


 Cite this: *Chem. Commun.*, 2022, 58, 13891

 Received 5th November 2022,  
Accepted 21st November 2022

DOI: 10.1039/d2cc05993b

rsc.li/chemcomm

## *In situ* formation of inorganic healing overlayer for interface-stabilized all-inorganic CsPbIBr<sub>2</sub> perovskite solar cells†

 Junshuai Zhang,<sup>a</sup> Qiyao Guo,<sup>a</sup> Yuanyuan Zhao,<sup>b</sup> Jialong Duan<sup>ib</sup>\*<sup>a</sup> and Qunwei Tang<sup>ib</sup><sup>a</sup>

**A perovskite layer functionalized to be an outermost screen can strongly affect the capacity of the underlying device to avoid becoming decomposed under external stimuli, and subsequently affect the photovoltaic performance as well. Herein, we report an interface-stabilization strategy for an all-inorganic CsPbIBr<sub>2</sub> film involving forming *in situ* an inorganic ZrO<sub>2</sub> layer to solidify the soft perovskite lattice. As a result of defect passivation and self-encapsulation, the best device achieved an enhanced efficiency of up to 10.12%, which was much higher than the 7.47% efficiency for the reference device tested, with prolonged stability under conditions of persistent light irradiation and exposure to air.**

Although organic–inorganic hybrid metal halide perovskite solar cells (PSCs) have shown impressive achievements in power conversion efficiency (PCE) with certified values up to 25.7%, their inferior long-term stability remains the largest challenge to their further commercialization.<sup>1–4</sup> Considering the volatile nature of their organic species, *e.g.*, CH<sub>3</sub>NH<sub>3</sub><sup>+</sup> (MA<sup>+</sup>), CH<sub>3</sub>(NH<sub>2</sub>)<sub>2</sub><sup>+</sup> (FA<sup>+</sup>), under high temperature, completely substituting the organic components with inorganic cesium (Cs<sup>+</sup>) to fabricate all-inorganic PSCs has recently become more prevalent.<sup>5–7</sup> Unfortunately, the I-rich perovskite also suffers from serious phase transition when exposed to air.<sup>8,9</sup> By contrast, mixed-halide CsPbIBr<sub>2</sub> with a bandgap of ~2.11 eV can effectively balance the often contradictory goals of efficiency and stability, making it one of the top candidates for use as a photovoltaic material in tandem PSCs.<sup>10–12</sup> However, its reported efficiency deficit is much higher than that of a traditional hybrid device, with this deficit mainly related to the detrimental nonradiative recombination loss.

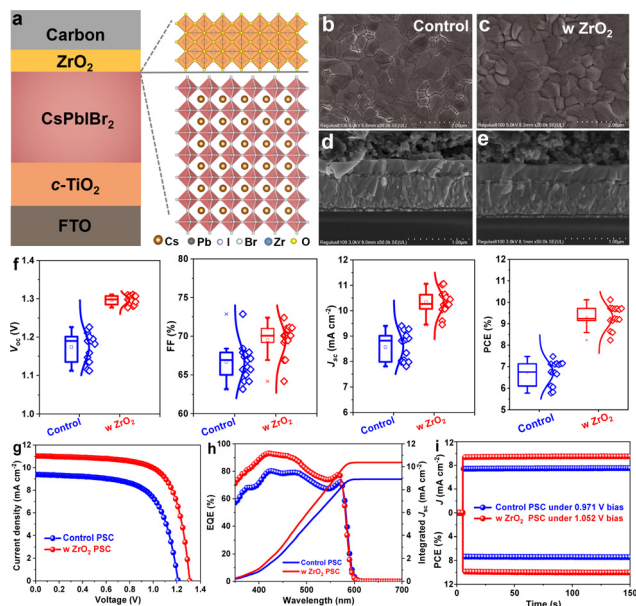
Due to the soft-lattice structure of perovskite, many crystallographic defects inevitably form during the rapid crystallization of its film, in particular point defects (such as vacancies, interstitials, and anti-site substitutions) and high-dimensional defects (such as grain boundaries and dislocations), which are regarded as detrimental charge recombination centers and initiation points to degrade the photoactive phase.<sup>13</sup> In particular, the top interface is extremely susceptible because of the direct invasion and accumulation of moisture and oxygen. Therefore, many physicochemical strategies for healing the defective layer have been reported, with examples of these strategies including mechanical surface polishing and Lewis acid-base passivation by organic molecules.<sup>14–18</sup> It is recognized that depositing a thin 2D perovskite onto a 3D-structured perovskite film can effectively stabilize a defective lattice in a hybrid species, with this stabilization mainly stemming from the incorporation of hydrophobic organic spacers and coordination of the dangling bonds to eliminate defects and block penetration of moisture/oxygen.<sup>19,20</sup> Such an approach, however, is difficult to implement for inorganic perovskites because the stronger binding of Cs<sup>+</sup> than of organic MA<sup>+</sup> or FA<sup>+</sup> hinders the cation exchange reaction.<sup>21,22</sup> Therefore, it is urgent to develop a new strategy to solidify the interface of all-inorganic PSCs.

