



Cite this: *Mater. Adv.*, 2020,
1, 2057

A bilayer $\text{TiO}_2/\text{Al}_2\text{O}_3$ as the mesoporous scaffold for enhanced air stability of ambient-processed perovskite solar cells†

Dong Wang,^a Qian Chen,^{ID} *^a Hongbo Mo,^b Janet Jacobs,^c Andrew Thomas^{a,c} and Zhu Liu^{*a}

Research on the air stability of perovskite solar cells (PSCs) has attracted great attention in the field of photovoltaics. Although devices based on lead halide perovskites show outstanding efficiency, their commercialization is hampered by the instability of the perovskite under exposure to moisture in ambient air. The use of mesoporous scaffolds has been considered as a promising method to protect the perovskite from moisture ingress. Here, we demonstrate a $\text{TiO}_2/\text{Al}_2\text{O}_3$ bilayer as the mesoporous scaffold to enhance the air stability of PSCs. PSCs based on the $\text{TiO}_2/\text{Al}_2\text{O}_3$ mesoporous scaffold were assembled in ambient air at a relative humidity (RH) of over 65%, delivering a champion power conversion efficiency (PCE) of 16.84%. More importantly, PSCs based on $\text{TiO}_2/\text{Al}_2\text{O}_3$ retained 82% of their initial PCE after storage in ambient air for 2000 h without any encapsulation. In comparison, PSCs based on a single layer of mesoporous TiO_2 retained only 57% of their initial PCE. In addition, the $\text{TiO}_2/\text{Al}_2\text{O}_3$ mesoporous scaffold shows no reduction in light transmission in the visible region in comparison to a single layer of TiO_2 , indicating its potential to be used for semi-transparent and tandem PSC applications.

Received 3rd August 2020,
Accepted 17th August 2020

DOI: 10.1039/d0ma00562b

rsc.li/materials-advances

Introduction

Hybrid organic–inorganic lead halide perovskite materials have achieved incredible progress in their application in low-cost and highly efficient photovoltaic devices in the past few years. Long-term stability, however, is still one of the main challenges that impede the commercialization of perovskite solar cells (PSCs), due to the instability of the perovskite in ambient air,^{1–3} particularly under exposure to moisture in the air.⁴ The ingress of water molecules induces the formation of unstable hydrates within MAPbI_3 lattices, which further degrade into PbI_2 and other products.⁵ PbI_2 is an intrinsic or a weak p-type semiconductor with a wide bandgap of 2.37 eV at room temperature.⁶ The position of its valence band edge relative to MAPbI_3 prevents the backflow of

electrons from the electron transport material (e.g. TiO_2) towards the perovskite. Several studies have utilized this passivation effect to eliminate defect states at the TiO_2 /perovskite interface and further to reduce the recombination rate.^{7,8} However, the PbI_2 -induced enhancement of device performance can only be achieved by precise morphological control of PbI_2 crystals.⁹ The formation of a large amount of PbI_2 at the interface between the perovskite and the charge transport layer hinders efficient charge extraction.⁹ The high conduction band edge and the low valence band edge of PbI_2 relative to the perovskite make it a barrier not only for the backflow but also for the forward flow of charge carriers.¹⁰ Furthermore, PbI_2 is a water-soluble, toxic material (solubility of 0.063 g cm^{-3} at 20°C),¹¹ which brings its own environmental concerns.

In order to minimize moisture-induced degradation of perovskite-based solar cells, a variety of approaches have been taken, such as compositional tailoring,^{12,13} interfacial engineering,^{14,15} encapsulation^{16,17} and so on. One attractive approach is the use of mesoscopic engineering.^{18–21} Mesoscopic structures have been extensively employed in PSCs, where they mainly function as a large-surface-area supporting scaffold and/or charge transport medium. Several studies have reported enhanced air stability offered by the mesoscopic structures in comparison to planar heterojunctions,^{22–24} despite the fact that planar devices are benefiting from simple architecture and a low-temperature process.²⁵

^a Department of Materials, the University of Manchester, Oxford Road, Manchester M13 9PL, UK. E-mail: qian.chen-2@manchester.ac.uk, zhu.liu@manchester.ac.uk

^b Department of Mechanical, Aerospace and Civil Engineering, the University of Manchester, Oxford Road, Manchester M13 9PL, UK

^c Photon Science Institute, the University of Manchester, Oxford Road, Manchester M13 9PL, UK

† Electronic supplementary information (ESI) available: Relative humidity (RH); SEM images of mesoporous scaffolds; PV performances of the PSCs based on different mesoporous scaffolds with varying thickness; air stability of the PSCs based on different mesoporous scaffolds; hysteresis behavior shown in the aged PSCs based on different mesoporous scaffolds. See DOI: 10.1039/d0ma00562b



TiO₂ is one of the most widely used materials for mesoscopic electron transport, with suitable band alignment relative to the perovskite²⁶ and a high electron mobility rate,²⁷ which facilitates the extraction of photogenerated electrons. Owing to the released polarization in perovskite lattices by improved charge injection from the perovskite to mesoporous TiO₂, PSCs structured on the mesoporous TiO₂ scaffold exhibit reduced hysteresis relative to planar heterojunctions.^{28,29} Furthermore, mesoscopic TiO₂-based PSCs exhibit enhanced moisture resistance, which can also be ascribed to the rapid charge extraction process.³⁰ Inspired by the excellent photovoltaic performance and improved air stability of devices based on mesoscopic TiO₂, many efforts have been made to fabricate stable and efficient PSCs by design and control of mesoscopic architecture. Some groups have demonstrated a simple method to adjust the mesoscopic TiO₂ structure to enhance air stability by increasing its thickness, unfortunately, with the sacrifice of power conversion efficiency (PCE).^{22,23} This is mainly ascribed to increased series resistance and recombination rate, as a result of long electron diffusion length in TiO₂ and potential exposure of TiO₂ nanoparticles to the hole transporter when a thickened TiO₂ scaffold is employed.^{31,32} In addition, the defect states below the conduction band edge of TiO₂ possibly trap the injected photoelectrons and thus cause recombination at the perovskite/TiO₂ interface when the solar cell is operated under UV light exposure.³³ Generally, loss of photovoltaic performance caused by increased thickness of the mesoporous TiO₂ scaffold occurs in the charge transfer/collection process in TiO₂. Al₂O₃, as an insulator, was pioneered by Snaith³⁴ to fabricate mesoscopic Al₂O₃-based PSCs, where Al₂O₃ only functions as the support, resulting in higher open-circuit voltage (V_{oc}) than that of the TiO₂ counterparts. In addition, a tenfold increase in the rate of charge collection is observed in Al₂O₃-based devices compared to that observed in the TiO₂ counterparts, indicating faster electron diffusion through the perovskite phase, than through the n-type TiO₂.³⁴ Further developments based on this strategy employed insulating metal oxides, such as Al₂O₃ or ZrO₂, as a spacer between the electron transport and hole transport materials to form a triple-^{19,20} or a quadruple-²¹ layer mesoscopic architecture, with a ~ 10 μ m thick mesoscopic carbon electrode for better air stability. However, the opaque carbon layer does not allow incident light to go through it, leading to a low light transmittance of the overall device. In addition, PSCs based on the triple-layer TiO₂/ZrO₂/carbon architecture commonly suffer from severe charge recombination than the conventional mesoscopic PSCs based on a TiO₂ scaffold (thickness < 200 nm), due to the longer charge carrier transport distance.³⁵ Recently, Wang *et al.* have demonstrated the addition of a super thin Al₂O₃ layer (thickness < 2.4 nm) between the mesoporous TiO₂ and ZrO₂ layers to reduce the charge recombination for the PSCs based on the TiO₂/ZrO₂/carbon architecture. As a result, the thickness of the ZrO₂ layer for the PSCs with the Al₂O₃ layer can be reduced from 3 μ m to 1.2 μ m and the modified devices also showed reduced charge recombination.³⁵

In this work, we demonstrate a bilayer mesoporous scaffold with enhanced air stability for PSCs by depositing a mesoporous Al₂O₃ layer on top of the mesoporous TiO₂ layer, upon which MAPbI_xCl_{3-x} based mesoscopic PSCs were assembled.

