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On-surface magnetocaloric effect for a van der Waals Gd(III) 2D MOF grown on Si†

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Gd complexes have been studied as magnetic coolants due to their large magneto-caloric effect. In this work we present a Gd 2D metal–organic framework (MOF) of formula $[\text{Gd}(\text{MeCOO})(\text{PhCOO})_2]$ (**1**). We characterize the magnetic properties of **1**, showing that it displays slow relaxation of the magnetization by ac susceptibility, and single-ion magnetic anisotropy using high-field EPR. By heat capacity and magnetization vs. field at various temperatures we determine the magnetic entropy change of compound **1**. We then grow **1** on functionalized silicon, and show that the surface-deposited 2D MOF **1**Si can be used as an on-surface magnetic cryogenic coolant.

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Introduction

Two dimensional (2D) materials are fantastic platforms that can provide exceptional performance in many applications.¹ Fabrication of heterostructures or multifunctional devices is a topic that is already being exploited by many research groups, in particular using graphene.² Molecule-based systems have been proposed as a basis for next-generation technologies, with applications ranging from spintronics and quantum technologies^{3–5} to magnetic cooling.⁶ For this purpose, surface deposition and two-dimensional (2D) organization of the systems are key challenges. There are already several examples where molecular systems are combined with graphene.^{7–12} Most 2D materials are inorganic solids, like chalcogenides, MXenes or graphene. In the last few years, researchers have started to develop ways to obtain 2D

coordination polymers or metal organic frameworks, sometimes also called metal–organic nanosheets (MONs).^{13,14} 2D metal–organic frameworks (MOFs) offer the possibility of organizing a molecular entity in 2D, on a topologically defined network: 2D arrays of magnetic molecules can be fundamental for advances in molecular spintronics applications or devices.¹⁵

For example, 2D arrays of qubits can be designed to control dipolar interactions and to integrate qubits into devices.^{16,17} If the magnetic molecules are organized in a 2D coordination polymer, intrinsic magnetic dilution becomes possible, as heterometallic complexes can be easily achieved. In this landscape, 2D MOFs afford the possibility of organizing functional units in crystalline 2D arrays. In our previous work we demonstrated that we could prepare homometallic and heterometallic 2D MOFs using lanthanoid ions. The compound $[\text{Ln}(\text{MeCOO})(\text{PhCOO})_2]$ is one such van der Waals 2D MOF, where the $\text{Ln} = \text{Dy}$ analogue displays single-molecule magnet (SMM) properties that are enhanced when the system is diluted with $\text{Ln} : \text{Dy} = 9 : 1$.¹⁸ The Tb and TbEu analogues display both SMM and luminescent properties.^{19,20} The fact that the interactions between 2D layers are purely van der Waals forces of the order of those found in graphene allows an easy delamination of the material in flakes and monolayers.

The Gd(III) ion possess a large $S = 7/2$ spin, which favors large magnetic entropy changes in the magnetization–demagnetization cycles, since maximum entropy can be obtained from $S_{\max} = R \ln(2S + 1)$. The magnetocaloric effect (MCE) exploits the entropy changes that result in the heating or cooling of the material by applying an external magnetic field at constant pressure. Complexes with Gd(III) ions isolated (or very weakly coupled) and with low molecular weight ligands are good candidates for low temperature magnetic cooling.^{6,21,22} The 2D material of formula $[\text{Gd}(\text{MeCOO})(\text{PhCOO})_2]$ is expected to fulfill these conditions, given the low molecular weight of the ligands,

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and the large separation between the Gd ions in the structure. The MCE performance of 4f and 3d molecular compounds of various dimensionalities (0D, 1D, 2D and 3D) has been recently reviewed; in particular, several 2D Gd MOFs have large magnetic entropy changes ($-\Delta S_m$) of the order of 30–40 J K⁻¹ kg⁻¹ at 7 T.²³ 2D materials offer advantages over 3D solids, in that they can be used to functionalize a surface to act as on-surface coolants, or they can be dispersed as 2D flakes on a coating, thus providing new possibilities for integrating the MCE into devices. Low temperature cooling applications with low energy consumption, such as the MCE with molecular materials that can be easily integrated into devices,^{24,25} can be a key technology for the development of quantum computers, which require ultra-low temperatures to function.

Thus, we decided to obtain the Gd(III) analogue of the complex, [Gd(MeCOO)(PhCOO)₂] (**1**), in order to exploit its MCE properties as a bulk material and as crystals grown on a surface. In this work, we report the synthesis and comprehensive study of the magnetic properties of [Gd(MeCOO)(PhCOO)₂] (**1**), including the MCE. We also exploit the surface growth of microcrystals of the 2D MOF **1Si** and we show that the MCE properties are retained.

Materials and methods

All reagents are acquired from commercial sources and used as received. [Gd(MeCOO)(PhCOO)₂] (**1**): 2D MOF is prepared following the reported procedure using a microwave assisted reaction and Gd(MeCOO)₃·xH₂O instead of Dy(MeCOO)₃·xH₂O.^{18,19} 0.26 mmol of Gd(MeCOO)₃·xH₂O and 0.52 mmol of PhCOOH were placed in a microwave reactor tube with 4 mL of a MeOH : MeCN mixture in 1 : 1 proportion. A pulse of 150 W is applied and the reaction kept for 10 minutes at a maximum temperature of 125 °C. The reaction is cooled to room temperature, the microwave tube is removed from the reactor and the colorless solution is left undisturbed. Colorless plate-like crystals suitable for single-crystal X-ray diffraction (SCXRD) grow in 7–15 days. Yield 30–45%. Yield of crystals can be improved by growing the crystals at 40 °C, in a sealed vial in an oven. Elemental analyses for [Gd(MeCOO)(PhCOO)₂] (**1**): experimental (calculated), C 42.01 (41.83); H 2.53 (2.85); N 0 (0%). Relevant IR data: IR data (cm⁻¹): 3070 (w), 3039 (w), 1590 (s), 1436 (w), 1394 (s), 1331 (m), 1248 (m), 1164 (s), 1068 (w), 1022 (w), 993 (m), 930 (m), 868 (w), 847 (w), 830 (w), 788 (m), 709 (s), 700 (s), 687 (s), 667 (s), 617 (w), 604 (w).

Silanization of p-doped Si (Si-TSP)

We followed the procedure modified from Godoy-Gallardo *et al.*²⁶ reported by us.²⁷ Resultant layer thickness was measured with ellipsometry, the TSP layers were consistently between 1 and 1.5 nm. Hydrophobicity of the resultant surface was measured with a homemade contact angle meter (48°).

On-surface growth of **1** (**1Si**)

0.26 mmol of Gd(MeCOO)₃·xH₂O and 0.52 mmol of PhCOOH were placed in a microwave reactor tube with 4 mL of a MeOH :

MeCN mixture in 1 : 1 ratio. A pulse of 150 W is applied and the reaction kept for 10 minutes at a maximum temperature of 125 °C. The reaction is cooled to room temperature, the microwave tube is removed from the reactor and the colorless solution is decanted and diluted to 10 mL of MeCN : MeOH 1 : 1. The 10 mL of solution are divided into vials. Si-TSP wafers are submerged in the colorless solution and kept there for times between 15 minutes and 2 days. After the indicated time, the wafers are carefully removed from the solution and rinsed with MeCN : MeOH and acetone, and dried. This process results in **1** chemisorbed on silicon wafers, **1Si**. Resultant layer thickness was measured with ellipsometry, the MOF layers were consistently between 2 and 7 nm. Hydrophobicity of the resultant surface was measured with a homemade contact angle meter (67°).

Single crystal X-ray diffraction data for **1** were collected at the XALOC beamline 13 of ALBA-CELLS Synchrotron (Spain) (*T* = 100 K, λ = 0.7239 Å). The structure was solved by intrinsic phasing methods (SHELXT using the XIA package for the data collected on the beamline) and refined on F². Hydrogen atoms were included at the calculated positions, riding on their carrier atoms. A cif file can be obtained from <https://www.ccdc.cam.ac.uk/structures/> with CCDC number 2250164. Infrared spectra were collected on an FT-IR Thermo Scientific NicoletTM iS5 Spectrometer equipped with an iD7 ATR accessory. Elemental analyses were carried out at the Centres Científics i Tecnològics de la Universitat de Barcelona (CCiTUB).