To address the aforementioned issue, we demonstrated in the current work an *in situ* pyrolytic ZrO<sub>2</sub> as a multifunctional capping layer to stabilize the CsPbIBr<sub>2</sub> film by directly spin-coating zirconium *n*-propoxide precursor solution onto the film surface and then subjecting the resulting composite to an annealing treatment under high temperature. In contrast to sensitive organic passivators or low-dimensional perovskites, including an inorganic ZrO<sub>2</sub> overlayer benefits the stability of the film by rendering the film insensitive to environmental stresses as a result of the tighter lattice structure and covalent bonds produced. As schematically illustrated in Fig. 1a, the all-inorganic stack included fluorinated tin oxide (FTO), TiO<sub>2</sub>, CsPbIBr<sub>2</sub>, ZrO<sub>2</sub>, and a carbon electrode. Scanning electron

<sup>a</sup> Institute of Carbon Neutrality, College of Chemical and Biological Engineering, Shandong University of Science and Technology, Qingdao 266590, P. R. China.  
E-mail: duanjialong@sdu.edu.cn

<sup>b</sup> College of Mechanical and Electronic Engineering, Shandong University of Science and Technology, Qingdao 266590, P. R. China

† Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d2cc05993b>



**Fig. 1** (a) Illustration of the structure of the ZrO<sub>2</sub>-capped device. (b and c) Top-view SEM images of perovskite films with and without the ZrO<sub>2</sub> capping layer. (d and e) Cross-sectional SEM images of the PSCs. (f) Statistical photovoltaic parameters of control and optimal PSCs. (g) *J*-*V* curves, (h) EQE spectra and (i) steady power output curves of control and optimal PSCs.

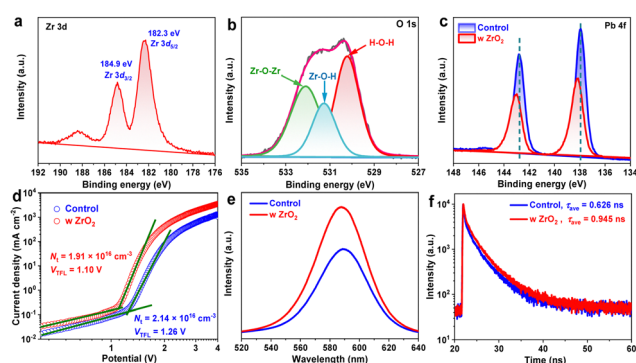
microscopy (SEM) images were acquired of the perovskite films to visualize their surface morphologies (Fig. 1b and c); a more compact topography occurred upon incorporation of the interface layer, and did so without destruction of the perovskite crystallization and absorption ability (Fig. S1, ESI<sup>†</sup>) owing to the filling of trenches. This feature was obvious from the cross-sectional morphologies of the PSCs (Fig. 1d, e and Fig. S2, ESI<sup>†</sup>), in which a ~10 nm-thick capping layer was observed. After optimizing the thickness by varying the precursor concentration (Fig. S3, ESI<sup>†</sup> and Table 1), the photovoltaic current density–voltage (*J*-*V*) curves of control and best PSCs were acquired, as shown in Fig. 1g with population statistics shown in Fig. 1f. Upon inserting the capping layer, increases were observed in all of the photovoltaic parameters including PCE, open-circuit voltage (*V*<sub>oc</sub>), short-circuit current density (*J*<sub>sc</sub>), and fill factor (FF). Specifically, the resultant carbon-based, all-inorganic CsPbI<sub>2</sub>Br<sub>2</sub> PSC delivered a significantly increased PCE of 10.12% with a *V*<sub>oc</sub> of 1.310 V, *J*<sub>sc</sub> of 11.027 mA cm<sup>-2</sup> and FF of 70.08%. The PCE value was much higher than the 7.47% value for the uncapped device, agreeing well with the integrated current densities from external quantum efficiency

