



Cite this: *Nanoscale*, 2025, **17**, 3152

Layer-number-dependent photoswitchability in 2D MoS₂-diarylethene hybrids†

Sewon Park,^a Jaehoon Ji,^a Srajan Pillai,^b Henry Fischer,^b Jean Rouillon,^c Carlos Benitez-Martin,^d Joakim Andréasson,^c Jeong Ho You^b and Jong Hyun Choi^a                                        <img alt="ORCID iD icon" data-bbox

dates for light-controlled charge transport at the interface between these molecules and TMDs. When exposed to light, photochromic molecules undergo isomerization to another structural form. In case of DAE, isomerization between the open and the closed isomeric forms causes shifts in their highest occupied molecular orbital (HOMO) and lowest unoccupied molecular orbital (LUMO) energy levels. Such photo-induced shifting can be exploited for the modulation of TMDs. Previous studies employed DAE to form hybrid heterostructures for various purposes; for example, on/off photoswitching effects to modulate monolayer MoSe_2 devices²⁵ and few-layer WSe_2 FETs.²⁶ Similarly, attempts were made to test photoresponsive MoS_2 -based FET systems using spiropyran²⁷ and azobenzene.²⁸ The charge transfer interaction at the hybrid heterointerfaces depends on TMD layer numbers as well as photochromic materials. Despite previous efforts, the fundamental understanding of layer-number-dependent photoswitchability in hybrid TMD heterostructures remains elusive.

In this study, we utilized DAE molecules to investigate photoswitchable charge transfer behaviors for various MoS_2 layer numbers from monolayer (1L) to trilayer (3L). The different layer number of MoS_2 flakes leads to distinct conduction band minimum (CBM) and valence band maximum (VBM) energies. We measured the PL emissions of MoS_2 after DAE isomerization using UV and visible light irradiation and observed distinct behaviors from the 1L, 2L, and 3L MoS_2 -DAE. Conductive atomic force microscopy (C-AFM) was also performed to evaluate light-induced changes in current in the 1L, 2L, and 3L hybrid systems. Our observation indicates monolayer MoS_2 interacts with DAE differently than bilayer or trilayer. To elucidate underlying mechanisms, we carried out DFT calculations to investigate the layer-number-dependent CBM and VBM energies of MoS_2 against the LUMO and HOMO levels of the open and the closed DAE isomers. We identified a hole transfer mechanism that modulated the properties of 1L MoS_2 , but not 2L or 3L, due to the energy level alignment. This work provides critical insights on the thickness-dependent characteristics of MoS_2 -organic heterojunctions, laying the groundwork for advanced photoswitchable hybrid devices.

Results and discussion

Fig. 1 illustrates the hybrid structures made of DAE and MoS_2 layers. DAE molecules are reversibly isomerized from the open isomeric form to the closed isomeric form (with UV light) and back (with visible light). We hypothesize that the accompanying changes in chemical structure may imply that the two forms interact differently with MoS_2 depending on the number of layers, as 1L to 3L MoS_2 exhibit significant variations in their optoelectronic properties.^{29–32} These variations affect the mechanisms and types of charge transport behaviors. To investigate the optical properties, we prepared MoS_2 flakes on SiO_2 /Si substrates by mechanical exfoliation from bulk crystals. The

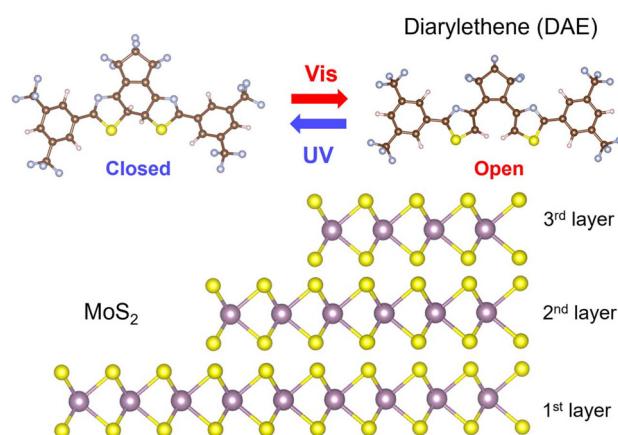


Fig. 1 Schematic illustrating layer-number-dependent charge transfer interactions between MoS_2 and DAE. UV and visible light trigger isomerization between the closed and the open form. UV irradiation results in a π -conjugated closed-form photoisomer, while visible light produces a cross-conjugated open-form. The photoisomerization changes the HOMO and LUMO levels of DAE. 1L, 2L, and 3L MoS_2 can interact differently with DAE, leading to layer-number-dependent photoswitchability.

sorting of the MoS_2 flakes into 1L to 3L layers was conducted using an optical microscope and reconfirmed *via* AFM height measurements. The samples were then annealed under Ar gas at 250 °C for an hour. After the annealing process, DAE molecules were deposited onto the MoS_2 flakes, and the samples were examined again. Fig. 2a shows a bright-field image of 1L MoS_2 , with its boundary marked by dotted lines. To confirm the thickness of the DAE layer, the height profile was measured using AFM along the dotted line in Fig. 2b. The 1L MoS_2 has a thickness of approximately 0.8 nm, while the DAE layer is about 2 nm thick. We examined a substantial number of samples (see also Fig. S1†), confirming highly consistent deposition of the organic layer with a uniform thickness.¹⁹

