



Cite this: *J. Mater. Chem. C*,
2024, 12, 13332

Alumina nanoparticles enable optimal spray-coated perovskite thin film growth on self-assembled monolayers for efficient and reproducible photovoltaics†

Elena J. Cassella,^{‡a} Robert D.J. Oliver,^{id} ^{‡b} Timothy Thornber,^a Sophie Tucker,^{ac} Rehmat Goodwin,^{ac} David G. Lidzey^{*a} and Alexandra J. Ramadan^{id} ^{*a}

The power conversion efficiencies of metal halide perovskite photovoltaics have increased rapidly over the past decade attracting significant academic and industrial interest. The ease with which high performance perovskite photovoltaics can be fabricated through solution processing routes has opened up significant possibilities for fabrication through existing, industrially mature high-throughput solution coating techniques such as spray-coating. The power conversion efficiencies of spray-coated metal halide perovskite photovoltaics are limited by non-radiative recombination at the interfaces with charge transport layers necessitating the implementation of new charge transport layers. The self-assembled monolayer (SAM) charge transport layers have resulted in record perovskite photovoltaic device performances, due to reduced non-radiative recombination. However, poor wettability associated with some SAMs significantly limits their applicability, this is exaggerated for droplet-based scalable technologies like spray-coating. Here we report an optimised aluminium oxide nanoparticle interlayer which enables spray-coating of triple cation metal halide perovskite thin films and devices onto Me-4PACz ([4-(3,6-dimethyl-9H-carbazol-9-yl)butyl]phosphonic acid). Our interlayer results in improved structural and optoelectronic properties of spray-coated perovskite thin films, compared to those fabricated through spin-coating. These improved properties enable the fabrication of p-i-n photovoltaic devices with efficiencies over 20% – some of the highest reported for both spray-coated devices in p-i-n architecture, and devices having a spray-coated “triple cation” perovskite active layer.

Received 18th April 2024,
Accepted 10th July 2024

DOI: 10.1039/d4tc01602e

rsc.li/materials-c

Introduction

Metal halide perovskite semiconductors have attracted considerable research interest as they possess excellent optoelectronic properties^{1–4} and can be fabricated through a range of solution and evaporation based processing routes.⁵ In combination this makes them incredibly versatile and attractive for a range of optoelectronic device applications.^{6–9} The solution processability of metal halide perovskites allows for their fabrication through a variety of deposition techniques compatible with high-throughput, roll-to-roll manufacturing processes. Of these

techniques, ultrasonic spray-coating is an industrially mature option for which exceptionally high speeds and material usage have been demonstrated.^{10,11} Ultrasonic spray-coating therefore represents an attractive route towards the realisation of high-throughput, low-cost, production of perovskite photovoltaics, a critical element in the drive towards commercialisation.

Interfacial recombination is a major energetic loss pathway for all perovskite photovoltaics, regardless of processing method, constraining their open-circuit voltage.¹² A key driver of interfacial recombination is mismatch in the energetic alignment at the interfaces between the perovskite semiconducting layer and electron and hole charge transport layers.^{13,14} In an effort to overcome this, new charge transporting layers with a range of energy level positions have been developed and explored in photovoltaic devices. A successful example of this is the carbazole-based self-assembled monolayers (SAMs) which have boosted reported performances of perovskite photovoltaics by significantly reducing the non-radiative recombination at the perovskite/hole transport layer interface in p-i-n photovoltaics.^{15,16} Whilst these SAM materials have resulted in a substantial increase in device performance and stability,^{17,18} several challenges remain.

^a Department of Physics and Astronomy, University of Sheffield, Hicks Building, Hounsfield Road, Sheffield S3 7RH, UK. E-mail: a.ramadan@sheffield.ac.uk, d.g.lidzey@sheffield.ac.uk

^b Department of Materials Science and Engineering, University of Sheffield, Sir Robert Hadfield Building, Mappin Street, Sheffield, S1 3JD, UK

^c Grantham Centre for Sustainable Futures, University of Sheffield, Sheffield S10 2TN, UK

† Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d4tc01602e>

‡ These authors contributed equally.



Solution processing of a perovskite layer atop these SAMs can prove extremely challenging due to inadequate wetting of the perovskite.^{19,20} This effect arises from the polarity mismatch between the underlying charge-transporting layer and the perovskite precursor solvent.^{21,22} This impacts the reproducibility of photovoltaic device fabrication and hinders the implementation of these SAMs particularly for scalable solution coating techniques.

To mitigate these wetting issues, a wide number of strategies have been adopted to improve the surface coverage of perovskite films onto these materials. These include solvent engineering,^{19,23,24} interfacial modification, co-deposition of two SAMs,^{20,25} co-deposition with a semiconducting polymer,²⁶ co-deposition with the perovskite precursor solution,²⁷ and thermal evaporation of the SAM layer.²⁸ Interfacial modification of the SAM layer through the introduction of an interlayer has proven particularly successful in its broad applicability to a range of perovskite compositions.^{29–32} Furthermore, a wide range of potential interlayers can be employed such as semiconducting polymers³³ or metal oxides^{34–36} and deposited through a range of fabrication routes.

Critically, few studies have explored the applicability of such wetting-enhancement strategies to commercially-relevant, large-area deposition technologies. This is particularly important for droplet-based methods such as spray-coating, which present unique film formation challenges compared to spin-coating or meniscus-assisted techniques such as slot-die or blade coating. Centrifugal forces distribute solutions evenly across substrates during spin-coating, whilst meniscus-assisted coating leverages capillary action to facilitate uniform thin film formation. Contrastingly, droplet-based methods deliver a droplet mist of the coating solution to the surface which must then coalesce. The surface free energy of the substrate is therefore crucial in determining the wetting behaviour and subsequent film formation.^{37,38} A high contact angle between the spray-coated droplets and the substrate means that

the droplets tend to bead up rather than spread out and coalesce. This can lead to the formation of gaps, pinholes, or discontinuities in the resulting perovskite thin film. Despite additional film formation challenges, spray-coating is a scalable solution processing technique capable of combining extremely high deposition speeds with the ability to coat three-dimensional, non-planar structures.^{39,40}

Here, we report an alumina nanoparticle surface modification interlayer to facilitate spray-coating of metal halide perovskite thin films onto one of the most commonly implemented carbazole-based self-assembled monolayers: [4-(3,6-dimethyl-9H-carbazol-9-yl)butyl]phosphonic acid (Me-4PACz). We present a comparison of the structural and optoelectronic properties of our spray-coated metal halide perovskite systems with equivalent spin-coated counterparts and demonstrate our optimised alumina interlayer results in spray-coated perovskite thin films which have improved optoelectronic properties relative to their spin-coated counterparts. With our fully optimised interlayer we report spray-coated *p-i-n* photovoltaics based on a 1.62 eV triple cation perovskite, $\text{Cs}_{0.05}\text{FA}_{0.81}\text{MA}_{0.14}\text{PbI}_{2.55}\text{Br}_{0.45}$ achieving a champion power conversion efficiency (PCE) of 20.4%. This represents the current state of the art for spray-coated *p-i-n* metal halide perovskite photovoltaics (Tables S1 and S2, ESI†).

Experimental

Materials

All materials were used as received without any further purification. PbI_2 and PbBr_2 were purchased from Tokyo Chemical Industry (TCI). Me-4PACz (> 99.99%) was purchased from Luminescence Technology Corp. Methylammonium bromide, formamidinium iodide, and 20×15 mm and 25×25 mm pre-patterned indium tin oxide substrates ($\approx 20 \Omega \square^{-1}$) were purchased from Ossila. Aluminium oxide nanoparticles (< 50 nm particle size (DLS), 20 wt% in isopropanol), C_{60} (sublimed, 99.95%), bathocuproine (BCP, sublimed, 99.99%), CsI (99.999% trace metals basis), and all solvents were purchased from Sigma Aldrich unless otherwise stated.