Alternating current (AC) and direct current (DC) magnetic data were collected using a Quantum Design MPMS SQUID magnetometer equipped with a 5 T magnet at SAI services at the University of Zaragoza. Experiments were conducted on powdered sample **1** embedded in Daphne oil to prevent grain orientation. AC measurements in the range between *f* = 0.1–1000 Hz, under an alternating field of *H*_{ac} = 4 Oe, were carried out in the same equipment. Additional AC measurements in the *f* = 100–10 000 Hz range were carried in a Quantum Design PPMS susceptometer at SAI. Magnetization *vs.* field data in the 2–10 K range were collected for **1** and **1Si** at the Unitat de Mesures Magnétiques at CCiTUB on a Quantum Design SQUID magnetometer equipped with a 7 T field.

Heat capacity measurements as a function of temperature were performed between 0.3–30 K at different applied fields between 0–30 kOe on a pressed powder pellet of **1** fixed with Apiezon N grease, using a Quantum Design PPMS equipped with a ³He refrigerator at SAI services.

X-ray absorption spectroscopy (XAS) and X-ray magnetic circular dichroism (XMCD) experiments of **1** flakes spin-coated onto silicon were performed at BOREAS beamline of the ALBA synchrotron. Measurements across the M_{4,5} edge of Gd were performed at 3.4 K. All spectra were recorded using Total Electron Yield (TEY) detection mode, with 90% circularly polarized light. The XMCD ($\mu^- - \mu^+$) and XAS ($\mu^+ + \mu^-/2$) spectra at 6 T were determined from eight X-ray absorption spectra measured under right-handed (μ^+) and left-handed (μ^-) circular polarizations. XMCD(*H*) cycles were performed by following the



resonant M_5 peak while sweeping the magnetic field between 6 T and -6 T at a rate of 2 T min $^{-1}$.

Ellipsometry data were collected with a J. A. Woollam Alpha-SE ellipsometer and analyzed with CompleteEase software. Scanning electron microscopy (SEM) images and energy dispersive X-ray spectroscopy (EDS) spectra were collected on a Jeol J-6510 scanning electron microscope (SEM). Contact angles were measured using home-made equipment at GMMF labs. Cleaning of the samples was performed using an ultrasound cleaning bath VWR USC600TH at room temperature. Raman spectra were collected at the CCiTUB.

Raman spectrum of **1** was computed using the Gaussian09 suite of programs. See ESI \dagger for details on the calculation.

High-frequency electron paramagnetic resonance (HF-EPR) spectra were collected using a home-built transmission spectrometer, as described in detail previously, operating in continuous-wave (CW) mode with field modulation producing a derivative spectrum. The sample was measured as a ground polycrystalline powder immobilized in an EPR sample holder. For the purposes of this investigation, all spectra were collected at a frequency of 104 GHz in applied fields up to 6 T and at temperatures in the range 5–300 K. All spectra were simulated using the EasySpin (version 6.0.0-dev.51) toolbox for Matlab.²⁸

Hard X-ray Photoelectron Spectroscopy (HAXPES) data were collected in ultra-high vacuum ($P = 1 \times 10^{-10}$ mbar UHV) conditions using a photon energy of 11 keV at the BM25-SpLine beamline at the European Synchrotron ESRF.²⁹ The HAXPES data was obtained using a specially developed high kinetic energy analyzer HV-CSA300 able to work up to kinetic energies of 15 keV.³⁰

Results and discussion

Synthesis and surface deposition

The 2D MOF reported in this work is a van der Waals 2D material with neutral nanosheets of formula $[\text{Gd}(\text{MeCOO})(\text{PhCOO})_2]$ (**1**). The material is prepared in a microwave reactor, as already reported. From the microwave assisted synthesis a colorless solution is obtained from which crystals suitable for single crystal X-ray diffraction can be grown in 5–15 days. The crystal structure obtained is shown in Fig. 1. It is analogous to that of the other reported compounds of the series.^{19,20} Compound **1** crystallizes in the monoclinic space group $P2_1/c$. van der Waals stacks of 2D layers of formula $[\text{Gd}(\text{MeCOO})(\text{PhCOO})_2]$ form along the a axis of the unit cell. The $\text{Gd}(\text{III})$ ion is octacoordinated in a very distorted square antiprism geometry. In fact SHAPE² analysis reveals that the geometry is best described as an intermediate between square antiprism and biaugmented trigonal prism. See ESI-1 \dagger for full SHAPE analysis on the Gd coordination environment. The 2D layers have a complicated topology of distorted triangles and hexagons, and are not planar (see ESI Fig. S1 \dagger). There are three distinct Gd–Gd distances in the layer: Gd–Gd at 3.947 Å (two monoatomic O-bridges from acetato ligands and two syn,syn-PhCOO); Gd–Gd at 4.700 Å (one monoatomic O-bridge from acetato and one syn,syn-PhCOO), and Gd–Gd at 6.961 Å (one anti,anti-acetato ligand).

In order to see if the material can be used as an effective on-surface magnetic coolant, we grafted nanocrystals of **1** onto functionalized Si wafers. The colorless solution obtained from the microwave assisted reaction can be used to graft the MOF onto a suitably functionalized substrate. In this work, we follow our use of a silanized silicon wafer with a succinate terminated silane. Hydrolysis of the succinate groups affords a surface of carboxylic acids. A simple ligand exchange reaction with one of the carboxylate ligands of the $[\text{Gd}(\text{MeCOO})(\text{PhCOO})_2]$ MOF should afford a chemically attached (chemisorbed) layer of $[\text{Gd}(\text{MeCOO})(\text{PhCOO})_2]$ on Si, **1Si**. By varying the time the **Si-TSP** wafer is submerged into the reaction mixture, one can control the amount of randomly oriented crystallites that grow on the surface of the Si wafer; for long deposition times, *e.g.*, 24 and 48 h, crystallites pile on the Si wafers as shown in the SEM images of Fig. 2. Ellipsometry analysis shows layer thicknesses of 2–6 nm for deposition times below 24 h. The contact angle changes from 48° for **Si-TSP** to 67° for **1Si**, confirming the more hydrophobic nature of the surface after MOF deposition. The image of the highest magnification (1 μm) in Fig. 2 shows a pair of crystals, clearly growing from the functionalized Si wafer. Using EDS analysis coupled to SEM we can ascertain that the thin layers and the crystallites deposited contain $\text{Gd}(\text{III})$ and organic ligands. Non-chemically attached crystallites can be removed by ultrasound cleaning of the **1Si** wafers. This shows that the first 1–5 nm layer is chemically bonded to the **Si-TSP**, but some of the larger crystals observed in Fig. 2 (48 h) are either more fragile and easily detached or not chemically attached to the surface and can be removed by sonication and washing. Raman spectra were collected on **1Si**. The spectra were compared to the calculated ones. There is a good agreement, as shown in ESI Fig. S2, \dagger confirming the chemical nature of the metal–organic layers grown on **1Si**.

The **1Si** sample was also chemically characterized by hard X-ray photoemission spectroscopy (HAXPES) at the BM25 beamline, SpLine at the European Synchrotron Radiation source (ESRF). Fig. 3 shows the photoemission spectra for the $\text{Gd} 2p_{1/2}$ and the $\text{O} 1s$ core levels. ESI Fig. S3 \dagger contains the $\text{Gd} 2p_{3/2}$ core level HAXPES spectra. The $\text{O} 1s$ signal shows a small peak at ~ 526 eV that can be attributed to oxygen bound to the $\text{Gd}(\text{III})$ ions of the MOF, while the bulk of the $\text{O} 1s$ signal arises from the SiO_2 layer.

