



## PAPER

[View Article Online](#)  
[View Journal](#) | [View Issue](#)
Cite this: *Nanoscale*, 2025, 17, 7908

# Ultrafast electron transfer at the $\text{ZnIn}_2\text{S}_4/\text{MoS}_2$ S-scheme interface for photocatalytic hydrogen evolution†

Himanshu Bhatt,<sup>a</sup> Mohammed Suleman Patel,<sup>a</sup> <sup>a</sup> Tanmay Goswami,<sup>a</sup>  
Dharmendra K. Yadav,<sup>a</sup> Atal Swathi Patra <sup>b</sup> and Hirendra N. Ghosh<sup>\*b</sup>

The performance of any photocatalyst relies on its solar harvesting and charge separation characteristics. Fabricating the S-scheme heterostructure is a proficient approach for designing next-generation photocatalysts with improved redox capabilities. Here, we integrated  $\text{ZnIn}_2\text{S}_4$  (ZIS) and  $\text{MoS}_2$  nanosheets to develop a unique S-scheme heterostructure through an *in situ* hydrothermal technique. The designed ZIS/ $\text{MoS}_2$  heterostructure showcased a 2.8 times higher photocatalytic  $\text{H}_2$  evolution rate than pristine ZIS nanosheets. The steady-state optical measurements revealed enhanced visible light absorption and reduced charge recombination in the heterostructure. Transient absorption (TA) spectroscopy revealed the interfacial electron transfer from ZIS to  $\text{MoS}_2$ . The X-ray photoelectron and electron/hole quenching TA spectroscopic measurements collectively confirmed the integration of both semiconductors in an S-scheme manner, facilitating enhanced  $\text{H}_2$  production in the case of the heterostructure. This study highlights the importance of in-depth spectroscopic investigations in advancing the photocatalytic performance of S-scheme heterostructure-based photocatalysts.

Received 30th November 2024,  
Accepted 8th February 2025

DOI: 10.1039/d4nr05043f

[rsc.li/nanoscale](https://rsc.li/nanoscale)

## Introduction

The utilization of semiconductor photocatalysis for hydrogen generation is increasingly recognized as a feasible approach to meet the higher demands of current energy consumption.<sup>1–5</sup> A significant obstacle in this domain is the inefficient spatial segregation of photogenerated electrons and holes in the photocatalysts. In this regard, the van der Waals heterostructure has been regarded as a potential candidate for overcoming this obstacle by providing vertical isolation to the charge carriers in different semiconductors.<sup>6</sup>

Following the discovery of graphene and its extraordinary properties, there has been renewed interest in the world of science for two-dimensional (2D) layered materials. This group includes metal chalcogenides, transition metal oxides, and other 2D materials.<sup>7</sup> One of the classes in this field is that of transition metal dichalcogenides or TMDCs. These are fascinating because they provide a wide range of structural and

optoelectronic properties, including high carrier mobility, tunable band positions, a large number of surface-active sites, and strong light-matter interactions.<sup>8,9</sup> These unique traits make 2D TMDCs ideal for constructing high-performance optoelectronic and photocatalytic devices.<sup>9–11</sup> Along with the strong light-matter interaction, an efficient photocatalyst should exhibit other essential properties like enhanced solar energy absorption, appropriate redox potential, slower carrier recombination and rapid carrier transport towards the photocatalytically active sites.<sup>12,13</sup> It is hard to have all these properties in a single 2D material. The construction of the van der Waals heterostructure with other 2D materials is one promising approach that utilizes the individual properties of its constituent materials and advances their photocatalytic performance. Recently, step scheme (S-scheme) heterojunctions, which provide enhanced photocatalytic activity by improving charge separation and retention of strong redox potentials, have gained significant attention in nanotechnology and materials science.<sup>14</sup> The S-scheme heterojunction offers an advanced charge transfer pathway by allowing electrons in the conduction band (CB) of one semiconductor, which is an oxidation photocatalyst (OP), to recombine with holes in the valence band (VB) of another semiconductor, called a reduction photocatalyst (RP), preserving the strong redox potentials of the respective charge carriers.<sup>15,16</sup> This unique mechanism distinguishes it from traditional type-II heterojunctions and has

<sup>a</sup>Institute of Nano Science and Technology, Knowledge City, Sector 81, SAS Nagar, Punjab-140306, India

<sup>b</sup>School of Chemical Sciences, National Institute of Science Education and Research (NISER), Bhubaneswar, Odisha 752050, India. E-mail: [hngghosh@niser.ac.in](mailto:hngghosh@niser.ac.in), [hngghosh2004@gmail.com](mailto:hngghosh2004@gmail.com)

† Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d4nr05043f>

been proven to facilitate more efficient charge separation, leading to improved photocatalytic performance. A series of S-scheme heterostructures has been utilized in photocatalytic applications.<sup>17–22</sup>