Thin film and device fabrication

ITO substrates were cleaned by sequential sonication in diluted Hellmanex solution ($\sim 1\%$ in deionised water), deionized water, acetone, and isopropanol (IPA). The substrates were UV ozone treated for 15 minutes immediately prior to thin film and device processing. All thin film and device processing was performed within an N_2 filled glovebox.

Me-4PACz solution (1 mmol, ethanol) was statically spin-coated (60 μL and 100 μL for small- and large-area substrates, respectively) onto the substrate for 30 s at 3000 rpm in an N_2 filled glovebox before annealing on a hotplate at 100°C for 10 minutes. No subsequent rinsing steps were applied. The Al_2O_3 nanoparticle solution (~ 0.1 wt%, isopropanol) was made by diluting a stock 20 wt% Al_2O_3 nanoparticle solution in isopropanol in a ratio of either 1 stock:200 neat IPA, 1 stock:100 neat IPA or 1 stock:50 neat IPA. The diluted solution was



Alexandra J. Ramadan

Alexandra Ramadan is a Lecturer in the Department of Physics and Astronomy at the University of Sheffield. She received her PhD from Imperial College London in 2016 where she worked under the supervision of Professor Sandrine Heutz (2016). In 2017 she received a UK-US Fulbright Scholarship to study highly crystalline organic solar cells at Princeton University. She also worked as a postdoctoral researcher at the University of Oxford with Professor Henry Snaith FRS on perovskite semiconductors and their optoelectronic devices. In 2022 she was awarded the Rank Prize Research Visionary Award in optoelectronics. Her research focusses on perovskite semiconductors for next generation optoelectronic devices.

Alexandra J. Ramadan is a Lecturer in the Department of Physics and Astronomy at the University of Sheffield. She received her PhD from Imperial College London in 2016 where she worked under the supervision of Professor Sandrine Heutz (2016). In 2017 she received a UK-US Fulbright Scholarship to study highly crystalline organic solar cells at Princeton University. She also worked as a postdoctoral researcher at the University of Oxford with Professor Henry Snaith FRS on perovskite semiconductors and their optoelectronic devices. In 2022 she was awarded the Rank Prize Research Visionary Award in optoelectronics. Her research focusses on perovskite semiconductors for next generation optoelectronic devices.



extensively vortex-mixed prior to further use. The Al_2O_3 nanoparticles were spin-coated (60 μL and 100 μL for small- and large-area substrates, respectively) onto SAM thin films at 2000 rpm for 30 s, before annealing on a hotplate at 100 $^\circ\text{C}$ for 1 minute. A 1.2 M solution of $\text{Cs}_{0.05}\text{FA}_{0.81}\text{MA}_{0.14}\text{PbI}_{2.55}\text{Br}_{0.45}$ was prepared in a 4:1 mixture of dimethylformamide (DMF):dimethyl sulfoxide (DMSO). The resulting solution was filtered through a 0.2 μm PTFE filter. The precursor solution (60 μL and 100 μL for small- and large-area substrates, respectively) was spin-coated at 2000 rpm for 10 s (200 rpm s^{-1} acceleration) then 4000 rpm for 20 s. An anisole anti-solvent quench (100 μL and 300 μL for small- and large-area substrates, respectively) was employed 10 s before the end of the full spin program. The films were then annealed on a hotplate at 110 $^\circ\text{C}$ for 30 min. A Sonotek Exactacoat system was used for spray-coating the perovskite precursor solution following our previously reported procedure.⁴¹ Briefly, the spray head moved linearly over the substrate at a speed of 80 mm s^{-1} at a separation distance of ~ 10 cm. The precursor solution was delivered at 1 mL min^{-1} and sheared into a droplet mist by an ultrasonically-vibrating tip driven at 2 W. N_2 shaping gas set to 3 psi directed the droplet mist to the substrate surface, with the hotplate-mounted substrate heated to 40 $^\circ\text{C}$. Around 30 s after spray deposition, the still-wet film was transferred to the antechamber of the glovebox and exposed to a vacuum of ~ 80 Pa for 1 min. After 1 min of vacuum exposure, the yellow precursor film had become uniformly dark brown. The film was annealed on a hotplate at 110 $^\circ\text{C}$ for 30 min.

A C_{60}/BCP (20 nm and 6 nm, respectively) electron-transporting layer was thermally evaporated (Angstrom Engineering) at a chamber base pressure of at least 2.4×10^{-6} mbar from alumina crucible sources (RADAK, Luxel Corp.) at a constant rate of 0.3 $\text{\AA} \text{s}^{-1}$. Silver pellets (Lesker) were deposited from a resistive boat source at a ramped rate of 0.1–1.0 $\text{\AA} \text{s}^{-1}$ through a shadow mask to form the Ag back-electrode (80 nm).

Materials characterization. Ultraviolet-visible (UV-vis) transmission measurements were recorded over the spectral range 300–900 nm using a Fluoromax-4 fluorometer (Horiba). Absorbance values were calculated according to $A = -\log_{10}(I/I_0)$ where I and I_0 are the intensity of light transmitted through and incident upon the sample respectively. The absorption coefficient, α , was then calculated according to $\alpha = A \ln(10)/L$, where L is the sample thickness.

Photoluminescence measurements were carried out by photo-exciting samples from the perovskite side using a 405 nm CW laser (Edmund Optics #35-069). The intensity was modulated with a neutral density filter wheel. The laser was focused onto the sample spot size $\sim 0.25 \text{ mm}^2$, and emission was collected through the same lens. A beamsplitter directed the emission through a 455 nm longpass filter (Thorlabs FGL455) after which it was fibre coupled to a spectrometer (OceanOptics Flame).

Film thicknesses were measured using a Bruker DekTak XT system. A razor blade was used to scratch the samples before scanning the profilometer tip (12.5 μm diameter) across the “scratch” at a stylus force of 3 mg. Vision64 (Bruker) was used to level the 1D line scan and extract the step height between the thin film surface and the bottom of the scratched valley.

Atomic force microscopy measurements were carried out in air in PeakForce Tapping mode with a Bruker Dimension Icon AFM system using a Bruker scanasyst-air silicon nitride probe (spring constant 0.4 N m^{-1}). Imaging was carried out with peak forces in the range 500–800 pN and tapping amplitudes in the range of 70–100 nm, at an actuation rate of 2 kHz. Images were recorded at scan rates of 0.3–0.6 Hz at a resolution of 512×512 pixels. Images were processed using the plane levelling function in Gwyddion.⁴² Contact angle measurements of water droplets were performed using a contact angle goniometer (Ossila). The Ossila contact angle software was used to extract the contact angle between water droplets and the sample surface at the point of initial contact. X-ray diffraction (XRD) measurements were performed using a Panalytical X'pert³ diffractometer equipped with a Cu line focus X-ray tube operating at a voltage of 45 kV and a current of 40 mA ($\lambda = 0.154 \text{ nm}$). Data was collected *via* a 1D-detector in Bragg-Brentano geometry. Energy-dispersive X-ray (EDX) spectroscopy measurements, were obtained using a field emission scanning electron microscope (FEI Nova NanoSEM 450, ThermoFisher) equipped with an EDX silicon drift detector (Oxford Instruments X-Max^N 80, Oxford Instruments). For all measurements a working distance between 4.9 and 5.2 mm was used. An accelerating voltage of 5 keV was used for all samples. Spectral maps were collected over an area of around 115 μm^2 at a resolution of 2048×1408 pixels per scan. For each pixel, the electron beam dwell time was 300 ns, with the processing time set to 6 s. Each scan comprised two recorded frames. All data were processed using Aztec (Oxford Instruments) software, with the binning factor set to 2.

