^{-1}$ ) and bandgap of 1.41 eV. In the study by Hu *et al.*,<sup>72</sup> CBS thin films were prepared *via* low-temperature (100–150 °C) hydrothermal (HD) synthesis using copper and bismuth-based compounds and thiourea (Tu). The material consisted of whisker-like particles. The best-quality CBS was obtained at the growth temperature of 120 °C for <8 h. Better homogeneity was realized by liquid-solution diffusion due to the formation of  $[\text{Cu}(\text{Tu})_x]^+$  and  $[\text{Bi}(\text{Tu})_y]^{3+}$  complexes, which reduced the dissociation of ions, and thus inhibited the formation of secondary phases ( $\text{Cu}_2\text{S}$  and  $\text{Bi}_2\text{S}_3$ ) in the solution.

In another study by Koskela *et al.*,<sup>27</sup> thiol-amine (1,2-ethanedithiol, EDT) was used to dissolve CuO and  $\text{Bi}_2\text{S}_3$  precursors to form a colloidal ink (yellowish) for the preparation of CBS. The ink was solution-processed on a fluorine-doped tin oxide (FTO) substrate and annealed at 400 °C under  $\text{N}_2$  to form a pure-phase CBS film (Fig. 7a). The as-produced film exhibited a high absorption coefficient ( $\alpha = 5 \times 10^4 \text{ cm}^{-1}$ ) and optical absorption bandgap of 1.47 eV, as shown in the inverse logarithmic derivative (ILD) plot in Fig. 7b. Besides, the resistivity of CBS follows typical semiconducting characteristics with temperature, as shown in Fig. 7c.

Chen *et al.*<sup>73</sup> prepared an Au- $\text{In}_2\text{S}_3$ /CBS thin film heterostructure using a combination of simple spray pyrolysis deposition (SPD) on molybdenum-coated glass (Mo-SLG) and chemical bath deposition. Firstly, a CBS precursor solution was prepared using a simple solution of  $\text{Cu}(\text{NO}_3)_2 \cdot 3\text{H}_2\text{O}$  (99.9%) and  $\text{Bi}(\text{NO}_3)_3 \cdot 5\text{H}_2\text{O}$  (99.9%) in  $\text{HNO}_3$ . Then, the precursor was sprayed (controlled at 7 min) on an Mo/SLG substrate (reaction controlled at 380 °C). After this step, the precursor film was sulfurized in a tube furnace at 550 °C for 1 h. The  $\text{In}_2\text{S}_3$  layer was prepared *via* the chemical bath deposition method, using  $\text{InCl}_3$ , thioacetamide and acetic acid in water. The bath was heated at

65 °C for 30 min, followed by annealing at 100 °C for 20 min to obtain  $\text{In}_2\text{S}_3$ /CBS. Mott–Schottky analysis confirmed the p-type semiconductor characteristics. The band gap was calculated to be 1.58 eV.

Fazal *et al.*<sup>46</sup> reported the preparation of a CBS thin film by the complexation of  $\text{Cu}^{2+}$  and  $\text{Bi}^{3+}$  ions with EDTA to form  $[\text{Cu}(\text{EDTA})]^{2-}$  and  $[\text{Bi}(\text{EDTA})]^{-1}$ . Then, the  $\text{S}^{2-}$  ions released from thioacetamide reacted with  $\text{Bi}^{3+}$ , accompanied by a change in the color of the solution to grayish brown, indicating the successful synthesis of CBS by ion consumption and integration into the  $\text{Bi}_2\text{S}_3$  lattice. The slower kinetics posed a major limitation in obtaining a thicker film. This issue was solved by increasing the deposition time from 6 to 12 h. The proposed reaction mechanism is shown in eqn (5) and (6).



The spray pyrolysis technique is very attractive to produce thin films given that it is a versatile method that enables facile control of the thickness. CBS was successfully deposited on a glass substrate using  $\text{BiCl}_3$ ,  $\text{CuCl}_2$  and thiourea precursor solutions<sup>74</sup> with varying substrate temperatures in the range of 250–400 °C. The optical bandgap was estimated to be 1.72 and 1.65 eV for the films synthesized at 250 °C to 300 °C, and 350 °C to 400 °C, respectively. All the thin-films showed p-type conductivity. In another approach, spray pyrolysis was performed using bismuth tris(4-methylbenzodithioate)  $[\text{Bi}(\text{4-MBDT})_3]$  mixed with copper(i) acetate and 1,2-ethanedithiol precursors.<sup>75</sup> In this work, CBS was synthesized on a glass substrate maintained at  $150 \text{ °C} \leq T_{\text{dep}} \leq 200 \text{ °C}$  (at higher temperatures, unknown phases appeared, and at lower temperatures, the complete conversion of the precursor to CBS was not obtained). This technique allowed the formation of films with an easily controlled and reproducible thickness. The optimum condition to spray the precursor solution was



Fig. 7 (a) X-ray diffraction pattern for pure-phase  $\text{Cu}_3\text{BiS}_3$ . (b) Inverse logarithmic derivative plot from absorption spectra derived from transmittance. The dashed black line corresponds to the linear regression used to estimate the optical band gap. (c) Resistivity *versus* temperature for CBS synthesized by solution processing. (Inserted) picture of semiconductor ink containing CuO and  $\text{Bi}_2\text{S}_3$  dissolved and  $\text{Cu}_3\text{BiS}_3$  thin film. Reproduced with permission from ref. 27 Copyright 2021, the American Chemical Society.



identified to be 175–200 °C with 12 deposition cycles. The authors observed a broad light absorption up to 1000 nm. However, it was not possible to determine the bandgap of the material due to the unclear absorption onset energy. Attempts to fabricate a functioning photovoltaic device remain unsuccessful. The authors highlighted the necessity to identify suitable hole-transporting materials for CBS.

The choice of synthesis method depends on the desired properties of the material being synthesized. Table 3 shows the correlation between some targeted properties (particle, compositional, thickness control, and uniformity) and the synthesis method that could be more suitable to achieve it.

## 4 Applications

### 4.1 Thin-film photovoltaic device

Thin-film photovoltaics (PV) is excellent technology with a thinner light absorber layer, which makes it lighter, flexible and reduces the material consumption. The potential of CBS material in the field of solar PV was only recognized in 1997.<sup>76</sup> There are promising findings on the structural, morphology, synthesis, and optoelectronic properties of the CBS material; however, its extension to the device level is still sporadic, with limited architectures explored in the literature.