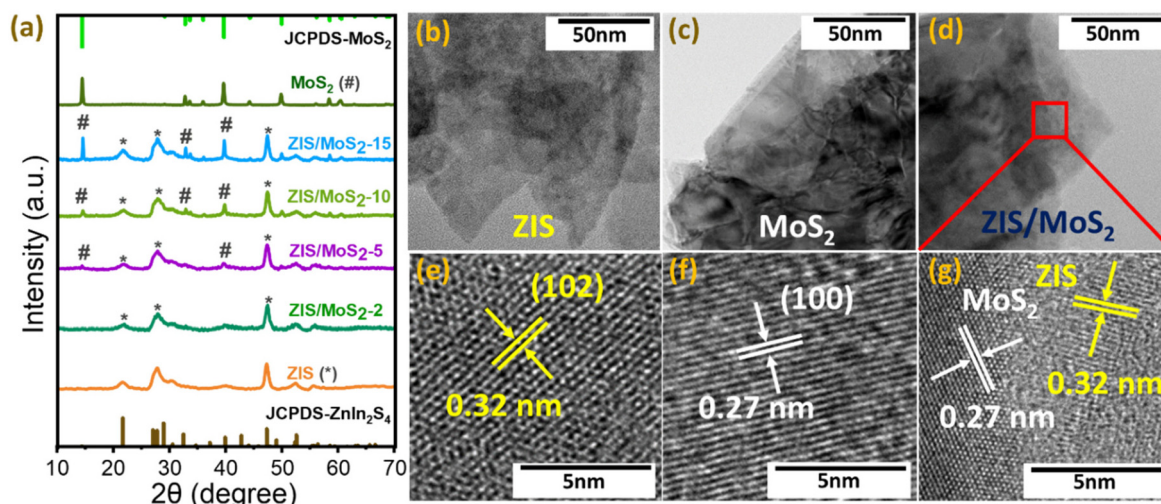
Among the materials used for constructing S-scheme heterojunctions, ZnIn<sub>2</sub>S<sub>4</sub> (ZIS) and molybdenum disulfide (MoS<sub>2</sub>) are two promising candidates due to their favourable band structures and photocatalytic properties.<sup>23,24</sup> Ternary metal chalcogenides, especially II–III<sub>2</sub>–VI<sub>4</sub> semiconductors like ZIS, are gaining recognition in photocatalysis due to their suitable band gaps, easy synthesis routes, varied catalytic functions, and cost-effectiveness. ZIS, a ternary metal chalcogenide, has gained attention for its stable physicochemical properties, tunable bandgap, and excellent visible light absorption.<sup>25,26</sup> However, despite the high photocatalytic potential of ZIS, its performance is often limited by the fast recombination of charge carriers. On the other side, MoS<sub>2</sub>, a 2D TMDC, offers high conductivity and an optimal electronic structure for photocatalytic hydrogen evolution.<sup>27</sup> By integrating these two semiconductors in an S-scheme heterojunction, synergistic effects can enhance light absorption, facilitate efficient charge separation, and boost overall photocatalytic performance. Previous studies on the ZIS/MoS<sub>2</sub> heterostructure have indicated their potential in hydrogen evolution reaction (HER) application.<sup>28,29</sup> Despite these advancements, several challenges remain, including the need for better control over the morphology of the heterojunction and the optimization of interfacial contact between the semiconductors to minimize energy losses during the charge transfer process. One promising direction for further enhancing the performance of S-scheme photocatalysts lies in using transient absorption (TA) spectroscopy to study the ultrafast dynamics of charge transfer at the semiconductor interface. Ultrafast spectroscopy can provide crucial insights into the lifetime and pathways of photogenerated charge carriers, allowing researchers to fine-tune the material's properties for maximum efficiency.<sup>30–32</sup> To date, very limited reports are available in the existing literature, where TA spectroscopy has been to investigate the responsible photophysical processes in S-scheme heterostructures for their enhanced photocatalytic performance.<sup>33,34</sup> Deng *et al.* fabricated the In<sub>2</sub>O<sub>3</sub>/Nb<sub>2</sub>O<sub>5</sub> heterostructure and performed TA spectroscopic measurements. This study revealed the generation of strong interfacial electric fields (IEFs) at the S-scheme interface, which promoted rapid charge transfer, significantly reducing recombination rates and enhancing photocatalytic CO<sub>2</sub> reduction.<sup>33</sup> Similarly, Li *et al.* designed an S-scheme heterostructure comprising organic (g-C<sub>3</sub>N<sub>4</sub>) and inorganic (CdS) semiconductors. They explored the S-scheme charge transfer pathways responsible for the improved H<sub>2</sub> production in the heterostructure by performing various experimental and theoretical measurements.<sup>34</sup> To the best of our knowledge, no report in the literature discusses the ultrafast charge carrier dynamics in the S-scheme ZIS/MoS<sub>2</sub> heterostructure. A detailed spectroscopic investigation is needed to explore the charge transfer characteristics in this unique heterojunction and its effective utilization in photocatalytic H<sub>2</sub> production.

In this work, we have provided an in-depth discussion of the charge transfer dynamics in a novel S-scheme ZIS/MoS<sub>2</sub> heterostructure using pump–probe TA spectroscopy. A simple and cost-effective hydrothermal synthetic route was followed to construct the heterostructure. TEM and HRTEM measurements demonstrated the formation of a van der Waal heterostructure between these two semiconductors. The ZIS/MoS<sub>2</sub> heterostructure displayed substantially enhanced photocatalytic H<sub>2</sub> production compared to pristine systems. Steady-state spectroscopy disclosed enhanced absorption and reduced charge recombination in the case of the heterostructure. XPS studies suggested the formation of the S-scheme interface, which exhibited the directional electric field. The underlying mechanisms of charge separation and transfer within the heterojunction were elucidated using TA spectroscopy. The ultrafast study revealed the interfacial electron transfer and confirmed the S-scheme interface. Our findings highlight the role of interface engineering and morphology control in optimizing the performance of S-scheme heterojunctions, paving the way for further progress in sustainable hydrogen production technologies.

## Results and discussion

The ZIS nanosheets and ZIS/MoS<sub>2</sub> heterostructures (ZIS/MoS<sub>2</sub>-2, ZIS/MoS<sub>2</sub>-5, ZIS/MoS<sub>2</sub>-10 and ZIS/MoS<sub>2</sub>-15) with different weight % of MoS<sub>2</sub> (2, 5, 10 and 15 wt%) were synthesized using a robust and cost-effective hydrothermal method. The pristine MoS<sub>2</sub> nanosheets were obtained using the liquid-exfoliation technique. We performed XRD to study the structural properties of all the pristine and heterostructures. Fig. 1a shows the XRD patterns of the as-synthesized ZIS, MoS<sub>2</sub>, ZIS/MoS<sub>2</sub>-2, ZIS/MoS<sub>2</sub>-5, ZIS/MoS<sub>2</sub>-10 and ZIS/MoS<sub>2</sub>-15 heterostructures. The XRD pattern of ZIS shows a good resemblance with the hexagonal crystal structure of ZIS (JCPDS 01-089-3963).<sup>13,29</sup> On the other side, the MoS<sub>2</sub> XRD pattern exhibits peaks at 2θ values of 14.5°, 32.8°, 33.6°, 35.9°, 39.6°, 44.3°, 49.9°, 58.4° and 60.5°, corresponding to the reflection planes of (002), (100), (101), (102), (103), (006), (105), (110) and (008), respectively.<sup>35,36</sup> The presence of these peaks confirms the hexagonal phase of MoS<sub>2</sub>. In the case of heterostructures, we observed the presence of the planes related to both ZIS and MoS<sub>2</sub>. However, the intensity of the MoS<sub>2</sub> planes increased with the increasing amount of MoS<sub>2</sub> in heterostructures.

Next, we studied the surface morphology of the synthesized materials using transmission electron microscopy (TEM) and high-resolution TEM (HR-TEM). Fig. 1b and c show the TEM image of ZIS and MoS<sub>2</sub>, respectively. Both the pristine systems exhibit a nanosheet-like morphology, as also observed in the field emission scanning electron microscopy (FESEM) images of both the pristine systems (Fig. S1†). As shown in Fig. 1d, the growth of the ZIS nanosheets was observed on top of MoS<sub>2</sub> nanosheets in ZIS/MoS<sub>2</sub>-10, which indicated the formation of the van der Waals heterostructure. Fig. 1(e–g) shows the HR-TEM images of the ZIS, MoS<sub>2</sub>, and ZIS/MoS<sub>2</sub> hetero-



