

**Silicene, silicene derivatives, and their device applications**

Journal:	<i>Chemical Society Reviews</i>
Manuscript ID	CS-REV-04-2018-000338.R1
Article Type:	Review Article
Date Submitted by the Author:	27-Jun-2018
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REVIEW

Silicene, silicene derivatives, and their device applications

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Accepted 00th January 20xx

DOI: 10.1039/x0xx00000x

www.rsc.org/

Silicene, the ultimate scaling of silicon atomic sheet in a buckled honeycomb lattice, represents a mono-elemental class of two-dimensional (2D) materials similar to graphene but with unique potential for a host of exotic electronic properties. Nonetheless, there is a lack of experimental studies largely due to the interplay between material degradation and process portability issues. This Review highlights state-of-the-art experimental progress and future opportunities in synthesis, characterization, stabilization, processing and experimental device example of monolayer silicene and thicker derivatives. Electrostatic characteristics of Ag-removal silicene field-effect transistor exhibits ambipolar charge transport, corroborating with theoretical predictions on Dirac Fermions and Dirac cone in band structure. The electronic structure of silicene is expected to be sensitive to substrate interaction, surface chemistry, and spin-orbit coupling, holding great promise for a variety of novel applications, such as topological bits, quantum sensing, and energy devices. Moreover, the unique allotropic affinity of silicene with single-crystalline bulk silicon suggests a more direct path for the integration with or revolution to ubiquitous semiconductor technology. Both materials and process aspects of silicene research also provide transferable knowledge to other Xenes like stanene, germanene, phosphorene, and so forth.

1. Silicon at the two-dimensional level

The latest findings in condensed matter physics have witnessed a curious trend, where farsighted theoretical predictions have come to reality sometimes with unexpected turns. Although, on one hand, for instance, the isolation of graphene clearly contradicted the well-established Mermin's theorem,¹ conversely, on the other hand, the rise of topology as a physical (and not only mathematical) concept to interpret the new solid state phases, *e.g.* topological insulators (TIs),² is the most astonishing example to prove that theory is often able to see further on yet to be realized materials. Notably, the initial theoretical investigation subsequently flooded in a large experimental effort. Quite similarly, it can happen that, due to the periodic table kinship, one may wonder why there would not exist two-dimensional (2D) honeycomb lattices made of silicon, germanium, or tin, as they are placed just below carbon and, like a *gedanken* experiment, conceive a brand new

material. This is what occurred in 1994 when Takeda and Shiraishi proposed the aromatic stages of silicon and germanium. Indeed, as they wrote in their pioneering work, "the present infinite 2D Si aromatic stage (*i.e.* silicene) is a hypothetical material, but it is an important model for investigating the aromaticity of Si elements".³ In this humble way and perhaps with modest purpose, the concept of silicene made its debut as a 2D honeycomb lattice of silicon atoms even earlier than the rise of graphene in 2004.⁴ More recently, silicene has attracted a substantial interest as a material for nanotechnology owing to its unique solid-state properties including quantum spin Hall (QSH) effect (a 2D TI state), strong spin-orbit coupling (SOC), giant magnetoresistance, field-tunable bandgap, non-linear electro-optic effects, and piezomagnetism.^{5–11} Silicene can be considered a prototype of the elemental analogues of graphene, including germanene, stanene, *etc.* referred to as Xenes.¹² Importantly, Xenes feature non-planar atomic structure, which enhances certain physical properties such as the SOC, out-of-plane phonon scattering and other related physics. However, the metastable silicene structure also reflects in generally poor air-stability that can be addressed by proper encapsulation or passivation of reactive surfaces. The latter aspect paves the way to design 2D materials with passivated surfaces due to functionalization, including hydrogenation which is predicted to both stabilize silicene (similar to the passivation of bulk silicon)¹³ and also boost its bandgap while maintaining the atomically-thin profile.^{14,15}

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Electronic Supplementary Information (ESI) available: [details of any supplementary information available should be included here]. See DOI: 10.1039/x0xx00000x

TUTORIAL REVIEW

In this light, a major driver for silicene (and generally the Xenes) research is the potential device benefits for ultra-scaled silicon technology for conventional, flexible, and quantum or topology-based information processing. Silicon has been at the heart of complementary metal-oxide semiconductor (CMOS) technology for over five decades, especially because of the lucky combination with its oxide establishing the probably most studied interface ever.¹⁶ With Moore's law as the driving engine (namely, the temporal expansion of the number of transistors in a chip), significant performance and cost-related gains have been achieved from one technology node to the next, as the lateral size of the semiconductor devices is reduced along with the scaling of the silicon channel thickness.¹⁷ More specifically, an order of magnitude reduction in the silicon channel thickness has resulted in approximately an order of magnitude power reduction in continuously scaled devices thus promising a solid advance in the device miniaturization and densification in a chip. However, the conventional *Smart-Cut* approach (*i.e.* a top-down technology for the transfer a thin layer of silicon from a donor substrate to another substrate, mainly used for silicon-on-insulator wafer manufacturing) to produce an ultra-thin silicon body, suffers from intrinsic physical limitation to scale beyond 5 nm.¹⁸ On the other hand, starting from a bottom-up approach, it is possible to synthesize 2D mono-elemental sheets of silicon, resulting in the ultimate thickness scaling and the associated reduction in electronic device power consumption. Moreover, while carrier mobility deteriorates with scaling in the traditional ultrathin body silicon, this is not the case with silicene or 2D silicon.^{19,20} 2D silicon can also be stacked vertically to enable three-dimensional integrated chips in a layer-by-layer fashion.

In addition to its conventional semiconducting properties, certain epitaxial phases of silicene that feature the QSH effect can be employed for making topological devices such as a gate-controlled topological field-effect transistor (FET) that is based on topological phase change.¹² However, for all the novel concepts based on Xenes to be realized for applications, the air-stability issue and process integration onto technologically relevant substrates will have to be decisively tackled. In this Review, our intent is to briefly retrace the silicene's thread by providing an overview of the chemistry and physics of silicene, including the most relevant experimental achievements in terms of synthesis and chemical functionalization (Section 2), methodologies for silicene processing (Section 3), and for silicene integration into nanotechnology platforms (Section 4).

2. Epitaxial silicene

2.1 Chemistry and physics of silicene

The theoretical background of silicene and germanene dates back to the seminal paper by Takeda and Shiraishi on the hexagonal rings of silicon and germanium with reference to

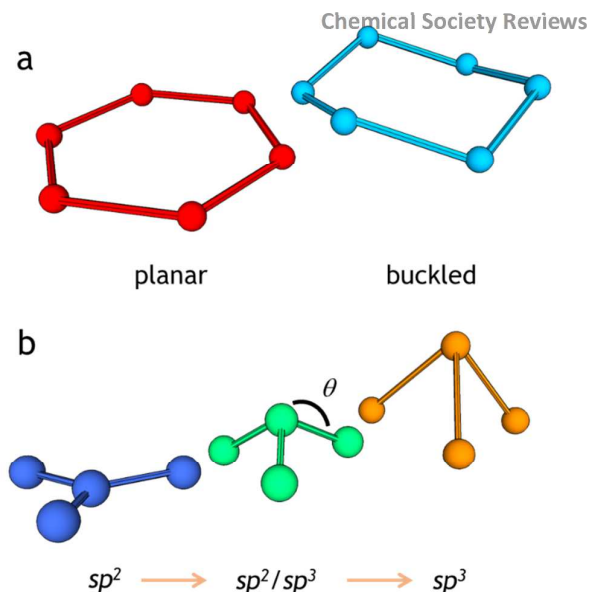


Figure 1 Chemical bonds in silicene. a) planar and buckled hexagonal rings stages, b) evolution from sp^2 to sp^3 hybridization of the four-bonded group 14 Xenes.

the carbon counterpart.³ These aromatic stages were predicted to be more stable in a regularly corrugated structure (so-called buckled with a D_{3d} group symmetry) than the flat state (D_{6h} group symmetry) stage that is specific of carbon. This picture was then confirmed by total energy minimization *ab initio* calculations.²¹ This intrinsic puckering, hereafter referred to as buckling, is the origin of many peculiar properties (Figure 1a). The mirror symmetry breaking in silicene by buckling removes the instability associated with the planar high-symmetry structure. Buckling arises from the pseudo Jahn-Teller distortion associated with the coupling between occupied and unoccupied molecular orbitals where σ - π mixing allows the system to gain stability.²² It is noteworthy, conversely, that the low buckling does not break the hexagonal symmetry, *i.e.* for the honeycomb lattice the two atoms in the unit cell are equivalent and electrons can hop among the nearest-neighbour atoms, thus retaining the existence of Dirac fermions.²³ In this framework, the Si-Si bond length in the buckled state (2.247 Å) is longer than the flat stage (2.226 Å), thereby resulting in a larger effective volume of the lattice unit size.^{3,22} Nonetheless, electrons in the buckled state are more delocalized (in and out the stage plane) producing a smaller inter-electron repulsive energy. This deformation also causes a partial sp^3 hybridization, which is a second marked difference with respect to graphene. While silicon prefers sp^3 hybridization in its bulk form, an interplay between sp^3 and sp^2 hybrid bonds is characteristic of silicene (Figure 1b). The mixed sp^2 - sp^3 hybridization state can be understood as evolving with the buckling and, more specifically, with the θ angle, *i.e.* the angle between the Si-Si bond and the direction normal to the plane. Hence, the sp^2 (planar), low-buckled (mixed sp^2 - sp^3), and sp^3 configurations correspond to $\theta=90^\circ$, $\theta=101.73^\circ$, and $\theta=109.47^\circ$, respectively

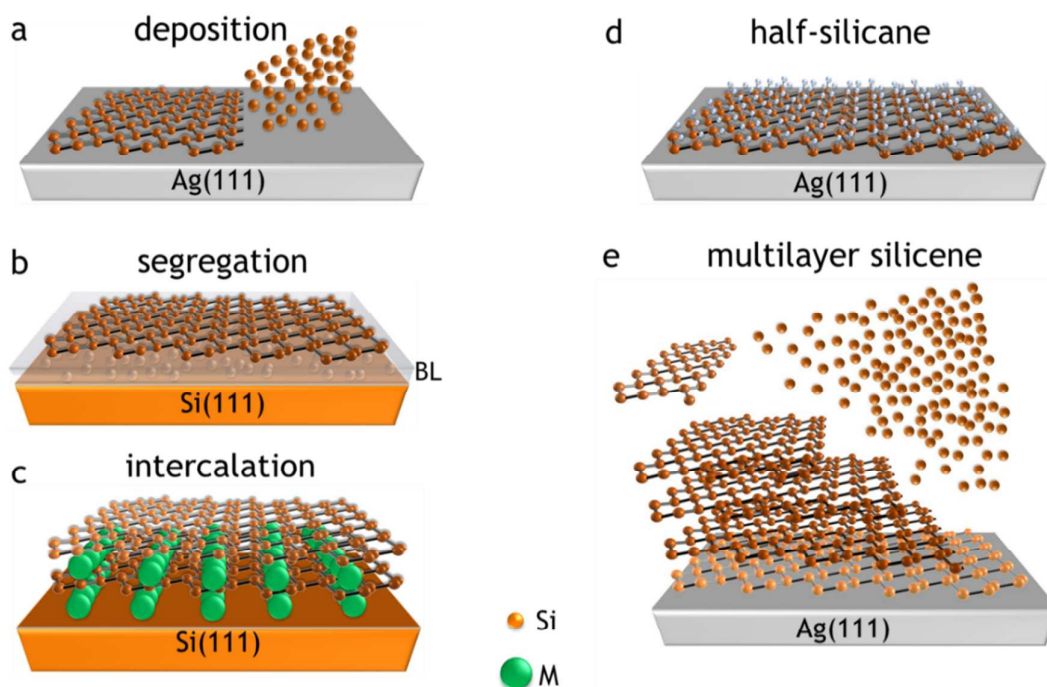


Figure 2 Silicene growth methodologies. a) Epitaxial silicene by deposition where the hot incoming atoms are condensated onto a supporting template: the case of epitaxial silicene on Ag(111),⁴⁸ b) Epitaxial silicene by segregation where the supporting substrate acting as reservoir through a buffer layer: the case of epitaxial silicene on ZrB₂ as buffer layer grown on Si(111),⁷⁵ c) intercalated silicene where foreign atoms are intercalated in a pre-established network: M (Eu, Sr, or Gd) atom are supplied by a film pre-grown on a Si(111) substrate and then intercalate through a layered silicide network,^{81–83} d) Functionalization of silicene exemplified in the half-silicene after H₂ exposure to Ag(111)-supported silicene,¹⁰³ e) silicene multilayer by deposition on a Ag(111) substrate.¹⁰⁸

(Figure 1b).⁵ Despite the stability issues that may arise from unsaturated sp^3 bonding, buckling represents an additional degree of freedom to manipulate the silicene properties such as bandgap opening, electronic structure, incorporated staggered field, selective chemical reactivity. On the other hand, buckling exposes silicene to have mixed sp^2 and sp^3 hybrid bonds, which results in a high environmental reactivity eventually leading to material degradation (see Sections 3.1 and 3.2).

