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Coordination-induced reversible electrical conductivity variation in the MOF-74 analogue Fe₂(DSBDC)[†]

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Inner-sphere changes at the open Fe centers in Fe₂(DSBDC) (DSBDC⁴⁻ = 2,5-disulfidobenzene-1,4-dicarboxylate), as caused by coordination and release of solvent molecules, lead to reversible structural and electrical conductivity changes. Specifically, coordination of *N,N*-dimethylformamide (DMF) to the open Fe sites improves the room-temperature electrical conductivity by three orders of magnitude. Supported by additional density functional theory calculations, we attribute the electrical conductivity enhancement to partial electron transfer from Fe to DMF, which generates hole carriers and improves the charge carrier density in Fe₂(DSBDC).

Introduction

Traditionally raising interest due to their highly porous nature,¹ metal-organic frameworks (MOFs) offer an excellent platform for investigating and tuning electrical, magnetic, and optical processes in organic-based materials, some of which have already led to applications in lithium-ion batteries,²⁻⁴ thermoelectrics,^{5,6} and supercapacitors,⁷ among others. The electronic properties may be modulated either inherently, through changes in chemical composition, or through host-guest interactions. Guest molecules may cause structural variation,^{8,9} tune electrical conductivity,^{10,11} induce spin state transitions,^{12,13} modulate luminescence wavelength,^{14,15} and even participate in guest-guest interactions.¹⁶ These phenomena are of fundamental interest and form the basis of various applications including chemiresistive sensors,^{11,17,18} fluorescence sensors,¹⁹ and cooperative adsorption.^{13,16}

Although several strategies are effective for increasing the electrical conductivity in MOFs,^{20,21} among the most promising, and certainly more tractable synthetically, are those that involve post-synthetic doping. Guest molecules introduced after the formation of a MOF may tune either the charge mobility or the charge density in the skeleton of a given material, allowing in either case continuous enhancement of electrical conductivity over several orders of magnitude. Chemically, the dopant molecules may engage with the MOF either through outer-sphere electron transfer (*i.e.* redox reactivity)²²⁻²⁵ or through inner-sphere reactivity. The latter involves binding of guest molecules to coordinatively unsaturated metal centers to form charge transport pathways and/or to inject charge carriers.¹⁰ The inner sphere mechanism is proposed to be operative, for instance, in TCNQ-infiltrated Cu₃(benzene-1,3,5-tricarboxylate)₂ (TCNQ = 7,7,8,8-tetracyanoquinodimethane), wherein TCNQ molecules coordinate to Cu sites and generate increasingly efficient charge percolation pathways and million-fold increase in the electrical conductivity of the parent MOF.¹⁰

Herein, we show that simple coordination of solvent molecules such as DMF modulates the electrical conductivity of Fe₂(DSBDC)²⁶ by three orders of magnitude. Fe₂(DSBDC) is a structural analogue of the well-known MOF-74 series M₂(2,5-dihydroxybenzene-1,4-dicarboxylate) (M₂(DOBDC), M = Mg, Mn, Fe, Co, Ni, Cu, Zn)²⁷⁻³⁴ where phenoxide groups on the DOBDC ligand have been replaced by thiophenoxide. Each Fe atom in Fe₂(DSBDC) is coordinated by two *trans* thiophenoxide groups, three meridionally coordinated carboxylate groups, and one DMF molecule. The latter can be removed to yield coordinatively unsaturated Fe sites. The Fe and thiophenoxide S atoms form (-Fe-S-)_∞ chains, which are bridged by DSBDC⁴⁻ ligands to form a three-dimensional framework containing one-dimensional hexagonal pores (Fig. 1). Previously, we reported that the regular hexagonal pores of DMF-filled Fe₂(DSBDC) (Fe₂(DSBDC)(DMF)₂·x(DMF)) distort significantly when the unbound guest DMF molecules are replaced by dichloromethane (DCM) and the material is activated 80 °C to give Fe₂(DSBDC)(DMF)₂.²⁶ The distortion is reversible: soaking