Device characterisation. Current–voltage (J/V) measurements were recorded under ambient conditions using a Newport 92251A-1000 solar simulator. No preconditioning of devices was used. Prior to testing, the Air Mass 1.5 (AM1.5) spectrum was adjusted to 100 mW cm^{-2} at the substrate holder location using a national renewable energy laboratory (NREL) certified silicon reference cell. The active measurement area was defined using metal aperture masks with areas of 0.025 cm^2 , 0.20 cm^2 , or 1.0 cm^2 . A Keithley 237 source-measure unit swept devices between -0.1 and 1.2 V and back to -0.1 V at 0.4 V s^{-1} with a delay time of 50 ms. Stabilised power output measurements were performed by holding the device at a bias defined by the average voltage at maximum power (V_{mpp}) determined from the forward and reverse sweeps. External quantum efficiency measurements were recorded over a 325–900 nm range using a Newport QuantX-300 Quantum Efficiency Measurement System. The system was equipped with a 100 W Xenon arc lamp focused through an Oriel Monochromator (CS130B) and chopped at 25 Hz. An Oriel LSH-7320 ABA LED solar simulator was used to perform light intensity-dependent J/V measurements by adjusted the output power between 0.1 and 1.1 suns as determined by an NREL-certified silicon reference cell.

Results and discussion

Optimisation of spray-coated perovskite coverage

Alumina nanoparticles have previously been reported to enable spin-coating of metal halide perovskite thin films from the



commonly used mixed dimethylformamide:dimethylsulfoxide solvent system onto Me-4PACz layers^{34,36} (Fig. S1, ESI†). However, this interfacial modification does not enable coalescence of spray-coated perovskite precursor droplets from the same solvent system and results in incomplete thin films unsuitable for photovoltaic devices (Fig. S2, ESI†). The separation of the spray-coated perovskite solution droplets suggests that the

surface energy of the Me-4PACz surface remains too high to facilitate the formation of a continuous thin film.

Therefore we investigated the impact of aluminium oxide nanoparticle ($\text{Al}_2\text{O}_3\text{-np}$) solution concentration on the surface energy of Me-4PACz thin films. We report the $\text{Al}_2\text{O}_3\text{-np}$ concentrations investigated as a ratio of nanoparticles to isopropanol ($\text{Al}_2\text{O}_3\text{-nps}$:isopropanol) and we investigated the

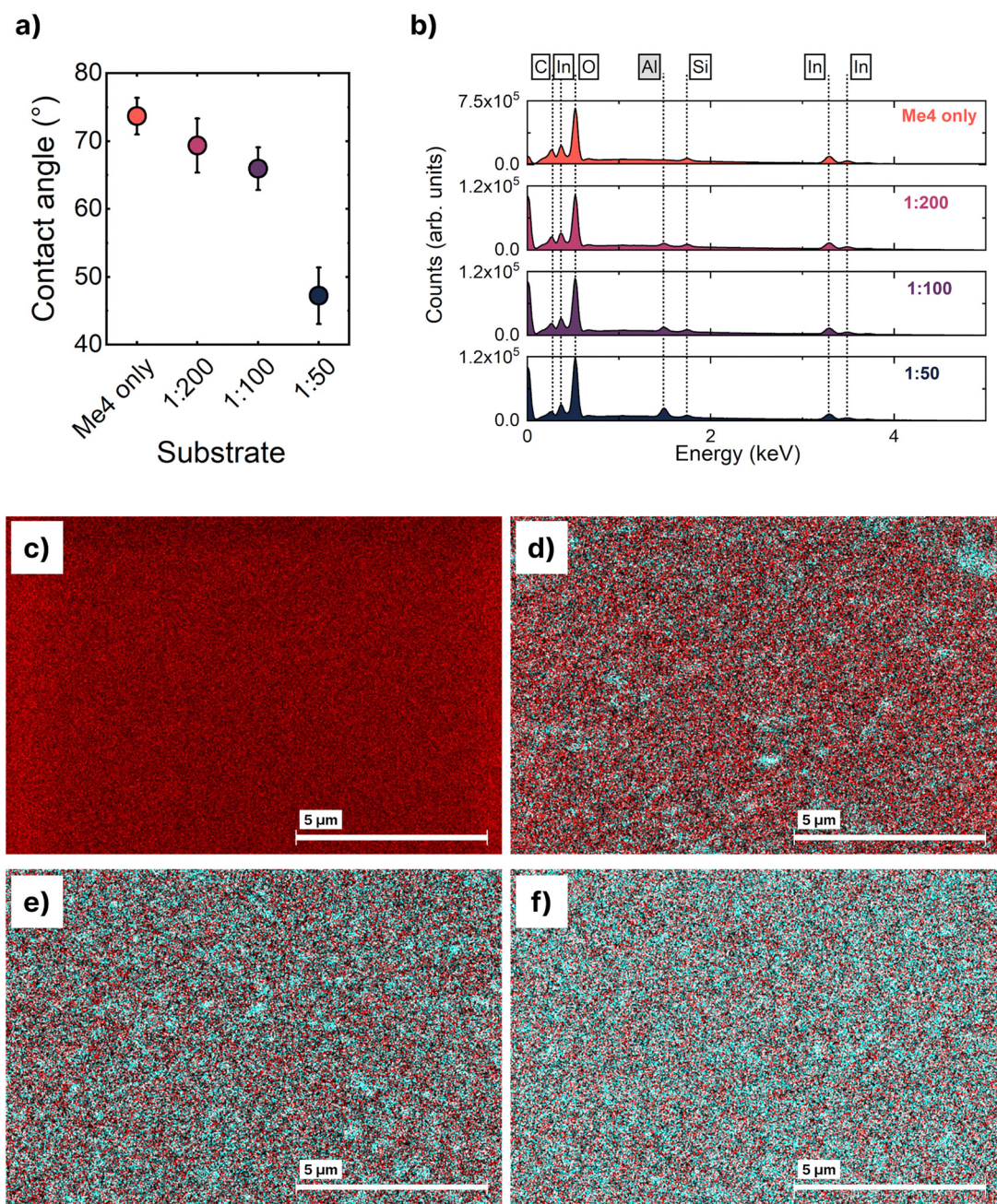


Fig. 1 Characterisation of Al_2O_3 treatment of Me-4PACz thin films on indium tin oxide (ITO) coated glass. (a) Summary of the measured contact angles of a droplet of water on neat Me-4PACz thin films and Al_2O_3 treated with a range of nanoparticle concentrations. (b) Energy dispersive X-ray (EDX) spectroscopy measurements and corresponding elemental images measured using scanning electron microscopy of thin films of Me-4PACz spin-coated onto indium tin oxide (ITO) glass and treated with Al_2O_3 nanoparticle solutions diluted with isopropanol in the following ratios (nanoparticles: isopropanol) (c) no nanoparticles, (d) 1: 200, (e) 1: 100, and (f) 1: 50. Red indicates carbon and cyan indicates aluminium.