We measured PL to probe interfacial charge transfer behaviors in the MoS_2 -DAE hybrids using a confocal Raman microscope. To explore wavelength-dependent photoswitchability, we initially exposed the samples to visible light (using a 530 nm long-pass filter) for 10 minutes to trigger DAE isomerization to the open isomeric form. After the exposure, we recorded the PL spectra using a 633 nm HeNe laser for excitation. Note that DAE molecules exhibit minimal absorbance at 633 nm.²⁵ Subsequently, we irradiated the samples with UV light (~312 nm) for 2 minutes to switch it to the closed form and compared the resulting PL intensities. The alternating UV-visible light exposure was performed three times, and the results are presented in Fig. 2d–f (also see Fig. S2†). The sharp Raman peak at ~650 nm indicates the fingerprint of A_{1g} band in MoS_2 , which is consistent with previous reports.^{33,34} Comparing the PL characteristics of 1L, 2L, and 3L MoS_2 , distinct patterns emerge. The 1L MoS_2 -DAE sample shows a distinct PL peak at approximately 660 nm, indicative of a direct bandgap. The PL peaks of 2L and 3L hybrids shift to about 670 and 680 nm, respectively. These layer-number-dependent PL



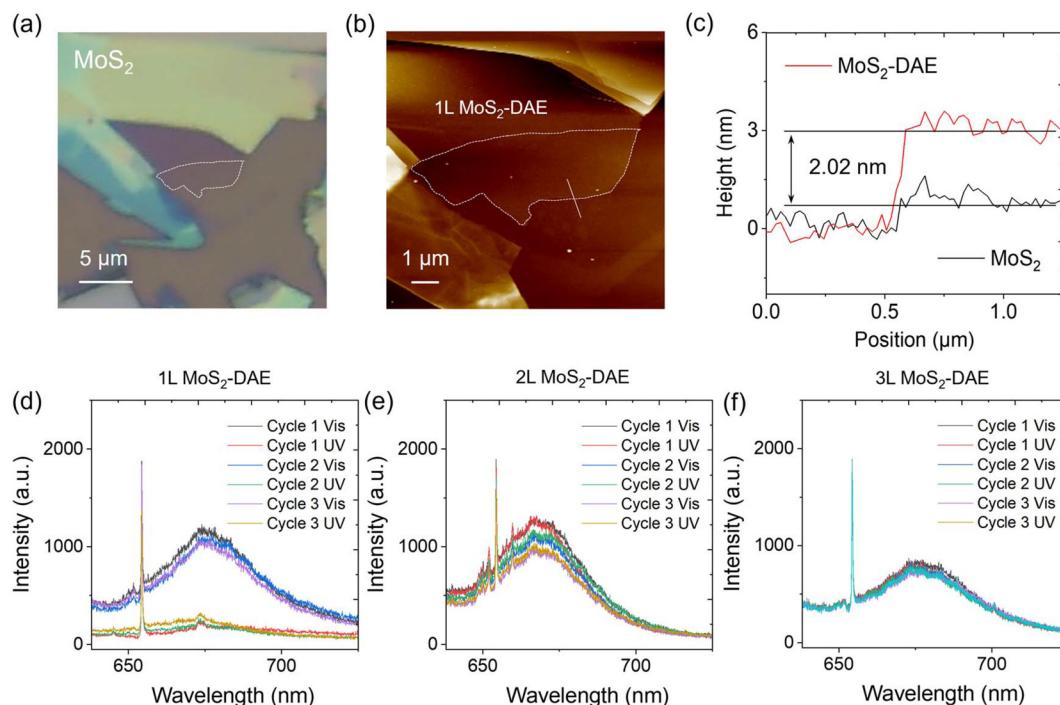


Fig. 2 (a) Optical and (b) AFM images of 1L MoS₂ indicated by dotted lines. (c) The AFM height profile shows that the thickness of the DAE layer is approximately 2 nm, while the monolayer MoS₂ has a thickness of about 0.8 nm. (d)–(f) PL intensity of 1L, 2L, and 3L MoS₂-DAE samples measured under alternating UV-visible irradiation. A 633 nm HeNe laser was used for excitation. (d) 1L MoS₂ shows a drastic PL quenching after UV irradiation. However, the signal recovers fully under visible light. This photo-modulation is repeated three times. In contrast, no significant PL quenching is observed from the hybrids with (e) 2L and (f) 3L MoS₂.

emissions are in good agreement with the measurements in literature.^{35,36} Fig. 2d shows a substantial PL quenching in 1L MoS₂-DAE after UV irradiations. The emission signals are recovered after exposure to visible light. This behavior must be associated with DAE's isomeric forms since 1L MoS₂ without a DAE layer does not show any significant differences between UV and visible irradiation (Fig. S3a†). It is attributed to photo-induced charge transfer at the hybrid interfaces that suppresses excitonic recombination of electron–hole pairs in MoS₂. Three consecutive cycles of visible and UV exposure result in consistent PL quenching and recovery, in line with what is expected from robust photoswitching (no significant degradation in performance due to undesired formation of irreversible photo-products).^{37,38} Fig. 2e and f displays no significant quenching effect in both 2L and 3L MoS₂-DAE samples after each cycle. The experimental results suggest that the 1L MoS₂-DAE hybrid experiences notable charge transfer interactions, while there are no such interactions between DAE molecules and 2L or 3L MoS₂.