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[†]Electronic supplementary information (ESI) available: Experimental details, computation details, thermogravimetric analysis profiles, powder X-ray diffraction patterns, BET surface area analysis. See DOI: 10.1039/c8dt02197j



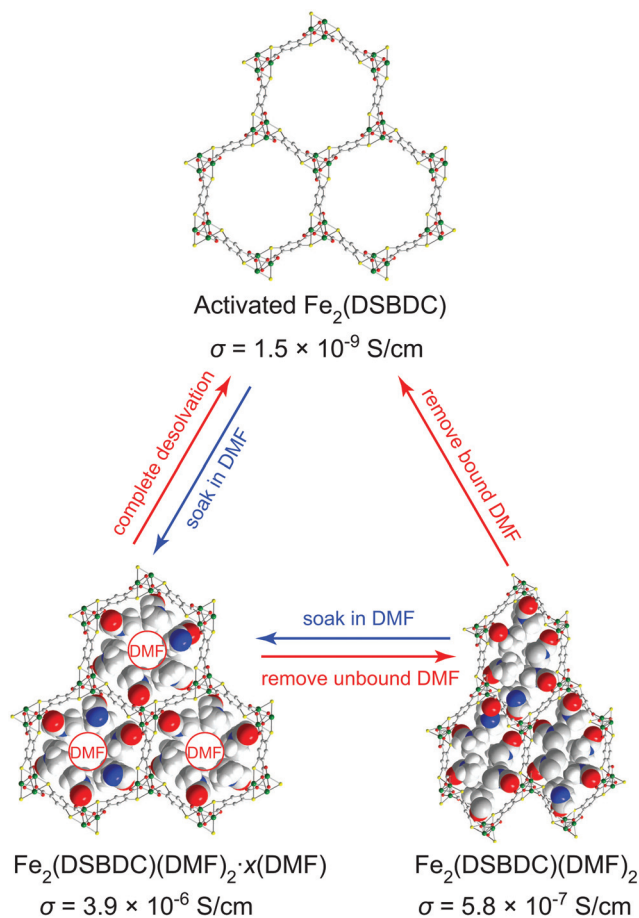


Fig. 1 Structural changes induced in $\text{Fe}_2(\text{DSBDC})$ by coordination and release of DMF, with respective electrical conductivity (σ) values. The structures of $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2 \cdot x(\text{DMF})$ and $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2$ were determined by single-crystal XRD and supported by DFT simulations, respectively.²⁶ The structure of the activated $\text{Fe}_2(\text{DSBDC})$ is proposed based on its PXRD pattern. Green, yellow, red, blue, grey, and white spheres represent Fe, S, O, N, C, and H atoms, respectively. Unbound DMF molecules and framework H atoms are omitted for clarity.

$\text{Fe}_2(\text{DSBDC})(\text{DMF})_2$ in DMF regenerates the regular hexagonal pores and original structure. We extend the structural and electrical properties study of this system here by investigating the fully desolvated material, $\text{Fe}_2(\text{DSBDC})$, which shows divergent behaviour from its solvated congeners.

Results and discussion

$\text{Fe}_2(\text{DSBDC})(\text{DMF})_2 \cdot x(\text{DMF})$ was synthesized according to a previously reported procedure.²⁶ Owing to the sensitivity of this material to both water and air, all subsequent manipulations were performed in a dry, oxygen-free atmosphere. Thermogravimetric analysis (TGA) of this material revealed two consecutive weight loss events at 210–335 °C and 320–560 °C, respectively (Fig. S1†). The first weight loss step can be correlated to the loss of both bound and unbound DMF molecules, while the second likely corresponds to framework