In terms of the electronic band structure, buckled silicene appears as a graphene-like Dirac crystal where π and π^* bands cross linearly at the Dirac points, namely high symmetry K and K' points in the hexagonal Brillouin zone.^{21,24,25} Additional theoretical efforts on the freestanding silicene focused on its optical and topological properties. Indeed, similar to graphene, the low-frequency absorbance of freestanding silicene and other Xenes is basically determined by the fine-structure constant (or Sommerfeld's constant) $\alpha = e^2/\hbar c$ (e , \hbar , and c being the electron charge, the reduced Planck constant, and the speed of light, respectively), irrespective to the X atom or the buckling character of its bonds, whereas differences occur at higher frequencies due to interband transitions related to the van Hove singularities of the joint density of states (DOS).²⁶ On the other hand, when the lattice gains weight from the (planar) graphene to (buckled) Xenes for X spanning from silicon to tin, a topologically nontrivial electronic structures sets in due to the increasing SOC, which results in the QSH

effect, namely the physical hallmark of a 2D TI phase of matter.^{5,9,25,27} Based on the Kane-Mele model,²⁸ the effective SOC in the buckled Xenes opens a gap at the Dirac points that is topological in character, and hence endowed with conductive edge states at the geometrical borders, *e.g.* at the boundaries of a silicene ribbon. Furthermore, the stronger the effective SOC is at the Dirac points, the larger is the gap opening. That is why topological properties are expected to be more robust in Xenes with increasing atomic mass such as germanene and stanene.¹² In this scenario, an even richer topological phase diagram can be envisaged which includes quantum anomalous Hall effect states, valley-polarized metal phases, and chiral superconducting states as a function of an externally applied electromagnetic field or intrinsic exchange field.^{5,12,29–31}

Driven by this perspective, the recently developed paradigm of epitaxial silicene on substrate has been extended to other Xenes^{12,32} with a nearly silicene-like approach to the synthesis and identification. In this respect, an increasing theoretical and experimental effort has been recently devoted to template engineering, namely designer weakly or non-interacting substrates serving as templates for silicene. Promising substrates of this kind include insulators,³³ semiconductors,^{34–38} and metals.³⁹

2.2 Main achievements in the synthesis of silicene

Due to the absence of a graphite-like form of silicon in nature, silicene as well as all the Xenes, is synthesized by means of a bottom-up approach, namely epitaxial growth on a substrate. Despite the relatively high cost, this approach offers the potential wafer scale production targeting technology applications. Physical identification of silicene (and Xenes) usually relies on interplay between an experimental tool such as scanning tunnelling microscopy/spectroscopy (STM/S) or angle-resolved photoelectron spectroscopy (ARPES), and density functional theory (DFT) calculations. Three general routes have been implemented for the silicene epitaxy so far (see **Figure 2**): a) thermal evaporation on substrate (*epitaxial silicene by deposition*); b) surface segregation from a substrate (*epitaxial silicene by segregation*); c) intercalation through a silicene network (*intercalated silicene*).

a) *Epitaxial silicene by deposition*. The process stems from the condensation and self-organization of thermally evaporated silicon atoms onto a substrate (**Figure 2a**). Although the number of substrates on which silicene can be synthesized is rapidly expanding, the (111)-terminated silver substrate is the only one enabling both the synthesis and device integration as yet. In this respect, it is important to note that while Ag(110) was originally shown to accommodate silicon nanoribbons with hexagonal⁴⁰ or pentagonal⁴¹ structure, Ag(111) has proven to be a quite universal template for other Xenes such as borophene,^{42,43} germanene,⁴⁴ stanene,⁴⁵ and antimonene.⁴⁶ Moreover, the Ag(111) template can be suitably reduced to an epitaxial film on mica or Si(111)⁴⁷ substrates, thus avoiding the unpractical use of highly expensive monocrystalline silver substrate and enabling delamination in subsequent processing (see Section 3). Lattice commensuration with the freestanding form of silicene and a narrow growth parametric window (mostly temperature and deposition flux) are the key prerequisites of the sought-after recipe for the synthesis of silicene on Ag(111) by epitaxial deposition.^{48,49} This methodology has been successfully replicated on other substrates, including iridium,⁵⁰ molybdenum disulphide (MoS₂),⁵¹ zirconium carbide (ZrC),⁵² ruthenium,⁵³ graphite.⁵⁴ In most cases, epitaxial silicene by deposition takes place on a large-scale through the emergence of differently reconstructed domains. Reconstructions arise as lattice-distortion of the freestanding silicene driven by the commensurability relationship with the surface lattice of the substrate.^{55–58} Based on STM investigations, specifically on Ag(111)-supported silicene, periodic reconstructions mainly include 4×4, √13×√13, 2√3×2√3 surface phases (terminologies refer to the coincidence of silicon atom with respect to the surface atoms of the substrate, see **Figure 3a**) of the epitaxial silicene opposed to the alternating sequence of up and down adjacent atoms in freestanding silicene, namely a 1×1 phase. As such, the buckling distribution, *i.e.* the periodic arrangement of non-planar buckled bonds in each surface phase, and hence, the Si-Si bond length can vary from one superstructure to another. In case of the Ag-supported silicene, a comprehensive phase diagram can be sketched up as a function of the substrate temperature and silicon coverage.^{56–61} To complement STM studies, ARPES provides evidences of the electronic band

structure of silicene. In these kind of studies, some interference is produced by the inevitable orbital interaction of silicene with the substrate, which makes the silicene metallic in character.^{62,63} As such, isolation of the true characteristics of silicene when its supported on silver is made difficult by the bandfolding of the silicene reconstructions and the linear dispersion of silver bands near the Fermi level.^{62–73} The case of Ag-supported silicene is then qualified by the suppression of the π molecular orbitals as a consequence of the hybridization of the silicon and silver electronic states, whereas the σ orbitals are clearly identified.⁶⁵ Nonetheless, evidence of a low-energy plasmonic excitation was recently deduced from electron energy loss spectroscopy, and it is consistent with characteristic π -like plasmon in freestanding silicene.⁷⁴

b) *Epitaxial silicene by segregation*. In this configuration, the reservoir of silicon atoms is supplied by the underlying substrate instead of the deposition flux (**Figure 2b**). An example in this respect is the case of the silicon atoms thermally diffusing from a (111)-terminated silicon substrate up to the surface of a thin crystalline zirconium diboride (ZrB₂) film.⁷⁵ Similar to the epitaxial silicene by deposition, the self-organization of the silicene lattice is dictated by the commensurability with the ZrB₂ surface lattice. Interaction with the substrate determines the electronic band structure of the overall surface system but contrary to the metallicity of the metal-supported silicene, the ZrB₂-supported silicene results in a gapped electronic structure.

c) *Intercalated silicene*. Complementary to the physical deposition techniques usually adopted in the silicene framework, the synthesis of silicon nanosheets can be carried out via chemical methods. The chemistry versatility allows for synthesizing silicon nanoparticles/dots, silicon fullerenes, silicon nanowires, silicon nanotubes, silicon nanoribbons (for details see Ref.⁷⁶ and references therein). A paradigmatic example is represented by silicon nanosheets naturally embedded in calcium disilicide (CaSi₂).^{77,78} The idea here is to find pre-formed silicene-like structures to be synthesized via the topochemical deintercalation of the silicide.⁷⁹ CaSi₂ is a Zintl silicide, where 2D silicon puckered sheets made of Si₆ rings are separated from each other by planar monolayers of Ca²⁺, *i.e.* intercalated multilayer silicene. Chemical manipulation through fluoride diffusion into CaSi₂ results in the formation of silicon layered structures in between calcium fluoride (CaF₂) planes with a reduced number of unsaturated silicon bonds and bandgap opening.⁷⁸ Starting from a similar framework, Tokmachev *et al.* have reported on the *ad hoc* tailored silicidation where silicon atoms from a bulk substrate thermally intercalated through epitaxially deposited strontium, europium, or gadolinium so as to form Zintl phase silicide (*e.g.* SrSi₂, EuSi₂, GdSi₂) where silicene nanosheets are incorporated as network constituents (see **Figure 2c**).^{81–83}

Although the quest for silicene still justifies the survey for a special and easy route of its reliable synthesis, other silicon-based nanostructures have been recently investigated as well. Indeed, every kind of nanoscaled silicon compatible with the

electronics industry represents a possible solution to replace cubic silicon in scalable devices and fulfil the technology scopes as outlined in Section 1. In this framework, theoretical efforts were mainly driven to explore allotropic phases of silicon other than silicene,^{84–86} whereas the experimental activity successfully provided evidences of hexagonal silicon, so called lonsdaleite phase,⁸⁷ or the new orthorhombic allotrope of silicon called Si₂₄.⁸⁸ Another flourishing field is related to the low-dimensional silicon nanosheets that are endowed with unconventional properties. Within this latter class, the generic term of silicon nanosheets refers to silicon nanomaterials and nanostructures, which exhibit properties different from those of bulk silicon due to the quantum confinement effects.⁷⁶ For instance, cubic silicon nanosheets in the thickness range 1–13 nm display strong thickness-dependent photoluminescence in visible range with the bandgap energies ranging from 1.6 to 3.2 eV.⁸⁹ In nearly the same thickness range (2–18 nm), the excitation energy-dependent Raman intensity of ultrathin silicon nanomembranes relies on the combined effects of interference and resonance from the band structure modulation.⁹⁰

Apart from the epitaxial methods, chemical processing is rapidly emerging as a route to achieve large-scale and low-cost production of silicon nanosheets with thickness of ~4 nm and lateral size of several micrometres. In this respect, Lang *et al.* developed a scalable synthesis of ultra-low friction silicon nanosheets by means of the intrinsic delithiation process of the Li₁₃Si₄ alloy, that possesses an orthorhombic structure markedly different from that of CaSi₂.⁹¹

2.3 Silicene derivatives and functionalization

Two main follow-up topics tail the epitaxial synthesis of silicene: chemical functionalization and pile-up stacking of multi-layered silicene. Both routes are conceived to take silicene single layer as a precursor for new artificial material with expanded functionalities. On the former aspect, a number of theoretical works has conceptualized the absorption of foreign chemical atoms/molecules to yield functionalized silicene,^{92–97} but the experimental activity has been mainly focused on the oxygen and hydrogen functionalization thus far. Oxidation is probably the easiest route to modify the silicene bandgap.⁹⁸ Silicene can be exploited as precursor stage for silicene oxides,^{99,100} alternative to hexagonal boron nitride (hex-BN) and in parallel with current effort to synthesize 2D honeycomb-like silica.¹⁰¹ Interestingly, oxygen intercalation allowed the formation of a freestanding-like silicene.¹⁰² The oxygen atoms intercalate into a bilayer silicene on the Ag(111) surface, resulting in the isolation of the top layer of silicene. This layer exhibits the signature of a 1×1 honeycomb lattice and hosts massless Dirac fermions due to the weaker interaction with the substrate, similar to the intercalation approach described in Section 2.2. Hydrogen absorption into silicene is also appealing as a molecular dissociation path for

hydrogen evolution reaction.¹⁰³ As full hydrogenation is inhibited by the substrate at the bottom surface, final product is the half-silicane, *i.e.* just the top surface undergoes hydrogenation therein showing an asymmetric out-of-plane structural configuration^{103,104} (see **Figure 2d** as a pictorial sketch for half-silicane). Interestingly, the silicene hydrogenation process has been demonstrated to be reversible.¹⁰⁵

Retracing back to the route of silicene's synthesis, one may wonder whether graphite-like silicon can be artificially created by piling up single layers of silicene. Such a crystal, termed silicite, was predicted to set in as a new thermodynamically stable layered-phase of silicon, characterized by a stacking of dumbbell patterned silicene sheets with strong directionality in the electronic and structural properties.^{106,107} Additionally, the increased absorption in the visible range would make silicite,¹⁰⁶ as well as other allotropic silicon phases^{84,88} very interesting for a silicon-based photonics. Experimentally, multilayer silicene is referred to as the sequential deposition of individual silicene layers by epitaxy starting from the first one as a template (see **Figure 2e** for a pictorial sketch for multilayer silicene). The as-defined multilayer silicene on Ag(111) shows an island growth mode with a characteristic termination, which is independent on the number of layers (≥ 2).^{108–111} Similar to the single-layer silicene, the presence of Dirac fermions and the layered nature of this artificial crystal are highly debated in literature^{111–113} and the true character of the surface termination (being made of silicon or segregated silver from the substrate) is still controversial.^{114–116} Nonetheless, increasing the growth temperature was proven to govern the transition from the as-defined multilayer silicene to cubic silicon.¹¹⁷ Although the multilayer silicene on Ag(111) is yet to be fully understood, we emphasize on the role played by seed single layer silicene in determining its structural and electronic properties.¹¹⁷ For this reason, multilayer silicene can be regarded as a derivative of silicene. Alternative to the epitaxy, multilayer silicene growth has been also reported by ion implantation.¹¹⁸

2.4 The class of epitaxial Xenexes

The route of silver templating not only drove the epitaxial synthesis of silicene but also inspired an intense effort to synthesize akin material made of elements other than silicon. These are the epitaxial Xenexes. The term Xenexes, originally restricted to group 14 elements,¹² is being gradually expanded to include the ever-growing class of 2D mono-elemental crystals spanning from group 13 to 16 of the periodic table. Experimentally realized Xenexes, today, stretch out from the 2D icosagens (from group 13 elements) like borophene^{42,43} and recently reported gallenene,¹¹⁹ sought-after, to 2D pnictogens (*e.g.* phosphorene, antimonene, and bismuthene)^{120,121} and 2D chalcogens (selenene, tellurene).¹²² Leveraging on the advances from synthesizing silicene, molecular beam epitaxy is gaining the largest popularity, mainly because of its ability to offer atomic-scale precision. Therefore, we focus on Xenexes