The additional mesoporous Al₂O₃ layer is expected to increase the thickness of the infiltrating perovskite, thus enhancing the air stability of the devices without sacrificing the PCE. As a result, fully ambient-processed PSCs based on the TiO₂/Al₂O₃ mesoporous scaffold were found to deliver the highest PCE of 16.84%, which is slightly higher than the PCE of PSCs based on a single layer of TiO₂ (16.43%). After storing in highly humidified air (>65% relative humidity, RH) for over 2000 h without any encapsulation, the PSCs based on TiO₂/Al₂O₃ retained 82% of their initial PCE, in comparison to PSCs based on TiO₂ alone which exhibited only 57% of their initial PCE. In addition, the TiO₂/Al₂O₃ mesoporous scaffold shows no reduction in light transmission in the visible region, indicating its potential to be used for semi-transparent and tandem PSC applications.

Experimental section

Device fabrication

Pre-patterned indium tin oxide (ITO) glass substrates (Ossila, 20 × 15 mm, 20 Ω sq⁻¹) were ultrasonically cleaned using 3 vol% Hellmanex III solution, ethanol, and deionized water for 10 min each, and then treated with UV-ozone for 15 min. The compact and mesoporous TiO₂ layers were produced in the same way as we have previously reported.³⁶ For the TiO₂/Al₂O₃ bilayer mesoporous scaffold-based devices, a home-made Al₂O₃ paste composed of 1.5 g dispersion of 20 wt% Al₂O₃ in isopropanol (Sigma-Aldrich, <50 nm) and an organic binder (4 g terpineol and 0.5 g ethyl cellulose mixed in isopropanol) was spin-coated on top of the mesoporous TiO₂ layer at 6000 rpm for 30 s to form a mesoporous Al₂O₃ layer, followed by drying at 125 °C for 10 min. The substrate was then annealed at 500 °C for 30 min in a muffle furnace. For the single mesoporous TiO₂ or Al₂O₃ scaffold-based devices, the mesoporous TiO₂ or Al₂O₃ layer was directly produced on top of the compact TiO₂ layer. The thickness of the mesoporous Al₂O₃ layer was controlled by adjusting the mass ratio of Al₂O₃ to organic binder in the paste used for spin coating. The thickness of the mesoporous TiO₂ layer was controlled by adjusting the dilution rate of the commercial TiO₂ paste used for spin coating. The perovskite film was generated *via* the widely used anti-solvent method: 460 mg of PbI₂ (Sigma-Aldrich, 99.9985%) and 28 mg of PbCl₂ (Sigma-Aldrich, 99.999%) were dissolved in 200 μ l of DMSO (Sigma-Aldrich, $\geq 99.7\%$) and 850 μ l of DMF (Sigma-Aldrich, $\geq 99.8\%$) at 70 °C. After cooling down to room temperature, 190 mg of MAI (Ossila, 98%) was added into the lead solution to form the perovskite precursor. 90 μ l of the prepared perovskite precursor was then spin-coated onto the preheated substrate (preheated at 70 °C for 10 min) at 4000 rpm for 30 s, during which 200 μ l of ethyl acetate (Sigma-Aldrich, 99.8%) was dripped onto the substrate from the 6th second to the 9th second after the start of spinning.³⁷ The substrate was then annealed at 115 °C for 10 min. Subsequently, 80 μ l of spiro-MeOTAD solution, which consisted of 43 mg of spiro-MeOTAD (Sigma-Aldrich, 99%), 10 μ l of Li-TSFI solution (520 mg Li-TSFI in 1 ml acetonitrile (Sigma-Aldrich, 99.8%)) and 15 μ l of 4-*tert*-butyl



pyridine in 0.5 ml of chlorobenzene, was spin-coated on the perovskite film at 4000 rpm for 30 s. Finally, Au electrodes were deposited *via* thermal evaporation. The relative humidity was measured using a digital humidity meter (Fisher Scientific). All processes were conducted under ambient air with a RH of over 65% without sealing (ESI,† Fig. S1).

Materials and device characterization

The morphologies of mesoporous scaffolds and the perovskite films coated on them were observed through a field-emission scanning electron microscope (Ultra-55, Carl Zeiss). The phase transformations in perovskite films were characterized using X-ray diffraction (PANalytical) with Cu K α radiation in the 2θ range from 10° to 30° with the glancing incidence fixed at 2° . The light absorption spectra of perovskite films were acquired using a UV-vis spectrophotometer (Shimadzu UV-2401PC). The current density–voltage (J – V) characteristics for the fresh and the aged solar cell devices were measured using an Oriel solar simulator equipped with a Keithley 2420 source meter under 100 mW cm^{-2} (AM 1.5G) power calibrated using an NREL certified reference cell. A square metal mask with an area of 0.024 cm^2 was used to determine the effective area under illumination. The transient photo-current and transient-voltage measurements were performed using a Thorlabs solid-state laser diode of wavelength 405 nm (DL5146-101S) as the light source. The transient response was recorded using a Keysight DSOS604A digital storage oscilloscope, which, in the case of

photocurrent measurements, was connected across a small short circuit sampling resistor of 100Ω . For open-circuit photovoltage measurements, the oscilloscope was connected directly across the terminals of the device. Averages were recorded typically over many tens of transients for each measurement.