decomposition. Indeed, powder X-ray diffraction (PXRD) confirmed that heating $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2 \cdot x(\text{DMF})$ at 200 °C under vacuum for 2 days causes amorphization (Fig. S2†). In an attempt to activate the material without decomposition, we exchanged DMF by soaking the as-synthesized MOF in lower-boiling methanol (MeOH), to produce $\text{Fe}_2(\text{DSBDC})(\text{MeOH})_2 \cdot x(\text{MeOH})$. Near-complete displacement of DMF by MeOH was confirmed by infrared (IR) spectroscopy, which revealed significantly decreased intensity of the C=O stretch at 1654 cm^{-1} and appearance of a MeOH C–O stretching band at 1008 cm^{-1} as well as a broad O–H stretching band at ~3300 cm^{-1} (Fig. 2). Although $\text{Fe}_2(\text{DSBDC})(\text{MeOH})_2 \cdot x(\text{MeOH})$ remains crystalline when kept in liquid MeOH (Fig. 3), it also decomposes upon desolvation (Fig. S2†), possibly because of a stronger interaction between MeOH and the framework. We therefore sought to replace methanol with a weaker interacting solvent and soaked the methanol-exchanged material in tetrahydrofuran (THF) to produce $\text{Fe}_2(\text{DSBDC})(\text{THF})_2 \cdot x(\text{THF})$. IR spectroscopy once again confirmed near-quantitative exchange, with complete disappearance of the bands at 1008 and 3300 cm^{-1} and rise of characteristic C–H stretching bands for THF at 2969 and 2866 cm^{-1} , which are blue-shifted from the C–H stretching bands for MeOH at 2938 and 2831 cm^{-1} . Note that direct replacement of DMF by THF is not successful, presumably because THF is too poorly coordinating and cannot displace Fe-bound DMF, thus requiring the intermediacy of MeOH.

TGA of $\text{Fe}_2(\text{DSBDC})(\text{THF})_2 \cdot x(\text{THF})$ confirmed that THF can be removed even below 100 °C, much lower than the decomposition temperature of approximately 270 °C (Fig. S3†). Therefore, to remove all solvent and produce completely activated MOF, $\text{Fe}_2(\text{DSBDC})(\text{THF})_2 \cdot x(\text{THF})$ was heated at 170 °C under dynamic vacuum (<3 mTorr) for 36 h. This yielded a dark red powder whose elemental microanalysis revealed a

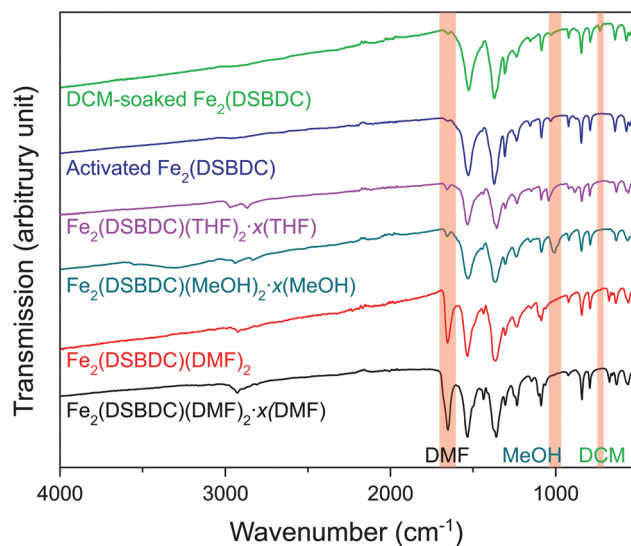


Fig. 2 IR spectra of various forms of $\text{Fe}_2(\text{DSBDC})$. Characteristic peaks of DMF, MeOH, and DCM are highlighted.



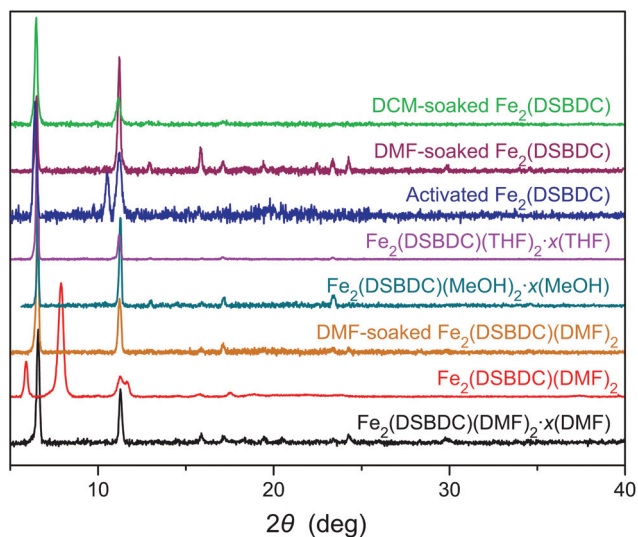


Fig. 3 PXRD patterns of various phases of $\text{Fe}_2(\text{DSBDC})$.