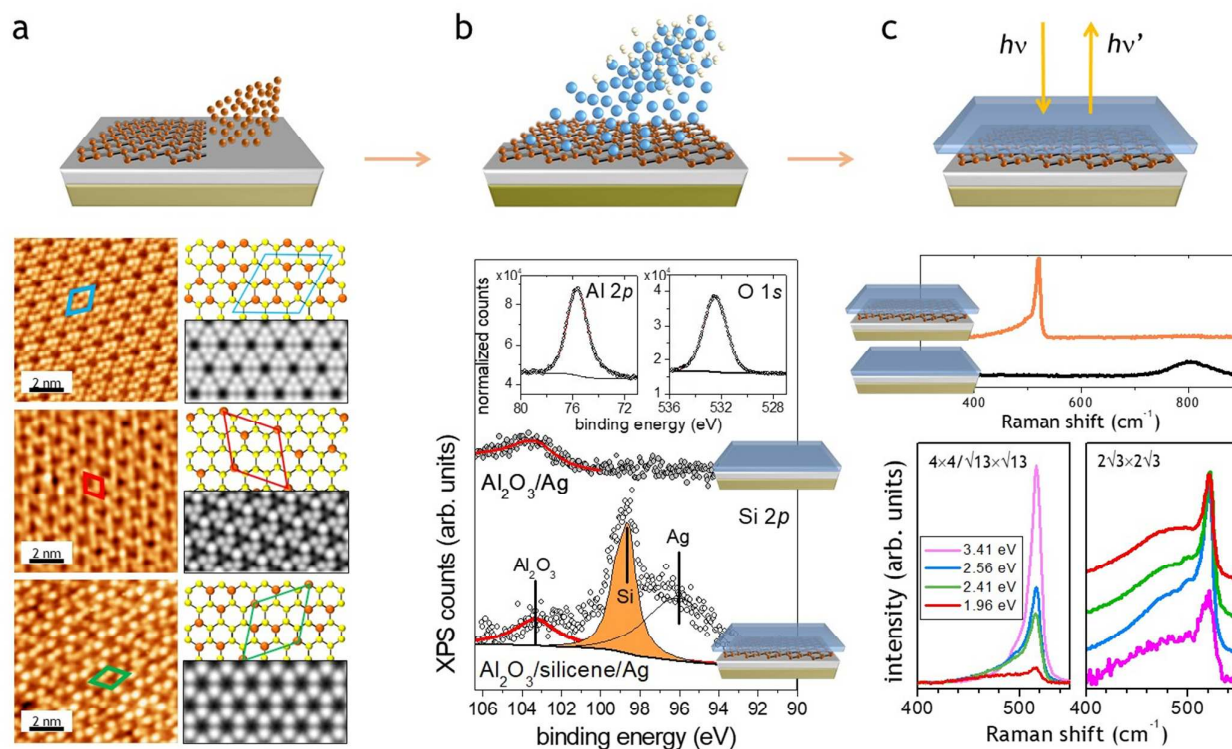


Figure 3 Route for silicene encapsulation. a) Epitaxial growth of silicene on Ag(111)/mica substrates. The three different silicene phases (4×4 , $\sqrt{13} \times \sqrt{13}$, and $2\sqrt{3} \times 2\sqrt{3}$, from top to bottom) are identified by high-resolution STM (left side), explained in terms of atomic lattice sketches (right side, top) where yellow and orange spheres correspond to lower and upper atoms, respectively, rhombi highlight the unit cells, and matched with the respective DFT simulations. Adapted from Ref.¹⁸⁰ with permission from the IOP publishing, copyright 2012. b) Encapsulation by means of reactive co-deposition of Al and O₂. Probing the Si 2p core-level (bottom spectra) in the silicene-free control Al₂O₃/Ag structure and in the Al₂O₃/silicene/Ag stacked structure makes evidence of the integrity of the elemental silicon bonding (orange peak) retained in the encapsulated silicene (left side feature at higher binding energy is assigned to Al₂O₃ background after comparing the two reported structures); formation of a stoichiometric Al₂O₃ capping layer is confirmed by the observation of the Al 2p and O 1s core-level lines (top spectra). Data are adapted from Ref.¹⁴¹ with permission from the John Wiley and Sons, copyright 2013. c) Raman spectroscopy (excitation energy: 2.41 eV) of encapsulated silicene. Raman spectra are acquired with and without silicene in between an Ag-on-mica substrate and the Al₂O₃ capping layer (top panel); when silicene is in, a well-defined peak is measured as a characteristic signature. The multi-wavelength (excitation laser energies from the ultraviolet down to the visible blue, green, and red are listed in the legend) Raman spectroscopy study (bottom panel) shows a resonant and non-resonant behaviour for the encapsulated $4 \times 4/\sqrt{13} \times \sqrt{13}$ (left) and $2\sqrt{3} \times 2\sqrt{3}$ (right) phases, respectively. Adapted from Ref.¹⁴⁸ with permission from the American Chemical Society publishing, copyright 2013.

grown by epitaxy in this section. Experimental facts for the Xene identification are generally based on STM/S and ARPES investigation supported by DFT models. Epitaxial growth was successfully applied to the following groups of elements in the periodic table:

- *Group 13 (icosagens)*: borophene was demonstrated independently by two groups on Ag(111) surface^{42,43} bringing evidence of a two different unbuckled phases with metallic nature;
- *Group 14*: growth of metallic germanene has been demonstrated on several metallic (111)-terminated surfaces (so far including aluminium, copper, gold, platinum, antimony),^{123–127} and a limited number of non-metallic substrates [MoS₂ and hexagonal aluminium nitride (hex-AlN)],^{128,129} being of metallic character in both cases; evidences of epitaxial stanene are limited to substrates where stanene originates from a single-layer of the α -tin phase, such as bismuth telluride (Bi₂Te₃)¹⁰⁹ and InSb(111),¹³⁰ or as a second layer on top of the Ag(111) surface;⁴⁵

- *Group 15 (pnictogens)*: the monolayer of hexagonal phosphorus, namely phosphorene (sometimes termed blue phosphorus),¹³¹ was reported on Au(111) surface¹³² with a semiconducting character (gap of 1.12 eV from the local density of states) and a silicene-like buckled structure (opposed to the more stable black phosphorus);¹³² other epitaxially grown 2D pnictogens in this number include bismuthene on Bi₂Te₃¹³³ and silicon carbide (SiC)¹³⁴ substrates, and antimonene on germanium, palladium telluride (PdTe₂), and silver substrates.^{46,135,136}
- *Group 16 (chalcogens)*: it is the case of van der Waals (vdW) epitaxy of 2D tellurium, termed tellurene, on highly oriented pyrolytic graphite (HOPG)¹³⁷ and graphene/6H-SiC substrates,¹³⁸ and 2D selenium, termed selenene, on a freshly cleaned Si(111) substrate.¹³⁹ Unlike previously mentioned Xenes, the case of tellurene and selenene differs in that they are inherently made of an array of parallel atomic chains arranged on a 2D hexagonal framework. As such, they are referred to as one-dimensional vdW crystal.¹⁴⁰

Aiming at a universal approach to device fabrication (see Section 3.5), it deserves notice that most of the epitaxial Xenos listed above are hosted by (111)-terminated silver or gold templates that in turn bear epitaxy on cleavable mica substrates. As such, they are prone to be treated by delamination in a subsequent processing stage as detailed in Section 3.

3. Silicene processing

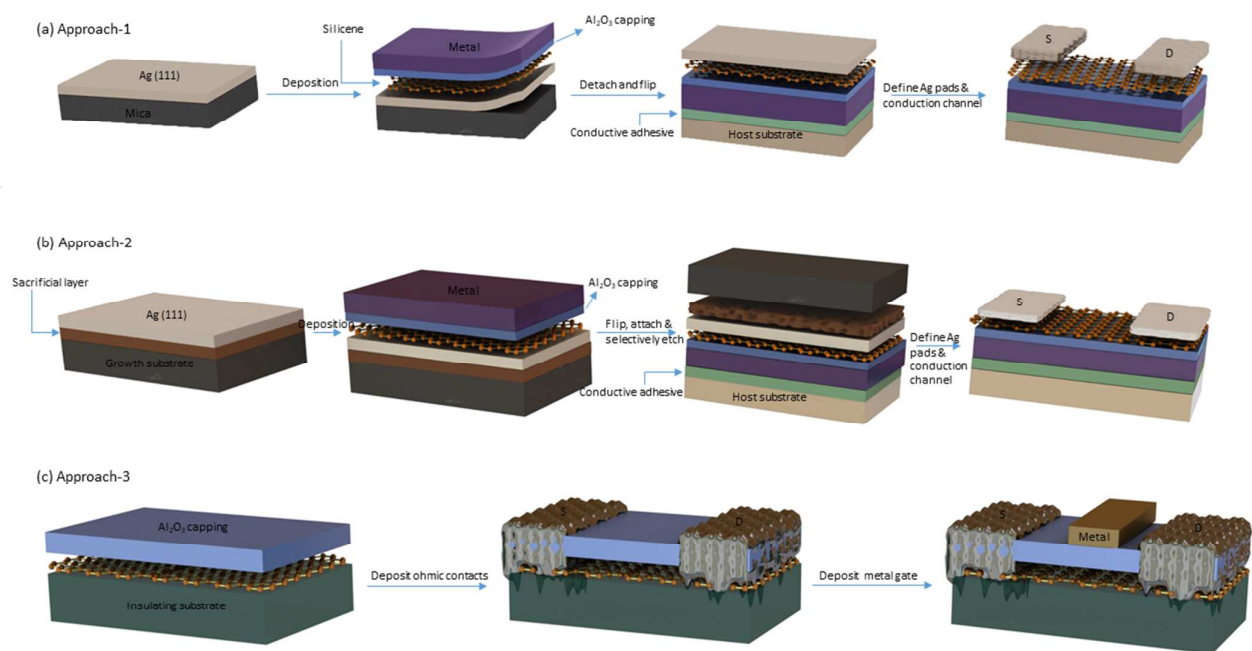
3.1 The route to silicene devices: interface engineering

The oxidation of silicene is quite limited when exposed to pure molecular oxygen (O_2) under ultra-high vacuum (UHV) conditions,^{141,142} but its unavoidable degradation under ambient conditions is widely reported.^{100,143,144} Even though the exact oxidation mechanisms deserve further insights (see Section 2.3), the extreme surface sensitivity of silicene, arising from its metastable nature, requires strategic interface engineering to preserve its properties right from its synthesis, all the way down to device fabrication and post-fabrication characterization. There are three technical milestones, all related to strategic interface engineering, in silicene processing towards a functional device. The first milestone is the synthesis of silicene on cleavable Ag(111)/mica substrates with *in situ* STM and *ex situ* Raman spectroscopy as front-end characterization (**Figure 3**). The second milestone is the design of aluminium oxide (Al_2O_3)/silicene interface with Al_2O_3 as a capping layer to preserve silicene when exposed to post-deposition in ambient condition. As seen in Section 3.2, these two progresses addressed the accessibility issues (large-scale silicene sheet with affordable separation process from growth substrate) and stability issues (pristine silicene after deposition but *before* transfer or decoupling from growth substrates to a device substrate) in silicene processing. The latest milestone is the design of transfer and device fabrication process with top and bottom interfaces of silicene sandwiched in between capping Al_2O_3 layer¹⁴¹ and underline growth catalyst Ag(111), known as silicene encapsulated delamination with native electrodes (SEDNE) process.¹⁴⁵ SEDNE process (details and development in Section 3.3) addressed both stability and portability (intact silicene *during* transfer/decoupling and device fabrication) issues. These two aspects share a common cause: silicene is a metastable phase originated from its mixed sp^2 - sp^3 silicon sheet (see Section 1.2).