following dilutions: 1:200, 1:100, and 1:50. Contact angle measurements were used to approximate the change in surface free energy of the SAM surface upon interfacial modification with Al_2O_3 -nps. The contact angle between water droplets and the SAM/ Al_2O_3 -np surfaces was measured, and the resulting data reported in Fig. 1a. Without modification, the hydrophobic Me-4PACz surface had a water contact angle of $73.7 \pm 2.7^\circ$ (mean $\pm$ standard deviation, $n = 9$). When treated with 1:200, 1:100, and 1:50 dilutions of Al_2O_3 -nps, this contact angle decreased to $69.4 \pm 4.0^\circ$ ($n = 6$), $65.9 \pm 3.1^\circ$ ($n = 9$), and $47.2 \pm 4.2^\circ$ ($n = 8$), respectively (see Table S3 for full data, ESI†). This decrease in contact angle is suggestive of a lower surface energy which should facilitate deposition of high-quality perovskite thin films. The decrease in contact angle with increasing Al_2O_3 -np solution concentration is concomitant with an increase in aluminium concentration. Fig. 1b shows the energy dispersive X-ray (EDX) spectra of bare Me-4PACz and Me-4PACz treated with each dilution of Al_2O_3 -nps. The wt% of Al increased from 0% (bare Me-4PACz) to 1%, 1.5%, and 3.4% for 1:200, 1:100, and 1:50 dilutions, respectively. EDX maps of C and Al distribution (Fig. 1c–f) show the Al_2O_3 -nps are randomly distributed across the sample surface.

Subsequently we investigated the effect of increased Al_2O_3 -np concentration on the morphology, crystal structure, and optoelectronic properties of perovskite thin films. A “triple cation” $\text{Cs}_{0.05}\text{FA}_{0.81}\text{MA}_{0.14}\text{PbI}_{2.55}\text{Br}_{0.45}$ perovskite thin film was deposited onto Me-4PACz/ Al_2O_3 -np coated ITO glass substrates by either spin-coating or ultrasonic spray-coating. As with

spin-coating it was not possible to fabricate uniform perovskite thin films using spray-coated on Me-4PACz without Al_2O_3 -nps due to severe dewetting of the perovskite precursor solution. Hereafter, perovskite thin films which are fabricated through spray-coating on either a 1:100 or 1:50 dilution of Al_2O_3 -nps are referred to as “spray 1:100” or “spray 1:50”, respectively. Spray-coating of perovskite thin films atop the 1:200 dilution of nanoparticles resulted in incomplete film formation due to an insufficient modification of the surface energy of the Me-4PACz thin film (Fig. S2, ESI†) and were not further investigated. We compare the properties of our spray-coated perovskite thin films with optimised spin-coated perovskite thin films with a 1:100 dilution of Al_2O_3 -nps on Me-4PACz, these are referred to as “spin 1:100”.

Fig. 2a–c and Fig. S3 (ESI†) show atomic force microscopy (AFM) images of $\text{Cs}_{0.05}\text{FA}_{0.81}\text{MA}_{0.14}\text{PbI}_{2.55}\text{Br}_{0.45}$ thin films deposited on to Al_2O_3 treated Me-4PACz thin films. The corresponding trend in root-mean-squared (RMS) roughness (R_q) is shown in Fig. 2d with full data reported in Table S4 (ESI†). Whilst the perovskite grain size remains relatively consistent we observe a considerable increase in the surface roughness between the spin- and spray-coated thin films and also with increasing concentration of nanoparticles. Whilst the spin-coated thin films are annealed immediately following film deposition our spray-coated films undergo an intermediate vacuum annealing step prior to thermal annealing. Previous reports have shown that the inclusion of DMSO as a co-solvent in perovskite processing onto hydrophobic surfaces can result



Fig. 2 Atomic force microscopy (AFM) images of $\text{Cs}_{0.05}\text{FA}_{0.81}\text{MA}_{0.14}\text{PbI}_{2.55}\text{Br}_{0.45}$ thin films on Al_2O_3 treated Me-4PACz on indium tin oxide (ITO) glass. The perovskite thin films were fabricated through spin- or spray-coating and the ratios represent dilution of nanoparticles : isopropanol. (a) 1:100 spin-coating (b) 1:100 spray-coating and (c) 1:50 spray-coating. (d) root-mean-squared roughness determined from AFM images, error bars represent the mean standard deviation ($n = 5$) (e) and corresponding X-ray diffraction patterns.



in dewetting of the thin films during annealing and an increase in the surface roughness of the resulting perovskite thin films.²² The combination of solvent system, the higher vapour pressure of DMF (relative to DMSO), and the vacuum annealing processing step is likely to result in a thin film which has a higher proportion of DMSO present when undergoing the thermal annealing process. This could therefore lead to an increase in film roughness due to the perovskite thin film dewetting from the SAMs. Furthermore, in the spray-coated 1:50 films there are regions that are substantially higher relative to the rest of the film (Fig. S4, ESI†). The EDX data suggests a non-uniform distribution of the nanoparticles and we propose these higher regions are the result of the perovskite film forming atop regions of these clustered nanoparticles. The film thicknesses were measured using a profilometer (Fig. S5 and Table S5, ESI†) and the spray-coated films were observed to be significantly thinner than those fabricated through spin-coating ~ 540 nm (spray 1:100) and ~ 590 nm (spray 1:50) vs. 670 nm (spin 1:100). Fig. 2e shows X-ray diffraction patterns corresponding to the perovskite thin films on Al_2O_3 treated Me-4PACz. Additional plots can be found in Fig. S6 and S7 (ESI†). There is a consistent shift in peak positions between the spin and spray-coated samples with the spray-coated samples both exhibiting a shift of $2\theta = \sim 0.05^\circ$ towards lower angles. This is typically indicative of an increase in the unit cell size suggesting that there is a structural difference between the spin- and spray-coated films. We note a small peak at $\sim 12.7^\circ$ within the spray

1:50 sample, indicative of PbI_2 developing during film formation.⁴³

To investigate the influence of this structural change on the optoelectronic properties of these thin films, we conducted photoluminescence (PL) spectroscopy as a function of illumination intensity. As can be seen in Fig. 3a, the centre of the PL peak redshifts substantially between the spin-coated and spray-coated samples for the same Al_2O_3 concentration. This indicates that the emission is originating from a different structural phase in each sample, suggesting that the different film fabrication routes and subsequent film-formation mechanisms have a strong influence on the optical properties. The magnitude of the red-shift is reduced for the spray-coated 1:50 sample, suggesting that the resulting thin film is more similar to the spin-coated counterpart. Furthermore, the spray-coated samples exhibit an asymmetric PL peak shape. This asymmetry in combination with the red-shifting of the peaks is suggestive of emission from iodide-rich structures⁴⁴ which is consistent with the peak shifting also observed in the X-ray diffraction measurements. This is further evidence that the metal halide perovskite thin films are not uniform, despite the advances made through the alumina interlayer.

Fig. 3a also shows a substantial increase in the PL intensity when moving to spray-coated thin films, and with an increased Al_2O_3 concentration. This is strong evidence for a reduction in the severity of interfacial recombination in these systems. To understand whether the recombination behaviour was



Fig. 3 Optoelectronic properties of $\text{Cs}_{0.05}\text{FA}_{0.81}\text{MA}_{0.14}\text{PbI}_{2.55}\text{Br}_{0.45}$ thin films on Al_2O_3 treated Me-4PACz on indium tin oxide (ITO) glass. Ratios represent dilution of nanoparticles: isopropanol. (a) Intensity dependent steady-state photoluminescence spectra, the integrated area under the peak was then used to calculate (b) the ideality factor, and (c) the absorption coefficient as determined through ultraviolet-visible spectroscopy measurements. The shaded area represents the uncertainty in these measurements.



changing in these systems, we performed emission measurements across a wide excitation power range which we present in Fig. 3a and summarise in Fig. 3b. The absolute intensity of the peaks is increased by a factor of ~ 2.5 and ~ 3.5 for spray 1:100 and spray 1:50 over spin 1:100 respectively. However, this increase is not dependent on illumination intensity, giving rise to near identical ideality factors for these systems. This indicates that the underlying recombination mechanisms are unchanged in these systems.