Yin and Jia *et al.*<sup>77</sup> successfully fabricated CBS nanosheet films (30 nm thick) on TiO<sub>2</sub> nanorod arrays on FTO substrate *via* the solvothermal route, forming a p–n heterojunction. The photoelectric properties of the CBS/TiO<sub>2</sub> composite film exhibit a fill factor of 60.8%,  $J_{sc}$  of 4.449 mA cm<sup>-2</sup>,  $V_{oc}$  of 442 mV and energy conversion efficiency of 1.281%, which is 2.5-times higher than the TiO<sub>2</sub> standalone device. The reported energy conversion efficiency of the CBS/TiO<sub>2</sub> thin-film solar cell is very promising.

Li *et al.*<sup>45</sup> reported a facile and low-toxic solution route for the deposition of CBS films without post sulfurization. They also proposed a prototype glass/Mo/CBS/CdS/ZnO/ITO/Al solar cell. This device showed a short-circuit current density ( $J_{sc}$ ) of 2.37 mA cm<sup>-2</sup>, open-circuit voltage ( $V_{oc}$ ) of 0.19 V, fill factor (FF) of 37.56% and conversion efficiency of 0.17% (Fig. 8a and b). The authors mentioned that the low  $J_{sc}$  may be due to the poor collection of photogenerated charges carriers, which is probably

limited by the improper junction between CBS and CdS (more information can be found in the report by Whittles *et al.*<sup>15</sup>).

Yang *et al.*<sup>58</sup> studied the effect of the atomic ratio of the Cu/Bi precursors (1.0–3.0) and the annealing temperature in the range of 350 °C to 475 °C for CBS thin films prepared *via* the co-evaporation of bismuth and CuS materials in a vacuum system, followed by an annealing process. In this work, the samples annealed at 450 °C showed better results. The Cu/Bi atomic ratio of 1.5 showed improved mobility and carrier density (1.5 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> and 2.8 × 10<sup>16</sup> cm<sup>-3</sup>, respectively). The CBS thin film showed a high absorption coefficient of 1.5 × 10<sup>5</sup> cm<sup>-1</sup> at 600 nm. The device structure of Glass/Mo/CBS/CdS/ZnO/ZnO:Al/Ni/Al yielded a  $V_{oc}$  of 0.32 V and FF of 18.2 mA cm<sup>-2</sup> with 1.7% power conversion efficiency, which is the highest performance reported to date.

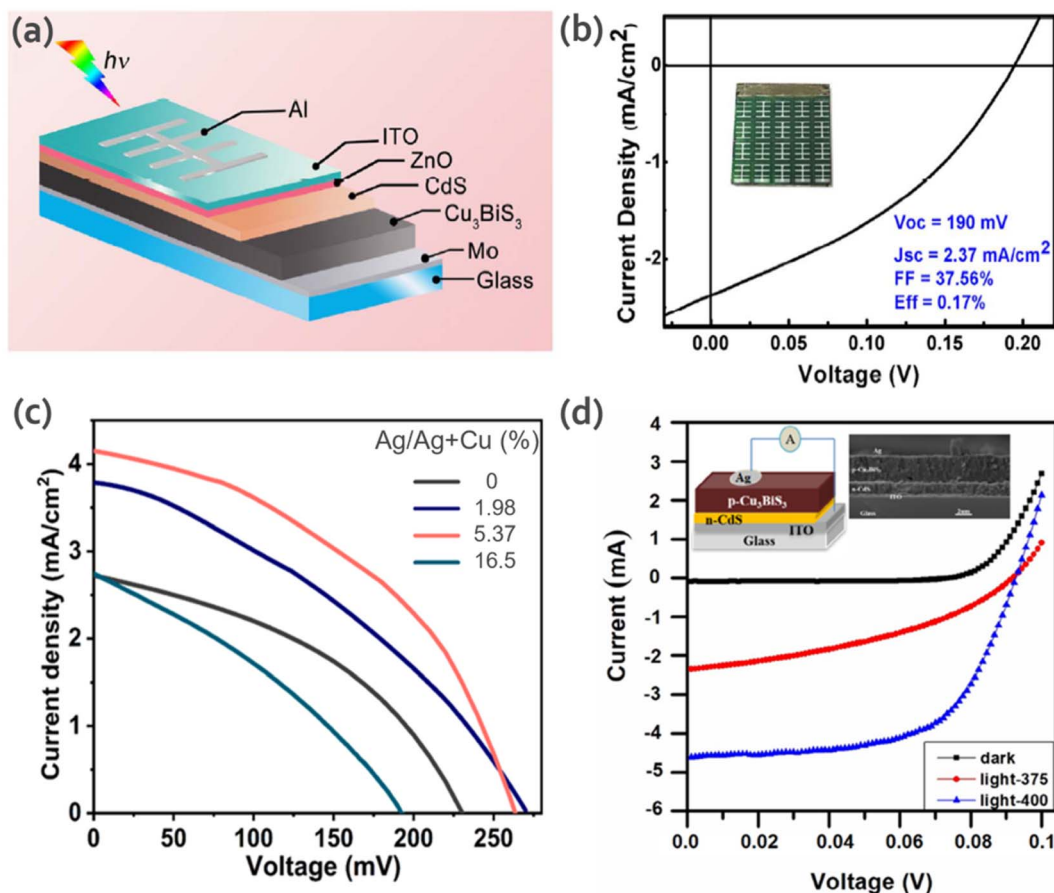
Zhao and collaborators<sup>60</sup> studied the impact of Ag doping on CBS, aiming to optimize the material properties for photovoltaic applications. The CBS-Ag absorber layer was prepared by solution processing. To control the Ag concentration in individual precursor solutions, the mole ratio of Ag/(Ag + Cu) was varied to be 0%, 1.98%, 5.37%, 16.5% and 35.8%, which was denoted as A0, A1, A2, A3 and A4, respectively. The devices were fabricated based on the structure of SLG/Mo/CBS/CdS/ZnO/ITO/Al. Fig. 8c shows the illuminated current density–voltage ( $J$ – $V$ ) curves for the A0–A3 samples. The sample with 5.37% Ag/(Ag + Cu) showed a  $V_{oc}$  of 263 mV,  $J_{sc}$  of 4.15 mA cm<sup>-2</sup> and FF of 44.1%. A power energy conversion of 0.48% was obtained. Compared to A0 (no doping), an improvement of ~85% by 5.37% Ag doping was clearly evident.