3.2 Encapsulation and Raman spectroscopy of silicene

Here we will focus on key enabling technologies for accessibility and stability of silicene. Accessibility of a silicene sheet was enabled by the successful use of a cleavable mica substrate as host for the Ag-supported silicene (**Figure 3a**). This is basically due to the fact that mica can epitaxially accommodate Ag(111) film with high structural quality. Based on STM investigations in **Figure 3a**, silicene phases as detailed in Section 2.2 can be nicely reproduced on an epi-Ag(111) film supported by mica with the extra-value of being cleavable from the background substrate compared to conventionally used (111)-terminated silver monocrystals. On the other hand, handling of silicene is hurdled by the stability of the top (exposed) surface

while the bottom one is protected by the interface with the silver catalyst substrate. Stabilization of Ag-supported silicene was first reported by means of *in situ* encapsulation with an ultra-thin (few nm) Al_2O_3 layer. Al_2O_3 is sequentially grown after silicene by means of reactive co-deposition of an aluminium flux in O_2 rich pressure (partial pressure 1×10^{-6} mbar) at room temperature in an UHV background environment (**Figure 3b**).¹⁴¹ Chemical evidences of stoichiometric Al_2O_3 as well as silicene integrity after co-deposition is given by *in situ* X-ray photoemission spectroscopy (XPS) that rules out any intermixing or compound formation involving silicene. A sketch of the Al_2O_3 encapsulation process and related XPS survey are illustrated in **Figure 3b**. The same methodology has been successfully implemented in highly-buckled silicene grown on MoS_2 where the chemical integrity of the silicon nanosheet is again demonstrated by *in situ* XPS.¹⁴⁶ This is not the case of the silicene grown on ZrB_2 where dissociative chemisorption is observed to take place in between silicon and O_2 thus resulting in an Al-mediated oxidation of the ZrB_2 -supported silicene; AlN is therein proposed as alternative solution bypassing the silicene oxidation.¹⁴⁷ Nonetheless, the effectiveness of the Al_2O_3 capping layer on the Ag-supported silicene has been validated by the observation of a characteristic Raman spectrum of silicene,¹⁴⁸ as discussed in the following. Raman spectroscopy offers a quick and effective identification tool that is conventionally used to characterize graphene (and in general, carbon-based nanosystems) and other 2D layered materials in terms of the atomic structure, disorder, defects, and electronic properties.¹⁴⁹ Despite the presence of the metal substrate (that is intrinsically Raman silent), the Raman spectrum of the Ag-supported silicene exhibits an intense peak located at 516 cm^{-1} (for the 4×4 and $\sqrt{13} \times \sqrt{13}$ phases) or 521 cm^{-1} (for the $2\sqrt{3} \times 2\sqrt{3}$ phase) presenting an asymmetric and broad shoulder at lower frequency ($440\text{--}500\text{ cm}^{-1}$) as reported in **Figure 3c** (top panel).¹⁴⁸ DFT calculations rationalize the intense peak as the Raman active E_{2g} mode for the silicene superstructures.^{150,151} Similar to graphene, the zone-center E_{2g} vibrational mode represents an in-plane displacement and is due to the bond stretching of all pairs of silicon atoms lying in six-atom rings with frequency being strictly dependent on the Si-Si bond length. For freestanding silicene the E_{2g} mode is expected at higher frequency ($\sim 570\text{ cm}^{-1}$) because of the shorter bond length (2.28 \AA , when calculated as *infinite* sheet).²⁴ Hence, the softer E_{2g} mode of epitaxial silicene superstructures is related to the slightly longer Si-Si distance in the 4×4 phase ($2.34\text{--}2.39\text{ \AA}$) and the $\sqrt{13} \times \sqrt{13}$ phase ($2.31\text{--}2.36\text{ \AA}$) with respect to the cubic silicon. Conversely, the E_{2g} mode frequency of the $2\sqrt{3} \times 2\sqrt{3}$ silicene phase is blueshifted with respect to both the 4×4 and the $\sqrt{13} \times \sqrt{13}$, according to a shorter mean bond length ($2.28\text{--}2.37\text{ \AA}$). Along with the purely in-plane E_{2g} mode, additional out-of-plane modes contribute to the overall Raman spectrum of the Ag-supported silicene.¹⁵² Basically, these consist of *breathing-like* displacements of the (planar and non-planar) hexagon rings constituting the silicene lattice.¹⁴⁸ An additional point of interest for the Raman spectroscopy of the Ag-supported silicene is the phase-



dependent resonance behaviour (**Figure 3c**, bottom). Resonant effects are usually measured in semiconductors as a function of the excitation wavelength. This is the case of the mixed $4 \times 4 - \sqrt{13} \times \sqrt{13}$ silicene phase where an increasing Raman intensity is measured with increasing excitation wavelength. Conversely, no resonance is observed in the $2\sqrt{3} \times 2\sqrt{3}$ phase. This discrepancy is rationalized in terms of the different electronic band structure that are expected to take place with varying silicene phases, the 4×4 phase being qualified by a semiconducting bandgap and parabolic bands, whereas the $2\sqrt{3} \times 2\sqrt{3}$ one by Dirac bands with a semi-metallic character.^{144,150,151} Although the resonant Raman scattering appears to be sensitive of the inner electronic structure of the Ag-supported silicene, this is not the case of photoemission and optical spectroscopies where silicene turns out to be strongly interacting with the substrate thus resulting in the suppression of the (expected) Dirac cones, the emergence of metallic hybrid bands arising from silicon *p* and silver *d* states,^{62,65} and a metal-like carrier dynamics.⁶³ The interpretation of the Raman spectrum of the encapsulated silicene is consistent with subsequent *in situ* investigations of the vibrational spectrum of the freshly grown (*i.e.* uncapped) Ag-supported silicene.^{153,154} In these studies, not only the asymmetric E_{2g} peak is recognized but also additional features

in the low frequency spectral range (below 220 cm^{-1}) are reported that were initially associated with defect-induced D mode in analogy with graphene.¹⁵³ A closer insight into the latter aspect was gained by matching the experimental low-frequency datum with the phonon spectrum of the freestanding silicene¹⁵⁴ as well as by local inspection via tip-enhanced Raman spectroscopy (TERS).¹⁵⁵ Both approaches assign the low-frequency features to zone-centered modes A1 and A2 related to out-of-plane optical (ZO) phonons. Similar to the case of a single layer, the multilayer silicene (see Section 1.4) is also characterized by a sharp Raman mode that is blueshifted up to 526 cm^{-1} with respect to that of bulk silicon.¹¹⁷ A similar shift is observed starting from the early growth stages of the $\sqrt{3} \times \sqrt{3}$ silicene phase that is a precursor for the multilayer silicene, in addition to a characteristic multipeak profile in the low-frequency range associated with edge-induced scattering.¹⁵³ Overall, though not as direct and explicative as the Raman spectrum of graphene, monitoring the main Raman-active mode on the encapsulated silicene proved to be a fast, non-destructive, and versatile tool to check the silicene status throughout a process sequence outside the vacuum ambient. Optical absorbance techniques are also emerging as an alternative option to probe the characteristic electronic DOS in silicene.^{26,63}

3.3 Delamination transfer and device fabrication

Subsequent to the aforementioned encapsulated deposition of silicene, it is critical to keep sandwich style $\text{Al}_2\text{O}_3/\text{silicene}/\text{Ag}(111)$ film stack intact due to stability and portability concerns during operational approaches in delamination transfer and device fabrication steps. Unlike graphene and other stable 2D materials, substrate etching based or wet transfer methods^{156–158} cannot be readily applied to the case of silicene. The current SEDNE process entails mechanical delamination with a two-tape method to separate the $\text{Ag}/\text{silicene}/\text{Al}_2\text{O}_3$ stack from the mica substrate,¹⁴⁵ and then the application of the stack onto the device substrate (Approach-1 in **Figure 4a**). The first tape picks up $\text{Al}_2\text{O}_3/\text{silicene}/\text{Ag}$ stack out from the bulk mica substrate, thanks to mica being a layered cleavable substrate. However, this delamination may not be perfect since some mica residues may remain in certain areas leading to a partial exposure of the $\text{Ag}(111)$ surface. A second tape, such as thermal releasing tape, then steps in to flip the $\text{Ag}/\text{silicene}/\text{Al}_2\text{O}_3$ film stack with Al_2O_3 dielectric facing down before applying to a back-gate substrate. Both a blue tape with the least amount of adhesive and a thermal releasing tape (120–150 °C) are used in the delamination and transfer step. It is then feasible to obtain films, which are a few cm^2 in area with the delamination and transfer, with the uniformity of the films being monitored by Raman mapping. SEDNE process sets a milestone in silicene device study, as a significant breakthrough to enable silicene transistor makes its debut, opening up exciting avenues of research and technology development. Below we envisage some perspectives to extend or upgrade SEDNE processing for the upcoming generation of silicene devices.

The delamination and transfer steps, as described above, significantly affect the device performance and yield. For the delamination step, it demands an interface engineering between supporting substrate and catalyst film. For the transfer step, it is critical to have a seamless contact between Al_2O_3 capping layer and device substrate. Both aspects call for innovative ideas to optimize silicene transistor design in future generation devices. In transfer aspect, a promising process entails deposition of a thicker Al_2O_3 gate dielectric (Al_2O_3 capping in **Figure 4a**) followed by metal gate stack directly on top of silicene stack. Subsequent to mica delamination, the whole stack can then be mounted onto a conductive host substrate (see the process flow in **Figure 4a**) as a common back gate setup for electrical characteristics. Unlike the original SEDNE process employs vacuum annealing steps to form firm contact between Al_2O_3 dielectric layer and back gate device substrate, the modified SEDNE approach leads to a seamless interface engineering between silicene and gate dielectric plus metal stack with simplified steps and reduced costs in time and energy.

In terms of delamination, it could also be possible to etch away the growth substrate, whilst the stack stays supported on the

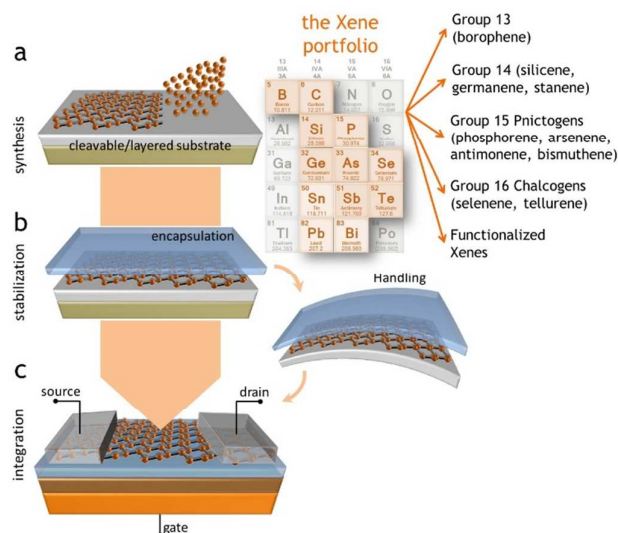


Figure 5 Conceptual flow for the UXEDO process as a universal approach to treat epitaxial Xene that are grown on cleavable or detachable substrates. Xenens of this kind are displayed in the extract from the periodic table around group 13–16 elements, functionalized Xenens are also listed as an engineered material option. Highlighted elements are those where reduction to the epitaxial Xene state has been shown. Process steps consists of: a) epitaxial synthesis of the Xene on substrate and further functionalization if any, b) stabilization via encapsulation and handling, c) integration into an operational device (the field-effect transistor is taken as an example).

host substrate. Dry or wet etching with protected side-walls and back side of the stack could be explored. In this case, it is highly desirable to make use of a selective etchant, which attacks only mica, but does not affect the underlying silver layer. As mica etching could be challenging, it is worthwhile to explore other substrates that are lattice-matched with silver, and relatively easy to etch away selectively. One can also propose that the effort to remove the mica substrate (*e.g.* by etching) may be facilitated by hetero-integration of a sacrificial layer, which is epitaxially compatible and easily removable, in between the substrate and silver (as shown in Approach-2 in **Figure 4b**). Removal of the sacrificial layer through certain selective etching or other techniques would then simultaneously delaminate silicene from the growth substrate. Selective etching of silver for patterning the source-drain contacts and the channel region is another crucial part of device fabrication and instant measurement. Silver electrode patterning is performed by means of two lithography steps: the first electron beam lithography (EBL) on polymethyl methacrylate (PMMA) followed by wet-etching to define source and drain contact pads and the last lithography and etching to open the silicene channel. For silver etching, rapid degradation of silicene was observed with commonly used etchants like nitric acid. A potassium iodide- and iodine-based etchant has been developed in-house for etching to form Ag-free silicene channel.¹⁴⁵ While under-etching leaves silver residue showing metallic I-V response, over-etch instantly degrades silicene. It is therefore essential to perform calibration of etching rate.

TUTORIAL REVIEW

In SEDNE process, a remaining challenge stems from the requirement of developing a post-fabrication passivating layer for silicene. This is a critical step to achieve a robust device enabling any thorough characterization of the electronic transport through silicene because of an extremely rapid (~2 minutes for monolayer) degradation at ambient. It is beneficial to integrate a passivation layer with a high resistance to air and moisture and chemical inertness with robustness at cryogenic temperatures to allow for transport studies across a wide range of conditions. Possible candidates in this respect are non-interacting polymers or metal oxide dielectrics. Another approach to improve the durability of silicene devices is a symmetric dual-gate device structure, where silicene is sandwiched between two dielectric layers. Although this conceptual design is originally for better revealing of charge transport behaviour under doping and gate control,⁸ the sandwiched encapsulation is a good passivation strategy for experimental investigation on silicene layer as well.

As the existing fabrication technique continues to evolve, it is important to note three key technical nodes: passivation, conceptual engineering of device structure, and direct synthesis on insulating substrate. As per the last aspect, processing of silicene-based devices would be much simplified if the silver template could be bypassed by use of insulating substrates. In this case, subsequent to growth, silicene could be capped with a thin layer of Al₂O₃ *in situ*, followed by deposition and annealing, if necessary, of Ohmic source, drain contacts and deposition of gate metal all in the same *ex situ* process flow (Approach-3 in Figure 4c). In this respect, having silicene synthesized on device-friendly substrate can significantly facilitate the integration flow and strengthen the process reliability. A requirement for this purpose would rely on the use of non-interacting substrates where silicene can preserve its structural and electronic integrity during device fabrication stages. Reported cases of relevance are currently limited to a designer approach. For instance, transition metal dichalcogenides like MoS₂ and molybdenum telluride (MoTe₂) were proposed as templates for a vdW epitaxy of silicene (and other Xenes)³⁶ therein taking benefit from the intrinsic scalability (either lateral or vertical). Attempts to grow silicene and germanene in these substrates result in a metallic character. Nonetheless, MoS₂-supported silicene proved to bear transistor processing after Al₂O₃ encapsulation.¹⁴⁶ Alternatively, more recently Al₂O₃(0001), namely sapphire, was proposed as a candidate to accommodate silicene and germanene thus paving the way to the exploration of matter-light interaction otherwise inaccessible in metal substrates.³³ This would likely extend the silicene's applications to the photonics field, as pointed out in Section 1.2.