Results and discussion

The thickness of the fabricated single mesoporous TiO_2 and bilayer mesoporous $\text{TiO}_2/\text{Al}_2\text{O}_3$ films is $\sim 150 \text{ nm}$ and $\sim 300 \text{ nm}$, respectively (ESI,† Fig. S2). These films were used as supporting scaffolds to prepare perovskite solar cells in ambient air at $>65\%$ RH (thermal evaporation of the gold electrodes was carried out in vacuum). Fig. 1a shows the schematics of the device configurations based on different mesoporous scaffolds. The perovskite absorber is infiltrated within the mesoporous scaffold and forms a confined mesoscopic morphology covered by a compact capping layer. As illustrated by the band diagram in Fig. 1a, the Al_2O_3 film acts as both a supporting scaffold and a spacer between the electron and hole transporting layers (ETL, HTL). The Al_2O_3 spacer may block the penetration of spiro-MeOTAD or gold, thereby avoiding direct contact with TiO_2 , which will be further discussed below. Here, however, the main purpose of introducing the additional mesoporous Al_2O_3 layer is to enhance the air stability of the device by increasing the thickness of the perovskite that permeates the mesoporous scaffold. As can be

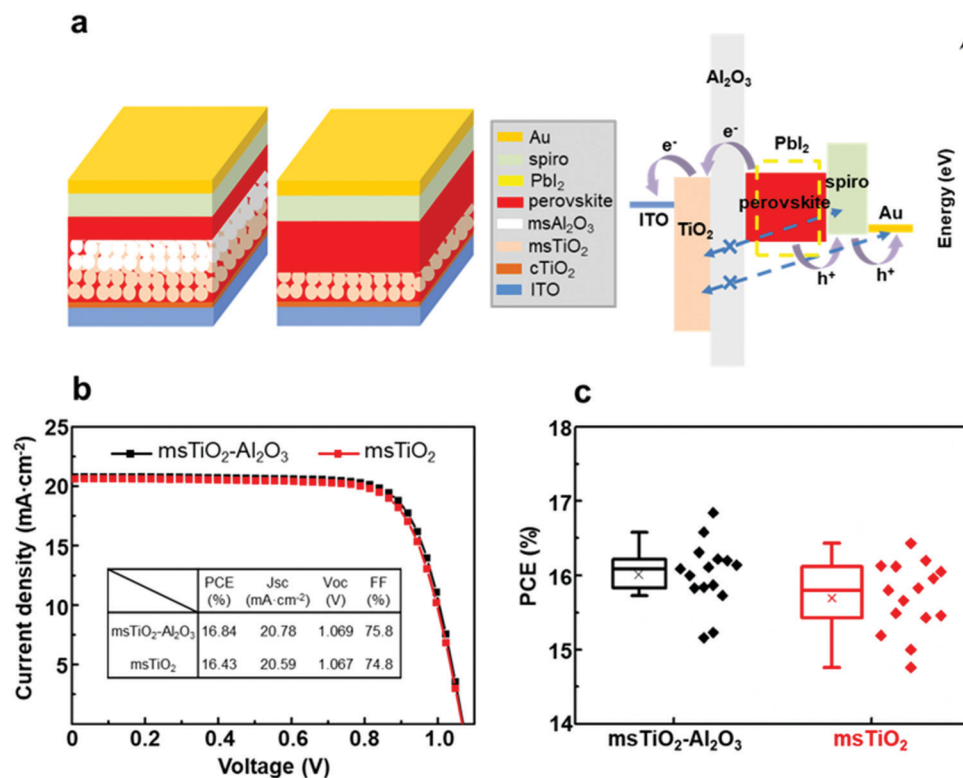


Fig. 1 (a) Schematics of device configurations of the PSCs based on mesoporous $\text{TiO}_2/\text{Al}_2\text{O}_3$ and mesoporous TiO_2 scaffolds, together with the band diagram of the PSC based on the mesoporous $\text{TiO}_2/\text{Al}_2\text{O}_3$ scaffold. (b) Current density–voltage curves of the best cells in their respective counterparts. The inset shows the photovoltaic parameters. (c) PCE distributions of the PSCs based on different mesoporous scaffolds.



seen from the band alignment diagram between perovskite and the degradation product PbI_2 , PbI_2 has a wider bandgap than the perovskite. This implies impeded charge transfer possibly occurring as a result of PbI_2 formation.^{9,10} We turn now to discuss the relationship between photovoltaic performance and air stability in the mesoscopic single layer- and bilayer-based PSCs.

Firstly, we measured the PCE of the PSCs based on bilayer mesoporous scaffolds with a varying Al_2O_3 layer thickness of 50–100 nm, 150 nm, 300 nm, 500 nm and 700 nm (see ESI,† Fig. S3 and Table S1). It was found that the PCE decreases dramatically with the thickness of the Al_2O_3 layer >150 nm. This result is possibly due to a longer charge carrier transport distance for the PSCs based on the thicker Al_2O_3 layer. Due to the difficulty in preparing a uniform ultra-thin mesoscopic layer, we only obtained a mesoporous Al_2O_3 layer with uneven thickness ranging from 50 nm to 100 nm. The PSCs based on $\text{TiO}_2/\text{Al}_2\text{O}_3$ with the Al_2O_3 thickness fluctuating between 50 nm and 100 nm show lower PCE than those based on the uniform bilayer mesoporous scaffold with an Al_2O_3 thickness of 150 nm. This is possibly due to the negative effect of the wave-like surface morphology of the mesoporous scaffold on the growth of perovskite crystals. The highest PCE for the fully ambient-processed PSC is based on the bilayer $\text{TiO}_2/\text{Al}_2\text{O}_3$ mesoporous scaffold with 150 nm-thick Al_2O_3 , with a champion PCE of 16.84% and a mean value of 16.01%. This result is slightly

higher than the PCE for single TiO_2 devices (champion = 16.43%, mean = 15.70%) (Fig. 1b and c).

The stability of PSCs was examined by storing the devices under ambient air with a RH > 65% for 12 weeks (~2000 h) during which the photovoltaic parameters (PCE, open-circuit voltage V_{oc} , short-circuit current density J_{sc} , fill factor FF) were measured weekly. Normalized PCE, V_{oc} , J_{sc} , and FF of the PSCs, based on different mesoporous scaffolds, are shown in Fig. 2a–d as a function of time. In order to obtain reliable results, each point was plotted as an average over 8 cells in the same counterparts. Fig. 2b and d show that the V_{oc} and FF of both single and bilayer mesoporous scaffold-based PSCs decrease slightly and roughly linearly. However, a distinguishable difference for the J_{sc} between the two types of devices is more marked in Fig. 2c: the single mesoporous scaffold-based PSC shows a sudden drop after 4 weeks and shows a reduction of ~30% in the J_{sc} after 12 weeks' storage. The bilayer mesoporous scaffold-based PSC, on the other hand, retains ~95% of its initial J_{sc} . A stable J_{sc} over time suggests more uniform perovskite loading in the bilayer mesoporous scaffold, which is further confirmed in our study. As a result, the PCE of the bilayer $\text{TiO}_2/\text{Al}_2\text{O}_3$ mesoporous scaffold-based solar cell after 12 weeks of storage in ambient air retains 82% of its initial value, while in single TiO_2 mesoporous scaffold-based device the PCE falls to 57% of the original (ESI,† Fig. S4). In addition, we also studied the stability of the single Al_2O_3 mesoporous