In another study by Hussain *et al.*,<sup>59</sup> CBS thin films were deposited on a soda-lime glass substrate *via* a one-stage co-evaporation process using Cu<sub>2</sub>S and Bi<sub>2</sub>S<sub>3</sub> sources at different deposition temperatures (room temperature and in the range of 275–400 °C). The crystalline quality of the films increased with an increase in the deposition temperature. The CBS films were applied in solar cell structures (glass/ITO/n-CdS/p-CBS/Al). Fig. 8d shows the performance of the CBS thin film solar cell devices deposited at 375 °C and 400 °C. The results showed photovoltaic behavior and an increase in the photocurrent from 1.0 to 3.0 mA under illumination at 0.09 V. To date, CBS has shown a relatively poor solar cell performance with a maximum of 1.7%, which is way far from the Shockley–Queisser limit

**Table 3** Correlation between targeted properties (composition, phase, particle size, and scalability) and their corresponding synthesis method. \*NCs = Nanocrystals and TF = thin films

Synthesis Method	Composition and phase Type control	Particle size control	Scalability for	Suitable	Comment
Hot injection	NCs ✓	✓	✗	PEC	Limited yield
Chemical bath deposition	TF ✗	✗	✓	PV, PEC	Slower growth kinetics, crystallinity control is difficult
Hydrothermal synthesis	NCs ✓	✓	✗	PEC	Considerations required for substrate
Spray pyrolysis	TF ✗	✗	✓	PV, PEC	Less control overgrowth kinetics
Co-evaporation	TF ✓	✗	✓	PV, PEC	High throughput is a challenge





**Fig. 8** (a) Schematic diagram of the CBS solar cell device. (b)  $J$ - $V$  curve of the CBS solar cell under AM 1.5G simulated irradiation. The inset shows an image of the fabricated device. Reproduced with permission from ref. 45 Copyright 2017, Elsevier B.V. (c)  $J$ - $V$  curves for different Ag/(Ag + Cu) ratio values. Reproduced with permission from ref. 60 Copyright 2021, Elsevier Ltd (d) Typical  $I$ - $V$  curves for the CBS thin film solar cell device at 375 °C and 400 °C obtained under illumination of 100 mW cm<sup>-2</sup>. The inserted figure shows a schematic diagram and cross-sectional SEM image of the device. Reproduced with permission from ref. 59 Copyright 2019, Elsevier B.V.

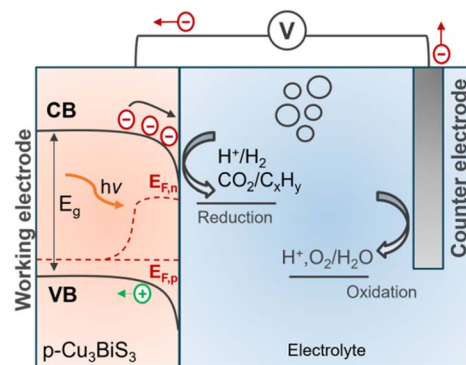
(~33.7% efficiency, 1.3 eV).<sup>17</sup> The theoretical device simulation study by Mesa *et al.*<sup>78</sup> showed a  $V_{oc}$  of 0.712 V,  $J_{sc}$  of 36.25 mA cm<sup>-2</sup>, FF of 79.54% and efficiency of 19.86%, which are quite far from the experimental studies.

Thus, the performance-limiting factors of CBS devices need to be identified. The role of intrinsic defects in carrier recombination and energy band alignment requires attention to realize the potential of CBS-based devices we believe that this material continues to generate interest and more exploration is required in this front.

## 4.2 Photoelectrocatalysis

Photoelectrocatalysis (PEC) has been used in many different applications such as disinfection, degradation of contaminant compounds, CO<sub>2</sub> reduction, and water splitting. Typically, a p-type semiconductor is used as a photocathode to drive photoelectrochemical reactions. Briefly, when a p-type semiconductor is immersed in an electrolyte, charge transfer occurs due to the chemical potential equilibrium (Fermi level,  $E_F$ ) of the semiconductor and the redox potential of the electrolyte. This process results in an electric field, manifested by band bending

(downward for p-type) of the semiconductor, as shown in the scheme presented in Fig. 9. Upon illumination, the excited electrons are driven towards the interface due to the gradient of the quasi-Fermi levels, where they can reduce the different



**Fig. 9** Schematic of the PEC setup showing the band bending and quasi-Fermi levels for electrons ( $E_{F,n}$ ) and holes ( $E_{F,p}$ ) for a typical p-type semiconductor.



redox species at the surface, such as water or  $\text{CO}_2$ , to hydrogen or hydrocarbons. Simultaneously, the other half reaction is balanced by the oxidation reaction at the anode (counter electrode). The photovoltage required to drive the electrochemical reaction can be achieved either by a single material with the optimum bandgap or tandem systems (*i.e.*, heterojunction and Z-scheme). Readers are encouraged to refer to Walter *et al.*<sup>79</sup> for a focused discussion on photoelectrochemistry.

Cu–Bi–S compounds have not been well investigated as photocatalysts although they present outstanding optoelectronic

properties. There are few reports on water splitting and/or hydrogen evolution reaction employing CBS. Kamimura and collaborators<sup>80</sup> fabricated CBS on a molybdenum-coated glass substrate employing a simple electrodeposition method using  $\text{Cu}(\text{NO}_3)_2 \cdot 3\text{H}_2\text{O}$  and  $\text{Bi}(\text{NO}_3)_3 \cdot 5\text{H}_2\text{O}$  as precursors. An n-type buffer layer,  $\text{In}_2\text{S}_3$ , was deposited by chemical bath atop CBS films, followed by the deposition of Pt on  $\text{In}_2\text{S}_3/\text{CBS}$  films to improve the photocatalytic response given that Pt nanoparticles present a low overpotential. A band gap of 1.63 eV and shift in the onset potential towards a positive potential were observed for the Pt– $\text{In}_2\text{S}_3/\text{CBS}$  material