3.4 A universal approach to Xene processing

(111)-terminated silver or gold, that can be supported by mica substrates, are extensively used as templates for a number of epitaxial Xenes such as borophene, germanene, stanene, phosphorene, and antimonene (see Section 2.4). As such, the SEDNE process has merits to recast as a universal path to process this kind of Xenes with a silicene-like approach

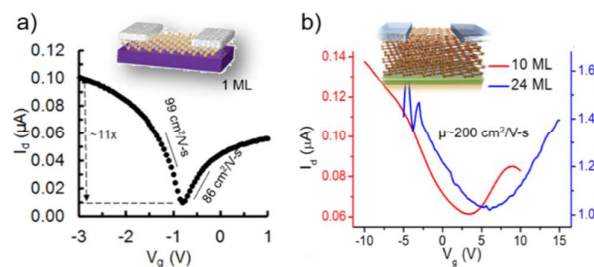


Figure 6 Electrical characterization of silicene devices: a) transfer characteristic curve of monolayer silicene with ~11x gate modulation, exhibiting a gate control over a dozen with extracted low-field carrier mobilities 99 and 86 cm²/Vs for electrons and holes, respectively, b) multilayer silicene with broader ambipolar I_d-V_g curve and mobility ~200 cm²/Vs (stable up to 48 hrs). Panel a) is adapted from Ref.¹⁴⁵ with permission from Springer Nature, copyright 2015. Panel b) is adapted from Ref.¹¹⁷ with permission from the American Chemical Society, copyright 2017.

provided that portability and stability are ensured for each specific material (see Section 2.3). Figure 5 depicted a main flow chart for a Universal Xene Encapsulation, Decoupling and Operation (UXEDO) process. Two key steps are as follows:

All-around encapsulation. It is necessary for air-sensitive Xenes, such as silicene and phosphorene, to have sandwiched or surrounded media to prevent from oxidation or degradation. For instance, a conformal atomic layer deposition of Al₂O₃ could keep the investigated Xene intact for device fabrication. Otherwise, side wall¹⁵⁹ or one-side exposed surface¹⁴¹ will be the leak source for hazardous molecules compromising the 2D materials. A complete encapsulation is inevitable to keep pristine property.

Intact decoupling. Xene needs to transfer from growth substrate to device substrates without sacrificing electrical properties through either mechanical or chemical method. The key is to keep the relevant electronic states of Xene states approximately at the Fermi level, *i.e.* with minimum unintentional extrinsic doping induced. In practice, pre-defined structures (source, drain and gate pads along with the dielectrics) are desired to achieve an intact decoupling of the Xene layer. Another common feature of this universal approach is to have pre-defined pads for optimized surface integration. This provides not only a cleaner interface by direct contact of Xenes with dielectrics or metals without media residue, but also lower contact resistance for better device performance.

Additional merit of this universal Xene process is its versatility to several substrate types: either rigid (silicon, III-V compounds) or flexible substrates (plastic, fabrics, willow glass). This new universal approach addresses a major challenge on material preservation during transfer and device fabrication for silicene, and is applicable to other air-sensitive Xenes such as germanene, stanene, phosphorene, and so on (Section 2.4).

4. Silicene-based technology applications

4.1 Silicene transistors

Silicene FET is an effective vehicle to understand the electronic property of silicene and fulfil various electronic, sensing or even energy devices. After a briefing of simulation studies on silicene transistors, we will focus on experimental performance and its correlation with bandgap engineering through device physics.

Some works summarized simulation studies on dual-gated silicene FET,^{8,160} surface modified silicene FET, silicene thin film FET, silicene nanomesh FET, and silicene nanoribbon FET *etc.* The gate modulation of these silicene FETs varies sharply from 4.2 to 4×10^8 , while bandgap has a narrow distribution of 160 – 680 meV, with more details available elsewhere.^{161,162} Nevertheless, there are new developments in simulation on silicene transistors. Salimian *et al.* proposed a silicene nanotube FET by transfer matrix method.¹⁶³ They investigated the effect of channel length, chirality and diameter of a tube to the channel current, and concluded that I_{\max}/I_{\min} ratio varies by chirality, diameter of silicene nanotubes, and perpendicular electric field, while OFF current strongly depends on characteristics of silicene nanotubes. Patel *et al.* reported that, in a dual-gate silicene FET, decreasing channel length degrades device parameters due to increased leakage whereas decreasing oxide thickness improves these parameters due to increased gate control over the silicene channel uniformly.¹⁶⁴ The evaluation of silicene FET requires considering characteristic parameters, such as mobility, gate modulation (I_{\max}/I_{\min} ratio), subthreshold swing, and transconductance *etc.*, which all connect to one key phrase: bandgap engineering, either through external electrical field or through chemical or physical surface modification.

The Quantum TI-FET

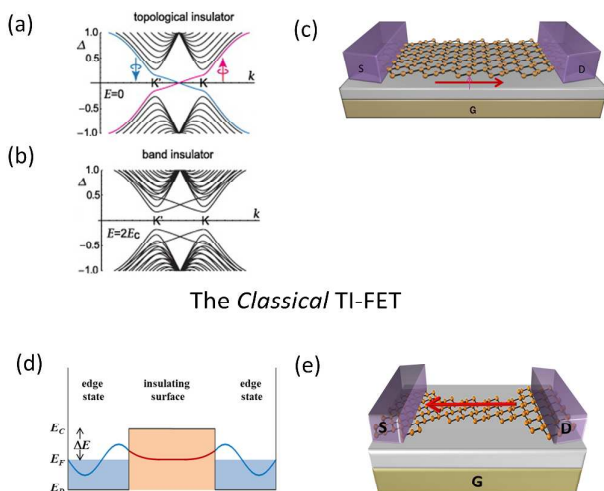


Figure 7 Topology as a technology driver for 2D silicon. a) Non-trivial topological edge states crossing the gap at pristine silicene (ON/topological insulator). There are two edge states since a nanoribbon has two edges (red and blue lines for the left and right edges). b) Edge states disappear with the application of out-of-plane electric field greater than critical field (OFF/band insulator). Panel a) and b) are adapted from Ref.¹⁸¹ with permission from IOP publishing, copyright 2012. c) Quantum TI-FET driven by an out-of-plane electric field d) A simplified double quantum well model (along the transverse direction) for understanding the nanoribbon width and energy barrier dependence of the tunnelling and overlap of the wavefunction of edge electrons. e) Sketch of a classical TI-FET device structure predicated on inter-edge scattering in 2D nanoribbon owing to the nanoribbon width and controllable by the field-effect on E_F .

The experimental investigation on silicene transistors was falling behind theoretical or simulation study due to air stability and fabrication portability issues as previously discussed in Section 3.3. Electrostatic transfer and output measurements, such as drain current (I_d) response to gate voltage (V_g), on monolayer silicene (Ag-free) transistors at ambient condition (Figure 6a) revealed device behaviour similar to graphene, corroborating theoretical expectations on ambipolar Dirac charge transport.¹⁵⁰ This work also supports DFT on p - d hybridization of Si-Ag stabilizes silicene grown on Ag(111).¹⁶⁷ In addition, we discovered that material attributions, such as the number of layers, could play a role in stability and electrical characteristics of silicene transistors. According to electrostatic measurement data, single-layer silicene could yield field-effect mobility $\mu \sim 100$ cm²/Vs at residual carrier density of $\sim 5 \times 10^9$ cm⁻² with gate modulation around $11\times$, whereas multilayer silicene showed a broader I_D - V_g curve, similar mobility ~ 200 cm²/Vs at residual carrier density 2×10^{12} cm⁻² (Figure 6b). Multilayer silicene devices exhibited a notably longer lifetime up to 48 hrs compared to 2 minutes for monolayer counterparts.¹¹⁷ The field-effect mobility (μ) and residual carrier density (n_o) data provide abundant information about band structure of silicene. Though pristine freestanding silicene is predicted to offer intrinsic mobility ~ 1000 cm²/Vs,¹⁶⁸ the substrate effect on silicene's acoustic phonon energies and electron-phonon coupling is likely perturbed negatively in our experimental studies and result in much lower mobility value than calculated. We recall here that the thermally generated n_o of a Dirac semiconductor with zero bandgap inherently depends on the Fermi velocity (v_F), with $n_o \propto (1/v_F^2)$.¹⁶⁹ Given v_F of silicene is comparable to graphene^{48,168} the most plausible scenario to understand silicene's low n_o necessitates a small bandgap opening. In the limit of a weak perturbation to the Dirac dispersion of Ag-free silicene, the small bandgap that yields $n_o \sim 8 \times 10^9$ cm⁻² is calculated to be ~ 210 meV,¹⁴⁵ falling into theoretically calculated range.¹⁶²

Silicene exhibits higher SOC than graphene and bandgap opening is relatively more practical.¹⁷⁰ The latter one is a unique property of Xenes like silicene and other akin materials such as exfoliated phosphorene, enabling potential applications based on their transistors that graphene cannot afford (Table 1). Core research on silicene transistors is about bandgap engineering. There are four approaches to engineer bandgap of silicene: chemical modification, topography (*e.g.* nanoribbons), coupling to a substrate and external electric field.¹⁷¹ Above-mentioned experimental work on back-gate silicene transistors demonstrated that coupling to a substrate and applying an external vertical electric field could effectively tune or engineer the band structure. It also provides transferable knowledge to explore the other two approaches.

Xene sheets	band gap (eV)	I_{\max}/I_{\min}	field-Effect mobility (cm ² /Vs)	residual carrier density (cm ⁻²)
graphene ¹⁶⁵	0	5-10	18,000*	10^{11-12}
silicene ^{117,145}	0-0.2	10-12	100-200	10^{9-12}
phosphorene ^{8,166}	0.3-2	10^{2-4}	200-1560	-

Table 1. Comparison of Key Device Parameters in Xenos

*measured at room temperature on SiO₂, †exfoliated from black phosphorus

4.2 Silicene for a topology-based electronics

As discussed in Section 2, silicene is predicted to be a 2D-TI that host QSH effect.^{5,172} 2D TIs are characterized by bulk insulating states with gapless helical edge states, which are protected against backscattering by time-reversal symmetry.^{28,173} Transport through the helical edge states is ideally dissipationless even in the presence of (non-magnetic) defects, meaning that silicene could be a very promising material for applications in low-energy electronics. With recent advances in research on topological states of matter, the QSH state has emerged as a potential candidate for building novel quantum mechanical switches, the so-called *topological insulator field-effect transistors* (TI-FETs) (Figure 7c), which can be turned ON/OFF by a topological phase transition, as opposed to the conventional charge accumulation/depletion.¹² The topological phase can be switched (Figures 7a and b) between a non-trivial, ballistic QSH edge state (ON state) and a trivial insulating bulk state (OFF state) with the application of gate voltage (strain). The critical field required ($E_c = 2\lambda_{SO}/\delta$, where $2\lambda_{SO}$ is the SOC induced bandgap), for topological phase transition in silicene is approximately 0.05eV/nm, almost an order of magnitude less than the breakdown field of conventional solid-state-dielectrics ($E_{SiO_2} \sim 0.7$ eV/nm). Nonetheless, higher (lower) SOC results in room (low) temperature operation with comparatively high (low) critical electric field for realization of topological phase transition. Hence, silicene, with its moderate SOC (1.55 meV)⁵ that limits its operation in relatively low temperature, is still a potential candidate for realization of gate tunable TI-FET for its practically realizable critical field value for realization of topological phase transition.

A different approach for realization of classical TI-FET (Figure 7e) which is free from the critical field operation temperature trade-off of quantum TI-FET, has been proposed.¹⁷⁴ This scheme is based on engineering inter-edge elastic scattering of the edges of 2D TI materials for modulating the conductivity of the channel and hence puts a limitation on the width of nanoribbon just to make sure the edge-scattering is sufficient to play role in device operation. According to the model, the conductivity decreases ($\sigma \sim e^{\Delta E/kT}$, ΔE is the energy barrier between conduction band and Fermi energy) with electric-field due to the reduction in energy barrier facilitated by phonon-mediated scattering of edge electrons into bulk states. This is quite opposite to the operation of conventional thermionic FETs, where field-effect enhances carrier density (lowers barrier) and hence enhances conductivity ($\sigma \sim e^{-\Delta E/kT}$). The device physics can be better understood by using double quantum well model (Figure 7d) where quantum wells are represented as edge states in an insulating bulk. In this model, the basic parameter is the field-dependent tunnelling length scale given by

$$l_t = \frac{\hbar}{\sqrt{2m^*\Delta E}} \approx \frac{2 \text{ \AA}}{\sqrt{\left(\frac{m^*}{m_0}\right)\Delta E}}$$

where ΔE is in eV, and m^* and m_0 are the effective mass of bulk electron and the electron rest mass, respectively. This amounts to ~ 10 nm for silicene and to ensure sufficient inter-edge interactions, nanoribbon silicene TI-FETs should have ribbon widths within an order of magnitude of l_t ($\lesssim 100$ nm).