Next, bulk MoS₂ crystals were mechanically exfoliated onto ITO glass to investigate the hybrids using C-AFM (Fig. S4†). The measurements were conducted by placing the samples between the ITO and Pt-Ir probe while applying a DC bias voltage of ± 2 V. Fig. 3 and S5† present current *vs.* bias voltage plots of 1L–3L MoS₂-DAE hybrids after light exposure. In Fig. 3a, 1L MoS₂ with the closed-form DAE (after UV exposure)

demonstrates significantly higher current than with the open-form isomers (after visible irradiation) at the same applied voltage. For example, approximately 7.12 nA was measured at 0.60 V after UV irradiation compared to \sim 0.07 nA after exposure to visible light. It is thus evident that light conditions drastically influence the hybrid, and the conductivity of 1L MoS₂ is greatly enhanced with closed-form DAE. It should be noted that the photo-modulation is not observed from MoS₂ without DAE molecules (Fig. S3b†). The two different isomeric forms of DAE also change the threshold voltage significantly. It decreases from \sim 0.50 V (open-form) to \sim 0.14 V (closed-form) under positive potential and increases from \sim –0.60 V to \sim –0.30 V under negative potential. This could be due to either p- or n-doping effects. Either way, the number of majority carriers will increase, which will accordingly improve the conductivity and promote greater current flow under a given DC bias voltage.³⁹

In contrast, the C-AFM results of 2L and 3L MoS₂-DAE samples in Fig. 3b and c show that there are no significant changes in current and conductivity after UV-induced isomerization from the DAE open-form to the closed-form. The small difference in current of 2L MoS₂-DAE around 1 V DC bias voltage may be attributed to a local electric field effect, photo-thermal effect, or minor doping.^{1,40} Since only 1L hybrid demonstrated enhanced conductivity and current flow while 2L and 3L samples did not, we performed DFT calculations to



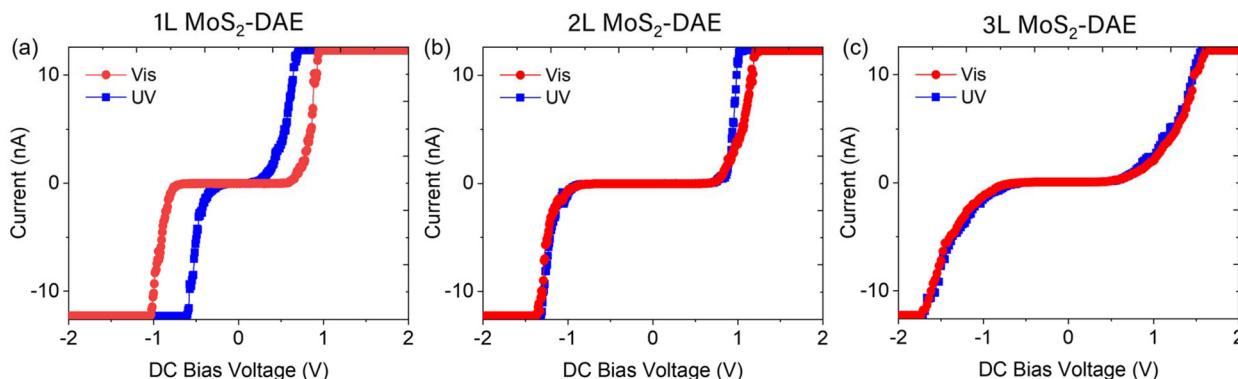


Fig. 3 Conductive AFM measurements on photoswitched hybrid samples under UV and visible irradiation: (a) 1L MoS₂-DAE, (b) 2L MoS₂-DAE, and (c) 3L MoS₂-DAE. Bias DC voltage was applied from -2 V to 2 V and currents were recorded within ± 12 nA. UV light enhances the current significantly in 1L MoS₂-DAE, whereas no significant changes are observed from the hybrids with 2L and 3L MoS₂.

elucidate potential doping mechanisms in 1L MoS₂-DAE that account for the difference.

Densities of states (DOS) were obtained from DFT calculations for the 1L, 2L, and 3L MoS₂-DAE hybrids as shown in Fig. 4 (also see Fig. S1†). The energy states of an isolated DAE photoisomer were calculated following the methodology from our previous work,²⁵ setting the vacuum level to zero. DFT calculations have proven to be effective determining relative energy level alignments in 2D heterolayers.^{41,42} The structural transformation of DAE between its open and closed states leads to considerable shifts in the HOMO and LUMO levels. The closed-form DAE exhibits HOMO and LUMO energy levels of approximately -5.75 and -4.31 eV, respectively; around -6.32 and -3.80 eV are obtained with the open isomeric form. The gap between the HOMO and LUMO levels is ~ 1.4 eV for

closed-form isomer and ~ 2.5 eV for open DAE, resulting in a gap change of around 1.1 eV. Thus, the energy level window becomes narrowed after UV irradiation.

It can be observed that the VBM level of MoS₂ rises drastically as the layer number increases from 1L to 3L, while the CBM remains similar. The CBM of 1L, 2L, and 3L MoS₂ are approximately -4.33 , -4.34 , and -4.36 eV, respectively; the VBM levels are about -5.91 , -5.49 , and -5.39 eV. Investigating the energy differences between LUMO and CBM and between HOMO and VBM is crucial for determining energetically favorable charge transfer mechanisms and identifying which carrier will be transferred between MoS₂ and DAE. Electron transfer from MoS₂ to DAE may occur if the DAE's LUMO is lower than the CBM of MoS₂, and hole transfer will take place if the DAE's HOMO is higher than the VBM of MoS₂.