**Table 1** Photovoltaic parameters of various PSCs. The concentration refers to the zirconium *n*-propoxide precursor solution

Devices	<i>V</i> <sub>oc</sub> (V)	<i>J</i> <sub>sc</sub> (mA cm <sup>-2</sup> )	FF (%)	PCE (%)
Control	1.211	9.395	65.66	7.47
0.01 mg mL <sup>-1</sup>	1.282	10.502	69.33	9.33
0.02 mg mL <sup>-1</sup>	1.310	11.027	70.08	10.12
0.05 mg mL <sup>-1</sup>	1.254	10.045	68.27	8.64

(EQE) spectra in Fig. 1h and steady power outputs in Fig. 1i. It should be noted that the efficiency gradually decreased with increasing thickness owing to the insulation of the ZrO<sub>2</sub> layer. Given the photogenerated carrier extraction-transfer-recombination behaviors, the performance improvement was mainly ascribed to the boosted charge collection rather than back-recombination at the interface.

To study the origin of the performance enhancement, specifically to reveal the positive effect of the ZrO<sub>2</sub> capping on the perovskite film, we first acquired X-ray photoelectron spectroscopy (XPS) spectra. Two characteristic peaks centered at 182.3 and 184.9 eV in the Zr 3d region of the spectrum of the capped perovskite film demonstrated the successful formation of ZrO<sub>2</sub> after the annealing treatment at 200 °C,<sup>23</sup> consistent with the detection of an O 1s signal (denoting the Zr–O–Zr and Zr–O–H bonding) (Fig. 2a and b). Moreover, the locations of these Zr 3d peaks at higher binding energies than the Zr 3d<sub>5/2</sub> and 3d<sub>3/2</sub> peaks of tetragonal ZrO<sub>2</sub> species at 182.0 and 184.3 eV<sup>23</sup> demonstrated the occurrence of electron donation. Together with the peak difference between the Pb 4f XPS spectra of uncoated and ZrO<sub>2</sub>-coated perovskite films, shown in Fig. 2c, and the lack of any considerable difference when comparing the corresponding Cs 3d, I 3d and Br 3d spectra, shown in Fig. S4 (ESI<sup>†</sup>), our results provided confirmatory evidence for the strong interaction between the CsPbI<sub>2</sub>Br<sub>2</sub> film and ZrO<sub>2</sub> owing to the presence of spontaneous transfer of electrons from ZrO<sub>2</sub> to the undercoordinated Pb<sup>2+</sup> ions.<sup>24</sup> As a result, the defective surface was well healed, significantly eliminating the charge recombination centers. To examine the defect density (*N*<sub>t</sub>) levels in capped and uncapped perovskite films, we took space-charge-limited current (SCLC) measurements based on the structure of FTO/c-TiO<sub>2</sub>/perovskites/PCBM/carbon. As shown in Fig. 2d, the defect densities were calculated to be 2.14 × 10<sup>16</sup> cm<sup>-3</sup> and 1.91 × 10<sup>16</sup> cm<sup>-3</sup> for control and capped perovskite films according to the relationship<sup>25</sup>  $V_{TFL} = (eNL^2)/(2\epsilon\epsilon_0)$ , where *V*<sub>TFL</sub> represents the trap-filled limit voltage, *e* the elementary charge, *L* the thickness of the perovskite film (~200 nm from SEM images),  $\epsilon$  the relative dielectric constant,