formula that matched that of fully desolvated $\text{Fe}_2(\text{DSBDC})$, whose IR spectrum also indicates the complete absence of DMF, MeOH, or THF characteristic bands (Fig. 2 and ESI†). Most tellingly, an N_2 adsorption of $\text{Fe}_2(\text{DSBDC})$ exhibits a type I isotherm with a saturation uptake of approximately 160 cm^3 of N_2 per g at 77 K (Fig. 4) and an apparent Brunauer–Emmett–Teller (BET) surface area of $624 \text{ m}^2 \text{ g}^{-1}$ ($211 \text{ m}^2 \text{ mmol}^{-1}$) (Fig. S4†). Although smaller than the value observed for related $\text{Mn}_2(\text{DSBDC})$ ($329 \text{ m}^2 \text{ mmol}^{-1}$)³⁵ and MOF-74 analogues ($287\text{--}416 \text{ m}^2 \text{ mmol}^{-1}$),^{30,32,33} possibly due to a further structural distortion (*vide infra*), the apparent surface area of $\text{Fe}_2(\text{DSBDC})$ is significantly larger than that of $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2$ ($83 \text{ m}^2 \text{ g}^{-1}$, $40 \text{ m}^2 \text{ mmol}^{-1}$) (Fig. 4),³⁶ as would be expected upon removal of pore-blocking DMF molecules.

Although $\text{Fe}_2(\text{DSBDC})$ is crystalline, its PXRD pattern diverges from those of both $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2 \cdot x(\text{DMF})$ and

$\text{Fe}_2(\text{DSBDC})(\text{DMF})_2$: peaks at 6.6° and 11.3° corresponding to the $[2\text{--}10]$ and $[300]$ planes in the original regular hexagonal lattice are intact, but a new peak at 10.5° appears (Fig. 3). Qualitatively, this confirms that the hexagonal symmetry and pore shape is retained in $\text{Fe}_2(\text{DSBDC})$, but the lattice suffers a distortion along the crystallographic c direction, which runs parallel to the $(\text{--Fe--S--})_\infty$ chains. Unfortunately, the relatively low crystallinity of the activated material does not offer sufficient quality for further structural refinement. However, soaking $\text{Fe}_2(\text{DSBDC})$ in DMF or DCM produces materials whose PXRD patterns are identical to that of original $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2 \cdot x(\text{DMF})$ (Fig. 3).† This confirms that the distortion associated with solvent exchange and evacuation does not affect the connectivity of the framework. Similarly, any changes in electrical properties stemming from these distortions should also be reversible.

The electrical properties of $\text{Fe}_2(\text{DSBDC})$ were investigated with two-contact-probe devices in air-free conditions at 297 K using pressed pellets made by *in situ* press, an apparatus that was introduced previously.^{37,38} Single-crystal electrical conductivity measurements are unfeasible for all phases of $\text{Fe}_2(\text{DSBDC})$ because of the small crystallite size and extreme sensitivity to air. As shown in Fig. 5, current–voltage curves of all samples are linear, allowing the use of Ohm’s Law for resistance calculations. Notably, these measurements revealed that the electrical conductivity increases with the degree of solvation: it is $3.9 \times 10^{-6} \text{ S cm}^{-1}$ for $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2 \cdot x(\text{DMF})$, it drops to $5.8 \times 10^{-7} \text{ S cm}^{-1}$ for $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2$ upon removal of unbound DMF, and decreases further upon complete removal of DMF to only $1.5 \times 10^{-9} \text{ S cm}^{-1}$ in fully desolvated $\text{Fe}_2(\text{DSBDC})$. In line with the reversible structural distortions discussed above, the electrical properties are also recovered upon soaking $\text{Fe}_2(\text{DSBDC})$ or $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2$ in DMF,