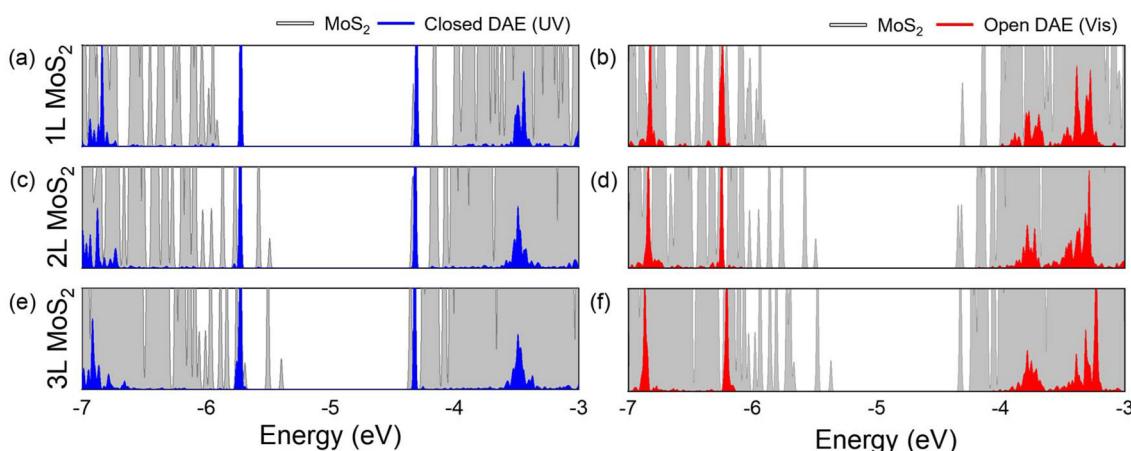


Fig. 4 DFT-calculation of density of electronic states (DOS), showing the alignment of VBM and CBM energies of 1L–3L MoS₂ with the HOMO and LUMO levels of closed- and open-form DAE under UV and visible light, respectively. The LUMO and HOMO levels of closed-form DAE are approximately -4.31 and -5.75 eV. In the open state, these values are around -3.80 and -6.32 eV, respectively. For closed-form DAE interacting with 1L MoS₂ (a), its HOMO level is higher than the MoS₂ VBM by approximately 0.16 eV. However, it is lower than the VBM levels of 2L and 3L MoS₂ by approximately 0.26 eV (c) and 0.36 eV (e). From (b), (d), and (f), the LUMO level is slightly higher than the CBM energies of all 1L, 2L, 3L MoS₂ by approximately 0.02 – 0.05 eV. In contrast, the HOMO level of open-form isomer is always lower than the MoS₂ VBM energies, and the LUMO is significantly higher than the CBM of 1L, 2L, and 3L MoS₂.



For the hybrids with the closed-form DAE, the LUMO level is slightly higher than the CBM of 1L, 2L, and 3L MoS₂ by approximately 0.02, 0.03, and 0.05 eV, respectively. Thus, there may not be a tendency of electron transfer from MoS₂ to the closed DAE. The HOMO is higher than the VBM of 1L MoS₂ by ~0.16 eV, but it becomes lower than the VBM of 2L and 3L MoS₂ by 0.26 eV and 0.36 eV, respectively. This suggests the possible hole transfer to the closed DAE from 1L MoS₂, but not from 2L or 3L MoS₂. This hole transfer generates the doping effect in the 1L MoS₂ flake.

In contrast, both the HOMO and LUMO levels of open-form DAE are outside the bandgap of all MoS₂ layers as shown in Fig. 4b, d, and f. The HOMO level of the open-form isomer is always lower than the VBM of MoS₂, while its LUMO is significantly higher than the CBM of 1L, 2L, and 3L MoS₂ by about 0.53, 0.54, and 0.56 eV, respectively. Therefore, charge transfer from MoS₂ layers to the open DAE isomers is unlikely.

The energy band diagram in Fig. 5 illustrates the fundamental mechanisms involved in MoS₂-DAE hybrids based on DFT calculations. Since the LUMO levels of both the open and the closed isomeric forms of DAE are higher than the CBM levels of the MoS₂ layers, electron transfer from MoS₂ to DAE cannot take place. On the other hand, the HOMO level of closed-form DAE is higher than the VBM energy of 1L MoS₂, which may facilitate the hole transfer from the 1L MoS₂ to

DAE. However, such transfer is unfavorable with 2L and 3L MoS₂. Overall, the photoswitchable hole transfer is a dominant mechanism that quenches PL and enhances conductivity of 1L MoS₂. Furthermore, there is a strong layer-number-dependence that alters the interactions between MoS₂ and photochromic DAE molecules.

This study is distinct and different from 1L MoSe₂-DAE system in our previous report, where the interfacial interactions in the hybrid structure were dominated by electron transfer mechanisms.²⁵ 1L MoSe₂ has a VBM that is higher than the HOMO of both open- and closed-form DAE by ~0.97 eV and 0.40 eV, respectively. The CBM energy is approximately 0.05 eV lower than the LUMO level of open-form DAE while ~0.46 eV higher than that of the corresponding closed isomer. This alignment facilitated electron transfer from the CBM of 1L MoSe₂ to the LUMO of closed-ring DAE while suppressing any interaction with the open isomeric form. Another similar study was reported which studied electron transfer in few-layer WSe₂ interfacing a different type of DAE.²⁶ The report presented that the CBM of few-layer WSe₂ (~3.5 eV) is higher than the LUMO of the closed form of DAE (~3.7 eV) but lower than that of the open form of DAE (~2.7 eV), resulting in electron charge transfer. In parallel, hole transfer has been observed in TMDs with organic layers. Huang and co-workers investigated interfacial charge transfer between 1L WS₂ and a layer of tetracene molecules by using an ultrafast transient absorption method.⁴³ When excited with a photon energy of ~2.1 eV which is below tetracene's absorption edge, a hole transfer took place from the WS₂ flake to tetracene. Building on these studies, our current work elucidates the layer-number-dependent photoswitchability in TMD-photochromic organic hybrid systems. We show that effective photo-modulated charge transfer in hybrid devices can be achieved by using DAE derivatives with energy levels specifically aligned to 2D semiconductors.