The absorption coefficient of the perovskite thin films was characterised using UV-vis spectroscopy which we show in Fig. 3c. Broadly, the absorption coefficient is similar for each sample with only minor differences. These differences include an apparent increase in below bandgap absorption in the spray-coated samples. However, this is likely due to the spray-coated samples exhibiting an increased roughness (Fig. 2d) resulting in increased scattering.

The absorption onset remains unchanged for each sample, which demonstrates that the bulk perovskite is comparable for each condition. This suggests that the red-shifting of the PL peaks is not a 'bulk' effect, but instead originates from a volume of the film that is too small to be detected in absorption measurements. This is likely due to differences in the growth mechanism between the different thin film fabrication routes

and due to the differences in the Al_2O_3 surface modification. This is supported by the clear surface roughness differences between the different systems, determined by AFM measurements, which suggest differences in film formation.

Photovoltaic device performance

To corroborate our structural, morphological and optoelectronic investigations, we fabricated a series of *p-i-n* photovoltaic devices. All devices had the configuration: ITO-coated glass/Me-4PACz/ Al_2O_3 -nps/ $\text{Cs}_{0.05}\text{FA}_{0.81}\text{MA}_{0.14}\text{PbI}_{2.55}\text{Br}_{0.45}/\text{C}_{60}$ /bathocuproine (BCP)/Ag, where the perovskite deposition technique and/or Al_2O_3 -nps concentration was altered. We show a statistical summary of the power conversion efficiency (PCE), short-circuit current density (J_{SC}), open-circuit voltage (V_{OC}), and fill factor (FF) in Fig. 4a–d respectively for the different conditions investigated in this study. Additionally, we summarise the device performance statistics in Table 1.

The average PCE of spin-coated cells (18.7%, $n = 48$ devices) is higher than the PCE of spray-coated cells (17.7%, $n = 48$) for a 1:100 dilution of Al_2O_3 -nps. Whilst the average J_{SC} and V_{OC} of spray 1:100 cells are reduced in comparison to spin 1:100 (22.5 mA cm^{-2} vs. 21.8 mA cm^{-2} , and 1.10 V vs. 1.08 V), device FF remains similar (75.6% vs. 75.4%). The reduction in J_{SC} can be ascribed to the reduced thickness of the spray-coated

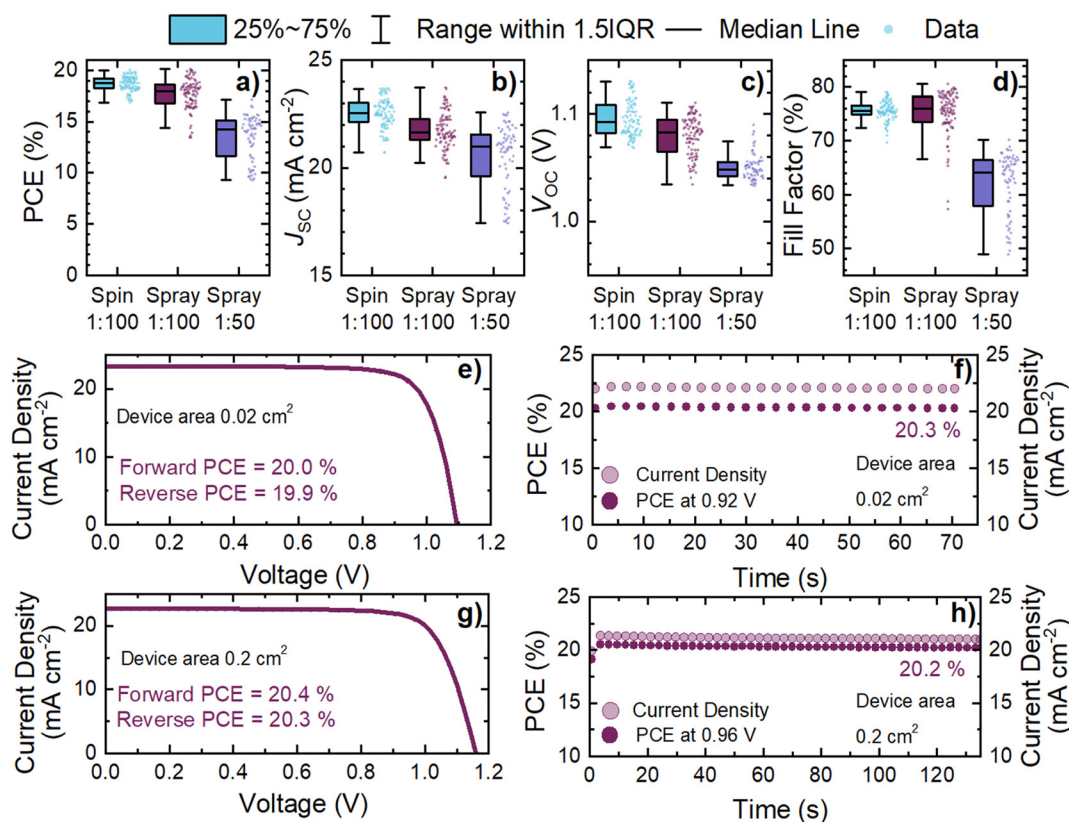


Fig. 4 Photovoltaic device performances for *p-i-n* $\text{Cs}_{0.05}\text{FA}_{0.81}\text{MA}_{0.14}\text{PbI}_{2.55}\text{Br}_{0.45}$ devices incorporating optimised Al_2O_3 nanoparticle interlayer. Box plots for device parameters showing data points relating to (a) power conversion efficiency (PCE), (b) short circuit current density (J_{SC}), (c) open-circuit voltage (V_{OC}), and (d) fill factor. Data relating to champion device performances is shown in (e) J-V characteristics for spray (1:100) condition for 0.02 cm^2 device active area and (f) corresponding stabilised power conversion efficiency at maximum power point, (g) J-V characteristics for spray (1:100) condition for 0.2 cm^2 device active area and (h) corresponding stabilised power conversion efficiency at maximum power point.



perovskite. The reduced V_{OC} is surprising, however, since the photoluminescence data in Fig. 3 suggested a reduced degree of non-radiative recombination at the perovskite/SAM interface for spray-coated devices. However, the increased roughness of the spray-coated perovskite has likely impacted the quality of the perovskite/ C_{60} interface, which has resulted in the reduction of the open-circuit voltage. Encouragingly, the maximum device performance of spray 1:100 devices is comparable to the maximum PCE of spin 1:100 cells (20.1% vs. 20.0%, respectively). The current–voltage (JV) curve of the best-performing spray 1:100 device (active area of 0.02 cm²) is shown in Fig. 4e. Fig. S9 and S10 (ESI†) shows the external quantum efficiency (EQE) response for the champion cells of each condition. The integrated short-circuit current density (J_{SC-EQE}) calculated using the EQE data is in good agreement with the J_{SC} extracted from the current–voltage characteristics. For the spray 1:100 device the J_{SC-EQE} = 22.6 mA cm^{−2} compared to 23.3 mA cm^{−2} for the measured J_{SC} . The mismatch is just 3% providing confidence to our solar simulator measurements. We calculated the photovoltaic bandgap,^{45–47} shown in Fig. S11 (ESI†), using the derivative of the EQE with respect to energy. This yields a bandgap of 1.62 eV, consistent with the optical data presented in Fig. 3.