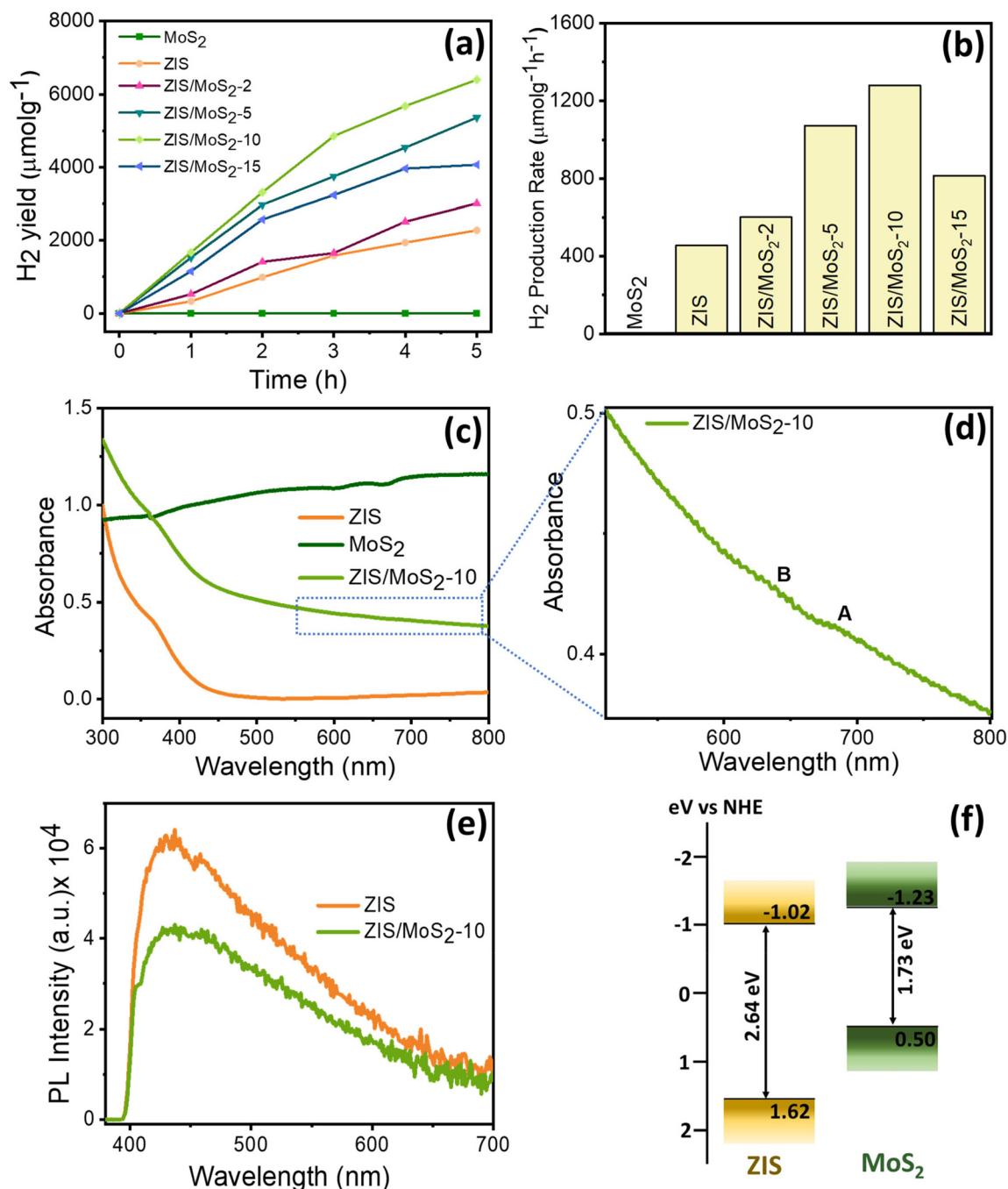
**Fig. 1** (a) XRD pattern of the as-synthesised ZIS, MoS<sub>2</sub>, ZIS/MoS<sub>2</sub>-2, ZIS/MoS<sub>2</sub>-5, ZIS/MoS<sub>2</sub>-10 and ZIS/MoS<sub>2</sub>-15 heterostructures. TEM and HRTEM images of the as-synthesized (b and e) ZIS, (c and f) MoS<sub>2</sub>, (d and g) the ZIS/MoS<sub>2</sub>-10 heterostructure.

structure. For ZIS, the interplanar spacing was found to be 0.32 nm, corresponding to the (102) plane of hexagonal ZIS. On the other hand, the HR-TEM image of MoS<sub>2</sub> shows the lattice fringes with a *d*-spacing of 0.27 nm corresponding to the (100) plane observed at a  $2\theta$  value of  $32.8^\circ$  in the XRD pattern of MoS<sub>2</sub> (Fig. 1a). The HR-TEM image of the ZIS/MoS<sub>2</sub> heterostructure (Fig. 1g) disclosed the two different lattice fringes with *d*-spacing of 0.32 nm and 0.27 nm, corresponding to the (102) and (100) planes of ZIS and MoS<sub>2</sub>, respectively, which portrays the clear interface between the ZIS and MoS<sub>2</sub> nanosheets in the heterostructure.

Furthermore, all the pristine and heterostructures were utilized in the photocatalytic reaction, in which the H<sub>2</sub> production was monitored under solar light illumination. Fig. 2a shows the photocatalytic H<sub>2</sub> yield for all the synthesized materials. The pristine MoS<sub>2</sub> did not show any photocatalytic activity, which was also observed earlier.<sup>28,29</sup> The ZIS nanosheets displayed a H<sub>2</sub> generation yield of  $2276 \mu\text{mol g}^{-1}$  under 5 h solar irradiation, resulting in a H<sub>2</sub> production rate of  $455 \mu\text{mol g}^{-1} \text{h}^{-1}$  (Fig. 2b). This H<sub>2</sub> production rate increased to  $602 \mu\text{mol g}^{-1} \text{h}^{-1}$  when 2 wt% of MoS<sub>2</sub> was integrated with ZIS nanosheets. The highest evolution rate was measured for the ZIS/MoS<sub>2</sub>-10 heterostructure as  $1280 \mu\text{mol g}^{-1} \text{h}^{-1}$ , which is nearly 2.8 times higher than that of pristine ZIS. Two other important photocatalytic parameters, solar-to-hydrogen conversion efficiency and apparent quantum efficiency (irradiation wavelength = 400 nm), were also found to be improved in the case of the ZIS/MoS<sub>2</sub>-10 heterostructure as compared to those of ZIS nanosheets. Detailed calculations are provided in the ESI.† Also, the ZIS/MoS<sub>2</sub>-10 heterostructure shows better H<sub>2</sub> evolution performance in comparison with many previously reported heterostructure-based photocatalysts, as listed in Table S3.† However, further increment of MoS<sub>2</sub> (15 wt%) decreased the evolution rate to  $813 \mu\text{mol g}^{-1} \text{h}^{-1}$ ; this is possibly due to the suppression of the ZIS active

sites after the introduction of excessive MoS<sub>2</sub>. After that, we also performed the recyclability and stability test for the ZIS/MoS<sub>2</sub>-10 heterostructure. Fig. S2a† shows the recycling H<sub>2</sub> evolution data for the ZIS/MoS<sub>2</sub>-10 heterostructure. The photocatalyst was found to be stable and retained its H<sub>2</sub> evolution performance (92%) after three cycles of 5 h reaction time. The stability of the ZIS/MoS<sub>2</sub>-10 heterostructure photocatalysts was also illustrated by XRD measurement, which was recorded after the photocatalysis (Fig. S2b†). No significant change was observed in the XRD pattern after the photocatalysis, which further confirmed the photostability of the photocatalyst.