**Fig. 2** (a) Zr 3d and (b) O 1s regions of the XPS spectrum of the ZrO<sub>2</sub>-capped perovskite film. (c) Pb 4f XPS spectra of perovskite films with and without a ZrO<sub>2</sub> layer. (d) Dark *J*-*V* curves of electron-only devices for SCLC analysis. (e) Steady PL spectra and (f) TRPL decay curves of the indicated perovskite films.

and  $\epsilon_0$  the vacuum permittivity. The reduced  $N_t$  was undoubtedly responsible for the higher  $V_{oc}$  and  $J_{sc}$  in the PSC, in accordance with the  $J$ - $V$  results. Steady photoluminescence (PL) spectra, shown in Fig. 2e, also revealed the remarkably increased PL intensity for the perovskite film coated with the  $ZrO_2$  layer, demonstrating the preferred radiative recombination rather than nonradiative recombination as a result of the reduction of detrimental defects.<sup>26,27</sup> In other words, the photo-generated carriers survived longer, *i.e.*, with a greater lifetime ( $\tau = 0.945$  ns), in the optimal perovskite film than in the control ( $\tau = 0.626$  ns) according to recorded time-resolved PL (TRPL) decay curves (Fig. 2f, with the corresponding parameters summarized in Table S1, ESI†). All these results validated the positive effect of the  $ZrO_2$  capping layer on the  $CsPbIBr_2$  film and the suppression of defect-induced carrier recombination.

Considering the dependence of photovoltaic performance on the perovskite film quality, the charge recombination dynamics in the PSC was then investigated. As shown in Fig. 3a, capacitance–voltage ( $C$ - $V$ ) measurements were taken to determine the built-in potential ( $V_{bi}$ ) evolution, which determines the level of charge extraction.<sup>28</sup> As expected, a larger  $V_{bi}$  was measured for the  $ZrO_2$ -treated film. This larger  $V_{bi}$  ensured a greater driving force for carrier extraction-transfer. In other words, the nonconductive inorganic  $ZrO_2$  has no effect on charge extraction owing to the tunnelling effect.<sup>29</sup> Based on the above characterizations, we speculatively proposed an intrinsic mechanism behind the performance enhancement. In the presence of the  $ZrO_2$  layer, two undesirable electron recombination channels were closed in the optimal PSC: (i) occurrence of the defect-electron coupling at the perovskite surface was inhibited due to defect passivation by  $ZrO_2$ ; and (ii) the larger  $V_{bi}$  occurring upon  $ZrO_2$  treatment hindered the migration of electrons from the perovskite to the carbon electrode to recombine with the hole. We further utilized the

dependences of  $J_{sc}$  and  $V_{oc}$  on light intensity ( $I$ ) to investigate the internal recombination mechanism. A linear  $I \sim J_{sc}$  relationship was found for both devices, as shown in Fig. S5 (ESI†), demonstrating the smooth charge extraction and transport at the PSC interface. However, compared to the 1.72 value of the ideality factor of the control device under the open-circuit condition, that of the device containing the  $ZrO_2$  capping was significantly lower, with a value of 1.42, as shown in Fig. 3b, an indicator of suppressed defect-assisted Shockley–Read–Hall (SRH) recombination;<sup>30</sup> that is, the corresponding recombination resistance was increased (Fig. 3c) and the lifetime of photogenerated carrier located at the conduction band in the PSC was increased with a slower  $V_{oc}$  decay (Fig. S6, ESI†).<sup>31</sup> As a result, at voltages  $< 1.1$  V, the overall leakage current density of the capped PSC in the dark condition was one order of magnitude lower than that of the control device (Fig. 3d), while higher current density values were found at voltages  $> 1.5$  V, mainly attributed to the boosted charge injection,<sup>32</sup> and again confirming the positive effect of  $ZrO_2$  on healing the defective perovskite surface.