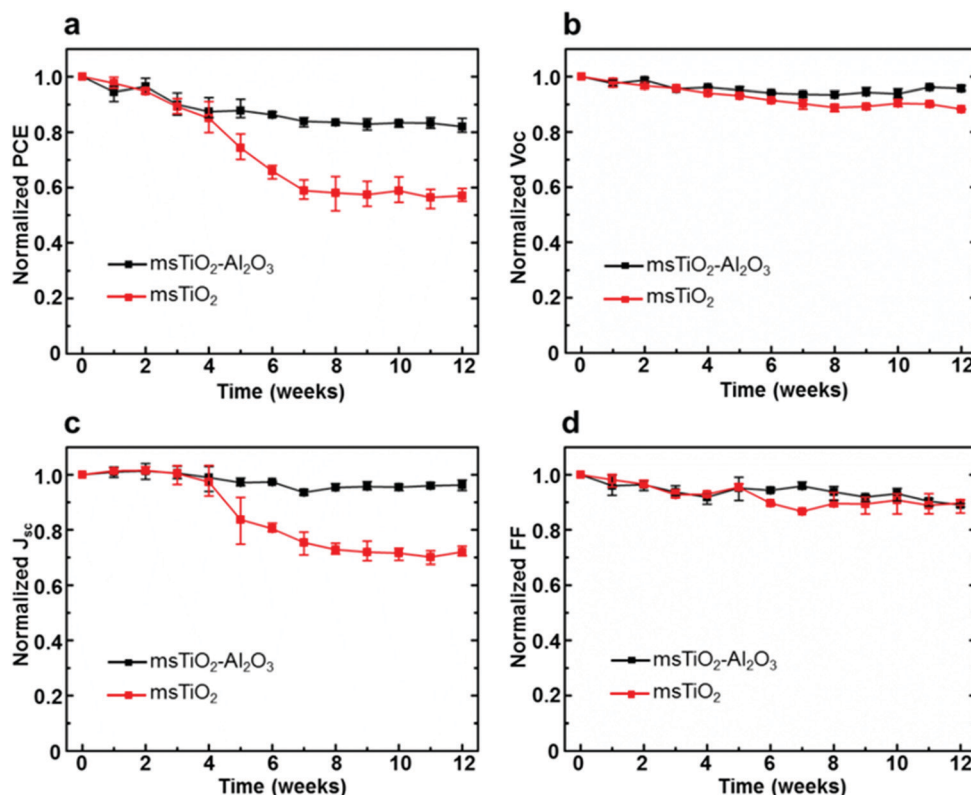


Fig. 2 Evolution of photovoltaic parameters: (a) PCE, (b) V_{oc} , (c) J_{sc} , and (d) FF of the PSCs based on different mesoporous scaffolds stored in ambient air at >65% relative humidity. The measurement was carried out immediately after fabrication and once a week in the following 12 weeks. Each value is the average over 8 cells in the same counterparts.



scaffold-based PSCs, as shown in Fig. S5 (ESI[†]), where moisture instability even more significant than that for the single layer TiO₂ is observed. More detailed information regarding the PCE decay of different types of PSCs can be found in Table S2 (ESI[†]).

The measurement of the decay of photovoltaic parameters in PSCs reveals that the additional mesoporous Al₂O₃ layer allows *J*_{sc} values in devices to be maintained over time, and therefore, significantly improves the long-time air stability of solar cells. It is also noteworthy that we have improved the air stability of the PSCs without any sacrifice of their initial PCE using a bilayer mesoporous scaffold, rather than a thickened single mesoporous scaffold. As a comparison, we fabricated PSCs based on the single mesoporous scaffold with various thicknesses, ranging from 100 nm to 800 nm for TiO₂, and 150 nm to 700 nm for Al₂O₃. As shown in Fig. S6 and S7 (ESI[†]), all devices based on thick single mesoporous scaffolds (>150 nm) suffer from a reduced photovoltaic performance, which is consistent with previous studies.^{22,23} The reduced photovoltaic performance, particularly low *V*_{oc} and FF, of PSCs based on the mesoporous TiO₂ with thickness over 150 nm is mainly caused by the long electron transport distance in the TiO₂ lattices and possible exposure of TiO₂ nanoparticles to the hole transporter. In the case of PSCs based on the single mesoporous Al₂O₃, although electrons can move rapidly in the perovskite itself, a lack of sufficient charge extraction leads to poor photovoltaic performance.

In summary, PSCs based on the TiO₂/Al₂O₃ bilayer mesoporous scaffold yield comparable PCE and improved air stability compared to the conventional PSCs based on the single mesoporous TiO₂ substrate. In order to investigate the impact of bilayer mesoscopic architecture on PSC device stability, the evolution of the phase, composition, and morphology of the perovskite layers generated on different mesoporous scaffolds was characterized and analyzed as a function of time.

XRD measurements of the perovskite films deposited on different mesoporous scaffolds were carried out immediately after fabrication and once a week in the following three weeks. As can be seen from Fig. 3a and b, the perovskite films are composed of a combined phase of perovskite and PbI₂. The intensity of the main peak (14.2°), corresponding to (110) planes of the perovskite, decreases over time, while the peak at 12.7° corresponding to PbI₂(001) planes has an increasing intensity. This can be observed for both bilayer and single layer mesoporous scaffold-based perovskite films, indicating the occurrence of perovskite degradation to form PbI₂ over time for both perovskite films. However, the ratio of peak intensities of (001) PbI₂ to (110) perovskite based on the bilayer mesoporous scaffold grows significantly slower than that of the perovskite film based on the single mesoporous scaffold. This agrees with the observation from the colour change of the perovskite films: the aged perovskite film based on the single

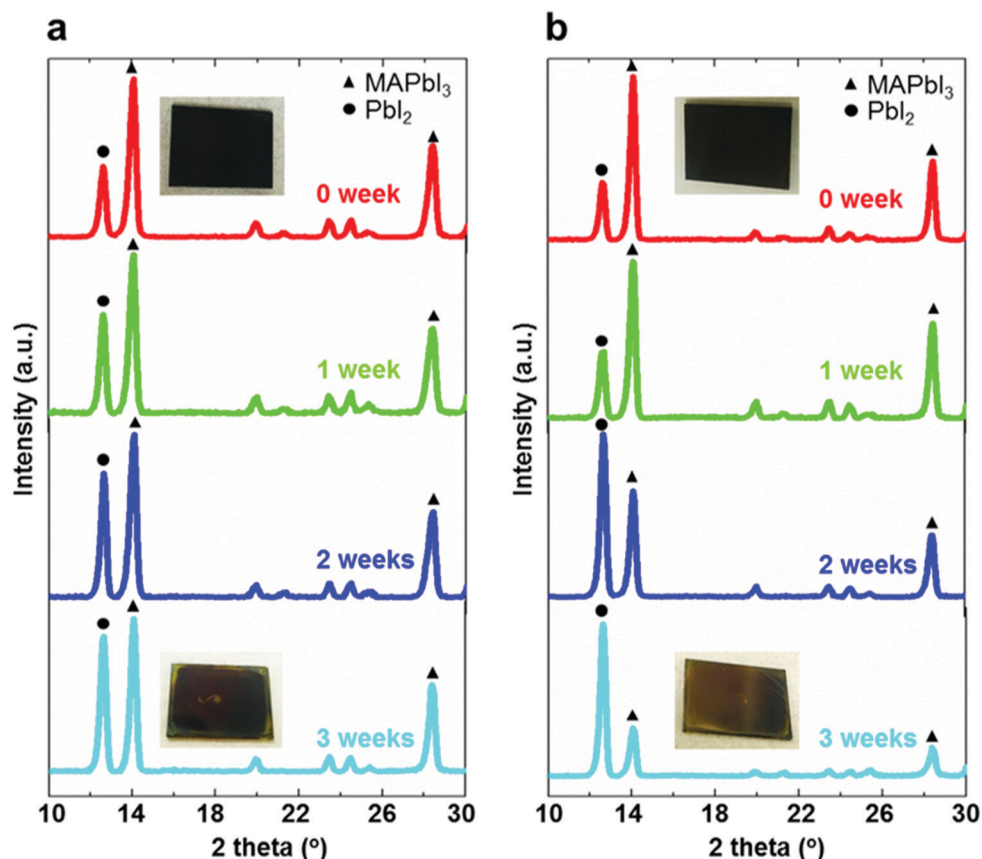


Fig. 3 XRD spectra in the 2θ range from 10° to 30° of the perovskite films based on (a) the bilayer mesoporous scaffold and (b) the single mesoporous scaffold after different aging times in ambient air at >65% relative humidity. The insets are the top views of the corresponding perovskite films.