To investigate the reason behind the superior photocatalytic performance of the ZIS/MoS<sub>2</sub>-10 heterostructure as compared to those of their pristine systems, we conducted an in-depth spectroscopic investigation using a wide range of spectroscopic techniques, including steady-state absorption and photoluminescence (PL) spectroscopy, X-ray photoelectron spectroscopy (XPS) and femtosecond transient absorption (TA) spectroscopy. Fig. 2c displays the steady-state absorption spectra of ZIS, MoS<sub>2</sub> and the ZIS/MoS<sub>2</sub>-10 heterostructure. The ZIS absorption spectra exhibit a peak maximum of around 390 nm, which originates from the higher energy transition.<sup>13,26</sup> In the heterostructure, the absorption was observed to be enhanced as compared to that of pristine ZIS due to the presence of MoS<sub>2</sub>. The characteristic A and B excitonic peaks were also observed in the bare MoS<sub>2</sub> flakes and heterostructure, as shown in Fig. 2d. A and B excitons originate from the transition between the conduction band minimum and spin-split valence band maxima at the K point in the momentum space.<sup>37,38</sup> The enhanced visible range absorption is beneficial for the heterostructure in photocatalytic hydrogen production. The PL spectra portray the recombination characteristics of charge carriers in the bare ZIS and ZIS/MoS<sub>2</sub>-10 heterostructure after photoexcitation at



**Fig. 2** H<sub>2</sub> production (a) yields and (b) rates of ZIS, MoS<sub>2</sub>, ZIS/MoS<sub>2</sub>-2, ZIS/MoS<sub>2</sub>-5, ZIS/MoS<sub>2</sub>-10 and ZIS/MoS<sub>2</sub>-15 heterostructures. (c) Steady-state absorption spectra of ZIS, MoS<sub>2</sub>, and ZIS/MoS<sub>2</sub>-10. (d) Enlarged view for ZIS/MoS<sub>2</sub>-10. (e) Photoluminescence spectra of ZIS and ZIS/MoS<sub>2</sub>-10 after photoexcitation at 350 nm. (f) Band position estimation for ZIS and MoS<sub>2</sub>.

350 nm, as shown in Fig. 2e. The ZIS spectra exhibit a broad PL emission extending beyond 700 nm with a peak maximum of around 435 nm. The 435 nm peak originated from the direct recombination of the conduction band electrons and valence band holes, whereas the extended broad emission in the red region is attributed to the trap state mediated recombination.<sup>26</sup> The PL intensity decreased in the heterostructure, which suggests the reduction in the recombination between

the CB electrons and VB holes of ZIS, possibly due to the interfacial charge transfer in the heterostructure. The band gaps of ZIS and MoS<sub>2</sub> were calculated as 2.64 and 1.73 eV, respectively (Fig. S3b and S3c†), with the help of DRS spectra, whereas the XPS VB spectra determine the respective VB of both semiconductors as 1.62 and 0.50 eV, respectively, as shown in Fig. S3a.† By combining these two pieces of information, the respective band positions were estimated and are shown in



Fig. 2f. ZIS and MoS<sub>2</sub> show a staggered type band alignment in the heterostructure.

Next, we conducted XPS studies to explore the surface chemical states of the synthesized bare and heterostructures. The high-resolution XPS (HR-XPS) spectra (Fig. 3) confirm the presence of Zn, In, S, and Mo in the pristine and heterostructures. For ZIS, the HR-XPS spectrum of Zn 2p (Fig. 3a) exhibits two peaks at 1021.9 and 1045.0 eV, corresponding to 2p<sub>3/2</sub> and 2p<sub>1/2</sub>, respectively. In Fig. 3b, two peaks at 444.9 and 452.4 eV are attributed to the 3d<sub>5/2</sub> and 3d<sub>3/2</sub> of In 3d.<sup>29</sup> The S 2p spectrum for ZIS (Fig. 3d) showed two characteristic peaks at 161.6 eV (2p<sub>3/2</sub>) and 162.7 eV (2p<sub>1/2</sub>),<sup>28</sup> whereas for MoS<sub>2</sub>, S 2p peaks were observed at 162.0 and 163.1 eV. Additionally, two peaks of Mo 3d<sub>5/2</sub> and Mo 3d<sub>3/2</sub> were found at 228.5 and 231.6 eV, respectively (Fig. 3c), for pristine MoS<sub>2</sub>.<sup>39</sup> It is worth noting that the peaks related to ZIS (Zn 2p, In 3d and S 2p) were observed to be shifted to lower binding energy in the heterostructure. Conversely, Mo 3d of ZIS/MoS<sub>2</sub>-10 experienced a positive shift to higher binding than pristine MoS<sub>2</sub>. These outcomes are attributed to the electron migration from MoS<sub>2</sub> to ZIS when they are integrated to form a heterostructure, resulting in a generation of the interfacial electric field at the interface with the direction from MoS<sub>2</sub> to ZIS. The above results indicate a formation of the S-scheme heterojunction between both nanosheets.<sup>20,33</sup>

To validate the S-scheme formation and investigate the interfacial photophysical processes involved in the superior photocatalytic performance of the ZIS/MoS<sub>2</sub> heterostructure, we employed state-of-the-art Femtosecond TA spectroscopy for pristine nanosheets and their heterostructure. Fig. 4a shows the TA spectra of ZIS as a function of different pump-probe time delays following the photoexcitation at 370 nm. The spectra exhibit a strong photo-induced bleach signal in the spectral range of 400–600 nm, with the peak maxima around 420 nm. In our previous work, we have already assigned this bleach maxima to the transition from the CBM to the VB states.<sup>13,26</sup> In addition to a strong bleach maximum, we observed a weak and broad bleach signal in the red region of the spectra, which originated from the mid-band gap trap state transitions.<sup>26</sup> The TA kinetics of ZIS, portrayed in Fig. 4c, reflects the time-dependent carrier population at the corresponding state monitored at the bleach maxima (~420 nm). The TA kinetics was fitted using a multiexponential function with one growth ( $\tau_g = 0.2$  ps) component and three decay components ( $\tau_1 = 5.6$  ps,  $\tau_2 = 52.4$  ps and  $\tau_3 \geq 1$  ns), which are tabulated in Table S4.† The growth of the bleach signal signifies the time involved in the cooling of the hot carrier to the particular energy state that is being monitored. Meanwhile, the decay time constants correspond to the trapping and recombination processes. The first two decay time constants ( $\tau_1 =$