The stabilised power output of the champion spray 1:100 device at maximum power point was 20.3% (Fig. 4f) which constitutes a record performance for both spray-coated triple cation and spray-coated p – i – n photovoltaic devices.^{24,41,48–50} This demonstrates the effectiveness of the Al_2O_3 -np interlayer approach as compared to other techniques. Indeed, our champion device outperforms the spin-coated counterpart (with no additional passivation) Me-4PACz devices in the literature.^{19,20,26,51–53} This highlights the potential of this approach for demonstrating spray-coated high-performance perovskite photovoltaics enabling high-throughput fabrication. Photovoltaic devices based on the spray 1:50 condition were fabricated, but the improvements implied by the optoelectronic measurements were not present in the device performance. The deterioration in device performance can be attributed to several factors. The increased roughness of the perovskite will lead to increased interfacial recombination limiting the V_{OC} and reducing the shunt resistance (Fig. S12, ESI†). The increased coverage of the insulating Al_2O_3 interlayer partially inhibits charge transfer to the SAM, which simultaneously constrains the J_{SC} and FF of these devices.

We investigated the ideality factor of photovoltaic devices based on each condition through light intensity dependent JV measurements. Fig. S13(a, c and e) (ESI†) show a double

logarithmic plot of J_{SC} versus light intensity and (b, d and e) show the light-dependent V_{OC} response for spray 1:50, spray 1:100, and spin 1:100 cells, respectively. The ideality factor (n) of the devices was found to be 1.46 (spray 1:50), 1.10 (spray 1:100), and 1.13 (spin 1:100) $k_B T/q$ respectively. Contrastingly, the ideality factors determined in Fig. 3(b) from perovskite thin-films on Al_2O_3 -np-coated Me-4PACz (without an electron-transporting layer) were near identical for each device type. This change in ideality factor for spray 1:50 devices supports our suggestion that the increased roughness at the C_{60} /perovskite interface in these samples is influencing the recombination processes that are occurring.⁵⁴ Additionally, this inferior morphology and increased surface recombination would simultaneously constrain the FF, as observed in our devices.

To explore the applicability of our Al_2O_3 -np interlayer for large scale deposition of spray-coated perovskite layers onto Me-4PACz, we fabricated spray-coated devices with larger active areas (0.2 cm² and 1 cm²) incorporating the best performing interlayer condition (1:100). The champion 0.2 cm² device demonstrated a stabilised PCE of 20.2% (Fig. 4g and h), equivalent to smaller area devices. Fig. S14 (ESI†) shows histogram plots of the device PCE distribution for each active area of spray 1:100 cells (0.02 cm² and 0.2 cm²) which shows comparable performance for each condition. All device parameters for these 0.02 cm² and 0.2 cm² devices are summarised in Fig. S15 (ESI†). We show champion JV characteristics and stabilised PCE for 1 cm² devices of both the spin 1:100 and spray 1:100 conditions in Fig. S16 (ESI†). For these conditions we achieve a maximum stabilised PCE of ~17% for both conditions.

Interestingly, upon scale-up to 0.2 cm² we observe an increased median J_{SC} and V_{OC} . This indicates a reduced severity of non-radiative recombination in these devices as well as efficient charge-extraction at short-circuit conditions. This suggests an improvement in the formation of the thin film layers resulting in a reduction in structural defects and hence recombination sites. However, the median PCE is reduced slightly, due to an inferior FF for the 0.2 cm² devices, highlighting a charge transport limitation in these systems. From inspecting the JV curve in Fig. 4g, this reduced FF originates from an increased series resistance within the scaled-up photovoltaics. We attribute this to the sheet resistance of our ITO substrates (~20 $\Omega \square^{-1}$) becoming significant in these larger devices. This phenomenon is exacerbated for 1 cm² devices where we observe the FF being constrained to ~66–68%. This arises from the gradient in the JV curves near V_{OC} becoming increasingly shallow, characteristic of an increased series resistance.

Table 1 Photovoltaic device performance parameters, champion device performance is reported in bold and brackets. Device area and number of devices per condition is presented in the table

	J_{SC} (mA cm ^{−2})	V_{OC} (V)	FF (%)	PCE (%)	n_{devices}
Spin 1:100 0.025 cm ²	22.5 ± 0.7 (23.7)	1.10 ± 0.02 (1.13)	75.6 ± 1.7 (79.1)	18.7 ± 0.7 (20.0)	48
Spray 1:100 0.025 cm ²	21.8 ± 0.9 (23.7)	1.08 ± 0.02 (1.11)	75.4 ± 4.1 (80.6)	17.7 ± 1.5 (20.1)	48
Spray 1:50 0.025 cm ²	20.5 ± 1.5 (22.6)	1.05 ± 0.01 (1.09)	62.1 ± 5.7 (70.2)	13.5 ± 2.2 (17.2)	32
Spray 1:100 0.2 cm ²	22.0 ± 1.5 (22.8)	1.09 ± 0.03 (1.16)	69.0 ± 5.6 (77.7)	16.7 ± 2.3 (20.4)	31



Therefore, this limitation is not intrinsic to the Al₂O₃-np layer we report here, highlighting that with sufficiently low sheet resistance, this approach is applicable for even larger area scale-up.

Conclusions

In summary, we have demonstrated an optimised Al₂O₃ interlayer which enables scalable, spray-coating of high performance 1.62 eV *p-i-n* metal halide perovskite photovoltaics onto the Me-4PACz SAM charge transporting layer. We demonstrated that the incorporation of this interlayer between the Me-4PACz and perovskite thin films results in spray-coated perovskite layers with improved structural and optoelectronic properties relative to their spin-coated counterparts. As a result we report champion photovoltaic devices with efficiencies of up to 20.4% which represents the state-of-the-art for spray-coated *p-i-n* PSCs. The improvements in device performance resulting from the incorporation of the optimised Al₂O₃-np interlayer remained effective as the device active area was scaled from 0.02 cm² to 0.2 cm², suggesting that this approach warrants further investigation for large area photovoltaic devices fabricated through high-throughput solution processing routes.

Author contributions

Elena J. Cassella: conceptualization, methodology, investigation, formal analysis, writing – original draft, review and editing. Robert D.J. Oliver: conceptualization, methodology, investigation, formal analysis, supervision, writing – original draft, review and editing. Timothy Thornber: methodology, investigation, formal analysis, writing – original draft, review and editing. Sophie Tucker: investigation, formal analysis. Rehmat Goodwin: investigation, formal analysis. David G. Lidzey: conceptualization, supervision, funding acquisition, writing – review and editing. Alexandra J. Ramadan: conceptualization, methodology, supervision, funding acquisition, writing – original draft, writing – review and editing.

Data availability

Data for this article, including raw data and processed data are available at ORDA at <https://doi.org/10.15131/shef.data.26180623>.

Conflicts of interest

D. G. L. is co-chairman of the materials science company Ossila Ltd. which retails materials and equipment for photovoltaic research, including perovskite photovoltaics.

Acknowledgements

This work was funded by the Engineering and Physical Sciences Research Council (EPSRC) grants EP/V027131/1, EP/X039285/1

and EP/T012455/1. T. T. acknowledges the Faculty of Science, University of Sheffield for a studentship. S. T. and R. G. acknowledge the EPSRC and the Grantham Centre for Sustainable Futures for funding. R. G. also acknowledges the University of Sheffield's Energy Institute for funding.