Magnetic properties

The magnetic properties of **1** are shown in Fig. 4a as the susceptibility times the temperature product (χT) vs. T and magnetization as a function of field plots. The χT product at high temperatures reaches a value of 7.73 emu K mol $^{-1}$, approaching the expected free-ion value of 7.875 emu K mol $^{-1}$ for one $\text{Gd}(\text{III})$ with $S = 7/2$ and $g = 2.0$. As the temperature decreases, the χT value remains nearly constant down to 50 K, and then drops to a value of 6.2 emu K mol $^{-1}$ at 2 K. The fit of the high temperature $1/\chi(T)$ data yields a negative Curie–Weiss constant of $\theta = -0.62$ K (Fig. 4, inset). The decrease of the $\chi T(T)$ curve is caused by the combined effect of magnetic anisotropy (zero field splitting, ZFS) and a weak antiferromagnetic (AF)



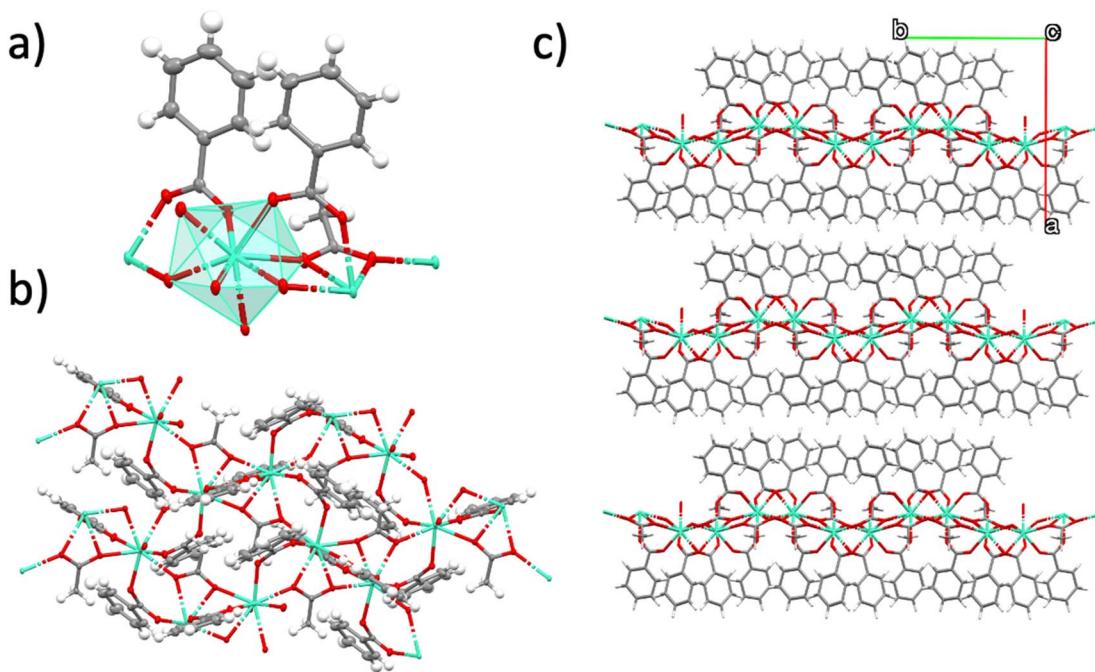


Fig. 1 Crystal structure of **1**. (a) A fragment of **1** showing thermal ellipsoids (50%) and with a polyhedral representation for Gd1; (b) top view of the 2D layer of **1**; (c) crystal packing along the *a*-axis of the unit cell. The crystal packing shows a stack of three layers of **1** with only van der Waals interactions between them. Gd: green; C: grey; O: red; H: white.

exchange; as shown in Fig. 4a, the data can be nicely fit within a dimer model for Gd(III) ions with $g = 1.985$, uniaxial anisotropy $D/k_B = -0.167$ K, $E/k_B = -0.0043$ K and weak AF coupling constant, $J/k_B = -0.036$ K, as obtained by EPR (*vide infra*). The magnetization for **1** saturates at a value of 7.0 at 2 K and 5 T

(Fig. 4b). The $M(H)$ curve falls below the theoretical Brillouin prediction for Gd(III), pointing again to the existence of AF interactions, and the data are well described within the dimeric model with the same g , anisotropy D , E , and coupling J values as above. The carboxylato ligands that bridge the Gd(III) ions in **1**

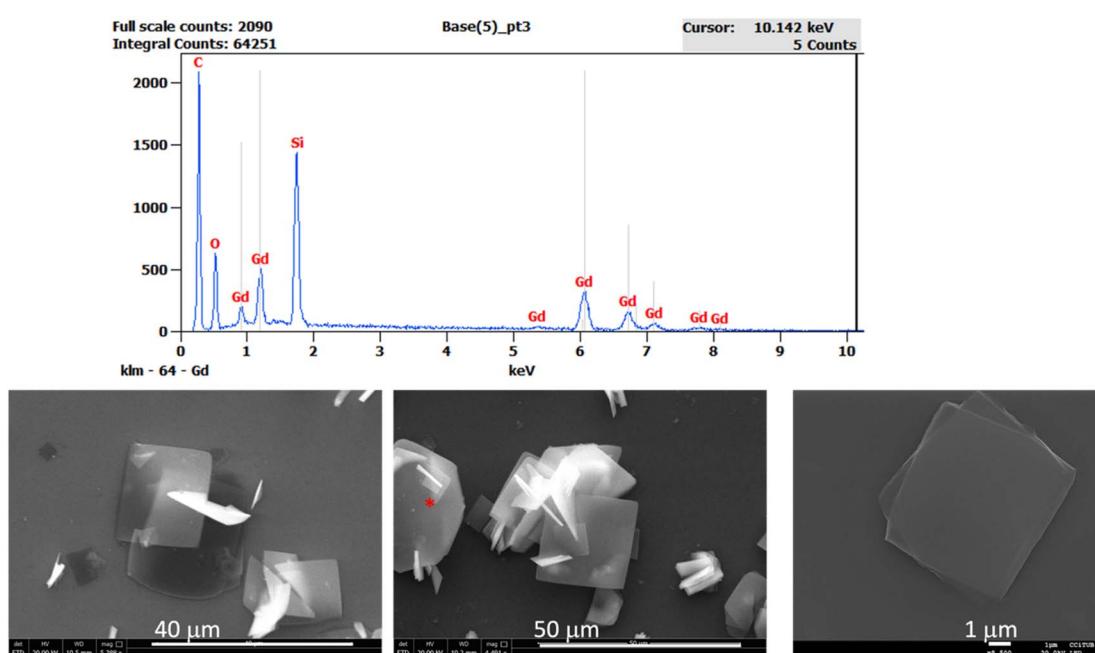


Fig. 2 SEM images for **1Si** (48 h deposition time) showing the growth of microcrystals from the functionalized Si. The EDS analysis corresponds to the flake with the red asterisk on the central image.



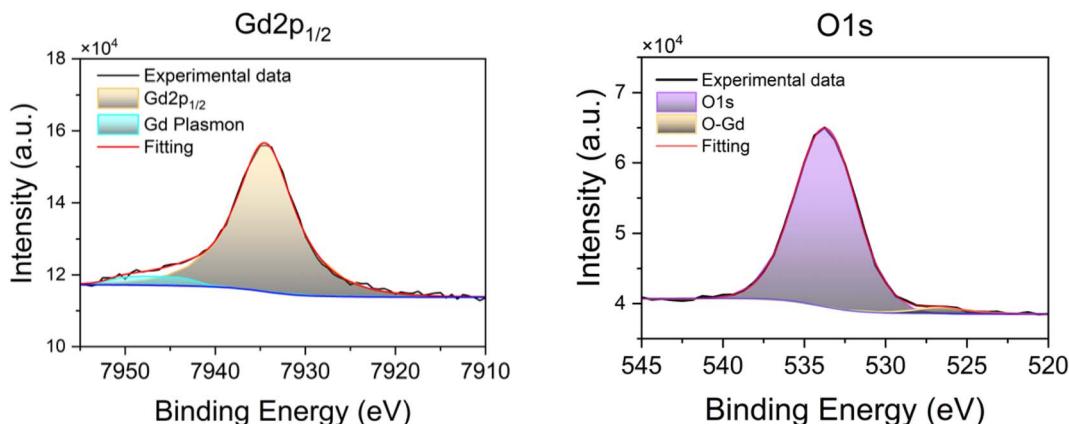


Fig. 3 HAXPES spectra for the Gd 2p_{1/2} and the O 1s core levels of 1Si.

provide the pathways for magnetic exchange, and this is typically weak for lanthanoid ions.