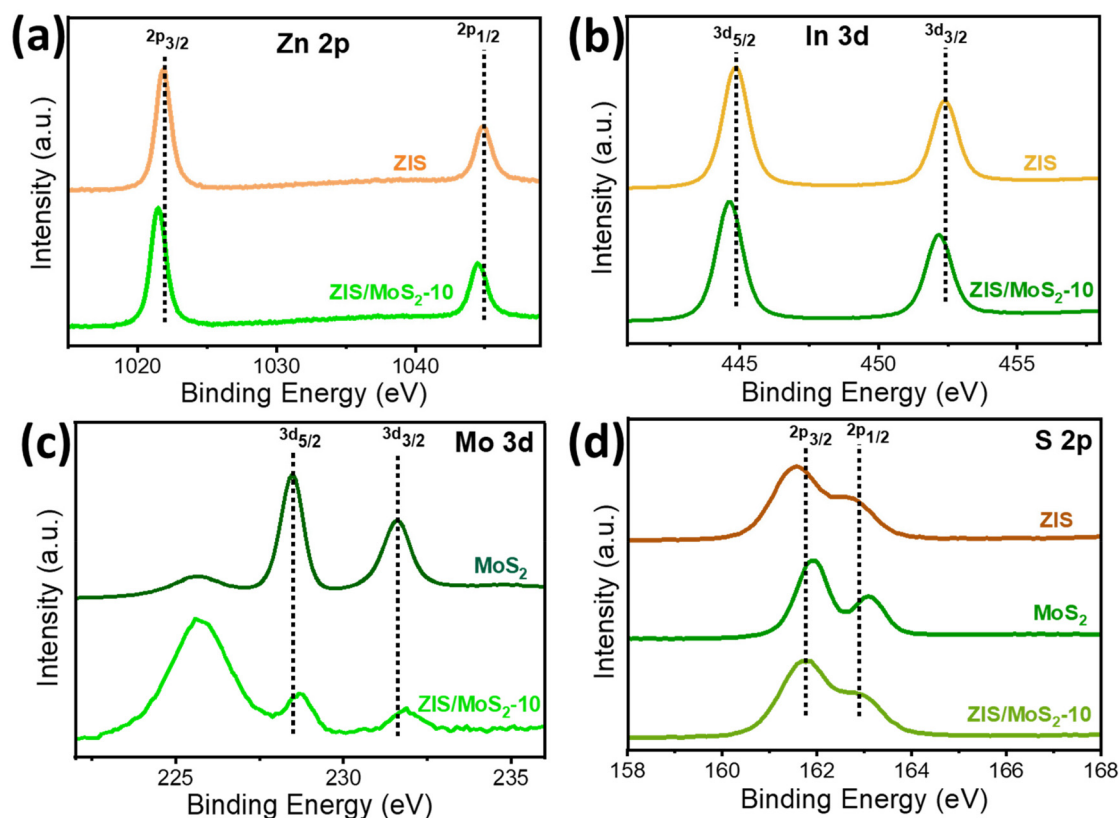
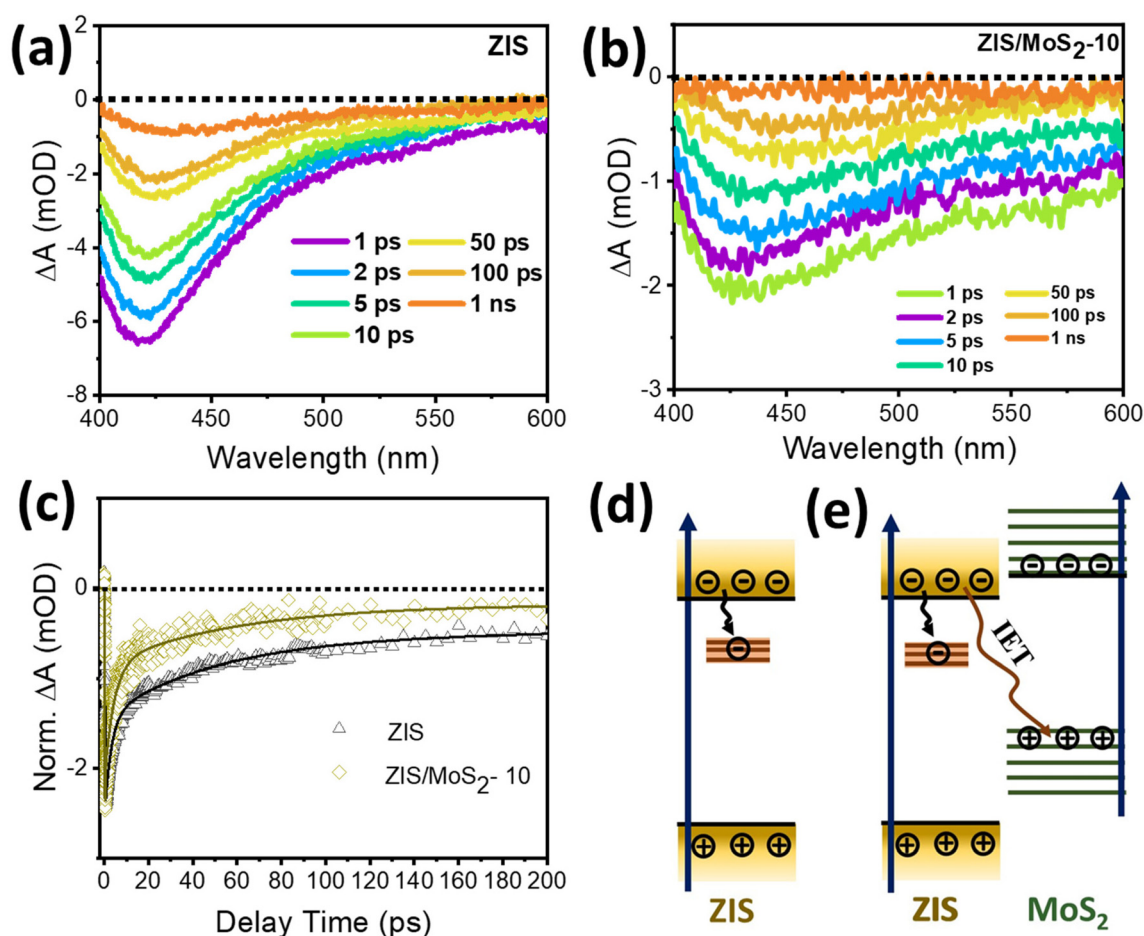


Fig. 3 High-resolution XPS spectra of (a) Zn 2p and (b) In 3d for ZIS and ZIS/MoS<sub>2</sub>-10 and (c) Mo 3d for MoS<sub>2</sub> and ZIS/MoS<sub>2</sub>-10. (d) S 2p for ZIS, MoS<sub>2</sub> and ZIS/MoS<sub>2</sub>-10.



**Fig. 4** TA spectra of (a) ZIS nanosheets and (b) the ZIS/MoS<sub>2</sub>-10 heterostructure after photoexcitation at 370 nm. (c) Comparative TA kinetics of ZIS and ZIS/MoS<sub>2</sub>-10, monitored at 420 nm. Ultrafast photophysical processes, such as electron trapping and interfacial electron transfer (IET) processes in (d) bare ZIS and (e) the ZIS/MoS<sub>2</sub> heterostructure.

5.6 ps and  $\tau_2 = 52.4$  ps) were assigned to the trapping and de-trapping of the charge carriers in the mid-band gap trap states, respectively (Fig. 4d). The longest-time component ( $\tau_3 \geq 1$  ns) is attributed to the electron-hole recombination. The role of individual carriers (electrons and holes) in the relaxation dynamics of the semiconductors is usually determined with the help of electron/hole quenching experiments.<sup>33,40</sup> In our study, we used benzoquinone (BQ) and TEOA as electron and hole quenchers, respectively, to find out the contribution of the respective charge carriers in the transient bleach signals of the pristine systems. As shown in Fig. S2a and S2b,† the comparative kinetic profiles monitored at bleach maxima of ZIS validate the domination of electrons in the ZIS bleach signal, as the bleach signal completely quenched around 100 ps with the addition of BQ, whereas the kinetics was observed to be similar in the case of TEOA.