TiO₂ mesoporous scaffold exhibits a more yellowish appearance, indicating the existence of more PbI₂,³⁸ than the perovskite film based on the TiO₂/Al₂O₃ bilayer mesoporous scaffold. It is therefore apparent that the TiO₂/Al₂O₃ bilayer mesoporous scaffold may reduce the degradation rate of the perovskite films during the aging test in ambient air.

To further understand the underlying mechanisms, the morphological changes of the perovskite films were investigated *via* SEM. Fig. 4 shows the top and cross-sectional views of the perovskite films deposited on different mesoporous scaffolds as a function of time. It can be observed that a capping layer covers the top of the perovskite-infiltrated mesoporous scaffold. For freshly prepared perovskite films, the capping layer based on the single mesoporous scaffold is clearly much thicker than that based on the bilayer mesoporous scaffold. However, the perovskite films based on the different scaffolds have a similar overall thickness of ~ 400 nm. After three weeks' aging in ambient air, the fresh perovskite capping layer which is composed of uniform grains deforms into a stack of irregularly distributed large flakes under exposure to humid air, indicating degradation of perovskite crystals. The perovskite which penetrates the mesoporous scaffold, however, is protected from degradation by the mesoporous scaffold and maintains its initial morphology during the aging test. This is consistent with our result from the XRD measurement that the perovskite films based on the TiO₂/Al₂O₃ bilayer mesoporous scaffold show less degradation than that of the single mesoporous TiO₂ counterparts, due to a thicker mesoporous scaffold and a thinner perovskite capping layer.

In order to study the influence of the mesoporous scaffold on charge transportation and recombination properties of the aged PSCs, photo-current and photo-voltage decay features of the aged PSCs based on different mesoporous scaffolds were studied by measuring the transient electrical response of solar cells excited using a pulsed laser. The photo-current decay is shown in Fig. 5a. The PSC based on the bilayer mesoporous scaffold exhibits an obvious fast photo-current response, whose

time constant $\tau_{\text{tr-bilayer}}$ is calculated to be 2.64 μs , while a longer photo-current decay time of $\tau_{\text{tr-single}} = 16.5 \mu\text{s}$ is obtained from the single mesoporous layer counterpart. Since the photo-current decay feature of photovoltaic devices is highly dependent on the charge transfer rate, a prolonged photo-current decay time for the single mesoporous scaffold-based devices indicates impeded charge extraction at the ETL/perovskite/HTL interfaces. As aforementioned, the perovskite capping layer degrades rapidly in ambient air. In the case of single mesoporous scaffold-based PSCs, a large amount of PbI₂ forms at the perovskite/HTL interface, producing a high density of defect states that block hole transport, thus resulting in a poorer photovoltaic performance after aging. For the photo-voltage measurements (Fig. 5b), a much faster decay ($\tau_{\text{re-single}} = 0.456$ ms) in the single mesoporous scaffold configuration is observed. In contrast, the photo-voltage of the bilayer mesoporous scaffold-based solar cell exhibits a slow decay process with $\tau_{\text{re-double}} = 299$ ms. As the transient photo-voltage measurement was carried out under open-circuit conditions, photo-generated electrons and holes recombine with each other, resulting in decay of the built-in voltage. Therefore, faster photo-voltage decay found in the aged PSCs based on the single mesoporous scaffold suggests a higher recombination rate, which can be understood as a shorter lifetime of the free charge carriers in the perovskite film.³⁹ The electron lifetimes (τ_n) are obtained by differentiating the photovoltage decay curve according to⁴⁰

$$\tau_n = \frac{k_B T}{q} \left(\frac{dV_{\text{oc}}}{dt} \right)^{-1}$$

where k_B is the Boltzmann constant, T is the absolute temperature, and q is the positive elementary charge. Fig. 5c plots the electron lifetimes as a function of photovoltage for the aged PSCs based on different mesoporous scaffolds. Shorter electron lifetimes at equal potentials are observable in the aged PSC based on the single TiO₂ mesoporous scaffold compared to the bilayer counterpart. The high



Fig. 4 SEM images of the perovskite film generated on the bilayer mesoporous scaffold after aging for (a) 0 week, (b) 1 week, (c) 2 weeks and (d) 3 weeks; SEM images of the perovskite film generated on the single mesoporous scaffold after aging for (e) 0 week, (f) 1 week, (g) 2 weeks and (h) 3 weeks. All aging tests were carried out in ambient air at $>65\%$ relative humidity. The insets are the top views of the corresponding perovskite films.





Fig. 5 Transient (a) photo-current and (b) photo-voltage decay of the PSCs based on different mesoporous scaffolds after storing for 12 weeks in ambient air at > 65% relative humidity. (c) Electron lifetimes for the above PSC devices. (d)–(f) J – V curves under the dark condition of different PSCs after storing for 0 week, 6 weeks, and 12 weeks, respectively.

recombination rate for the aged PSC devices based on the single TiO₂ mesoporous scaffold can be ascribed to the rapid degradation of the perovskite capping layer. In contrast, the Al₂O₃ layer in the bilayer mesoporous scaffold acts as a spacing interlayer between ETL and HTL, potentially suppressing short-circuit taking place at the interface, regardless of the absence of the perovskite capping layer. To evaluate this function of the TiO₂/Al₂O₃ bilayer scaffold, the dark currents of different PSC devices were measured during the aging test, as shown in Fig. 5d–f. The devices present a diode-like behaviour.⁴¹ Reduced current leakage in the interfacial charge transfer process is found in the aged PSCs based on the TiO₂/Al₂O₃ bilayer mesoporous scaffold relative to the single TiO₂ counterparts (aged), suggesting a reduced recombination rate at the interface and stronger heterojunctions.³⁹

Different J – V characteristics dependent on the voltage sweep direction indicate the hysteretic behavior of the PSCs,⁴² which

can be associated with the degradation process and interfacial contacts of PSC devices.^{42–45} To find the role of the additional Al₂O₃ mesoporous layer in the hysteretic behavior in J – V characteristics of the PSCs during aging, we performed J – V scans by reverse and forward sweep directions of the PSCs based on different mesoporous scaffolds at different stages of aging, as shown in Fig. 6. The hysteresis index (HI) for PCE or for V_{oc} is defined⁴⁶ as $(PCE(V_{oc})_R - PCE(V_{oc})_F)/PCE(V_{oc})_R$, where R or F represents reverse or forward scan, respectively. HIs for different devices are given in Fig. S8 and Table S3 (ESI[†]) as a function of aging time. The initial HIs for both bilayer mesoporous scaffold- and single mesoporous scaffold-based freshly prepared devices remain at a relatively low value. After aging for 12 weeks, the devices based on the single mesoporous scaffold show a more remarkable hysteresis behavior compared to the bilayer counterparts, with an HI_{PCE} of 0.68. On the other hand,