References

- 1 S. D. Stranks, G. E. Eperon, G. Grancini, C. Menelaou, M. J. P. Alcocer, T. Leijtens, L. M. Herz, A. Petrozza and H. J. Snaith, *Science*, 2013, **342**, 341–344.
- 2 S. De Wolf, J. Holovsky, S.-J. Moon, P. Löper, B. Niesen, M. Ledinsky, F.-J. Haug, J.-H. Yum and C. Ballif, *J. Phys. Chem. Lett.*, 2014, **5**, 1035–1039.
- 3 A. Miyata, A. Mitoglu, P. Plochocka, O. Portugall, J. T.-W. Wang, S. D. Stranks, H. J. Snaith and R. J. Nicholas, *Nat. Phys.*, 2015, **11**, 582–587.
- 4 C. Wehrenfennig, G. E. Eperon, M. B. Johnston, H. J. Snaith, L. M. Herz, C. Wehrenfennig, G. E. Eperon, M. B. Johnston, H. J. Snaith and L. M. Herz, *Adv. Mater.*, 2014, **26**, 1584–1589.
- 5 Z. Li, T. R. Klein, D. H. Kim, M. Yang, J. J. Berry, M. F. A. M. Van Hest and K. Zhu, *Nat. Rev. Mater.*, 2018, **3**, 1–20.
- 6 L. Zhang, L. Mei, K. Wang, Y. Lv, S. Zhang, Y. Lian, X. Liu, Z. Ma, G. Xiao, Q. Liu, S. Zhai, S. Zhang, G. Liu, L. Yuan, B. Guo, Z. Chen, K. Wei, A. Liu, S. Yue, G. Niu, X. Pan, J. Sun, Y. Hua, W.-Q. Wu, D. Di, B. Zhao, J. Tian, Z. Wang, Y. Yang, L. Chu, M. Yuan, H. Zeng, H.-L. Yip, K. Yan, W. Xu, L. Zhu, W. Zhang, G. Xing, F. Gao and L. Ding, *Nano-Micro Lett.*, 2023, **15**, 177.
- 7 Z.-Y. Chen, N.-Y. Huang and Q. Xu, *Coord. Chem. Rev.*, 2023, **481**, 215031.
- 8 A. Fakharuddin, M. K. Gangishetty, M. Abdi-Jalebi, S.-H. Chin, A. R. bin Mohd Yusoff, D. N. Congreve, W. Tress, F. Deschler, M. Vasilopoulou and H. J. Bolink, *Nat. Electron.*, 2022, **5**, 203–216.
- 9 Y. Wu, J. Feng, Z. Yang, Y. Liu and S. (Frank) Liu, *Adv. Sci.*, 2023, **10**, 2205536.
- 10 N. Rolston, W. J. Scheideler, A. C. Flick, J. P. Chen, H. Elmaraghi, A. Sleugh, O. Zhao, M. Woodhouse and R. H. Dauskardt, *Joule*, 2020, **4**, 2675–2692.
- 11 L.-H. Chou, J. M. W. Chan and C.-L. Liu, *Sol. RRL*, 2022, **6**, 2101035.
- 12 W. Tress, *Adv. Energy Mater.*, 2017, **7**, 1602358.
- 13 C. M. Wolff, P. Caprioglio, M. Stolterfoht and D. Neher, *Adv. Mater.*, 2019, **31**, 1902762.
- 14 M. Stolterfoht, P. Caprioglio, C. M. Wolff, J. A. Márquez, J. Nordmann, S. Zhang, D. Rothhardt, U. Hörmann, Y. Amir, A. Redinger, L. Kegelmann, F. Zu, S. Albrecht, N. Koch, T. Kirchartz, M. Saliba, T. Unold and D. Neher, *Energy Environ. Sci.*, 2019, **12**, 2778–2788.
- 15 A. Al-Ashouri, A. Magomedov, M. Roß, M. Jošt, M. Talaikis, G. Chistiakova, T. Bertram, J. A. Márquez, E. Köhnen, E. Kasparavičius, S. Levenco, L. Gil-Escrig, C. J. Hages, R. Schlattmann, B. Rech, T. Malinauskas, T. Unold,