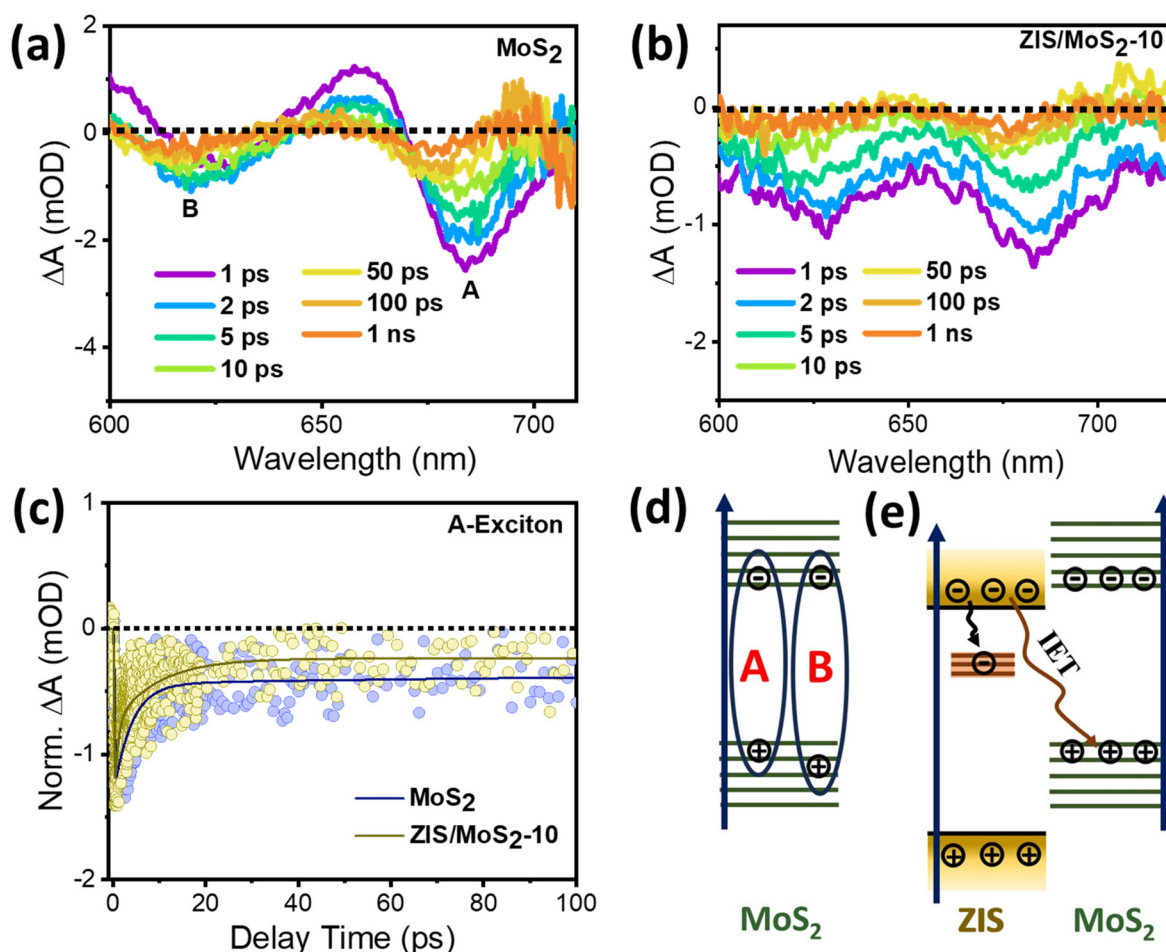
Next, we carried out our transient spectroscopic investigation for the ZIS/MoS<sub>2</sub>-10 heterostructure to understand the interfacial charge transfer phenomena. It is worth mentioning that we maintained similar experimental conditions throughout the TA measurements. The TA spectra of the hetero-

structure (Fig. 4b), dominated by a broad bleach signal, peaking around  $\sim 423$  nm, show a spectral similarity to the pristine ZIS nanosheets. Fig. 4c shows the comparative TA kinetics of ZIS before and after the integration of MoS<sub>2</sub> nanosheets. The kinetic profile was observed to be accelerated in the case of ZIS/MoS<sub>2</sub>-10, with reduced decay time constants ( $\tau_1 = 2.7$  ps and  $\tau_2 = 30.2$  ps) as compared to that of pristine ZIS ( $\tau_1 = 5.6$  ps and  $\tau_2 = 52.4$  ps), which affirms the interfacial charge transfer from the ZIS to MoS<sub>2</sub> side. This charge transfer reduced the charge recombination in the ZIS side as the contribution of the  $\tau_3$  reduced by half (Table S4†) in the case of the heterostructure (25% for ZIS and 12.5% for ZIS/MoS<sub>2</sub>-10). As the bleach signal is dominantly contributed by the electrons, the faster TA kinetics can be evidently assigned to the thermodynamically feasible interfacial electron transfer (IET) from the ZIS CB to MoS<sub>2</sub> VB, as shown schematically in Fig. 4e. This is an important characteristic evident for the S scheme interface.<sup>33,41</sup> In the case of the S-scheme heterojunction, the generated directional electric field from MoS<sub>2</sub> to ZIS allows electron migration in the opposite direction, *i.e.*, from ZIS to MoS<sub>2</sub> after the light irradiation.<sup>19</sup>

Furthermore, we conducted TA spectroscopy for another constituent of the heterostructure, *i.e.*, MoS<sub>2</sub>, in the spectral range of 600–720 nm. Fig. 5a shows the TA spectra of pristine MoS<sub>2</sub> nanosheets after 370 nm photoexcitation. The spectrum comprises two bleach signals with maxima positioned at ~684 nm and ~620 nm, assigned to the A and B excitonic peaks, respectively. The formation of these excitons in MoS<sub>2</sub> is schematically illustrated in Fig. 5d. These excitonic features retained their spectral identity in the case of ZIS/MoS<sub>2</sub>-10, as portrayed in Fig. 5b. The TA decay kinetics of A excitons for bare MoS<sub>2</sub> and the heterostructure (Fig. 5c) were fitted multi-exponentially in three-time constants, as listed in Table S4† ( $\tau_1 = 3.02$  ps,  $\tau_2 = 69.9$  ps and  $\tau_3 \geq 1$  ns for MoS<sub>2</sub>,  $\tau_1 = 0.94$  ps,  $\tau_2 = 15.7$  ps and  $\tau_3 \geq 1$  ns for ZIS/MoS<sub>2</sub>-10). Faster kinetics was observed for the heterostructure compared to individual MoS<sub>2</sub> kinetics, which was expected because of the participation of MoS<sub>2</sub> holes in the recombination processes with the interfacial transferred electrons from ZIS (Fig. 5e). The stronger influence of the TEOA (hole quencher) as compared to that of BQ (electron quencher) in the kinetic profile of pristine MoS<sub>2</sub>

(Fig. S4c and S4d†) also supports the above findings, which suggest the recombination of the ZIS CB electrons and MoS<sub>2</sub> VB holes due to the IET process. The involvement of the MoS<sub>2</sub> holes in the S-scheme recombination process suppressed the recombination on the MoS<sub>2</sub> side, as indicated by the reduced proportion of the longest decay component related to charge recombination ( $\tau_3 \geq 1$  ns) in the heterostructure (16.7%) as compared to that of pristine MoS<sub>2</sub> (33.4%). Consequently, the electrons of the MoS<sub>2</sub> CB are isolated, and they actively participate in the photocatalytic H<sub>2</sub> reduction process.