X-band CW EPR spectra were collected for **1** at room temperature, 80 K and 10 K (ESI Fig. S4†). The signals are broad and relatively featureless, as typical for Gd(III), since high fields are needed to resolve the full spectrum. HF-EPR experiments were therefore performed on a bulk powder sample of **1**. At the lowest temperature (5 K), multiple fine structures appear on either side of the central portion of the spectrum (Fig. S5†), indicative of measurable magnetic anisotropy. Simulations were then performed both without (Fig. S6†) and with (Fig. S7†) the inclusion of Gd···Gd exchange interactions in order to accurately determine the *g*-tensor components and the 2nd order zero-field splitting (ZFS) interactions. Either approach yielded the same parameters for the individual Gd(III) ions according to the following spin Hamiltonian, assuming a 4f⁷ electronic configuration with total spin $S = 7/2$ and orbital angular momentum $L = 0$: $\hat{H} = \mu_B \mathbf{B}_0 \cdot \tilde{\mathbf{g}} \cdot \hat{\mathbf{S}} + D\hat{S}_z^2 + E(\hat{S}_x^2 - \hat{S}_y^2)$

(eqn (1)) where μ_B is the Bohr magneton, \mathbf{B}_0 the externally applied magnetic field vector, $\tilde{\mathbf{g}}$ the Landé-tensor, and the \hat{S}_i ($i = x, y, z$) denote components of the total spin operator, $\hat{\mathbf{S}}$ with D and E parameterizing the axial and rhombic 2nd-order ZFS interactions, respectively. The parameters employed in the best simulations are: $g = 1.985 \pm 0.005$, $D/k_B = -0.167 \pm 0.003$ K, $E/k_B = -0.0043 \pm 0.0015$ K, *i.e.*, $E/D = 0.026$. In-depth discussion of the simulation procedures is given in the ESI.† Introduction of a weak Gd···Gd interaction *via* a simple dimer model (eqn S(1)†) improves the simulations. Both ferromagnetic (FM) and antiferromagnetic (AF) cases were considered, with coupling constants $J/k_B = \pm 0.036$ K ($-2J(S_i \cdot S_j)$ convention chosen for consistency with the magnetic measurements). The FM simulations give somewhat better agreement with the experimental spectra, however, given the simplicity of the employed dimer model, it is not possible to rule out an AF coupling on the basis of the EPR results, and $\chi T(T)$ data, which are more sensitive to interactions, point towards the predominance of AF

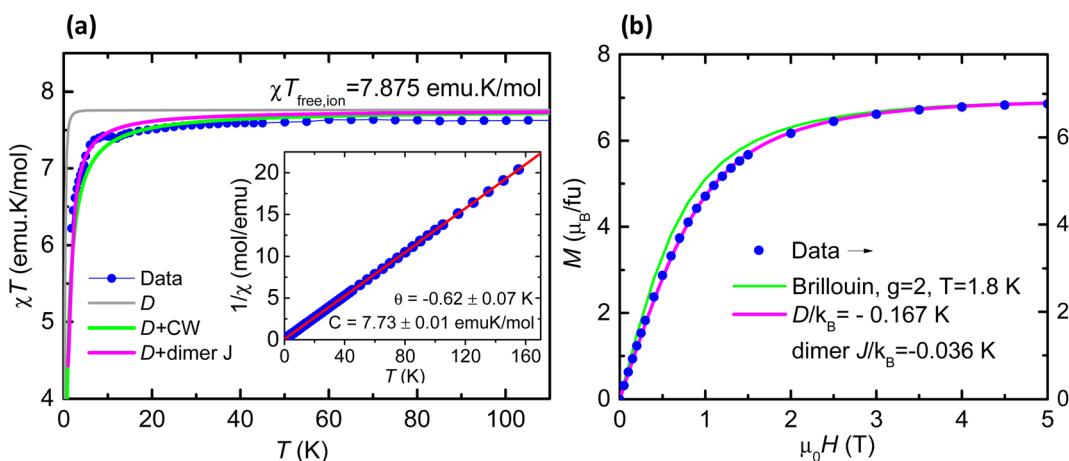


Fig. 4 (a) Susceptibility times the temperature product, χT , as a function of the temperature at 0.1 T measured for complex **1** and theoretical curves using a model: (grey line) including only ZFS, $D/k_B = -0.167$ K and $E/k_B = -0.0043$ K; (green line), ZFS and interactions within a Curie-Weiss mean-field approximation ($\theta = -0.62$ K); and (pink line), ZFS and interactions within a dimer model with AF coupling constant $J/k_B = -0.036$ K. Inset: Curie plot; (b) magnetization vs. magnetic field data plot at 1.8 K, (green line) theoretical Brillouin function for $S = 7/2$ and $g = 2$, and (pink line) theoretical prediction for ZFS and dimeric $J/k_B = -0.036$ K.



interactions. Likewise, inclusion of three (or even two) *g*-tensor components leads to an over-parameterization of the simulations. Therefore, an isotropic *g*-tensor was assumed, a choice that is justified on the grounds that the ZFS interaction is expected to dominate the low-temperature spectral fine structure pattern (see ESI†). HF-EPR has been used to constrain the single-ion spin Hamiltonian parameters and to establish magnetic anisotropy parameters.

Once the anisotropic nature of the Gd(III) ions in **1** was confirmed by HF-EPR, with an uniaxial value $D/k_B = -0.167$ K, the dynamic magnetic properties of **1** were characterized by measuring the AC magnetic susceptibility. Fig. 5a shows the out-of-phase component of the susceptibility, $\chi''(f, T)$, as a function of the frequency at constant applied magnetic field, $\mu_0 H = 0.2$ T, at different temperatures between 1.8–20 K. In addition, results of $\chi''(f, H)$ measurements performed at constant $T = 2$ K and different applied magnetic fields in the 0–1 T range are shown Fig. S8.† From these data, three main spin relaxation processes can be discerned, with relaxation times denoted τ_1 , τ_2 and τ_3 . The dependence of the relaxation time with the inverse of the temperature, $\tau(1/T)$ for the different relaxation processes is shown in Fig. 5b. The relaxation behavior of Gd(III) is very different from orthodox mononuclear SMMs or single-ion magnets (SIMs), since the anisotropy energy due to ZFS is very small, in this case $\Delta_{\text{ZFS}}/k_B \approx |D|(S^2 - 1/4) \approx 2.0$ K with $S = 7/2$, $D/k_B = -0.167$ K (-0.117 cm⁻¹), and the calculated energy difference between the lower and uppermost levels of two interacting Gd ions ($J/k_B = -0.036$ K) at $\mu_0 H = 0.2$ T is just 5.8 K.

The first slow process has relaxation times on the order of $\tau_1 \approx 0.1$ to 1 s. The $\tau_1(1/T)$ dependence at 0.2 T could be fitted to an Arrhenius law with an activation energy of $U/k_B = 6.1$ K, of the same order as the above calculated energy gap of 5.8 K. Previously, similar field-induced relaxation processes in Gd compounds were reported when subjected to comparable DC fields.^{31,32} An additional shoulder in the $\chi''(f, T)$ with a practically T -independent relaxation time of ≈ 0.02 s is observed, which may be tentatively ascribed to a direct process, although with large uncertainty as the complete χ'' peak associated with this process cannot be discriminated from the main τ_2 peak. The second relaxation process (τ_2) is temperature-dependent, but it

cannot be assigned to an Orbach mechanism, since the energy value determined from $\ln[\tau_2(1/T)]$ at 10 K, $U/k_B = 30.4$ K, is an order of magnitude larger than the anisotropy barrier. However, noting that this process follows a temperature dependence of the type $\tau_2^{-1} \propto T^2$ (see Fig. 5c), it can be explained by a Resonant Phonon Trapping (RPT) mechanism.^{31,33} This process takes place when a phonon-bottleneck effect sets in. Then the energy of the lattice modes created by the relaxing spins cannot be released into a thermal reservoir sufficiently fast. These phonons may be reabsorbed by other spins, so that the effective relaxation time becomes longer. It is noted that, for RPT to take place, the system must fulfill certain conditions.³⁰ In particular, the wavelength of the trapped phonons should be considerably larger than the typical separation between the spins, *i.e.*, $k_0 r_{\text{av}} \ll 1$, where r_{av} is the average distance between spin sites, and k_0 is the wave vector of the trapped phonon with energy $h\nu_0$. Indeed, the difference between the Zeeman split $m_S = 7/2$ and $m_S = -7/2$ energy levels in a magnetic field of 0.2 T is ~ 1.9 K. Therefore, using the Debye approximation, $h\nu_0 = c_s k_0$, and the average sound velocity obtained from the specific-heat data ($c_s = 4.77 \times 10^3$ m s⁻¹), we obtain $k_0 = 1.1 \times 10^7$ m⁻¹. Taking into account that $r_{\text{av}} \approx 7.16$ Å in this compound, considering both intra- and inter-layer distances, the condition $k_0 r_{\text{av}} \approx 8 \times 10^{-3} \ll 1$ is fulfilled. Thus, it is concluded that τ_3 may be explained in terms of the RPT mechanism. The relaxation time τ_3 is of the same order as previously reported for RPT processes in the $[\text{Gd}_2(\text{-fum})_3(\text{H}_2\text{O})_4] \cdot 3\text{H}_2\text{O}$ ³⁴ and $[\text{Gd}_2(\text{CNCH}_2\text{COO})_6(\text{H}_2\text{O})_4] \cdot 2\text{H}_2\text{O}$.³¹

Finally, the very fast relaxation process at the highest frequencies (Fig. 5a and b), with a temperature-independent relaxation time of $\tau_3 \approx 4 \times 10^{-5}$ s, may be assigned to a Quantum Tunneling of the Magnetization (QTM) process.