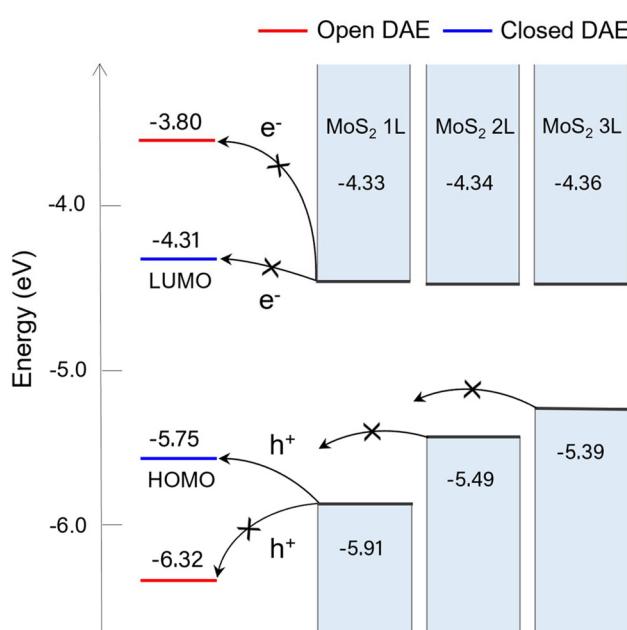


Fig. 5 The energy diagrams of hybrid structures involving 1L to 3L MoS₂ and DAE show the alignment of MoS₂'s CBM and VBM energies against the HOMO and LUMO levels of open- and closed-form DAE. In 1L MoS₂, hole transfer from MoS₂ to closed-form DAE is favorable since the HOMO is positioned higher than the CBM. In contrast, the CBM energies of 2L and 3L MoS₂ are higher than the HOMO levels of both open and closed states, prohibiting hole transfer mechanisms. Photoinduced electron transfer is energetically not favored in all MoS₂ configurations, because the CBM energies are lower than the LUMO of both open and closed DAE.

Conclusions

In this study, we demonstrated the effective photo-modulation of MoS₂ flakes using DAE molecules. The interaction between DAE and MoS₂ varies significantly with the number of MoS₂ layers, as evidenced by PL and C-AFM measurements. Notably, monolayer MoS₂-DAE displays significant PL quenching and increased conductivity after exposure to UV light (closed-ring isomers). DFT calculations support the observation showing that hole transfer is favorable in 1L MoS₂ when interacting with the closed-form DAE. On the other hand, 2L and 3L MoS₂ will not have any transfer interactions with DAE due to the energy level alignment.

This work highlights the importance of precise layer control in developing optoelectronic devices and the potential of photoswitchable TMD devices. Our strategy may be extended to other 2D materials by utilizing other photochromic molecules (e.g., spiropyran and azobenzene) or a combination of multiple such molecules. A detailed understanding of interfacial interactions between TMDs and photochromic mole-



cules will be required to design such systems. The photoswitchability and layer-number-dependence could open new possibilities for designing light-responsive nanoelectronics and optoelectronics. The advancements will usher in a new era of electronics, incorporating innovative functions within high-performing 2D materials toward light-driven technologies.

Materials and methods

Preparation of MoS₂

Bulk MoS₂ crystals (2D Semiconductors) were mechanically exfoliated onto cleaned p-doped SiO₂/Si substrates and indium tin oxide glass (ITO, MTI KJ Group) using the Scotch tape method.^{44,45} They were subsequently heated on a hot plate at 50 °C for 10 or 30 minutes to enhance the adhesion of the exfoliated flakes. Before deposition, the SiO₂/Si substrates were cleaned by ultrasonication in acetone, methanol, and deionized water (DI) for 15 minutes each to remove any contaminants or impurities from the surface. The residual fluids were blown off using compressed air, followed by soft baking on a hot plate at 110 °C for 5 minutes to ensure complete evaporation. The annealing process was conducted in a controlled Ar gas atmosphere at 250 °C for 1 hour to ensure optimal contact between the MoS₂ flakes and the substrate, reduce surface contamination, enhance the quality of the interface, and improve the electronic properties.^{45–47}

Synthesis of DAE molecules

To synthesize the DAE compound, we followed a synthetic approach for the Negishi coupling reaction. This method was based on the Knochel modification⁴⁸ and subsequent references,³⁷ as recommended by previous reports.^{25,49} Briefly, we utilized 2,4-dibromo-5-methylthiazole and 3,5-bis(trifluoromethyl)-1-bromobenzene as reactants.²⁵ This approach yielded 2-(3,5-bis(trifluoromethyl)phenyl)-4-bromo-5-methylthiophene as an intermediate product. The intermediate was converted *in situ* to the corresponding boronic pinacol ester, which was crucial for the Suzuki coupling reaction with 1,2-dichlorohexa-fluorocyclopent-1-ene, enabling the synthesis of the targeted DAE. The final DAE molecules were confirmed by nuclear magnetic resonance (NMR) measurement.^{25,50} They have zero absorption at 633 nm and do not exhibit fluorescent properties.^{25,51}

Functionalization of MoS₂ with DAE

Both SiO₂/Si substrates and ITO glass with MoS₂ flakes were immersed in a DAE solution (0.02 mg mL⁻¹ in chloroform) for 12 hours at 4 °C.²⁶ After drying the DAE solution from the samples with compressed air, they underwent soft baking on a hot plate at 50 °C for 30 minutes. The photoswitching behaviors were studied by exposing the MoS₂-DAE samples to UV light for 2 minutes, followed by exposure to visible light for 10 minutes. UV light was emitted by a Spectroline™ TE lamp at approximately 312 nm, while visible light was produced from a Xe-arc lamp with a 530 nm long-pass filter.