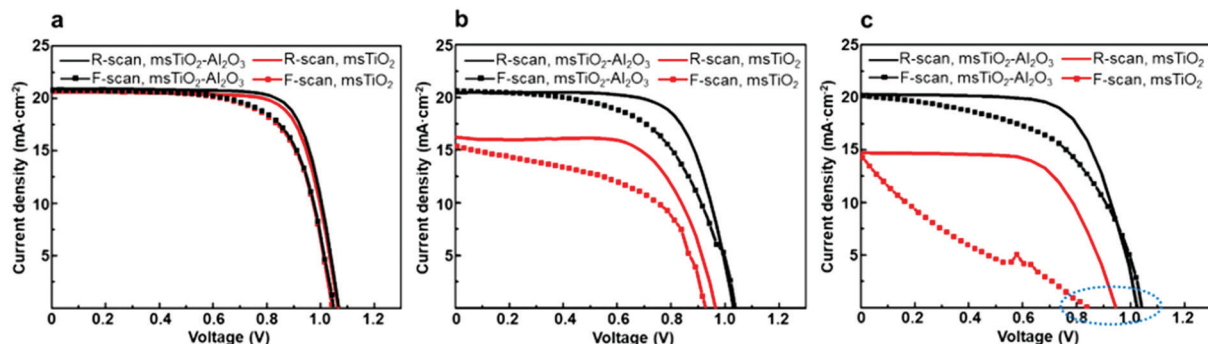


Fig. 6 J – V characteristics of different PSCs showing hysteresis after aging for (a) 0 week, (b) 6 weeks, and (c) 12 weeks. All aging tests were carried out in ambient air at > 65% relative humidity. The hysteresis for photovoltage of the aged devices is highlighted by the blue dotted line circle.

PSCs based on the $\text{TiO}_2/\text{Al}_2\text{O}_3$ bilayer mesoporous scaffold show an HI_{PCE} of only 0.17. Obvious hysteresis for PCE and for photovoltage is observed in the aged single TiO_2 mesoporous scaffold-based devices, which is possibly triggered by pre-biasing conditions in the presence of ion migration and interfacial charge trapping.⁴⁷ Decomposition of the perovskite into PbI_2 and other products under exposure to moisture could induce high concentration of ionic defects.^{45,48} Consequently, increased density of mobile ions or vacancies existing in the degraded perovskite amplifies the detrimental effect of ion migration toward the electrodes. Furthermore, the presence of interfacial charge trap sites mediated by the degradation-induced ionic defects worsens the hysteresis effect in the aged devices.^{42,46,47} In addition, interfacial recombination mediated by TiO_2 may also contribute to the notable hysteretic behavior observed in the aged devices based on the single TiO_2 mesoporous scaffold. In a typical configuration of an n-i-p mesoscopic PSC, a perovskite capping layer is sandwiched between mesoporous scaffold and HTL. As the perovskite capping layer degrades, the mesoporous scaffold is gradually exposed to the perovskite/HTL interface. In the case of single TiO_2 devices (ESI,† Fig. S9), the appearance of TiO_2 nanoparticles in the p-type contact area increases the possibility of interfacial charge recombination between free holes and the trapped electrons in the defect states in TiO_2 .³³ In turn, less severe hysteresis observed in the

aged $\text{TiO}_2/\text{Al}_2\text{O}_3$ bilayer mesoporous scaffold-based devices compared to that of the aged single TiO_2 devices could indicate, to some extent, reduced formation of ionic defects and suppressed interfacial recombination by adding a mesoporous Al_2O_3 layer on top of mesoporous TiO_2 . Preventing TiO_2 -assisted interfacial recombination by building an interlayer of insulator, such as Al_2O_3 or ZrO_2 , between TiO_2 and HTL (or electrode) has also been reported in other studies.^{14,20,49,50}

The moisture-induced mechanism for PSCs based on different mesoporous scaffolds is illustrated in Fig. 7. Since water vapor in humidified air causes decomposition of the capping layer first, the “unprotected” perovskite capping layer degrades rapidly under exposure to ambient air. Defects such as cracks or voids existing at perovskite grain boundaries allow water molecules to penetrate, accelerating the perovskite degradation. Finally, the uniform perovskite grains are completely transformed into disorganized flake-like grains covered by PbI_2 . After this fast decomposition of the perovskite capping layer, water penetration slows down due to the presence of the mesoscopic layer, as shown in Fig. 7a. The mesoscopic framework substantially promotes the nucleation and growth of the infiltrated perovskite crystals and thus the pore-filling factor.⁵¹ Therefore, the perovskite layer within the mesoscopic framework remains in a dense morphology for a longer period of time than the capping layer, and acts as a barrier against water ingress during the



Fig. 7 Schematics of the moisture-induced degradation processes of the perovskite films generated on (a) the bilayer mesoporous scaffold and (b) the single mesoporous scaffold. Routes 1, 2, and 3 represent hole transport from the perovskite to the HTL, recombination at the ETL/perovskite, and recombination at the ETL/HTL, respectively.



Fig. 8 Absorbance and transmittance of the perovskite films generated on (a) the bilayer mesoporous scaffold and (b) the single mesoporous scaffold before and after aging for 3 weeks in ambient air at >65% relative humidity.

aging test. In addition, the water-absorbing mesoscopic surface can retain water²⁰ and prevent further water penetration into the deep sites. Due to the slower degradation rate occurring in the perovskite within the mesoporous scaffold, the bilayer mesoporous scaffold with increased thickness relative to the single mesoporous scaffold can load more perovskite, ensuring a more stable photovoltaic performance of the PSCs. On the other hand, as shown in Fig. 7b, a ~250 nm-thick perovskite capping layer generated on the single mesoporous scaffold degrades into a large amount of PbI_2 accumulating at the external surface, which may seriously hinder the transport of the holes to the HTL (route 1). While the TiO_2 nanoparticles are exposed to the perovskite/HTL interface, the accumulated holes at the interface or the free holes in the HTL may recombine with the trapped electrons in the electronic trap states of TiO_2 (route 2 or 3).³³ In the case of a bilayer mesoporous scaffold, the charge transfer rate is less likely to suffer a significant reduction because of the reduced PbI_2 formation at the interface. In addition, the existence of a layer of Al_2O_3 insulator effectively prevents direct contact at the interface between ETL and HTL (or charge carrier accumulations), particularly when the perovskite capping layer has failed to do so due to its severe degradation. Therefore, for the aged PSCs, devices based on the $\text{TiO}_2/\text{Al}_2\text{O}_3$ bilayer mesoporous scaffold show a reduced possibility of charge recombination and current leakage at the interface, as well as less severe hysteric behavior, in comparison to the PSCs based on the single TiO_2 mesoporous scaffold.