4.3 Perspectives for silicene-based junction

Besides homogeneous silicene sheets or nanotubes, there is a rising research interests in silicene based junction or heterostructures that hold great promise for exotic electronic,^{160,164} magnetic, and thermal applications.^{175–177} Zhou *et al.* theoretically investigated the spin transport in the silicene channel with a Fe(111)/silicene stack injector.¹⁷⁶ The partial DOS of Fe layer in this combination shows that spin-down states dominate above the Fermi level, resulting in a negligible spin-up current and high spin injection efficiency. Thus, they present Fe(111)/silicene heterostructure as a good candidate for achieving efficient spin injection devices. Interestingly, thermo-spin is another option for silicene-junction based spintronic devices. Zhai *et al.* proposed a heterojunction of silicene (or germanene) intercalated between two ferromagnetic dielectric layers, and such configure with a proximity-induced asymmetric magnetic field could yield an attractive phenomenon named valley-locked spin-dependent Seebeck effect (VL-SSE) driven by a thermal gradient.¹⁷⁵ The VL-SSE operates in a way that charge carriers from only one valley get thermally excited, having opposite spin polarization counter-propagating along the direction of thermal gradient, whereas nearly zero carrier excited from the other insulating valley due to relatively wide bandgap. It is worth noting that thermal conductivity is a critical parameter still lacking in research for silicene-based junction or heterostructure applications. Zhang *et al.* utilizes a multiscale modelling approach (molecular dynamics plus finite element analysis) to investigate the heat dissipation in 2D transistors based on phosphorene and silicene.¹⁷⁷ They found that the heat dissipation ability of 2D transistors improves by increasing thermal conductivities of the channel and substrate, *i.e.* forming phosphorene/silicene heterostructure. Energy storage is another field of interest for application involving silicene-based junctions. In this respect, multilayer silicene was recently proposed as anode for Li-ion batteries instead of conventionally used graphite.¹⁷⁸ This direction is driven by the exceptionally high specific capacity of silicon (4200 mAhg⁻¹) against graphite (≈ 371 mAhg⁻¹). Similar to graphite, multilayer silicene would enable intercalation/deintercalation of lithium atoms during charge/discharge of the cell despite an expected capacity scaling down to ≈ 954 mAhg⁻¹. This functionality is basically supported by the high versatility of silicene to recast in multiple lithiated configurations.¹⁷⁹ To this purpose, multilayer silicene grown on Ag(111) (see Section 1.4) can be readily processed via encapsulation free decoupling as described in

Section 2.4 for further integration into a battery-functional junction. However, thorough investigation of the oxidation mechanism in multilayer silicene is demanding to assess the environmental stability in such a configuration.

Overall, silicene-based junctions have plenty perspective applications not limited to aforementioned thermal or spintronic (magnetic-spin or thermos-spin) or energy cases, as silicene is an ideal channel material with prominent characters, such as tunable band gap and compatibility with the ubiquitous semiconductor industry.

5. Conclusions

Since the rise of the epitaxial silicene in 2012, not only the literature about it has been tremendously expanding but also new research forefronts on silicene-like Xenes have been consequently triggered. As a result, silicene and its derivatives could potentially offer an intriguing platform for both fundamental research and device applications. For the former aspect, Xenes share the common virtue of buckled elemental atomic sheets, which encloses a rich variety of fundamental properties of matter and particle physics. Concomitantly, as for the latter aspect, previous challenges in accessibility, stability, and portability of silicene have been addressed, with promising experimental device study supporting theoretical prediction on Dirac cone existing in its electronic band structure. Witnessed by the recent experimental research progress in epitaxial growth, Raman characterization, interface engineering, and device fabrication, silicene and its derivatives are one step closer to integration of QSH effect, topological bits, flexible electronics, and energy devices, to name a few. Nevertheless, there is still a lack of experimental study on these exotic device ideas that may inspire or foster on-going work on a silicene impact on societal challenges. Specifically, several material and processing correlated issues are yet to be resolved, such as interface matching between Xenes and dielectrics or metals contacts on device substrates, and post-fabrication passivation. Inspiringly, recent synthetic (superlattice and number of layer control) and processing (sandwich encapsulation) advances are encouraging towards further adventure on exotic quantum and topological phenomena in silicene and similar Xenes for potentially innovative device concepts revolutionizing nowadays semiconductor technology.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

A.M. acknowledges funding support from H2020 ERC CoG 2017 Grant N. 772261 "XFab". A.M. and C.G. acknowledge funding support from CNR grant Laboratori Congiunti "SFET", and Fondazione CARIPOLO – Regione Lombardia for the project "Crystal", grant N. 2016-0978. D.A. acknowledges support from the Army Research Office (ARO), the Presidential Early

Career Award for Engineers and Scientists (PECASE), and the Gordon and Betty Moore Foundation. L.T. acknowledges support from Young 1000-Talent Award #12 (2016), National Science Foundation of China (51602051), Fundamental Research Funds for the Central Universities (2242017R30008).

Notes and references

- 1 M. I. Katsnelson, *Mater. Today*, 2007, **10**, 20–27.
- 2 J. E. Moore, *Nature*, 2010, **464**, 194–198.
- 3 K. Takeda and K. Shiraishi, *Phys. Rev. B*, 1994, **50**, 14916–14922.
- 4 K. S. Novoselov, A. K. Geim, S. V. Morozov, D. Jiang, Y. Zhang, S. V. Dubonos, I. V. Grigorieva and A. A. Firsov, *Science*, 2004, **306**, 666–9.
- 5 C.-C. Liu, W. Feng and Y. Yao, *Phys. Rev. Lett.*, 2011, **107**, 076802.
- 6 S. Rachel and M. Ezawa, *Phys. Rev. B*, 2014, **89**, 195303.
- 7 C. Xu, G. Luo, Q. Liu, J. Zheng, Z. Zhang, S. Nagase, Z. Gao and J. Lu, *Nanoscale*, 2012, **4**, 3111.
- 8 Z. Ni, Q. Liu, K. Tang, J. Zheng, J. Zhou, R. Qin, Z. Gao, D. Yu and J. Lu, *Nano Lett.*, 2012, **12**, 113–118.
- 9 N. D. Drummond, V. Zolyomi and V. I. Fal'ko, *Phys. Rev. B*, 2012, **85**, 075423.
- 10 H. Bao, W. Liao, J. Guo, H. Zhao and G. Zhou, *Laser Phys. Lett.*, 2015, **12**, 095902.
- 11 H. M. Le, T.-T. Pham, T. S. Dinh, Y. Kawazoe and D. Nguyen-Manh, *J. Phys. Condens. Matter*, 2016, **28**, 135301.
- 12 A. Molle, J. Goldberger, M. Houssa, Y. Xu, S. C. Zhang and D. Akinwande, *Nat. Mater.*, 2017, **16**, 163–169.
- 13 N. M. Johnson and M. D. Moyer, *Appl. Phys. Lett.*, 1985, **46**, 787–789.
- 14 N. M. Johnson, D. K. Biegelsen and M. D. Moyer, *Appl. Phys. Lett.*, 1982, **40**, 882–884.
- 15 S. Trivedi, A. Srivastava and R. Kurchania, *J. Comput. Theor. Nanosci.*, 2014, **11**, 781–788.
- 16 M. M. Atalla, E. Tannenbaum and E. J. Scheibner, *Bell Syst. Tech. J.*, 1959, **38**, 749–783.
- 17 H. N. Khan, D. A. Hounshell and E. R. H. Fuchs, *Nat. Electron.*, 2018, **1**, 14–21.
- 18 M. Bruel, *MRS Bull.*, 1998, **23**, 35–39.
- 19 K. Uchida and S. Takagi, *Appl. Phys. Lett.*, 2003, **82**, 2916–2918.
- 20 Z.-G. Shao, X.-S. Ye, L. Yang and C.-L. Wang, *J. Appl. Phys.*, 2013, **114**, 093712.
- 21 S. Cahangirov, M. Topsakal, E. Aktürk, H. Şahin and S. Ciraci, *Phys. Rev. Lett.*, 2009, **102**, 236804.
- 22 D. Jose and A. Datta, *J. Phys. Chem. C*, 2012, **116**, 24639–24648.
- 23 S. Ryu and Y. Hatsugai, *Phys. Rev. Lett.*, 2002, **89**, 077002.
- 24 M. Houssa, G. Pourtois, V. V. Afanas'ev and A. Stesmans, *Appl. Phys. Lett.*, 2010, **97**, 112106.
- 25 G. G. Guzmán-Verri and L. C. Lew Yan Voon, *Phys. Rev. B*, 2007, **76**, 075131.
- 26 L. Matthes, O. Pulci and F. Bechstedt, *New J. Phys.*, 2014, **16**, 105007.
- 27 M. Ezawa, *EPL*, 2013, **104**, 27006.