Sample characterization

A Renishaw inVia Raman confocal microscope was used for MoS₂ flake inspection and PL measurements with 10×, 20×, 50×, and 100× objective lenses. PL emissions were recorded using a 633 nm HeNe laser and a 100× objective lens. The layer numbers were first identified based on characteristic color contrasts caused by differences in transparency and interference effects^{33,47} and confirmed by AFM imaging. AFM and C-AFM were performed on both MoS₂ and MoS₂-DAE samples using the Bruker Dimension Icon AFM. Topological imaging to measure flake thickness utilized a SCANASYST-AIR probe, while an SCM-PIT-V2 probe was employed in C-AFM measurements. While measuring C-AFM, a current range of ±12 nA with a sensitivity of 1 nA V⁻¹ was set, and bias voltage ranging from -2 V to 2 V was applied to avoid potential damage to both the probe tip and the samples.⁵²

DFT calculation

Vienna *Ab initio* Simulation Package (VASP) was used for DFT calculations with the projector augmented wave (PAW) method.^{53,54} To handle electron exchange and correlation, the Perdew–Burke–Ernzerhof (PBE) functional⁵⁵ within the generalized gradient approximation was applied. The hybrid structures consisted of a DAE isomer placed on top of a (5 × 5) MoS₂ slab as illustrated in Fig. S6.† A 14 Å vacuum layer was added. This setup resulted in an areal density of 4.5×10^{13} molecules per cm². The MoS₂ bi- and tri-layers were constructed with the AA' type stacking order where the Mo(S) atom in one layer is aligned over the S(Mo) atom in the adjacent layer.⁵⁶ To accurately model the van der Waals interactions between the MoS₂ layers and the DAE, a dispersion correction was incorporated using the DFT-D2 method with Grimme.³² The electronic minimization was performed with a tolerance level of 10⁻⁵ eV. Ionic positions were relaxed until the Hellmann–Feynman forces were below 10 meV Å⁻¹ with a 2 × 2 × 1 Γ -centered k -point mesh for the DAE-(5 × 5) MoS₂ hybrid systems. After the ionic relaxation, a finer k -point mesh of 3 × 3 × 1 Γ -centered mesh was used for generating precise density of states.

Data availability

The data supporting this article have been included as part of the ESI document.† Data are available upon request from the authors.

Conflicts of interest

The authors declare no competing financial interest.

Acknowledgements

This work was funded by the U.S. National Science Foundation under award no. 2151869 and 2151887.



References

1 Q. H. Wang, K. Kalantar-Zadeh, A. Kis, J. N. Coleman and M. S. Strano, Electronics and optoelectronics of two-dimensional transition metal dichalcogenides, *Nat. Nanotechnol.*, 2012, **7**, 699–712.

2 J. Yuan, S. Najmaei, Z. Zhang, J. Zhang, S. Lei, P. M. Ajayan, B. I. Yakobson and J. Lou, Photoluminescence quenching and charge transfer in artificial heterostacks of monolayer transition metal dichalcogenides and few-layer black phosphorus, *ACS Nano*, 2015, **9**, 555–563.

3 D. Yang, S. J. Sandoval, W. Divigalpitiya, J. Irwin and R. Frindt, Structure of single-molecular-layer MoS₂, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1991, **43**, 12053.

4 V. Yadav, S. Roy, P. Singh, Z. Khan and A. Jaiswal, 2D MoS₂-based nanomaterials for therapeutic, bio-imaging, and biosensing applications, *Small*, 2019, **15**, 1803706.

5 X. Zhuang, Y. Mai, D. Wu, F. Zhang and X. Feng, Two-dimensional soft nanomaterials: a fascinating world of materials, *Adv. Mater.*, 2015, **27**, 403–427.

6 S. Balendhran, S. Walia, H. Nili, J. Z. Ou, S. Zhuiykov, R. B. Kaner, S. Sriram, M. Bhaskaran and K. Kalantar-zadeh, Two-dimensional molybdenum trioxide and dichalcogenides, *Adv. Funct. Mater.*, 2013, **23**, 3952–3970.

7 K. F. Mak, C. Lee, J. Hone, J. Shan and T. F. Heinz, Atomically thin MoS₂: a new direct-gap semiconductor, *Phys. Rev. Lett.*, 2010, **105**, 136805.

8 D. Jariwala, V. K. Sangwan, L. J. Lauhon, T. J. Marks and M. C. Hersam, Emerging device applications for semiconducting two-dimensional transition metal dichalcogenides, *ACS Nano*, 2014, **8**, 1102–1120.

9 M. Naqi, M. Kaniselvan, S. Choo, G. Han, S. Kang, J. Kim, Y. Yoon and S. Kim, Ultrasensitive multilayer MoS₂-based photodetector with permanently grounded gate effect, *Adv. Electron. Mater.*, 2020, **6**, 1901256.

10 H. Kwon, P. J. Jeon, J. S. Kim, T.-Y. Kim, H. Yun, S. W. Lee, T. Lee and S. Im, Large scale MoS₂ nanosheet logic circuits integrated by photolithography on glass, *2D Mater.*, 2016, **3**, 044001.

11 S. Barua, H. S. Dutta, S. Gogoi, R. Devi and R. Khan, Nanostructured MoS₂-based advanced biosensors: a review, *ACS Appl. Nano Mater.*, 2017, **1**, 2–25.

12 Z. Yu, Z. Y. Ong, S. Li, J. B. Xu, G. Zhang, Y. W. Zhang, Y. Shi and X. Wang, Analyzing the carrier mobility in transition-metal dichalcogenide MoS₂ field-effect transistors, *Adv. Funct. Mater.*, 2017, **27**, 1604093.

13 J. Lin, J. Zhong, S. Zhong, H. Li, H. Zhang and W. Chen, Modulating electronic transport properties of MoS₂ field effect transistor by surface overlayers, *Appl. Phys. Lett.*, 2013, **103**, 063109.