Besides long-term air stability, transmittance is another important indicator to evaluate perovskite solar cells.⁵² One of the advantages of PSCs is their semi-transparent feature over the visible light range (380–740 nm), which enables the possibility of fabricating tandem solar cells to utilize solar energy more fully. UV-visible spectroscopy of perovskite films generated on the bilayer mesoporous scaffold (Fig. 8a) and on the single mesoporous scaffold (Fig. 8b) was performed under incident light wavelengths ranging from 300 nm to 1100 nm. The average visible transmittance (AVT) of the perovskite film generated on the bilayer mesoporous scaffold slightly decreases

from 9.48% to 7.97% after aging, still preserving the semi-transparent feature. In contrast, a distinct decrease of the AVT from 9.52% to 4.64% is found in the aged perovskite film based on the single mesoporous scaffold. As Fig. 4 demonstrates, the perovskite film based on the single mesoporous scaffold suffers severe degradation. During the degradation process, the smooth perovskite capping layer composed of neatly arranged large grains is transformed into a stack of unorganized flakes. The high volume of “new-born” grain boundaries and rough surfaces scatter incident light, leading to poor light transmittance. For the perovskite film based on the bilayer scaffold, the Al_2O_3 layer introduces no reduction in light transmittance due to the similar overall thickness between fresh perovskite films based on different mesoporous scaffolds. During the aging test, the perovskite film based on the bilayer mesoporous scaffold shows a smaller change in morphology as well as a more stable light transmission over time compared to the single mesoporous scaffold.

Conclusions

In this work, we have fabricated a $\text{TiO}_2/\text{Al}_2\text{O}_3$ bilayer mesoporous scaffold for perovskite solar cells with a specific focus on enhancing the air stability of the devices. Based on the $\text{TiO}_2/\text{Al}_2\text{O}_3$ bilayer mesoporous scaffold, stable and efficient perovskite solar cells were produced and stored under ambient conditions at a RH of over 65%, delivering 16.84% initial PCE. Over 80% of the initial PCE was retained after 2000 h. The improved air stability of PSCs is mainly attributed to a reduced degradation rate of the perovskite absorber by introducing an additional mesoporous Al_2O_3 scaffold, which can protect more of the infiltrated perovskite from attack by moisture. More stable perovskite within the mesoscopic layer ensures an adequate yield of photo-generated charge carriers and thereby a more stable photovoltaic performance of the PSCs over time. In addition, the existence of an additional Al_2O_3 layer can mitigate the charge transportation property deterioration due to the perovskite degradation. Furthermore, the perovskite film based on the $\text{TiO}_2/\text{Al}_2\text{O}_3$ bilayer mesoporous



scaffold shows comparable and more stable light transmission than the perovskite film based on a single TiO₂ mesoporous scaffold, indicating its potential to be used in tandem PSC applications. Our research proposes a promising approach using a well-designed TiO₂/Al₂O₃ bilayer mesoscopic architecture towards the fabrication of stable and efficient PSCs.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

Dong Wang acknowledges the financial support by the State Scholarship Fund under the auspices of the China Scholarship Council (No. 201706120028). The authors particularly thank Dr John E. Warren for his help with XRD characterization.

Notes and references

- 1 J. Cao, C. Li, X. Lv, X. Feng, R. Meng, Y. Wu and Y. Tang, *J. Am. Chem. Soc.*, 2018, **140**, 11577–11580.
- 2 N. Aristidou, I. Sanchez-Molina, T. Chotchuanchutaval, M. Brown, L. Martinez, T. Rath and S. A. Haque, *Angew. Chem., Int. Ed.*, 2015, **54**, 8208–8212.
- 3 B. Conings, J. Drijkoningen, N. Gauquelin, A. Babayigit, J. D'Haen, L. D'Olieslaeger, A. Ethirajan, J. Verbeeck, J. Manca, E. Mosconi, F. D. Angelis and H.-G. Boyen, *Adv. Energy Mater.*, 2015, **5**, 1500477.
- 4 J. A. Christians, P. A. Miranda Herrera and P. V. Kamat, *J. Am. Chem. Soc.*, 2015, **137**, 1530–1538.
- 5 J. C.-R. Ke, A. S. Walton, D. J. Lewis, A. Tedstone, P. O'Brien, A. G. Thomas and W. R. Flavell, *Chem. Commun.*, 2017, **53**, 5231–5234.
- 6 J. Zhang, T. Song, Z. Zhang, K. Ding, F. Huang and B. Sun, *J. Mater. Chem. C*, 2015, **3**, 4402–4406.
- 7 L. Wang, C. McCleese, A. Kovalsky, Y. Zhao and C. Burda, *J. Am. Chem. Soc.*, 2014, **136**, 12205–12208.
- 8 Q. Chen, H. Zhou, T.-B. Song, S. Luo, Z. Hong, H.-S. Duan, L. Dou, Y. Liu and Y. Yang, *Nano Lett.*, 2014, **14**, 4158–4163.
- 9 Y. H. Lee, J. Luo, R. Humphry-Baker, P. Gao, M. Grätzel and M. K. Nazeeruddin, *Adv. Funct. Mater.*, 2015, **25**, 3925–3933.
- 10 H. Xu, Y. Wu, J. Cui, C. Ni, F. Xu, J. Cai, F. Hong, Z. Fang, W. Wang, J. Zhu, L. Wang, R. Xu and F. Xu, *Phys. Chem. Chem. Phys.*, 2016, **18**, 18607–18613.
- 11 W. Veiga and C. M. Lepienski, *Mater. Sci. Eng., A*, 2002, **335**, 6–13.
- 12 M. Wang, X. Jiang, J. Bian, Y. Feng, C. Wang, Y. Huang, Y. Zhang and Y. Shi, *ACS Appl. Mater. Interfaces*, 2019, **11**, 2989–2996.
- 13 M. Z. Mokhtar, Q. Chen, Q. Lian, D. J. Lewis, B. R. Saunders, A. S. Walton, C.-R. Ke, E. Whittaker, B. Hamilton and S. Haque, *ACS Appl. Energy Mater.*, 2018, **1**, 5567–5578.
- 14 S. Guarnera, A. Abate, W. Zhang, J. M. Foster, G. Richardson, A. Petrozza and H. J. Snaith, *J. Phys. Chem. Lett.*, 2015, **6**, 432–437.
- 15 M. M. Tavakoli, P. Yadav, R. Tavakoli and J. Kong, *Adv. Energy Mater.*, 2018, **8**, 1800794.
- 16 E. Y. Choi, J. Kim, S. Lim, E. Han, A. W. Y. Ho-Baillie and N. Park, *Sol. Energy Mater. Sol. Cells*, 2018, **188**, 37–45.
- 17 Y. I. Lee, N. J. Jeon, B. J. Kim, H. Shim, T.-Y. Yang, S. I. Seok, J. Seo and S. G. Im, *Adv. Energy Mater.*, 2018, **8**, 1701928.
- 18 Y. Wang, W.-Y. Rho, H.-Y. Yang, T. Mahmoudi, S. Seo, D.-H. Lee and Y.-B. Hahn, *Nano Energy*, 2016, **27**, 535–544.
- 19 Y. Wang, T. Mahmoudi, W.-Y. Rho and Y.-B. Hahn, *Nano Energy*, 2019, **64**, 103964.
- 20 A. Mei, X. Li, L. Liu, Z. Ku, T. Liu, Y. Rong, M. Xu, M. Hu, J. Chen, Y. Yang, M. Grätzel and H. Han, *Science*, 2014, **345**, 295–298.
- 21 K. Cao, Z. Zuo, J. Cui, Y. Shen, T. Moehl, S. M. Zakeeruddin, M. Grätzel and M. Wang, *Nano Energy*, 2015, **17**, 171–179.
- 22 J. Yang, K. M. Fransishyn and T. L. Kelly, *Chem. Mater.*, 2016, **28**, 7344–7352.
- 23 J. Yin, J. Cao, X. He, S. Yuan, S. Sun, J. Li, N. Zheng and L. Lin, *J. Mater. Chem. A*, 2015, **3**, 16860–16866.
- 24 A. Fakharuddin, F. Di Giacomo, I. Ahmed, Q. Wali, T. M. Brown and R. Jose, *J. Power Sources*, 2015, **283**, 61–67.
- 25 M. H. Kumar, N. Yantara, S. Dharani, M. Grätzel, S. Mhaisalkar, P. P. Boix and N. Mathews, *Chem. Commun.*, 2013, **49**, 11089–11091.
- 26 W. Geng, C.-J. Tong, J. Liu, W. Zhu, W.-M. Lau and L.-M. Liu, *Sci. Rep.*, 2016, **6**, 20131.
- 27 P. Tiwana, P. Docampo, M. B. Johnston, H. J. Snaith and L. M. Herz, *ACS Nano*, 2011, **5**, 5158–5166.
- 28 H.-S. Kim, I.-H. Jang, N. Ahn, M. Choi, A. Guerrero, J. Bisquert and N.-G. Park, *J. Phys. Chem. Lett.*, 2015, **6**, 4633–4639.
- 29 Y. Zhang, M. Liu, G. E. Eperon, T. C. Leijtens, D. McMeekin, M. Saliba, W. Zhang, M. de Bastiani, A. Petrozza, L. M. Herz, M. B. Johnston, H. Lin and H. J. Snaith, *Mater. Horiz.*, 2015, **2**, 315–322.
- 30 F. T. F. O'Mahony, Y. H. Lee, C. Jelliet, S. Dmitrov, D. T. J. Bryant, J. R. Durrant, B. C. O'Regan, M. Grätzel, M. K. Nazeeruddin and S. A. Haque, *J. Mater. Chem. A*, 2015, **3**, 7219–7223.
- 31 H.-S. Kim, C.-R. Lee, J.-H. Im, K.-B. Lee, T. Moehl, A. Marchioro, S.-J. Moon, R. Humphry-Baker, J.-H. Yum, J. E. Moser, M. Grätzel and N.-G. Park, *Sci. Rep.*, 2012, **2**, 591.
- 32 G. S. Han, H. S. Chung, D. H. Kim, B. J. Kim, J.-W. Lee, N.-G. Park, I. S. Cho, J.-K. Lee, S. Lee and H. S. Jung, *Nanoscale*, 2015, **7**, 15284–15290.
- 33 T. Leijtens, G. E. Eperon, S. Pathak, A. Abate, M. M. Lee and H. J. Snaith, *Nat. Commun.*, 2013, **4**, 2885.
- 34 M. M. Lee, J. Teuscher, T. Miyasaka, T. N. Murakami and H. J. Snaith, *Science*, 2012, **338**, 643–647.
- 35 Q. Wang, S. Liu, Y. Ming, Y. Guan, D. Li, C. Zhang, Z. Wang, Y. Rong, Y. Hu and H. Han, *Sustainable Energy Fuels*, 2018, **2**, 2412–2418.
- 36 Q. Chen, M. Z. Mokhtar, J. C.-R. Ke, A. G. Thomas, A. Hadi, E. Whittaker, M. Curioni and Z. Liu, *Sustainable Energy Fuels*, 2018, **2**, 1216–1224.



