

Cite this: *Chem. Commun.*, 2012, **48**, 6544–6546

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# Prussian blue: a new framework of electrode materials for sodium batteries†

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Received 10th March 2012, Accepted 4th May 2012

DOI: 10.1039/c2cc31777j

Prussian blue and its analogues consisting of different transition-metal ions (Fe, Mn, Ni, Cu, Co and Zn) have been synthesized at room temperature. Insertion of Na into  $\text{KFe}_2(\text{CN})_6$  in a carbonate electrolyte exhibited a reversible capacity near  $100 \text{ mA h g}^{-1}$  with no capacity fade in 30 cycles. The data indicate that a Na-ion battery with a Prussian blue framework as a cathode will be feasible.

The Li-ion battery has played a dominant role in portable electronic devices and power tools; it is being developed actively for plug-in hybrid vehicles. Its success places an increasing demand on sources of lithium, which has raised a desire to build a Na-ion battery as an alternative.<sup>1,2</sup> The sodium–sulfur battery of Kummer and Weber<sup>3</sup> operates at 300 to 350 °C, which creates maintenance problems and elevated operational costs. Therefore, there is motivation to find a room-temperature rechargeable Na-ion battery. Present Li-ion rechargeable batteries rely on an oxide host of the cathode from/into which Li can be extracted/inserted reversibly.<sup>4–11</sup> The  $\text{Li}^+$  ion is small enough to have an acceptable mobility and a solid-solution range in close-packed oxide-ion arrays; but the larger  $\text{Na}^+$  ion requires a more open framework in which to move reversibly with an acceptable mobility. The hexagonal framework of  $\text{Na}_{1+3x}\text{Zr}_2(\text{P}_{1-x}\text{Si}_x\text{O}_4)_3$ , commonly referred to as NASICON, exhibits fast  $\text{Na}^+$ -ion conductivity;<sup>12</sup> its framework has the hexagonal  $\text{Fe}_2(\text{SO}_4)_3$  structure. Unfortunately, the  $\text{Na}^+$  mobility is greatly reduced where the  $\text{M}_2(\text{XO}_4)_3$  units of the framework are rotated with respect to one another to give a monoclinic structure, and substitution of a transition-metal of interest for Zr(IV) gives the monoclinic structure if  $\text{Na}^+$  is the guest alkali ion.<sup>13</sup> Therefore, there is interest in identifying another open host framework containing a suitable transition-metal ion.

$\text{KFe}_2(\text{CN})_6$ , Prussian blue, has a cubic framework (space group  $Fm\bar{3}m$ ) with Fe(II) and Fe(III) on alternate corners of a cube of corner-shared octahedra bridged by linear  $(\text{C}\equiv\text{N})^-$  anions (Fig. 1); the low-spin Fe(III) bond only with C atoms, the high-spin Fe(II) only with N atoms,<sup>14–16</sup> and the  $\text{C}\equiv\text{N}$



Fig. 1 Framework of Prussian blue analogues.

bond opens the faces of the elementary cubes for  $\text{Na}^+$  to move between half-filled body-center positions.

Other transition-metal M(II) ions can replace the Fe(II) of Prussian blue in the family  $\text{KMFe}(\text{CN})_6$ . Widman *et al.*<sup>17</sup> and later Wessells *et al.*<sup>18–20</sup> have investigated Na and K insertion into  $\text{KMFe}(\text{CN})_6$  with  $\text{M} = \text{Cu}$  and  $\text{Ni}$  using a Pt counter electrode and an Ag–AgCl reference electrode with an aqueous electrolyte. The cell had low potentials owing to the limitation of the aqueous electrolyte and the reference electrode. Nevertheless, their data invite exploration of reversible insertion of  $\text{Na}^+$  into a Prussian blue analogue with an aprotic electrolyte. Therefore, we have synthesized  $\text{KMFe}(\text{CN})_6$  compounds (MFHC) with  $\text{M} = \text{Mn}, \text{Fe}, \text{Co}, \text{Ni}$  and  $\text{Zn}$ ; and we show that in an organic liquid-carbonate electrolyte and a Na anode,  $\text{KFe(II)Fe(III)(CN)}_6$  can exhibit a reversible capacity of *ca.*  $100 \text{ mA h g}^{-1}$ .

In this communication, MFHC are synthesized at room temperature. Briefly, the M(II) salt solution of 0.1 M was slowly added into the  $\text{K}_3\text{Fe}(\text{CN})_6$  solution of 0.05 M with a strong magnetic stirring. The total mole of the M(II) salt is twice that of  $\text{K}_3\text{Fe}(\text{CN})_6$ . After 30 minutes of magnetic stirring, the mixed solution was moved to ultrasonic stirring. The process was kept for 2 hours. The products of  $\text{KM(II)FHC}$  were filtered and washed with DI-water three times. The final products were dried in a vacuum box at 70 °C overnight. These compounds were stored in a vacuum desiccator.