**Fig. 10** (a) Schematic illustration of the band diagram of  $\text{TiO}_2$  and CBS heterojunction towards carrier separation and transfer process. (b) LSV polarization curves for  $\text{TiO}_2$ , CBS nanorods and C@T heterojunction. (Inset) Magnified polarization curves for  $\text{TiO}_2$ . Reproduced with permission from ref. 81 Copyright 2018, The Royal Society of Chemistry. (c) Schematic illustration of the fabrication of flower-like CBS and A@C composite. (d) LSV polarization curves of CBS and A@C samples. (Inset): magnified polarization curves for CBS. (e) Rate of  $\text{H}_2$  production by photoelectrocatalysis for CBS and A@C (Au/CBS) samples with different mass ratios. Reproduced with permission from ref. 87 Copyright 2021, The Royal Society of Chemistry. (f) Energy diagram for Au– $\text{In}_2\text{S}_3/\text{CBS}$  electrodes depicting the process of  $\text{H}_2\text{O}_2$  formation. (g) Current–potential curves for Pt– $\text{CBS}_3$  and Pt– $\text{In}_2\text{S}_3/\text{CBS}$ .  $\text{Na}_2\text{SO}_4$  solution (pH 6) was used as the electrolyte, under chopped illumination from AM 1.5 G solar simulator. (Inset): picture of the fabricated  $\text{Cu}_3\text{BiS}_3$  material. Reproduced with permission from ref. 73 Copyright 2022, Elsevier B.V. (h) SEM of CBS film sprayed on substrate at the temperature of  $380^\circ\text{C}$ , and working diagram of CBS– $\text{BiVO}_4$  tandem cell. (i)  $J$ – $V$  curves for CBS photoanode,  $\text{BiVO}_4$  photoanode and CBS photocathode behind the  $\text{BiVO}_4$  photoanode. (Inset:  $J_{\text{op}} =$  operation point photocurrent density and  $\text{STH} =$  solar to hydrogen efficiency). Reproduced under the terms of the CC-BY-4.0 license from ref. 19 (j) Schematic for BS/CBS ( $\text{Bi}_2\text{S}_3/\text{CBS}$ ) photocathode and band diagram calculated from UPS. (k) Photocurrent density versus potential for BS–CBS photocathode,  $\text{BiVO}_4$  photoanode, and BS–CBS photocathode behind the  $\text{BiVO}_4$  photoanode. Reproduced under the terms of the CC-BY license from ref. 20.



compared to Pt-CBS. However, the applied potential necessary for PEC hydrogen evolution was +0.55 ( $\sim$ +0.6 V) vs. RHE in 0.1 M Na<sub>2</sub>SO<sub>4</sub> (pH 6) aqueous solution. At a higher potential, the film suffered from degradation, which was ascribed to the self-reduction and/or reduction of the In<sub>2</sub>S<sub>3</sub> buffer layer. The faradaic efficiency of 90% was maintained after 60 min of light illumination.

Li *et al.*<sup>81</sup> reported the fabrication of a CBS nanorod/TiO<sub>2</sub> (C@T) heterojunction (Fig. 10a) *via* a solution dip-coating method. To analyze the photo-response of CBS nanorods/C@T towards water-splitting potential, PEC measurements were performed under illumination for 10 s. The device showed a higher photocurrent response ( $\sim$ 10<sup>-4</sup> A cm<sup>-2</sup>) for the C@T heterojunction compared to bare CBS nanorods (Fig. 10b). The improvement was due to the formation of a heterojunction, which enhanced the separation and transfer of photogenerated charge carriers, decreased the recombination, and consequently increased the photocurrent.

In another study by Li *et al.*,<sup>45</sup> they investigated the HER. They fabricated flower-like CBS using a Cu and Bi-based solution mixed with Tu. Nanocomposites of Au-CBS (A@C) were also prepared *via* the direct growth of Au NPs on the surface of CBS *via* a one-step solvothermal method (Fig. 10c). Pure CBS presented inactive photoelectrocatalytic behavior (0.35  $\mu$ mol h<sup>-1</sup> of H<sub>2</sub>), as shown in Fig. 10d. In contrast, all the A@C composites exhibited enhanced catalytic activity for H<sub>2</sub> evolution with an increasing load of Au NPs (Fig. 10e) due to their higher contact area and active sites. Besides, the difference in work function between CBS and Au and higher conductivity of Au accelerated the separation of photogenerated charges.

Recently, Chen and collaborators<sup>73</sup> fabricated a uniform polycrystalline Au/In<sub>2</sub>S<sub>3</sub>/CBS thin film *via* spray pyrolysis (more details of the synthesis presented in Section 3.2) for the generation of hydrogen peroxide (H<sub>2</sub>O<sub>2</sub>). Fig. 10f shows the energy band diagram for the studied material. The selectivity for O<sub>2</sub> and H<sub>2</sub>O<sub>2</sub> photoelectrochemical formation improved due to the presence of Au nanoparticles, where 5.5 mg L<sup>-1</sup> h<sup>-1</sup> cm<sup>-1</sup> of H<sub>2</sub>O<sub>2</sub> was produced with a Faraday efficiency of 71%. To the best of our knowledge, this is the first time that CBS was used to produce H<sub>2</sub>O<sub>2</sub>, which shows its versatility for different photoelectrochemical applications. The current–voltage curves for Pt–Cu<sub>3</sub>BiS<sub>3</sub> and Pt–In<sub>2</sub>S<sub>3</sub>/Cu<sub>3</sub>BiS<sub>3</sub> are shown in Fig. 10g.

In a recent study by Huang *et al.*,<sup>82</sup> they fabricated CBS employing one-step spray pyrolysis and tested it for the PEC hydrogen evolution reaction (HER). The influence of the substrate temperature (360–420 °C) and Cu : Bi molar ratio (1 : 1, 2 : 1, 4 : 1 and 5 : 1) in the precursor solution was investigated. Among them, the Cu : Bi molar ratio of 3 : 1 matched well with the reference pattern for CBS without showing secondary phases, while the other Cu : Bi ratios presented secondary phases (Cu<sub>2-x</sub>S and Cu<sub>3</sub>Bi<sub>3</sub>S<sub>7</sub>). The 3 : 1 Cu : Bi thin-film at 380 °C showed better photocatalytic behavior upon further coating with CdS, TiO<sub>2</sub> and Pt nanoparticles. The device was coupled to a BiVO<sub>4</sub> photoanode to build a tandem cell (CBS–BiVO<sub>4</sub>). For this device, they achieved high stability (over 60 h) and STH efficiency (2.04%) for the bias-free standalone system. An illustration of the cell and the *J–V* curves of the materials are presented in Fig. 10h and I, respectively.