Fig. 6 shows a schematic representation of the proposed photocatalytic H<sub>2</sub> production mechanism based on the spectroscopic investigation. In the ZIS/MoS<sub>2</sub> heterostructure, the charge carriers (electrons and holes) were generated in both the semiconductors after solar irradiation. Subsequently, the CB electrons of ZIS migrated towards the VB of MoS<sub>2</sub> due to the generation of the interfacial electric field in the S-scheme manner (IET) and recombined with the VB holes of MoS<sub>2</sub>. This interfacial recombination isolated the CB electrons of MoS<sub>2</sub> and VB holes of ZIS. The ZIS holes interacted with a sacrificial



**Fig. 5** TA spectra of (a) MoS<sub>2</sub> and (b) the ZIS/MoS<sub>2</sub>-10 heterostructure after photoexcitation at 370 nm. (c) Comparative TA kinetics of MoS<sub>2</sub> and ZIS/MoS<sub>2</sub>-10, monitored at the A excitonic position (687 nm). Ultrafast photophysical processes like the formation of A and B excitons and IET in (d) MoS<sub>2</sub> and (e) the ZIS/MoS<sub>2</sub> heterostructure.

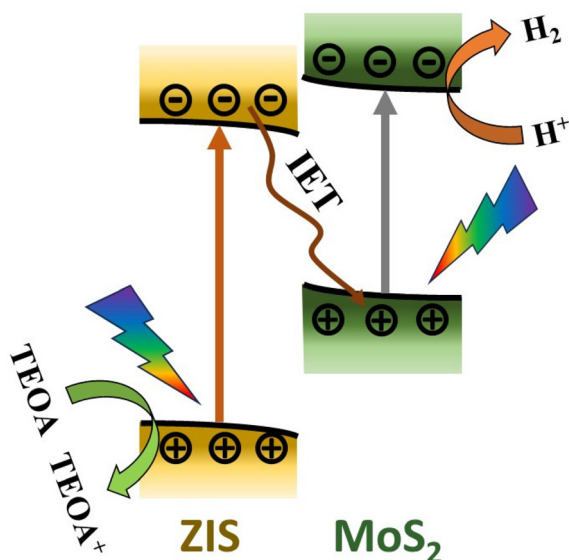


Fig. 6 Schematic representation of the plausible mechanism of the photocatalytic hydrogen production in the ZIS/MoS<sub>2</sub> heterostructure.

agent, TEOA. In contrast, the electrons of the MoS<sub>2</sub> CB, having a higher reduction potential, efficiently participated in the photocatalytic reduction process and resulted in a higher H<sub>2</sub> production for the ZIS/MoS<sub>2</sub> heterostructure in comparison with both of its constituents solely.

## Conclusion

In summary, we synthesized a unique S-scheme heterostructure between ZIS and MoS<sub>2</sub> using a hydrothermal technique. The TEM measurements confirmed the formation of a van der Waals heterostructure between these two nanosheets. Under solar energy irradiation, the hydrogen evolution rate through photocatalysis in the ZIS/MoS<sub>2</sub> heterostructure was quantified at 1280  $\mu\text{mol g}^{-1} \text{h}^{-1}$ , showcasing an increase of almost 2.8 times in comparison with the hydrogen evolution rate of pristine ZIS nanosheets. Steady-state absorption spectroscopy showed enhanced absorption in the visible range for the heterostructure. The lower PL intensity of ZIS/MoS<sub>2</sub> is attributed to the lesser recombination in the heterostructure. XPS indicated the creation of a directional electric field from MoS<sub>2</sub> to ZIS owing to the S scheme interface characteristics. We further conducted ultrafast TA measurements to investigate the interfacial charge transfer between these two 2D materials. TA study demonstrated the faster decay kinetics for the heterostructure than for both the pristine systems. The quenching experiment revealed the dominance of electrons and holes in the ZIS and MoS<sub>2</sub> bleach signals, respectively. TA analysis concluded that the interfacial electron transfer from the conduction band of ZIS to MoS<sub>2</sub> enhanced the S-scheme charge separation, which resulted in higher photocatalytic activity in the heterostructure. Overall, this work provides a comprehensive spectroscopic understanding of the interfacial charge carrier

dynamics in a 2D/2D ZIS/MoS<sub>2</sub> heterostructure. This detailed exploration would be beneficial in designing state-of-the-art van der Waal's heterostructure-based S-scheme photocatalysts and advancing their utilization for industrial H<sub>2</sub> production.

## Data availability

The data supporting this article have been included in the ESI.† It contains synthesis techniques, instrumental details, DRS spectra, XPS valence band spectra, AQE and STH calculations, a performance comparison table, and transient absorption spectra and dynamics.

## Conflicts of interest

The authors declare no conflicts of interest.

## Acknowledgements

H. B., M. S. P. and D. K. Y. thank the Institute of Nano Science and Technology (INST), Mohali, India, for providing the research fellowship. T. G. acknowledges the CSIR (09/1129 (0009)/2017-EMR-I), India, for providing fellowship during his research tenure. A. S. P. thanks NISER Bhubaneswar for the research fellowship. H. N. G. is grateful to DST, Govt. of India, for the J C Bose Fellowship (JCB/2018/000047) and SERB/DST project (CRG/2019/000938). The authors acknowledge INST, Mohali, India, for providing instrumental facilities and supporting this research work.

## References

- 1 Z. Yan, K. Yin, M. Xu, N. Fang, W. Yu, Y. Chu and S. Shu, *Chem. Eng. J.*, 2023, **472**, 145066.
- 2 B. K. Patra, S. Khilari, D. Pradhan and N. Pradhan, *Chem. Mater.*, 2016, **28**, 4358–4366.
- 3 Y. Liu, W. Yang, Q. Chen, Z. Xie and T. Lian, *J. Chem. Phys.*, 2023, **159**, 104706.
- 4 Y. Magazov, V. Kudryashov, K. Moldabekov, M. Amze, A. Nurmanova, A. Aliyev and N. Nuraje, *Commun. Mater.*, 2024, **5**, 126.
- 5 F. Qureshi and M. Tahir, *Int. J. Hydrogen Energy*, 2024, **69**, 760–776.
- 6 X. Chen, W. Pan, R. Guo, X. Hu, Z. Bi and J. Wang, *J. Mater. Chem. A*, 2022, **10**, 7604–7625.
- 7 M. Zeng, Y. Xiao, J. Liu, K. Yang and L. Fu, *Chem. Rev.*, 2018, **118**, 6236–6296.
- 8 Z. Hu, Z. Wu, C. Han, J. He, Z. Ni and W. Chen, *Chem. Soc. Rev.*, 2018, **47**, 3100–3128.
- 9 R. Yang, Y. Fan, Y. Zhang, L. Mei, R. Zhu, J. Qin, J. Hu, Z. Chen, Y. Hau Ng, D. Voiry, S. Li, Q. Lu, Q. Wang, J. C. Yu and Z. Zeng, *Angew. Chem., Int. Ed.*, 2023, **62**, e202218016.
- 10 Q. H. Wang, K. Kalantar-Zadeh, A. Kis, J. N. Coleman and M. S. Strano, *Nat. Nanotechnol.*, 2012, **7**, 699–712.