14 S. Aftab, Samiya, H. M. Ul Haq, S. Yousuf, M. U. Khan, Z. Ahmed, J. Aziz, M. W. Iqbal, A. ur Rehman and M. Z. Iqbal, van der Waals multi-heterostructures (PN, PIN, and NPN) for dynamic rectification in 2D materials, *Adv. Mater. Interfaces*, 2020, **7**, 2001479.

15 K. S. Novoselov, A. Mishchenko, A. Carvalho and A. H. Castro Neto, 2D materials and van der Waals heterostructures, *Science*, 2016, **353**, aac9439.

16 J. Ji and J. H. Choi, Layer-number-dependent electronic and optoelectronic properties of 2D WSe₂-organic hybrid heterojunction, *Adv. Mater. Interfaces*, 2019, **6**, 1900637.

17 F. Pizzocchero, L. Gammelgaard, B. S. Jessen, J. M. Caridad, L. Wang, J. Hone, P. Bøggild and T. J. Booth, The hot pick-up technique for batch assembly of van der Waals heterostructures, *Nat. Commun.*, 2016, **7**, 11894.

18 X. Yang and M. Yan, Removing contaminants from transferred CVD graphene, *Nano Res.*, 2020, **13**, 599–610.

19 J. Ji and J. H. Choi, Recent progress in 2D hybrid heterostructures from transition metal dichalcogenides and organic layers: Properties and applications in energy and optoelectronics fields, *Nanoscale*, 2022, **14**, 10648–10689.

20 Y. L. Huang, Y. J. Zheng, Z. Song, D. Chi, A. T. Wee and S. Y. Quek, The organic-2D transition metal dichalcogenide heterointerface, *Chem. Soc. Rev.*, 2018, **47**, 3241–3264.

21 M. Irie, T. Fukaminato, K. Matsuda and S. Kobatake, Photochromism of diarylethene molecules and crystals: memories, switches, and actuators, *Chem. Rev.*, 2014, **114**, 12174–12277.

22 F. A. Jercia, V. V. Jercia and R. Hoogenboom, Advances and opportunities in the exciting world of azobenzenes, *Nat. Rev. Chem.*, 2022, **6**, 51–69.

23 R. Klajn, Spiropyran-based dynamic materials, *Chem. Soc. Rev.*, 2014, **43**, 148–184.

24 C. L. Fleming, S. Li, M. Grøtli and J. Andréasson, Shining new light on the spiropyran photoswitch: a photocage decides between cis-trans or spiro-merocyanine isomerization, *J. Am. Chem. Soc.*, 2018, **140**, 14069–14072.

25 S. Park, J. Ji, C. Cunningham, S. Pillai, J. Rouillon, C. Benitez-Martin, M. Fang, E. H. Yang, J. Andreasson, J. H. You and J. H. Choi, Photoswitchable optoelectronic properties of 2D MoSe₂/Diarylethene hybrid structures, *Sci. Rep.*, 2024, **14**, 7325.

26 H. Qiu, Y. Zhao, Z. Liu, M. Herder, S. Hecht and P. Samorì, Modulating the Charge Transport in 2D Semiconductors via Energy-Level Phototuning, *Adv. Mater.*, 2019, **31**, 1903402.

27 M. Gobbi, S. Bonacchi, J. X. Lian, A. Vercouter, S. Bertolazzi, B. Zyska, M. Timpel, R. Tatti, Y. Olivier, S. Hecht, M. V. Nardi, D. Beljonne, E. Orgiu and P. Samorì, Collective molecular switching in hybrid superlattices for light-modulated two-dimensional electronics, *Nat. Commun.*, 2018, **9**, 2661.

28 Y. Zhao, S. Ippolito and P. Samorì, Functionalization of 2D Materials with Photosensitive Molecules: From Light-Responsive Hybrid Systems to Multifunctional Devices, *Adv. Opt. Mater.*, 2019, **7**, 1900286.

29 C. Lee, H. Yan, L. E. Brus, T. F. Heinz, J. Hone and S. Ryu, Anomalous lattice vibrations of single-and few-layer MoS₂, *ACS Nano*, 2010, **4**, 2695–2700.



30 Y. Zhang, J. Ye, Y. Matsuhashi and Y. Iwasa, Ambipolar MoS₂ thin flake transistors, *Nano Lett.*, 2012, **12**, 1136–1140.

31 T. Zheng, P. Valencia-Acuna, P. Zereshki, K. M. Beech, L. Deng, Z. Ni and H. Zhao, Thickness-dependent interlayer charge transfer in MoSe₂/MoS₂ heterostructures studied by femtosecond transient absorption measurements, *ACS Appl. Mater. Interfaces*, 2021, **13**, 6489–6495.

32 S. Grimme, Semiempirical GGA-type density functional constructed with a long-range dispersion correction, *J. Comput. Chem.*, 2006, **27**, 1787–1799.

33 J. Ji, C. M. Delehey, D. N. Houpt, M. K. Heighway, T. Lee and J. H. Choi, Selective chemical modulation of interlayer excitons in atomically thin heterostructures, *Nano Lett.*, 2020, **20**, 2500–2506.

34 H. Li, Q. Zhang, C. C. R. Yap, B. K. Tay, T. H. T. Edwin, A. Olivier and D. Baillargeat, From bulk to monolayer MoS₂: evolution of Raman scattering, *Adv. Funct. Mater.*, 2012, **22**, 1385–1390.

35 S. Mouri, Y. Miyauchi and K. Matsuda, Tunable photoluminescence of monolayer MoS₂ via chemical doping, *Nano Lett.*, 2013, **13**, 5944–5948.