The presence of mixed phases can result in structural instability, which can lead to a loss in the photocatalytic activity and overall performance. However, recently, in a report by Moon *et al.*,<sup>20</sup> it was shown that mixed phases can actually allow better band alignment with the n-type layer (band diagram showing the relative energy positions of each of the semiconductors is shown in Fig. 10j). This changes the band structure of the material, which leads to an improvement in charge transport. The Bi<sub>2</sub>S<sub>3</sub>–Cu<sub>3</sub>BiS<sub>3</sub> (BS-CBS) mixed phase with elongated 1D nanorods achieved a high onset potential of 0.9 V and showed three-times higher performance compared to the single-phase CBS. Thus, the BS-CBS internal heterojunction facilitated charge separation. A standalone PEC–PEC tandem device was proposed by coupling the BS-CBS photocathode with a BiVO<sub>4</sub>-based photoanode. Fig. 10k shows the *J–V* curves for BS-CBS, Mo:BiVO<sub>4</sub>, and BS-CBS photocathode behind the BiVO<sub>4</sub> photoanode. The BS-CBS-based tandem cell showed an efficiency of 2.33% and stability of over 21 h.

The photo(electro)reduction of CO<sub>2</sub> using CBS has not been demonstrated to date. The first work on CO<sub>2</sub> reduction using photo electrocatalysis was reported by Halmann *et al.*,<sup>83</sup> in which CO<sub>2</sub> was converted to methanol, formic acid and formaldehyde by a GaP electrode. Since then, various studies have been proposed. The complex dependence among the reaction products and intermediates, supporting electrolyte, applied potential, pH gradient, and reaction kinetics is not well understood. Besides, in the CO<sub>2</sub>RR, losses are caused by band bending (kinetic), resistance (electrolyte) and the overvoltage potential, and therefore an energy input of higher than 1.7 eV is necessary to overcome this limitation. In this case, a material presenting a bandgap of >1.8–2.0 eV is required for CO<sub>2</sub> photo electroreduction application, and therefore CBS is a potential candidate for CO<sub>2</sub> reduction, given that it fulfills most of the above-mentioned criteria. However, the low photovoltage and aqueous stability of CBS require detailed studies in this front. It is well known that copper-based materials suffer from photo corrosion (photo instability). For instance, copper oxide suffers from redox instabilities and photodegradation due to the oxidation/reduction of Cu<sup>+</sup> (Cu<sub>2</sub>O) to form Cu<sup>2+</sup> (CuO) or Cu<sup>0</sup>, respectively.<sup>84–86</sup> There are many strategies to overcome this limitation, which includes the creation of a p–n heterojunction to avoid the recombination of photogenerated charge carriers and multifunctional coatings to protect the surface and improve the long-term durability.

## 5 Future perspectives

Bi-based chalcogenides have attracted significant attention, presenting interesting optoelectronic properties and multifunctionalities, which make them suitable for application in solar cells and PECs. We observed that while the low *V*<sub>oc</sub> of CBS-based photocathodes in solar cells is a major limitation, their show an excellent photoelectrochemical performance with a high onset potential. The defect tolerance from the electronic configuration needs to be experimentally scrutinized. Although antibonding VBM states generally lead to shallow defect levels in the band gap for specific point defects (*e.g.*, Cu vacancy in



CuInSe<sub>2</sub> forms a shallow acceptor level due to the antibonding p-d states), it is paramount to consider all the possible point defects to examine the recombination centers or traps, given that some of them can act as recombination centers or carrier traps. Deep defects have been observed in Cu<sub>3</sub>N (interstitial defect Cu<sub>i</sub>) and CIGSe(S)<sub>2</sub> (antisite (In,Ga)<sub>Cu</sub>) semiconductors despite having a defect tolerant electronic configuration.<sup>88–90</sup> Addressing the  $V_{oc}$  deficit in CBS requires an in-depth understanding of the factors responsible for the  $V_{oc}$  losses, as follows: (i) phase purity and doping control, (ii) identification of deep defects, (iii) determination of carrier lifetime and mobility limiting mechanism, (iv) role of intrinsic disorder and band tails, and (v) knowledge of energy band positions to construct heterojunctions with minimized interface recombination.

Here, we provide possible solutions towards the development of CBS.

(i) Phase purity and stoichiometric control: although many reports claim the synthesis of phase-pure CBS based on XRD analysis, it is possible that a low concentration of impurity phases may be excluded in the traditional XRD measurements. In this case, the more sensitive Raman spectroscopy will be useful to rule out the presence of secondary phases. Moreover, the complex carbon-containing precursors in the commonly used solution processing and relatively low processing temperatures can lead to nanoscale amorphous phase impurities, which remain undetected in XRD and Raman measurements. Accordingly, high-resolution transmission electron microscopy and sensitive magnetic measurements will be useful to identify these defects and amorphous phase domains. It is worth investigating single crystals and vacuum processed polycrystalline thin films to make a comparative study on the inevitability of certain defects and phase impurities in CBS.

(ii) Open questions on intrinsic properties and defects: the fundamental properties such as the electronic band structure,

nature of the optical bandgap, charge carrier diffusion lengths (carrier mobility and lifetimes), doping, intrinsic defects and their respective energies in the bandgap are not well understood. The lack of control of the electronic properties explains the poor performance of photovoltaic devices. It is important to establish whether the carrier mobilities are limited by phonon scattering or impurity scattering. A recent study reported a low carrier lifetime of 0.7 ns, indicating high recombination losses.<sup>20</sup> This clearly suggests the active role of deep defects in the material.

Theoretical studies will be highly useful to gain insights into the optoelectronic properties of CBS and validate experimental conclusions. For instance, experimental results show direct bandgap behavior for CBS, while theoretical calculations predict indirect behavior. Thus, this discrepancy needs to be addressed. Furthermore, the role of deep defects and their importance is still lacking. Atomic-level insights into defects can be obtained by advanced characterization techniques such as scanning transmission microscopy (STM) and X-ray photoelectron spectroscopy (XPS). STM allows the direct visualization of lattice defects and XPS enables the measurement of the chemical and electronic states of the elements. Spectroscopic techniques such as photoluminescence and femtosecond transient spectroscopy will be useful to analyze the performance-limiting deep defects and carrier dynamics in CBS. The knowledge of performance-limiting defects is necessary to engineer high-performing devices and develop passivation strategies.