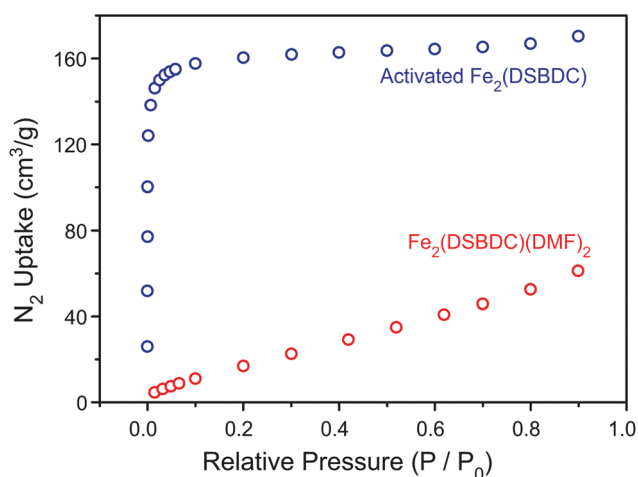


Fig. 4 N_2 sorption isotherms of $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2$ and activated $\text{Fe}_2(\text{DSBDC})$.

- $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2 \cdot x(\text{DMF})$
- $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2$
- DMF-soaked $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2$
- Activated $\text{Fe}_2(\text{DSBDC})$
- DMF-soaked $\text{Fe}_2(\text{DSBDC})$
- DCM-soaked $\text{Fe}_2(\text{DSBDC})$

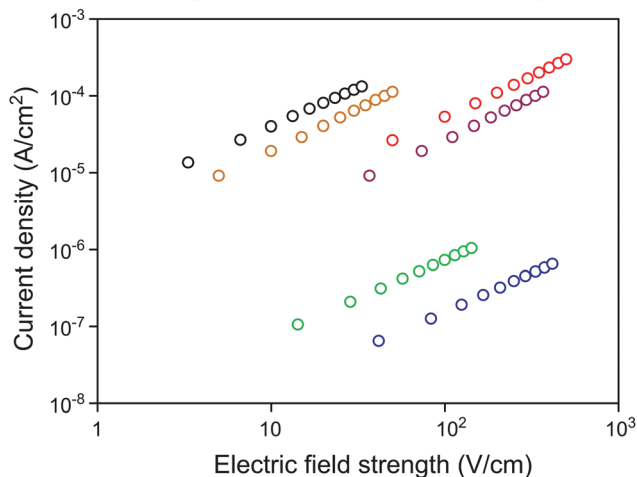


Fig. 5 Plots of current density versus electric field strength of various forms of $\text{Fe}_2(\text{DSBDC})$.



with the recovery values of $3.1 \times 10^{-7} \text{ S cm}^{-1}$ and $2.3 \times 10^{-6} \text{ S cm}^{-1}$, respectively, in good agreement with the conductivity of pristine, as-synthesized $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2 \cdot x(\text{DMF})$ (Scheme S1†).

It should be noted, however, that *coordination* to the Fe center is critical for improved electrical conductivity. Indeed, although soaking fully desolvated $\text{Fe}_2(\text{DSBDC})$ in DCM recovers the skeleton structure observed in $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2 \cdot x\text{DMF}$ (Fig. 3), DCM treatment does not recover electrical conductivity, which remains at $7.3 \times 10^{-9} \text{ S cm}^{-1}$ (Fig. 5), similar to that of activated $\text{Fe}_2(\text{DSBDC})$ itself. In other words, swelling with DCM (or presumably other non-coordinating solvents) recovers the original structure, but does not recover the electrical properties, which are more critically dependent on metal coordination rather than structural deformations.