Finally, we aged the unencapsulated control and  $ZrO_2$ -capped PSCs under persistent light irradiation and dry air at a temperature ( $T$ ) of 25 °C and relative humidity (RH) of 10%, which was controlled by using a drybox with an automatic adjustment function. As shown in Fig. 4a, b and Fig. S7 (ESI†), the efficiency level of the latter device remained at 50% and 90% of the initial efficiency after operation over 65 hours and storage over 60 days, respectively, while the performance of the control device degraded more sharply. Considering the higher contact angle measured for the capped film than for the uncapped one (Fig. 4c and d), the improved stability of the capped device can be attributed to the self-encapsulation effect of the inert  $ZrO_2$  layer on the underlying perovskite film,<sup>33–35</sup> inhibiting the direct penetration of external moisture/oxygen

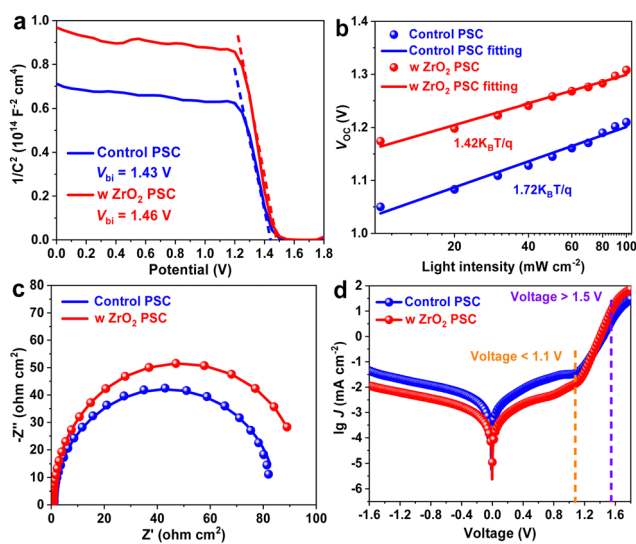


Fig. 3 (a)  $C$ - $V$  curves, (b) dependence of  $V_{oc}$  on light intensity, (c) EIS spectra, and (d) dark  $J$ - $V$  curves of control and  $ZrO_2$ -capped  $CsPbIBr_2$  PSCs.

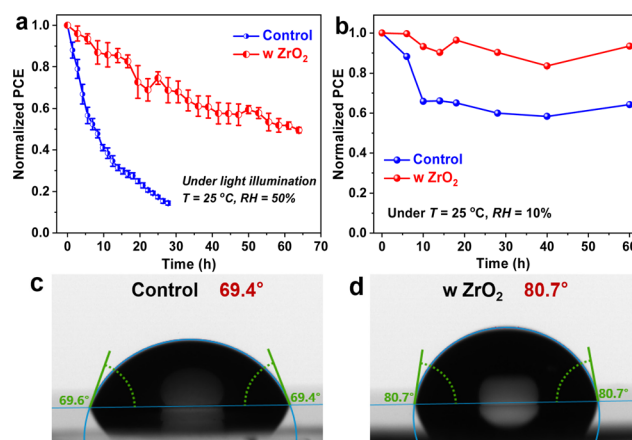


Fig. 4 (a and b) Stability levels of unencapsulated control and  $ZrO_2$ -capped PSCs under (a) persistent light irradiation and (b) dry air conditions. The errors in the data (indicated by error bars) originate from the fluctuation of the light intensity of the xenon lamp. (c and d) Contact angles of water on perovskite films before and after they were capped with a  $ZrO_2$  layer at  $T = 25$  °C and  $RH = 50\%$ .

into the perovskite lattice that would have otherwise deteriorated the photoactive phase.

In summary, we demonstrated an effective and feasible strategy to heal the defective surface of an inorganic perovskite film by *in situ* forming an inorganic ZrO<sub>2</sub> covering layer to overcome the bottleneck in converting the 3D crystal into a 2D structure. As a result of the migration of electrons from ZrO<sub>2</sub> to the perovskite surface, the detrimental defects were passivated, leading to suppressed nonradiative recombination. Finally, the best device achieved a significantly increased PCE of up to 10.12% for a carbon-based, fully inorganic CsPbIBr<sub>2</sub> PSC. In addition to efficiency, more importantly, the robust interface separated perovskite from environmental moisture/oxygen stresses owing to the self-encapsulation of the inert ZrO<sub>2</sub> layer, leading to a remarkable enhancement of the durability of the PSC.