To further characterize the material, a suspension of flakes of **1** was spin-coated onto silicon substrates. For the exfoliation, we followed the procedure we reported for the Dy and Tb analogue.^{35,36} Optical images show partial coverage of the surface with flakes of average lateral size of approx. 3.4 μm (ESI Fig. S9a†). The magnetic properties of the flakes of **1** on a flat substrate were characterized by X-ray absorption spectroscopy (XAS) and X-ray magnetic circular dichroism (XMCD) at the Boreas beamline of the ALBA synchrotron. Fig. 6 shows the XAS

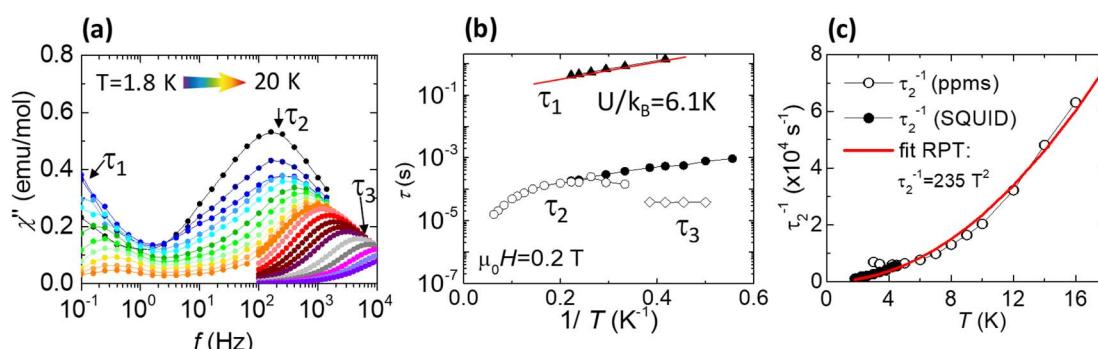


Fig. 5 (a) Imaginary component of the ac susceptibility as a function of the frequency, $\chi''(f)$, of complex **1** at $\mu_0 H = 0.2$ T and different temperatures ranging from 1.8 K to 20 K; (b) magnetic relaxation time as a function of the inverse temperature ($1/T$) at $\mu_0 H = 0.2$ T; (c) temperature dependence of $\tau_2^{-1}(T)$, and a fit to a power law $\tau_2^{-1} = KT^2$, with $K = 235$ s⁻¹ K⁻² between 2–16 K.



and XMCD spectra measured at a magnetic field of 6 T and 3.4 K at the $M_{4,5}$ edges of Gd. The XAS spectral features are typical of Gd(III): the M_5 edge shows a shoulder at 1182.4 eV, a main peak at 1184.8, and two satellites at 1189.3 and 1191.9 eV, whereas the M_4 edge exhibits two peaks at 1213.9, 1215.6 eV and a satellite at 1218.9 eV.³⁷ We determined the spin (m_s), orbital (m_L), and total magnetic moment ($m_{\text{tot}} = m_L + m_s$) per Gd(III) ion in the sample from the XAS-XMCD spectra using the corrected sum rules for lanthanides.^{38,39} We obtained values of, $m_s = 6.31 \pm 0.09 \mu_B/\text{ion}$, $m_L = 0.19 \pm 0.27 \mu_B/\text{ion}$ and $m_{\text{tot}} = 6.50 \pm 0.18 \mu_B/\text{ion}$. As expected for the $4f^7$ configuration of Gd(III) the orbital magnetic moment is vanishing though not exactly zero, as previously observed for Gd(III) complexes on surfaces.⁴⁰ The field-dependence of the total magnetic moment, $m_{\text{tot}}(H)$, was determined by following the intensity of the XMCD(H) peak at the M_5 edge between -6 T and 6 T, and scaling the curve with the value of m_{tot} at 6 T obtained from the sum rules. Fig. S9b† compares the experimental curve, $m_{\text{tot}}(H)$, for the flakes of **1** on silicon with the SQUID magnetization data for the bulk material at the same temperature, 3.4 K. The two curves coincide within the experimental error. These results demonstrate the preservation of magnetic properties of **1** upon exfoliation and surface deposition.

Magnetocaloric effect

To investigate the magnetic entropy in **1**, we studied the heat capacity (HC) as a function of temperature in the 0.3 – 30 K range, at different applied magnetic fields up to 3 T (Fig. 7a).

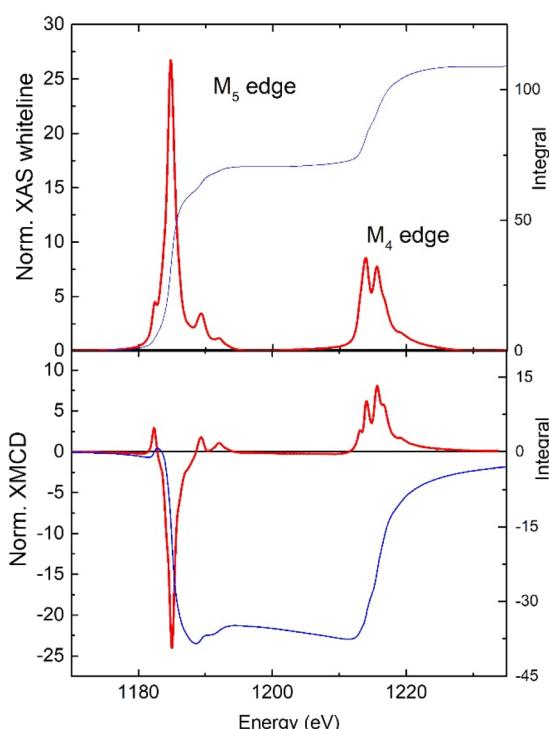


Fig. 6 Normalized, background-subtracted XAS and XMCD spectra at the $M_{4,5}$ edges of Gd measured at 3.4 K and 6 T, and the associated integrated signals (right scale).

The magnetic contribution $c_m(T,H)$ to the specific heat $c_p(T,H)$, was found after subtracting the superimposed lattice contribution, $c_L = AT^n$. The specific heat curves show a broad Schottky-like contribution that shifts towards higher temperatures with increasing field. No sign of antiferromagnetic (AF) ordering is detected down to the lowest measured temperature of 0.35 K. The complex topology of the 2D layer in **1** supports a model where Gd···Gd dinuclear units are coupled in the 2D layer. The HC experimental data are compared to the calculated HC for a randomly oriented system of Gd···Gd coupled dinuclear units, with the ZFS parameters found by EPR ($D/k_B = -0.167$ K, $E = -0.0043$ K, $g = 1.985$) and using the Hamiltonian in eqn (1) with a term introducing Gd···Gd coupling ($J/k_B = -0.036$ K). The calculated curves with these parameters are compatible with the experimental data, and nicely reproduce the observed field evolution (see Fig. 7b). However, the experimental curve at zero field, as compared to the theoretical prediction, exhibits a broader contribution shifted to lower temperatures, suggesting that in this low temperature range the model used may be too simple and that additional magnetic interactions are involved.

The magnetocaloric effect (MCE) performance was evaluated from the heat capacity data.⁴¹ From the HC data of Fig. 7a we firstly determine by the total entropies at applied magnetic fields $H_l = 0$ and H_h as a functions of temperature using the equation:

$$S(T)_H = \int_0^T \frac{C(T)_H}{T} dT \quad (1)$$

Then we calculate the magnetic entropy change and adiabatic temperature change for that field change, $\Delta B = (H_l - 0)$ T, as $\Delta S_m(T)_{\Delta H} = [S(T)_{Hh} - S(T)_0]_T$ and $\Delta T_{\text{ad}}(T)_{\Delta H} = [T(S)_{Hh} - T(S)_0]_S$, respectively. Fig. 7c shows the magnetic entropy change, $-\Delta S_m(T, \Delta B)$, for selected field changes, $\mu_0 \Delta H = 0.5, 1, 2, 3$ T. For practical purposes it is useful to consider the value of the magnetic entropy change for a magnetic field of 2 T that is easily achieved with permanent magnets. Thus, for $\mu_0 \Delta H = 2$ T, the maximum value is found to be $-\Delta S_m^{\text{max}} = 1.22$ R (22.1 J K $^{-1}$ kg $^{-1}$) at $T_{\text{max}} = 0.79$ K. This is 59% of the maximum magnetic entropy available for the system: $S_{\text{max}} = R \ln(2S + 1) = 38$ J K $^{-1}$ kg $^{-1}$. The relative cooling power (RCP) is defined as $\text{RCP} = |\Delta S_m - \Delta S_m^{\text{max}}| / \delta T_{\text{FWHM}}$.⁴² For **1**, The full width at half maximum of the ΔS_m curve is $\delta T_{\text{FWHM}} = 2.48$ K and $\text{RCP}(1) = 54.8$ J kg $^{-1}$. The determined adiabatic temperature change as a function of the final temperature after the cooling, $\Delta T_{\text{ad}}(T)$, is shown in Fig. 7d. For $\mu_0 \Delta H = 2$ – 0 T this temperature change reaches a maximum of $\Delta T_{\text{ad}}^{\text{max}} = 6.45$ K at $T = 1.76$ K and drops to zero above 19 K, thus the material would be best suited as coolant at cryogenic temperatures.