Density functional theory (DFT) calculations were further used to probe the origin of the coordination-induced electrical conductivity enhancement (Fig. 6). The crystal structure of the activated $\text{Fe}_2(\text{DSBDC})$ was allowed to reach geometric and electronic equilibrium, resulting in regular hexagonal pores that resemble the structure of the DCM-soaked $\text{Fe}_2(\text{DSBDC})$. The projected density of states (PDOS) shows that the valance band of activated $\text{Fe}_2(\text{DSBDC})$ is dominated by Fe and S orbitals, suggesting that perturbation to either Fe or S valence will directly modify the valance band character. We hence conjecture that the major charge transport mechanism in activated $\text{Fe}_2(\text{DSBDC})$ is hole hopping between Fe and S within the $(-\text{Fe}-\text{S}-)_{\infty}$ chains, similar to the situation proposed in $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2$.^{26,36}

Because the skeletons of $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2$ and the DCM-soaked $\text{Fe}_2(\text{DSBDC})$ are isostructural, DMF likely does not affect charge mobility. Instead, it modulates charge density in the skeleton of the framework. Indeed, both a comparison of solid-state and truncated cluster calculations show that the work function of the activated $\text{Fe}_2(\text{DSBDC})$ is smaller than the ionization potential of DMF by 1 eV (Fig. 6). As a result, partial electron transfer occurs when DMF binds to Fe centers in activated $\text{Fe}_2(\text{DSBDC})$, as evidenced by approximately 1% reduction in local electron density of Fe, observed in both the solid-state and cluster calculations. Electron transfer thus generates holes as charge carriers, induces effectively mixed valency in $(-\text{Fe}-\text{S}-)_{\infty}$ chains, and improves charge density and electrical conductivity.³⁶

Conclusions

DMF binding to the coordinatively unsaturated Fe centers in $\text{Fe}_2(\text{DSBDC})$ leads to thousand-fold increase in the electrical conductivity at 297 K. DFT calculations identify electron transfer from Fe centers to the bound DMF molecules as a possible mechanism for improving charge density in the skeleton of the MOF. These results reinforce the importance of redox-matching between the framework and guest molecules to achieve electrically conductive MOFs, and point out the possibility of tuning electrical conductivity *via* coordinating guest molecules as well as applying MOFs with open metal sites for chemiresistive sensing of coordinating molecules.

Conflicts of interest

There are no conflicts to declare.

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Notes and references

† The structural reversibility and crystallinity of activated $\text{Fe}_2(\text{DSBDC})$ rule out desolvation-induced local destruction because it tends to be irreversible and leads to amorphization.

§ The calculation results of $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2$ reported in ref. 26 were used to analyze the electron density on Fe centers in $\text{Fe}_2(\text{DSBDC})(\text{DMF})_2$.

- 1 H. Zhou, J. R. Long and O. M. Yaghi, *Chem. Rev.*, 2012, **112**, 673.

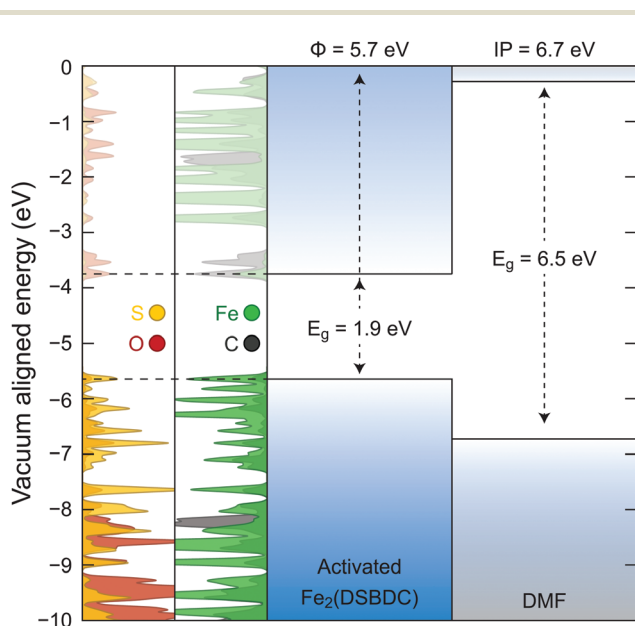


Fig. 6 Electronic band structure of activated $\text{Fe}_2(\text{DSBDC})$ and molecular orbital energy levels of DMF. The left two panels are projected density of states in $\text{Fe}_2(\text{DSBDC})$. E_g represents the band gap for $\text{Fe}_2(\text{DSBDC})$ and the HOMO–LUMO gap for DMF. ϕ and IP represent work function and ionization potential, respectively.