The authors gratefully acknowledge financial support provided by the National Natural Science Foundation of China (22109053), the Guangdong Basic and Applied Basic Research Foundation (2020A1515110548), the Guangzhou Science and Technology Planning Project (202102020775, 202102010091), and National Key Research and Development Program of China (2021YFE0111000).

## Conflicts of interest

There are no conflicts to declare.

## Notes and references

- 1 A. Kojima, K. Teshima, Y. Shirai and T. Miyasaka, *J. Am. Chem. Soc.*, 2009, **131**, 6050–6051.
- 2 Y. Zhao, F. Ma, Z. Qu, S. Yu, T. Shen, H.-X. Deng, X. Chu, X. Peng, Y. Yuan, X. Zhang and J. You, *Science*, 2022, **377**, 531–534.
- 3 N. Li, X. Niu, Q. Chen and H. Zhou, *Chem. Soc. Rev.*, 2020, **49**, 8235–8286.
- 4 Y. Zhou, L. M. Herz, A. K.-Y. Jen and M. Saliba, *Nat. Energy*, 2022, **7**, 794–807.
- 5 J. Zhang, G. Hodes, Z. Jin and S. (Frank) Liu, *Angew. Chem., Int. Ed.*, 2019, **58**, 15596–15618.
- 6 Y. Cui, J. Shi, F. Meng, B. Yu, S. Tan, S. He, C. Tan, Y. Li, H. Wu, Y. Luo, D. Li and Q. Meng, *Adv. Mater.*, 2022, **34**, 2205028.
- 7 J. Liang and Y. B. Qi, *Mater. Today Nano*, 2021, **16**, 100143.
- 8 H. Yao, Z. Li, C. Shi, Y. Xu, Q. Wang, Z. Li, G. Peng, Y. Lei, H. Wang, Z. Ci and Z. Jin, *Adv. Funct. Mater.*, 2022, **32**, 2205029.
- 9 Y. Lei, Z. Li, H. Wang, Q. Wang, G. Peng, Y. Xu, H. Zhang, G. Wang, L. Ding and Z. Jin, *Sci. Bull.*, 2022, **67**, 1352–1361.
- 10 W. Zhu, Z. Zhang, D. Chen, W. Chai, D. Chen, J. Zhang, C. Zhang and Y. Hao, *Nano-Micro Lett.*, 2020, **12**, 87.
- 11 J. Wang, X. Wu, Y. Liu, Q. Xue, H.-L. Yip, A. K. Y. Jen and Z. Zhu, *Energy Technol.*, 2021, **9**, 2100562.
- 12 Q. Liu, J. Qiu, X. Yan, Y. Fei, Y. Qiang, Q. Chang, Y. Wei, X. Zhang, W. Tian, S. Jin, Z. Yu and L. Sun, *J. Energy Chem.*, 2022, **74**, 387–393.
- 13 Y. Zheng, X. Yang, R. Su, P. Wu, Q. Gong and R. Zhu, *Adv. Funct. Mater.*, 2020, **30**, 2000457.
- 14 Y. Xu, L. Zhai, L. Sun, J. Wang, X. Tan, H. Huang, Y. Wang, G. Yang, K. Jiang, Y. Yang, L. Zhang, Z. Tan and C. Zou, *Chem. Commun.*, 2022, **58**, 7132–7135.
- 15 S. Chen, Y. Liu, X. Xiao, Z. Yu, Y. Deng, X. Dai, Z. Ni and J. Huang, *Joule*, 2020, **4**, 2661–2674.
- 16 A. Yusoff, M. Vasilopoulou, D. G. Georgiadou, L. C. Palilis, A. Abate and M. K. Nazeeruddin, *Energy Environ. Sci.*, 2021, **14**, 2906–2953.
- 17 R. Zhao, L. Xie, R. Zhuang, T. Wu, R. Zhao, L. Wang, L. Sun and Y. Hua, *ACS Energy Lett.*, 2021, **6**, 4209–4219.