- 37 J. Troughton, K. Hooper and T. M. Watson, *Nano Energy*, 2017, **39**, 60–68.
- 38 G. Niu, W. Li, F. Meng, L. Wang, H. Dong and Y. Qiu, *J. Mater. Chem. A*, 2014, **2**, 705–710.
- 39 Z. Liu, B. Sun, T. Shi, Z. Tang and G. Liao, *J. Mater. Chem. A*, 2016, **4**, 10700–10709.
- 40 X. Gao, J. Li, J. Baker, Y. Hou, D. Guan, J. Chen and C. Yuan, *Chem. Commun.*, 2014, **50**, 6368–6371.
- 41 J. Cui, F. Meng, H. Zhang, K. Cao, H. Yuan, Y. Cheng, F. Huang and M. Wang, *ACS Appl. Mater. Interfaces*, 2014, **6**, 22862–22870.
- 42 R. T. Ginting, M.-K. Jeon, K.-J. Lee, W.-Y. Jin, T.-W. Kim and J.-W. Kang, *J. Mater. Chem. A*, 2017, **5**, 4527–4534.
- 43 P. Calado, A. M. Telford, D. Bryant, X. Li, J. Nelson, B. C. O'Regan and P. R. F. Barnes, *Nat. Commun.*, 2016, **7**, 13831.
- 44 Y. Zhang, M. Liu, G. E. Eperon, T. C. Leijtens, D. McMeekin, M. Saliba, W. Zhang, M. de Bastiani, A. Petrozza, L. M. Herz, M. B. Johnston, H. Lin and H. J. Snaith, *Mater. Horiz.*, 2015, **2**, 315–322.
- 45 W. Tress, N. Marinova, T. Moehl, S. M. Zakeeruddin, M. K. Nazeeruddin and M. Grätzel, *Energy Environ. Sci.*, 2015, **8**, 995–1004.
- 46 J. Hong, H. Kim and I. Hwang, *J. Mater. Chem. C*, 2019, **7**, 1173–1181.
- 47 S. van Reenen, M. Kemerink and H. J. Snaith, *J. Phys. Chem. Lett.*, 2015, **6**, 3808–3814.
- 48 E. C. Smith, C. L. C. Ellis, H. Javai, L. A. Renna, Y. Liu, T. P. Russel, M. Bag and D. Venkataraman, *J. Phys. Chem. C*, 2018, **122**, 13986–13994.
- 49 M. Xu, Y. Rong, Z. Ku, A. Mei, T. Liu, L. Zhang, X. Li and H. Han, *J. Mater. Chem. A*, 2014, **2**, 8607–8611.
- 50 Y. Wang, T. Mahmoudi, H.-Y. Yang, K. S. Bhat, J.-Y. Yoo and Y.-B. Hahn, *Nano Energy*, 2018, **49**, 59–66.
- 51 N. Islam, M. Yang, K. Zhu and Z. Fan, *J. Mater. Chem. A*, 2015, **3**, 24315–24321.
- 52 C. Dokkhan, M. Z. Mokhtar, C.-R. Ke, A. S. Walton, Q. Chen, N. W. Hodson, Q. Lian and B. R. Saunders, *ACS Appl. Energy Mater.*, 2019, **2**, 6624–6633.