Alternatively, the MCE can be characterized by collecting the magnetization data at various temperatures. Maxwell's equations can be used to extract the magnetic entropy change, $-\Delta S_m$.^{43,44} Magnetization isothermal curves at temperatures from 2 K to 10 K and fields from 0 to 7 T were collected for **1**, and used to determine $-\Delta S_m(T)$ (see ESI Fig. S10†). The magnetic entropy change value at 2 K and 7 T is $-\Delta S_m(7$ T) = 35 J K $^{-1}$



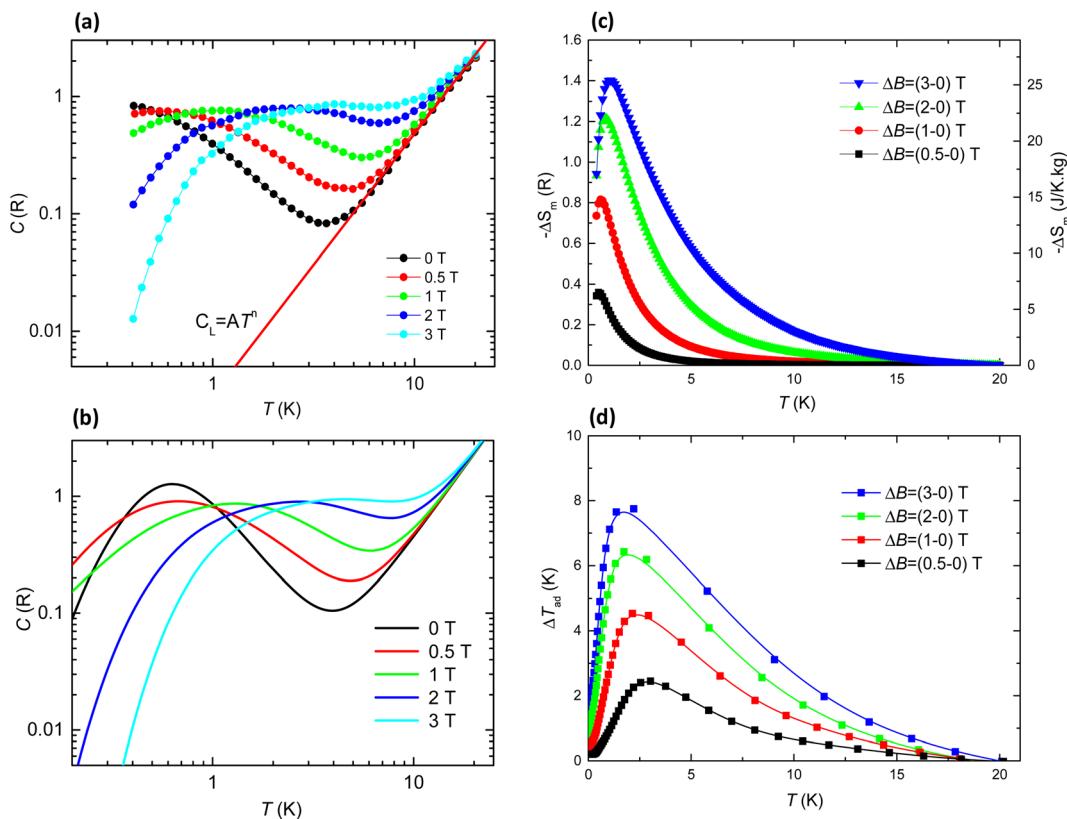


Fig. 7 (a) Specific heat as a function of temperature at the indicated magnetic fields, $c_p(T,H) = c_L + c_m$, for Gd 2D MOF 1 . Red line: lattice contribution to the HC, $c_L = AT^n$, with $A = 1.9 \times 10^{-3} \text{ R K}^{-2}$,³⁷ $n = 2.27$; (b) simulated $c_p(T,H)$ curves within a model of interacting Gd ions with EPR-obtained values $g = 1.985$, $D/k_B = -0.167 \text{ K}$, $E/k_B = -0.0043 \text{ K}$, and $\text{Gd}\cdots\text{Gd}$ coupling constant $J/k_B = -0.036 \text{ K}$; (c) temperature dependence of the magnetic entropy change, $-\Delta S_m$, and (d) temperature dependence of the adiabatic temperature change, $\Delta T_{\text{ad}}(T)$, as obtained from the heat capacity measurements for the indicated applied magnetic field changes, $\mu_0\Delta H$, where T denotes the final temperature of the adiabatic cooling.

kg^{-1} , which is 98% of the maximum theoretical entropy, S_{max} . At 2 K and 2 T the determined value, $-\Delta S_m(2 \text{ T}) = 22.3 \text{ J K}^{-1} \text{ kg}^{-1}$, is in good agreement with the value found from heat capacity. The peak in the $-\Delta S_m(T)$ curve that was observable in the data obtained from HC is not detected here, since magnetization measurements were performed only down to 2 K. The MCE performance of 4f and 3d molecular compounds of various dimensionalities (0D, 1D, 2D and 3D) has been recently reviewed,²³ and the value of $-\Delta S_m$ at 7 T in units of $\text{J K}^{-1} \text{ kg}^{-1}$ is proposed as a figure of merit. The lowest entropy values correspond to compounds with large organic ligands that result in a small density of Gd(III) ions in the material, while the large entropy values are for materials with formate, oxalate or other ligands such as sulfate.⁴⁵⁻⁴⁹ The $-\Delta S_m(7 \text{ T})$ of **1** is similar to the reported 2D compounds, with values that range from $27 \text{ J K}^{-1} \text{ kg}^{-1}$ (ref. 50) to $49.29 \text{ J K}^{-1} \text{ kg}^{-1}$.⁵¹

A great advantage of compound **1** is that it can be exfoliated and deposited or grafted onto a surface, as we have shown here. This is one requirement for the application of molecular materials as magnetic refrigerators. Evangelisti and co-workers showed already that molecular Gd(III) species could be confined on a surface and retain their effective magnetic cooling properties. They propose the use of Gd MOFs in on-chip micro-refrigerators, by depositing the MOF on a functionalized silicon

membrane. As indicated, aliphatic linkers used to graft the MOF on the silicon chip should not hinder the conduction of heat.^{24,44} Our XAS and XMCD data already shows that the flakes of **1** have unchanged magnetic properties. Thus, we studied the magnetocaloric effect of **1Si**. Magnetization *vs.* field curves at various temperatures were collected on a SQUID magnetometer in order to ascertain the magnetic entropy change of **1Si** (Fig. 8b and S11†). The Si wafer with crystals of **1** grown over it, **1Si**, was positioned perpendicular and parallel to the field in the SQUID. When compared to the bulk **1**, the $-\Delta S_m(T)$ curves show slightly lower values of entropy at 2 T, as well as a the maximum shifted to higher temperatures than for **1**. The magnetic entropy change for **1Si** was determined to be $-\Delta S_m(7 \text{ T}) = 31.8 \text{ J K}^{-1} \text{ kg}^{-1}$, or $-\Delta S_m(2 \text{ T}) = 17.2 \text{ J K}^{-1} \text{ kg}^{-1}$ (perpendicular) and $-\Delta S_m(7 \text{ T}) = 35.6 \text{ J K}^{-1} \text{ kg}^{-1}$, or $-\Delta S_m(2 \text{ T}) = 16.6 \text{ J K}^{-1} \text{ kg}^{-1}$ (parallel). By fitting the experimental entropy data to a peak using a logarithmic peak function (solid line in Fig. 8), we can obtain a rough estimate of δT_{FWHM} at 2 T as 4.4 K (perpendicular) or 4.7 K (parallel), thus $\text{RCP}(\mathbf{1Si}) = 77.7 \text{ J kg}^{-1}$ (perpendicular) or 75.0 J kg^{-1} (parallel). These approximate values are of the same order of the RCP for the bulk sample, **1**, and one order of magnitude lower than the reported RCP for bulk Gd.⁵² The small differences between the parallel and perpendicular **1Si** entropy data can be attributed to the anisotropic shape of the sample. The thin layer