- 18 G. Liu, Y. Zhong, W. Feng, M. Yang, G. Yang, J.-X. Zhong, T. Tian, J.-B. Luo, J. Tao, S. Yang, X.-D. Wang, L. Tan, Y. Chen and W.-Q. Wu, *Angew. Chem., Int. Ed.*, 2022, **61**, e202209464.
- 19 D. Zhao, D. Gao, X. Wu, B. Li, S. Zhang, Z. Li, Q. Wang, Z. Wu, C. Zhang, W. C. H. Choy, X. Zhong, Q. He and Z. Zhu, *Adv. Energy Mater.*, 2022, **34**, 2204661.
- 20 M. A. Mahmud, T. Duong, J. Peng, Y. Wu, H. Shen, D. Walter, H. T. Nguyen, N. Mozaffari, G. D. Tabi, K. R. Catchpole, K. J. Weber and T. P. White, *Adv. Funct. Mater.*, 2022, **32**, 2009164.
- 21 X. Zhao, T. Liu, Q. C. Burlingame, T. Liu, R. Holley III, G. Cheng, N. Yao, F. Gao and Y.-L. Loo, *Science*, 2022, **377**, 307–310.
- 22 Y. Wang, T. Zhang, M. Kan and Y. Zhao, *J. Am. Chem. Soc.*, 2018, **140**, 12345–12348.
- 23 H. Li, Y. Wu, C. Li, Y. Gong, L. Niu, X. Liu, Q. Jiang, C. Sun and S. Xu, *Appl. Catal., B*, 2019, **251**, 305–312.
- 24 Y. Wang, H. Yang, K. Zhang, M. Tao, M. Li and Y. Song, *ACS Energy Lett.*, 2022, **7**, 3646–3652.
- 25 S. Wang, H. Guo, J. Wu, Y. Lei, X. Li, Y. Fang, Y. Dai, W. Xiang and Y. Lin, *Chem. Commun.*, 2022, **58**, 8384–8387.
- 26 S. Wang, X. Zhang, W. Zhu, Z. Tang, J. Liu, H. Zhang, L. Ding and F. Hao, *Appl. Surf. Sci.*, 2022, **602**, 154393.
- 27 Q. Guo, J. Duan, J. Zhang, Q. Zhang, Y. Duan, X. Yang, B. He, Y. Zhao and Q. Tang, *Adv. Mater.*, 2022, **34**, 2202301.
- 28 H. Wang, F. Ye, J. Liang, Y. Liu, X. Hu, S. Zhou, C. Chen, W. Ke, C. Tao and G. Fang, *Joule*, 2022, DOI: [10.1016/j.joule.2022.10.001](https://doi.org/10.1016/j.joule.2022.10.001).
- 29 J. Li, F. Yan, P. Yang, Y. Duan, J. Duan and Q. Tang, *Sol. RRL*, 2021, **6**, 2100791.
- 30 B. Li, Z. Li, X. Jiang, Z. Wang, Y. Rao, C. Zhang, M. Zhu, L. Zhang, L. Wen, S. K. So, Y. Zhou, S. Pang and Z. Zhou, *Appl. Surf. Sci.*, 2022, **603**, 154410.
- 31 Y. Xie, J. Feng, M. Chen, X. Zhu, Y. Zhou, Z. Li, D. Yang and S. F. Liu, *ACS Appl. Energy Mater.*, 2022, **5**, 8034–8041.
- 32 F. Yan, P. Yang, J. Li, Q. Guo, Q. Zhang, J. Zhang, Y. Duan, J. Duan and Q. Tang, *Chem. Eng. J.*, 2022, **430**, 132781.
- 33 T. Tian, J.-X. Zhong, M. Yang, W. Feng, C. Zhang, W. Zhang, Y. Abdi, L. Wang, B.-X. Lei and W.-Q. Wu, *Angew. Chem., Int. Ed.*, 2021, **60**, 23735–23742.
- 34 Y. Zhan, J. Peng, C. Cao and Q. Cheng, *Nano Energy*, 2022, **101**, 107575.
- 35 T. Liu, Y. Zhou, Z. Li, L. Zhang, M.-G. Ju, D. Luo, Y. Yang, M. Yang, D. H. Kim, W. Yang, N. P. Padture, M. C. Beard, X. C. Zeng, K. Zhu, Q. Gong and R. Zhu, *Adv. Mater.*, 2018, **8**, 1800232.