36 J. Shi, P. Yu, F. Liu, P. He, R. Wang, L. Qin, J. Zhou, X. Li, J. Zhou and X. Sui, 3R MoS₂ with broken inversion symmetry: a promising ultrathin nonlinear optical device, *Adv. Mater.*, 2017, **29**, 1701486.

37 M. Herder, B. M. Schmidt, L. Grubert, M. Pätz, J. Schwarz and S. Hecht, Improving the fatigue resistance of diarylethene switches, *J. Am. Chem. Soc.*, 2015, **137**, 2738–2747.

38 L. Hou, R. Ringström, A. B. Maurer, M. Abrahamsson, J. Andréasson and B. Albinsson, Optically Switchable NIR Photoluminescence of PbS Semiconducting Nanocrystals using Diarylethene Photoswitches, *J. Am. Chem. Soc.*, 2022, **144**, 17758–17762.

39 S. Zhang, H. M. Hill, K. Moudgil, C. A. Richter, A. R. Hight Walker, S. Barlow, S. R. Marder, C. A. Hacker and S. J. Pookpanratana, Controllable, wide-ranging n-doping and p-doping of monolayer group 6 transition-metal disulfides and diselenides, *Adv. Mater.*, 2018, **30**, 1802991.

40 W. Sakong, H. Z. Gul, B. Ahn, S. Oh, G. Kim, E. Sim, J. Bahng, H. Yi, M. Kim and M. Yun, Optical duality of molybdenum disulfide: metal and semiconductor, *Nano Lett.*, 2022, **22**, 5207–5213.

41 J. Kunstmann, F. Mooshammer, P. Nagler, A. Chaves, F. Stein, N. Paradiso, G. Plechinger, C. Strunk, C. Schüller and G. Seifert, Momentum-space indirect interlayer excitons in transition-metal dichalcogenide van der Waals heterostructures, *Nat. Phys.*, 2018, **14**, 801–805.

42 R. Besse, J. F. Silveira, Z. Jiang, D. West, S. Zhang and J. L. Da Silva, Beyond the Anderson rule: importance of interfacial dipole and hybridization in van der Waals heterostructures, *2D Mater.*, 2021, **8**, 041002.

43 T. Zhu, L. Yuan, Y. Zhao, M. Zhou, Y. Wan, J. Mei and L. Huang, Highly mobile charge-transfer excitons in two-dimensional WS₂/tetracene heterostructures, *Sci. Adv.*, 2018, **4**, eaao3104.

44 S. Park, J. Ha, M. F. Khan, C. Im, J. Y. Park, S. H. Yoo, M. A. Rehman, K. Kang, S. H. Lee and S. C. Jun, Pronounced Optoelectronic Effect in n-n Res₂ Homostructure, *ACS Appl. Electron. Mater.*, 2022, **4**, 4306–4315.

45 J. Ji and J. H. Choi, Understanding the Effects of Dielectric Property, Separation Distance, and Band Alignment on Interlayer Excitons in 2D Hybrid MoS₂/WSe₂ Heterostructures, *ACS Appl. Electron. Mater.*, 2021, **3**, 3052–3059.

46 J. Choi, H. Zhang, H. Du and J. H. Choi, Understanding solvent effects on the properties of two-dimensional transition metal dichalcogenides, *ACS Appl. Mater. Interfaces*, 2016, **8**, 8864–8869.

47 H. Zhang, J. Ji, A. A. Gonzalez and J. H. Choi, Tailoring photoelectrochemical properties of semiconducting transition metal dichalcogenide nanolayers with porphyrin functionalization, *J. Mater. Chem. C*, 2017, **5**, 11233–11238.

48 S. Sase, M. Jaric, A. Metzger, V. Malakhov and P. Knochel, One-pot Negishi cross-coupling reactions of in situ generated zinc reagents with aryl chlorides, bromides, and triflates, *J. Org. Chem.*, 2008, **73**, 7380–7382.

49 G. Naren, W. Larsson, C. Benitez-Martin, S. Li, E. Pérez-Inestrosa, B. Albinsson and J. Andréasson, Rapid amplitude-modulation of a diarylethene photoswitch: en route to contrast-enhanced fluorescence imaging, *Chem. Sci.*, 2021, **12**, 7073–7078.

50 C.-W. Hsu, C. Sauvée, H. Sundén and J. Andréasson, Writing and erasing multicolored information in diarylethene-based supramolecular gels, *Chem. Sci.*, 2018, **9**, 8019–8023.

51 L. Hou, W. Larsson, S. Hecht, J. Andréasson and B. Albinsson, A general approach for all-visible-light switching of diarylethenes through triplet sensitization using semiconducting nanocrystals, *J. Mater. Chem. C*, 2022, **10**, 15833–15842.

52 J. Choi, H. Zhang and J. H. Choi, Modulating optoelectronic properties of two-dimensional transition metal dichalcogenide semiconductors by photoinduced charge transfer, *ACS Nano*, 2016, **10**, 1671–1680.

53 G. Kresse and J. Furthmüller, Efficient iterative schemes for ab initio total-energy calculations using a plane-wave basis set, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1996, **54**, 11169.

54 G. Kresse and D. Joubert, From ultrasoft pseudopotentials to the projector augmented-wave method, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1999, **59**, 1758.

55 J. P. Perdew, K. Burke and M. Ernzerhof, Generalized gradient approximation made simple, *Phys. Rev. Lett.*, 1996, **77**, 3865.

56 J. He, K. Hummer and C. Franchini, Stacking effects on the electronic and optical properties of bilayer transition metal dichalcogenides MoS₂, MoSe₂, WS₂, and WSe₂, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2014, **89**, 075409.