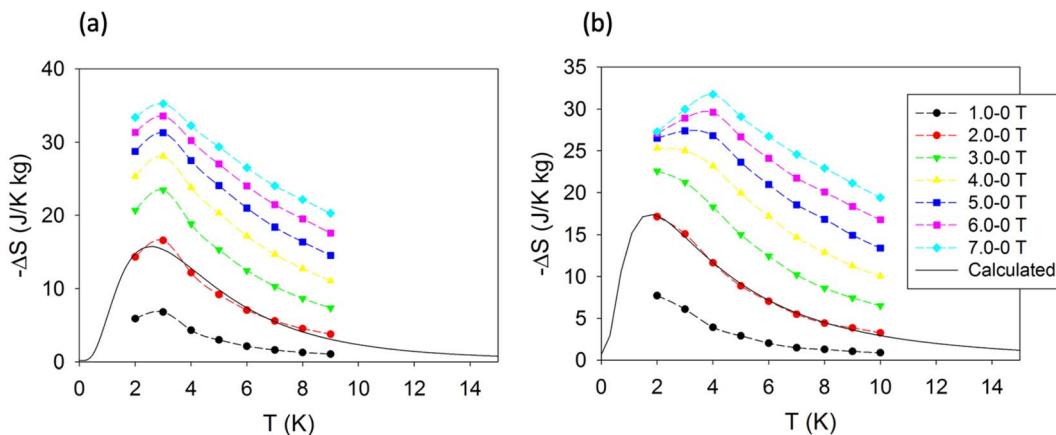


Fig. 8 Magnetic entropy change, $-\Delta S_m$, for **1** grafted on a TSP functionalized Si wafer, **1Si**, with the wafer parallel (a) and perpendicular (b) to the SQUID magnetic field. The solid line is a mathematical fit to a logarithmic peak function.

1Si is formed by micro and nanocrystals of **1** grafted on Si. The crystals are anisotropic in shape and orientation with respect to the field in the SQUID and polydisperse in size. The covalent bond to the Si-COO⁻ groups is not equal for all crystals, as evidenced by the SEM images in Fig. 2 and by AFM (ESI Fig. S12†), as a large number of small crystals grow flat on the substrate (ESI Fig. S11c†) but other, larger crystals grow at random angles from the substrate (SEM images in Fig. 2). The attachment of the crystals to the substrate can cause different strains in the crystals that along with the polydispersity in size can be at the origin of the broader maximum and the shift to higher temperatures of the maximum of entropy change.

Previously, thin films of oxides and metals had been proposed for their potential for cryogenic magnetic cooling of microdevices.^{42,53} We have shown that 2D MOF [Gd(MeCOO)(PhCOO)₂] (**1**) can be grown on functionalized silicon wafers (**1Si**), and it shows a MCE similar to that of the bulk compound, opening up applications of this material as an effective cryogenic magnetic cooler for microdevices.

Conclusions

The family of [Ln(MeCOO)(PhCOO)₂] van der Waals 2D MOFs offers a truly versatile system that can be used to access properties like SMM behavior (Dy, Tb), and a combination of SMM behavior and luminescence, including Ln-Ln' energy transfer (TbEu). With the Gd analogue [Gd(MeCOO)(PhCOO)₂] **1** reported here, a large magnetic entropy change and strong magnetocaloric effect are added to the properties of this versatile family of materials. The 2D MOF **1** shows a large magnetic entropy change at the technologically relevant field of 2 T, as we show by heat capacity and magnetization measurements. The paramagnetic Gd complex is anisotropic, as demonstrated by HF-EPR. Under the application of a magnetic field, different slow spin relaxation pathways were observed, the two main ones ascribed, respectively, to a thermally activated process between ground and excited states, favored by the lifting of the Kramers degeneracy by the external field, and a resonant phonon trapping (RTP) mechanism.

Furthermore, **1** can be exfoliated into monolayers and flakes, by using a simple ultrasound bath. This makes these materials amenable to the preparation of heterostructures. The flakes can be physiosorbed on surfaces by spin coating without affecting the magnetic properties of the material, as we demonstrate here by XAS and XMCD. We show how these 2D MOFs can be integrated in devices, **1** can be chemisorbed on suitably functionalized surfaces (**1Si**) to attain thin layers of the material that are functional in showing strong magnetic entropy change and MCE performance with $RCP(1Si) = 77.7 \text{ J kg}^{-1}$. This opens the possibility of using this material as an on-surface cryogenic coolant for microdevices.

Author contributions

SK: investigation, formal analysis, data curation (synthesis and characterization, SEM); GGR: investigation, formal analysis, data curation (synthesis, AFM, HAXPES, CW-EPR); AA: investigation, formal analysis, data curation, writing – review and edit (heat capacity, magnetic data); JH, SH: investigation, formal analysis, data curation, writing – review and edit (HFEPR); LB: formal analysis and resources (CW-EPR); JRZ: investigation, writing – review and edit (HAXPES); JJ: investigation, writing – review and edit (theoretical calculations SHAPE and RAMAN); EB: investigation, formal analysis, data curation, visualization, writing – original draft, review and edit (heat capacity, magnetic data, XMCD); ECS: conceptualization, supervision, funding acquisition, resources, investigation, formal analysis, data curation, visualization, writing – original draft, review and edit (SCXRD, HAXPES, magnetism).

Conflicts of interest

There are no conflicts to declare.