- 28 C. L. Kane and E. J. Mele, *Phys. Rev. Lett.*, 2005, **95**, 226801.
- 29 M. Ezawa, *Phys. Rev. Lett.*, 2012, **109**, 055502.
- 30 C. J. Tabert and E. J. Nicol, *Phys. Rev. Lett.*, 2013, **110**, 197402.
- 31 M. Ezawa, *J. Phys. Soc. Japan*, 2015, **84**, 121003.
- 32 C. Grazianetti, E. Cinquanta and A. Molle, *2D Mater.*, 2016, **3**, 012001.
- 33 M. X. Chen, Z. Zhong and M. Weinert, *Phys. Rev. B*, 2016, **94**, 075409.
- 34 A. Bhattacharya, S. Bhattacharya and G. P. Das, *Appl. Phys. Lett.*, 2013, **103**, 123113.
- 35 S. Kokott, P. Pflugradt, L. Matthes and F. Bechstedt, *J. Phys. Condens. Matter*, 2014, **26**, 185002.
- 36 E. Scalise, M. Houssa, E. Cinquanta, C. Grazianetti, B. van den Broek, G. Pourtois, A. Stesmans, M. Fanciulli and A. Molle, *2d Mater.*, 2014, **1**, 11010.
- 37 H. Liu, J. Gao and J. Zhao, *J. Phys. Chem. C*, 2013, **117**, 10353–10359.
- 38 M. Houssa, B. van den Broek, E. Scalise, G. Pourtois, V. V. Afanas'ev and A. Stesmans, *Phys. Chem. Chem. Phys.*, 2013, **15**, 3702.
- 39 P. Pflugradt, L. Matthes and F. Bechstedt, *New J. Phys.*, 2014, **16**, 075004.
- 40 P. De Padova, P. Perfetti, B. Olivieri, C. Quaresima, C. Ottaviani and G. Le Lay, *J. Phys. Condens. Matter*, 2012, **24**, 223001.
- 41 S. Sheng, R. Ma, J. Wu, W. Li, L. Kong, X. Cong, D. Cao, W. Hu, J. Gou, J.-W. Luo, P. Cheng, P.-H. Tan, Y. Jiang, L. Chen and K. Wu, *Nano Lett.*, 2018, acs.nanolett.8b00289.
- 42 A. J. Mannix, X.-F. Zhou, B. Kiraly, J. D. Wood, D. Alducin, B. D. Myers, X. Liu, B. L. Fisher, U. Santiago, J. R. Guest, M. J. Yacamán, A. Ponce, A. R. Oganov, M. C. Hersam and N. P. Guisinger, *Science*, 2015, **350**, 1513–6.
- 43 B. Feng, J. Zhang, Q. Zhong, W. Li, S. Li, H. Li, P. Cheng, S. Meng, L. Chen and K. Wu, *Nat. Chem.*, 2016, **8**, 563–568.
- 44 C.-H. Lin, A. Huang, W. W. Pai, W.-C. Chen, T.-Y. Chen, T.-R. Chang, R. Yukawa, C.-M. Cheng, C.-Y. Mou, I. Matsuda, T.-C. Chiang, H.-T. Jeng and S.-J. Tang, *Phys. Rev. Mater.*, 2018, **2**, 024003.
- 45 J. Yuhara, Y. Fujii, K. Nishino, N. Isobe, M. Nakatake, L. Xian, A. Rubio and G. Le Lay, *2D Mater.*, 2018, **5**, 025002.
- 46 Y. Shao, Z.-L. Liu, C. Cheng, X. Wu, H. Liu, C. Liu, J.-O. Wang, S.-Y. Zhu, Y.-Q. Wang, D.-X. Shi, K. Ibrahim, J.-T. Sun, Y.-L. Wang and H.-J. Gao, *Nano Lett.*, 2018, **18**, 2133–2139.
- 47 H.-C. Hsu, Y.-H. Lu, T.-L. Su, W.-C. Lin and T.-Y. Fu, *Semicond. Sci. Technol.*, DOI:10.1088/1361-6641/aaad88.
- 48 P. Vogt, P. De Padova, C. Quaresima, J. Avila, E. Frantzeskakis, M. C. Asensio, A. Resta, B. Ealet and G. Le Lay, *Phys. Rev. Lett.*, 2012, **108**, 155501.
- 49 C.-L. Lin, R. Arafune, K. Kawahara, N. Tsukahara, E. Minamitani, Y. Kim, N. Takagi and M. Kawai, *Appl. Phys. Express*, 2012, **5**, 045802.
- 50 L. Meng, Y. L. Wang, L. Z. Zhang, S. X. Du, R. T. Wu, L. F. Li, Y. Zhang, G. Li, H. T. Zhou, W. A. Hofer and H.-J. Gao, *Nano Lett.*, 2013, **13**, 685–690.
- 51 D. Chiappe, E. Scalise, E. Cinquanta, C. Grazianetti, B. van den Broek, M. Fanciulli, M. Houssa and A. Molle, *Adv. Mater.*, 2014, **26**, 2096–2101.
- 52 T. Aizawa, S. Suehara and S. Otani, *J. Phys. Chem. C*, 2014, **118**, 23049–23057.
- 53 L. Huang, Y.-F. Zhang, Y.-Y. Zhang, W. Xu, Y. Que, E. Li, J.-B. Pan, Y.-L. Wang, Y. Liu, S.-X. Du, S. T. Pantelides and H.-J. Gao, *Nano Lett.*, 2017, **17**, 1161–1166.
- 54 M. De Crescenzi, I. Berbezier, M. Scarselli, P. Castrucci, M. Abbarchi, A. Ronda, F. Jardali, J. Park and H. Vach, *ACS Nano*, 2016, **10**, 11163–11171.
- 55 D. Chiappe, C. Grazianetti, G. Tallarida, M. Fanciulli and A. Molle, *Adv. Mater.*, 2012, **24**, 5088–5093.
- 56 B. J. Feng, Z. J. Ding, S. Meng, Y. G. Yao, X. Y. He, P. Cheng, L. Chen and K. H. Wu, *Nano Lett.*, 2012, **12**, 3507–3511.
- 57 H. Jamgotchian, Y. Colignon, N. Hamzaoui, B. Ealet, J. Y. Hoarau, B. Aufray and J. P. Bibérian, *J. Phys. Condens. Matter*, 2012, **24**, 172001.
- 58 R. Arafune, C.-L. Lin, K. Kawahara, N. Tsukahara, E. Minamitani, Y. Kim, N. Takagi and M. Kawai, *Surf. Sci.*, 2013, **608**, 297–300.
- 59 C. Grazianetti and A. Molle, in *GraphiTA*, eds. V. Morandi and L. Ottaviano, Springer, Cham, 2017, pp. 137–152.
- 60 Z.-L. Liu, M.-X. Wang, J.-P. Xu, J.-F. Ge, G. Le Lay, P. Vogt, D. Qian, C.-L. Gao, C. Liu and J.-F. Jia, *New J. Phys.*, 2014, **16**, 075006.
- 61 C. Grazianetti, D. Chiappe, E. Cinquanta, M. Fanciulli and A. Molle, *J. Phys. Condens. Matter*, 2015, **27**, 255005.
- 62 D. Tsoutsou, E. Xenogiannopoulou, E. Golias, P. Tsipas and A. Dimoulas, *Appl. Phys. Lett.*, 2013, **103**, 231604.
- 63 E. Cinquanta, G. Fratesi, S. dal Conte, C. Grazianetti, F. Scotognella, S. Stagira, C. Vozzi, G. Onida and A. Molle, *Phys. Rev. B*, 2015, **92**, 165427.
- 64 J. Avila, P. De Padova, S. Cho, I. Colombo, S. Lorcy, C. Quaresima, P. Vogt, a. Resta, G. Le Lay and M. C. Asensio, *J. Phys. Condens. Matter*, 2013, **25**, 262001.
- 65 S. K. Mahatha, P. Moras, V. Bellini, P. M. Sheverdyeva, C. Struzzi, L. Petaccia and C. Carbone, *Phys. Rev. B*, 2014, **89**, 201416.
- 66 L. Chen, C.-C. Liu, B. Feng, X. He, P. Cheng, Z. Ding, S. Meng, Y. Yao and K. Wu, *Phys. Rev. Lett.*, 2012, **109**, 056804.
- 67 Y. Feng, D. Liu, B. Feng, X. Liu, L. Zhao, Z. Xie, Y. Liu, A. Liang, C. Hu, Y. Hu, S. He, G. Liu, J. Zhang, C. Chen, Z. Xu, L. Chen, K. Wu, Y.-T. Liu, H. Lin, Z.-Q. Huang, C.-H. Hsu, F.-C. Chuang, A. Bansil and X. J. Zhou, *Proc. Natl. Acad. Sci. U. S. A.*, 2016, **113**, 14656–14661.
- 68 P. Gori, O. Pulci, F. Ronci, S. Colonna and F. Bechstedt, *J. Appl. Phys.*, 2013, **114**, 113710.
- 69 Z.-X. Guo, S. Furuya, J. Iwata and A. Oshiyama, *Phys. Rev. B*, 2013, **87**, 235435.
- 70 Y.-P. Wang and H.-P. Cheng, *Phys. Rev. B*, 2013, **87**, 245430.
- 71 C. L. Lin, R. Arafune, K. Kawahara, M. Kanno, N. Tsukahara, E. Minamitani, Y. Kim, M. Kawai and N. Takagi, *Phys. Rev. Lett.*, 2013, **110**, 76801.
- 72 S. Cahangirov, M. Audiffred, P. Tang, A. Iacomino, W. Duan, G. Merino and A. Rubio, *Phys. Rev. B*, 2013, **88**, 035432.
- 73 M. Satta, S. Colonna, R. Flammini, A. Cricenti and F. Ronci, *Phys. Rev. Lett.*, 2015, **115**, 026102.

- 74 C. Vacacela Gomez, M. Pissarra, M. Gravina, P. Riccardi and A. Sindona, *Phys. Rev. B*, 2017, **95**, 085419.
- 75 R. Friedlein and Y. Yamada-Takamura, *J. Phys. Condens. Matter*, 2015, **27**, 203201.
- 76 H. Okamoto, Y. Sugiyama and H. Nakano, *Chem. - A Eur. J.*, 2011, **17**, 9864–9887.
- 77 E. Noguchi, K. Sugawara, R. Yaokawa, T. Hitosugi, H. Nakano and T. Takahashi, *Adv. Mater.*, 2015, **27**, 856–860.
- 78 R. Yaokawa, T. Ohsuna, T. Morishita, Y. Hayasaka, M. J. S. Spencer and H. Nakano, *Nat. Commun.*, 2016, **7**, 10657.
- 79 S. M. Pratik, A. Nijamudheen and A. Datta, *Chem. - A Eur. J.*, 2015, **21**, 18454–18460.
- 80 S. Yamanaka, *Dalt. Trans.*, 2010, **39**, 1901–1915.
- 81 A. M. Tokmachev, D. V. Averyanov, I. A. Karateev, O. E. Parfenov, A. L. Vasiliev, S. N. Yakunin and V. G. Storchak, *Nanoscale*, 2016, **8**, 16229–16235.
- 82 A. M. Tokmachev, D. V. Averyanov, I. A. Karateev, O. E. Parfenov, O. A. Kondratev, A. N. Taldenkov and V. G. Storchak, *Adv. Funct. Mater.*, 2017, **27**, 1606603.
- 83 A. M. Tokmachev, D. V. Averyanov, O. E. Parfenov, A. N. Taldenkov, I. A. Karateev, I. S. Sokolov, O. A. Kondratev and V. G. Storchak, *Nat. Commun.*, 2018, **9**, 1672.
- 84 S. Botti, J. A. Flores-Livas, M. Amsler, S. Goedecker and M. A. L. Marques, *Phys. Rev. B - Condens. Matter Mater. Phys.*, 2012, **86**, 121204.
- 85 Y. He, H. Li, Y. Sui, J. Qi, Y. Wang, Z. Chen, J. Dong and X. Li, *Sci. Rep.*, 2015, **5**, 14792.
- 86 F. Gimbert, C.-C. Lee, R. Friedlein, A. Fleurence, Y. Yamada-Takamura and T. Ozaki, *Phys. Rev. B*, 2014, **90**, 165423.
- 87 H. I. T. Hauge, M. A. Verheijen, S. Conesa-Boj, T. Etzelstorfer, M. Watzinger, D. Kriegner, I. Zardo, C. Fasolato, F. Capitani, P. Postorino, S. Kölling, A. Li, S. Assali, J. Stangl and E. P. A. M. Bakkers, *Nano Lett.*, 2015, **15**, 5855–5860.
- 88 D. Y. Kim, S. Stefanoski, O. O. Kurakevych and T. A. Strobel, *Nat. Mater.*, 2014, **14**, 169–173.
- 89 S. W. Kim, J. Lee, J. H. Sung, D. Seo, I. Kim, M.-H. Jo, B. W. Kwon, W. K. Choi and H.-J. Choi, *ACS Nano*, 2014, **8**, 6556–6562.
- 90 S. Lee, K. Kim, K. P. Dhakal, H. Kim, W. S. Yun, J. Lee, H. Cheong and J.-H. Ahn, *Nano Lett.*, 2017, **17**, 7744–7750.
- 91 J. Lang, B. Ding, S. Zhang, H. Su, B. Ge, L. Qi, H. Gao, X. Li, Q. Li and H. Wu, *Adv. Mater.*, 2017, **29**, 1701777.
- 92 V. Q. Bui, T.-T. Pham, H.-V. S. Nguyen and H. M. Le, *J. Phys. Chem. C*, 2013, **117**, 23364–23371.
- 93 N. Gao, G. Y. Lu, Z. Wen and Q. Jiang, *J. Mater. Chem. C*, 2017, **5**, 627–633.
- 94 B. Huang, H. J. Xiang and S.-H. Wei, *Phys. Rev. Lett.*, 2013, **111**, 145502.
- 95 H. Sahin and F. M. Peeters, *Phys. Rev. B*, 2013, **87**, 085423.
- 96 J. Sivek, H. Sahin, B. Partoens and F. M. Peeters, *Phys. Rev. B*, 2013, **87**, 085444.
- 97 T. H. Osborn, A. A. Farajian, O. V. Pupyshcheva, R. S. Aga and L. C. Lew Yan Voon, *Chem. Phys. Lett.*, 2011, **511**, 101–105.
- 98 Y. Du, J. Zhuang, H. Liu, X. Xu, S. Eilers, K. Wu, P. Cheng, J. Zhao, X. Pi, K. W. See, G. Peleckis, X. Wang and S. X. Dou, *ACS Nano*, 2014, **8**, 10019–10025.
- 99 R. Wang, X. Pi, Z. Ni, Y. Liu, S. Lin, M. Xu and D. Yang, *Sci. Rep.*, 2013, **3**, 3507.
- 100 T. Morishita and M. J. S. Spencer, *Sci. Rep.*, 2015, **5**, 17570.
- 101 V. O. Özçelik, S. Cahangirov and S. Ciraci, *Phys. Rev. Lett.*, 2014, **112**, 246803.
- 102 Y. Y. Du, J. Zhuang, J. Wang, Z. Li, H. Liu, J. Zhao, X. Xu, H. Feng, L. Chen, K. Wu, X. Wang and S. X. Dou, *Sci. Adv.*, 2016, **2**, e1600067–e1600067.
- 103 J. Qiu, H. Fu, Y. Xu, Q. Zhou, S. Meng, H. Li, L. Chen and K. Wu, *ACS Nano*, 2015, **9**, 11192–11199.
- 104 W. Wang, W. Olovsson and R. I. G. Uhrberg, *Phys. Rev. B*, 2016, **93**, 081406.
- 105 J. Qiu, H. Fu, Y. Xu, A. I. Oreshkin, T. Shao, H. Li, S. Meng, L. Chen and K. Wu, *Phys. Rev. Lett.*, 2015, **114**, 126101.
- 106 S. Cahangirov, V. O. Özçelik, A. Rubio and S. Ciraci, *Phys. Rev. B*, 2014, **90**, 085426.
- 107 S. Cahangirov, V. O. Özçelik, L. Xian, J. Avila, S. Cho, M. C. Asensio, S. Ciraci and A. Rubio, *Phys. Rev. B*, 2014, **90**, 035448.
- 108 P. Vogt, P. Capiod, M. Berthe, A. Resta, P. De Padova, T. Bruhn, G. Le Lay and B. Grandidier, *Appl. Phys. Lett.*, 2014, **104**, 021602.
- 109 F. Zhu, W. Chen, Y. Xu, C. Gao, D. Guan, C. Liu, D. Qian, S.-C. Zhang and J. Jia, *Nat. Mater.*, 2015, **14**, 1020–1025.
- 110 P. De Padova, A. Generosi, B. Paci, C. Ottaviani, C. Quaresima, B. Olivieri, E. Salomon, T. Angot and G. Le Lay, *2D Mater.*, 2016, **3**, 031011.
- 111 H. Fu, L. Chen, J. Chen, J. Qiu, Z. J. Ding, J. Zhang, K. H. Wu, H. Li and S. Meng, *Nanoscale*, 2015, **7**, 15880–15885.
- 112 P. De Padova, P. Vogt, A. Resta, J. Avila, I. Razado-Colambo, C. Quaresima, C. Ottaviani, B. Olivieri, T. Bruhn, T. Hirahara, T. Shirai, S. Hasegawa, M. Carmen Asensio and G. Le Lay, *Appl. Phys. Lett.*, 2013, **102**, 163106.
- 113 Z. Li, J. Zhuang, L. Chen, Z. Ni, C. Liu, L. Wang, X. Xu, J. Wang, X. Pi, X. Wang, Y. Du, K. Wu and S. X. Dou, *ACS Cent. Sci.*, 2016, **2**, 517–521.
- 114 A. J. Mannix, B. Kiraly, B. L. Fisher, M. C. Hersam and N. P. Guisinger, *ACS Nano*, 2014, **8**, 7538–7547.
- 115 S. K. Mahatha, P. Moras, P. M. Sheverdyaeva, R. Flammini, K. Horn and C. Carbone, *Phys. Rev. B*, 2015, **92**, 245127.
- 116 A. Curcella, R. Bernard, Y. Borenzstein, M. Lazzeri, A. Resta, Y. Garreau and G. Prévot, *2D Mater.*, 2017, **4**, 025067.
- 117 C. Grazianetti, E. Cinquanta, L. Tao, P. De Padova, C. Quaresima, C. Ottaviani, D. Akinwande and A. Molle, *ACS Nano*, 2017, **11**, 3376–3382.
- 118 H.-S. Tsai, C.-H. Hsiao, C.-W. Chen, H. Ouyang and J.-H. Liang, *Nanoscale*, 2016, **8**, 9488–9492.
- 119 V. Kochat, A. Samanta, Y. Zhang, S. Showmick, P. Manimunda, S. A. S. Asif, A. S. Stender, R. Vajtai, A. K. Singh, C. S. Tiwary and P. M. Ajayan, *Sci. Adv.*, 2018, **4**, e1701373.
- 120 M. Pumera and Z. Sofer, *Adv. Mater.*, 2017, **29**, 1605299.
- 121 S. Zhang, Z. Yan, Y. Li, Z. Chen and H. Zeng, *Angew. Chemie Int. Ed.*, 2015, **54**, 3112–3115.
- 122 L. Xian, A. Pérez Paz, E. Bianco, P. M. Ajayan and A. Rubio, *2D Mater.*, 2017, **4**, 041003.
- 123 M. Derivaz, D. Dentel, R. Stephan, M. C. Hanf, A.