- 11 D. Monga, S. Sharma, N. P. Shetti, S. Basu, K. R. Reddy and T. M. Aminabhavi, *Mater. Today Chem.*, 2021, **19**, 100399.
- 12 S. Saeed, E. M. L. D. de Jong, K. Dohnalova and T. Gregorkiewicz, *Nat. Commun.*, 2014, **5**, 4665.
- 13 H. Bhatt, T. Goswami, D. K. Yadav, N. Ghorai, A. Shukla, G. Kaur, A. Kaur and H. N. Ghosh, *J. Phys. Chem. Lett.*, 2021, **12**, 11865–11872.
- 14 B. Zhu, J. Sun, Y. Zhao, L. Zhang and J. Yu, *Adv. Mater.*, 2024, **36**, 2310600.
- 15 L. Zhang, J. Zhang, H. Yu and J. Yu, *Adv. Mater.*, 2022, **34**, 2107668.
- 16 J. Wang, R. Pan, S. Yan, R. Wang, X. Niu, Q. Hao, J. Ye, Y. Wu and H. Ying Yang, *Chem. Eng. J.*, 2023, **463**, 142489.
- 17 H. Tran Huu, M. D. N. Thi, V. P. Nguyen, L. N. Thi, T. T. T. Phan, Q. D. Hoang, H. H. Luc, S. J. Kim and V. Vo, *Sci. Rep.*, 2021, **11**, 14787.
- 18 A. Xu, Y. Zhang, H. Fan, X. Liu, F. Wang, X. Qu, L. Yang, X. Li, J. Cao and M. Wei, *ACS Appl. Nano Mater.*, 2024, **7**, 3488–3498.
- 19 X. Luan, Z. Yu, J. Zi, F. Gao and Z. Lian, *Adv. Funct. Mater.*, 2023, **33**, 2304259.
- 20 X. Hu, Y. Xu, S. Tang, W. Shi, X. Wang, Y.-X. Yu and W.-D. Zhang, *Small*, 2024, **20**, 2306563.
- 21 F. Zhao, X. Yao, Y. Zhao, J. Yu, J. Dong, X. Liu, J. Cao, D. Zhang and X. Pu, *J. Colloid Interface Sci.*, 2024, **675**, 471–480.
- 22 Q. Liu, X. Yao, X. Zhao, F. Zhang, D. Zhang, J. Liu, P. Cai and X. Pu, *Surf. Interfaces*, 2024, **54**, 105213.
- 23 N. Li, J. Ma, W. Wang, Q. Chang, L. Liu, C. Hao, H. Zhang, H. Zhang, S. Hu and S. Wang, *J. Colloid Interface Sci.*, 2024, **676**, 496–505.
- 24 D. Zhang, D. Zhang, F. Zhao, Y. Zhao, H. Li, J. Liu, X.-Y. Ji, X. Pu and H. Zhang, *J. Mater. Chem. A*, 2024, **12**, 33546–33558.
- 25 W.-K. Chong, B.-J. Ng, Y. J. Lee, L.-L. Tan, L. K. Putri, J. Low, A. R. Mohamed and S.-P. Chai, *Nat. Commun.*, 2023, **14**, 7676.
- 26 T. Goswami, D. K. Yadav, H. Bhatt, G. Kaur, A. Shukla, K. J. Babu and H. N. Ghosh, *J. Phys. Chem. Lett.*, 2021, **12**, 5000–5008.
- 27 Z. Li, X. Meng and Z. Zhang, *J. Photochem. Photobiol. C*, 2018, **35**, 39–55.
- 28 Z. Zhang, L. Huang, J. Zhang, F. Wang, Y. Xie, X. Shang, Y. Gu, H. Zhao and X. Wang, *Appl. Catal., B*, 2018, **233**, 112–119.
- 29 W. Li, Z. Lin and G. Yang, *Nanoscale*, 2017, **9**, 18290–18298.
- 30 K. L. Corp and C. W. Schlenker, *J. Am. Chem. Soc.*, 2017, **139**, 7904–7912.
- 31 T. Goswami, H. Bhatt, D. K. Yadav and H. N. Ghosh, *J. Phys. Chem. B*, 2022, **126**, 572–580.
- 32 C. Feng, T. Bo, P. Maity, S. Zuo, W. Zhou, K.-W. Huang, O. F. Mohammed and H. Zhang, *Adv. Funct. Mater.*, 2024, **34**, 2309761.
- 33 X. Deng, J. Zhang, K. Qi, G. Liang, F. Xu and J. Yu, *Nat. Commun.*, 2024, **15**, 4807.
- 34 R. Li, H. Li, X. Zhang, B. Liu, B. Wu, B. Zhu, J. Yu, G. Liu, L. Zheng and Q. Zeng, *Adv. Funct. Mater.*, 2024, **34**, 2402797.
- 35 T. Goswami, H. Bhatt, D. K. Yadav, R. Saha, K. J. Babu and H. N. Ghosh, *J. Chem. Phys.*, 2022, **156**, 034704.
- 36 A. Kaushik, J. Singh, R. Soni and J. P. Singh, *ACS Appl. Nano Mater.*, 2023, **6**, 9236–9246.
- 37 L. Wang, Z. Wang, H.-Y. Wang, G. Grinblat, Y.-L. Huang, D. Wang, X.-H. Ye, X.-B. Li, Q. Bao, A.-S. Wee, S. A. Maier, Q.-D. Chen, M.-L. Zhong, C.-W. Qiu and H.-B. Sun, *Nat. Commun.*, 2017, **8**, 13906.
- 38 T. Goswami, H. Bhatt, K. J. Babu, G. Kaur, N. Ghorai and H. N. Ghosh, *J. Phys. Chem. Lett.*, 2021, **12**, 6526–6534.
- 39 W. Yan, Y. Xu, S. Hao, Z. He, L. Wang, Q. Wei, J. Xu and H. Tang, *Inorg. Chem.*, 2022, **61**, 4725–4734.
- 40 S. Kaniyankandy, S. N. Achary, S. Rawalekar and H. N. Ghosh, *J. Phys. Chem. C*, 2011, **115**, 19110–19116.
- 41 M. Zhu, Z. Sun, M. Fujitsuka and T. Majima, *Angew. Chem., Int. Ed.*, 2018, **57**, 2160–2164.