The structures of the Prussian blue compounds were characterized by Powder X-ray diffraction (XRD) obtained with a Philips X-ray diffractometer equipped with Cu  $\text{K}\alpha$  radiation ( $\lambda = 1.5418 \text{ \AA}$ ). The angular resolution in  $2\theta$  scans was  $0.02^\circ$  over a  $2\theta$  range of  $10\text{--}70^\circ$ . The powder X-ray diffraction (XRD) patterns of Fig. 2a for the  $\text{KM(II)Fe(III)(CN)}_6$  compounds show that  $\text{KZnFe}(\text{CN})_6$  forms a hexagonal lattice (space group  $R\bar{3}c$ ) as has been reported.<sup>21</sup> Whereas all the

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† Electronic supplementary information (ESI) available. See DOI: 10.1039/c2cc31777j

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**Fig. 2** (a) XRD patterns of Prussian blue compounds. (b)  $\chi_m^{-1}$ - $T$  plot of Prussian blue compounds.

other M(II) cations simply change the cubic lattice parameter;  $a_0 = 10.170 \text{ \AA}$ ,  $10.151 \text{ \AA}$ ,  $10.095 \text{ \AA}$ ,  $10.119 \text{ \AA}$ ,  $10.112 \text{ \AA}$ , respectively, for  $M = \text{Mn, Fe, Co, Ni, and Cu}$ . It is noticed that with  $M = \text{Co}$ , the lattice parameter is quite close to the high temperature phase of cobalt-iron cyanide, indicating that the compound mostly consists of  $\text{Fe}^{\text{III}}(\text{t}_{2g}^5\text{e}_g^0, S = 1/2)$ - $\text{CN-Co}^{\text{II}}(\text{t}_{2g}^5\text{e}_g^2, S = 3/2)$ .<sup>22</sup> In order to further determine the electron configuration, magnetic susceptibility of all products was studied. Fig. 2b shows the inverse molar magnetic susceptibility ( $\chi_m^{-1}$ ) as a function of temperature;  $\mu_B$  calculated from this plot indicates that Fe(III) remains in the low-spin with  $S = 1/2$  while  $M = \text{Mn, Co, Ni, Cu, Zn}$  have, respectively, spins  $S \approx 5/2, 3/2, 1, 1/2, 0$ .

Whether a localized-electron configuration on an octahedral-site cation is in a high-spin or a low-spin state depends on the relative magnitudes of the cubic-field splitting  $\Delta_c$  of the  $\pi$ -bonding  $t_2$  and  $\sigma$ -bonding  $e$  orbitals of the d-state manifold *versus* the exchange energy  $\Delta_{\text{ex}}$  responsible for Hund's highest-spin

rule in the free atom; a low-spin configuration has a  $\Delta_c > \Delta_{\text{ex}}$ , a high-spin configuration has  $\Delta_c < \Delta_{\text{ex}}$ . From second-order perturbation theory appropriate for the covalent interaction to the bonding of localized 3d electrons,  $\Delta\varepsilon = |b|^2/\Delta E$ , where  $b = \langle \Psi_d | H' | \Psi_p \rangle \approx \varepsilon(\Psi_d, \Psi_p)$  varies as the overlap integral ( $\Psi_d, \Psi_p$ ) of the cation-d and anion-p orbitals and  $\Delta E$  is the energy difference of the lowest empty cation-d and highest filled anion-p orbitals in the absence of covalent bonding. The  $(\text{C} \equiv \text{N})^-$  ion has its  $p_\pi$  orbitals primarily active in the internal triple bond whereas there is a strong  $p_\sigma$  overlap with the cation-d orbitals. Therefore, since  $\Delta E$  is the same for both the  $\pi$ -bonding  $t$  and  $\sigma$ -bonding  $e$  orbitals, the difference in the magnitude of the  $p_\sigma$  and  $p_\pi$  overlap integrals makes  $\Delta_c$  larger for a  $(\text{CN})^-$  ion than that for an  $\text{O}^{2-}$  ion. In addition,  $\Delta E$  for the M-C bond is smaller than that for the M-N bond, so  $\Delta_c$  for the Fe(III) bonding to C is much larger than  $\Delta_c$  for the M(II) bonding to N. As a result, the Fe(III) is in its low-spin state despite a large  $S = 5/2$  in the high-spin state whereas the Fe(II), which bond to N, remains in its high-spin state despite a smaller  $S = 2$ . Nevertheless, we should anticipate a  $\Delta_c$  only a little smaller than  $\Delta_{\text{ex}}$  for the Fe(II) of Prussian blue.