- Mehdaoui, P. Sonnet and C. Pirri, *Nano Lett.*, 2015, **15**, 2510–2516.
- 124 J. Gou, Q. Zhong, S. Sheng, W. Li, P. Cheng, H. Li, L. Chen and K. Wu, *2D Mater.*, 2016, **3**, 045005.
- 125 M. E. Dávila, L. Xian, S. Cahangirov, A. Rubio and G. Le Lay, *New J. Phys.*, 2014, **16**, 095002.
- 126 L. Li, S. Lu, J. Pan, Z. Qin, Y. Wang, Y. Wang, G. Cao, S. Du and H.-J. Gao, *Adv. Mater.*, 2014, **26**, 4820–4824.
- 127 Z. Qin, J. Pan, S. Lu, Y. Shao, Y. Wang, S. Du, H.-J. Gao and G. Cao, *Adv. Mater.*, 2017, **29**, 1606046.
- 128 L. Zhang, P. Bampoulis, A. N. Rudenko, Q. Yao, A. van Houselt, B. Poelsema, M. I. Katsnelson and H. J. W. Zandvliet, *Phys. Rev. Lett.*, 2016, **116**, 256804.
- 129 F. d'Acapito, S. Torrenco, E. Xenogiannopoulou, P. Tsipas, J. Marquez Velasco, D. Tsoutsou and A. Dimoulas, *J. Phys. Condens. Matter*, 2016, **28**, 045002.
- 130 C.-Z. Xu, Y.-H. Chan, P. Chen, X. Wang, D. Flötotto, J. A. Hlevyack, G. Bian, S.-K. Mo, M.-Y. Chou and T.-C. Chiang, *Phys. Rev. B*, 2018, **97**, 035122.
- 131 Z. Zhu and D. Tománek, *Phys. Rev. Lett.*, 2014, **112**, 176802.
- 132 J. L. Zhang, S. Zhao, C. Han, Z. Wang, S. Zhong, S. Sun, R. Guo, X. Zhou, C. D. Gu, K. Di Yuan, Z. Li and W. Chen, *Nano Lett.*, 2016, **16**, 4903–4908.
- 133 T. Hirahara, G. Bihlmayer, Y. Sakamoto, M. Yamada, H. Miyazaki, S. Kimura, S. Blügel and S. Hasegawa, *Phys. Rev. Lett.*, 2011, **107**, 166801.
- 134 F. Reis, G. Li, L. Dudy, M. Bauernfeind, S. Glass, W. Hanke, R. Thomale, J. Schäfer and R. Claessen, *Science*, 2017, **357**, 287–290.
- 135 X. Wu, Y. Shao, H. Liu, Z. Feng, Y.-L. Wang, J.-T. Sun, C. Liu, J.-O. Wang, Z.-L. Liu, S.-Y. Zhu, Y.-Q. Wang, S.-X. Du, Y.-G. Shi, K. Ibrahim and H.-J. Gao, *Adv. Mater.*, 2017, **29**, 1605407.
- 136 M. Fortin-Deschênes, O. Waller, T. O. Mentş, A. Locatelli, S. Mukherjee, F. Genuzio, P. L. Levesque, A. Hébert, R. Martel and O. Moutanabbir, *Nano Lett.*, 2017, **17**, 4970–4975.
- 137 J. Chen, Y. Dai, Y. Ma, X. Dai, W. Ho and M. Xie, *Nanoscale*, 2017, **9**, 15945–15948.
- 138 X. Huang, J. Guan, Z. Lin, B. Liu, S. Xing, W. Wang and J. Guo, *Nano Lett.*, 2017, **17**, 4619–4623.
- 139 J. Qin, G. Qiu, J. Jian, H. Zhou, L. Yang, A. Charnas, D. Y. Zemlyanov, C.-Y. Xu, X. Xu, W. Wu, H. Wang and P. D. Ye, *ACS Nano*, 2017, **11**, 10222–10229.
- 140 Y. Du, G. Qiu, Y. Wang, M. Si, X. Xu, W. Wu and P. D. Ye, *Nano Lett.*, 2017, **17**, 3965–3973.
- 141 A. Molle, C. Grazianetti, D. Chiappe, E. Cinquanta, E. Cianci, G. Tallarida and M. Fanciulli, *Adv. Funct. Mater.*, 2013, **23**, 4340–4344.
- 142 X. Xu, J. Zhuang, Y. Du, H. Feng, N. Zhang, C. Liu, T. Lei, J. Wang, M. Spencer, T. Morishita, X. Wang and S. X. Dou, *Sci. Rep.*, 2015, **4**, 7543.
- 143 H. Van Bui, F. B. Wiggers, R. Friedlein, Y. Yamada-Takamura, A. Y. Kovalgin and M. P. de Jong, *J. Chem. Phys.*, 2015, **142**, 064702.
- 144 M. Houssa, B. Van Den Broek, E. Scalise, B. Ealet, G. Pourtois, D. Chiappe, E. Cinquanta, C. Grazianetti, M. Fanciulli, A. Molle, V. V. Afanas'ev and A. Stesmans, *Appl. Surf. Sci.*, DOI:10.1016/j.apsusc.2013.09.062.
- 145 L. Tao, E. Cinquanta, D. Chiappe, C. Grazianetti, M. Fanciulli, M. Dubey, A. Molle and D. Akinwande, *Nat. Nanotechnol.*, 2015, **10**, 227–231.
- 146 A. Molle, A. Lamperti, D. Rotta, M. Fanciulli, E. Cinquanta and C. Grazianetti, *Adv. Mater. Interfaces*, 2016, **3**, 1500619.
- 147 R. Friedlein, H. Van Bui, F. B. Wiggers, Y. Yamada-Takamura, A. Y. Kovalgin and M. P. de Jong, *J. Chem. Phys.*, 2014, **140**, 204705.
- 148 E. Cinquanta, E. Scalise, D. Chiappe, C. Grazianetti, B. Van Den Broek, M. Houssa, M. Fanciulli and A. Molle, *J. Phys. Chem. C*, 2013, **117**, 16719–16724.
- 149 A. C. Ferrari and D. M. Basko, *Nat. Nanotechnol.*, 2013, **8**, 235–246.
- 150 E. Scalise, E. Cinquanta, M. Houssa, B. van den Broek, D. Chiappe, C. Grazianetti, G. Pourtois, B. Ealet, A. Molle, M. Fanciulli, V. V. Afanas'ev and A. Stesmans, *Appl. Surf. Sci.*, 2014, **291**, 113–117.
- 151 E. Scalise, M. Houssa, G. Pourtois, B. van den Broek, V. Afanas'ev and A. Stesmans, *Nano Res.*, 2013, **6**, 19–28.
- 152 E. Cinquanta, E. Scalise, D. Chiappe, C. Grazianetti and B. Van Den, *J. Phys. Chem. C*, 2013, 1–24.
- 153 J. Zhuang, X. Xu, Y. Du, K. Wu, L. Chen, W. Hao, J. Wang, W. K. Yeoh, X. Wang and S. X. Dou, *Phys. Rev. B*, 2015, **91**, 161409.
- 154 D. Solonenko, O. D. Gordan, G. Le Lay, H. Şahin, S. Cahangirov, D. R. T. Zahn and P. Vogt, *2D Mater.*, 2016, **4**, 015008.
- 155 S. Sheng, J. Wu, X. Cong, W. Li, J. Gou, Q. Zhong, P. Cheng, P. Tan, L. Chen and K. Wu, *Phys. Rev. Lett.*, 2017, **119**, 196803.
- 156 X. Li, W. Cai, J. An, S. Kim, J. Nah, D. Yang, R. Piner, A. Velamakanni, I. Jung, E. Tutuc, S. K. Banerjee, L. Colombo and R. S. Ruoff, *Science*, 2009, **324**, 1312–4.
- 157 X. Li, Y. Zhu, W. Cai, M. Borysiak, B. Han, D. Chen, R. D. Piner, L. Colombo and R. S. Ruoff, *Nano Lett.*, 2009, **9**, 4359–4363.
- 158 J. W. Suk, A. Kitt, C. W. Magnuson, Y. Hao, S. Ahmed, J. An, A. K. Swan, B. B. Goldberg and R. S. Ruoff, *ACS Nano*, 2011, **5**, 6916–6924.
- 159 J.-S. Kim, Y. Liu, W. Zhu, S. Kim, D. Wu, L. Tao, A. Dodabalapur, K. Lai and D. Akinwande, *Sci. Rep.*, 2015, **5**, 8989.
- 160 Z. Zhuo, X. Wu and J. Yang, *Nanoscale*, 2018, **10**, 1265–1271.
- 161 R.-G. Quhe, Y.-Y. Wang and J. Lü, *Chinese Phys. B*, 2015, **24**, 088105.
- 162 J. Zhao, H. Liu, Z. Yu, R. Quhe, S. Zhou, Y. Wang, C. C. Liu, H. Zhong, N. Han, J. Lu, Y. Yao and K. Wu, *Prog. Mater. Sci.*, 2016, **83**, 24–151.
- 163 F. Salimian and D. Dideban, *ECS J. Solid State Sci. Technol.*, 2018, **7**, M1–M5.
- 164 N. Patel and S. Choudhary, *Superlattices Microstruct.*, 2017, **110**, 155–161.

Journal Name

ARTICLE

- 165 A. Reina, X. Jia, J. Ho, D. Nezich, H. Son, V. Bulovic, M. S. Dresselhaus and J. Kong, *Nano Lett.*, 2009, **9**, 30–35.
- 166 L. Tao, W. Zhu, J.-S. Kim and D. Akinwande, in *2015 45th European Solid State Device Research Conference (ESSDERC)*, IEEE, 2015, pp. 168–171.
- 167 J. Gao, G. Zhang and Y.-W. Zhang, *Sci. Rep.*, 2016, **6**, 29107.
- 168 X. Li, J. T. Mullen, Z. Jin, K. M. Borysenko, M. Buongiorno Nardelli and K. W. Kim, *Phys. Rev. B - Condens. Matter Mater. Phys.*, 2013, **87**, 1–9.
- 169 P. J. Dobson, *Contemp. Phys.*, 2012, **53**, 379–380.
- 170 M. Vali, D. Dideban and N. Moezi, *J. Comput. Electron.*, 2016, **15**, 138–143.
- 171 H. J. W. Zandvliet, *Nano Today*, 2014, **9**, 691–694.
- 172 M. Tahir, A. Manchon, K. Sabeeh and U. Schwingenschlögl, *Appl. Phys. Lett.*, 2013, **102**, 162412.
- 173 X.-L. Qi, T. L. Hughes and S.-C. Zhang, *Phys. Rev. B*, 2010, **81**, 134508.
- 174 W. G. Vandenberghe and M. V. Fischetti, *J. Appl. Phys.*, 2014, **116**, 173707.
- 175 X. Zhai, Y.-T. Wang, R. Wen, S.-X. Wang, Y. Tian, X. Zhou, W. Chen and Z. Yang, *Phys. Rev. B*, 2018, **97**, 085410.
- 176 J. Zhou, A. Bournel, Y. Wang, X. Lin, Y. Zhang and W. Zhao, *Appl. Phys. Lett.*, 2017, **111**, 182408.
- 177 H. Zhang, H. Wang, S. Xiong, H. Han, S. Volz and Y. Ni, *J. Phys. Chem. C*, 2018, **122**, 2641–2647.
- 178 J. Zhuang, X. Xu, G. Peleckis, W. Hao, S. X. Dou and Y. Du, *Adv. Mater.*, 2017, **29**, 1606716.
- 179 G. A. Tritsarlis, E. Kaxiras, S. Meng and E. Wang, *Nano Lett.*, 2013, **13**, 2258–2263.
- 180 H. Enriquez, S. Vizzini, A. Kara, B. Lalmi and H. Oughaddou, *J. Phys. Condens. Matter*, 2012, **24**, 314211.
- 181 M. Ezawa, *New J. Phys.*, 2012, **14**, 033003.