- 2 Z. Zhang, H. Yoshikawa and K. Awaga, *J. Am. Chem. Soc.*, 2014, **136**, 16112.
- 3 Z. Zhang, H. Yoshikawa and K. Awaga, *Chem. Mater.*, 2016, **28**, 1298.
- 4 M. L. Aubrey and J. R. Long, *J. Am. Chem. Soc.*, 2015, **137**, 13594.
- 5 K. J. Erickson, F. Léonard, V. Stavila, M. E. Foster, C. D. Spataru, R. E. Jones, B. M. Foley, P. E. Hopkins, M. D. Allendorf and A. A. Talin, *Adv. Mater.*, 2015, **27**, 3453.
- 6 L. Sun, B. Liao, D. Sheberla, D. Kraemer, J. Zhou, E. A. Stach, D. Zakharov, V. Stavila, A. A. Talin, Y. Ge, M. D. Allendorf, G. Chen, F. Léonard and M. Dincă, *Joule*, 2017, **1**, 168.
- 7 D. Sheberla, J. C. Bachman, J. S. Elias, C.-J. Sun, Y. Shao-Horn and M. Dincă, *Nat. Mater.*, 2017, **16**, 220.
- 8 K. Barthelet, J. Marrot, D. Riou and G. Férey, *Angew. Chem., Int. Ed.*, 2002, **41**, 281.
- 9 C. R. Murdock, B. C. Hughes, Z. Lu and D. M. Jenkins, *Coord. Chem. Rev.*, 2014, **258–259**, 119.
- 10 A. A. Talin, A. Centrone, A. C. Ford, M. E. Foster, V. Stavila, P. Haney, R. A. Kinney, V. Szalai, F. El Gabaly, H. P. Yoon, F. Léonard and M. D. Allendorf, *Science*, 2014, **343**, 66.
- 11 M. G. Campbell, D. Sheberla, S. F. Liu, T. M. Swager and M. Dincă, *Angew. Chem., Int. Ed.*, 2015, **54**, 4349.
- 12 X. Bao, H. J. Shepherd, L. Salmon, G. Molnár, M.-L. Tong and A. Bousseksou, *Angew. Chem., Int. Ed.*, 2013, **52**, 1198.
- 13 D. A. Reed, B. K. Keitz, J. Oktawiec, J. A. Mason, T. Runčevski, D. J. Xiao, L. E. Darago, V. Crocellà, S. Bordiga and J. R. Long, *Nature*, 2017, **550**, 96.
- 14 M. D. Allendorf, C. A. Bauer, R. K. Bhakta and R. J. T. Houk, *Chem. Soc. Rev.*, 2009, **38**, 1330.
- 15 N. B. Shustova, B. D. McCarthy and M. Dincă, *J. Am. Chem. Soc.*, 2011, **133**, 20126.
- 16 T. M. McDonald, J. A. Mason, X. Kong, E. D. Bloch, D. Gygi, A. Dani, V. Crocellà, F. Giordanino, S. O. Odoh, W. S. Drisdell, B. Vlasisavljevich, A. L. Dzubak, R. Poloni, S. K. Schnell, N. Planas, K. Lee, T. Pascal, L. F. Wan, D. Prendergast, J. B. Neaton, B. Smit, J. B. Kortright, L. Gagliardi, S. Bordiga, J. A. Reimer and J. R. Long, *Nature*, 2015, **519**, 303.
- 17 M. G. Campbell, S. F. Liu, T. M. Swager and M. Dincă, *J. Am. Chem. Soc.*, 2015, **137**, 13780.
- 18 M. Campbell and M. Dincă, *Sensors*, 2017, **17**, 1108.
- 19 K. Müller-Buschbaum, F. Beuerle and C. Feldmann, *Microporous Mesoporous Mater.*, 2015, **216**, 171.