Fig. 3 shows the electrochemical behavior of  $\text{KMFe}(\text{CN})_6\text{-Na}$  cells at the third cycle with a current of  $C/20$  in the range 2.0 to 4.0 V vs.  $\text{Na}^+/\text{Na}$ . In these cells, Na is being inserted reversibly into the  $\text{KMFe}(\text{CN})_6$  cathode and the low-spin Fe(III) is being converted to low-spin Fe(II). However, during charge, some  $\text{K}^+$  ions are removed as well as  $\text{Na}^+$  ions, so more than one  $\text{Na}^+$  ion per formula unit is inserted during the second and third discharge. Therefore, on the third discharge-charge cycle, the  $\text{KFe}_2(\text{CN})_6$  cathode shows two peaks at 2.97 and 3.69 V on its charge curves and two peaks at 2.92 and 3.58 V on its discharge curves. The peaks at 2.97/2.92 V correspond to oxidation-reduction of the high-spin Fe(III)-Fe(II) couple bonding to N and those at 3.69/3.58 V to the low-spin Fe(III)-Fe(II) couple bonding to C.<sup>17</sup> The  $M = \text{Co, Ni, and Cu}$  show that the voltages of the low-spin Fe(III)-Fe(II) couple are all similar, but lower than that of  $\text{KFe}_2(\text{CN})_6$ ; there is little, if any, removal of  $\text{K}^+$  ions on charge where the lattice parameter



**Fig. 3** The charge-discharge curves of Prussian blue analogues  $\text{KMFe}(\text{CN})_6\text{-Na}$  cells at the third cycle with a current of  $C/20$ . Insets show their corresponding chronoamperograms.



**Fig. 4** (a) The redox potentials of the low-spin Fe(II)–Fe(III) in the Prussian blue compounds. (b) Cycling performance of KFeFHC–Na battery cycled at C/20 between 2.0–4.0 V.

$a_0$  is smaller. Although more  $K^+$  is removed from  $KMnFe(CN)_6$  which has the largest lattice parameter, nevertheless relatively little participation of the Mn(III)–Mn(II) couple is found at 3.82/3.56 V; instead it exhibits a somewhat larger discharge capacity on the low-spin Fe(III)–Fe(II) couple. These features of the redox potentials and capacities of the low-spin redox energies are summarized in Fig. 4a. The  $KFe_2(CN)_6$ –Na battery in EC : DEC (1 : 1 v/v) with 1 M  $NaClO_4$  as an electrolyte was cycled at C/20 30 times at room temperature. Fig. 4b shows that the PB framework exhibits excellent capacity retention; more than 99% discharge capacity is retained after 30 cycles. Meanwhile, the cell shows a low coulombic efficiency at the initial cycling. A most probable cause for such a low coulombic efficiency is crystalline water in the PB compound. The large interstitial spaces in PB readily absorb water molecules.<sup>23</sup> It is doubtful that the water molecules can be totally removed from the lattice. Decomposition of the residual water during the charge process would cause a low efficiency. With subsequent cycles, the amount of residual water would decrease and the efficiency increase as shown in Fig. 4b.

We are left to ask why the voltage of the low-spin Fe(III)–Fe(II) couple is significantly larger in the  $KFe_2(CN)_6$  than in the other  $KMFe(CN)_6$  compounds. A larger voltage means a more stable redox energy and, therefore, a weaker Fe–C bond since the redox couple is an antibonding state with respect to the anion-p orbitals. Clearly a quantitative calculation of the effect of the strength of the M(II)–N bonding on the strength of the Fe–C bonding as well as a measure of the variation of the Fe–C bond length with changing of the M(II) cation is needed to clarify this question.

In summary, the hexacyanides with the Prussian blue structure,  $KMFe(CN)_6$ , are easily synthesized and have a low cost. We have shown that the linear  $(C\equiv N)^-$  molecule gives an M(II)–N $\equiv$ C–Fe(III) bond length that allows  $Na^+$  ions to be inserted reversibly into the empty large-ion sites. Moreover, the strong  $\pi$ -bond component in the linear  $C\equiv N$  bond enhances the crystal-field splitting  $\Delta_c$  of the  $\pi$ -bonding t and  $\sigma$ -bonding e orbitals of the 3d manifolds on the transition-metal cations, particularly for the Fe–C bond, which favors stabilization of low-spin state Fe(II):  $3d^5$  configuration in an octahedral site of C atoms whereas the Fe(II):  $3d^6$  configuration coordinated to N atoms is in a high-spin state. A flat discharge capacity of over  $70\text{ mA h g}^{-1}$  with good reversibility can be realized for a Na-ion battery using these cathodes, which warrants further investigation of exchanging  $Na^+$  for  $K^+$  in these structures.

## Notes and references

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