(iii) Improving the optoelectronic properties and charge transport: control of the doping and morphology is critical to improve the optoelectronic properties. An effective doping strategy may be guided by leveraging the knowledge of doping in common chalcogenide absorbers (CIGS and CZTS). Varying the growth conditions from Cu-poor (Bi-rich) to Cu-rich (Bi-



Fig. 11 CBS device architecture engineering for (a) p-i-n solar cell, (b) n-i-p solar cell, (c) CBS-sensitized solar cell, (d) mesoporous solar cell and (e) photoelectrochemical cell. ETL = electron transport layer; HTL = hole transport layer; and TCO = transparent conductive oxide.



poor), cation substitution with Ag/Na, anion substitution with oxygen/selenium, and alkali doping are promising routes in this regard. Besides doping, control of the morphology is desirable for specific applications. It is worth emphasizing that the impurity phases, which are otherwise considered detrimental, may have better optical and electrical properties than wittichenite CBS itself. Thus, the controllable synthesis of different phases may open new avenues for the Cu–Bi–S phase. The knowledge of energy band alignment with common ETL and HTL contact layers is essential to enhance the charge carrier transport and maximize the carrier extraction efficiency. The application of the conventionally used buffer layers, such as CdS, may not be suitable in terms of band alignment. In this case, alternative buffer layers such as Zn(O,S), In<sub>2</sub>S<sub>3</sub>, and (Zn,Sn)O may be interesting to explore. Besides the buffer layers, the band offsets can also be changed by tuning the energy band of CBS by alloying and/or cation/anion substitution. However, explicit observation *via* experimental measurement for tuning the band edge energies of CBS is lacking to date.

(iv) Interface-engineered device architectures: there are few reports on CBS photovoltaic devices, and most of them show a low efficiency. Thus, more studies on high-quality and defect-passivated CBS absorbers<sup>91</sup> are required. More device architectures need to be explored to assess the full potential of CBS. Fig. 11 shows the feasible device architectures.

(v) Long-term stability: the stability of the photocathode is a major concern in photoelectrochemical applications, which may be due to the environment (oxidation) and aqueous medium (redox instability). Additionally, photoinduced degradation (photocorrosion) is also a major cause of the degradation of many photocathodes, such as Cu<sub>2</sub>O. A protection barrier is required to prevent the degradation in PC and PEC devices. Laminating the surface with hybrid polymeric encapsulants embedded with catalysts provides higher operational stability. Atomic layer deposition (ALD) is also an option to suppress corrosion or add functionality to the device (*i.e.*, MoS<sub>2</sub> exhibits intrinsic catalytic activity). Finally, the construction of heterostructures or Z-schemes can also be an alternative to enhance the stability for (photo)electrocatalytic reactions.

## 6 Concluding remarks

In summary, the Cu<sub>x</sub>Bi<sub>y</sub>S<sub>z</sub> system remains unexplored, requiring further studies on this front. In particular, the Cu<sub>3</sub>BiS<sub>3</sub> compound exhibits excellent optical absorption characteristics and suitable bandgap for solar cell and photoelectrochemical applications. CBS offers new opportunities for both domains. Although photovoltage limitations have hampered the development of CBS-based solar cells, promising results on photoelectrochemical water splitting have opened new avenues. More complex photoelectrochemical reactions, such as CO<sub>2</sub> reduction can be targeted with modified interfaces and surfaces. Nevertheless, fundamental studies with innovative device designs are required to fully harness the potential of the CBS absorber. The development of an effective and stable visible light active photocatalyst is crucial to achieve high efficiency in (photo)electrochemical reactions.

## Author contributions

Daniely S. and Sudhanshu S.: conceptualization, writing – original draft. Sudhanshu S. and Bart V.: reviewing and editing, supervision. Bart V.: funding acquisition.

## Conflicts of interest

There are no conflicts to declare.

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## References

- 1 NREL, Best Research-Cell Efficiency Chart, <https://www.nrel.gov/pv/cell-efficiency.html>, accessed 20-07-2022.
- 2 B. Zohuri and P. McDaniel, in *Introduction to Energy Essentials*, ed. B. Zohuri and P. McDaniel, Academic Press, 2021, pp. 255–276.
- 3 S. Ikeda, *High Temp. Mater. Processes*, 2021, **40**, 446–460.
- 4 A. Redinger, D. M. Berg, P. J. Dale and S. Siebentritt, *J. Am. Chem. Soc.*, 2011, **133**, 3320–3323.
- 5 T. Tanaka, T. Sueishi, K. Saito, Q. Guo, M. Nishio, K. M. Yu and W. Walukiewicz, *J. Appl. Phys.*, 2012, **111**, 053522.
- 6 E. Deady, C. Moon, K. Moore, K. M. Goodenough and R. K. Shail, *Ore Geol. Rev.*, 2022, **143**, 104722.
- 7 R. Mohan, *Nat. Chem.*, 2010, **2**, 336.
- 8 U. S. G. Survey, *Bismuth. Mineral Commodity Summaries*, 2021.
- 9 U. S. G. Survey, *Bismuth Mineral Commodity Summaries*, 2014.
- 10 5NPlus, *2019 Annual Report*, 2019.
- 11 Y. Fang, W. Zhao, W. Li and X. Han, *Appl. Phys. A*, 2021, **127**, 9.
- 12 S. G. Deshmukh and V. Kheraj, *Nanotechnol. Environ. Eng.*, 2017, **2**, 15.
- 13 Y. Wang, S. R. Kavanagh, I. Burgués-Ceballos, A. Walsh, D. O. Scanlon and G. Konstantatos, *Nat. Photonics*, 2022, **16**, 235–241.
- 14 Y.-T. Huang, S. R. Kavanagh, M. Righetto, M. Rusu, I. Levine, T. Unold, S. J. Zelewski, A. J. Sneyd, K. Zhang, L. Dai, A. J. Britton, J. Ye, J. Julin, M. Napari, Z. Zhang, J. Xiao, M. Laitinen, L. Torrente-Murciano, S. D. Stranks, A. Rao, L. M. Herz, D. O. Scanlon, A. Walsh and R. L. Z. Hoyer, *Nat. Commun.*, 2022, **13**, 4960.
- 15 T. J. Whittles, T. D. Veal, C. N. Savory, P. J. Yates, P. A. E. Murgatroyd, J. T. Gibbon, M. Birkett, R. J. Potter, J. D. Major, K. Durose, D. O. Scanlon and V. R. Dhanak, *ACS Appl. Mater. Interfaces*, 2019, **11**, 27033–27047.
- 16 F. Mesa and G. Gordillo, *J. Phys.: Conf. Ser.*, 2009, **167**, 012019.