- 20 L. Sun, M. G. Campbell and M. Dincă, *Angew. Chem., Int. Ed.*, 2016, **55**, 3566.
- 21 V. Stavila, A. A. Talin and M. D. Allendorf, *Chem. Soc. Rev.*, 2014, **43**, 5994.
- 22 Y. Kobayashi, B. Jacobs, M. D. Allendorf and J. R. Long, *Chem. Mater.*, 2010, **22**, 4120.
- 23 L. S. Xie, L. Sun, R. Wan, S. S. Park, J. A. DeGayner, C. H. Hendon and M. Dincă, *J. Am. Chem. Soc.*, 2018, **140**, 7411.
- 24 M. L. Aubrey, B. M. Wiers, S. C. Andrews, T. Sakurai, S. E. Reyes-Lillo, S. M. Hamed, C.-J. Yu, L. E. Darago, J. A. Mason, J.-O. Baeg, F. Grandjean, G. J. Long, S. Seki, J. B. Neaton, P. Yang and J. R. Long, *Nat. Mater.*, 2018, **17**, 625.
- 25 J. G. Park, M. L. Aubrey, J. Oktawiec, K. Chakarawet, L. E. Darago, F. Grandjean, G. J. Long and J. R. Long, *J. Am. Chem. Soc.*, 2018, **140**, DOI: 10.1021/jacs.8b03696.
- 26 L. Sun, C. H. Hendon, M. A. Minier, A. Walsh and M. Dincă, *J. Am. Chem. Soc.*, 2015, **137**, 6164.
- 27 N. L. Rosi, J. Kim, M. Eddaoudi, B. Chen, M. O'Keeffe and O. M. Yaghi, *J. Am. Chem. Soc.*, 2005, **127**, 1504.
- 28 P. D. C. Dietzel, Y. Morita, R. Blom and H. Fjellvåg, *Angew. Chem., Int. Ed.*, 2005, **44**, 6354.
- 29 P. D. C. Dietzel, B. Panella, M. Hirscher, R. Blom and H. Fjellvåg, *Chem. Commun.*, 2006, 959.
- 30 W. Zhou, H. Wu and T. Yildirim, *J. Am. Chem. Soc.*, 2008, **130**, 15268.
- 31 P. D. C. Dietzel, R. Blom and H. Fjellvåg, *Eur. J. Inorg. Chem.*, 2008, 3624.
- 32 E. D. Bloch, L. J. Murray, W. L. Queen, S. Chavan, S. N. Maximoff, J. P. Bigi, R. Krishna, V. K. Peterson, F. Grandjean, G. J. Long, B. Smit, S. Bordiga, C. M. Brown and J. R. Long, *J. Am. Chem. Soc.*, 2011, **133**, 14814.
- 33 R. Sanz, F. Martínez, G. Orcajo, L. Wojtas and D. Briones, *Dalton Trans.*, 2013, **42**, 2392.
- 34 W. L. Queen, M. R. Hudson, E. D. Bloch, J. A. Mason, M. I. Gonzalez, J. S. Lee, D. Gygi, J. D. Howe, K. Lee, T. A. Darwish, M. James, V. K. Peterson, S. J. Teat, B. Smit, J. B. Neaton, J. R. Long and C. M. Brown, *Chem. Sci.*, 2014, **5**, 4569.
- 35 L. Sun, T. Miyakai, S. Seki and M. Dincă, *J. Am. Chem. Soc.*, 2013, **135**, 8185.
- 36 L. Sun, C. H. Hendon, S. S. Park, Y. Tulchinsky, R. Wan, F. Wang, A. Walsh and M. Dincă, *Chem. Sci.*, 2017, **8**, 4450.
- 37 L. Sun, S. S. Park, D. Sheberla and M. Dincă, *J. Am. Chem. Soc.*, 2016, **138**, 14772.
- 38 F. Wudl and M. R. Bryce, *J. Chem. Educ.*, 1990, **67**, 717.