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References

- 1 E. Coronado, *Nat. Rev. Mater.*, 2020, **5**, 87–104.
- 2 L. A. Benítez, J. F. Sierra, W. Savero Torres, A. Arrighi, F. Bonell, M. V. Costache and S. O. Valenzuela, *Nat. Phys.*, 2018, **14**, 303–308.
- 3 A. Gaita-Ariño, F. Luis, S. Hill and E. Coronado, *Nat. Chem.*, 2019, **11**, 301–309.
- 4 J. López-Cabrelles, L. Escalera-Moreno, Z. Hu, H. Prima-García, G. Mínguez- Espallargas, A. Gaita-Ariño and E. Coronado, *Inorg. Chem.*, 2021, **60**, 8575–8580.
- 5 F. Luis, P. J. Alonso, O. Roubeau, V. Velasco, D. Zueco, D. Aguilà, J. I. Martínez, L. A. Barrios and G. Aromí, *Commun. Chem.*, 2020, **3**, 176.
- 6 M. Evangelisti, A. Candini, A. Ghirri, M. Affronte, E. K. Brechin and E. J. L. McInnes, *Appl. Phys. Lett.*, 2005, **87**, 072504.
- 7 I. Cimatti, L. Bondì, G. Serrano, L. Malavolti, B. Cortigiani, E. Velez-Fort, D. Betto, A. Ouerghi, N. B. Brookes, S. Loth, M. Mannini, F. Totti and R. Sessoli, *Nanoscale Horiz.*, 2019, **4**, 1202–1210.
- 8 A. Nurbawono and C. Zhang, *Appl. Phys. Lett.*, 2013, **103**, 4.
- 9 S. Marocchi, A. Candini, D. Klar, W. Van den Heuvel, H. Huang, F. Troiani, V. Corradini, R. Biagi, V. De Renzi, S. Klyatskaya, K. Kummer, N. B. Brookes, M. Ruben, H. Wende, U. del Pennino, A. Soncini, M. Affronte and V. Bellini, *ACS Nano*, 2016, 9353–9360.
- 10 A. Candini, S. Klyatskaya, M. Ruben, W. Wernsdorfer and M. Affronte, *Nano Lett.*, 2011, **11**, 2634–2639.
- 11 L. Gragnaniello, F. Paschke, P. Erler, P. Schmitt, N. Barth, S. Simon, H. Brune, S. Rusponi and M. Fonin, *Nano Lett.*, 2017, **17**, 7177–7182.
- 12 V. Corradini, A. Candini, D. Klar, R. Biagi, V. De Renzi, A. L. Rizzini, N. Cavani, U. Del Pennino, S. Klyatskaya, M. Ruben, E. Velez-Fort, K. Kummer, N. B. Brookes, P. Gargiani, H. Wende and M. Affronte, *Nanoscale*, 2017, **10**, 277–283.
- 13 M. Zhao, Y. Wang, Q. Ma, Y. Huang, X. Zhang, J. Ping, Z. Zhang, Q. Lu, Y. Yu, H. Xu, Y. Zhao and H. Zhang, *Adv. Mater.*, 2015, **27**, 7372–7378.
- 14 H. Zhua and D. Liu, *J. Mater. Chem. A*, 2019, **7**, 21004–21035.
- 15 G. Gabarró-Riera, G. Aromí and E. C. Sañudo, *Coord. Chem. Rev.*, 2023, **475**, 214858.
- 16 A. Urtizberea, E. Natividad, P. J. Alonso, M. A. Andrés, I. Gascón, M. Goldmann and O. Roubeau, *Adv. Funct. Mater.*, 2018, **28**, 1–15.
- 17 A. Urtizberea, E. Natividad, P. J. Alonso, L. Pérez-Martínez, M. A. Andrés, I. Gascón, I. Gimeno, F. Luis and O. Roubeau, *Mater. Horiz.*, 2020, **7**, 885–897.
- 18 J. González, P. Sevilla, G. Gabarró-Riera, J. Jover, J. Echeverría, S. Fuertes, A. Arauzo, E. Bartolomé and E. C. Sañudo, *Angew. Chem., Int. Ed.*, 2021, **60**, 12001–12006.
- 19 E. Bartolomé, A. Arauzo, S. Herce, A. Palau, N. Mestres, S. Fuertes, P. Sevilla, N. S. Settineri, L. Navarro-Spreafico, J. González and E. C. Sañudo, *Molecules*, 2021, **26**(18), 5503.
- 20 E. Bartolomé, A. Arauzo, S. Fuertes, L. Navarro-Spreafico, P. Sevilla, H. F. Cortés, N. Settineri, S. J. Teat and E. C. Sañudo, *Dalton Trans.*, 2023, **52**, 7258–7270.
- 21 V. K. Pecharsky and K. G. Gschneider Jr, *J. Magn. Magn. Mater.*, 1999, **200**, 44–56.
- 22 J. Romero Gómez, R. Ferreiro Garcia, A. de Miguel Catoira and M. Romero Gómez, *Renewable Sustainable Energy Rev.*, 2013, **17**, 74–82.
- 23 P. Konieczny, W. Sas, D. Czernia, A. Pacanowska, M. Fitta and R. Pelka, *Dalton Trans.*, 2022, **51**, 12762–12780.
- 24 G. Lorusso, E. Natividad, M. Evangelisti and O. Roubeau, *Mater. Horiz.*, 2019, **6**, 144–154.
- 25 G. Lorusso, M. Jenkins, P. González-Monje, A. Arauzo, J. Sesé, D. Ruiz-Molina, O. Roubeau and M. Evangelisti, *Adv. Mater.*, 2013, **25**, 2984–2988.
- 26 M. Godoy-Gallardo, M. C. Manzanares-Céspedes, P. Sevilla, J. Nart, N. Manzanares, J. M. Manero, F. J. Gil, S. K. Boyd and D. Rodríguez, *Mater. Sci. Eng. C*, 2016, **69**, 538–545.
- 27 G. Gabarró-Riera, J. Jover, J. R. Zuazo, E. Bartolomé and E. C. Sañudo, *Inorg. Chem. Front.*, 2022, **9**, 4160–4170.
- 28 S. Stoll and A. Schweiger, *J. Magn. Reson.*, 2006, **178**, 42–55.
- 29 J. R. Rubio-Zuazo and G. R. Castro, *Nucl. Instrum. Methods Phys. Res., Sect. A*, 2005, **547**, 64–72.
- 30 J. Rubio-Zuazo, M. Escher, M. Merkel and G. R. Castro, *Rev. Sci. Instrum.*, 2010, **81**, 043304.
- 31 A. Arauzo, A. Lazarescu, S. Shova, E. Bartolomé, R. Cases, J. Luzón, J. Bartolomé and C. Turta, *Dalton Trans.*, 2014, **43**, 12342–12356.
- 32 B. Liu, B.-W. Wang, Z.-M. Wang and S. Gao, *Sci. China: Chem.*, 2012, **5**, 926.
- 33 R. R. Schenker, M. N. Leuenberger, G. Chaboussant, D. Loss and H. U. Güdel, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2005, **72**, 184403–184413.
- 34 M. Orendáč, L. Sedláková, E. Čižmár, a. Orendáčová, a. Feher, W. H. Zhu, Z. M. Wang and S. Gao, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2010, **81**, 1–8.
- 35 J. González, P. Sevilla, G. Gabarró-Riera, J. Jover, J. Echeverría, S. Fuertes, A. Arauzo, E. Bartolomé and E. C. Sañudo, *Angew. Chem., Int. Ed.*, 2021, **60**, 12001–12006.



36 E. Bartolomé, A. Arauzo, S. Herce, A. Palau, N. Mestres, S. Fuertes, P. Sevilla, N. S. Settineri, L. Navarro-Spreafico, J. González and E. C. Sañudo, *Molecules*, 2021, **26**, 1–14.

37 G. Champion, N. Lalioti, V. Tangoulis, M.-A. Arrio, P. Sainctavit, F. Villain, A. Caneschi, D. Gatteschi, C. Giorgetti, F. Baudelet, M. Verdaguer and C. Cartier dit Moulin, *J. Am. Chem. Soc.*, 2003, **125**, 8371–8376.

38 B. Thole, P. Carra, F. Sette and G. van der Laan, *Phys. Rev. Lett.*, 1992, **68**, 1943–1946.

39 P. Carra, B. T. Thole, M. Altarelli and X. Wang, *Phys. Rev. Lett.*, 1993, **70**, 694–697.

40 P. Stoll, M. Bernien, D. Rolf, F. Nickel, Q. Xu, C. Hartmann, T. R. Umbach, J. Kopprasch, J. N. Ladenthin, E. Schierle, E. Weschke, C. Czekelius, W. Kuch and K. J. Franke, *Phys. Rev. B*, 2016, **94**, 224426.

41 M. Evangelisti, A. Candini, A. Ghirri, M. Affronte, S. Piligkos, E. K. Brechin and E. J. L. McInnes, *Polyhedron*, 2005, **24**, 2573–2578.

42 J. H. Jia, Y. J. Ke, X. Li, H. R. Zhang, Z. P. Yu, Z. H. Cheng, K. Zhai, Z. Y. Liu and J. F. Wang, *J. Mater. Chem. C*, 2019, **7**, 14970–14976.

43 M. Evangelisti and E. K. Brechinb, *Dalton Trans.*, 2010, **39**, 4672–4676.

44 G. Lorusso, M. Jenkins, P. González-Monje, A. Arauzo, J. Sesé, D. Ruiz-Molina, O. Roubeau and M. Evangelisti, *Adv. Mater.*, 2013, **25**, 2984–2988.

45 S. P. Bera, A. Mondal and S. Konar, *Chem.-Asian J.*, 2019, **14**, 3702–3711.

46 T.-Q. Song, J. Dong, A.-F. Yang, X.-J. Che, H.-L. Gao, J.-Z. Cui and B. Zhao, *Inorg. Chem.*, 2018, **57**, 3144–3150.

47 S.-J. Liu, C. Cao, S.-L. Yao, T.-F. Zheng, Z.-X. Wang, C. Liu, J.-S. Liao, J.-L. Chen, Y.-W. Li and H.-R. Wen, *Dalton Trans.*, 2017, **46**, 64–70.

48 S.-D. Han, J.-H. Li, H.-H. Liua and G.-M. Wang, *Dalton Trans.*, 2017, **46**, 10023–10028.

49 G. Lorusso, M. A. Palacios, G. S. Nichol, E. K. Brechin, M. Evangelisti and O. Roubeau, *Chem. Commun.*, 2012, **48**, 7592.

50 W.-W. Wei, L.-P. Lu, S.-S. Feng, M.-L. Zhu and U. Englerta, *Acta Crystallogr., Sect. C: Struct. Chem.*, 2021, **77**, 591–598.

51 S. P. Bera, A. Mondal and S. Konar, *Chem.-Asian J.*, 2019, **14**, 3702–3711.

52 C. W. Miller, D. V. Williams, N. S. Bingham and H. Srikanth, *J. Appl. Phys.*, 2010, 107.

53 W. Sas and M. Fitta, *J. Magn. Magn. Mater.*, 2023, **575**, 170719.