- 17 W. Shockley and H. J. Queisser, *J. Appl. Phys.*, 1961, **32**, 510–519.
- 18 S. Rühle, *Sol. Energy*, 2016, **130**, 139–147.
- 19 D. Huang, L. Li, K. Wang, Y. Li, K. Feng and F. Jiang, *Nat. Commun.*, 2021, **12**, 3795.
- 20 S. Moon, J. Park, H. Lee, J. W. Yang, J. Yun, Y. S. Park, J. Lee, H. Im, H. W. Jang, W. Yang and J. Moon, *Adv. Sci.*, 2023, **10**, 2206286.
- 21 E. W. Nuffield, *Econ. Geol.*, 1947, **42**, 147–160.
- 22 V. Kocman and E. W. Nuffield, *Acta Crystallogr., Sect. B: Struct. Crystallogr. Cryst. Chem.*, 1973, **29**, 2528–2535.
- 23 K. W. Böer and U. W. Pohl, in *Semicond. Phys.*, ed. K. W. Böer and U. W. Pohl, Springer International Publishing, Cham, 2018, vol. 3, pp. 59–108.
- 24 J. Capistrán-Martínez, D. Loeza-Díaz, D. Mora-Herrera, F. Pérez-Rodríguez and M. Pal, *J. Alloys Compd.*, 2021, **867**, 159156.
- 25 M. Kumar and C. Persson, *J. Renewable Sustainable Energy*, 2013, **5**, 031616.
- 26 A. B. Kehoe, D. J. Temple, G. W. Watson and D. O. Scanlon, *Phys. Chem. Chem. Phys.*, 2013, **15**, 15477–15484.
- 27 K. M. Koskela, A. C. Taddeo, K. Chen and R. L. Brutchey, *ACS Appl. Energy Mater.*, 2021, **4**, 11026–11031.
- 28 A. Walsh, D. J. Payne, R. G. Egdell and G. W. Watson, *Chem. Soc. Rev.*, 2011, **40**, 4455–4463.
- 29 F. Miao, S. Wan, Q. Zhang, X. Lu, G. Gu and S. Yin, *Chin. J. Quantum Electron.*, 2011, **28**, 558.
- 30 M. F. Razmara, C. M. B. Henderson, R. A. D. Patrick, A. M. T. Bell and J. M. Charnock, *Mineral. Mag.*, 1997, **61**, 79–88.
- 31 D. Colombara, L. M. Peter, K. Hutchings, K. D. Rogers, S. Schäfer, J. T. R. Duffton and M. S. Islam, *Thin Solid Films*, 2012, **520**, 5165–5171.
- 32 W. Chamorro, F. Mesa, M. Hurtado and G. Gordillo, Study of structural and morphological properties of ZnS films deposited on Cu<sub>3</sub>BiS<sub>3</sub>, in *25th European Photovoltaic Solar Energy Conference and Exhibition*, Spain, 2010, pp. 575–579, DOI: [10.4229/25thEUPVSEC2010-1DV.3.8](https://doi.org/10.4229/25thEUPVSEC2010-1DV.3.8).
- 33 A. Kitakaze, *Phase Relations of Some Sulfide Systems*, Mem Fac Eng Yamaguchi Univ, 2017.
- 34 A. Zakutayev, C. M. Caskey, A. N. Fioretti, D. S. Ginley, J. Vidal, V. Stevanovic, E. Tea and S. Lany, *J. Phys. Chem. Lett.*, 2014, **5**, 1117–1125.
- 35 R. M. Córdova-Castro, M. Casavola, M. van Schilfgaarde, A. V. Krasavin, M. A. Green, D. Richards and A. V. Zayats, *ACS Nano*, 2019, **13**, 6550–6560.
- 36 H. Ben Abdallah and W. Ouerghui, *Opt. Quantum Electron.*, 2021, **54**, 20.
- 37 R. S. Silva, H. D. Mikhail, E. V. Guimarães, E. R. Gonçalves, N. F. Cano and N. O. Dantas, *Molecules*, 2017, **22**, 1142.
- 38 C. Kittel, *Introduction to Solid State Physics*, Wiley, 2004.
- 39 J. Y. Kim, J.-W. Lee, H. S. Jung, H. Shin and N.-G. Park, *Chem. Rev.*, 2020, **120**, 7867–7918.
- 40 W.-J. Yin, T. Shi and Y. Yan, *J. Phys. Chem. C*, 2015, **119**, 5253–5264.
- 41 R. v. d. Krol and M. Graetzel, *Electronic Materials, Science & Technology, Photoelectrochemical Hydrogen Production*, 2012.
- 42 F. Mesa, G. Gordillo, T. Dittrich, K. Ellmer, R. Baier and S. Sadewasser, *Appl. Phys. Lett.*, 2010, **96**, 082113.
- 43 B. Murali, M. Madhuri and S. B. Krupanidhi, *J. Appl. Phys.*, 2014, **115**, 173109.
- 44 B. B. Dumre, N. J. Szymanski, V. Adhikari, I. Khatri, D. Gall and S. V. Khare, *Sol. Energy*, 2019, **194**, 742–750.
- 45 J. Li, X. Han, Y. Zhao, J. Li, M. Wang and C. Dong, *Sol. Energy Mater. Sol. Cells*, 2018, **174**, 593–598.
- 46 T. Fazal, S. Iqbal, M. Shah, Q. Mahmood, B. Ismail, H. O. Alsaab, N. S. Awwad, H. A. Ibrahim and E. B. Elkaeed, *Results Phys.*, 2022, **36**, 105453.
- 47 D. Lee, H. Ahn, S. Park, H. Shin and Y. Um, *Nanosci. Nanotechnol. Lett.*, 2018, **10**, 1–4.
- 48 A. M. Russell and K. L. Lee, in *Structure-Property Relations in Nonferrous Metals*, 2005, pp. 18–27.
- 49 T. Wei, Y.-N. Zhu, X. An, L.-M. Liu, X. Cao, H. Liu and J. Qu, *ACS Catal.*, 2019, **9**, 8346–8354.
- 50 D. Maarisetty and S. S. Baral, *J. Mater. Chem. A*, 2020, **8**, 18560–18604.
- 51 Z. Wang, M. Xiao, J. You, G. Liu and L. Wang, *Acc. Mater. Res.*, 2022, **3**, 1127–1136.
- 52 S. Trivedi, D. Prochowicz, N. Parikh, A. Mahapatra, M. K. Pandey, A. Kalam, M. M. Tavakoli and P. Yadav, *ACS Omega*, 2021, **6**, 1030–1042.
- 53 J. Zhong, W. Xiang, Q. Cai and X. Liang, *Mater. Lett.*, 2012, **70**, 63–66.
- 54 M. V. Yakushev, P. Maiello, T. Raadik, M. J. Shaw, P. R. Edwards, J. Krustok, A. V. Mudryi, I. Forbes and R. W. Martin, *Thin Solid Films*, 2014, **562**, 195–199.
- 55 O. C. Olatunde and D. C. Onwudiwe, *Results Chem.*, 2022, **4**, 100273.
- 56 Y. Pan, H. Deng, D. Zheng, J. Chen, Y. Zhong, J. Tao, P. Yang and J. Chu, *J. Mater. Sci.: Mater. Electron.*, 2022, **33**, 585–595.
- 57 L. Zhang, X. Jin, C. Yuan, G. Jiang, W. Liu and C. Zhu, *Appl. Surf. Sci.*, 2016, **389**, 858–864.
- 58 Y. Yang, X. Xiong, H. Yin, M. Zhao and J. Han, *J. Mater. Sci.: Mater. Electron.*, 2019, **30**, 1832–1837.
- 59 A. Hussain, J. T. Luo, P. Fan, G. Liang, Z. Su, R. Ahmed, N. Ali, Q. Wei, S. Muhammad, A. R. Chaudhry and Y. Q. Fu, *Appl. Surf. Sci.*, 2020, **505**, 144597.
- 60 W. Zhao, Y. Fang, W. Li and X. Han, *Sol. Energy*, 2021, **221**, 109–113.
- 61 Y.-B. Lu, W.-Y. Cong, C. Guan, H. Sun, Y. Xin, K. Wang and S. Song, *J. Mater. Chem. A*, 2019, **7**, 27469–27474.
- 62 S. Siebentritt, M. Igalson, C. Persson and S. Lany, *Prog. Photovoltaics*, 2010, **18**, 390–410.
- 63 S. R. Kavanagh, A. Walsh and D. O. Scanlon, *ACS Energy Lett.*, 2021, **6**, 1392–1398.
- 64 S. U. Rahayu and M.-W. Lee, *AIP Conf. Proc.*, 2020, **2221**, 030005.
- 65 M. Chakraborty, R. Thangavel, P. Komninou, Z. Zhou and A. Gupta, *J. Alloys Compd.*, 2019, **776**, 142–148.
- 66 E. Verónica, M. T. S. Nair and P. K. Nair, *Semicond. Sci. Technol.*, 2003, **18**, 190.
- 67 N. J. Gerein and J. A. Haber, *Chem. Mater.*, 2006, **18**, 6297–6302.