- C. A. Kaufmann, L. Korte, G. Niaura, V. Getautis and S. Albrecht, *Energy Environ. Sci.*, 2019, **12**, 3356–3369.
- 16 A. Magomedov, A. Al-Ashouri, E. Kasparavičius, S. Strazdaite, G. Niaura, M. Jošt, T. Malinauskas, S. Albrecht and V. Getautis, *Adv. Energy Mater.*, 2018, **8**, 1801892.
 - 17 A. Al-Ashouri, E. Köhnen, B. Li, A. Magomedov, H. Hempel, P. Caprioglio, J. A. Márquez, A. B. Morales Vilches, E. Kasparavičius, J. A. Smith, N. Phung, D. Menzel, M. Grischek, L. Kegelmann, D. Skroblin, C. Gollwitzer, T. Malinauskas, M. Jošt, G. Matič, B. Rech, R. Schlatmann, M. Topič, L. Korte, A. Abate, B. Stannowski, D. Neher, M. Stölterfoht, T. Unold, V. Getautis and S. Albrecht, *Science*, 2020, **370**, 1300–1309.
 - 18 M. Roß, S. Severin, M. B. Stutz, P. Wagner, H. Köbler, M. Favin-Lévêque, A. Al-Ashouri, P. Korb, P. Tockhorn, A. Abate, B. Stannowski, B. Rech and S. Albrecht, *Adv. Energy Mater.*, 2021, **11**, 2101460.
 - 19 A. Kulkarni, R. Sarkar, S. Akel, M. Häser, B. Klingebiel, M. Wuttig, S. Wiegand, S. Chakraborty, M. Saliba and T. Kirchartz, *Adv. Funct. Mater.*, 2023, **33**, 2305812.
 - 20 A. Al-Ashouri, M. Marčinkas, E. Kasparavičius, T. Malinauskas, A. Palmstrom, V. Getautis, S. Albrecht, M. D. McGehee and A. Magomedov, *ACS Energy Lett.*, 2023, **8**, 898–900.
 - 21 F. H. Isikgor, A. S. Subbiah, M. K. Eswaran, C. T. Howells, A. Babayigit, M. De Bastiani, E. Yengel, J. Liu, F. Furlan, G. T. Harrison, S. Zhumagali, J. I. Khan, F. Laquai, T. D. Anthopoulos, I. McCulloch, U. Schwingenschlögl and S. De Wolf, *Nano Energy*, 2021, **81**, 105633.
 - 22 K. O. Brinkmann, J. He, F. Schubert, J. Malerczyk, C. Kreusel, F. van gen Hassend, S. Weber, J. Song, J. Qu and T. Riedl, *ACS Appl. Mater. Interfaces*, 2019, **11**, 40172–40179.
 - 23 J. Li, J. Dagar, O. Shargaieva, M. A. Flatken, H. Köbler, M. Fenske, C. Schultz, B. Stegemann, J. Just, D. M. Többsen, A. Abate, R. Munir and E. Unger, *Adv. Energy Mater.*, 2021, **11**, 2003460.
 - 24 E. J. Cassella, E. L. K. Spooner, T. Thornber, M. E. O’Kane, T. E. Catley, J. E. Bishop, J. A. Smith, O. S. Game and D. G. Lidzey, *Adv. Sci.*, 2022, **9**, 2104848.
 - 25 Q. Cao, T. Wang, X. Pu, X. He, M. Xiao, H. Chen, L. Zhuang, Q. Wei, H.-L. Loi, P. Guo, B. Kang, G. Feng, J. Zhuang, G. Feng, X. Li and F. Yan, *Adv. Mater.*, 2024, **36**(16), 2311970.
 - 26 K. Hossain, A. Kulkarni, U. Bothra, B. Klingebiel, T. Kirchartz, M. Saliba and D. Kabra, *ACS Energy Lett.*, 2023, **8**, 3860–3867.
 - 27 X. Zheng, Z. Li, Y. Zhang, M. Chen, T. Liu, C. Xiao, D. Gao, J. B. Patel, D. Kuciauskas, A. Magomedov, R. A. Scheidt, X. Wang, S. P. Harvey, Z. Dai, C. Zhang, D. Morales, H. Pruetz, B. M. Wieliczka, A. R. Kirmani, N. P. Padture, K. R. Graham, Y. Yan, M. K. Nazeeruddin, M. D. McGehee, Z. Zhu and J. M. Luther, *Nat. Energy*, 2023, **8**, 462–472.
 - 28 A. Farag, T. Feeney, I. M. Hossain, F. Schackmar, P. Fassl, K. Küster, R. Bäuerle, M. A. Ruiz-Preciado, M. Hentschel, D. B. Ritzer, A. Diercks, Y. Li, B. A. Nejjand, F. Laufer, R. Singh, U. Starke and U. W. Paetzold, *Adv. Energy Mater.*, 2023, **13**, 2203982.
 - 29 M. Pitaro, J. E. S. Alonso, L. Di Mario, D. G. Romero, K. Tran, J. Kardula, T. Zaharia, M. B. Johansson, E. M. J. Johansson, R. C. Chiechi and M. A. Loi, *Adv. Funct. Mater.*, 2023, 2306571.
 - 30 Y. Xu, F. Fei, X. Dong, L. Li, Y. Li, N. Yuan and J. Ding, *Sol. RRL*, 2023, **7**, 2300283.
 - 31 N. Phung, M. Verheijen, A. Todorova, K. Datta, M. Verhage, A. Al-Ashouri, H. Köbler, X. Li, A. Abate, S. Albrecht and M. Creatore, *ACS Appl. Mater. Interfaces*, 2022, **14**, 2166–2176.
 - 32 D. Takhellambam, L. A. Castriotta, G. Zanotti, L. Mancini, V. Raglione, G. Mattioli, B. Paci, A. Generosi, M. Guaragno, V. Campanari, G. Ammirati, F. Martelli, E. Calabrò, A. Cricenti, M. Luce, N. Yaghoobi Nia, F. Di Giacomo and A. Di Carlo, *Sol. RRL*, 2023, **7**, 2300658.
 - 33 S.-C. Liu, H.-Y. Lin, S.-E. Hsu, D.-T. Wu, S. Sathasivam, M. Daboczi, H.-J. Hsieh, C.-S. Zeng, T.-G. Hsu, S. Eslava, T. J. Macdonald and C.-T. Lin, *J. Mater. Chem. A*, 2024, **12**, 2856–2866.
 - 34 M. Taddei, J. A. Smith, B. M. Gallant, S. Zhou, R. J. E. Westbrook, Y. Shi, J. Wang, J. N. Drysdale, D. P. McCarthy, S. Barlow, S. R. Marder, H. J. Snaith and D. S. Ginger, *ACS Energy Lett.*, 2022, **7**, 4265–4273.
 - 35 W. Peng, K. Mao, F. Cai, H. Meng, Z. Zhu, T. Li, S. Yuan, Z. Xu, X. Feng, J. Xu, M. D. McGehee and J. Xu, *Science*, 2023, **379**, 683–690.
 - 36 W. H. K. Perera, M. G. Masteghin, H. Shim, J. D. Davies, J. L. Ryan, S. J. Hinder, J. S. Yun, W. Zhang, K. D. G. I. Jayawardena and S. R. P. Silva, *Sol. RRL*, 2023, **2300388**, 1–11.
 - 37 Y. Diao, L. Shaw, Z. Bao and S. C. B. Mannsfeld, *Energy Environ. Sci.*, 2014, **7**, 2145–2159.
 - 38 J. E. Bishop, J. A. Smith and D. G. Lidzey, *ACS Appl. Mater. Interfaces*, 2020, **12**, 48237–48245.
 - 39 T. Carey, C. Jones, F. Le Moal, D. Deganello and F. Torrisi, *ACS Appl. Mater. Interfaces*, 2018, **10**, 19948–19956.
 - 40 T. Thornber, O. S. Game, E. J. Cassella, M. E. O’Kane, J. E. Bishop, T. J. Routledge, T. I. Alanazi, M. Togay, P. J. M. Isherwood, L. C. Infante-Ortega, D. B. Hammond, J. M. Walls and D. G. Lidzey, *ACS Appl. Mater. Interfaces*, 2022, **14**, 37587–37594.
 - 41 J. E. Bishop, C. D. Read, J. A. Smith, T. J. Routledge and D. G. Lidzey, *Sci. Rep.*, 2020, **10**, 6610.
 - 42 D. Nečas and P. Klapetek, *Open Phys.*, 2012, **10**, 181–188.
 - 43 T. J. Jacobsson, J.-P. Correa-Baena, E. Halvani Anaraki, B. Philippe, S. D. Stranks, M. E. F. Bouduban, W. Tress, K. Schenk, J. Teuscher, J.-E. Moser, H. Rensmo and A. Hagfeldt, *J. Am. Chem. Soc.*, 2016, **138**, 10331–10343.
 - 44 A. J. Knight, J. Borchert, R. D. J. Oliver, J. B. Patel, P. G. Radaelli, H. J. Snaith, M. B. Johnston and L. M. Herz, *ACS Energy Lett.*, 2021, **6**, 799–808.
 - 45 P. K. Nayak, S. Mahesh, H. J. Snaith and D. Cahen, *Nat. Rev. Mater.*, 2019, **4**, 269–285.
 - 46 U. Rau, B. Blank, T. C. M. Müller and T. Kirchartz, *Phys. Rev. Appl.*, 2017, **7**, 044016.
 - 47 A. J. Ramadan, R. D. J. Oliver, M. B. Johnston and H. J. Snaith, *Nat. Rev. Mater.*, 2023, **8**, 822–838.
 - 48 J. Su, H. Cai, J. Yang, X. Ye, R. Han, J. Ni, J. Li and J. Zhang, *ACS Appl. Mater. Interfaces*, 2019, **12**, 3531–3538.



- 49 J. E. Bishop, J. A. Smith, C. Greenland, V. Kumar, N. Vaenas, O. S. Game, T. J. Routledge, M. Wong-Stringer, C. Rodenburg and D. G. Lidzey, *ACS Appl. Mater. Interfaces*, 2018, **10**, 39428–39434.
- 50 M. Park, S. C. Hong, Y. W. Jang, J. Byeon, J. Jang, M. Han, U. Kim, K. Jeong, M. Choi and G. Lee, *Int. J. Precis. Eng. Manuf.*, 2022, 1–12.
- 51 J. Siekmann, A. Kulkarni, S. Akel, B. Klingebiel, M. Saliba, U. Rau and T. Kirchartz, *Adv. Energy Mater.*, 2023, **13**, 2300448.
- 52 Y. Yuan, G. Yan, C. Dreessen, T. Rudolph, M. Hülsbeck, B. Klingebiel, J. Ye, U. Rau and T. Kirchartz, *Nat. Mater.*, 2024, **23**, 391–397.
- 53 Y. Wang, S. Akel, B. Klingebiel and T. Kirchartz, *Adv. Energy Mater.*, 2024, **14**, 2302614.
- 54 P. Caprioglio, C. M. Wolff, O. J. Sandberg, A. Armin, B. Rech, S. Albrecht, D. Neher and M. Stollerfoht, *Adv. Energy Mater.*, 2020, **10**, 2000502.