- 68 C. Yan, E. Gu, F. Liu, Y. Lai, J. Li and Y. Liu, *Nanoscale*, 2013, **5**, 1789–1792.
- 69 X. Ren, L. Yuan, Q. Liang, R. Xie, Z. Geng, Y. Sun, L. Wang, K. Huang, T. Wu and S. Feng, *ChemPlusChem*, 2018, **83**, 812–818.
- 70 S.-K. Han, C. Gu, M. Gong and S.-H. Yu, *J. Am. Chem. Soc.*, 2015, **137**, 5390–5396.
- 71 S. G. Deshmukh, S. J. Patel, K. K. Patel, A. K. Panchal and V. Kheraj, *J. Electron. Mater.*, 2017, **46**, 5582–5588.
- 72 J. Hu, B. Deng, C. Wang, K. Tang and Y. Qian, *Mater. Chem. Phys.*, 2003, **78**, 650–654.
- 73 C. Chen, M. Yasugi, L. Yu, Z. Teng and T. Ohno, *Appl. Catal., B*, 2022, **307**, 121152.
- 74 S. Liu, X. Wang, L. Nie, L. Chen and R. Yuan, *Thin Solid Films*, 2015, **585**, 72–75.
- 75 N. Pai, J. Lu, D. C. Senevirathna, A. S. R. Chesman, T. Gengenbach, M. Chatti, U. Bach, P. C. Andrews, L. Spiccia, Y.-B. Cheng and A. N. Simonov, *J. Mater. Chem. C*, 2018, **6**, 2483–2494.
- 76 P. K. Nair, L. Huang, M. T. S. Nair, H. Hu, E. A. Meyers and R. A. Zingaro, *J. Mater. Res.*, 1997, **12**, 651–656.
- 77 J. Yin and J. Jia, *CrystEngComm*, 2014, **16**, 2795–2801.
- 78 F. Mesa and D. Fajardo, *Can. J. Phys.*, 2014, **92**, 892–895.
- 79 M. G. Walter, E. L. Warren, J. R. McKone, S. W. Boettcher, Q. Mi, E. A. Santori and N. S. Lewis, *Chem. Rev.*, 2010, **110**.
- 80 S. Kamimura, N. Beppu, Y. Sasaki, T. Tsubota and T. Ohno, *J. Mater. Chem. A*, 2017, **5**, 10450–10456.
- 81 J. Li, X. Han, M. Wang, Y. Zhao and C. Dong, *New J. Chem.*, 2018, **42**, 4114–4120.
- 82 Y.-T. Huang, S. R. Kavanagh, D. O. Scanlon, A. Walsh and R. L. Z. Hoyer, *Nanotechnology*, 2021, **32**, 132004.
- 83 M. Halmann, *Nature*, 1978, **275**, 115–116.
- 84 T. Baran, A. Visibile, M. Busch, X. He, S. Wojtyla, S. Rondinini, A. Minguzzi and A. Vertova, *Molecules*, 2021, **26**, 7271.
- 85 S. Tiwari, S. Kumar and A. K. Ganguli, *J. Photochem. Photobiol., A*, 2022, **424**, 113622.
- 86 Q. Peng, S. Zhang, H. Yang, B. Sheng, R. Xu, Q. Wang and Y. Yu, *ACS Nano*, 2020, **14**, 6024–6033.
- 87 J. Li, Y. Zhao, X. Han and D. Xiao, *New J. Chem.*, 2021, **45**, 1231–1239.
- 88 Y. S. Yee, H. Inoue, A. Hultqvist, D. Hanifi, A. Salleo, B. Magyari-Köpe, Y. Nishi, S. F. Bent and B. M. Clemens, *Phys. Rev. B*, 2018, **97**, 245201.
- 89 Q. Cao, O. Gunawan, M. Copel, K. B. Reuter, S. J. Chey, V. R. Deline and D. B. Mitzi, *Adv. Energy Mater.*, 2011, **1**, 845–853.
- 90 S. Shukla, M. Sood, D. Adeleye, S. Peedle, G. Kusch, D. Dahliah, M. Melchiorre, G.-M. Rignanese, G. Hautier, R. Oliver and S. Siebentritt, *Joule*, 2021, **5**, 1816–1831.
- 91 F. Mesa, A. Dussan, B. A. Paez-Sierra and H. Rodriguez-Hernandez, *Univ. Sci.*, 2014, **19**, 99–105.

